

Kaon semileptonic transition form factors and the transverse spin density from the instanton vacuum

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Motivation

📌 Kaon semileptonic transition $K^0 \rightarrow \pi^- |V_{us}|$

📌 (semi)Exclusive weak processes

Weak generalized parton distributions

Vector & tensor generalized form factors at spacelike region

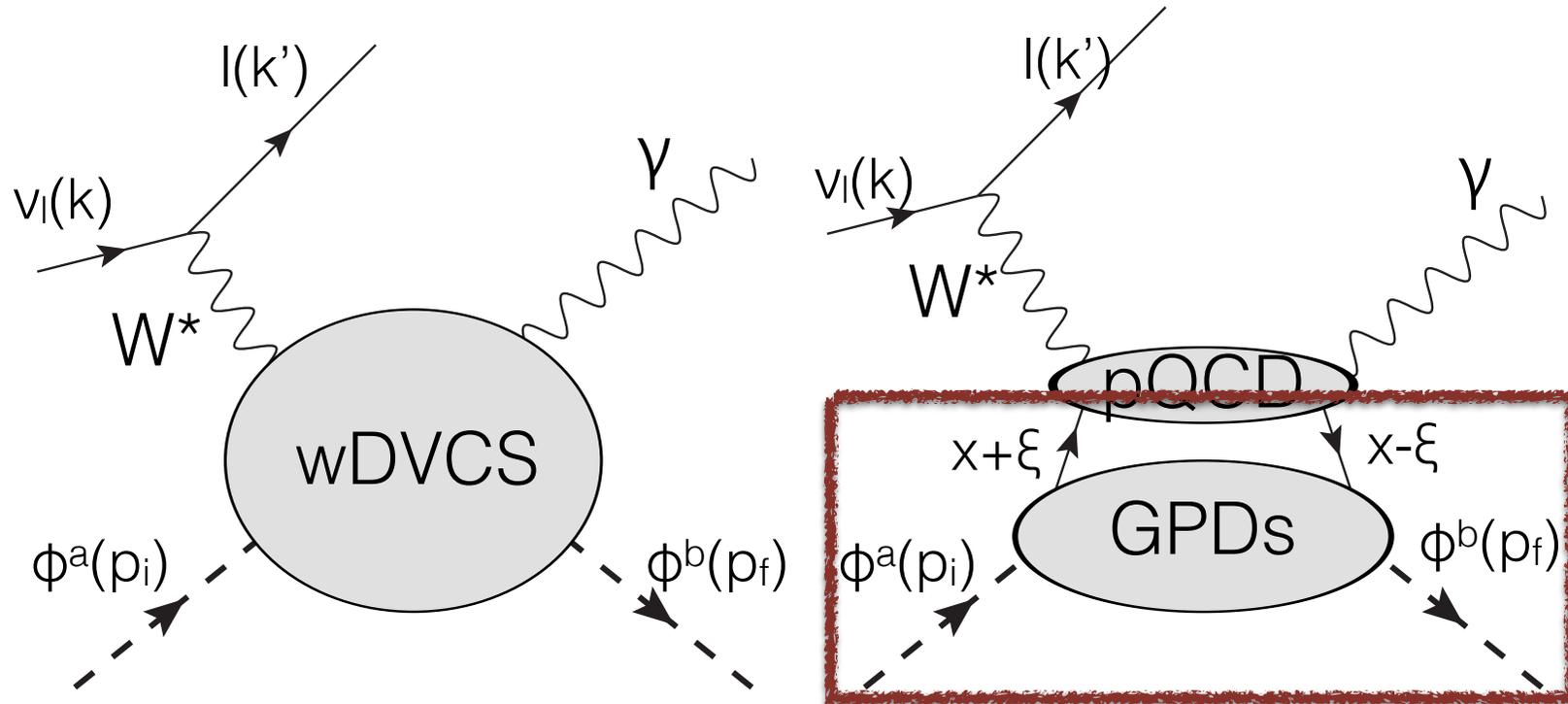
Generalized transverse distributions for the transition

📌 Tensor form factor for the kaon transitions

Relevant lattice study is done [[I. Baum et al, Phys.Rev.D 84, 074503 \(2011\)](#)]

Reveals quark spin structure

Weak DVCS & GPDs



- Light cone frame

$$P^\mu = \frac{1}{2} (p_i^\mu + p_f^\mu), \quad \Delta^\mu = p_f^\mu - p_i^\mu, \quad t = \Delta^2$$

$$n_\pm = \frac{1}{\sqrt{2}} (1, 0, 0, \pm 1)$$

$$v^\pm = \frac{1}{\sqrt{2}} (v^0 \pm v^3), \quad v_\perp = (v^1, v^2)$$

$$\xi = \frac{p_i^+ - p_f^+}{p_i^+ + p_f^+}$$

GPDs: Matrix elements

- Kaon Transition GPDs for $K^0 \rightarrow \pi^-$

$$2P^+ H^{K\pi}(x, \xi, t)$$

$$= \int \frac{d\lambda}{2\pi} e^{ix\lambda(P \cdot n)} \langle \pi^-(p_f) | \bar{s}(-\lambda n/2) \gamma^+ [-\lambda n/2, \lambda n/2] u(\lambda n/2) | K^0(p_i) \rangle$$

$$\frac{P^+ \Delta^j - \Delta^j P^+}{m_K} E_T^{K\pi}(x, \xi, t)$$

$$= \int \frac{d\lambda}{2\pi} e^{ix\lambda(P \cdot n)} \langle \pi^-(p') | \bar{s}(-\lambda n/2) i\sigma^{+j} [-\lambda n/2, \lambda n/2] u(\lambda n/2) | K^0(p) \rangle$$

- Polynomiality

Integrating over x with $x^n \rightarrow$ nth order Generalized Form Factors

GFFs for the Kaon Transition

- Generalized form factors for the kaon transition $K^0 \rightarrow \pi^-$

$$\langle \pi^-(p_f) | \mathcal{O}_V^{\mu\mu_1\cdots\mu_n} | K^0(p_i) \rangle = \mathcal{S} \left[2P^\mu P^{\mu_1} \cdots P^{\mu_n} A_{n+1,0}^{K\pi}(t) \right]$$

$$+ 2 \sum_{i=1, \text{odd}}^n \Delta^\mu \Delta^{\mu_1} \cdots \Delta^{\mu_i} P^{\mu_{i+1}} \cdots P^{\mu_n} A_{n+1,i+1}^{K\pi}(t)$$

$$+ 2 \sum_{i=0, \text{even}}^n \Delta^\mu \Delta^{\mu_1} \cdots \Delta^{\mu_i} P^{\mu_{i+1}} \cdots P^{\mu_n} C_{n+1,i+1}^{K\pi}(t)$$

$$\mathcal{O}_V^{\mu\mu_1\cdots\mu_n} = \mathcal{S} \left[\bar{s} (\gamma^\mu i \overleftrightarrow{D}^{\mu_1}) \cdots (i \overleftrightarrow{D}^{\mu_n}) u \right]$$

GFFs for the Kaon Transition

- Generalized form factors for the kaon transition $K^0 \rightarrow \pi^-$

$$\langle \pi^-(p_f) | \mathcal{O}_T^{\mu\nu\mu_1\cdots\mu_{n-1}} | K^0(p_i) \rangle =$$

$$\mathcal{AS} \left[\frac{(P^\mu \Delta^\nu - \Delta^\mu P^\nu)}{m_K} \sum_{i=\text{even}}^{n-1} \Delta^{\mu_1} \cdots \Delta^{\mu_i} P^{\mu_{i+1}} \cdots P^{\mu_{n-1}} B_{T n, i}^{K\pi}(t) \right]$$

.....

$$\mathcal{O}_T^{\mu\nu\mu_1\cdots\mu_{n-1}} = \mathcal{AS} \left[\bar{s} \sigma^{\mu\nu} (i \overleftrightarrow{D}^{\mu_1}) \cdots (i \overleftrightarrow{D}^{\mu_{n-1}}) u \right]$$

GFFs for the Kaon Transition

- n=0 generalized form factors for the kaon transition $K^0 \rightarrow \pi^-$

$$\langle \pi^-(p_f) | \bar{s} \gamma_\mu u | K^0(p_i) \rangle = 2P_\mu A_{1,0}^{K\pi}(t) + 2\Delta_\mu C_{1,1}^{K\pi}(t)$$

$$\langle \pi^-(p_f) | \bar{s} \sigma_{\mu\nu} u | K^0(p_i) \rangle = \left(\frac{P_\mu \Delta_\nu - P_\nu \Delta_\mu}{m_K} \right) B_T^{K\pi}(t)$$

.....

$$\langle \pi^-(p_f) | \bar{s} \gamma_\mu u | K^0(p_i) \rangle = (p_f + p_i)_\mu f_+(t) + (p_f - p_i)_\mu f_-(t)$$

$$A_{1,0}^{K\pi} = f_+, \quad C_{1,1}^{K\pi} = \frac{1}{2} f_-$$

Polynomiality

- Mellin moments of the $K^0 \rightarrow \pi^-$ transition GPDs

$$\int dx x^n H^{K\pi}(x, \xi, t) = A_{n+1,0}^{K\pi}(t) + \sum_{i=1,\text{odd}}^n (-2\xi)^{i+1} A_{n+1,i+1}^{K\pi}(t) + \sum_{i=1,\text{odd}}^{n+1} (-2\xi)^i C_{n+1,i}^{K\pi}(t)$$

$$\int dx x^n E_T^{K\pi}(x, \xi, t) = \sum_{i=0,\text{even}}^n (-2\xi)^i B_{T,n+1,i}^{K\pi}(t)$$

- $n=0$

$$\int dx H^{K\pi}(x, \xi, t) = A_{1,0}^{K\pi} - 2\xi C_{1,1}^{K\pi}, \quad \int dx E_T^{K\pi}(x, \xi, t) = B_{T,1,0}^{K\pi}(t)$$

Nonlocal Chiral Quark Model

$$S_{\text{eff}} = -N_c \text{Tr} \log \left[i\not{\partial} + i\hat{m} + i\sqrt{M(i\partial)}U\gamma^5\sqrt{M(i\partial)} \right]$$

- ❖ The chiral effective action
derived from the instanton vacuum

$$U\gamma^5 = \exp \left[\frac{i\gamma^5}{f_\phi} (\lambda \cdot \phi) \right]$$

- ❖ No free parameter
 - Average Instanton size & separation

$$\bar{\rho} \approx \frac{1}{3} \text{ fm} \quad \bar{R} \approx 1 \text{ fm}$$

- ❖ Nonlocality
 - Momentum-dependent dynamical quark mass

- ❖ Nicely reproduces pion properties: F_π , EMFF

- ❖ Explicit SU(3) symmetry breaking

$$\hat{m} = \text{diag}(m_u, m_d, m_s), \quad m_u = m_d = 5 \text{ MeV}, \quad m_s = 150 \text{ MeV}$$

[D. Diakonov, Instantons at work, arXiv:hep-ph/0212026v4]

Nonlocal Chiral Quark Model

- ❖ Momentum-dependent dynamical quark mass

$$\sqrt{M(i\partial)} = \sqrt{M_0 f(m) F^2(i\partial)}$$

$$F(k) = \frac{k}{\Lambda} \left[I_0\left(\frac{k}{2\Lambda}\right) K_1\left(\frac{k}{2\Lambda}\right) - I_1\left(\frac{k}{2\Lambda}\right) K_0\left(\frac{k}{2\Lambda}\right) - \frac{2\Lambda}{k} I_1\left(\frac{k}{2\Lambda}\right) K_1\left(\frac{k}{2\Lambda}\right) \right]$$

$$F_N(k) = \left(\frac{2N\Lambda^2}{2N\Lambda^2 + k^2} \right)^N \quad \Lambda = 1/\bar{\rho} = 600 \text{ MeV}$$

- ❖ Current quark mass correction $f(m)$

$$f(m) = \sqrt{1 + \frac{m^2}{d^2}} - \frac{m}{d}, \quad d \approx 198 \text{ MeV}$$

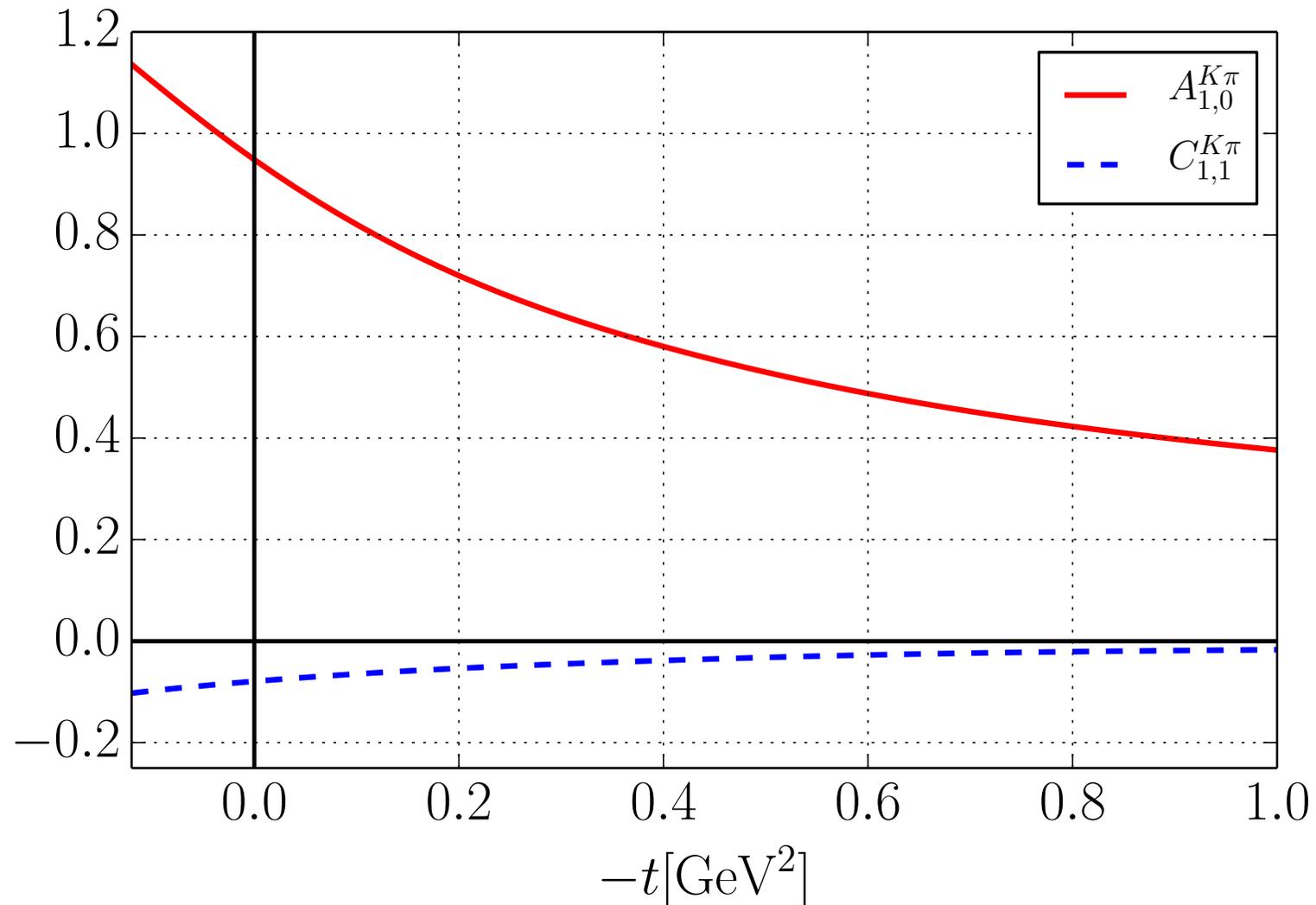
- ❖ Dynamical quark mass at $k=0$

$$M_0 \approx 350 \text{ MeV}$$

[M. Musakhanov Eur.Phys.J.C9,235(1999)]

Calculation of the Vector Form Factors

$$\langle \pi^-(p_f) | \bar{s} \gamma_\mu u | K^0(p_i) \rangle = 2P_\mu A_{1,0}^{K\pi}(t) + 2\Delta_\mu C_{1,1}^{K\pi}(t)$$



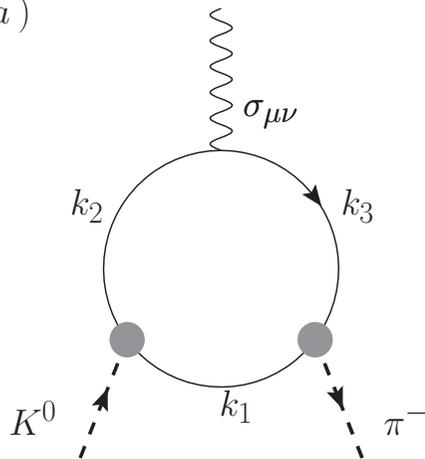
[Nam, S.-I., & Kim, H. C. (2007) Phys. Rev. D, 75(9), 094011.]

Calculation of the Tensor Form Factor

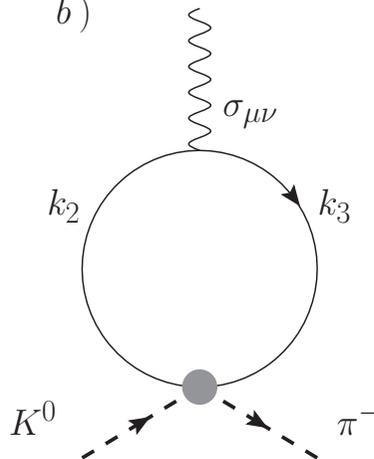
$$\langle \pi^-(p_f) | \bar{s} \sigma_{\mu\nu} u | K^0(p_i) \rangle =$$

$$\frac{8N_c}{f_\pi f_K} \int \frac{d^4l}{(2\pi)^4} \left[\frac{M_{2d} \sqrt{M_{1u}} \sqrt{M_{3s}}}{G_{u1} G_{d2} G_{s3}} \left(k_{1\mu} (k_{2\nu} \bar{M}_{3s} - k_{3\nu} \bar{M}_{2d}) + k_{2\mu} (k_{3\nu} \bar{M}_{1u} - k_{1\nu} \bar{M}_{3s}) \right. \right. \\ \left. \left. + k_{3\mu} (k_{1\nu} \bar{M}_{2d} - k_{2\nu} \bar{M}_{1u}) \right) - \frac{\sqrt{M_{1u}} \sqrt{M_{3s}}}{2G_{u1} G_{s3}} (k_{3\mu} k_{1\nu} - k_{3\nu} k_{1\mu}) \right]$$

a)



b)



$$G_{fi} = k_i^2 + \bar{M}_{if}^2$$

$$\bar{M}_{if} = m_f + M(k_i, m_f)$$

$$k_1 = l - \frac{p_i}{2} - \frac{q}{2} \quad k_2 = l + \frac{p_i}{2} - \frac{q}{2}$$

$$k_3 = l + \frac{p_i}{2} + \frac{q}{2}$$

Pseudoscalar meson decay
constants and masses

$$f_\pi = 93 \text{ MeV}, \quad f_K = 113 \text{ MeV}$$

$$m_\pi = 140 \text{ MeV}, \quad m_K = 495 \text{ MeV}$$

QCD RG Evolution for the tensor form factor

❖ Next-to-leading order

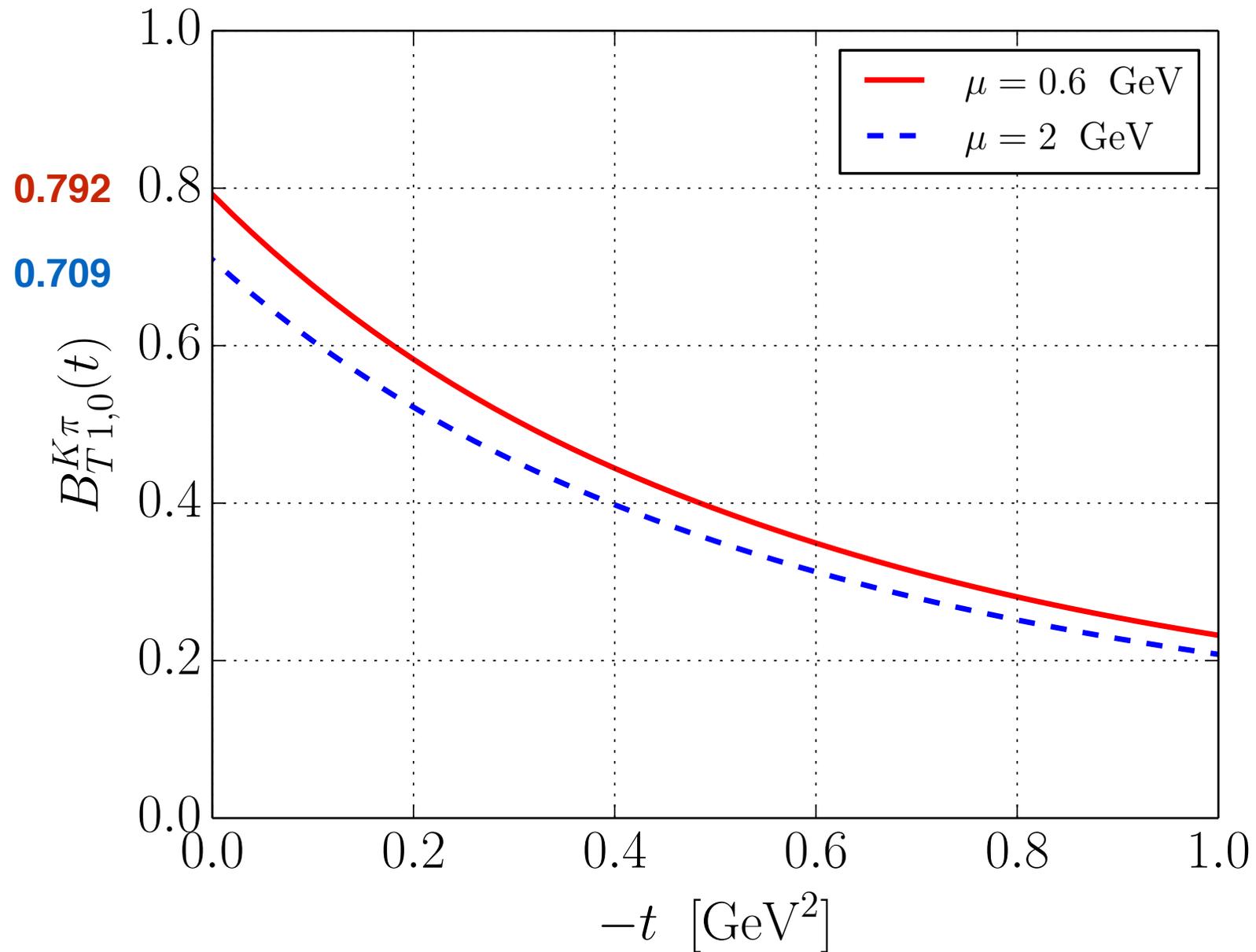
$$B_{T1,0}^{K\pi}(\mu^2) = \left(\frac{\alpha_s(\mu^2)}{\alpha_s(\mu_i^2)} \right)^{4/27} \left[1 - \frac{337}{486\pi} (\alpha_s(\mu_i^2) - \alpha_s(\mu^2)) \right] B_{T1,0}^{K\pi}(\mu_1)$$

$$\alpha_s^{\text{NLO}}(\mu^2) = \frac{4\pi}{9 \ln(\mu^2/\Lambda_{\text{QCD}}^2)} \left[1 - \frac{64}{81} \frac{\ln \ln(\mu^2/\Lambda_{\text{QCD}}^2)}{\ln(\mu^2/\Lambda_{\text{QCD}}^2)} \right]$$

$$\Lambda_{\text{QCD}} = 0.25 \text{ GeV}, \quad N_f = 3$$

[Glück et al. Zeits.Für.Phys. C.67, 433
Barone et al. Phys.Repts., 359, 1.]

Numerical Results



Comparison to Lattice Result

📌 Present work @ $\mu = 2 \text{ GeV}$

$$f_T^{K\pi}(t) = \frac{m_K + m_\pi}{2m_K} B_{T1,0}^{K\pi}(t)$$

$$B_{T1,0}^{K\pi}(0) = 0.71 \rightarrow f_T^{K\pi}(0) = 0.45$$

.....

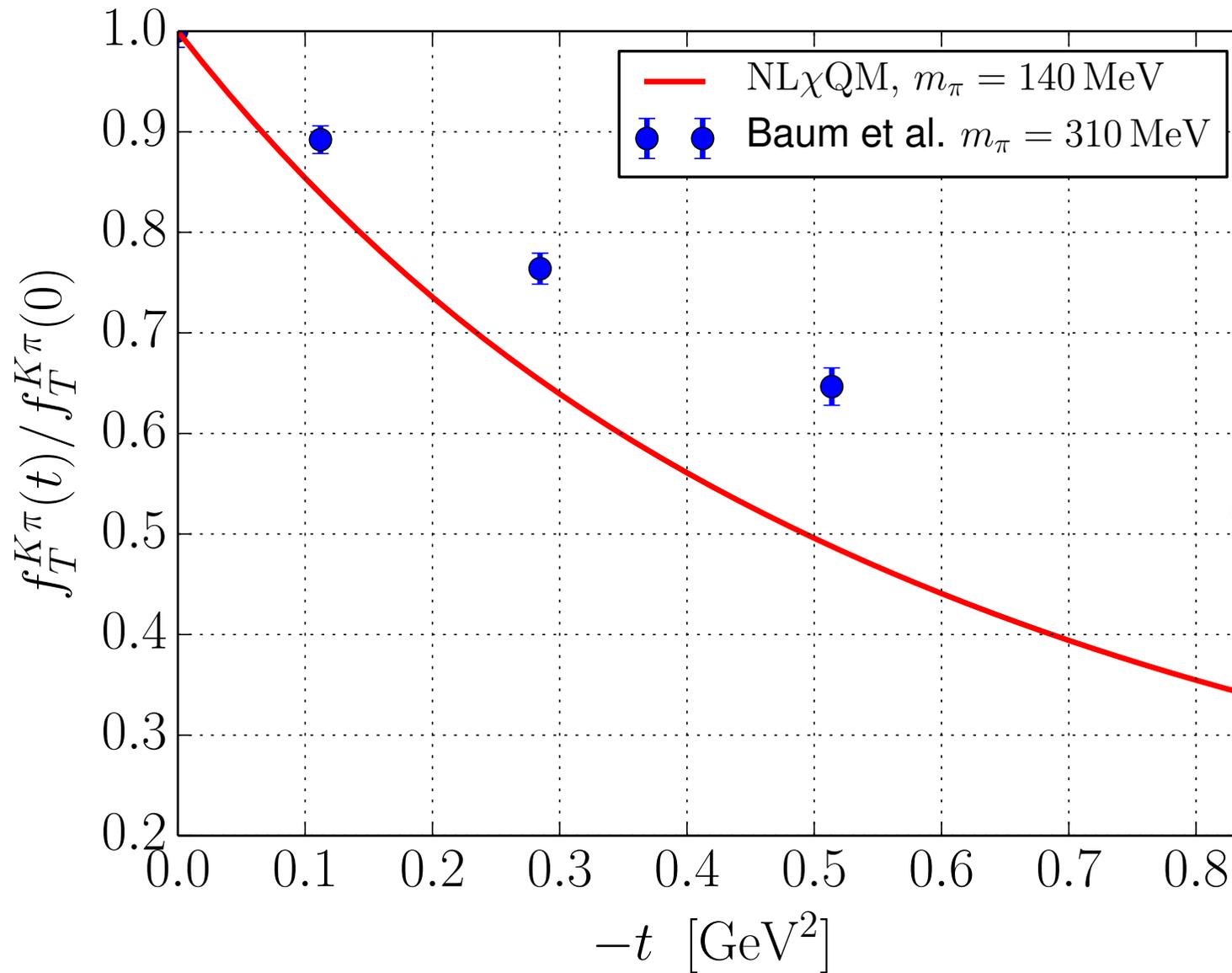
📌 I. Baum et al, $\mu = 2 \text{ GeV}$

$$\langle \pi^0 | \bar{s} \sigma^{\mu\nu} d | K^0 \rangle = (p_\pi^\mu p_K^\nu - p_\pi^\nu p_K^\mu) \frac{\sqrt{2} f_T^{K\pi}(q^2)}{M_K + M_\pi}$$

$$f_T^{K\pi}(0) = 0.417 (14_{\text{stat}}) (5_{\text{syst}}) = 0.417 (15)$$

(Extrapolated to Physical meson masses)

Comparison to Lattice Result



p-pole Parametrization

is this slide necessary?



p-pole parametrization

$$F(t) = \frac{F(0)}{\left(1 - \frac{t}{pM^2}\right)^p} \quad A_{1,0}^{K\pi}(t) \quad B_{T1,0}^{K\pi}(t)$$

p

1.31

2.2

M [GeV]

0.85

0.78

t=0

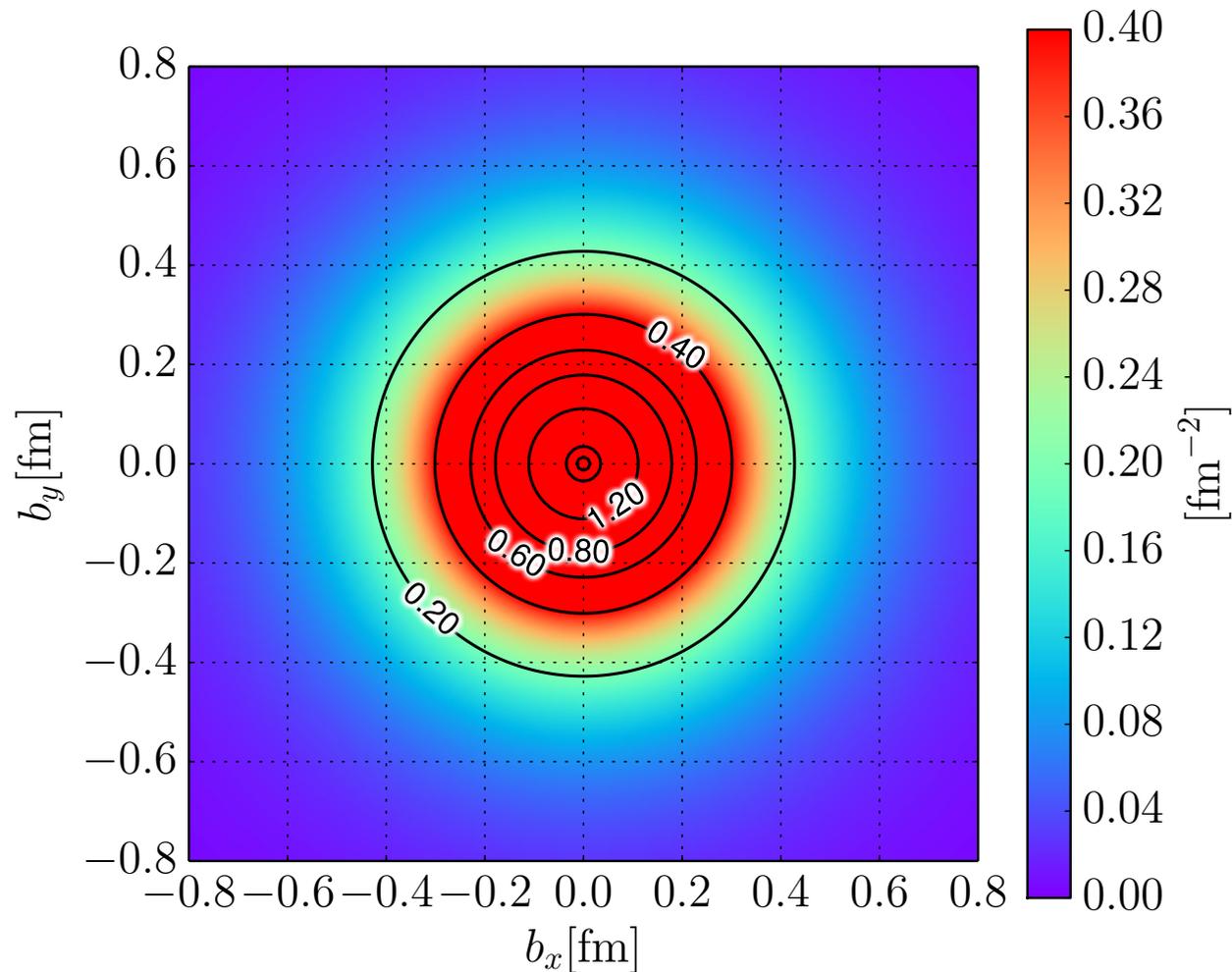
0.95

0.71

Transverse Charge Density

- Transverse charge density for the kaon transition $K^0 \rightarrow \pi^-$

$$\rho_1^{K\pi} = \int \frac{d^2 \Delta}{(2\pi)^2} e^{-i\mathbf{b}_\perp \cdot \Delta} \int dx H^{K\pi}(x, \xi = 0, t) = \frac{1}{(2\pi)^2} \int d^2 \Delta e^{-i\mathbf{b}_\perp \cdot \Delta} A_{1,0}^{K\pi}(t)$$



Transverse Quark Spin Density

- Quarks with definite transverse polarization \mathbf{s}

$$\frac{1}{2} \bar{\psi} [\gamma^+ - s^j i \sigma^{+j} \gamma_5] \psi \quad \vdots \quad \sigma^{\mu\nu} \gamma_5 = -\frac{1}{2} \epsilon^{\mu\nu\alpha\beta} i \sigma_{\alpha\beta}$$

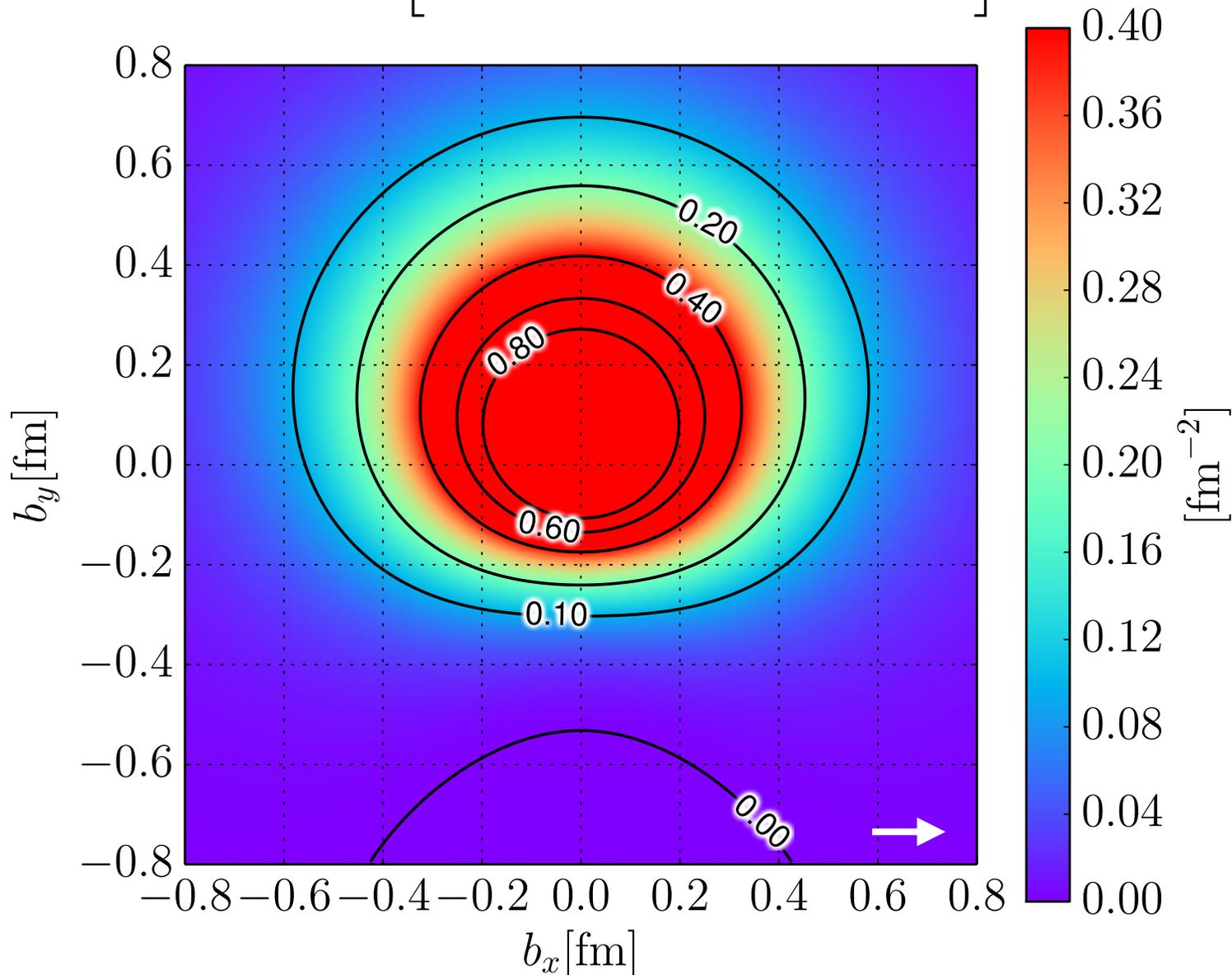
- Transverse quark spin density, $\xi=0$

$$\rho_1^{K\pi}(b, \mathbf{s}_\perp) = \frac{1}{2} \left[A_{1,0}^{K\pi}(b^2) - \frac{s_\perp^i \epsilon^{ij} b^j}{m_K} \frac{\partial B_{T1,0}^{K\pi}(b^2)}{\partial b^2} \right]$$

How the quark with polarized spin is distributed in the transverse plane during the K- π transition process

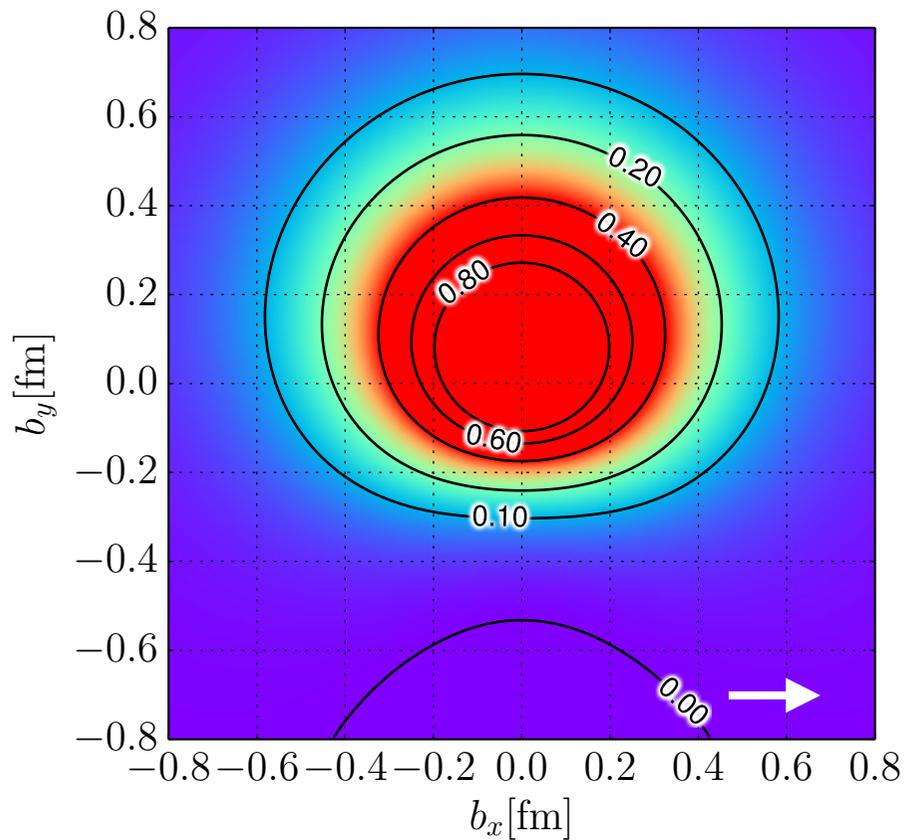
Transverse Quark Spin Density

$$\rho_1^{K\pi}(b, \mathbf{s}_\perp) = \frac{1}{2} \left[A_{1,0}^{K\pi}(b^2) - \frac{s_\perp^i \epsilon^{ij} b^j}{m_K} \frac{\partial B_{T1,0}^{K\pi}(b^2)}{\partial b^2} \right]$$

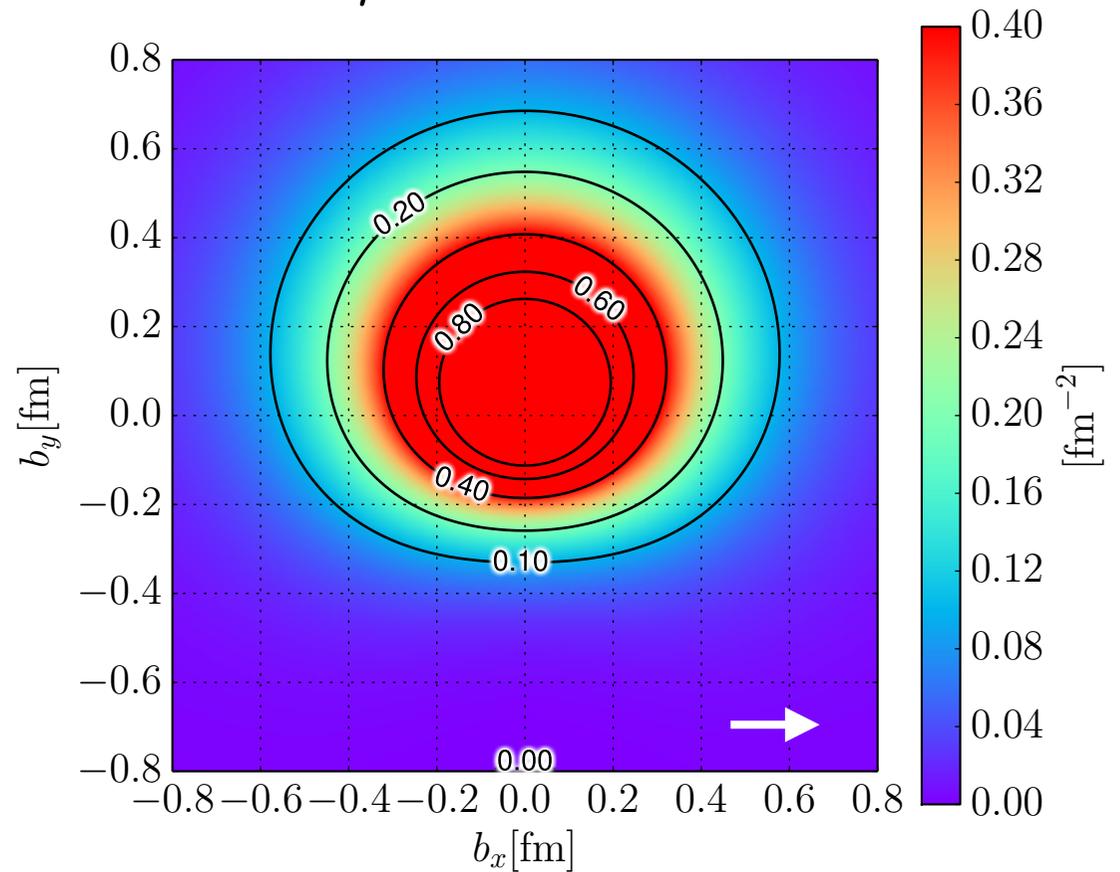


Transverse Quark Spin Density

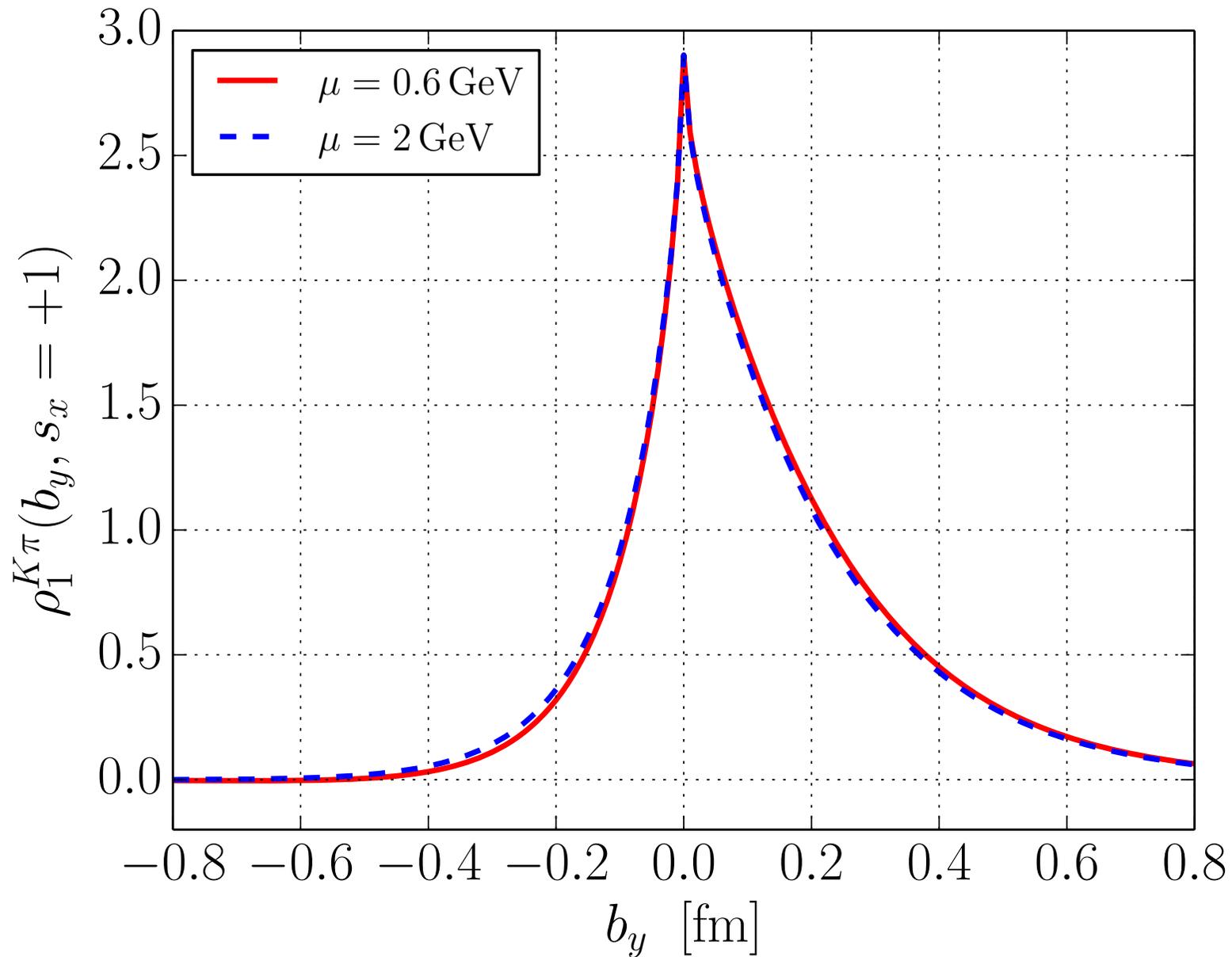
$\mu = 0.6 \text{ GeV}$



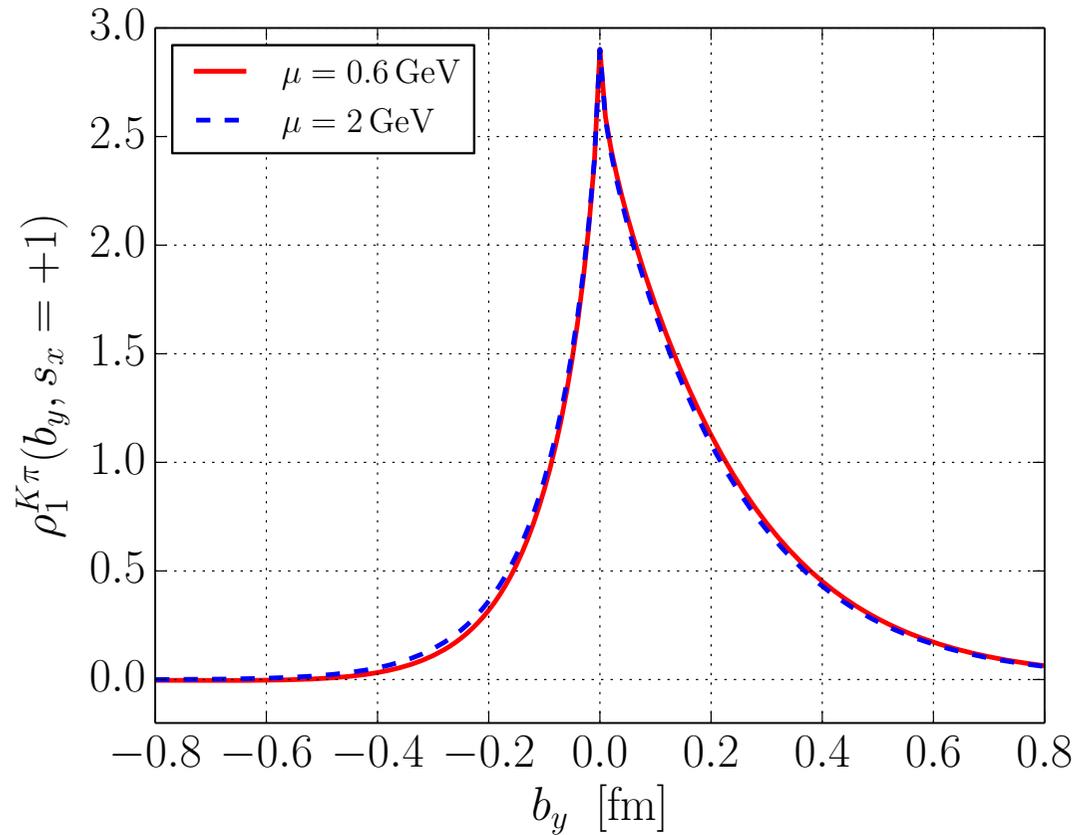
$\mu = 2.0 \text{ GeV}$



Transverse Quark Spin Density



Transverse Quark Spin Density



- Average Shift

$$\langle b_y \rangle^{K\pi} = \frac{\int d^2b b_y \rho_1^{K\pi}(b, s_\perp)}{\int d^2b \rho_1^{K\pi}(b, s_\perp)} = \frac{1}{2m_K} \frac{B_{T1,0}^{K\pi}(0)}{A_{1,0}^{K\pi}(0)}$$

$$= (0.17, 0.15) \text{ fm}$$

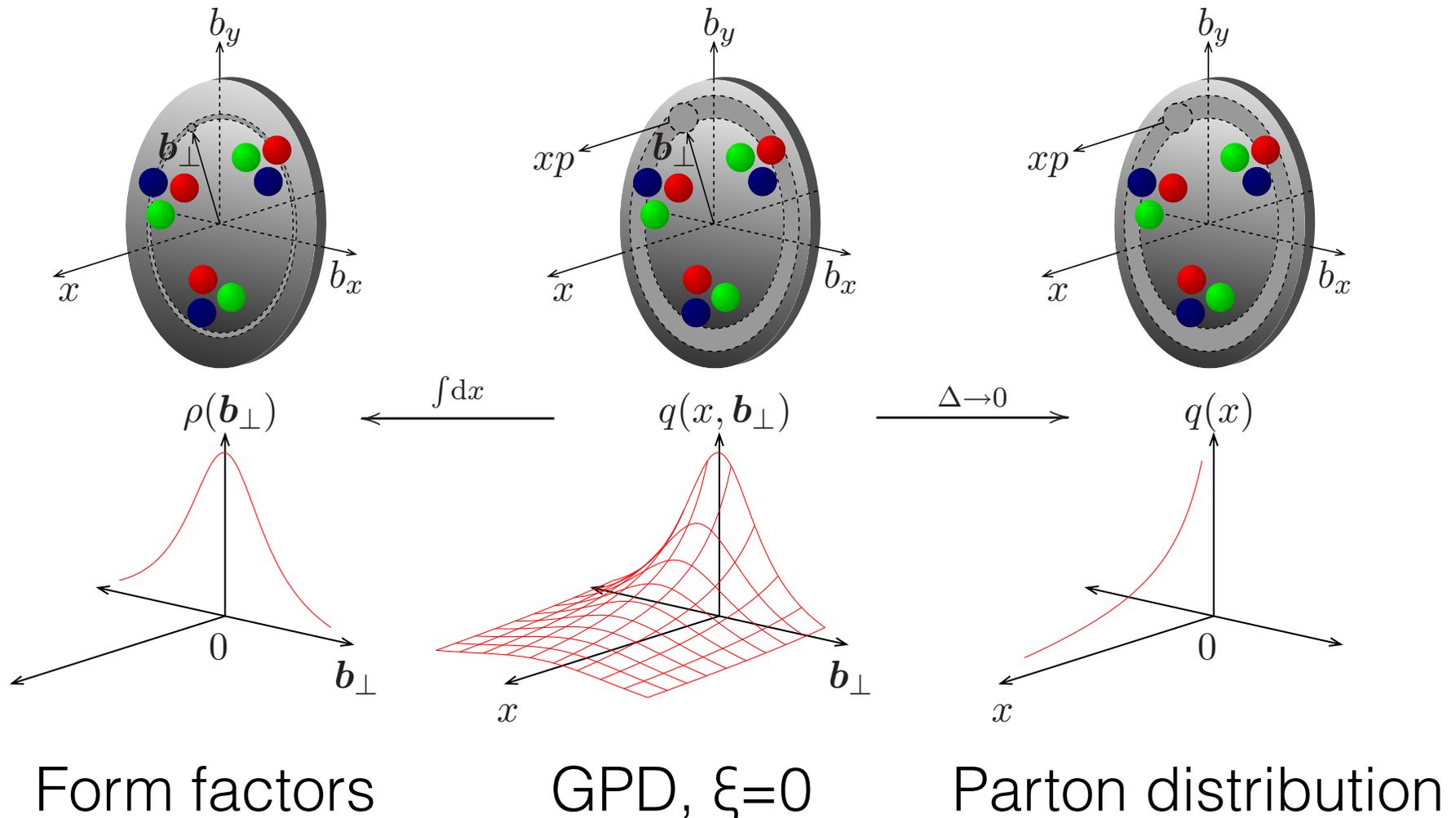
Summary & Outlook

- $K \rightarrow \pi$ generalized transition form factors ($n=0$)
- In good agreement with the lattice result
- Distorted quark spin structure when the quark spin is polarized
- Further studies on the wGPDs & GFFs

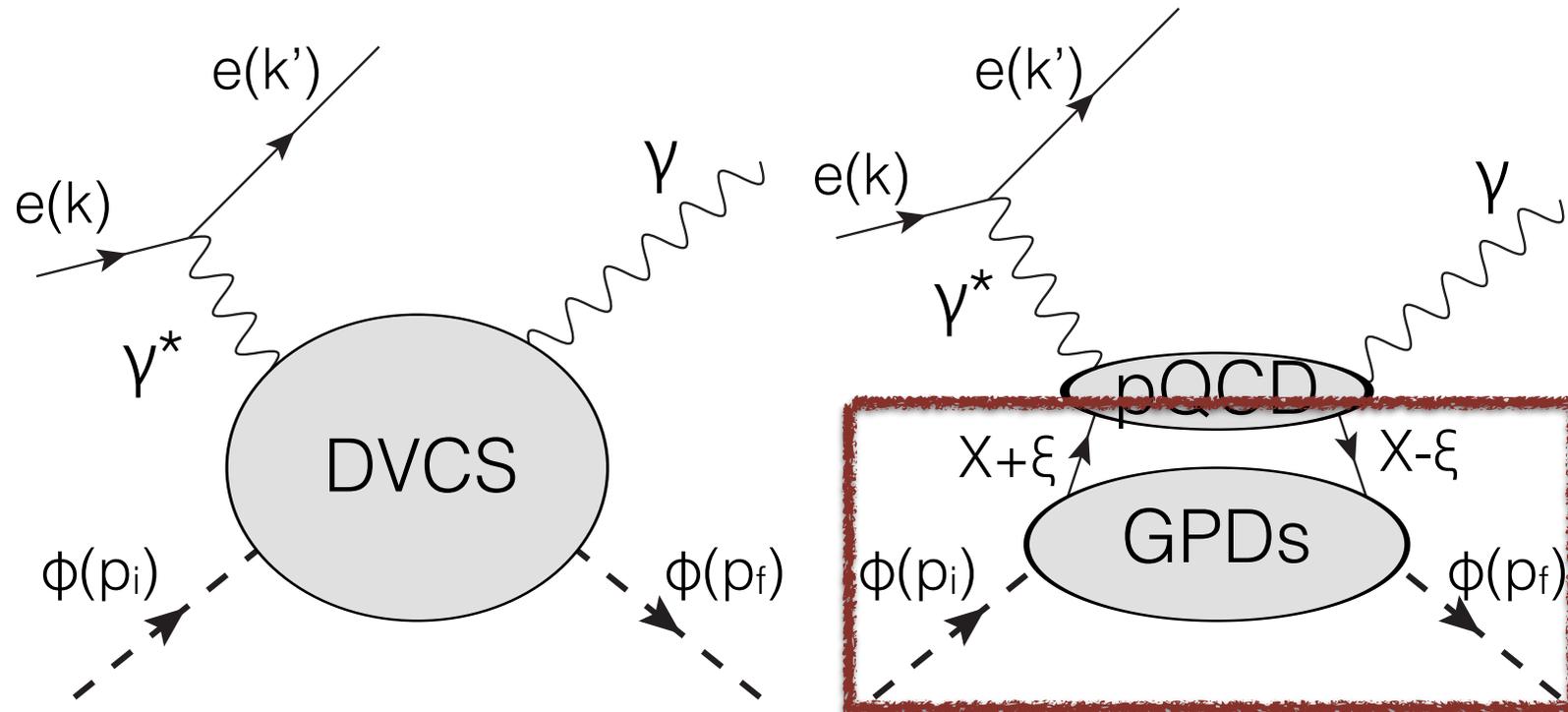
Thank you very much!

Generalized Parton Distributions

[D. Brömmel, Pion Structure From the Lattice, Regensburg Univ., Thesis]



DVCS & GPDs



Factorization : [Collins, J., & Freund, A. (1999). Phys. Rev. D, 59(7), 074009.]

- Light cone frame

$$P^\mu = \frac{1}{2} (p_i^\mu + p_f^\mu), \quad q^\mu = p_f^\mu - p_i^\mu, \quad t = q^2 \quad n_\pm = \frac{1}{\sqrt{2}} (1, 0, 0, \pm 1)$$

$$v^\pm = \frac{1}{\sqrt{2}} (v^0 \pm v^3), \quad v_\perp = (v^1, v^2) \quad \xi = \frac{p_i^+ - p_f^+}{p_i^+ + p_f^+}$$

Kaon l3 decay

The decay amplitude

$$T_{K \rightarrow l \mu \pi} = \frac{G_F}{\sqrt{2}} \sin \theta_c [W^\mu(p_l, p_\nu) F_\mu(p_\pi, p_K)]$$

$$G_F = 1.116 \times 10^{-5} \text{GeV}^{-2}$$

Weak leptonic element

$$W^\mu(p_l, p_\nu) = \bar{u}(p_\nu) \Gamma^\mu \nu(p_l)$$

Hadronic matrix element

$$F_\mu(p_\pi, p_K) = c \langle \pi(p_\pi) | \Gamma_\mu | K(p_K) \rangle$$

Hadronic Matrix Elements

❖ Vector transition

$$F_{\mu}^{K^0}(p_{\pi}, p_K) = \langle \pi^{-}(p_{\pi}) | \bar{s} \gamma_{\mu} u | K^0(p_K) \rangle = (p_K + p_{\pi})_{\mu} f_{l+}(t) + (p_K - p_{\pi})_{\mu} f_{l-}(t)$$

❖ Tensor transition

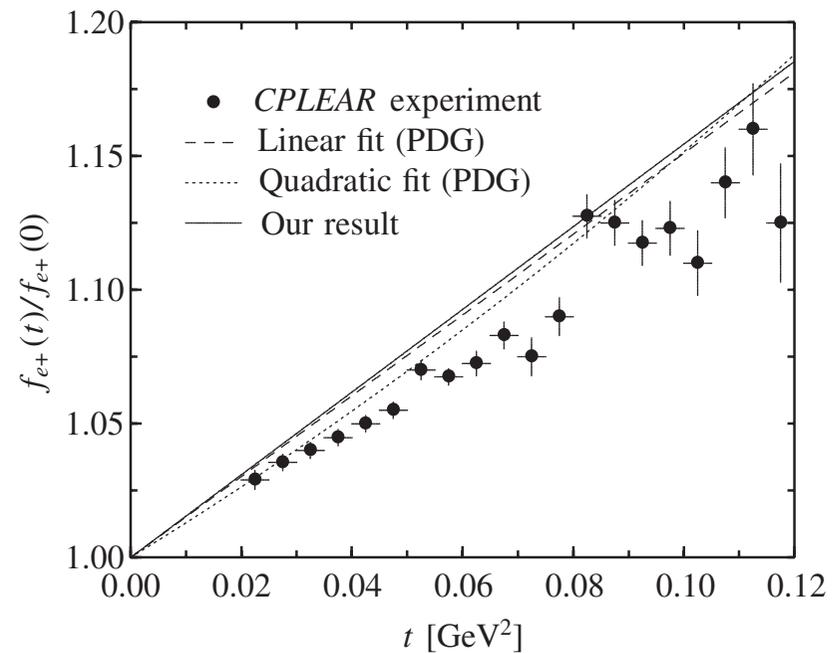
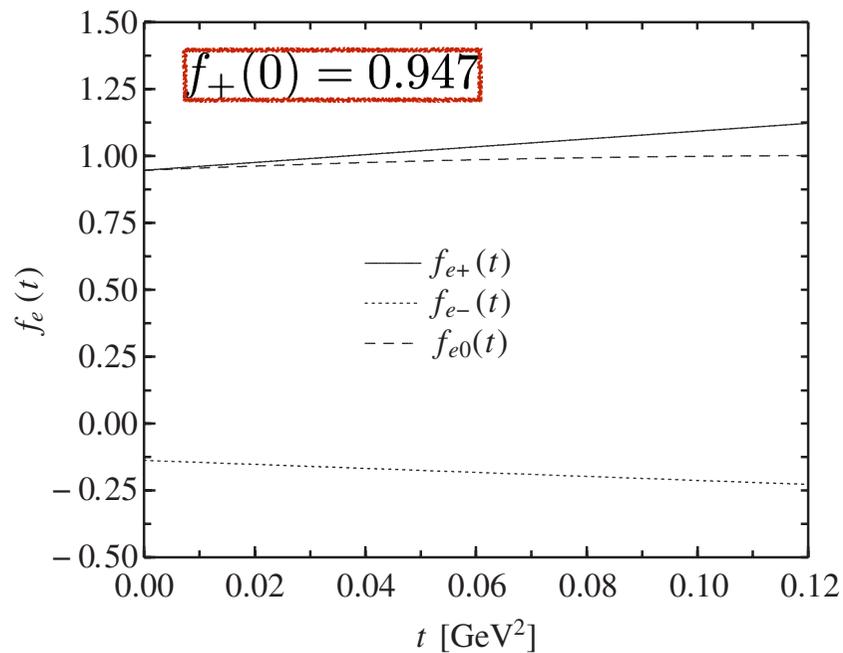
$$F_{\mu\nu}^{K^0}(p_{\pi}, p_K) = \langle \pi^{-}(p_{\pi}) | \bar{s} \sigma_{\mu\nu} u | K^0(p_K) \rangle = \frac{p_{K\mu} p_{\pi\nu} - p_{K\nu} p_{\pi\mu}}{m_K} B_T^{K\pi}(t)$$

❖ Scalar transition

$$F^{K^0}(p_{\pi}, p_K) = \langle \pi^{-}(p_{\pi}) | \bar{s} u | K^0(p_K) \rangle = -\frac{m_K^2 - m_{\pi}^2}{m_s - m_u} f_0(t)$$

Vector form factors

$$F_{\mu}^{K^0}(p_{\pi}, p_K) = \langle \pi^{-}(p_{\pi}) | \bar{s} \gamma_{\mu} u | K^0(p_K) \rangle = (p_K + p_{\pi})_{\mu} f_{l+}(t) + (p_K - p_{\pi})_{\mu} f_{l-}(t)$$



$$m_l^2 < t < (m_K - m_{\pi})^2 \approx 0.12 \text{ GeV}^2$$

[S.-i. Nam and H.-Ch. Kim, Phys. Rev. D 75, 094011 (2007).]

Generalized Form Factors

- Generalized form factors

Vector

$$\begin{aligned} & \langle \phi^a(p_f) | \psi^\dagger(0) \gamma_{\{\mu} i \overleftrightarrow{D}_{\mu_1} \dots \overleftrightarrow{D}_{\mu_{n-1}} \} \psi(0) | \phi^b(p_i) \rangle \\ & = 2P_{\{\mu} P_{\mu_1} \dots P_{\mu_{n-1}} \} A_{n0}(t) \\ & + \sum_{k=2 \text{ even}}^n q_{\{\mu} q_{\mu_1} \dots q_{\mu_{k-1}} P_{\mu_k} P_{\mu_{n-1}} \} 2^{-k} A_{nk}(t) \end{aligned}$$

Tensor

$$\begin{aligned} & \langle \phi(p_f) | \psi^\dagger(0) \sigma_{[\mu\nu} i \overleftrightarrow{D}_{\mu_1} \dots i \overleftrightarrow{D}_{\mu_{n-1}]} \psi(0) | \phi(p_i) \rangle \\ & = \frac{p_{[\mu} q_{\nu]} - q_{\mu} p_{\nu}}{m_\phi} \sum_{i=\text{even}}^{n-1} q_{\mu_1} \dots q_{\mu_i} P_{\mu_{i+1}} P_{\mu_{n-1}} B_{ni}(t) \end{aligned}$$

GFFs & Transverse Densities for the Kaon

K⁰ 스테이트에 (p_i)가 빠져있다 πππππ
π

- Generalised form factors for kaon transitions: n = 1

$$\langle \pi^-(p_f) | \bar{s}(0) \gamma_\mu u(0) | K^0 \rangle = 2P_\mu A_{10}^{K\pi}(t) + q_\mu C_{10}^{K\pi}(t)$$

$$\langle \pi^-(p_f) | \bar{s}(0) \sigma_{\mu\nu} u(0) | K^0 \rangle = \frac{p_{i\mu} p_{f\nu} - p_{i\nu} p_{f\mu}}{m_K} B_{10}^{K\pi}(t)$$

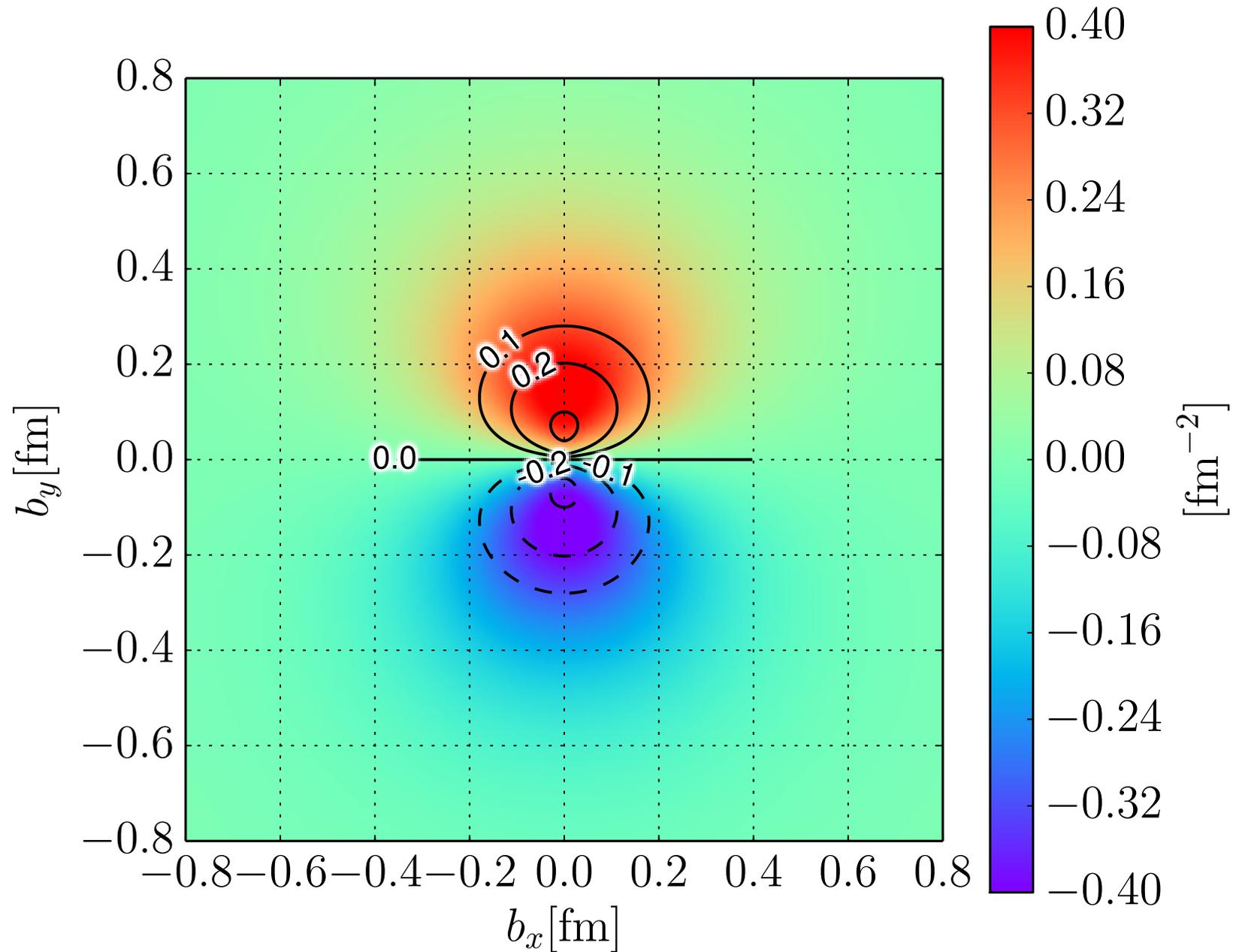
- Transverse densities:

2D Fourier transformation into the impact parameter b_\perp at $\xi=0$

$$\begin{aligned} F(b_\perp^2) &= \int dq_\perp^2 e^{-ib_\perp \cdot q_\perp} \int dX H(X, q_\perp^2, \xi = 0) \\ &= \int dq_\perp^2 e^{-ib_\perp \cdot q_\perp} F(q_\perp^2) \end{aligned}$$

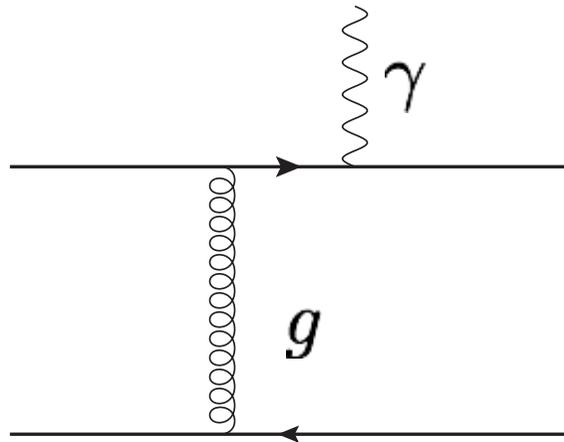
Probability distribution of the partons inside the hadrons
in the transverse impact parameter plane

Transverse Spin Density



Singular Behaviour of Transverse Densities

Perturbative QCD prediction: Form factor



$$\sim \frac{1}{Q^2}$$

Fourier transformation

$$A_{n0}(b_{\perp}) = \frac{1}{(2\pi)^2} \int d^2q_{\perp} e^{-ib_{\perp} \cdot q_{\perp}} A_{n0}(q_{\perp}^2)$$

Logarithmically divergent
at $b=0$ for $p=1$.
Singular when $p < 1.5$

[G.A. Miller, Phys.Rev.C 79, 055204(2009)]