## Critical behavior of a domain wall collapse and Oscillon with gravity Taishi Ikeda with Chul-Moon Yoo & Vitor Cardoso

## What is Oscillon?

- · Let's consider time evolution of scalar field.
  - scalar field + double well potential
    - ► EOM

$$\frac{\partial^2}{\partial t^2} \Phi - \nabla^2 \Phi = -V'(\Phi) \qquad V(\Phi) = \frac{1}{4} (\Phi^2 - \sigma^2)^2$$



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**1.About Oscillon** 

Oscillon = extremely long-lived time dependent solution of scalar field

- typical initial data
  - Gaussian bubble  $\Phi_c : \text{value of the scalar field} at the bubble's core$   $\Phi(t = 0, r) = (\Phi_c + \sigma)e^{-r^2/R_0^2} - \sigma$ • tanh bubble  $\Phi(t = 0, r) = \frac{1}{2}[(-\sigma - \Phi_c) \tanh(r - R_0) - \sigma + \Phi_c]$



## Life time vs bubble radius

Oscillon = extremely long-lived time dependent solution of scalar field

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Oscillon naturally appears during the collapse of bubble.



**1.About Oscillon** 

## Fine structure of Oscillon

- Fine structure of lifetime (Honda et al 2002)
  - When we tune the initial parameter, resonance structure of lifetime

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- Critical behavior appears around the fine structure. (critical solution



#### **1.About Oscillon**

## What we want to do?

- We want to construct Oscillon with gravity.
  - Dose it exist ( not only in week gravity case but also strong gravity case ) ?
  - its property ?
  - It may be one of the final state or intermediate state of gravitational collapse.
- What is interesting ?
  - Can Oscillon collapse to black hole ?
  - If critical behavior appears, there are special solutions associated with critical behavior.
- We need to develop the numerical code for spherically symmetric spacetime.

#### **1.About Oscillon**

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## System and initial data

#### · EOM

$$G_{\mu\nu} = 8\pi G T_{\mu\nu}$$
  
 $\nabla^2 \Phi = V'(\Phi)$   $V(\Phi) = \frac{1}{4}(\Phi^2 - 1)^2$   $G$ : Newton constant

- initial data
  - momentary static bubble

$$\begin{cases} \Phi(t = 0, r) = -1 + 2e^{-r^2/R_0^2} & \Phi \\ \Pi(t = 0, r) = 0 & \sigma \\ \begin{cases} \tilde{\gamma}_{ij} dx^i dx^j = dr^2 + r^2 d^2 \Omega \\ K_{ij} = 0 & -\sigma \end{cases}$$

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•  $\phi$  is given from the solution of Hamiltonian constraint.

#### **2.GBSSN formulation**

## **Our numerical code**

- Our numerical code
  - It is written in C++.
  - GBSSN formulation spherically symmetric case
  - free evolution
  - time integration : iterative Crank Nicolson shceme
  - spatial derivative : central difference
  - totally second order accuracy
  - We add 2nd order numerical dissipation term in each time evolution equation.
  - We use inhomogeneous grid
  - parallel computation by using Open MP

## Kodama mass = locally conserved energy in the spherically

Kodama mass

symmetric system

- definition
  - Kodama vector

$$K^A = \epsilon^{AB} \partial_B R$$

 $ds^{2} = \mathcal{G}_{AB} dx^{A} dx^{B} + R^{2} d^{2} \Omega$ (t, r) $\epsilon_{AB} = \sqrt{-\mathcal{G}} \varepsilon_{AB} \quad \varepsilon_{tr} = 1$ 

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(vector on 2-dim manifold charted by t and r)

- →  $K^{\mu}$  (vector on 4-din manifold)
- Kodama mass

$$S^{\mu} = T^{\mu}_{\ \nu} K^{\nu}$$
$$M(t; r_0) = \int_{\Sigma} S^t \alpha \sqrt{\gamma} dx^3$$



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#### **2.GBSSN formulation**

## $ds^2 = \mathcal{G}_{AB}dx^A dx^B + r^2 d^2 \Omega$

- proof

$$= -\alpha^{2}(t,r)dt^{2} + a^{2}(t,r)dr^{2} + r^{2}d^{2}\Omega$$

• component of  $K^A = \epsilon^{AB} \partial_B r$ 

$$K^t = \epsilon^{tr} \partial_r r = -\frac{1}{\alpha a}, \ K^r = 0$$

• component of 
$$S^{\mu} = T^{\mu}_{\ \nu} K^{\nu}$$

component of 
$$S^{\mu} = T^{\mu}_{\ \nu} K^{\nu}$$
  

$$\begin{cases} S^{t} = \frac{1}{\alpha^{3}a} T_{tt} = \frac{1}{\alpha^{3}a} G_{tt} = \frac{1}{\alpha^{3}a} (\frac{2\alpha^{2}\partial_{r}a}{ra^{3}} + \frac{\alpha^{2}}{r^{2}} - \frac{\alpha^{2}}{r^{2}a^{2}}) \\ S^{r} = \frac{1}{\alpha a^{3}} T_{rt} = \frac{1}{\alpha a^{3}} G_{rt} = -\frac{2\partial_{t}a}{r\alpha a^{4}} \\ \partial_{\mu} (\sqrt{-g} S^{\mu}) = 0 \end{cases}$$

$$\epsilon_{AB} = \sqrt{-\mathcal{G}}\varepsilon_{AB}$$
$$\varepsilon_{tr} = 1$$
$$\epsilon_{tr} = \alpha a$$

**2.GBSSN** formulation

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#### Kodama mass

#### **2.GBSSN formulation**

### Kodama mass

- conservation law of Kodama mass

$$\nabla_{\mu}S^{\mu} = \frac{1}{\sqrt{-g}}\partial_{\mu}(\sqrt{-g}S^{\mu}) = 0$$

$$M(t = 0; r_0) + P(t = 0, t; r_0) = M(t; r_0)$$

$$\begin{cases} M(t;r_0) = \int_{\Sigma} S^t \alpha \sqrt{\gamma} dx^3 \\ P(t=0,t;r_0) = \int_0^t dt \int_{\partial \Sigma} d^2 x \ n_\mu S^\mu \end{cases}$$

 We define the lifetime of Oscillon from Kodama mass in the case of Oscillon with gravity.





# Oscillon with gravity (preliminary result) 12

Some examples of numerical simulation



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#### 3.Oscillon with gravity

# Oscillon with gravity (preliminary result) 13



**3.Oscillon with gravity** 

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## Life time of Oscillon with gravity (preliminary result)

initial bubble radius vs life time of Oscillon with gravity



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**3.Oscillon with gravity** 

# Life time of Oscillon with gravity (preliminary result)



**3.Oscillon with gravity** 

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## Summary and future work

- Summary
  - Oscillon = long lifetime localized solution of scalar filed on Minkwski background
  - It is the product of nonlinear effect of scalar field.
  - We construct the Oscillon with gravity.
  - We get sign of its fine structure in the life time.
- Future work
  - Can Oscillon collapse to black hole ?
  - Can we check the critical behavior ?
  - What is the properties of the critical solution ?

