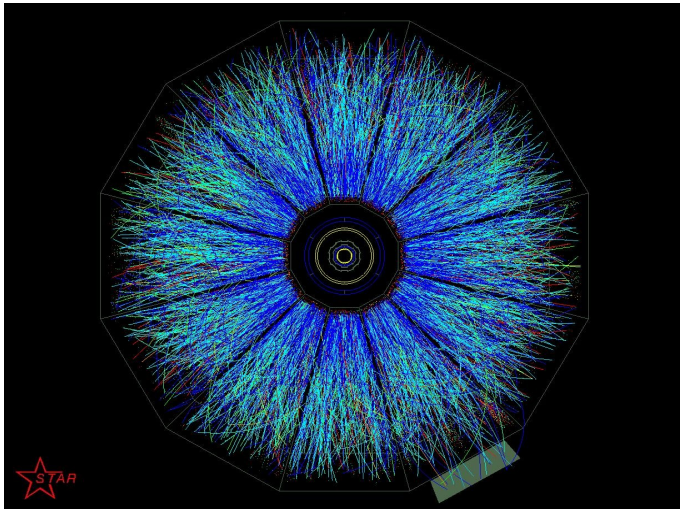


# The Strongly-Coupled Semi Quark-Gluon Plasma

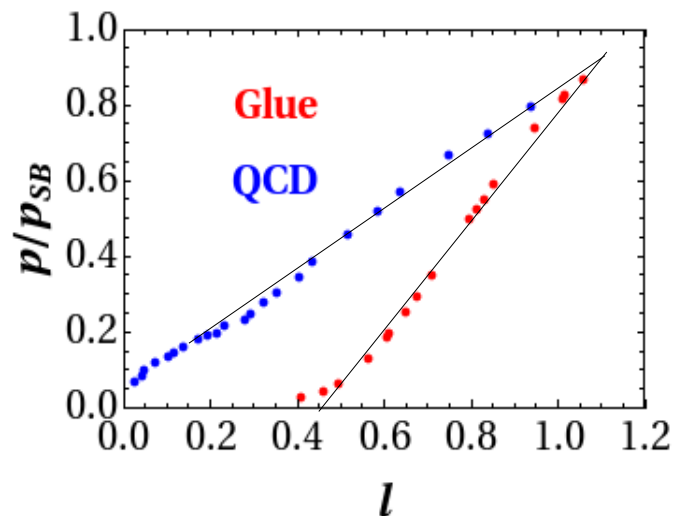


Jorge Noronha

- The semi QGP corresponds to the QCD state of matter where the Polyakov loop changes significantly with  $T$ , i.e., between  $0.8 - 2.5 T_c$  (Pisarski, Hidaka, 2008)

## THE SEMI QGP IS NOT APPROXIMATELY CONFORMAL !!!!!

- Physically this can be seen as a state where ionization of color is not complete.
- According to the lattice, most likely **the QGP created in heavy ions is a semi-QGP.**
- Lattice data suggests that it may be possible to find an effective description of the semi QGP where



Note that  $p/p_{SB} \sim \ell$

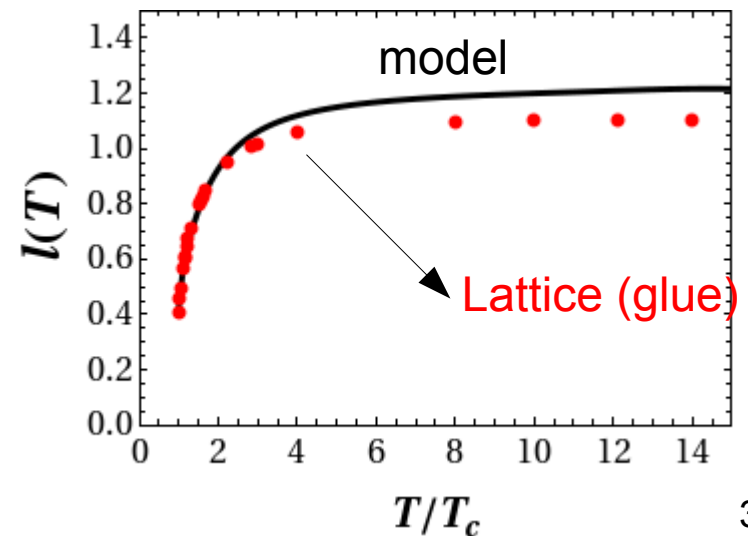
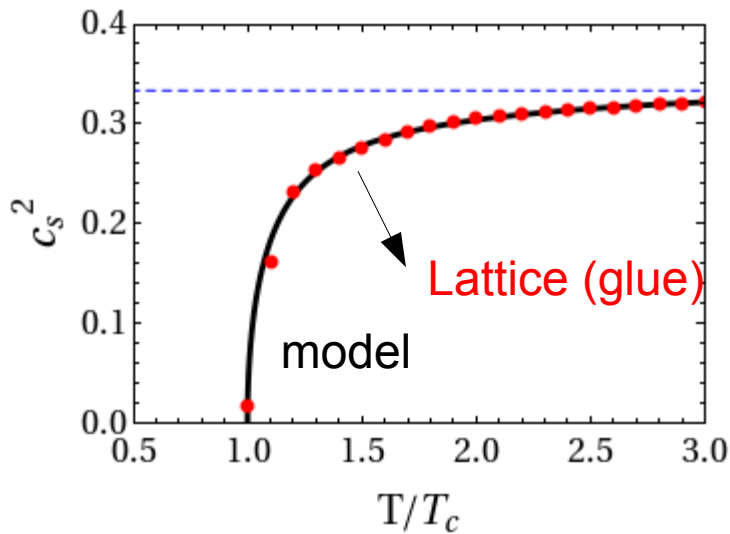
for glue, glue + quarks (even in the “hadronic phase”) !!!

How do we model this semi QGP ???

The **main idea** of this talk is to show you in a few slides that the semi QGP **can be successfully modeled using holography**

In fact, a “simple” 5d gravity dual involving the metric and a scalar field with dynamics given by

$$V(\phi) = -12(1 + a\phi^2)^{1/4} \cosh \sqrt{\frac{2}{3}}\phi + \dots$$

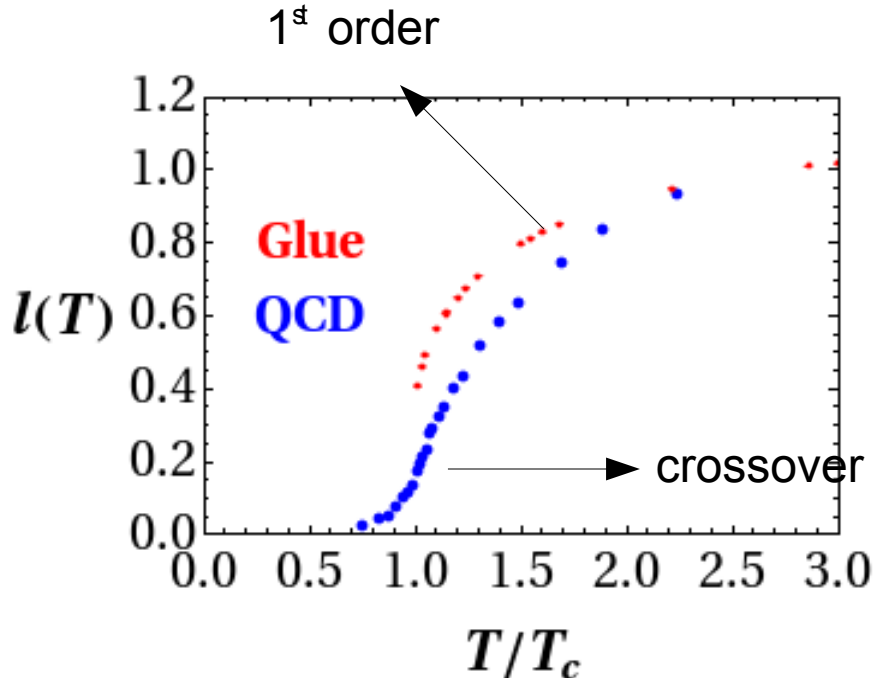


What does the lattice say about the semi Quark-Gluon Plasma?

Renormalized Polyakov loop in the fundamental representation

$$l \sim e^{-\mathcal{F}_Q/T}$$

Data from A. Bazavov et al (MILC Collaboration), 2009.  
 Gupta, Huebner, Kaczmarek, 2008.



In QCD, near  $T_c$  one does not have well defined hadrons or quarks and gluons ... (at least not in the way we have studied them before)

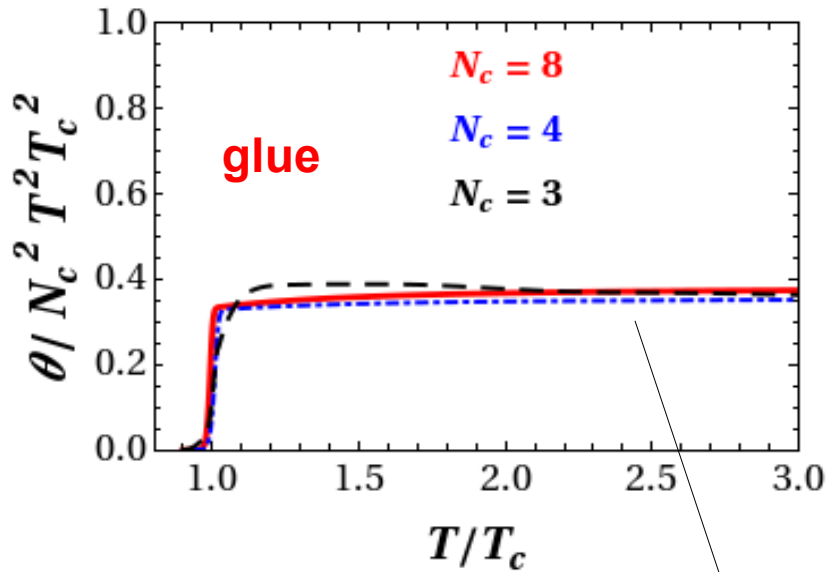
**Deconfinement as ionization of color charge**

Pisarski, Hidaka, 2008.

What does the lattice say about the semi Quark-Gluon Plasma?

Trace anomaly

$$\theta = \varepsilon - 3p$$



Plot made using the data from M. Panero, PRL 2009

Fuzzy bag?

Power like behavior near  $T_c$

Pisarski, 2006

Megias, Arriola, Salcedo, 2006

$$\frac{\varepsilon - 3p}{T^4} \sim c \frac{T_c^2}{T^2}$$

Qualitatively different than usual pQCD calculations (no logarithms) ...

# Holographic Model of Pure Glue at Large $N_c$

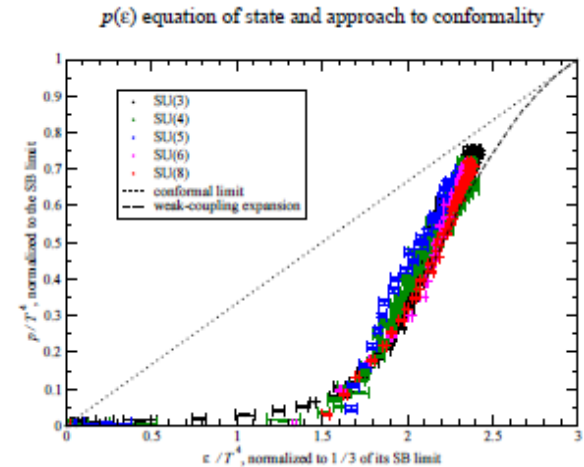
Lattice shows that **conformal invariance is badly violated near  $T_c$**

Assumption:  $1/3 \sim 1/N_c \sim 0$

From the gauge/string duality dictionary

$$\phi \sim \text{Tr} F_{\mu\nu}^2$$

Bulk            In the gauge theory at the boundary



M. Panero, arXiv:0907.3719 [hep-lat]

One should look for gravity duals with a nontrivial scalar field in the bulk !!!

Minimal extension of the good and old gravity setup (**a bottom-up approach**)

$$S = -\frac{1}{2\kappa_5^2} \int d^5x \sqrt{-G} \left[ \mathcal{R} - \frac{(\partial\phi)^2}{2} - V(\phi) \right]$$

Scalar potential

$$V(\phi)$$

Nontrivial fields in the 5d bulk:  $G_{\mu\nu}, \phi$

$$k_5^2 = c / N_c^2$$

Dual to a relevant deformation of a 4d CFT  $\mathcal{L}_{CFT} + \Lambda_\phi^{4-\Delta} \mathcal{O}_\phi$

Here  $\Lambda_\phi$  is the energy scale of the deformation and  $\Delta$  is the dimension of  $\mathcal{O}_\phi$  in the boundary, which is dual to  $\phi$  in the bulk.

An important property of these theories is that the thermodynamic quantities exhibit a power-like expansion in terms of

$$\left(\frac{\Lambda\phi}{T}\right)^{2(4-\Delta)}$$

Cherman, Nellore 2009  
Hohler, Stephanov, 2009

Thus, the power-like behavior seen on the data should be captured by these models for a convenient choice of parameters.

Here I will require that there is linear confinement below  $T_c$ .

A choice for the scalar potential that describes the lattice data is

$$V(\phi) = -12(1 + a\phi^2)^{1/4} \cosh \sqrt{\frac{2}{3}}\phi + b_2\phi^2 + b_4\phi^4 + b_6\phi^6$$

$$a = 1$$

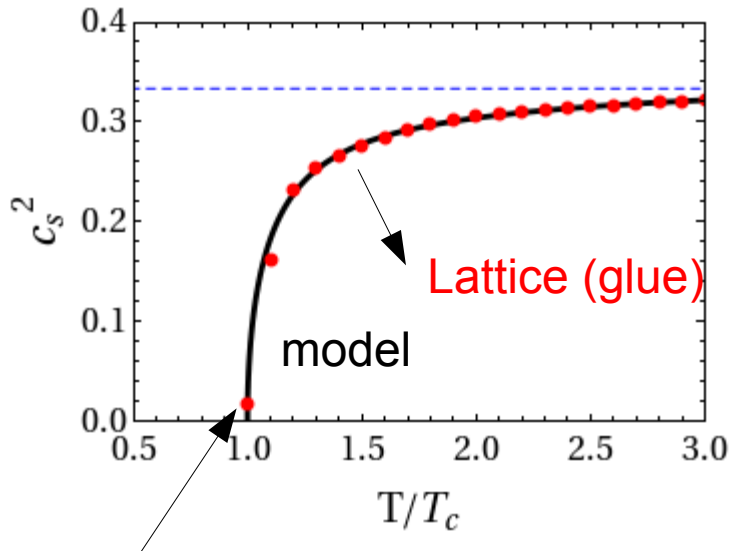
$$b_4 = 0.4$$

$$b_2 = 5$$

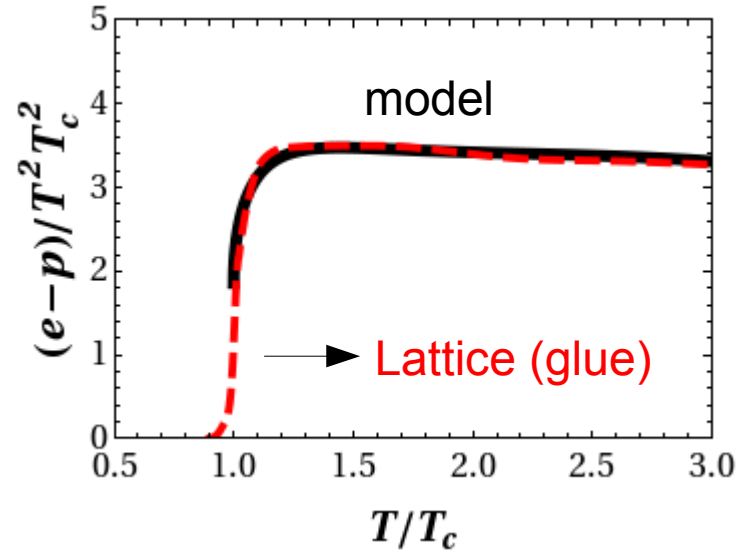
$$b_6 = 0.0098$$

$$\Delta = 2$$

This choice gives an amazing match to the lattice data ...



1<sup>st</sup> order transition



What about the Polyakov loop ???

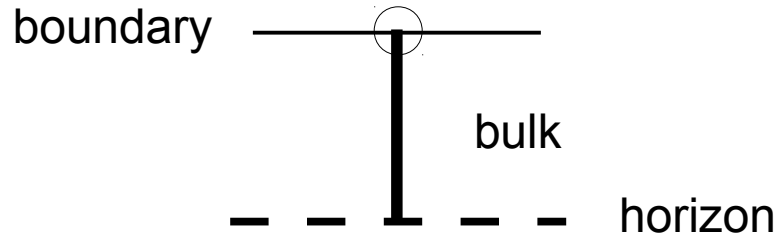
**(the ionization of color)**

Infinitely massive heavy quark  $\sim$  fundamental string in the bulk

Assuming

$N_c \gg 1$

$R^2/\alpha' \gg 1$



Maldacena, 1998  
 Rey et al. 1998  
 Brandhuber et al 1998

$$|\langle \ell(T) \rangle| \sim e^{-S_{NG}(\mathcal{D})}$$

$$S_{NG}(\mathcal{D}) = \frac{1}{2\pi\alpha'} \int_{\mathcal{D}} d^2\sigma q(\phi) \sqrt{\det h^{ab}}$$

Nambu-Goto action for the string in the bulk

$$h^{ab} = G_{\mu\nu} \partial^a X^\mu \partial^b X^\nu$$

Induced metric on the worldsheet

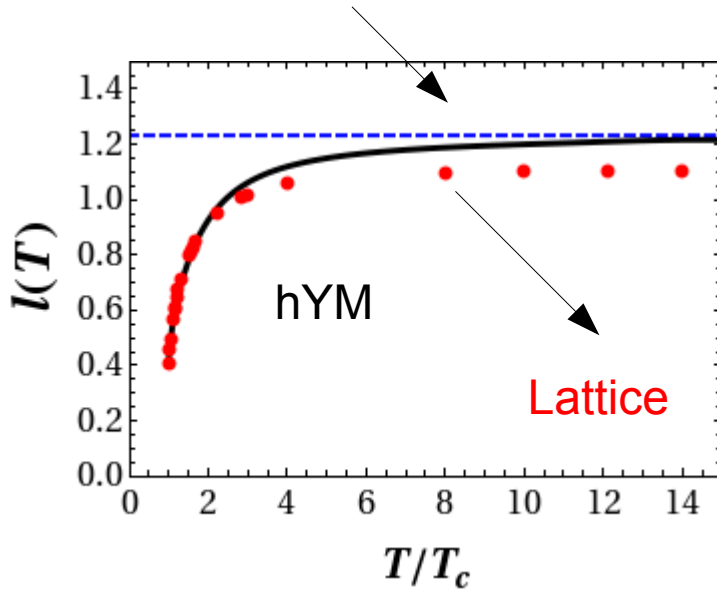
$q(\phi)$  coupling between string and the scalar field

JN, PRD (2010).

Let's assume a dilaton-like coupling function  $q(\phi) = e^{\sqrt{2/3}\phi}$

Kiritsis et al, 2008

Value at  $T/T_c \gg 1$



Good agreement below  $4T_c$  ...

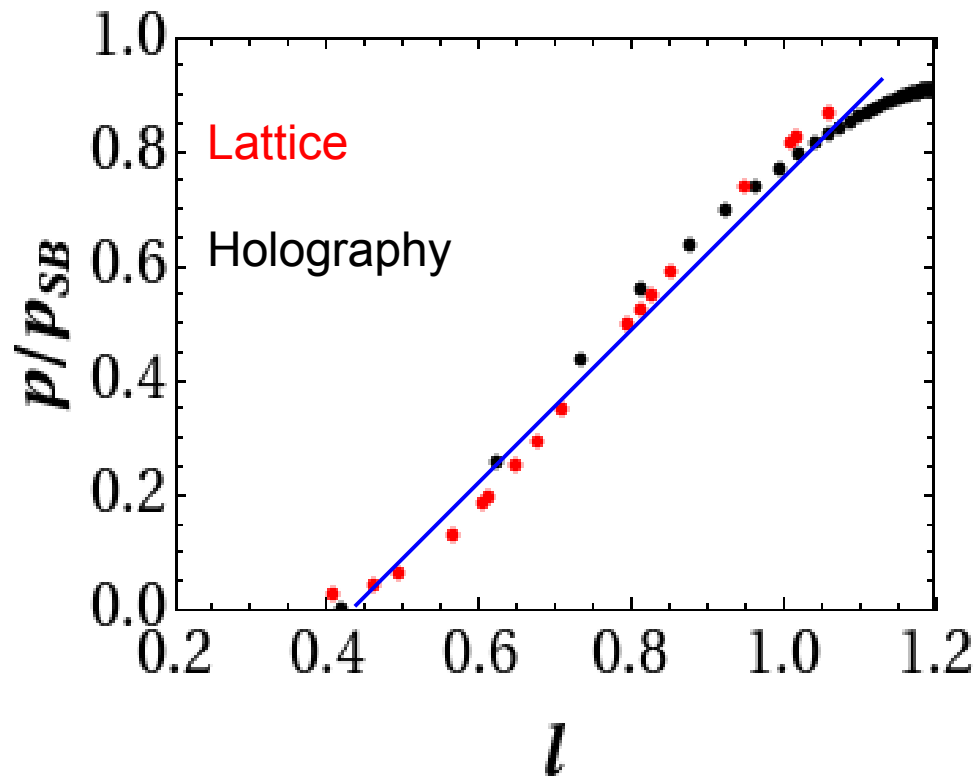
BUT

$R^2/\alpha' \sim 1$

Higher corrections to the gravity dual will be important here !!!

It should be expected because even here  $\eta/s = 1/4\pi$  at any  $T > T_c$

In this case the approximate linear relation between  $P$  and the loop is reproduced



Other interesting properties, **defined in real time**, that I have no time to discuss here in detail ...

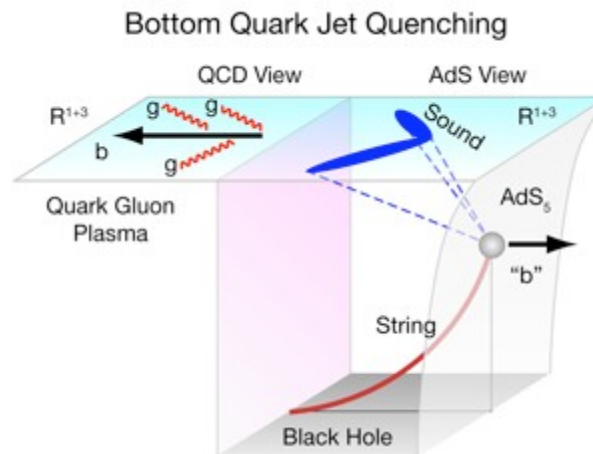
- $\eta/s = 1/(4\pi)$  above  $T_c$ . Below  $T_c$  one should expect  $\eta/s \sim N_c^2$

JN, arXiv:0912.4824 [hep-th]

- $\zeta/s$  has a peak at  $T_c$ .

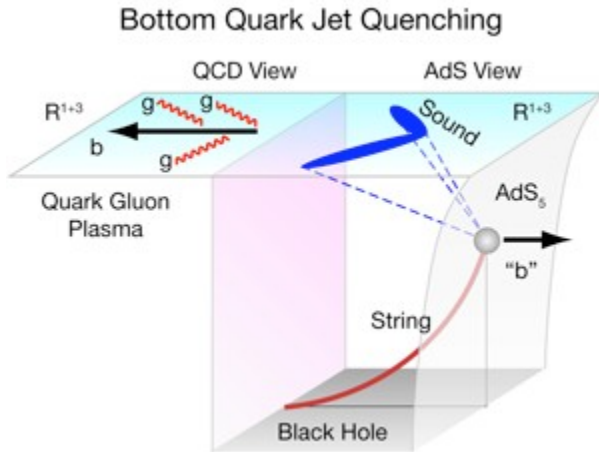
Gubser et al (2008).

## What about energy loss????



For the trailing string (heavy quarks) one can show that

Ficnar, JN, Gyulassy, to appear



$$v \rightarrow 1$$

Ultrarelativistic heavy quark

$$\frac{dE}{dx} = \left( \frac{dE}{dx} \right)_{CFT} \left( \frac{S}{S_{CFT}} \right)^{1/2}$$

Here we know exactly how confinement affects the heavy quark energy loss ...

↓

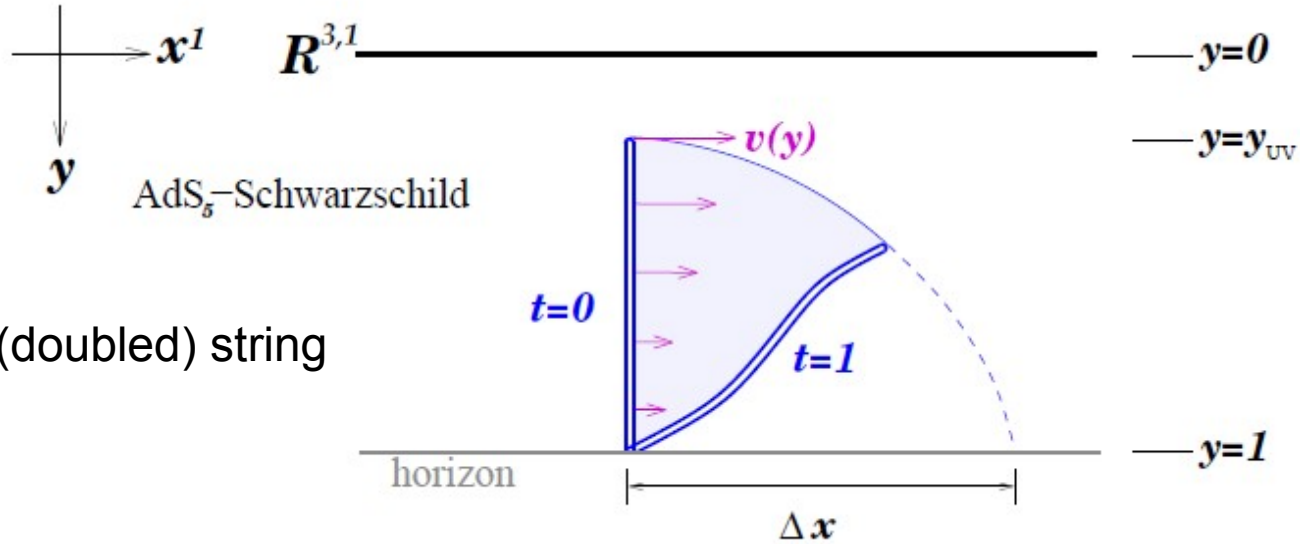
**DIRECT dependence on the medium properties (entropy density) !!!**

**No Tc enhancement !!!**

What happens to the charm to bottom ratio in R\_AA ???

**What about the glue energy loss ????**

Gubser et al, 2008



The “falling (doubled) string problem”

**For a CFT one finds**  $\Delta x \sim \frac{1}{T} \left( \frac{E}{T} \right)^{1/3}$

See also Chesler et al, 2009.

**BDMPs:**  $\Delta E \sim L^2$

## What about the glue energy loss in the semi QGP ????

Ficnar, JN, Gyulassy, to appear

The  $\Delta E \sim L^3$  may not survive when conformal invariance is badly broken near  $T_c$ . Power of  $L$  may be sensitive to the phase transition!

In fact, we know already that in geometries which lead to confinement, such as the Sakai-Sugimoto model,

$$\Delta E \sim L^{2.5} \quad \text{Yi Pang, 2008.}$$

**Confinement matters!!!**

**The gauge/string duality is the only framework where such questions can be reliably addressed.**

## Conclusions and Outlook

- The new state of QCD matter created in heavy ions is most likely a semi QGP, which is not conformal and is fundamentally different from the Gedanken free gas (used, for instance, in some jet quenching calculations).
- Gauge/string duality can be used to construct gravity duals that describe a chunk of hot strongly coupled matter that burns and flows like the semi QGP.
- These holographic theories naturally incorporate the non-perturbative power-like temperature terms observed on the lattice (perturbative calculations don't).
- It is possible, using the gauge/string duality, to understand how highly energetic probes lose energy near the deconfinement phase transition (predictions are on the way)

Back-up Slides