

# **Dynamical coupled-channels study of hyperon resonances using anti-kaon induced reactions**

**Hiroyuki Kamano  
(RCNP, Osaka U.)**

**Collaborators:**

**T.-S. H. Lee (Argonne )  
S. Nakamura (Osaka U.)  
T. Sato (Osaka U.)**

# Background & Motivation

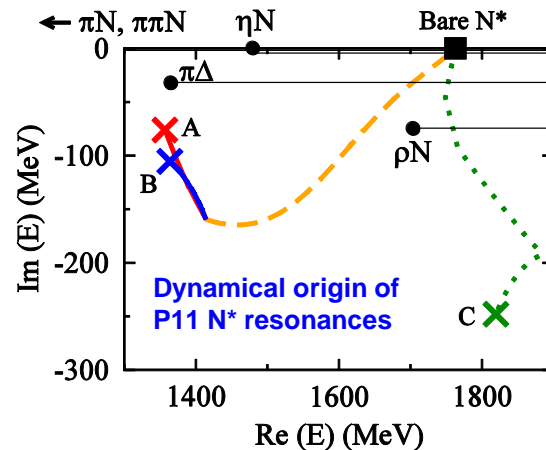
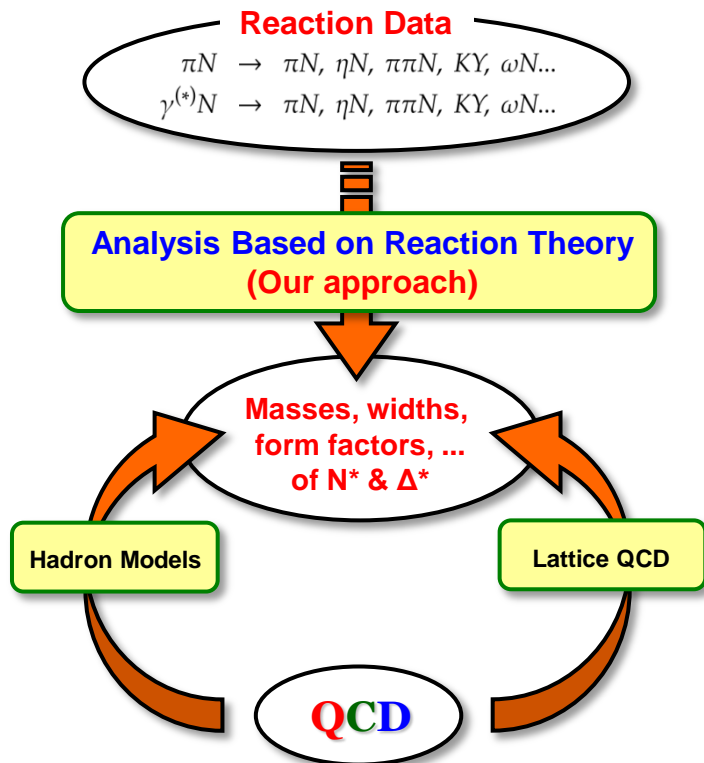
## $N^*$ spectroscopy via global analysis of $\pi N$ and $\gamma N$ reactions

- Based on **Dynamical Coupled-Channels (DCC)** approach [Matsuyama, Sato, Lee, Phys. Rep. 439, 193 (2007)]

## Latest published analysis (**ANL-Osaka**):

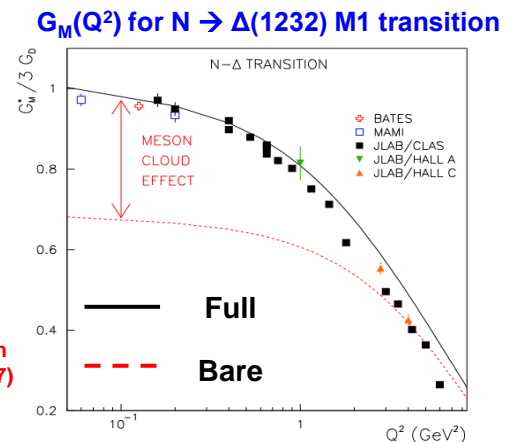
**Fully combined analysis** of  $\pi N, \gamma N \rightarrow \pi N, \eta N, K\Lambda, K\Sigma$  up to  $W = 2.1$  GeV. [HK, Nakamura, Lee, Sato, PRC88(2013)035209]

- Revealed role of **multichannel meson-baryon reaction dynamics** in understanding  $N^*$  and  $\Delta^*$  resonances:



Suzuki, Julia-Diaz, HK, Lee, Matsuyama, Sato  
PRL104 065203 (2010)

Julia-Diaz, Lee, Sato, Smith  
PRC75 015205 (2007)



# Background & Motivation

## $N^*$ spectroscopy via global analysis of $\pi N$ and $\gamma N$ reactions

- Based on **Dynamical Coupled-Channels (DCC)** approach [Matsuyama, Sato, Lee, Phys. Rep. 439, 193 (2007)]

### Reaction Data

$\pi N \rightarrow \pi N, \eta N, \pi\pi N, KY, \omega N \dots$   
 $\gamma^{(*)} N \rightarrow \pi N, \eta N, \pi\pi N, KY, \omega N \dots$

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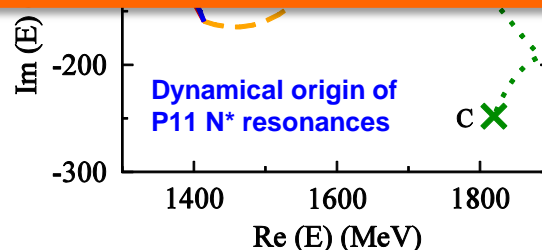
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Apply our DCC approach to the spectroscopy of hyperon resonances ( $Y^* = \Lambda^*, \Sigma^*$ ) !!

Hadron Models

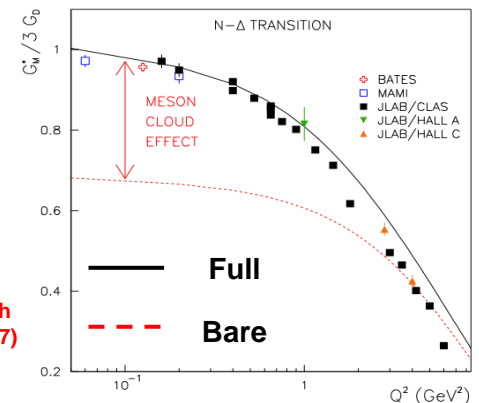
Lattice QCD

QCD



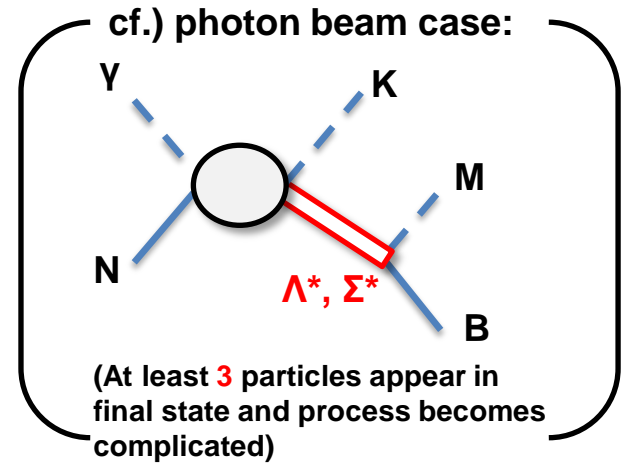
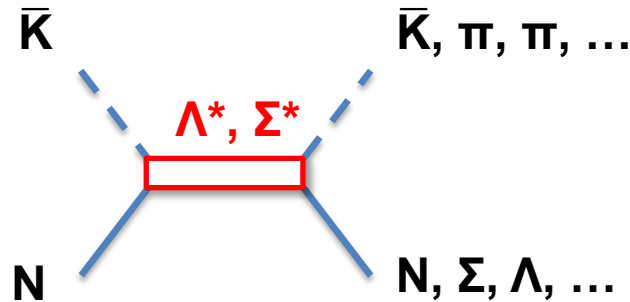
Julia-Diaz, Lee, Sato, Smith  
PRC75 015205 (2007)

$G_M(Q^2)$  for  $N \rightarrow \Delta(1232)$  M1 transition

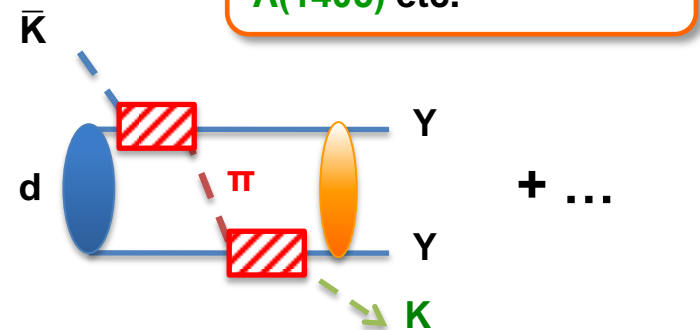
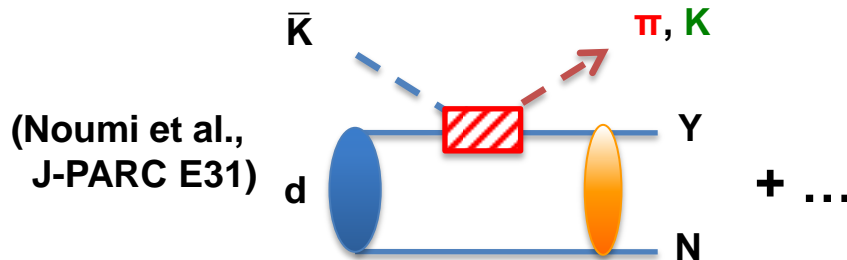


# Y\* spectroscopy using anti-kaon beam

- The **simplest** reactions for studying Y\* (=  $\Lambda^*$ ,  $\Sigma^*$ ).



- Deuteron reactions allow one to directly access  $\bar{K}N$  subthreshold region, and to study  $YN$  and  $YY$  interactions.



Necessary for studying  $\Lambda(1405)$  etc.

# Dynamical coupled-channels (DCC) approach for $Y^*$ production reactions

HK, Nakamura, Lee, Sato, PRC90(2014)065204

- ✓ Coupled-channels integral equations for partial-wave amplitudes of  $a \rightarrow b$  reaction:

$$T_{b,a}^{(LSJ)}(p_b, p_a; E) = V_{b,a}^{(LSJ)}(p_b, p_a; E) + \sum_c \int_0^\infty q^2 dq \underbrace{V_{b,c}^{(LSJ)}(p_b, q; E)}_{\text{CC effect}} \underbrace{G_c(q; E)}_{\text{off-shell effect}} T_{c,a}^{(LSJ)}(q, p_a; E)$$

- ✓ Reaction channels:

$$a, b, c = ( \bar{K}N, \pi\Sigma, \pi\Lambda, \eta\Lambda, K\Xi, \boxed{\pi\Sigma^*, \bar{K}^*N, \dots} )$$

quasi two-body channels

- ✓ Transition Potentials:

$$V_{a,b} = \underbrace{v_{a,b}}_{\text{Exchange potentials}} + \sum_{Y^*} \frac{\Gamma_{Y^*,a}^\dagger \Gamma_{Y^*,b}}{E - M_{Y^*}} \underbrace{\hspace{10em}}_{\text{Bare } Y^* \text{ states}}$$

# Dynamical coupled-channels (DCC) approach for $Y^*$ production reactions

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CC effect      off-shell effect

✓ Meson-Baryon Green functions

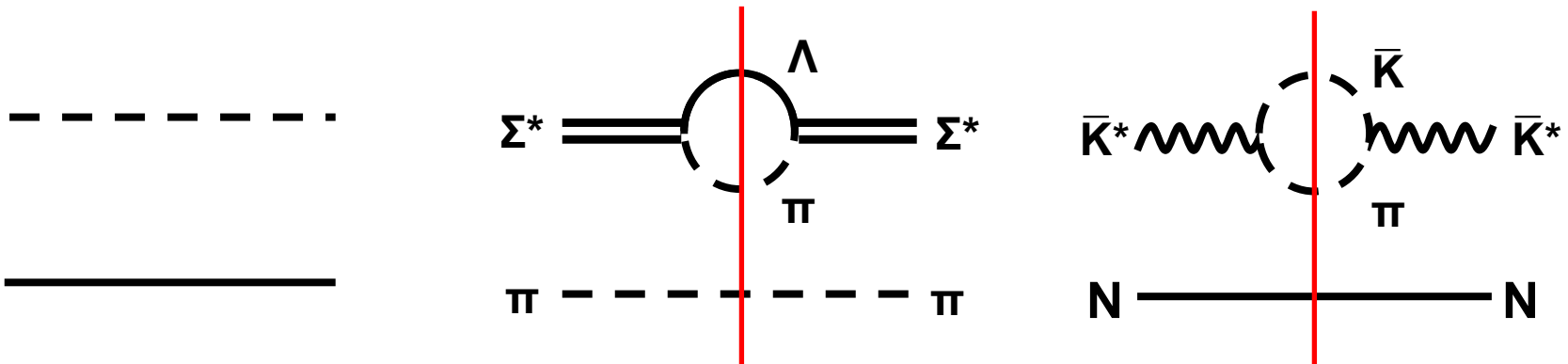
$MB = \bar{K}N, \pi\Sigma, \pi\Lambda, \eta\Lambda, K\Xi$

Stable channels

$MB = \pi\Sigma^*, \bar{K}^*N$

Quasi 2-body channels

Produces three-body unitary cuts !!



Exchange potentials

Bare T states

# Dynamical coupled-channels (DCC) approach for $Y^*$ production reactions

HK, Nakamura, Lee, Sato, PRC90(2014)065204

- ✓ Coupled-channels integral equations for partial-wave amplitudes of  $a \rightarrow b$  reaction:

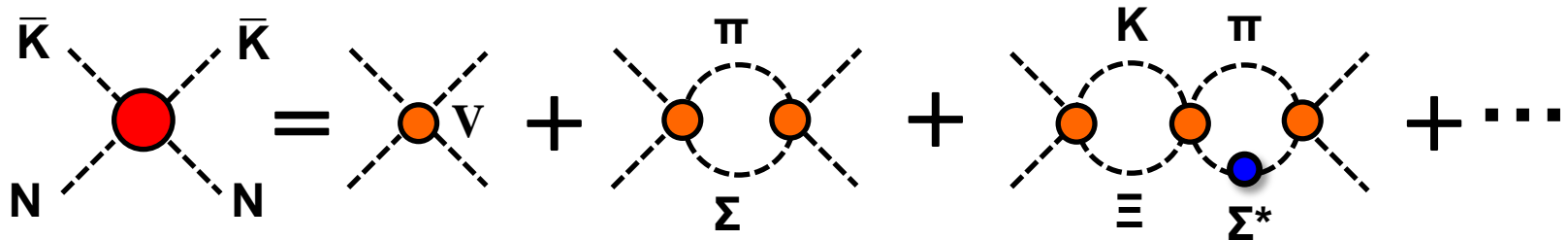
$$T_{b,a}^{(LSJ)}(p_b, p_a; E) = V_{b,a}^{(LSJ)}(p_b, p_a; E) + \sum_c \int_0^\infty q^2 dq V_{b,c}^{(LSJ)}(p_b, q; E) G_c(q; E) T_{c,a}^{(LSJ)}(q, p_a; E)$$

CC effect      off-shell effect

✓ Reaction channels

- ✓ Summing up all possible transitions between reaction channels !!  
 (→ satisfies **multichannel two-** and **three-body unitarity**)

e.g.)  $\bar{K}N$  scattering



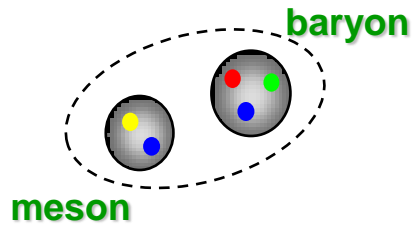
- ✓ **Momentum integral** takes into account **off-shell rescattering effects** in the intermediate processes.

# Dynamical coupled-channels (DCC) approach for $Y^*$ production reactions

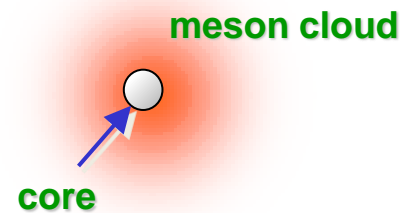
HK, Nakamura, Lee, Sato, PRC90(2014)065204

- ✓ Coupled-channels integral equations for partial-wave amplitudes of  $a \rightarrow b$  reaction:

Physical  $Y^*$ s will be a “mixture” of the two pictures:



$$|Y^*\rangle = |MB\rangle$$



$$|Y^*\rangle = |qqq\rangle + |m.c.\rangle$$

- ✓ Transition Potentials:

$$V_{a,b} = v_{a,b} + \sum_{Y^*} \frac{\Gamma_{Y^*,a}^\dagger \Gamma_{Y^*,b}}{E - M_{Y^*}}$$

Exchange potentials

Bare  $Y^*$  states



# What we have done so far

With the dynamical coupled-channels approach developed for the **S= -1 sector**, we made:

- ✓ Comprehensive analysis of *all* available data of  **$K^- p \rightarrow \bar{K}N, \pi\Sigma, \pi\Lambda, \eta\Lambda, K\Xi$**  up to  **$W = 2.1$  GeV**.  
[HK, Nakamura, Lee, Sato, PRC90(2014)065204]
  - Successfully determined the partial-wave amplitudes of  **$\bar{K}N \rightarrow \bar{K}N, \pi\Sigma, \pi\Lambda, \eta\Lambda, K\Xi$**  for **S, P, D, and F waves !!**
- ✓ Extraction of  **$\Lambda^*$  and  $\Sigma^*$**  mass spectrum defined by **poles of scattering amplitudes**.  
[HK, Nakamura, Sato, in preparation]

# Database of our analysis ( $W < 2.1\text{GeV}$ )

HK, Nakamura, Lee, Sato, PRC90(2014)065204

## Issues in the availability of data:

- ✓ Most data are from 60-70's.
- ✓ Kinematical coverage is rather scarce for most reactions.
- ✓ No data for spin rotations ( $\beta, R, A$ ).
- ✓ No data near the threshold for  $K^- p \rightarrow \bar{K}N, \pi\Sigma, \pi\Lambda$ .



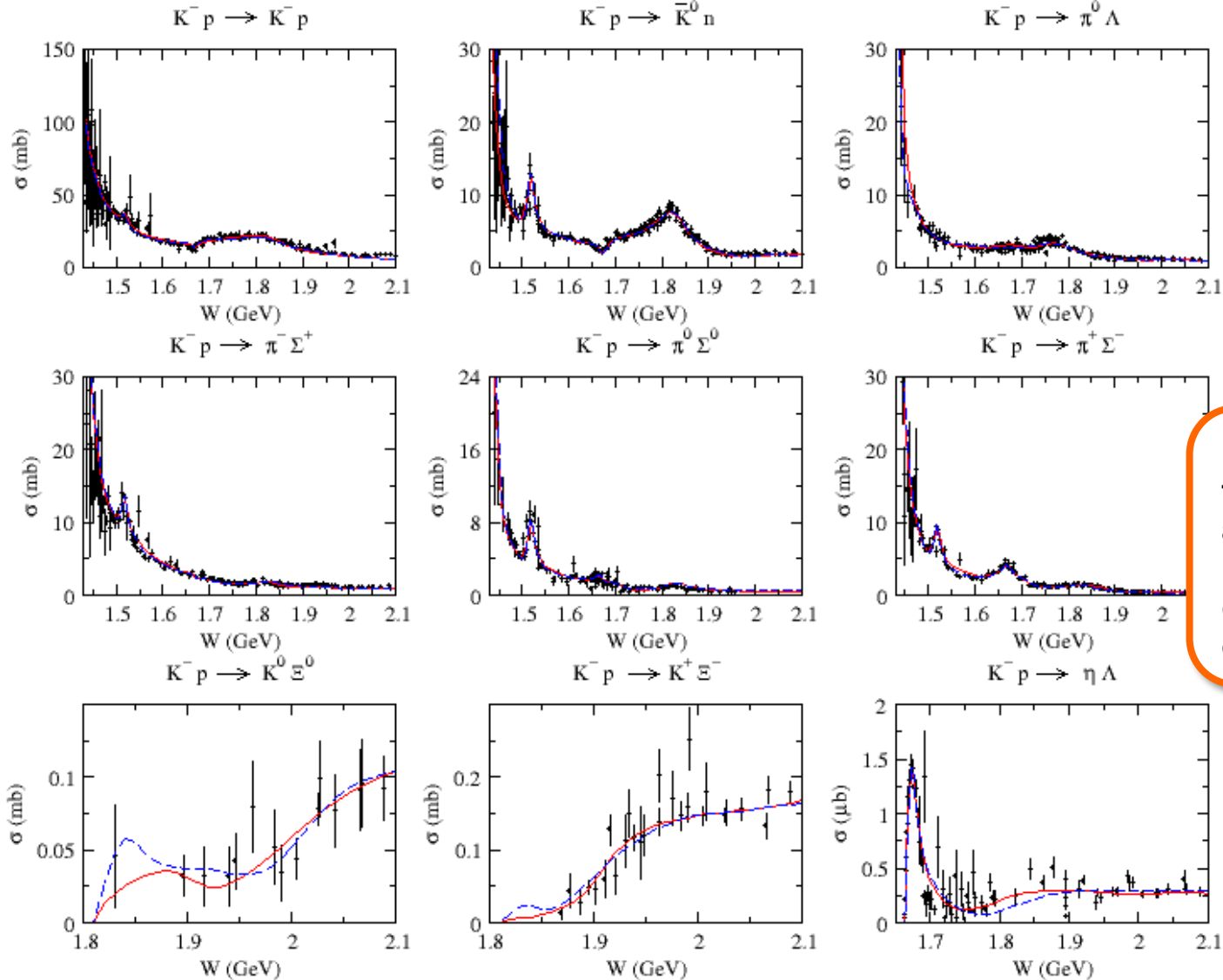
The  $K^- p$  reaction data are far from “complete”!!  
 → Need help of hadron beam facilities such as J-PARC !!

Reactions	Observables	Number of data	
$K^- p \rightarrow K^- p$	$d\sigma/d\Omega$	3962	} $d\sigma/d\Omega$ : 1465 MeV < W P : 1730 MeV < W $\beta, R, A$ : No data
	P	510	
	$\sigma$	253	
$K^- p \rightarrow \bar{K}^0 n$	$d\sigma/d\Omega$	2950	} $d\sigma/d\Omega$ : 1465 MeV < W P : No data $\beta, R, A$ : No data
	$\sigma$	260	
$K^- p \rightarrow \pi^- \Sigma^+$	$d\sigma/d\Omega$	1792	} $d\sigma/d\Omega$ : 1535 MeV < W P : 1535 MeV < W < 1967 MeV $\beta, R, A$ : No data
	P	418	
	$P \times d\sigma/d\Omega$	177	
	$\sigma$	173	
$K^- p \rightarrow \pi^0 \Sigma^0$	$d\sigma/d\Omega$	580	} $d\sigma/d\Omega$ : 1535 MeV < W < 1763 MeV P : 1535 MeV < W < 1696 MeV $\beta, R, A$ : No data
	P	196	
	$P \times d\sigma/d\Omega$	189	
	$\sigma$	125	
$K^- p \rightarrow \pi^+ \Sigma^-$	$d\sigma/d\Omega$	1786	} $d\sigma/d\Omega$ : 1536 MeV < W P : No data $\beta, R, A$ : No data
	$\sigma$	181	
$K^- p \rightarrow \pi^0 \Lambda$	$d\sigma/d\Omega$	2178	} $d\sigma/d\Omega$ : 1535 MeV < W P : 1535 MeV < W $\beta, R, A$ : No data
	P	693	
	$P \times d\sigma/d\Omega$	176	
	$\sigma$	207	
$K^- p \rightarrow \eta \Lambda$	$d\sigma/d\Omega$	160	} $d\sigma/d\Omega$ : 1664 MeV < W < 1696 MeV P : 1669 MeV < W < 1681 MeV $\beta, R, A$ : No data
	P	18	
	$\sigma$	78	
$K^- p \rightarrow K^0 \Xi^0$	$d\sigma/d\Omega$	33	} $d\sigma/d\Omega$ : 1970 MeV < W < 2070 MeV P : No data $\beta, R, A$ : No data
	$\sigma$	15	
$K^- p \rightarrow K^+ \Xi^-$	$d\sigma/d\Omega$	92	} $d\sigma/d\Omega$ : 1950 MeV < W < 2070 MeV P : No data $\beta, R, A$ : No data
	$\sigma$	27	
Total		17229	

# Results of the fits

$K^- p \rightarrow MB$  total cross sections

HK, Nakamura, Lee, Sato, PRC90(2014)065204



Red: Model A

Blue: Model B

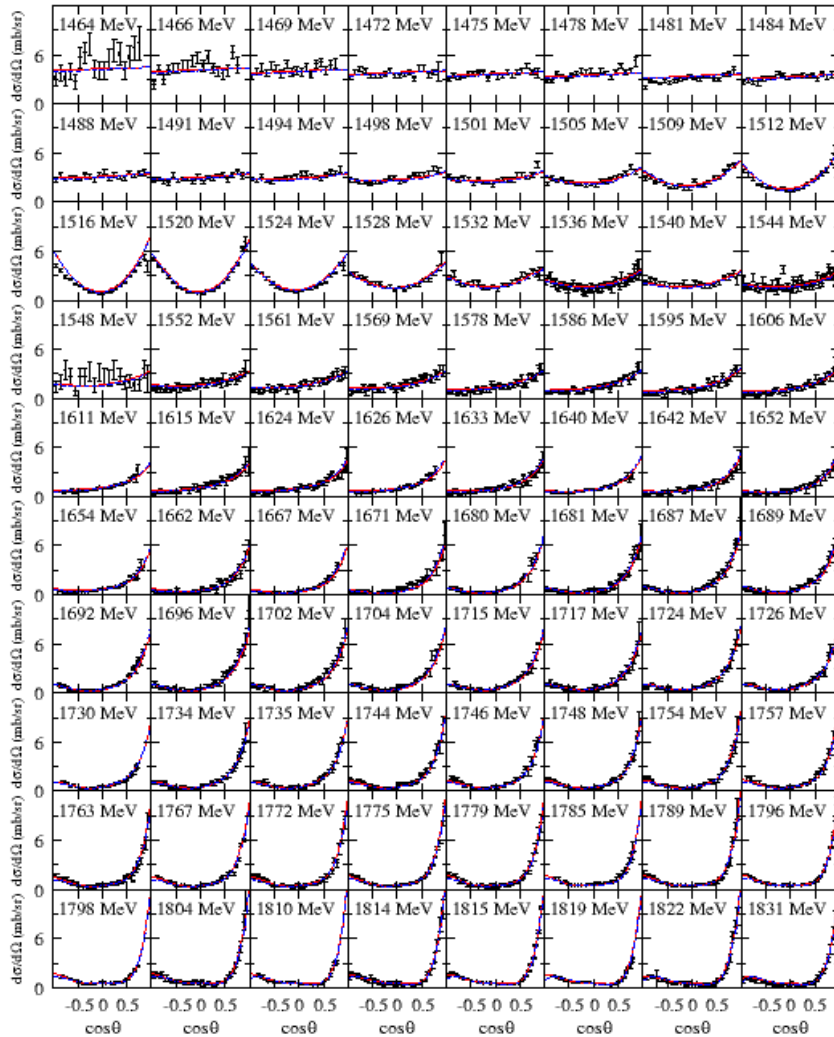
“Incompleteness” of the current database allows us to have two parameter sets that give similar quality of the fit.

# Results of the fits

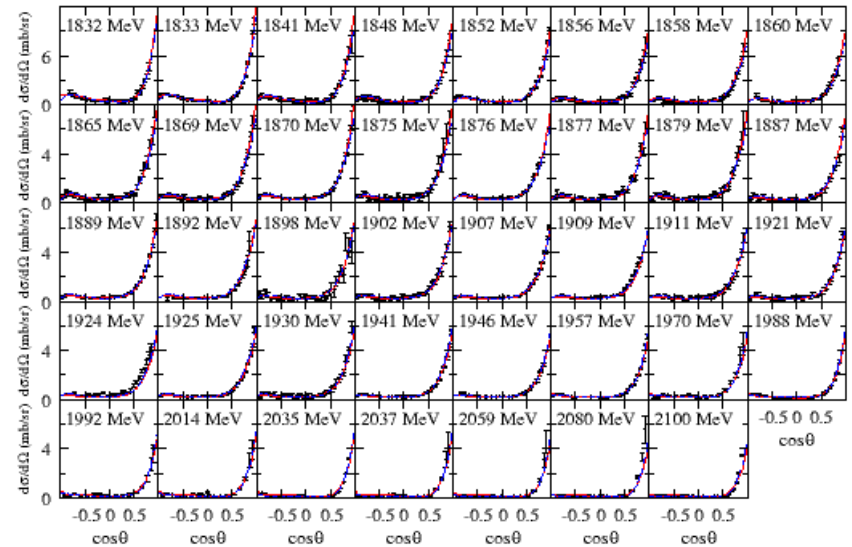
$K^- p \rightarrow K^- p$  scattering

HK, Nakamura, Lee, Sato, PRC90(2014)065204

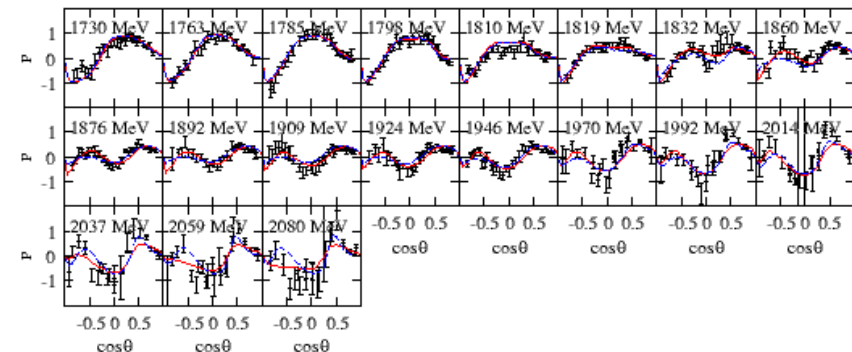
$d\sigma/d\Omega$  ( $1464 < W < 1831$  MeV)



$d\sigma/d\Omega$  ( $1832 < W < 2100$  MeV)



P



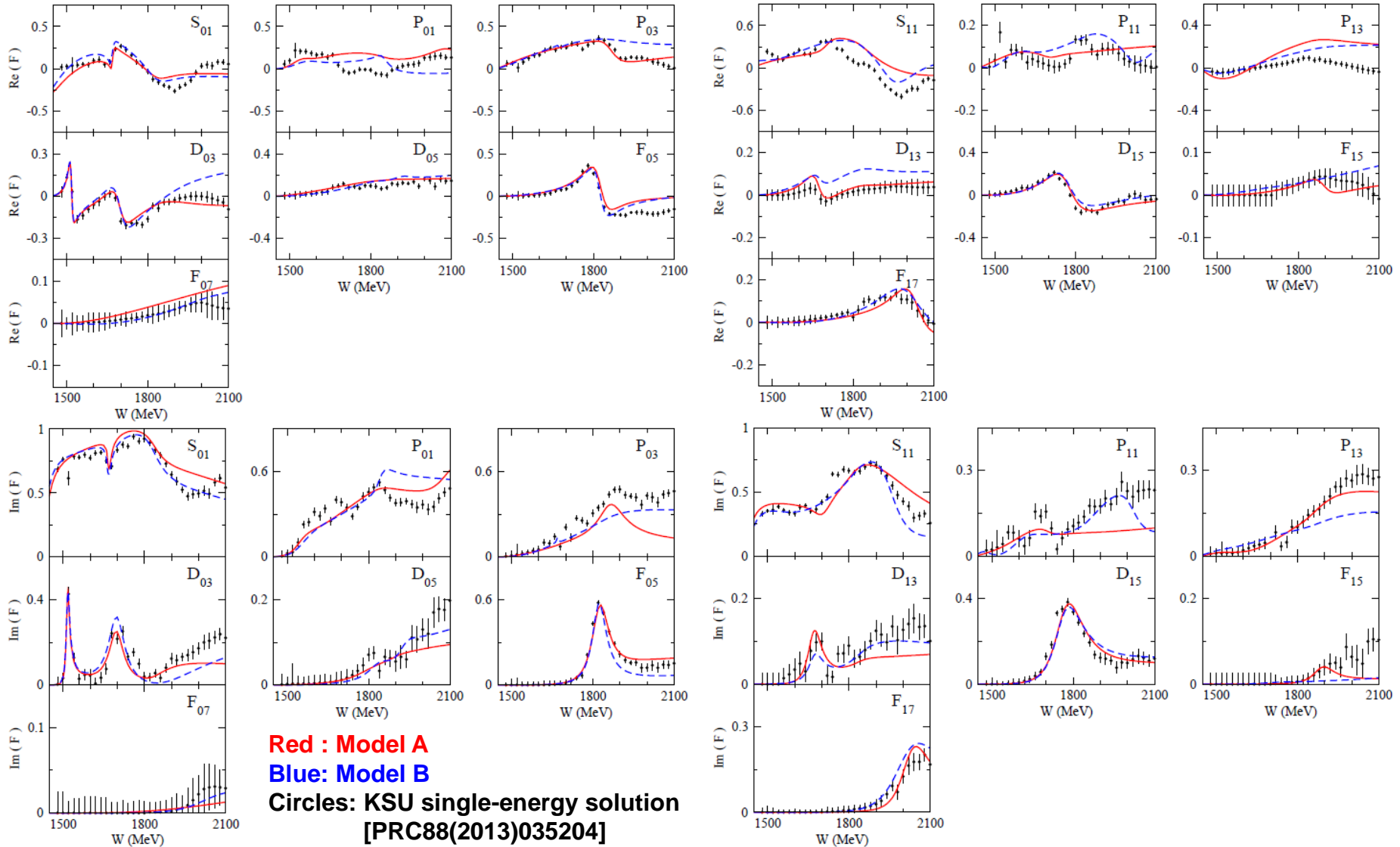
Red: Model A Blue: Model B

# Comparison of extracted partial-wave amplitudes

## Extracted $\bar{K}N$ scattering amplitudes

HK, Nakamura, Lee, Sato, PRC90(2014)065204

$L_{I2J}$  :  $L = S, P, \dots$  ;  $I =$  isospin;  $J =$  Total angular mom.

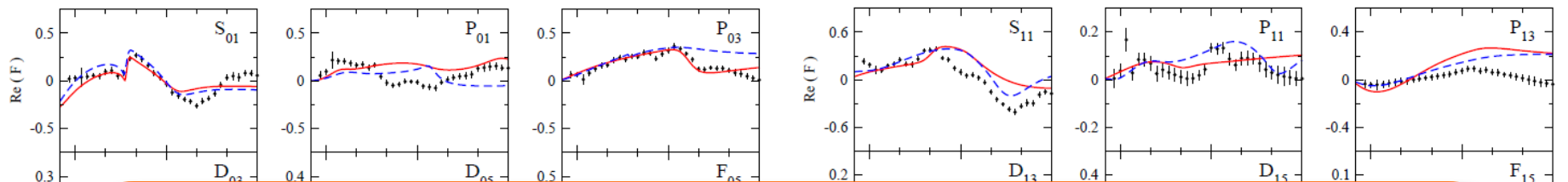


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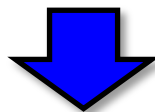
HK, Nakamura, Lee, Sato, PRC90(2014)065204

## Extracted $\bar{K}N$ scattering amplitudes

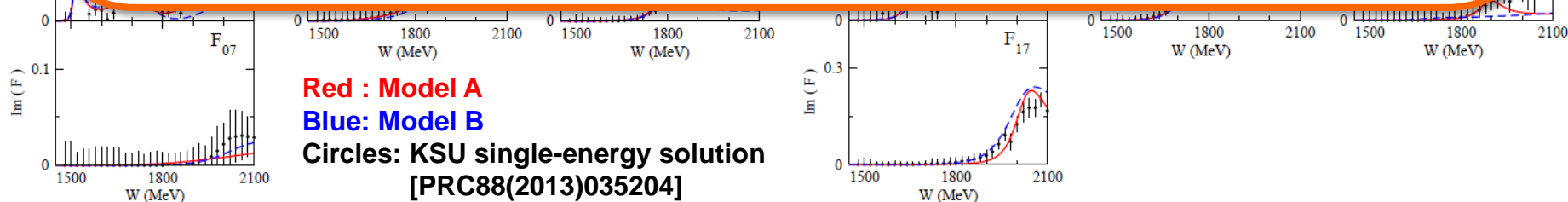
$L_{12J}$  :  $L = S, P, \dots$  ;  $I =$  isospin;  $J =$  Total angular mom.



**On- and off-shell** partial-wave amplitudes for  
 $\bar{K}N \rightarrow \pi\Sigma$ ,  $\bar{K}N \rightarrow \pi\Lambda$ ,  $\bar{K}N \rightarrow \eta\Lambda$ ,  $\bar{K}N \rightarrow K\Xi$   
are also available for **S, P, D, and F waves !!**



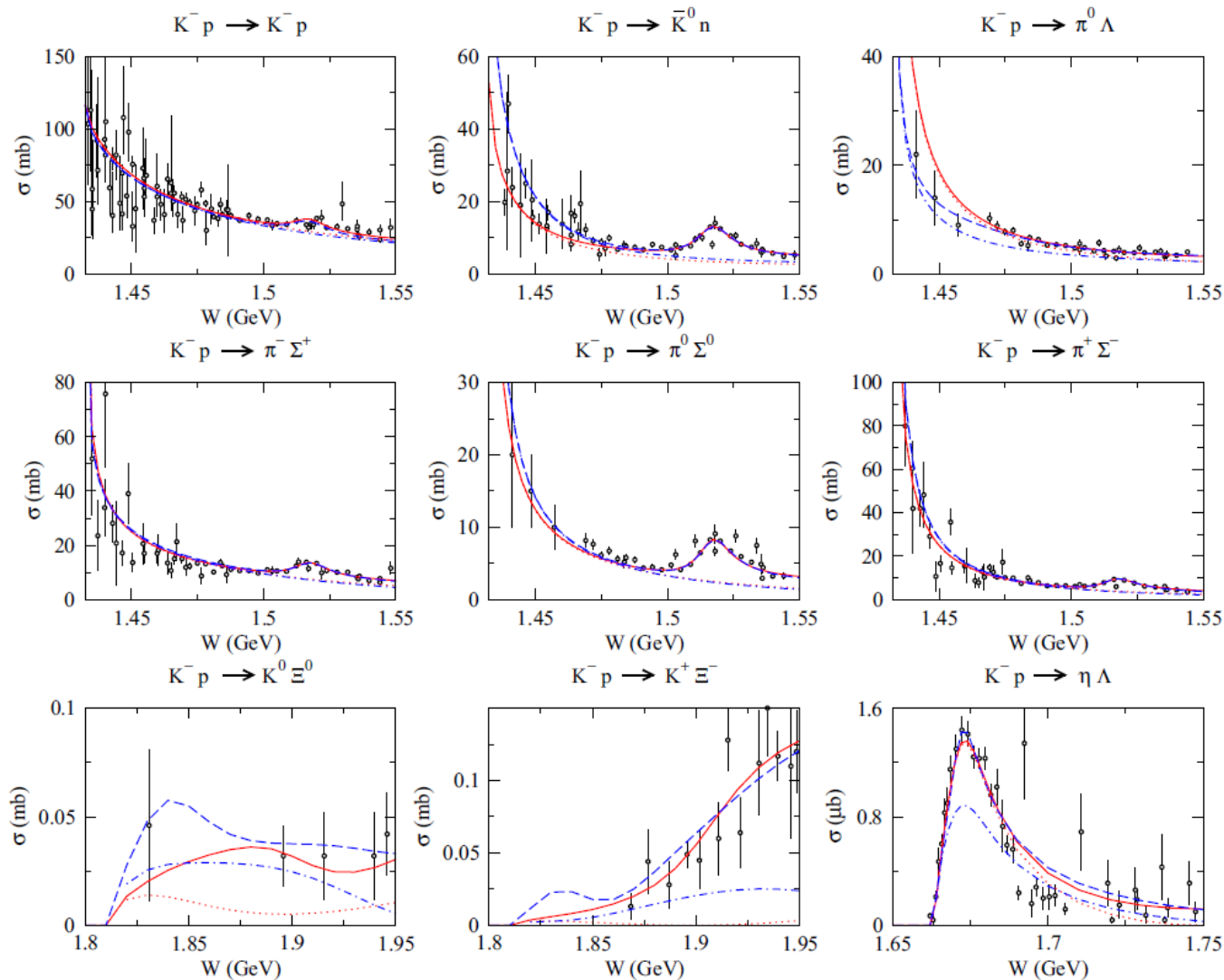
Input for elementary processes of nuclear  
target reactions (**hypernuclei, kaonic nuclei...**).



# S-wave contributions in the threshold region

$K^- p \rightarrow MB$  total cross sections

HK, Nakamura, Lee, Sato, PRC90(2014)065204

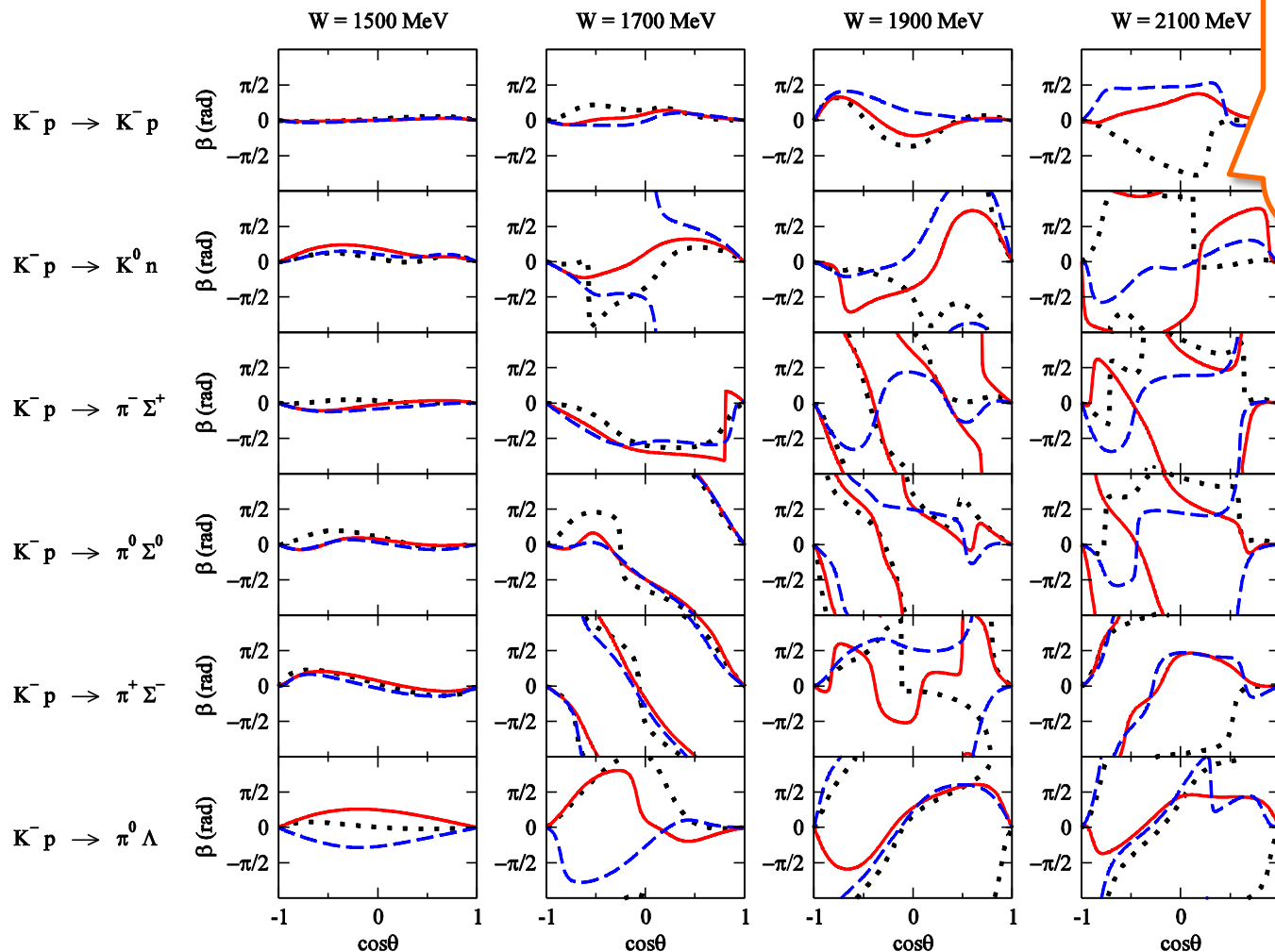


Higher partial waves  
can be significant  
at the region even  
close to the threshold !!

- Model A (full)
- ..... Model A (S wave only)
- - - Model B (full)
- . - Model B (S wave only)

# Predicted spin-rotation angle $\beta$

HK, Nakamura, Lee, Sato, PRC90(2014)065204



Analysis dependence is clearly seen in observables that are not yet measured.



Measurement of spin-rotation  $\beta$  will give strong constraints on the  $Y^*$  spectrum !!

**Red: Model A**

**Blue: Model B**

**Black: KSU**

The KSU results are computed by us using their amplitudes in PRC88(2013)035204.

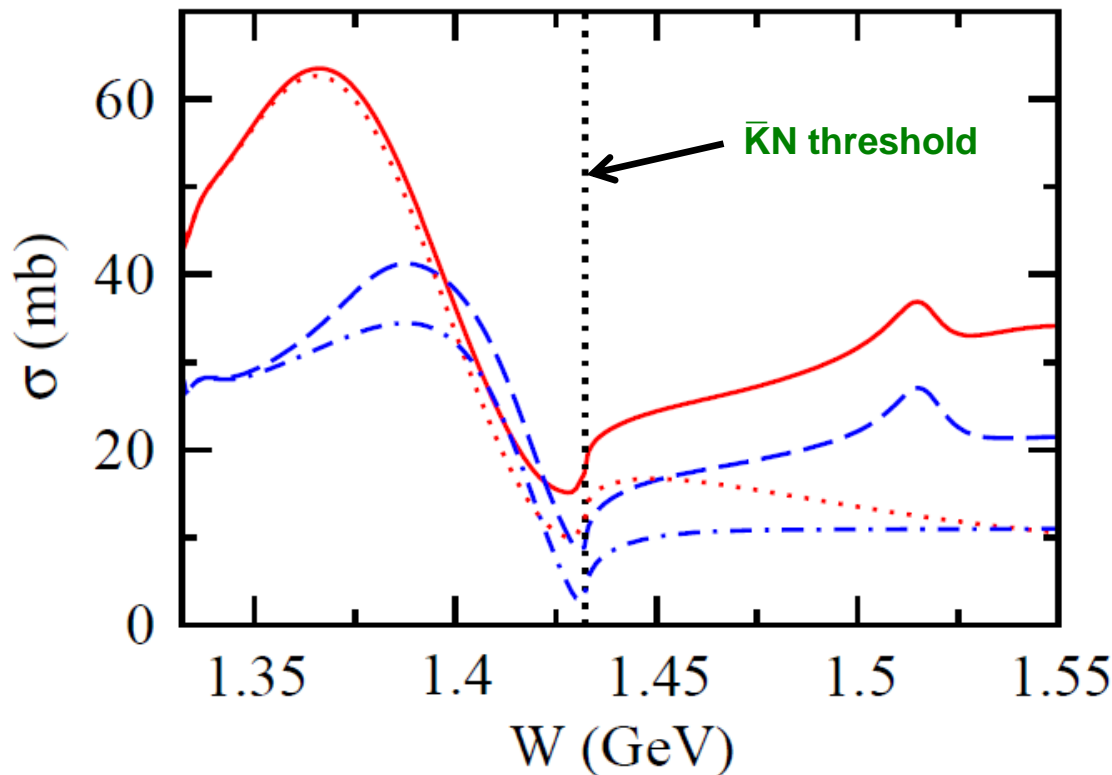
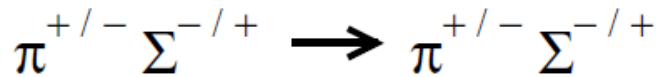
**## NOTE:**  
 $\beta$  is modulo  $2\pi$



# Predicted $\pi\Sigma$ scattering total cross section at low energies

HK, Nakamura, Lee, Sato, PRC90(2014)065204

- ✓ Predicted total cross section  $\sigma$  of  $\pi$ - $\Sigma$  scattering from the threshold up to  $W = 1.55$  GeV.



- Model A Full
- ⋯ Model A S-wave only
- - - Model B Full
- . - Model B S-wave only

Contribution from **higher partial waves** can be **sizeable** in the  $\bar{K}N$  subthreshold region !!

# Extracting $Y^*$ resonance parameters

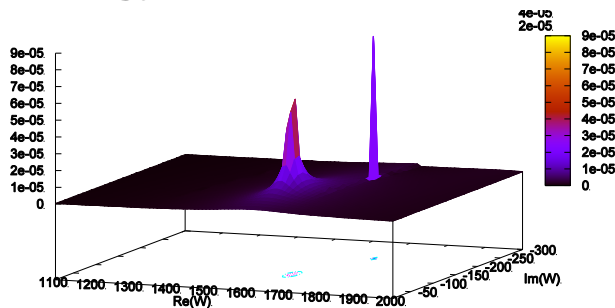
## Definitions of

- ✓  $Y^*$  masses (spectrum) → Pole positions of the amplitudes
- ✓  $Y^* \rightarrow$  MB coupling constants → Residues<sup>1/2</sup> at the pole

$Y^* \rightarrow b$   
coupling constant

$$\langle p_a | \hat{T}(E) | p_b \rangle \Big|_{E \rightarrow E_0} \rightarrow \frac{\bar{\Gamma}(E_0, p_a) \bar{\Gamma}(E_0, p_b)}{E - E_0} + (\text{regular terms})$$

Analytic continuation to (lower-half) complex energy plane.



$Y^*$  pole position  
( $\text{Im}(E_0) < 0$ )

(Multichannel) unitarity is a key to making “correct” analytic continuation !!

# Extracting $Y^*$ resonance parameters

## Definitions of

- ✓  $Y^*$  masses (spectrum) → Pole positions of the amplitudes
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$Y^* \rightarrow b$   
coupling constant

Consistent with the resonance theory based on **Gamow vectors**

G. Gamow (1928), R. E. Peierls (1959), ...

For a brief introduction of Gamow vectors, see, e.g., de la Madrid et al, quant-ph/0201091

→ Resonances are (**complex-energy**) **eigenstates** of the Hamiltonian of the underlying fundamental theory with **the purely outgoing boundary condition !!**

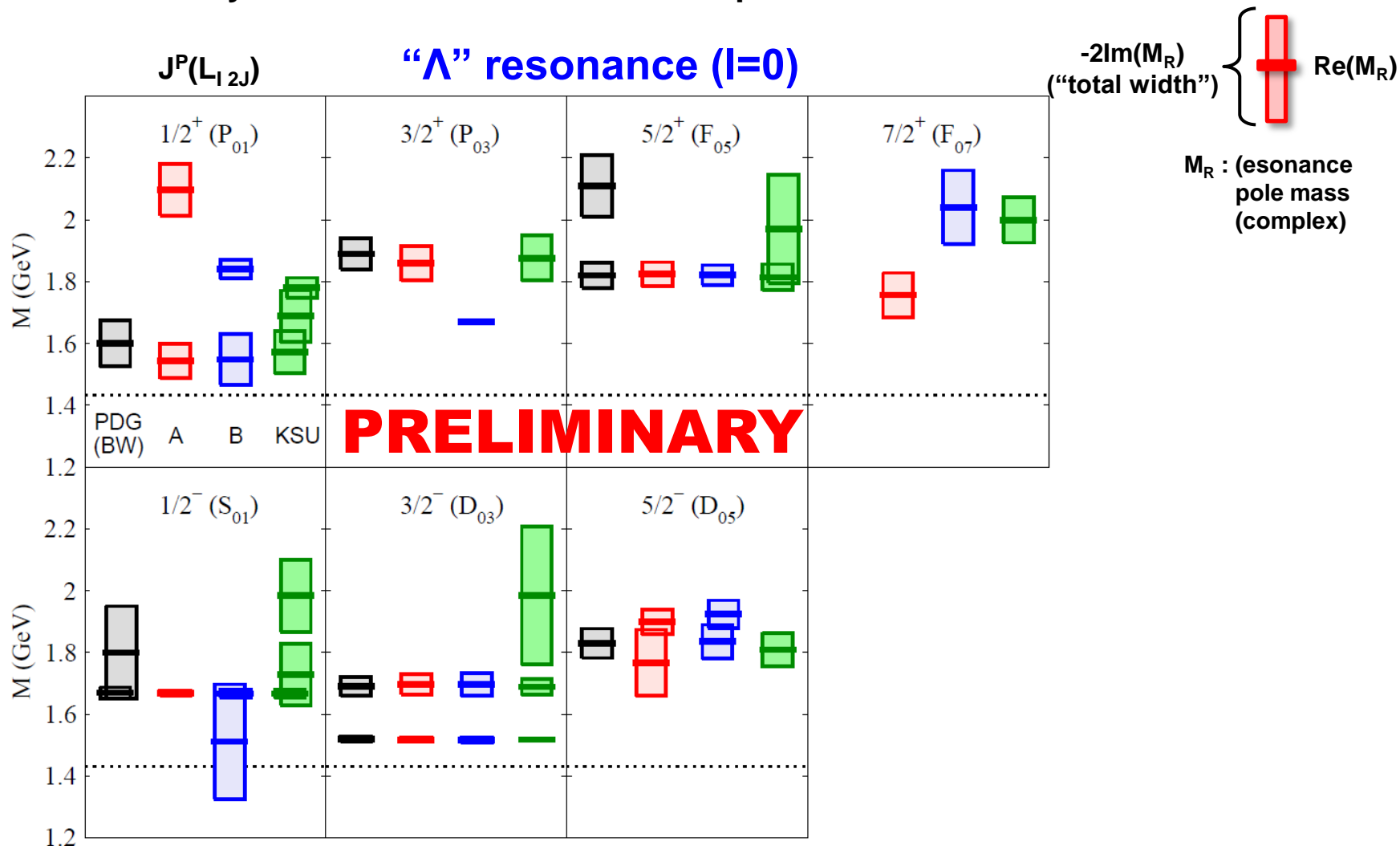
(**complex**) energy eigenvalues = pole values

transition matrix elements = (**residue**)<sup>1/2</sup> of the poles

# Comparison of $\Lambda^*$ spectrum between multichannel analyses

HK, Nakamura, Sato, in preparation

### Here only  $Y^*$ s ABOVE  $\bar{K}N$  threshold are presented.

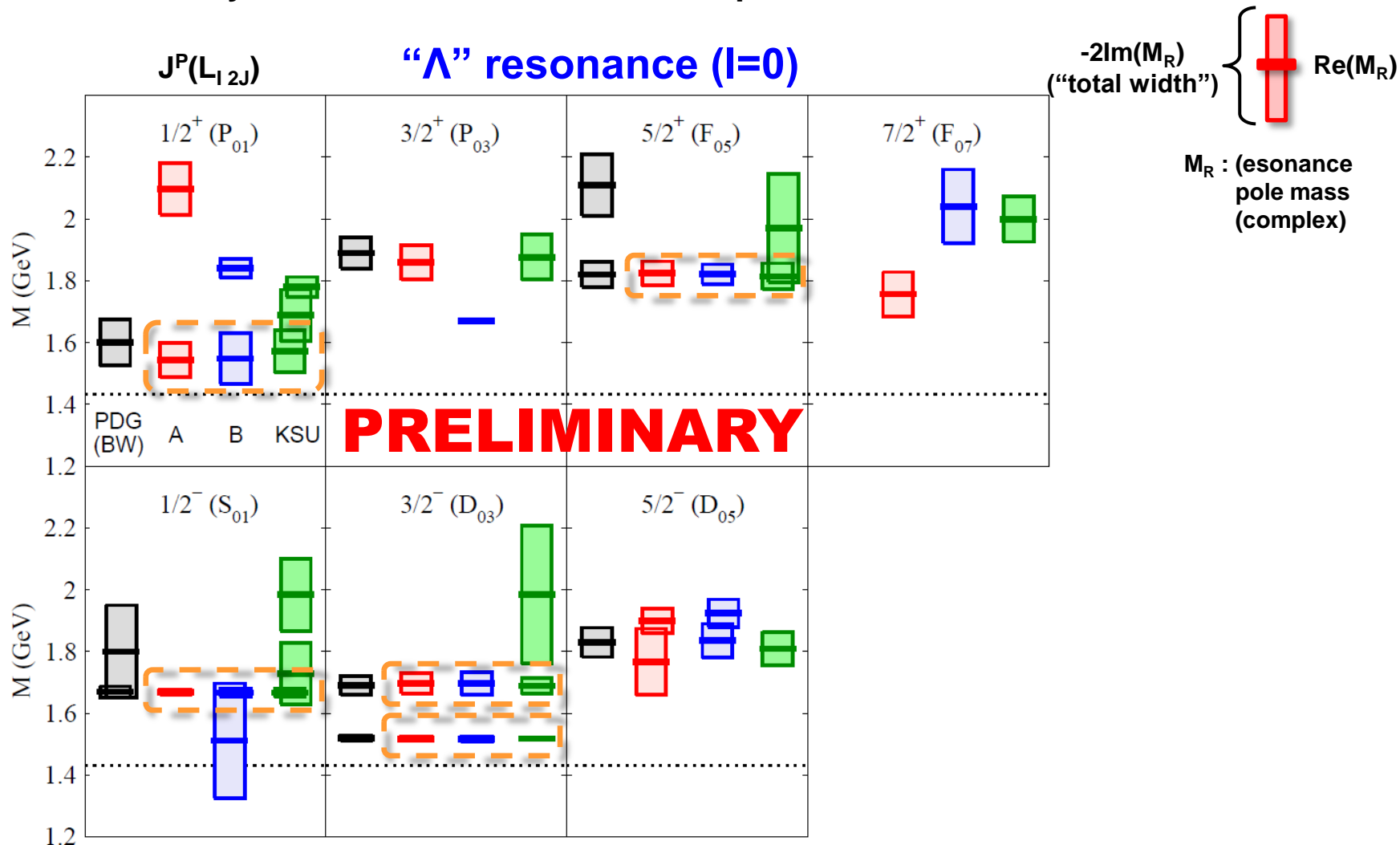


Red: Model A, Blue: Model B, Green: KSU[on-shell K-matrix,PRC88(2013)035205], Black: PDG(Breit-Wigner)

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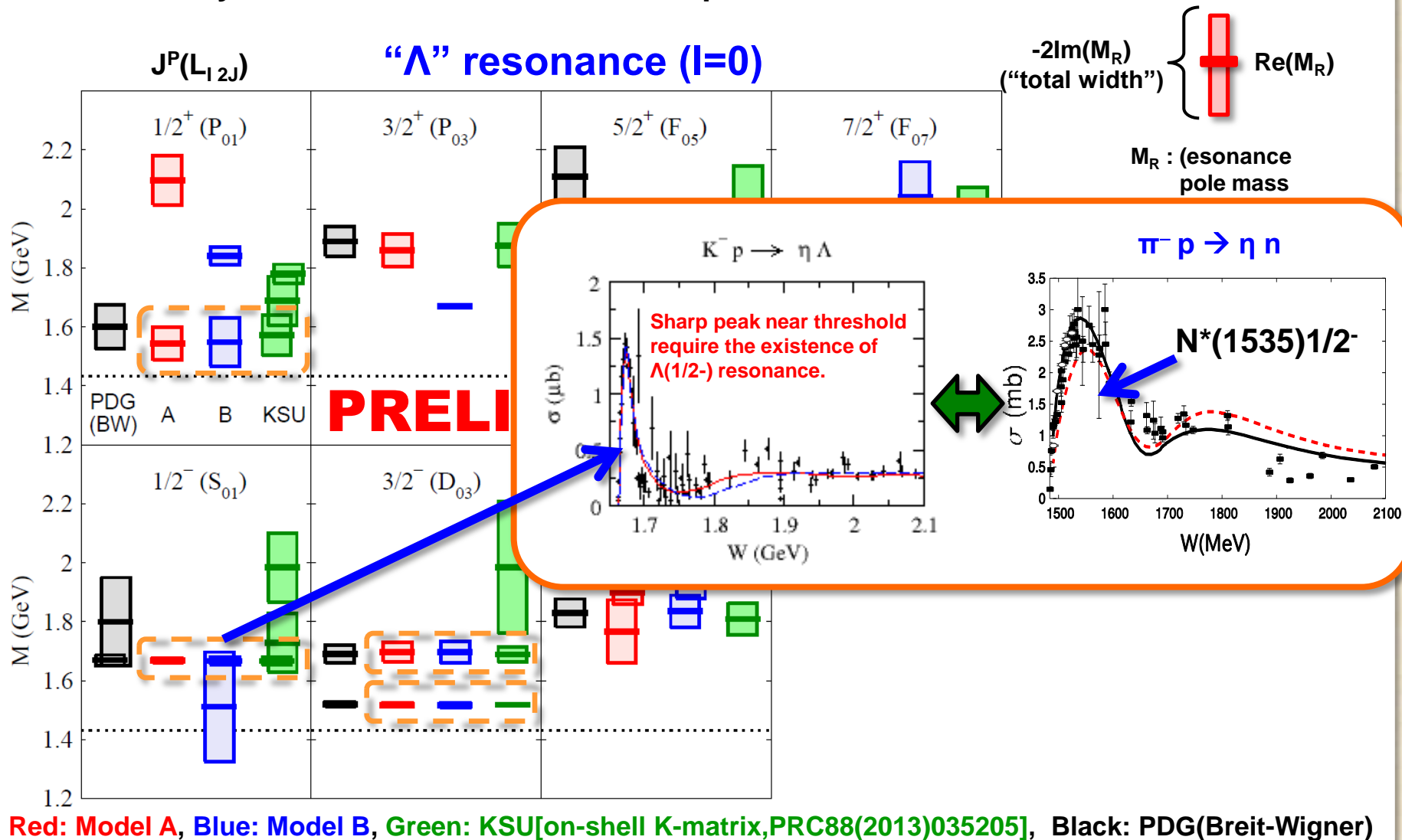


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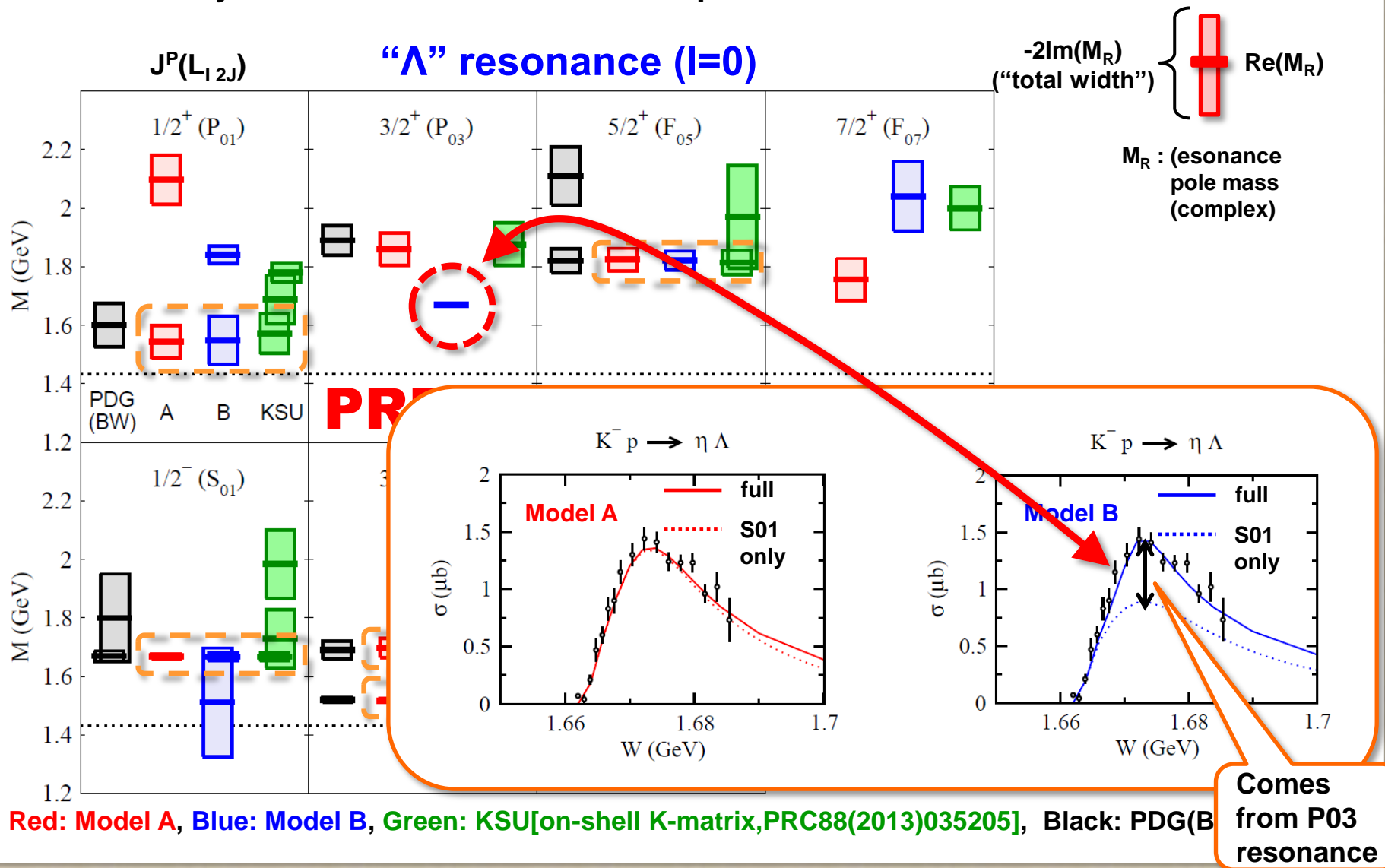
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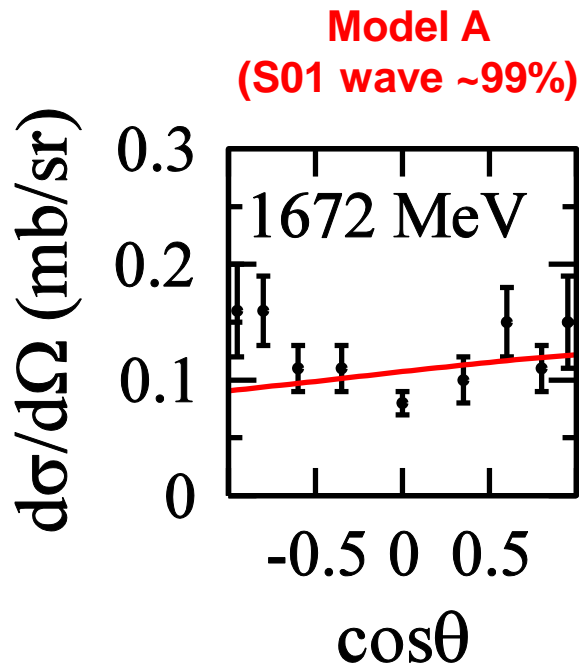


# P03 resonance just above the $\eta\Lambda$ threshold

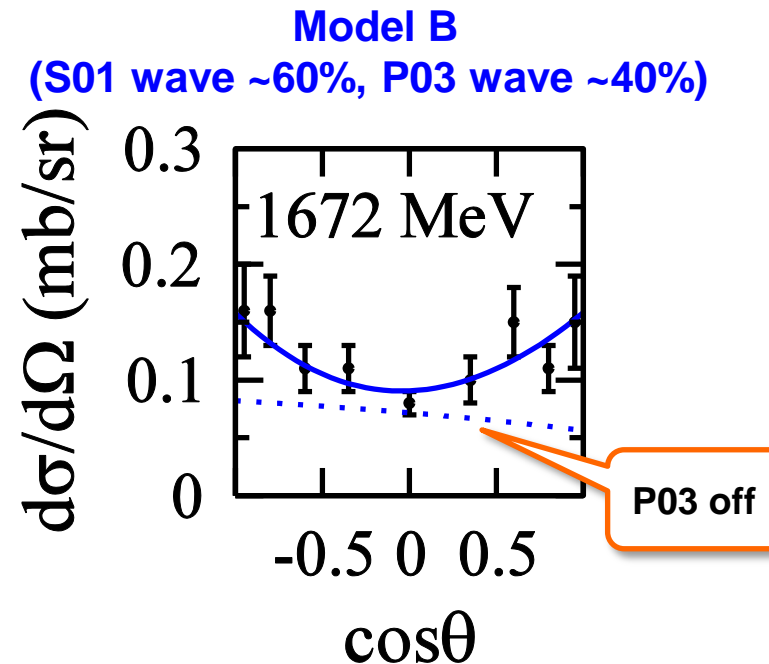
HK, Nakamura, Sato, in preparation

$d\sigma/d\Omega$  of  $K^- p \rightarrow \eta\Lambda$  @  $W=1672$  MeV (just 8 MeV above the threshold)

- Even in the region very close to the threshold, the  $d\sigma/d\Omega$  data show a clear angular dependence.



- ✓ Concave-up behavior of the data is not reproduced.



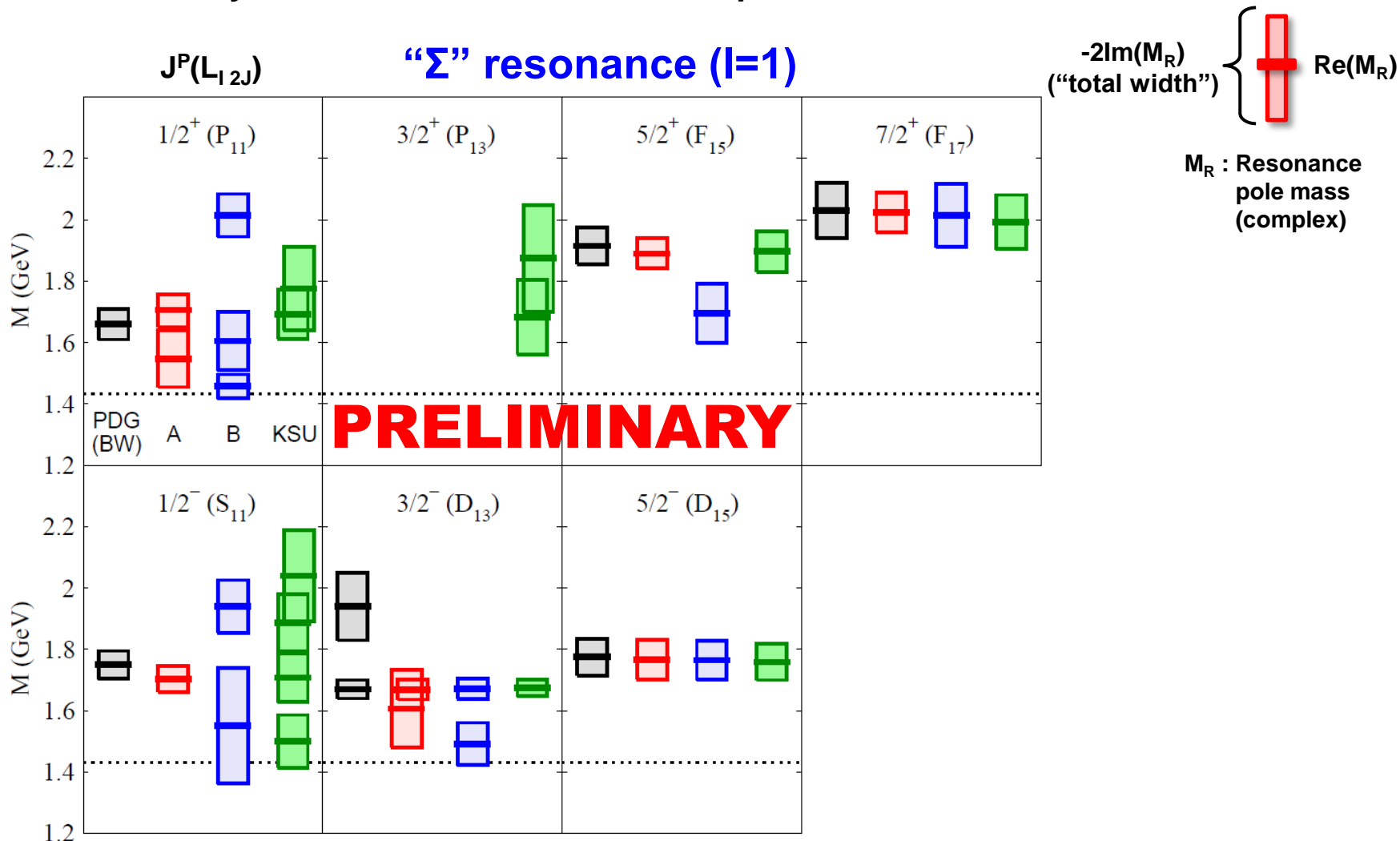
- ✓ **NEW narrow P03 resonance** is responsible for reproducing the angular dependence of  $d\sigma/d\Omega$  !!



# Comparison of $\Sigma^*$ spectrum between multichannel analyses

HK, Nakamura, Sato, in preparation

### Here only  $Y^*$ s ABOVE  $\bar{K}N$  threshold are presented.

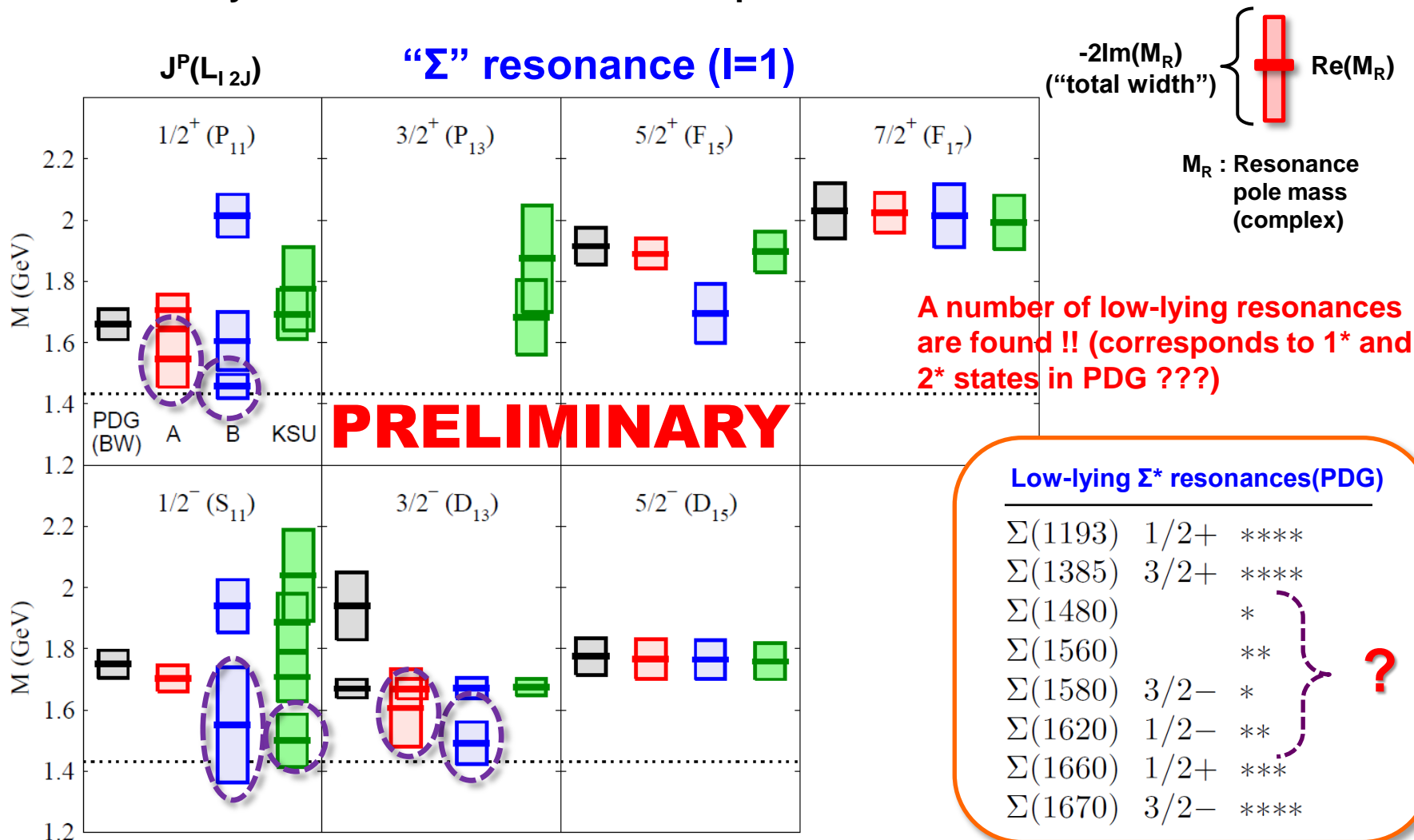


Red: Model A, Blue: Model B, Green: KSU[on-shell K-matrix,PRC88(2013)035205], Black: PDG(Breit-Wigner)

# Comparison of $\Sigma^*$ spectrum between multichannel analyses

HK, Nakamura, Sato, in preparation

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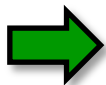


Red: Model A, Blue: Model B, Green: KSU[on-shell K-matrix,PRC88(2013)035205], Black: PDG(Breit-Wigner)

# Summary & Future works

- ✓ Comprehensive analysis of  $K^- p \rightarrow \bar{K}N, \pi\Sigma, \pi\Lambda, \eta\Lambda, K\Xi$  up to  $W = 2.1$  GeV has been accomplished **for the first time** within a dynamical coupled-channels approach.
- ✓ **Partial-wave (S, P, D, and F) amplitudes** &  $\Lambda^*$  and  $\Sigma^*$  mass spectrum have been successfully extracted.

Visible **analysis dependence** exits in extracted values.



Lack of the K- p reaction data for

{ spin-rotation observables ( $\beta, R, A$ )  
the  $\bar{K}N$  threshold region  
3-body ( $\pi\pi\Lambda, \pi \bar{K}N, \dots$ ) production reaction  
...

**J-PARC** is a **unique facility** to overcome this unsatisfactory situation !!

## ✓ Future works:

- $Y^*$  spectroscopy **below  $\bar{K}N$  threshold** with **K- d reactions** (J-PARC E31).
- Multi-strange baryon ( $\Xi^*$  and  $\Omega^*$ ) spectroscopy.
- Application to production reactions of **hypernuclei** and **kaonic nuclei**.

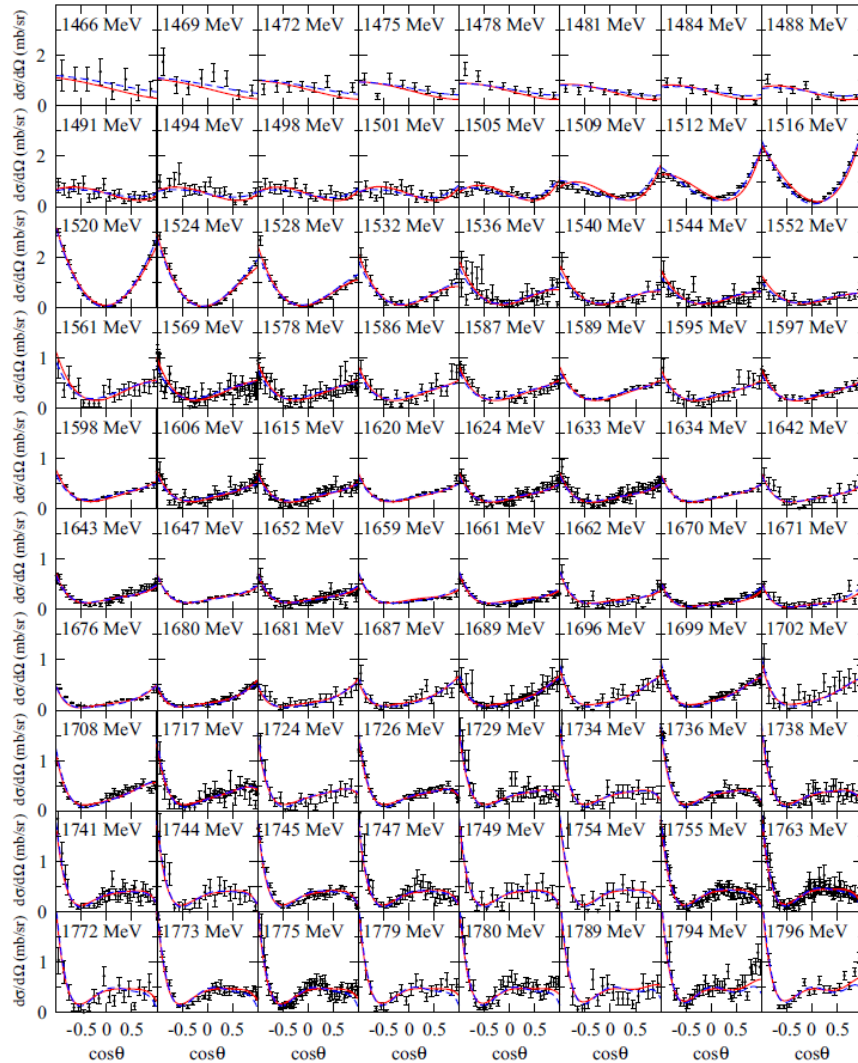
**Back up**

# Results of the fits

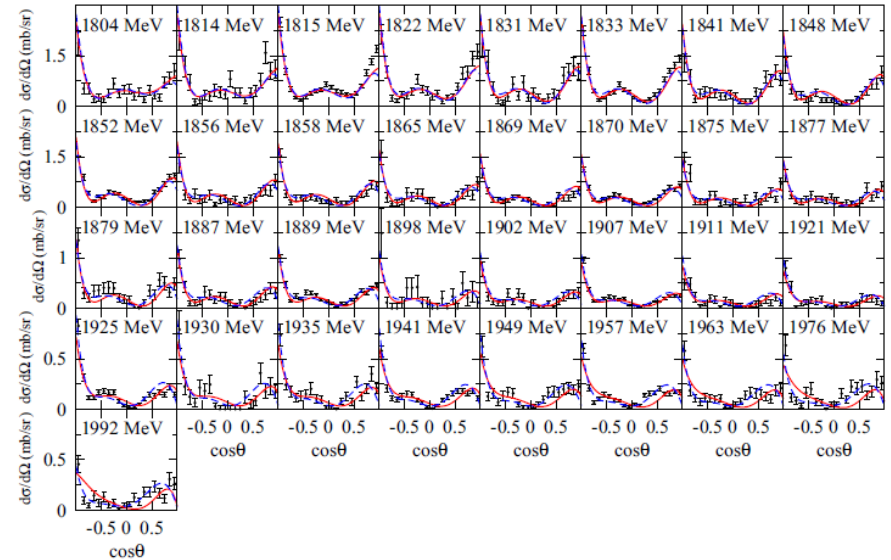
$K^- p \rightarrow K^0 n$  reaction

HK, Nakamura, Lee, Sato, PRC90(2014)065204

$d\sigma/d\Omega$  ( $1466 < W < 1796$  MeV)



$d\sigma/d\Omega$  ( $1804 < W < 1992$  MeV)



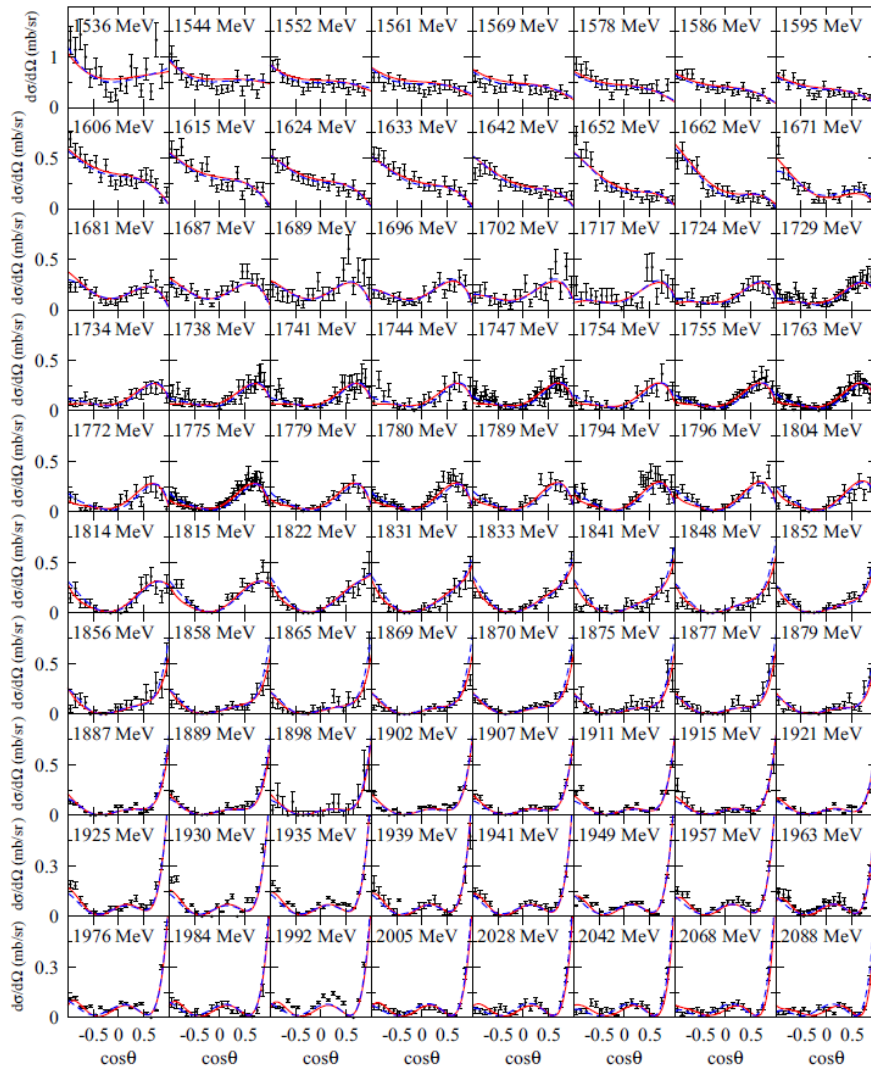
Red: Model A Blue: Model B

# Results of the fits

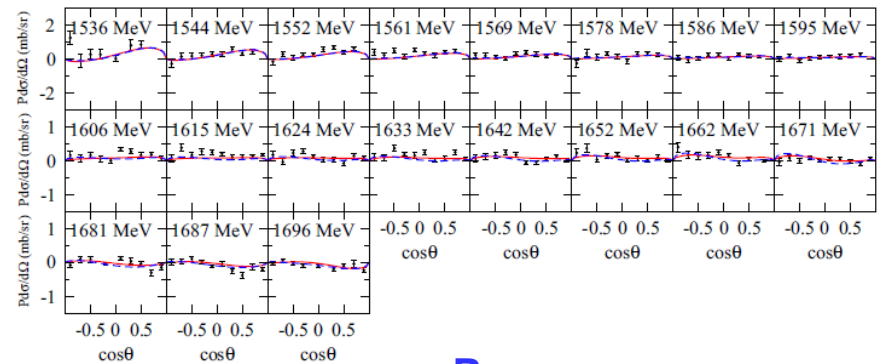
$K^- p \rightarrow \pi^- \Sigma^+$  reaction

HK, Nakamura, Lee, Sato, PRC90(2014)065204

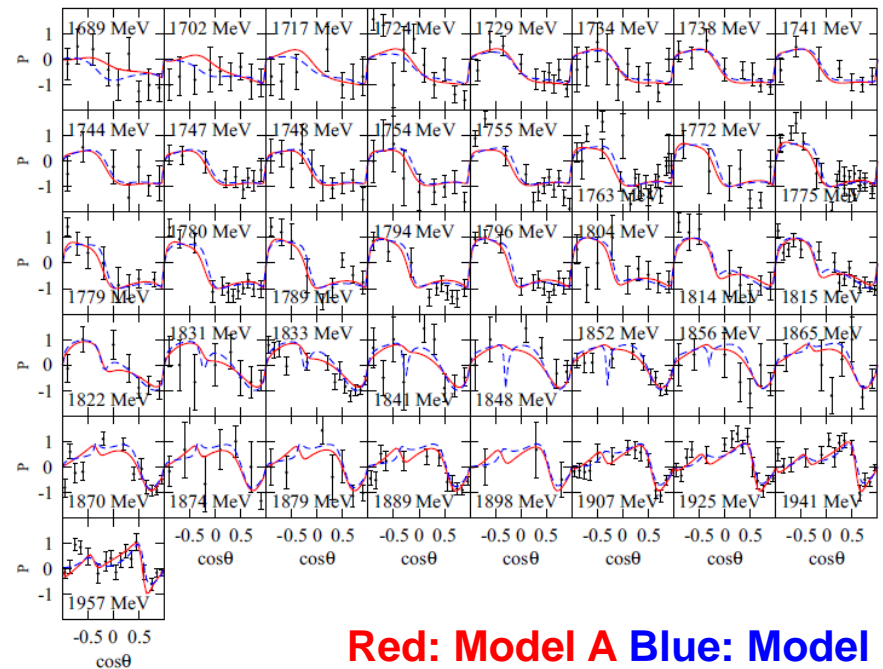
$d\sigma/d\Omega$



$P \times d\sigma/d\Omega$



$P$



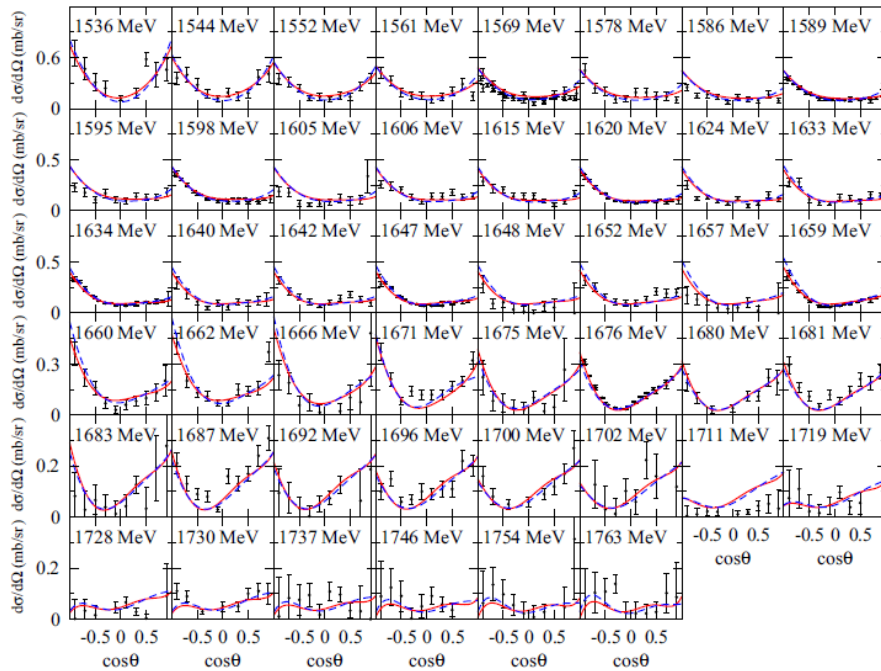
Red: Model A Blue: Model B

# Results of the fits

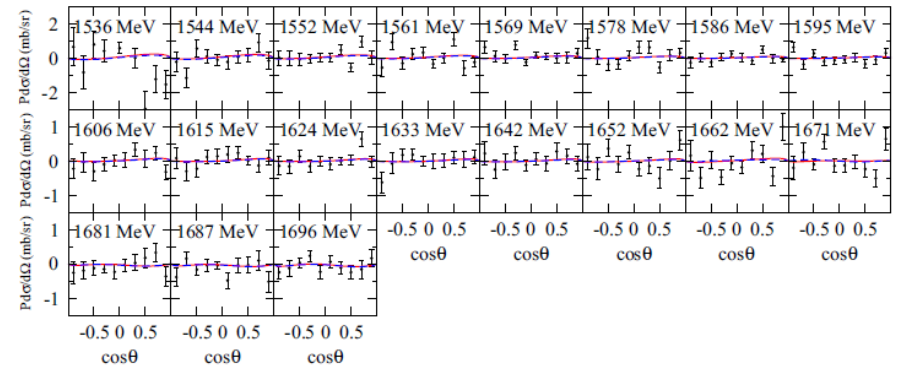
$K^- p \rightarrow \pi^0 \Sigma^0$  reaction

HK, Nakamura, Lee, Sato, PRC90(2014)065204

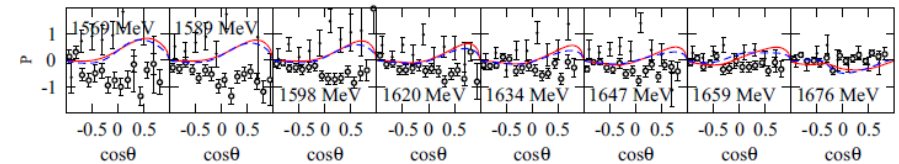
$d\sigma/d\Omega$



$P \times d\sigma/d\Omega$



$P$



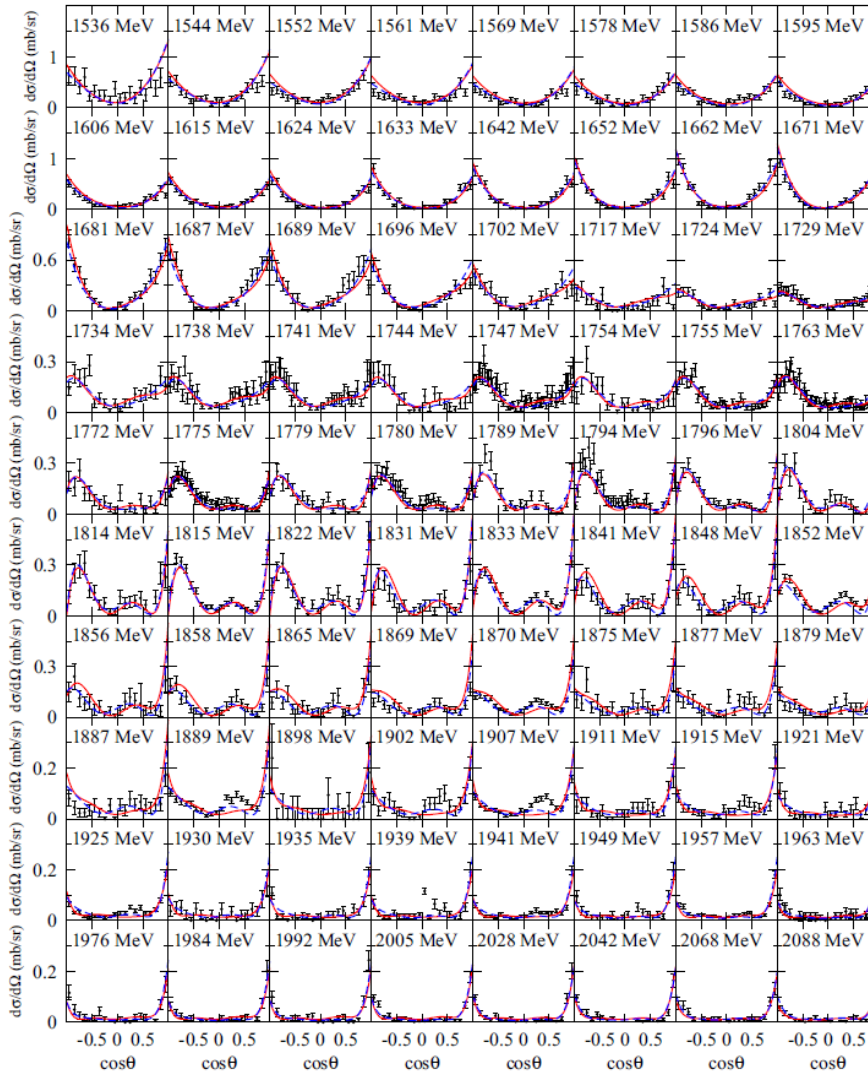
Red: Model A Blue: Model B

# Results of the fits

$K^- p \rightarrow \pi^+ \Sigma^-$  reaction

HK, Nakamura, Lee, Sato, PRC90(2014)065204

$d\sigma/d\Omega$



Red: Model A Blue: Model B

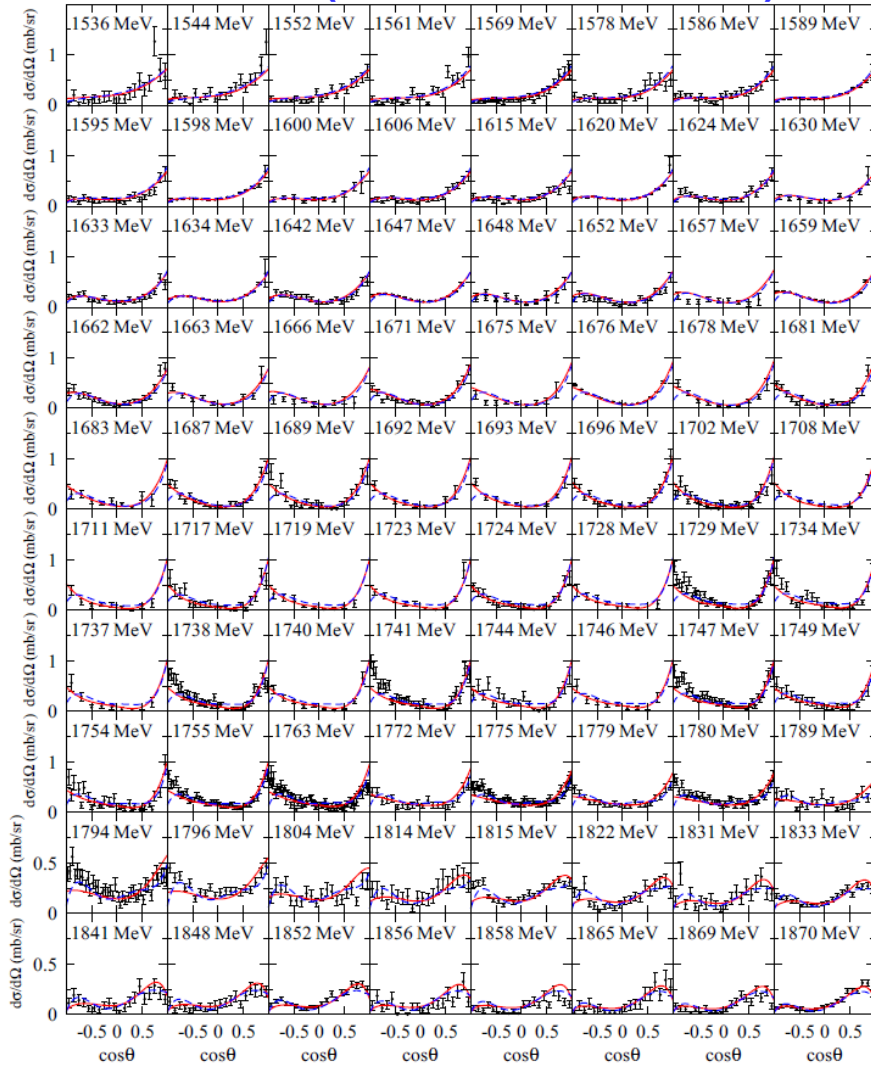


# Results of the fits

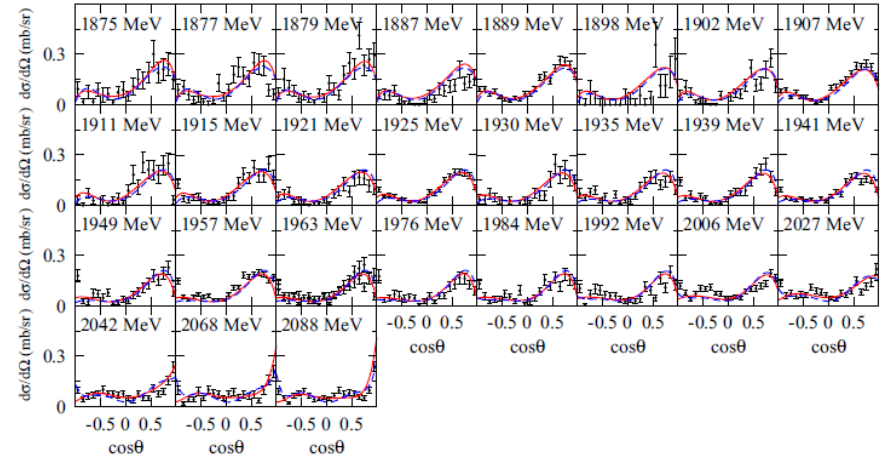
$K^- p \rightarrow \pi^0 \Lambda$  reaction

HK, Nakamura, Lee, Sato, PRC90(2014)065204

$d\sigma/d\Omega$  ( $1536 < W < 1870$  MeV)



$d\sigma/d\Omega$  ( $1875 < W < 2088$  MeV)



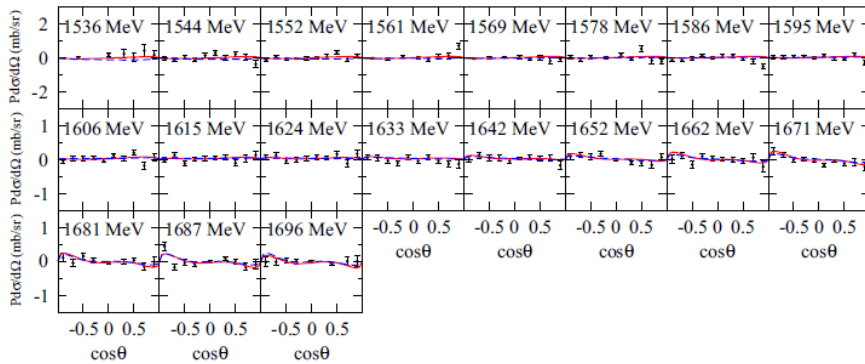
Red: Model A Blue: Model B

# Results of the fits

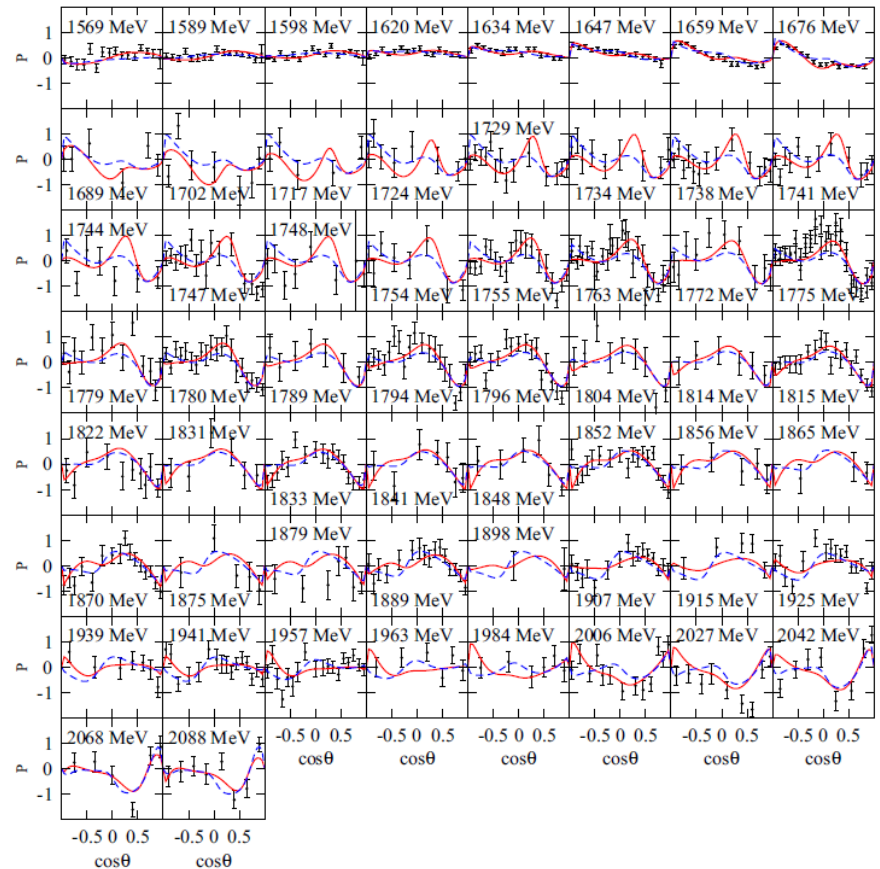
HK, Nakamura, Lee, Sato, PRC90(2014)065204

$K^- p \rightarrow \pi^0 \Lambda$  reaction (cont'd)

$P \times d\sigma/d\Omega$



$P$



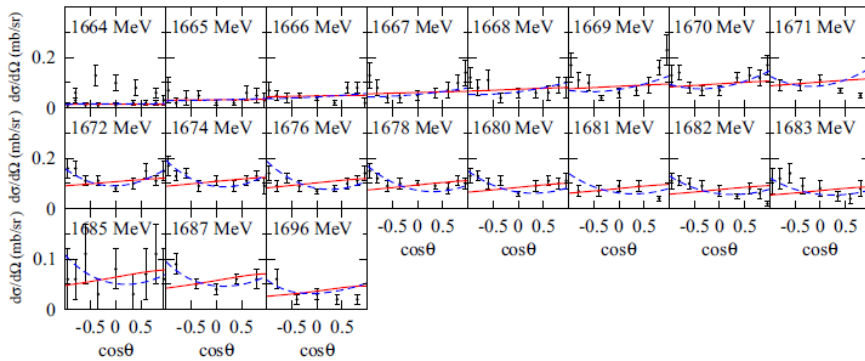
Red: Model A Blue: Model B

# Results of the fits

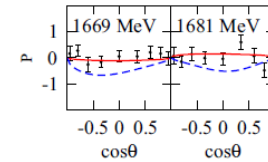
HK, Nakamura, Lee, Sato, PRC90(2014)065204

## $K^- p \rightarrow \eta \Lambda$ reaction

$d\sigma/d\Omega$

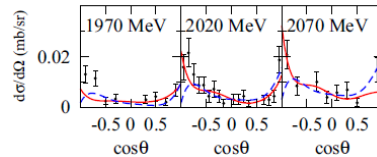


P



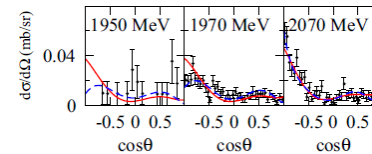
## $K^- p \rightarrow K^0 \Xi^0$ reaction

$d\sigma/d\Omega$



## $K^- p \rightarrow K^+ \Xi^-$ reaction

$d\sigma/d\Omega$



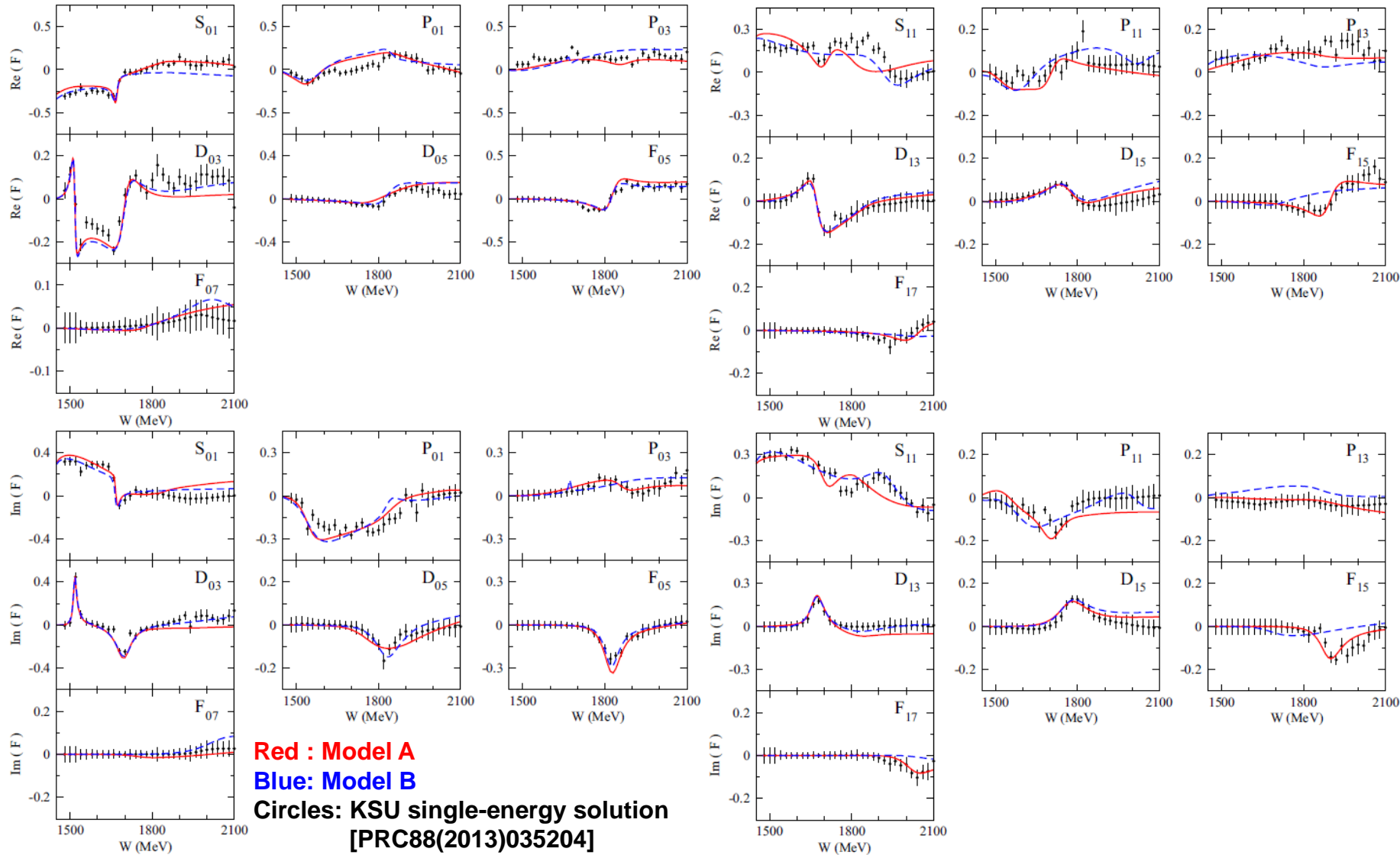
Red: Model A Blue: Model B

# Comparison of extracted partial-wave amplitudes

Extracted  $\bar{K}N \rightarrow \pi\Sigma$  amplitudes

HK, Nakamura, Lee, Sato, PRC90(2014)065204

$L_{I2J}$  :  $L = S, P, \dots$  ;  $I =$  isospin;  $J =$  Total angular mom.

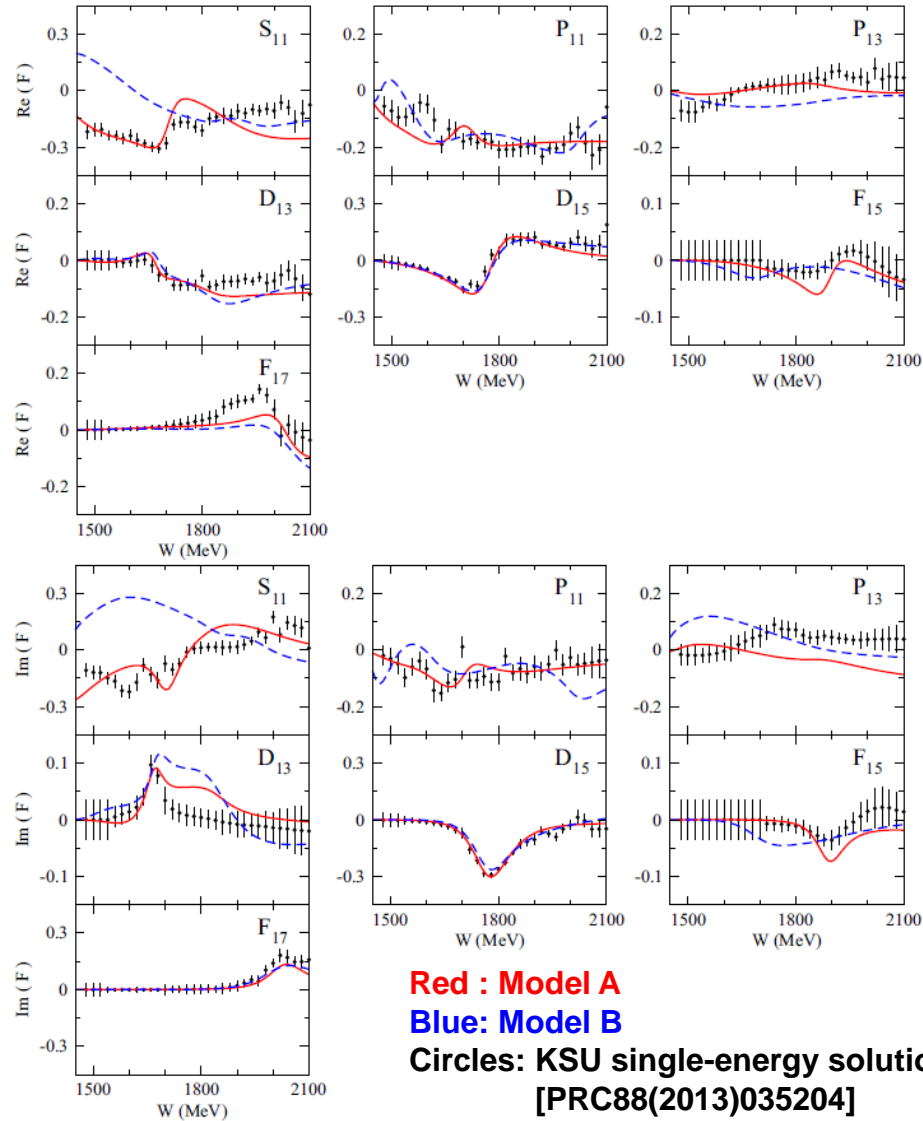


# Comparison of extracted partial-wave amplitudes

Extracted  $\bar{K}N \rightarrow \pi\Lambda$  amplitudes

HK, Nakamura, Lee, Sato, PRC90(2014)065204

$L_{I2J}$  : L = S,P,.. ; I = isospin; J = Total angular mom.

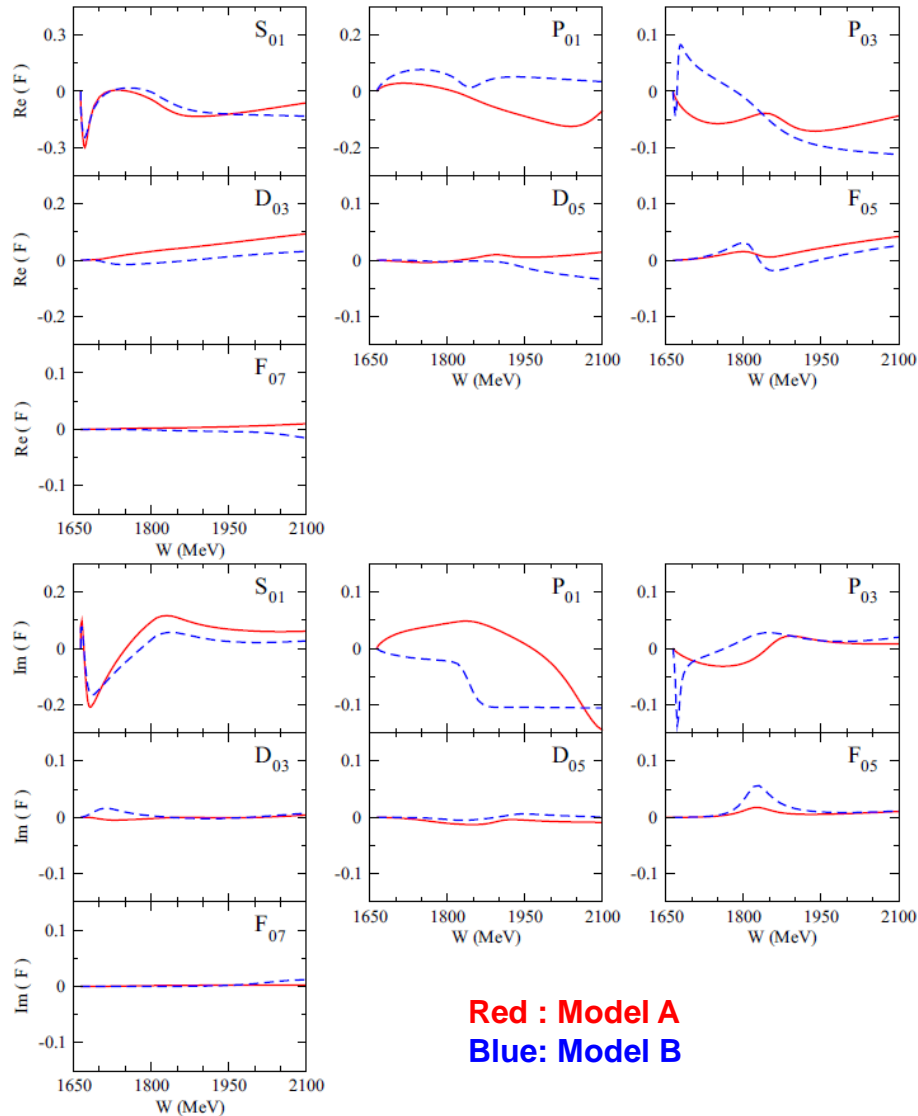


# Comparison of extracted partial-wave amplitudes

Extracted  $\bar{K}N \rightarrow \eta\Lambda$  amplitudes

HK, Nakamura, Lee, Sato, PRC90(2014)065204

$L_{I2J}$  :  $L = S, P, \dots$  ;  $I =$  isospin;  $J =$  Total angular mom.

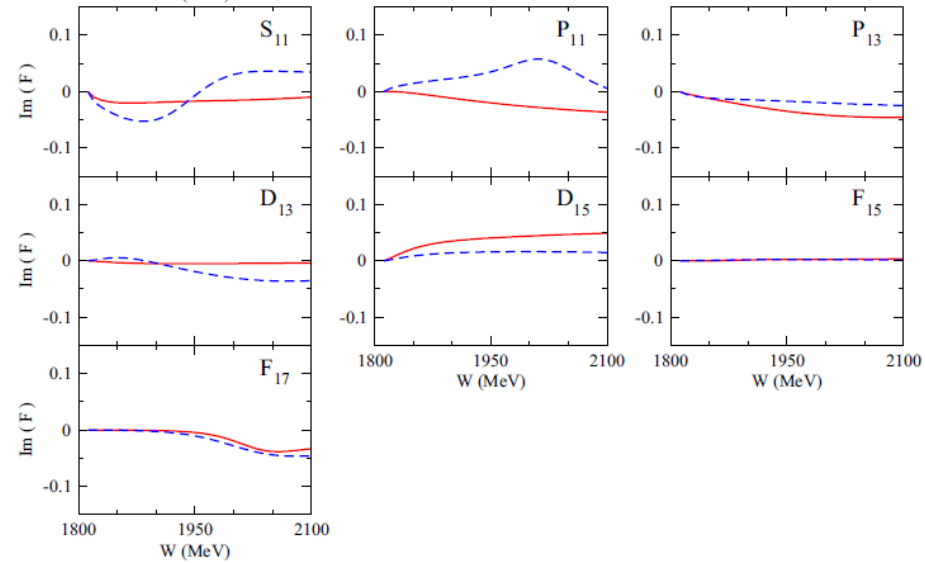
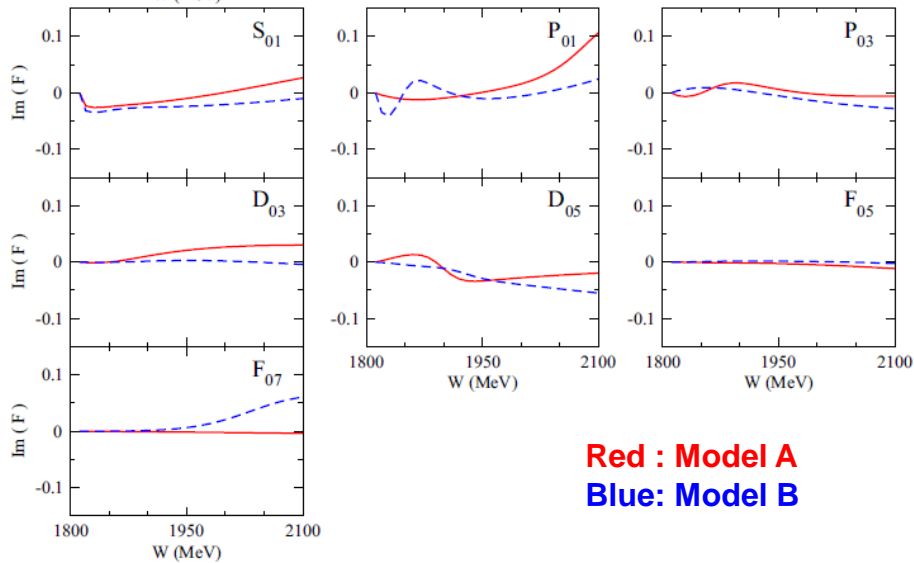
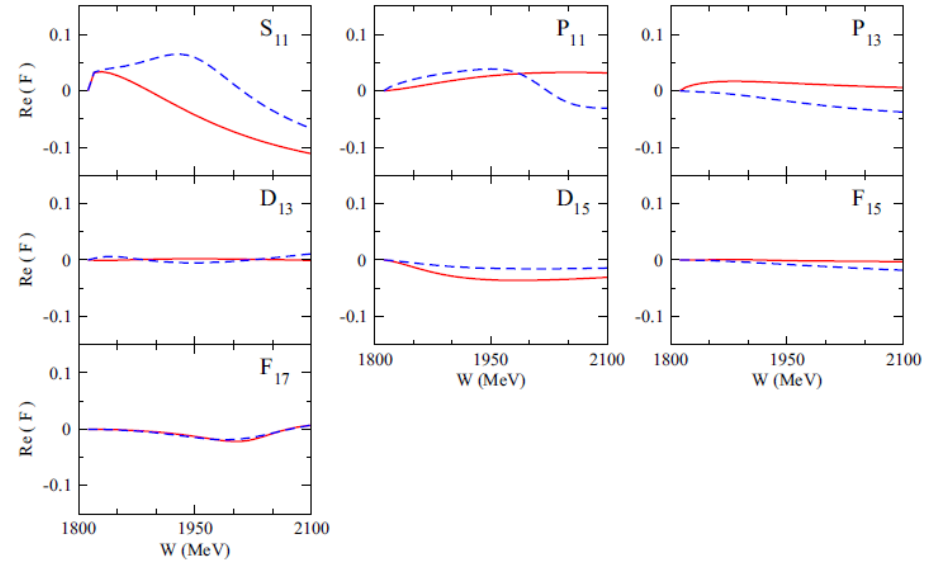
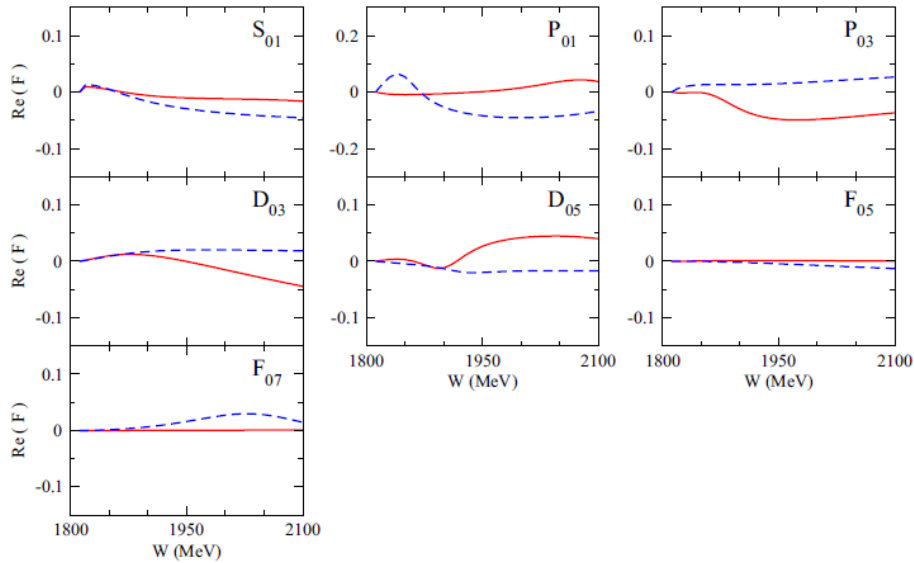


# Comparison of extracted partial-wave amplitudes

## Extracted $\bar{K}N \rightarrow K\bar{E}$ amplitudes

HK, Nakamura, Lee, Sato, PRC90(2014)065204

$L_{12J}$  :  $L = S, P, \dots$  ;  $I =$  isospin;  $J =$  Total angular mom.



Red : Model A  
Blue: Model B

# Extracted scattering lengths and effective ranges

HK, Nakamura, Lee, Sato, PRC90(2014)065204

## Scattering length and effective range

	Model A		Model B	
	$I = 0$	$I = 1$	$I = 0$	$I = 1$
$a_{\bar{K}N}$ (fm)	$-1.37 + i0.67$	$0.07 + i0.81$	$-1.62 + i1.02$	$0.33 + i0.49$
$a_{\eta\Lambda}$ (fm)	$1.35 + i0.36$	-	$0.97 + i0.51$	-
$a_{K\Xi}$ (fm)	$-0.81 + i0.14$	$-0.68 + i0.09$	$-0.89 + i0.13$	$-0.83 + i0.03$
$r_{\bar{K}N}$ (fm)	$0.67 - i0.25$	$1.01 - i0.20$	$0.74 - i0.25$	$-1.03 + i0.19$
$r_{\eta\Lambda}$ (fm)	$-5.67 - i2.24$	-	$-5.82 - i3.32$	-
$r_{K\Xi}$ (fm)	$-0.01 - i0.33$	$-0.42 - i0.49$	$0.13 - i0.20$	$-0.22 - i0.11$

$$a_{K-p} = -0.65 + i0.74 \text{ fm (Model A)}$$

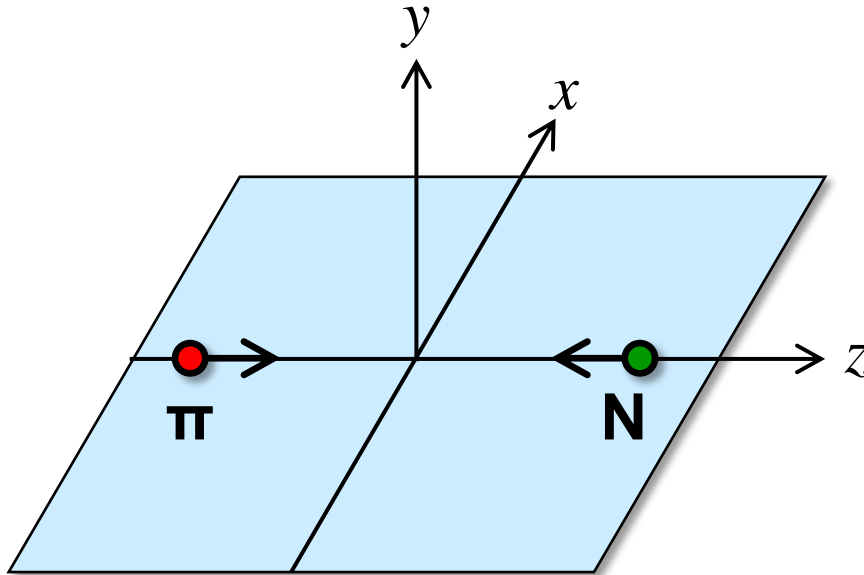
$$a_{K-p} = -0.65 + i0.76 \text{ fm (Model B)}$$



# Polarization observables for spin-0 + spin-1/2 $\rightarrow$ spin-0 + spin-1/2 reactions

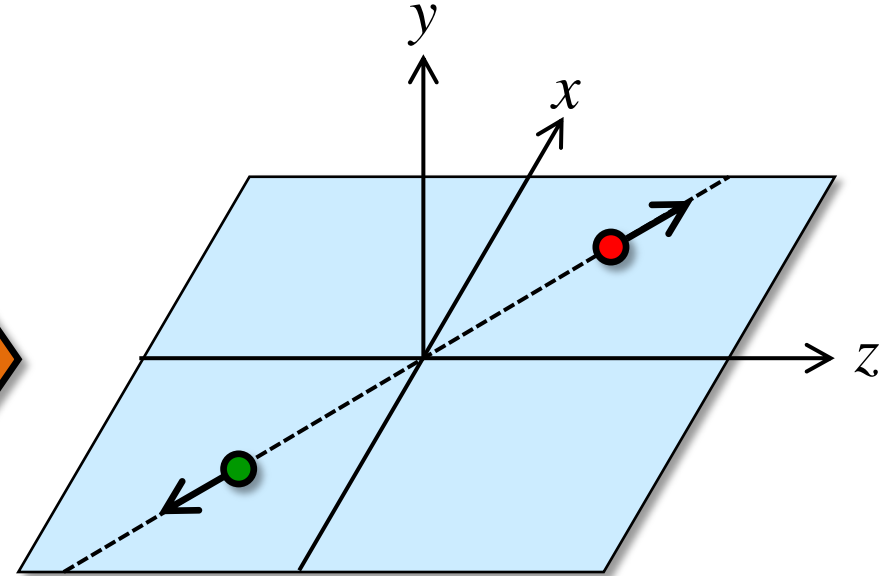
e.g.)  $\pi N$  scattering (in c.m. frame)

Wolfenstein, Phys. Rev. 96(1954)1654  
 Kelly, Sandusky, Cutkosky, PRD10(1974)2309  
 Saxon, Whittaker, Z. Phys. C9(1981)35  
 Supek et al., PRD47(1993)1762



Suppose target nucleon polarization points z-axis  
 ( $\rightarrow$  parallel to pion momentum):

$$\vec{P}_{N_T} = P_{N_T} \hat{z}$$



Polarization of the recoil nucleon is then expressed as

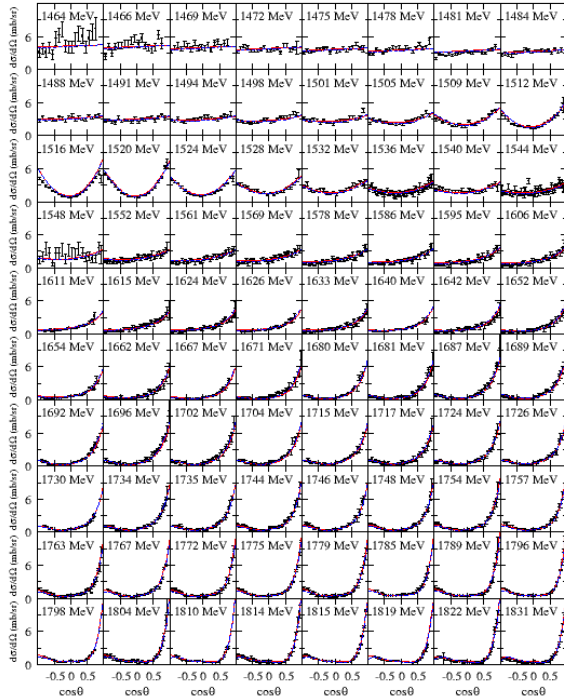
$$\begin{aligned} \vec{P}_{N_R} = & P \hat{y} \\ & + P_{N_T} \sqrt{1 - P^2} \sin \beta \hat{x} \\ & + P_{N_T} \sqrt{1 - P^2} \cos \beta \hat{z} \end{aligned}$$

$$(R = \sqrt{1 - P^2} \sin \beta, \quad A = \sqrt{1 - P^2} \cos \beta)$$

Note: Various conventions are taken for  $\beta$ ,  $R$ ,  $A$  in literatures.

# Extraction of resonance parameters

e.f.)  $d\sigma/d\Omega$  of  $K-p \rightarrow K-p$



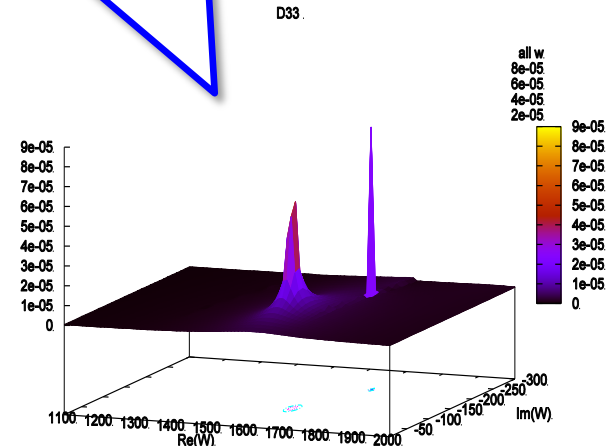
1. Construct model  
(determine amplitudes)  
by fitting data

3. Resonances are obtained as **poles of amplitudes** in complex energy plane !!

mass & width = pole value  
coupling constants = (residue)<sup>1/2</sup> at the pole

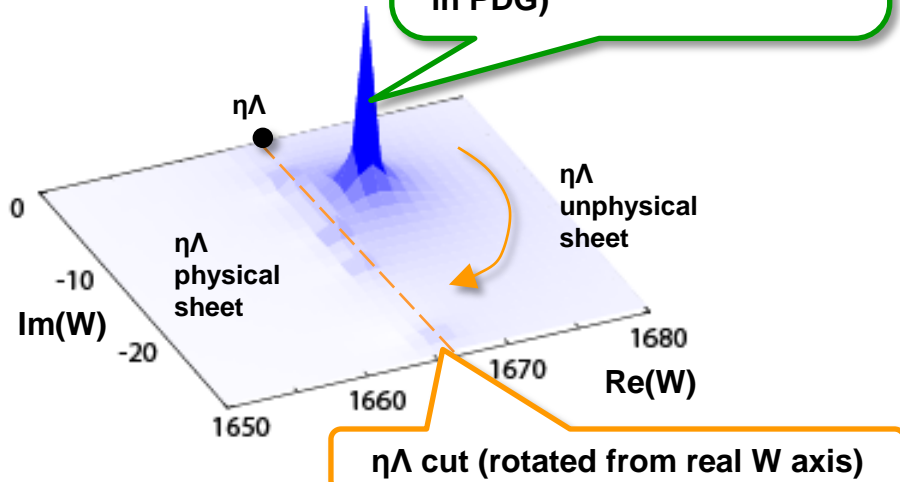
2. Analytic continuation  
of determined amplitudes  
to **complex energy plane** !!

Suzuki et al PRC79(2009)025205  
PRC82(2010)045206



# $\Lambda(1/2^-)$ resonance near the $\eta\Lambda$ threshold

$\Lambda(1/2^-)$  pole at  
**1669 – i9 MeV (Model A)**  
 (would correspond to  $\Lambda(1670)$   
 in PDG)



- ✓ Next higher excited states of  $\Lambda(1405)$ .
- ✓  $K^- p \rightarrow \eta\Lambda$  data make the existence of this  $\Lambda(1/2^-)$  **stable** and **model-independent**.
- ✓ **Dip at  $W \sim 1670$  MeV seen in  $K^- p \rightarrow MB$  total cross sections** is produced by this  $\Lambda(1/2^-)$ . [ $\eta\Lambda$  **cusplike effect is hindered** by the large contribution from  $\Lambda(1/2^-)$ .]

