### The Formation of Vortical Motion in Cosmic Large Scale Structure

#### **Ruth Durrer**

Université de Genève Départment de Physique Théorique and Center for Astroparticle Physics





• • • • • • • • • • • •

Gravity and Cosmology Workshop, Kyoto 2018

Ruth Durrer (Université de Genève, DPT & CAP)

Cosmological Vorticiy

Kvoto, Feb 27, 2018 1 / 20

## Outline

Introduction

- 2 Perturbative Results
- Overticity from N-body simulations
- Observation of vorticity

### 5 Conclusions

э

< □ > < □ > < □ > < □ > < □ >

### Introduction: rotational motion of galaxies



Most galaxies in the Universe rotate.

The rotation axes of neighboring galaxies are correlated.

New observations find alignments of jets in radio galaxies at z = 1 out to (10-20) Mpc (A. Taylor & P. Jagannathan (2016)).

イロト イヨト イヨト イヨー

#### NGC 5457

Can these vortical motions be explained within standard ACDM? Can we learn something about cosmology by observing them?

э

< □ > < □ > < □ > < □ > < □ >

### Introduction: vorticity in cosmology

At first order in perturbation theory the velocity field of dark matter is a gradient field with vanishing vorticity.

Helmholtz's (third) theorem implies that no vorticity is generated in a perfect fluid hence within the perfect fluid approximation the velocity remains a gradient field.

This theorem is valid also in General Relativity.

Helmholtz's (third) theorem implies that no vorticity is generated in a perfect fluid hence within the perfect fluid approximation the velocity remains a gradient field.

This theorem is valid also in General Relativity.

Therefore, within Lagrangian or Eulerian, Relativistic or Newtonian perturbation theory, no vorticity is generated at any order.

Helmholtz's (third) theorem implies that no vorticity is generated in a perfect fluid hence within the perfect fluid approximation the velocity remains a gradient field.

This theorem is valid also in General Relativity.

Therefore, within Lagrangian or Eulerian, Relativistic or Newtonian perturbation theory, no vorticity is generated at any order.

This is not true for the momentum (sometimes called 'mass weighted velocity') which acquires a rotational component at second order in perturbation theory.

Helmholtz's (third) theorem implies that no vorticity is generated in a perfect fluid hence within the perfect fluid approximation the velocity remains a gradient field.

This theorem is valid also in General Relativity.

Therefore, within Lagrangian or Eulerian, Relativistic or Newtonian perturbation theory, no vorticity is generated at any order.

This is not true for the momentum (sometimes called 'mass weighted velocity') which acquires a rotational component at second order in perturbation theory.

But CDM is not a (perfect) fluid. It is a collection of free streaming particles which can be accurately described with the Vlasov equation.

This distinction is important since a fluid assigns to a given point in space a fixed value of the velocity where as the distribution in phase space allows the full velocity space in each volume element.

Helmholtz's (third) theorem implies that no vorticity is generated in a perfect fluid hence within the perfect fluid approximation the velocity remains a gradient field.

This theorem is valid also in General Relativity.

Therefore, within Lagrangian or Eulerian, Relativistic or Newtonian perturbation theory, no vorticity is generated at any order.

This is not true for the momentum (sometimes called 'mass weighted velocity') which acquires a rotational component at second order in perturbation theory.

But CDM is not a (perfect) fluid. It is a collection of free streaming particles which can be accurately described with the Vlasov equation.

This distinction is important since a fluid assigns to a given point in space a fixed value of the velocity where as the distribution in phase space allows the full velocity space in each volume element.

In a fluid description shell (orbit)-crossing is a singular process while in phase space it is regular.

N-body simulations can accommodate shell crossing without problem, they are actually nothing else than a poor woman's Vlasov equation solver.

One might think that a perturbative approach to the Vlasov equation could be successful but...



... the flow of CMD is very cold. Contrary to the case of hot dark matter, a perturbative treatment using the Vlasov equation is not adequate for CDM.

• • • • • • • • • • • •

But one can go to higher moments of the Vlasov equation, beyond the 0th and first moments which yield the continuity and Euler equations for perfect fluids.

$$\partial_t \delta + \nabla((1+\delta)\mathbf{v}) = 0,$$
  

$$\left(\partial_t + v^i \partial_i\right) v_j + \mathcal{H} v_j + \partial_j \Phi + \frac{1}{\rho} \partial_i(\rho \sigma_{ij}) = 0,$$
  

$$\left(\partial_t + v^k \partial_k\right) \sigma^{ij} + 2\mathcal{H} \sigma^{ij} + \sigma^{ik} \partial_k v^i + \sigma^{jk} \partial_k v^i = 0$$
  

$$\sigma^{ijk} = 0$$

The curl of the Euler eqn. then gives,  $\boldsymbol{\omega}=\boldsymbol{\nabla}\wedge \mathbf{v},$ 

$$rac{\partial oldsymbol{\omega}}{\partial t} + \mathcal{H}oldsymbol{\omega} - 
abla \wedge [oldsymbol{v} \wedge oldsymbol{\omega}] = -
abla \wedge \left(rac{1}{
ho}
abla (
ho\sigma)
ight)\,.$$

To lowest order in perturbation theory, the velocity dispersion take the form  $\sigma_{ij} = \frac{\sigma_0}{3} a^{-2} \delta_{ij}$ .

We have solved the vorticity equation to lowest non-vanishing order (Cusin, Tansella & RD, 2017).

$$\langle \omega^{(2)}_i(\mathbf{k},t)\omega^{(2)\,*}_j(\mathbf{k}',t)
angle = (2\pi)^3\left(\delta_{ij}-\hat{k}_i\hat{k}_j
ight)\delta(\mathbf{k}-\mathbf{k}')\mathcal{P}_\omega(k,t)\,.$$

$$P_{\omega}(k) = \frac{1}{9} \frac{\sigma_0^2 D_+(t)}{\mathcal{H}_0^2 \Omega_m} \int \frac{d^3 \mathbf{w}}{(2\pi)^3} \left( \frac{\mathbf{w} \cdot (\mathbf{k} - \mathbf{w})}{w^2 |\mathbf{k} - \mathbf{w}|^2} \right)^2 |\mathbf{w} \wedge \mathbf{k}|^2 \left[ 2\mathbf{k} \cdot \mathbf{w} - k^2 \right]^2 P_{\delta}(w) P_{\delta}(|\mathbf{k} - \mathbf{w}|)$$

 $P_{\omega}(k,t) \stackrel{k \to 0}{\to} \frac{k^4 D_+(t)}{2}$ 

2

イロト イヨト イヨト イヨト

#### Perturbative Results: vorticity power spectrum

The rotational velocity spectrum,  $P_R = k^{-2}P_{\omega}$  compared to the gradient velocity spectrum  $P_G = k^{-2}P_{\theta}$ ,  $\theta = \nabla \cdot \mathbf{v}$ .



### Vorticity from N-body simulations: Pueblas et al.



The vorticity and divergence spectra,  $P_{\omega}$  and  $P_{\theta}$ . They find a slope  $P_{\omega} \propto k^{2.5}$  and time dependence  $P_{\omega} \propto D_{+}^{7}$ .

The results shown are from a L = 256Mpc simulations with  $N = 512^3$  particles using Gadget-2 with softening length 0.04.

## Vorticity from N-body simulations: Zhu etal.



The gradient and rotational velocity spectra,  $k^3 P_G(k)$  and  $k^3 P_R(k)$ . They find a slope  $P_R \propto k^0$ . It is not clear whether these are spectra are Fourier transforms from Eulerian or Lagrangian coordinates.

The results shown are from a L = 600Mpc simulations with  $N = 1024^3$  particles on a  $512^3$  grid using CUBEP<sup>3</sup>M using multi-grid techniques to compute the displacement field and velocity (in Lagrangian coordinates).

### Vorticity from N-body simulations: gevolution

We performed N-body simulations using gevolution.

Gevolution is a relativistic PM N-body code using a weak field approximation of the metric, which computes all 6 degrees of freedom of the gravitational field (Adamek, Daverio, RD, Kunz (2016)).



Ruth Durrer (Université de Genève, DPT & CAP)

We (Jelic-Cimek, Lepori & RD, in preparation) have tested different velocity reconstruction methods which are in good agreement.



4 6 1 1 4

### Vorticity with gevolution: vorticity power spectrum

The resulting spectrum behaves as  $k^{2.5}$  on large scales.



### Vorticity with gevolution: time dependence of vorticity



The resulting spectrum behaves as  $a^{\gamma}$  on large scales with  $\gamma \simeq 5$ .

э

< □ > < □ > < □ > < □ > < □ >

Can we do better?

э

< □ > < □ > < □ > < □ > < □ >

Can we do better?

Like gradient velocity, vortical velocity leads to redshift space distortions. But these only determine the radial component of the velocity, hence cannot distinguish between gradient and vortical motion.

Can we do better?

Like gradient velocity, vortical velocity leads to redshift space distortions. But these only determine the radial component of the velocity, hence cannot distinguish between gradient and vortical motion.

Gradient RSD:  $P(k,\mu) = \left(1 + \frac{f}{b}\mu^2\right)^2 P_g(k)$ 

Rotational RSD:  $P_{\text{RSD}\,\omega}(k,\mu) = \mathcal{H}^{-2}\mu^2(1-\mu^2)P_{\omega}(k)$ 

In real space:

$$\xi(r,\mu) = \xi_0(r) + \xi_2(r)P_2(\mu) + \xi_4(r)P_4(\mu)$$

$$\xi_n(r) = \frac{\mathcal{H}^{-2}}{2\pi^2} \int P_\omega(k) j_n(kr) k^2 dk$$

### RSD monopole & quadrupole

Constraints for vorticity from structure formation:

$$P_{\omega}(k,z) = A_V k^2 D_+^7(z) rac{(k/k_*)^{n_\ell}}{[1+(k/k_*)]^{n_\ell+n_s}}$$
  
 $n_\ell = 1.3, \quad n_s = 4.3, \quad k_* = 0.7 \ h/\mathrm{Mpc}, \quad A_V \simeq 10^{-5} (\mathrm{Mpc}/h)^3$ 



(from Bonvin, RD, Koshravi, Kunz, Sawicki, 2017) Red region: Scalar signal with error for SKA at  $\bar{z} = 0.35$ . Black dashed: including vorticity with  $A_V = 5 \times 10^{-3}$ .

### Limits from the RSD hexadecapole



(from Bonvin, RD, Koshravi, Kunz, Sawicki, 2017)

Left: Red region: Scalar signal with error for SKA at  $\bar{z} = 0.35$ . Black: including vorticity with  $A_V = 3 \times 10^{-5}$  (dotted line),  $A_V = 10^{-4}$  (dashed line) and  $A_V = 10^{-3}$  (dot-dashed line). Right: non-linear scalar, linear scalar, non-linear scalar+vector.

Constraints on  $A_V$  from an SKA like survey for  $A_V$  using data from  $x \in [x_{\min}, 40 \text{Mpc}/h]$ ,  $z \in [0.1, 2], \ \Delta z = 0.1$ .

| $x_{\min}  [Mpc/h]$ | mono                | quad              | hexa                | total             |
|---------------------|---------------------|-------------------|---------------------|-------------------|
| 2                   | $3.7	imes10^{-5}$   | $4.2	imes10^{-6}$ | $8.7 	imes 10^{-7}$ | $8.7	imes10^{-7}$ |
| 10                  | $9.4	imes10^{-4}$   | $2 	imes 10^{-3}$ | $7.1	imes10^{-5}$   | $7.1	imes10^{-5}$ |
| 20                  | $7.2 	imes 10^{-2}$ | $4.6	imes10^{-2}$ | $1.6	imes10^{-3}$   | $1.6	imes10^{-3}$ |

Ruth Durrer (Université de Genève, DPT & CAP)

2

イロン イロン イヨン イヨン

Its generation by non-linear gravitational dynamics should be observable in the near future.

3

< □ > < □ > < □ > < □ > < □ >

Its generation by non-linear gravitational dynamics should be observable in the near future.

It is sensitive to non-gravitational interactions, to the nature of dark matter and to modifications of gravity.

э

Its generation by non-linear gravitational dynamics should be observable in the near future.

It is sensitive to non-gravitational interactions, to the nature of dark matter and to modifications of gravity.

Most interesting would be to observe deviations from the pure N-body expectations!

Its generation by non-linear gravitational dynamics should be observable in the near future.

It is sensitive to non-gravitational interactions, to the nature of dark matter and to modifications of gravity.

Most interesting would be to observe deviations from the pure N-body expectations!

# Thank You !

< □ > < 同 > < 回 > < 回 >

### Shell crossing of a plane wave



Ruth Durrer (Université de Genève, DPT & CAP)

Cosmological Vorticiy

Kyoto, Feb 27, 2018 19 / 20

### Resolution dependence of vorticity



• • • • • • • • • •