

非可換ソリトンと保存量

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Ref: MH, “Commuting Flows and Conservation Laws for NC Lax Hierarchies,” [hep-th/0311206]

Cf : MH, “NC Solitons and Integrable Systems, Proc. of NCGP04 [hep-th/041mnnn]

NC Ward's observation: Almost all
NC integrable equations are
reductions of the NC ASDYM eqs.

MH&K.Toda, PLA316(03)77[hep-th/0211148]

NC ASDYM eq.

Successful

↓ Reductions

Reductions

NC KP eq. NC BCS eq.

NC KdV eq. NC Boussinesq eq.

NC NLS eq. NC mKdV eq.

NC sine-Gordon eq. NC Burgers eq. ...

(Almost all !?)

Successful?

A general framework is needed

2. NC Sato's Theory

- Sato's Theory : one of the most beautiful theory of solitons
 - Based on the existence of hierarchies and tau-functions
- Sato's theory reveals essential aspects of solitons:
 - Construction of exact solutions
 - Structure of solution spaces
 - Infinite conserved quantities
 - Hidden infinite-dim. symmetry

Let's discuss NC extension of Sato's theory

Derivation of soliton equations

- Give a Lax operator which is a pseudo-differential operator

$$L := \partial_x + u_2 \partial_x^{-1} + u_3 \partial_x^{-2} + u_4 \partial_x^{-3} + \dots$$

$$u_k = u_k(x^1, x^2, x^3, \dots)$$

- Define a differential operator

$$B_m := (L * \dots * L)_{\geq 0}$$

m times

Noncommutativity
is introduced here:

$$[x^i, x^j] = i\theta^{ij}$$

- Introduce NC (KP) hierarchy equations:

$$\frac{\partial L}{\partial x^m} = [B_m, L]_*$$

$$\begin{aligned} &\partial_m u_2 \partial_x^{-1} + \\ &\partial_m u_3 \partial_x^{-2} + \\ &\partial_m u_4 \partial_x^{-3} + \dots \end{aligned}$$

$$\begin{aligned} &F_{m2}(u) \partial_x^{-1} + \\ &F_{m3}(u) \partial_x^{-2} + \\ &F_{m4}(u) \partial_x^{-3} + \dots \end{aligned}$$



Each coefficient yields
a differential equation.

Star product:

$$f(x) * g(x) := f(x) \exp\left(\frac{i}{2} \theta^{ij} \vec{\partial}_i \vec{\partial}_j\right) g(x)$$

which realizes the
noncommutativity

NC KP hierarchy equations

Infinite kind of fields u_k can be represented in terms of $u_2 \equiv u$

$$\frac{\partial u}{\partial x^3} = \frac{1}{4} u_{xxx} + \frac{3}{4} (u_x * u + u * u_x) + \frac{3}{4} \partial_x^{-1} u_{yy} + \frac{3}{4} [u, \partial_x^{-1} u_{yy}]_*$$

$$u_x := \frac{\partial u}{\partial x} \quad \partial_x^{-1} := \int^x dx' \quad \text{etc.}$$

(2+1)-dim. NC KP equation

$$u = u(x^1, x^2, x^3, \dots)$$

$$\begin{array}{ccc} \updownarrow & \updownarrow & \updownarrow \\ x & y & t \end{array}$$

$$\frac{\partial u}{\partial x^4} = f(u, u_x, u_y, \dots)$$

$$\frac{\partial u}{\partial x^5} = g(u, u_x, u_y, \dots)$$

$$\frac{\partial u}{\partial x^6} = h(u, u_x, u_y, \dots)$$

...

NC KP hierarchy equations

^{reductions}
(KP hierarchy) \rightarrow (various hierarchies.)

- No-reduction \rightarrow NC KP $(x, y, t) = (x^1, x^2, x^3)$
- 2-reduction \rightarrow NC KdV $(x, t) = (x^1, x^3)$
- 3-reduction \rightarrow NC Boussinesq $(x, t) = (x^1, x^2)$
- 4-reduction \rightarrow NC Coupled KdV ...
- 5-reduction \rightarrow ...
- 3-reduction of BKP \rightarrow NC Sawada-Kotera
- 2-reduction of mKP \rightarrow NC mKdV
- Special 1-reduction of mKP \rightarrow NC Burgers
- ...



Noncommutativity should be introduced into space-time coords

4. Conservation Laws

- Conservation laws: $\partial_t \sigma = \partial_i J^i$ σ : Conserved density
time \nearrow ∂_t σ \leftarrow ∂_i J^i space

Then $Q := \int_{space} dx \sigma$ is a conserved quantity.

$$\because \partial_t Q = \int_{space} dx \partial_t \sigma = \int_{spatial} dS_i J^i = 0$$

infinity

Conservation laws for the hierarchies

$$\partial_m res_{-1} L^n = \partial_x J + \theta^{ij} \partial_j \Xi_i$$

time \nearrow ∂_m res_{-1} L^n \leftarrow ∂_x J \leftarrow ∂_j Ξ_i
space

I have found the explicit form !

$res_{-r} L^n$: coefficient
of ∂_x^{-r} in L^n

Noncommutativity should be introduced
in space-time directions only. \rightarrow

$$t \equiv x^m$$

∂_j should be space or time derivative
 \rightarrow ordinary conservation laws !

Hot Results

Infinite conserved densities for NC hierarchy eqs. ($n=1,2,\dots$)

$$\sigma = \text{res}_{-1} L^n + \theta^{im} \sum_{k=0}^{m-1} \sum_{l=0}^k (-1)^{k-l} \binom{k}{l} (\text{res}_{-(l+1)} L^n) \diamond (\partial_i \partial_x^{k-l} \text{res}_k L^m)$$

$t \equiv x^m$ $\text{res}_r L^n$: coefficient of ∂_x^r in L^n

\diamond : Strachan's product (commutative and non-associative)

$$f(x) \diamond g(x) := f(x) \left(\sum_{s=0}^{\infty} \frac{(-1)^s}{(2s+1)!} \left(\frac{1}{2} \theta^{ij} \vec{\partial}_i \vec{\partial}_j \right)^{2s} \right) g(x)$$

MH, [hep-th/0311206]

We can calculate the explicit forms of conserved densities for wide class of NC soliton equations.

- Space-Space noncommutativity:

NC deformation is slight: $\sigma = \text{res}_{-1} L^n$

- Space-time noncommutativity

NC deformation is drastical:

– Example: NC KP and KdV equations $([t, x] = i\theta)$

$$\sigma = \text{res}_{-1} L^n - 3\theta((\text{res}_{-1} L^n) \diamond u'_3 + (\text{res}_{-2} L^n) \diamond u'_2)$$

meaningful ?

4. Conclusion and Discussion

- We proved the existence of infinite conserved quantities for wide class of NC Lax hierarchies.
- We gave the infinite conserved densities explicitly from the viewpoint of Sato's theory, which suggests that infinite-dim. symmetry would be hidden in the NC (soliton) equations.
- The results show that NC soliton eqs. have very special properties though they include non-linear terms and infinite (time-!) derivatives.
- Of course, there are still many things to be seen in order to confirm the integrability.

Further directions

- Completion of NC Sato's theory
 - Hirota's bilinearization and tau-functions $u = \partial_x^2 \log \tau$
→ hidden symmetry (deformed affine Lie algebras?)
Cf. Dimakis&Mueller-Hoissen, Wang&Wadati...
 - Geometrical descriptions from NC extension of the theories of Krichever, Mulase and Segal-Wilson and so on.
- Confirmation of NC Ward's conjecture
 - NC twistor theory
e.g. Kapustin&Kuznetsov&Orlov, Hannabuss, Hannover group,...
 - D-brane interpretations → physical meanings
- Foundation of Hamiltonian formalism with space-time noncommutativity
 - Initial value problems, Liouville's theorem, Noether's thm,...