A ROBUST BAO EXTRACTOR: FROM THEORY TO DATA

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based on JCAP 1708 (2017) no.08, 007 JCAP 1801 (2018) no.01, 035 with Takahiro Nishimichi Eugenio Noda Marco Peloso k [h Mpc⁻¹]

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A TALK ABOUT LARGE SCALE STRUCTURES ...





... BUT ALSO ABOUT MODIFICATIONS OF GR!

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NGr

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NGr = (Newtonian Gravity)

COARSE-GRAINED VLASOV EQUATION

$$\begin{bmatrix} \frac{\partial}{\partial \tau} + \frac{p^{i}}{am} \frac{\partial}{\partial x^{i}} - am \frac{\partial \phi_{mic}}{\partial x^{i}} \frac{\partial}{\partial p^{i}} \end{bmatrix} f_{mic}(\mathbf{x}, \mathbf{p}, \tau) = 0$$
window function
$$f(\mathbf{x}, \mathbf{p}, \tau) = \int d^{3}y W [y/R] f_{mic}(\mathbf{x} + \mathbf{y}, \mathbf{p}, \tau)$$

$$\begin{bmatrix} \frac{\partial}{\partial \tau} + \frac{p^{i}}{am} \frac{\partial}{\partial x^{i}} - am \frac{\partial \phi}{\partial x^{i}} \frac{\partial}{\partial p^{i}} \end{bmatrix} f(\mathbf{x}, \mathbf{p}, \tau) =$$

$$am \begin{bmatrix} \langle \frac{\partial \phi_{mic}}{\partial x^{i}} \frac{\partial f_{mic}}{\partial p^{i}} \rangle_{R} - \frac{\partial \phi}{\partial x^{i}} \frac{\partial f}{\partial p^{i}} \end{bmatrix} (\mathbf{x}, \mathbf{p}, \tau)$$

No divergences from shell-crossing for R > 0 Physics at scale L insensitive to R in the R/L ->0 limit (to be checked!)

BEYOND PRESSURELESS PERFECTION

$$\begin{split} &\frac{\partial}{\partial \tau} \delta(\mathbf{x}) + \frac{\partial}{\partial x^{i}} \left[(1 + \delta(\mathbf{x})) v^{i}(\mathbf{x}) \right] = 0 \\ &\frac{\partial}{\partial \tau} v^{i}(\mathbf{x}) + \mathcal{H} v^{i}(\mathbf{x}) + v^{k}(\mathbf{x}) \frac{\partial}{\partial x^{k}} v^{i}(\mathbf{x}) = -\nabla_{x}^{i} \phi(\mathbf{x}) - J_{\sigma}^{i}(\mathbf{x}) - J_{1}^{i}(\mathbf{x}) \\ &\nabla^{2} \phi(\mathbf{x}) = \frac{3}{2} \Omega_{M} \mathcal{H}^{2} \delta(\mathbf{x}) \end{split}$$

Exact dynamics, including shell-crossing

Need input on the UV "sources"

$$J_{\sigma}^{i}(\mathbf{x}) \equiv \frac{1}{n(\mathbf{x})} \frac{\partial}{\partial x^{k}} (n(\mathbf{x})\sigma^{ki}(\mathbf{x}))$$
$$J_{1}^{i}(\mathbf{x}) \equiv \frac{1}{n(\mathbf{x})} \left(\langle n_{mic} \nabla^{i} \phi_{mic} \rangle(\mathbf{x}) - n(\mathbf{x}) \nabla^{i} \phi(\mathbf{x}) \right)$$

UV INFORMATION ?

Need input on the UV "sources"

$$J_{\sigma}^{i}(\mathbf{x}) \equiv \frac{1}{n(\mathbf{x})} \frac{\partial}{\partial x^{k}} (n(\mathbf{x})\sigma^{ki}(\mathbf{x}))$$
$$J_{1}^{i}(\mathbf{x}) \equiv \frac{1}{n(\mathbf{x})} \left(\langle n_{mic} \nabla^{i} \phi_{mic} \rangle(\mathbf{x}) - n(\mathbf{x}) \nabla^{i} \phi(\mathbf{x}) \right)$$

Measure them from N-body simulations

(MP, Mangano, Saviano, Viel 1108.5203, Manzotti, Peloso, MP, Viel, Villaescusa-Navarro 1407.1342)

EFToLSS: Expand in terms of long wavelength fields + power law expansion in momentum, with arbitrary coefficients to be fitted (Carrasco, Hertzberg, Senatore, 1206.2926)

Compute them from first principles. Shell-crossing! 1+1 dim attempts (Mc Quinn, White, 1502.07389; Taruya, Colombi, 1701.09088; Rampf, Frisch, 1705.08456; McDonald, Vlah, 1709.02834, Pajer, van der Woude, 1710.01736...)

UV SOURCES FROM N-BODY

$$\begin{split} \langle k^i J^i_{\sigma}(\mathbf{k}) \delta(-\mathbf{k}) \rangle'_{nb} &\propto \alpha_{nb} \frac{k^2}{k_m^2} \langle \delta(-\mathbf{k}) \delta(-\mathbf{k}) \rangle_{nb} & \text{shell crossing,} \\ \langle k^i J^i_{\sigma}(\mathbf{k}) \delta(-\mathbf{k}) \rangle'_{pt} &\propto \alpha_{pt} \frac{k^2}{k_m^2} \langle \delta(-\mathbf{k}) \delta(-\mathbf{k}) \rangle_{pt} & \text{no shell crossing,} \\ &\text{non-linearities up to PT order} \end{split}$$



1-LOOP PT + UV



no fitting on the PS

next order: 2-loop PT + $\langle J \delta \delta \rangle$ correlators

residual BAO's on the orange curves

BAO ARE LARGELY UV INDEPENDENT



CF: peak broadening, P(k): damping of the peaks

THE BAO DAMPING IS A CONTROLLED PHYSICAL EFFECT!



CONSISTENCY RELATIONS AND BAO EVOLUTION

evolution equation for P(k)

$$\langle \varphi_2(\mathbf{q})\varphi_a(\mathbf{k}-\mathbf{q})\varphi_b(-\mathbf{k})\rangle' \simeq -e^{\eta} \frac{\mathbf{k}\cdot\mathbf{q}}{q^2} P^0(q) \left(P_{ab}(k) - P_{ab}(|\mathbf{k}-\mathbf{q}|)\right) + O\left(\left(\frac{q}{k}\right)^0\right)$$

fully nonlinear!!

consistency relations from Galilean invariance (Peloso, MP 1310.7915)

CONSISTENCY RELATIONS AND BAO EVOLUTION

$$e^{\eta} I_{\mathbf{k},\mathbf{p_1},\mathbf{p_2}} \gamma_{acd} B_{bcd}(k,p_1,p_2) = -2e^{2\eta} \int^{\Lambda(k)} \frac{d^3q}{(2\pi)^3} \left(\frac{\mathbf{k} \cdot \mathbf{q}}{q^2}\right)^2 P^0(q) \left(P_{ab}(k) - P_{ab}(|\mathbf{k} - \mathbf{q}|)\right)$$
$$\simeq -2e^{2\eta} \frac{k^2}{(2\pi)^2} \int^{\Lambda(k)} dq P^0(q) \int_{-1}^1 dx \, x^2 \left(P_{ab}(k) - P_{ab}(k - qx)\right) \quad (\star)$$

sensitive to oscillatory features!

if the (nonlinear) P(k) has an oscillatory part: $P_{ab}(k) = P_{ab}^{nw}(k)(1 + A_{ab}(k)\sin(kr_{bao})) \equiv P_{ab}^{nw}(k) + P_{ab}^{w}(k)$

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hen (*) gives
$$-2e^{2\eta}k^2 \Xi(r_{bao})P^w_{ab}(k) + O\left(P^{nw''}_{ab}\right)$$

suppressed as $1/(k^2r^2_{bao})$

with
$$\Xi(r_{bao}) \equiv \frac{1}{6\pi^2} \int^{\Lambda(k)} dq P^0(q) \left(1 - j_0(q r_{bao}) + 2j_2(q r_{bao})\right)$$

Physically: large scale displacement act predominantly on P(k) (or CF) features

TWO BENEFITS

 A nonlinear evolution equation for the oscillating part of P(k) (including subdominant UV effects)

2) A tool to extract the BAO feature from a given (nonlinear) P(k)

$$R[P](k;\Delta, n) \equiv \frac{\int_{-\Delta}^{\Delta} dx \, x^{2n} \left(1 - \frac{P(k-x \, k_s)}{P(k)}\right)}{\int_{-\Delta}^{\Delta} dx \, x^{2n} \left(1 - \cos(2\pi x)\right)} \qquad k_s \equiv 2\pi/r_s, \qquad \text{BAO extractor}$$



IR+UV EVOLUTION OF BAO



BAO's match with N-body at all scales down to z=0!

IR+UV EVOLUTION OF BAO



BAO's match with N-body at all scales down to z=0!

compare with results for the broadband shape







ROBUSTNESS: REDSHIFT SPACE

shift of the BAO scale

 $\delta R[P_a](k_n;\alpha,\Delta) \equiv R[P_a](k_n/\alpha;\Delta) - R[P_{data}](k_n;\Delta)$



shift correctly reproduced by TRG equations $\frac{1}{2} \int_{0}^{\frac{1}{2}} \int_{0}^{\frac$

a

ROBUSTNESS: HALO BIAS





Extractor is only sensitive to scale-dependent bias

ROBUSTNESS: REDSHIFT SPACE+HALO BIAS



compare to "standard" analysis:

$$P_{\text{fit}}(k;\alpha) = P^{\text{smooth}}(k) \left\{ 1 + \left[O^{\text{linear}}\left(\frac{k}{\alpha}\right) - 1 \right] e^{-\frac{k^2 \Sigma_{\text{nl}}^2}{2}} \right\}$$

$$P_{\text{smooth}}(k) \equiv B_P^2 P^{0,nw}(k) + A_1 k + A_2 + \frac{A_3}{k} + \frac{A_4}{k^2} + \frac{A_5}{k^3} \qquad O^{\text{linear}}(k) \equiv \frac{P^0(k)}{P^{0,nw}(k)}$$

need: P^nw(k) + 7 nuisance parameters+ α

EXTRACIOR ON BOSS DATA (preliminary)



Notice, no TRG used here, only simple model+extractor

CONCLUSIONS

- IR and UV nonlinear effects clearly separated in TRG approach
- UV effects require external input (or nonperturbative breakthrough)
- IR effects are well understood and <u>contain physical information</u>
- BAO extractor is robust: can be efficiently computed and measured

SCALE DEPENDENCE OF SOURCES

$$\langle k^{i} J_{\sigma}^{i}(\mathbf{k}) \delta(-\mathbf{k}) \rangle_{nb}^{\prime} \propto \alpha_{nb} \frac{k^{2}}{k_{m}^{2}} \langle \delta(-\mathbf{k}) \delta(-\mathbf{k}) \rangle_{nb}$$

$$\langle k^{i} J_{\sigma}^{i}(\mathbf{k}) \delta(-\mathbf{k}) \rangle_{pt}^{\prime} \propto \alpha_{pt} \frac{k^{2}}{k_{m}^{2}} \langle \delta(-\mathbf{k}) \delta(-\mathbf{k}) \rangle_{pt}$$

