

Noncommutative Solitons and Integrable Systems

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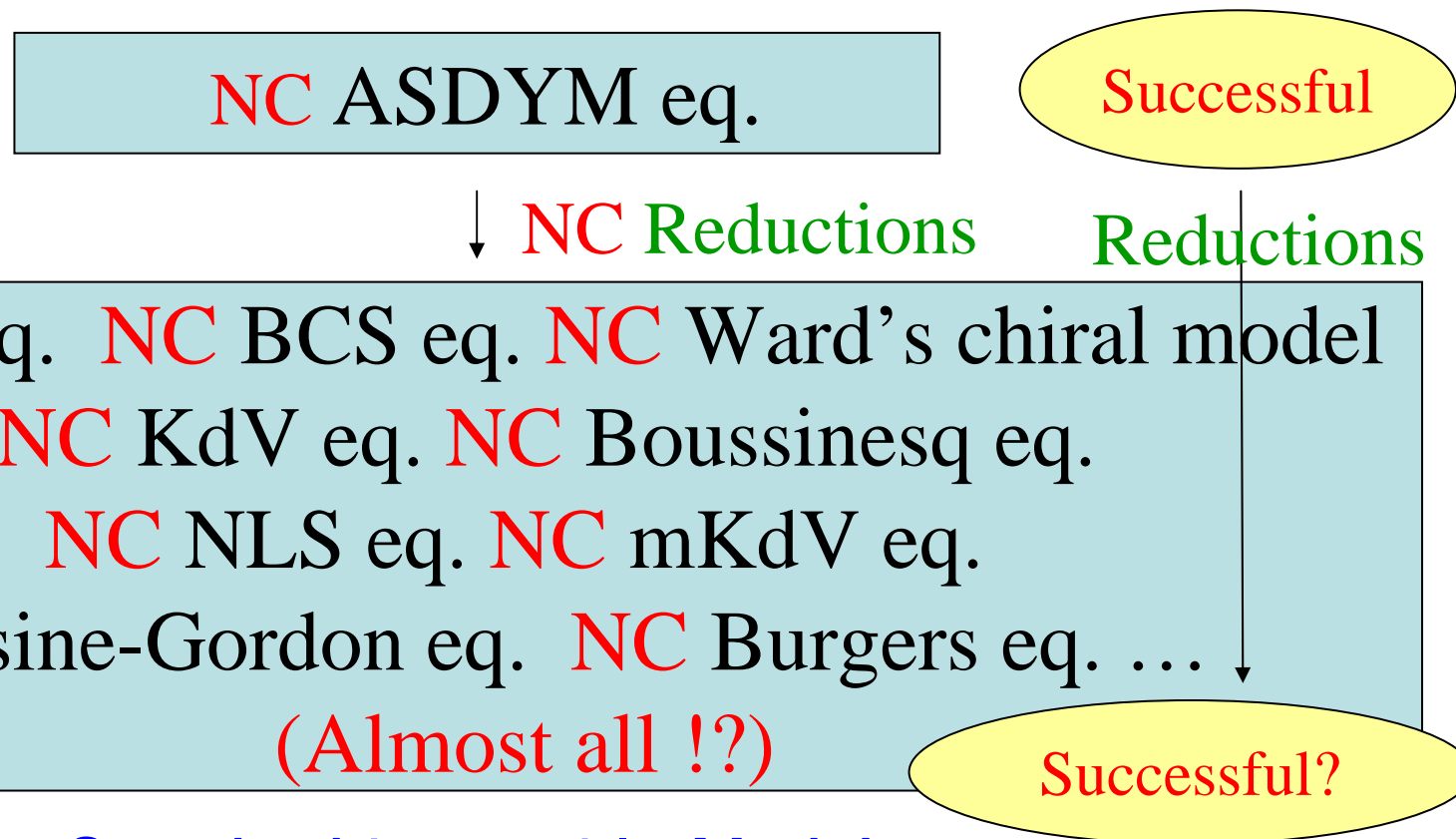
Ref: MH, “NC solitons and integrable systems,”
(to appear soon) [[hep-th/0504001](#)].

cf. MH, “Commuting flows and conservation laws
for NC Lax hierarchies,” [[hep-th/0311206](#)]

[J.Math.Phys.46 \(2005\)](#) to be published on April 1.

NC Ward's observation: Almost all
NC integrable equations are
reductions of the NC ASDYM eqs.

MH&K.Toda, PLA316(03)77[hep-th/0211148], Bad Honnef 03



``Beyond the Standard integrable Models

Program of NC extension of soliton theories

- (i) Confirmation of NC Ward's conjecture
 - NC twistor theory → geometrical origin
 - D-brane interpretations → applications to physics
 - (ii) Completion of NC Sato's theory
 - Existence of "hierarchies" → various soliton eqs.
 - Existence of infinite conserved quantities
 - infinite-dim. hidden symmetry
 - Construction of multi-soliton solutions
 - Theory of tau-functions → structure of the solution spaces and the symmetry
- (i),(ii) → complete understanding of the NC soliton theories

NC Lax hierarchy (m=1,2, ...)

$$\frac{\partial L}{\partial x^m} = [B_m, L]_*$$

$$L := \partial_x + u_2 \partial_x^{-1} + u_3 \partial_x^{-2} + u_4 \partial_x^{-3} + \dots, \quad B_m := (L * \dots * L)_{\geq 0}$$

This yields **wide class of NC soliton eqs.**

- No-reduction \rightarrow NC KP $u_t = \frac{1}{4}u_{xxx} + \frac{3}{4}(u * u)_x + \frac{3}{4}\partial_x^{-1}u_{yy} + \frac{3}{4}[u, \partial_x^{-1}u_{yy}]_*$
- 2-reduction \rightarrow NC KdV $u_t = \frac{1}{4}u_{xxx} + \frac{3}{4}(u_x * u + u * u_x)$
- 3-reduction \rightarrow NC Boussinesq
- 4-reduction \rightarrow NC Coupled KdV
- 2-reduction of mKP \rightarrow NC mKdV
- Special 1-reduction of mKP \rightarrow NC Burgers
- Matrix generalization \rightarrow NC AKNS, NLS, ...

We've solved
"hierarchy problem"

Infinite conserved densities for the NC soliton eqs. ($n=1,2,\dots$)

$$\sigma_n = \text{coef}_{-1} L^n + \theta^{im} \sum_{k=0}^{m-1} \sum_{l=0}^k \binom{k}{l} (\partial_x^{k-l} \text{coef}_{-(l+1)} L^n) \diamond (\partial_i \text{coef}_k L^m)$$

$t \equiv x^m$ $\text{coef}_r L^n$: coefficient of ∂_x^r in L^n

\diamond : Strachan's product (commutative and non-associative)

$$f(x) \diamond g(x) := f(x) \left(\sum_{s=0}^{\infty} \frac{(-1)^s}{(2s+1)!} \left(\frac{1}{2} \theta^{ij} \vec{\partial}_i \vec{\partial}_j \right)^{2s} \right) g(x)$$

MH, JMP46
[hep-th/0311206]

This suggests infinite-dimensional symmetries would be hidden.

Exact N-soliton solutions of the NC KP hierarchy

$L = \Phi \partial_x \Phi^{-1}$ solves the NC KP hierarchy !

$\Phi f := \left| W(y_1, \dots, y_N, f) \right|_{N+1, N+1}$ quasi-determinant of Wronski matrix

$$y_i = \exp \xi(x, \alpha_i) + a_i \exp \xi(x, \beta_i)$$

Etingof-Gelfand-Retakh

$$\xi(x, \alpha) = x_1 \alpha + x_2 \alpha^2 + x_3 \alpha^3 + \dots \quad [\text{q-alg/9701008}]$$

The exact solutions are actually N-soliton solutions !

Noncommutativity might affect the phase shift by $\theta^{ij} \omega_i k_j$

$\therefore \exp(\omega_i t - k_i x) * \exp(\omega_j t - k_j x)$ [MH, work in progress]

$$\cong \exp([t, x]_* \omega_{[i} k_{j]}) \exp((\omega_i + \omega_j)t - (k_i + k_j)x)$$

Exactly solvable!

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 - Existence of ``hierarchies'' → Solved!
 - Existence of infinite conserved quantities → infinite-dim. hidden symmetry Successful
 - Construction of multi-soliton solutions Successful
 - Theory of tau-functions → description of the symmetry and the soliton solutions

NC Ward's conjecture (NC NLS eq.)

- Reduced ASDYM eq.: $x^\mu \rightarrow (t, x)$

Legare,
[hep-th/0012077]

(i) $B' = 0$

(ii) $C' - \dot{A} + [A, C]_* = 0$

(iii) $A' - \dot{B} + [C, B]_* = 0$

A, B, C: 2 times 2
matrices (gauge fields)



Further
Reduction:

$$A = \begin{pmatrix} 0 & q \\ -\bar{q} & 0 \end{pmatrix}, B = \frac{i}{2} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}, C = i \begin{pmatrix} q * \bar{q} & \bar{q}' \\ q' & -\bar{q} * q \end{pmatrix}$$

$$(ii) \Rightarrow \begin{pmatrix} 0 & i\dot{q} - q'' - 2q * \bar{q} * q \\ i\dot{\bar{q}} + \bar{q}'' - 2\bar{q} * q * \bar{q} & 0 \end{pmatrix} = 0$$

NOT traceless

$$i\dot{q} = q'' + 2q * \bar{q} * q \quad : \text{NC NLS eq. !!!}$$

Note: $A, B, C \in u(2) \xrightarrow{\theta \rightarrow 0} su(2)$ **U(1) part is necessary !**

NC Ward's conjecture (NC Burgers eq.)

- Reduced ASDYM eq.: $x^\mu \rightarrow (t, x)$ MH & K.Toda, JPA36
[hep-th/0301213 v2]

$$(i) \quad \dot{A} + \underline{[B, A]}_* = 0$$

$$(ii) \quad \dot{C} - B' + \underline{[B, C]}_* = 0, \quad A, B, C \in \underline{u(1)}$$

$$\cong u(\infty)$$

should remain



Further

Reduction: $A = 0, B = u' - u^2, C = u$

$$(ii) \Rightarrow \dot{u} = u'' + 2u' * u \quad \xrightarrow[u = \tau^{-1} * \tau']{\quad} \quad \dot{\tau} = \tau''$$

: NC Burgers eq. !!!

Linear eq.
(Integrable !!!)

Note: Without the commutators $[,]$, (ii) yields:

$$\dot{u} = u'' + \underline{u' * u + u * u'} \quad : \text{neither linearizable nor Lax form}$$

Symmetric

NC Ward's conjecture (NC KdV eq.)

- **Reduced ASDYM eq.:** $x^\mu \rightarrow (t, x)$ MH, (HOT result)
[hep-th/0504001]

(i) $[A, B]_* = 0$

(ii) $A' - C' + [A, C]_* + [D, B]_* = 0$

(iii) $\dot{C} - D' + [C, D]_* = 0$

A, B, C: 2 times 2
matrices (gauge fields)

↓ Further Reduction (different type of Mason-Sparling!):

$$A = \frac{1}{2} \begin{pmatrix} 0 & 0 \\ -u & 0 \end{pmatrix}, B = \begin{pmatrix} 0 & 0 \\ -1 & 0 \end{pmatrix}, C = \begin{pmatrix} 0 & 1 \\ -u & 0 \end{pmatrix}, D = \frac{1}{4} \begin{pmatrix} -u' & 2u \\ -u'' - 2u^2 & u' \end{pmatrix}$$

(iii) $\Rightarrow \frac{1}{4} \begin{pmatrix} 0 & 0 \\ 4\dot{u} + u''' + 3(u' * u + u * u') & 0 \end{pmatrix} = 0$

$$\dot{u} + \frac{1}{4} u''' + \frac{3}{4} (u' * u + u * u') = 0 \quad : \text{ NC KdV eq. !!!}$$

Note: $A, B, C \in sl(2, R)$

Bad news for D-brane picture?
Good news for twistor picture !

Program of NC extension of soliton theories

- Confirmation of NC Ward's conjecture Going well
 - NC twistor theory → geometrical origin
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Work in progress | II continue studying at Oxford from Oct. 05
- Completion of NC Sato's theory
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