

Study on Chameleonic Dark Matter in $F(R)$ Gravity

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Reference

“Dark matter in modified gravity?” *Phys. Rev. D*95 044040 (2017)

“Cosmic History of Chameleonic Dark Matter in $F(R)$ Gravity” arXiv:1708.08702

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Introduction

Many kinds of Modified Gravity have been investigated.

- UV modification
 - effective theory for quantum gravity
- IR modification
 - **Dark energy** instead of cosmological constant
- How to test modified gravity theories?
 - Cosmology
 - Astrophysics
 - Gravitational Waves
 - **Dark matter**

→ To explore new application of modified gravity from viewpoint of **particle physics**!

N.B.) NOT modified Newtonian dynamics, but particle DM

New Scalar Field in F(R) Gravity

F(R) Gravity is one of modified gravity theories.


F(R) gravity in Jordan Frame : $g_{\mu\nu}$

$$S = \frac{1}{2\kappa^2} \int d^4x \sqrt{-g} F(R)$$

cf.) EH-action

$$\int d^4x \sqrt{-g} R$$

Replace: $R \rightarrow F(R)$

Weyl trans.  $\tilde{g}_{\mu\nu} = \Omega^2(x) g_{\mu\nu}$, $\Omega^2(x) = F_R(R) \equiv e^{2\sqrt{1/6\kappa}\varphi(x)}$

F(R) gravity in Einstein frame : $\tilde{g}_{\mu\nu}$

$$S = \int d^4x \sqrt{-\tilde{g}} \left[\frac{1}{2\kappa^2} \tilde{R} - \frac{1}{2} \tilde{g}^{\mu\nu} \partial_\mu \varphi \partial_\nu \varphi - V(\varphi) \right]$$

$$\text{where } V(\varphi) = \frac{1}{2\kappa^2} \frac{F_R(R)R - F(R)}{F_R^2(R)}$$

New scalar field $\varphi(x)$ appears (**Scalaron**)

Chameleon Mechanism

Viable $F(R)$ gravity possesses **Chameleon mechanism**

[Khoury and Weltman (2004)]

Potential of scalar field $V(\varphi)$ couples with trace of $T_{\mu\nu}$

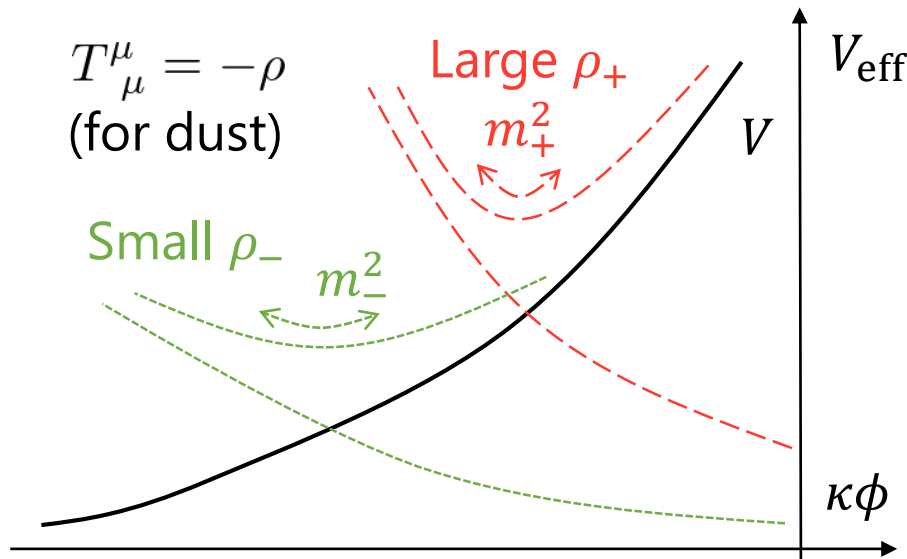
$$\tilde{\square}\varphi = \partial_\varphi V_{\text{eff}}(\varphi), \quad V_{\text{eff}}(\varphi) = V(\varphi) - \frac{1}{4}e^{-4\sqrt{1/6}\kappa\varphi}T^\mu{}_\mu$$



Chameleon Mechanism

$$T^\mu{}_\mu = -\rho$$

(for dust)



Scalaron mass

$$m_\varphi = V''_{\text{eff}}(\varphi_{\text{min}})$$

In high-density region,
scalar field is heavy
and suppressed.

The fifth force is screened
in Solar System

Dark Matter in F(R) Gravity ?

At classical level, scalar field is responsible for DE.

→ Particle picture of scalaron field?

→ "chameleon" particle [Burrage and Sakstein (2017)]

Scalaron's properties

- **SM singlet** scalar field from modified gravity
- **Massive** because of chameleon mechanism
- **Very weak interaction** suppressed by M_{pl}

Scalaron can be a DM candidate?

[Nojiri and Odintsov (2008)], [Cembranos (2009)] etc.

Scalaron Field = **Background** + **Oscillation**
Dark Energy Particle Picture = Dark Matter

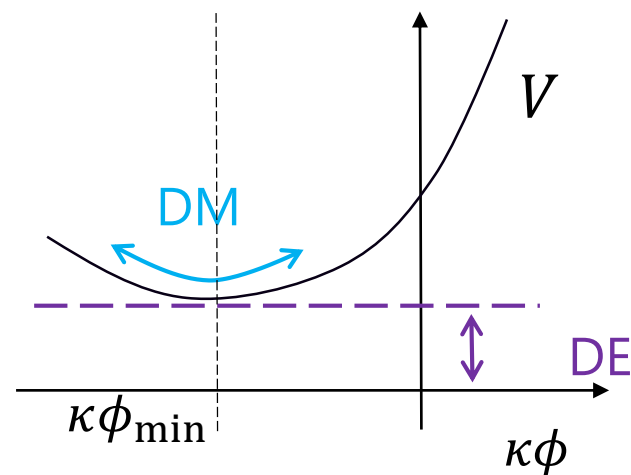
Chameleonic DM and Coincidence Problem

Scalar field with chameleon mechanism

- “Chameleonic” Dark Matter
- Environment-dependence (choice of T_{μ}^{μ})
- **Scalaron mass is NOT constant**, but determined by other ordinary matters
- Depends on cosmic history (time-dependent mass)

Scalar field for two dark components

- Unified treatment of DM & DE in one theory
- To estimate DM-DE ratio, and **address coincidence problem**
- To expect DM and DE densities are of same order



Cosmic Environment in Early Universe

To construct the time evolution of $T_{\mu}^{\mu} = -(\rho - 3p)$.

$$V_{\text{eff}}(\varphi) = V(\varphi) + \frac{1}{4} e^{-4\sqrt{1/6}\kappa\varphi} (\rho - 3p)$$

Trace of Energy-Momentum Tensor

$$\rho - 3p = \frac{gT^4}{2\pi^2} x^2 \int_0^{\infty} d\xi \frac{\xi^2}{\sqrt{\xi^2 + x^2}} \frac{1}{e^{\sqrt{\xi^2 + x^2}} \pm 1} \quad x = \frac{m}{T}, \quad \xi = \frac{p}{T}$$

At high temp. (relativistic)

$$\rho - 3p \approx \frac{g}{24} m^2 T^2 \begin{cases} 2 & \text{for bosons} \\ 1 & \text{for fermions} \end{cases}$$

At low temp. (non-relativistic)

$$\rho - 3p \approx \rho \approx mg \left(\frac{mT}{2\pi} \right)^{3/2} e^{-m/T}$$

For massless particles $\rho - 3p = 0$ (Radiation)

Model of F(R) Gravity

Starobinsky model with R^2 correction

$$F(R) = R - \beta R_c \left[1 - \left(1 + \frac{R^2}{R_c^2} \right)^{-n} \right] + \alpha R^2$$

where $R_c \sim \Lambda$ is constant curvature, and $\alpha, \beta, n > 0$

$$F(R) = R - \beta R_c \left[1 - \left(1 + \frac{R^2}{R_c^2} \right)^{-n} \right]$$

$$\alpha R^2$$

to cure singularity problem

Viable F(R) gravity model for DE

[Starobinsky (2007)]

[Frolov (2008)]

[Kobayashi and Maeda (2008)]

[Dev et al. (2008)]

In large-curvature limit $R > R_c$ (chameleon mechanism works in high-density region),

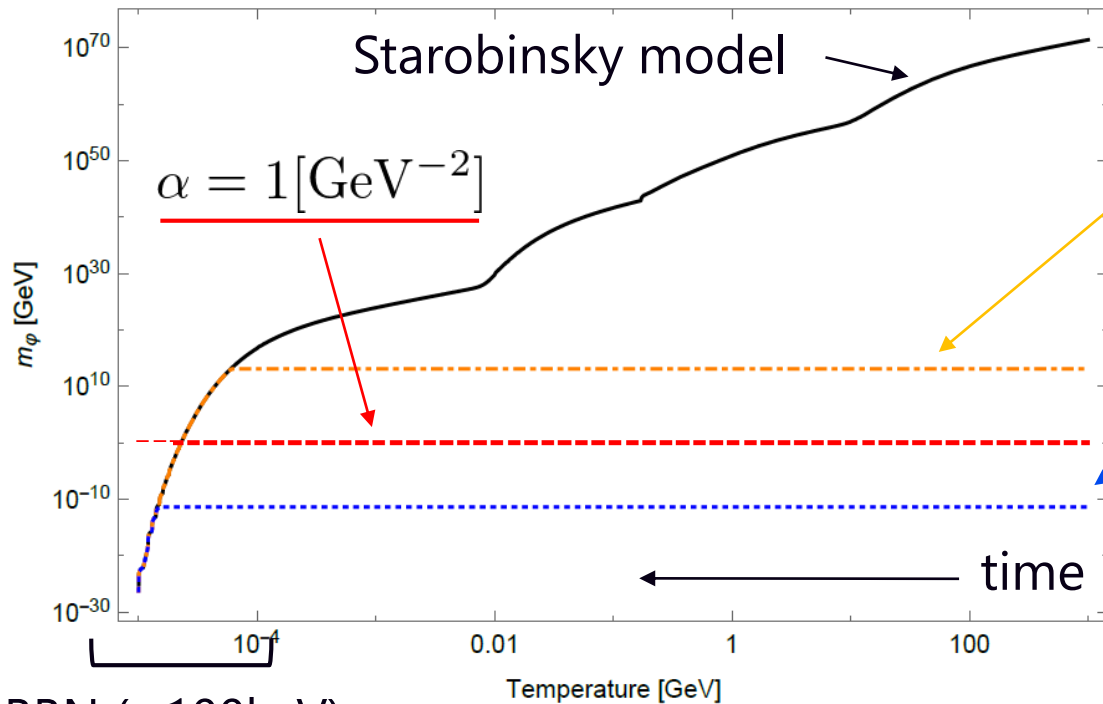
$$F(R) \approx R - \beta R_c + \beta R_c \left(\frac{R}{R_c} \right)^{-2n} + \alpha R^2 \quad \begin{array}{l} \beta R_c \approx 2\Lambda \\ \beta \gtrsim \mathcal{O}(1) \end{array}$$

Scalaron Mass in Early Universe

Mass

$$m_\phi^2 = \frac{R_c}{6n(2n+1)\beta} \left[\left(-\frac{\kappa^2 T}{R_c} \right)^{-2(n+1)} + \frac{\alpha R_c}{n(2n+1)\beta} \right]^{-1} \frac{1}{1 - 2\alpha\kappa^2 T}$$

For $n = 1, \beta = 2, R_c = \Lambda$ cf.) $\beta R_c = 2\Lambda$



$$\alpha = \frac{1}{6M^2}, \quad M = 10^{13} [\text{GeV}]$$

Inflation

$$\alpha = 10^{22} [\text{GeV}^{-2}]$$

bound from experiments

[Adelberger et. al (2006)]

[Berry and Gair (2011)]

BBN ($\sim 100\text{keV}$)

$m_\phi^2 \sim \alpha^{-1}$ in early Universe

cf.) R^2 -inflation

Scalaron Mass in Current Universe

Scalaron mass in the current Universe.

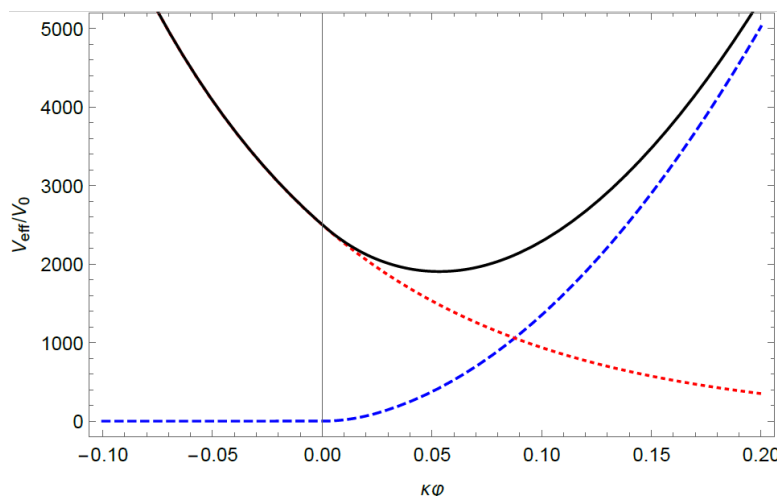
As an example, we study the environment in the galaxy

Typical density

$$-T^{\mu}_{\mu} = \rho \sim 3-5 \times 10^{-25} [\text{g}/\text{cm}^3]$$

Scalaron mass

$$m_{\varphi} = 10^{-24} \sim 10^{-23} [\text{eV}]$$



Scalaron is very light in the current Universe.

Ultralight axion $m \sim 10^{-23} \sim 10^{-22}$ [eV] for problems in **small-scale structure**

[Hu, Barkana, Gruzinov (2000)]

→ “Ultralight scalaron” also solves the problems?

Matter Coupling to SM Particles

Matter Sector

$$\begin{aligned}
 S_{\text{Matter}} &= \int d^4x \sqrt{-g} \mathcal{L}(g^{\mu\nu}, \Psi) \\
 &= \int d^4x \sqrt{-\tilde{g}} e^{-4\sqrt{1/6}\kappa\varphi(x)} \mathcal{L}\left(e^{2\sqrt{1/6}\kappa\varphi(x)} \tilde{g}^{\mu\nu}, \Psi\right)
 \end{aligned}$$

exponential form $e^{Q\kappa\varphi}$

↙

$$\varphi \rightarrow \varphi_{\text{min}} + \varphi$$

$$e^{Q\kappa\varphi(x)} \rightarrow e^{Q\kappa\varphi_{\text{min}}} e^{Q\kappa\varphi(x)}$$

$$\approx \underline{e^{Q\kappa\varphi_{\text{min}}}} \cdot \underline{(1 + Q\kappa\varphi + \mathcal{O}(\kappa^2\varphi^2))}$$

Frame-deference

Coupling to matter

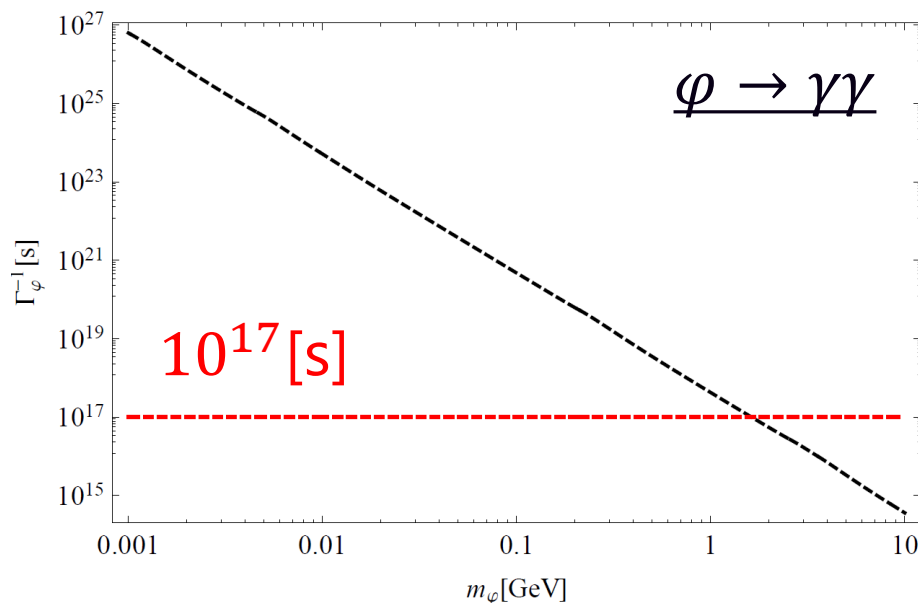
Massless vector field: $\mathcal{L} \supset g^2 \frac{\varphi}{M_{\text{pl}}} F_{\mu\nu}^2$ (induced from anomaly)

Massive fields: $\mathcal{L} \supset m^2 \frac{\varphi}{M_{\text{pl}}} \bar{\psi}\psi, m^2 \frac{\varphi}{M_{\text{pl}}} \tilde{g}^{\mu\nu} A_\mu A_\nu$

cf.) Coupling similar to Axion or Dilatonic DM

Scalaron Lifetime at late-time Universe

At late-time, the scalaron mainly decays into diphotons because scalaron mass becomes smaller in the cosmic history.



Lifetime $\Gamma_\phi^{-1} \geq 10^{17}$ [s]

$\rightarrow m_\phi \leq O(1)$ [GeV]

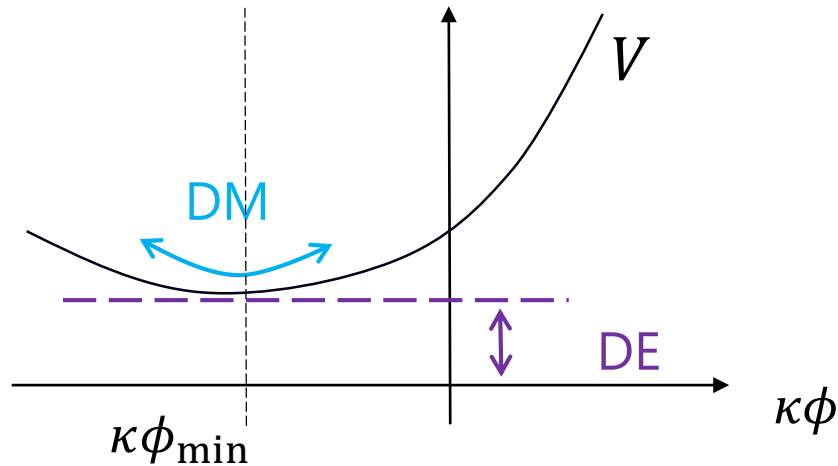
cf.) in galaxy at present

$m_\phi = 10^{-24} \sim 10^{-23}$ [eV]

cf.) Scalaron can be heavy in early Universe because it is in very short time

\rightarrow Small effect to total lifetime

Scalaron Relic Density



$$\rho = \frac{1}{2}\dot{\varphi}^2 + V(\varphi)$$

To estimate scalaron energy density at current Universe

To assume harmonic oscillation approximation is valid.

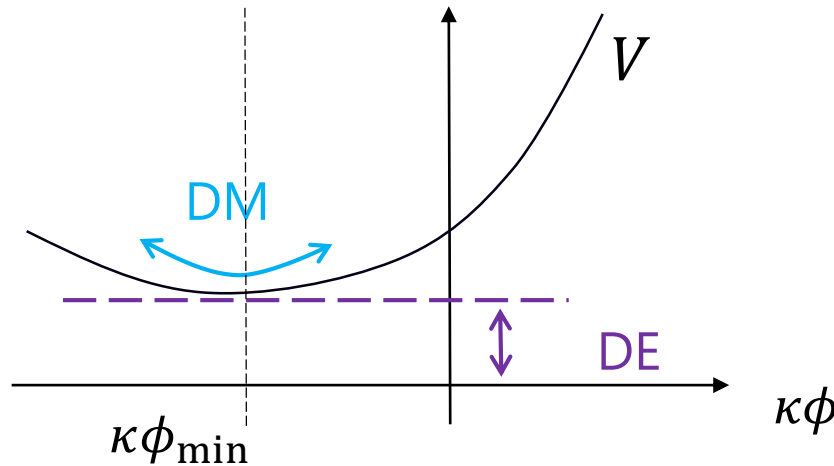
$$\kappa\varphi \approx \underline{\kappa\varphi_0} \cos(mt) + \kappa\varphi_{\min}, \quad V(\varphi) \approx V(\varphi_{\min}) + \frac{1}{2}m^2(\varphi - \varphi_{\min})^2$$

$$\text{Amplitude} \ll 1$$

$$V'(\varphi_{\min}) = 0$$

We obtain
$$\rho \approx \frac{1}{2}m^2\varphi_0^2 + V(\varphi_{\min})$$

Scalaron Relic Density



$$\rho = \frac{1}{2}\dot{\varphi}^2 + V(\varphi)$$

↓ harmonic oscillation approximation

$$\rho \approx \frac{1}{2}m^2\varphi_0^2 + V(\varphi_{\min})$$

$V(\varphi_{\min}) = \frac{\Lambda}{\kappa^2}$ for scalaron potential energy to be DE

$m_\varphi > 3H_0$ for scalaron to harmonically oscillate at present

If we input DM:DE \approx 3:7, we get $\kappa\varphi_0 < 0.3$

- Consistent with approximation, $\kappa\varphi_0 < 1$
- We need **all cosmic history to predict precise DM density** (= initial condition/value of scalaron)

Summary and Discussion

We studied scalaron as new dark matter candidate.

- Mass changes according to cosmic environment
- Very light in current Universe, possibly heavy in early Universe
- Long lifetime to be DM candidate at late-time
- Possibility to address the coincidence problem

Constraints on this scenario from (in-)direct detection

- Heavy scalaron at galactic center decays to photons?
- To discriminate from axion?
- Non-constant mass or chameleon mechanism are keys

Analysis in the early Universe

- Conditions for scalaron to survive in the early Universe?
- To include other BSM? (inflation, Particle creation etc.)
- Depends on our understanding of cosmic history