Constraints on Einsteinaether theory after GW170817

Shinji Mukohyama (YITP, Kyoto U)

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Implication of GW170817 on gravity theories @ late time

- $|c_{gw}^2 c_{\gamma}^2| < 10^{-15}$
- Horndeski theoy (scalar-tensor theory with 2nd-order eom): Among 4 free functions, $G_4(\phi, X) \& G_5(\phi, X)$ are strongly constrained. Still $G_2(\phi, X) \& G_3(\phi, X)$ are free.
- Generalized Proca theory (vector-tensor theory): Among 6 (or more) free functions, G₄(X) & G₅(X) are strongly constrained. Still G₂(X,F,Y,U), G₃(X), G₆(X), g₅(X) are free.
- Horava-Lifshitz theory (renormalizable quantum gravity): The coefficient of R⁽³⁾ is strongly constrained \rightarrow IR fixed point with $c_{gw}^2 = c_{\gamma}^2$? How to speed up the RG flow?
- Ghost condensation (simplest Higgs phase of gravity): No additional constraint
- Massive gravity (simplest modification of GR): Upper bound on graviton mass ≈ 10⁻²²eV Much weaker than the requirement from acceleration
- c.f. "All" gravity theories (including general relativity): The cosmological constant is strongly constrained ≈ 10⁻¹²⁰.

Einstein-aether theory

- LV gravity: GR + timelike unit vector u^μ
 → Useful for test of GR
- Low-E limit of non-projectable HL gravity

 predictions of a quantum gravity theory
- 4 parameters (c₁, c₂, c₃, c₄)

$$\begin{split} M^{\alpha\beta}_{\ \mu\nu} &= c_1 g^{\alpha\beta} g_{\mu\nu} + c_2 \delta^{\alpha}_{\mu} \delta^{\beta}_{\nu} + c_3 \delta^{\alpha}_{\nu} \delta^{\beta}_{\mu} - c_4 u^{\alpha} u^{\beta} g_{\mu\nu} \\ \mathcal{L}_{\&} &\equiv -M^{\alpha\beta}_{\ \mu\nu} \left(D_{\alpha} u^{\mu} \right) \left(D_{\beta} u^{\nu} \right) + \lambda \left(g_{\alpha\beta} u^{\alpha} u^{\beta} + 1 \right) \\ S_{\&} &= \frac{1}{16\pi G_{\&}} \int \sqrt{-g} \, d^4 x \Big[R(g_{\mu\nu}) + \mathcal{L}_{\&} \left(g_{\mu\nu}, u^{\lambda} \right) \Big] \\ S &= S_{\&} + \int \sqrt{-g} \, d^4 x \Big[\mathcal{L}_m \left(g_{\mu\nu}, \psi \right) \Big] \end{split}$$

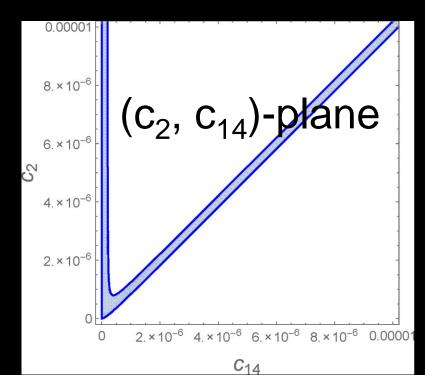
Constraints

- Stability $q_{S,V,T} > 0$ $c_{S,V,T}^2 \ge 0$
- Gravi-Cerenkov $c_{S,V,T}^2\gtrsim 1$
- GW170817 $-3 \times 10^{-15} < c_T 1 < 7 \times 10^{-16}$
- **BBN** $\left|\frac{G_{\cos}}{G_N} 1\right| \lesssim \frac{1}{8}$
- **PPN** $|\alpha_1| \le 10^{-4}$ $|\alpha_2| \le 10^{-7}$
- Pulsars $|\hat{\alpha}_1| \le 10^{-5}$ $|\hat{\alpha}_2| \le 10^{-9}$

Some formulae $c_{ij} \equiv c_i + \overline{c_j} \quad c_{ijk} = c_i + \overline{c_j} + \overline{c_k}$ Stability & Gravi-Cerenkov & GW170817 $q_{S} = \frac{(1-c_{13})(2+c_{13}+3c_{2})}{c_{123}} \qquad q_{V} = c_{14} \qquad q_{T} = 1-c_{13}$ $c_{S}^{2} = \frac{c_{123}(2-c_{14})}{c_{14}(1-c_{13})(2+c_{13}+3c_{2})} \qquad c_{V}^{2} = \frac{2c_{1}-c_{13}(2c_{1}-c_{13})}{2c_{14}(1-c_{13})} \qquad c_{T}^{2} = \frac{1}{1-c_{13}}$ BBN $G_N = \frac{G_{\varpi}}{1 - \frac{1}{2}c_{14}} \qquad G_{\rm COS} = \frac{G_{\varpi}}{1 + \frac{1}{2}(c_{13} + 3c_2)}$ $\alpha_1 = -\frac{8(c_3^2 + c_1c_4)}{2c_1 - c_1^2 + c_2^2}$ PPN $\alpha_2 = \frac{1}{2}\alpha_1 - \frac{(c_1 + 2c_3 - c_4)(2c_1 + 3c_2 + c_3 + c_4)}{c_{123}(2 - c_{14})}$ • Pulsars $\hat{\alpha}_1 = \alpha_1 + \frac{c_-(8+\alpha_1)\sigma_{\infty}}{2c_1}$ $\hat{\alpha}_2 = \alpha_2 + \frac{\hat{\alpha}_1 - \alpha_1}{2} - \frac{(c_{14} - 2)(\alpha_1 - 2\alpha_2)\sigma_{x}}{2(c_{14} - 2c_{13})}$

Constraints on (c_1, c_2, c_3, c_4) $c_{ij} \equiv c_i + c_j$ $c_{ijk} = c_i + c_j + c_k$

- Strongest constraint \leftarrow GW170817 $|c_{13}| < 10^{-15}$
- (c₁, c₁₄)-plane $0 < c_{14} \le 2.5 \times 10^{-5}$ $c_{14} \lesssim c_1$
- (c₂, c₁₄)-plane $c_{14} \le 2.5 \times 10^{-5}$ $c_2 \lesssim 0.095$
- Sensitivity σ_{lpha} not known in those range



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