

# Nucleon and photon structure functions at small $x$ in a holographic QCD model



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This talk is based on:

AW, Katsuhiko Suzuki, PRD89, 115015 (2014)

AW, Hsiang-nan Li, arXiv:1502.03894 [hep-ph]

AW, Hsiang-nan Li, (in preparation)

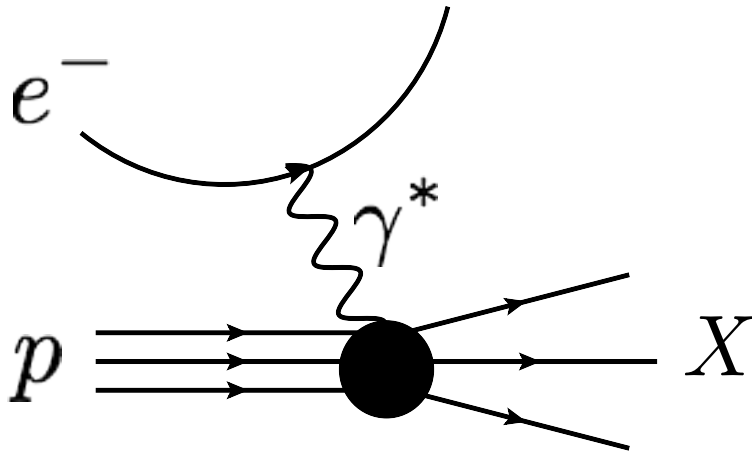
Hadrons and Hadron Interactions in QCD  
-- Effective theories and Lattice -- (HHIQCD2015)

February 19, 2015 @ Yukawa Institute for Theoretical Physics

Nucleon structure functions at small  $x$

# DIS structure functions

$$\frac{d^2\sigma}{dx dQ^2} = \frac{2\pi\alpha^2}{xQ^4} \left[ \left\{ 1 + (1-y)^2 \right\} F_2(x, Q^2) - y^2 F_L(x, Q^2) \right]$$



$$F_2(x, Q^2) = \frac{Q^2}{4\pi^2\alpha} \sigma_{tot}(x, Q^2)$$

$$F_L(x, Q^2) = \frac{Q^2}{4\pi^2\alpha} \sigma_L(x, Q^2)$$

- Structure functions are physical quantities which have information on the internal structure of hadrons.
- They depend on two kinematic variables, Bjorken- $x$  and photon 4-momentum squared  $Q^2$ .

# Longitudinal structure function $F_L$

- In the quark-parton model,  $F_2$  can be written as:

$$F_2 = x \sum_q e_q^2 q_i(x) \quad F_L = 0$$

- $F_L$  is expressed, for example by Altarelli-Martinelli equation, as:

Altarelli-Martinelli (1978)

$$F_L(x, Q^2) = \frac{\alpha_s(Q^2)}{2\pi} x^2 \int_x^1 \frac{dy}{y^3} \left[ \frac{8}{3} F_2(y, Q^2) + 4 \sum_q e_q^2 \left( 1 - \frac{x}{y} \right) y g(y, Q^2) \right]$$

At small  $x$ , the second term becomes dominant and  $F_L$  is approximately expressed by

Cooper-Sarkar et al. (1988)

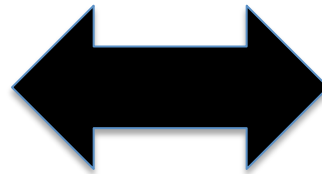
$$F_L \approx 0.3 \frac{4\alpha_s}{3\pi} x g(2.5x, Q^2)$$

Hence, studying  $F_L$  is related to the investigations on the gluon distribution and the higher order contribution in QCD.

# Holographic QCD

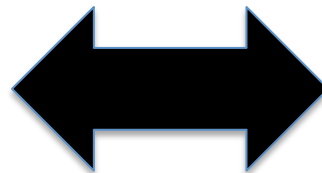
- Holographic QCD, which is constructed based on the AdS/CFT correspondence, has a potential to be a powerful tool for analyses on hadron physics.

type IIB  
supergravity theory  
on  $S^5 \times \text{AdS}_5$



strong coupling 4D N=4  
supersymmetric Yang-  
Mills (SYM) theory

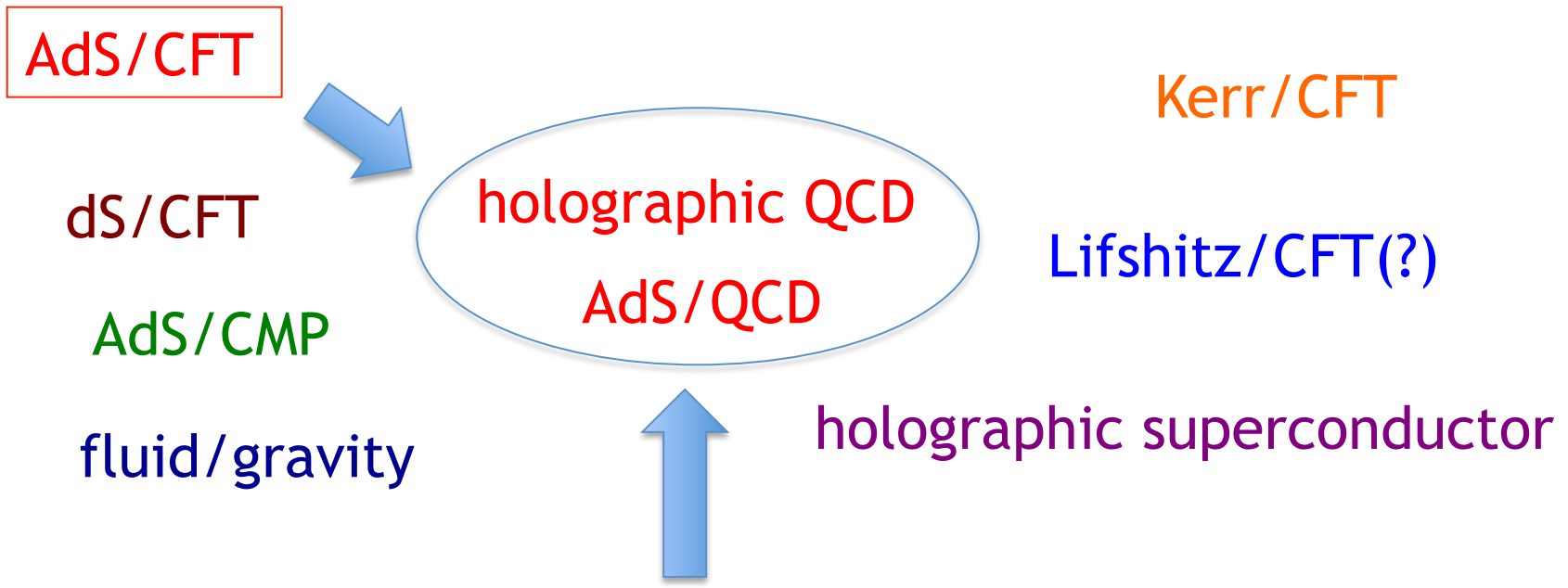
supergravity theory  
(classical theory)  
on  $\text{AdS}_5$



usual 4D QCD  
at strong coupling

# Gauge/string correspondence (holography)

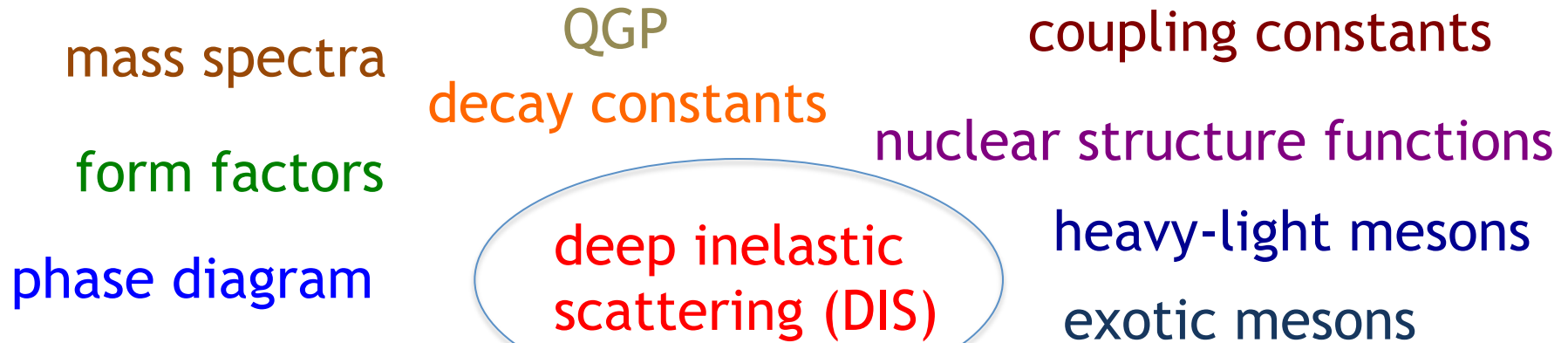
*many correspondences...*



*We are interested in hadron physics and QCD!*

# Holographic approaches to hadron physics

*many studies...*



- Polchinski-Strassler (2003)
- Brower-Polchinski-Strassler-Tan (2007)
- Cornalba-Costa-Penedones (2007, 2010)
- Hatta-Iancu-Mueller (2008)
- Bayona - Boschi-Filho - Braga (2008)
- Brower-Strassler-Tan (2009)
- Brower-Djuric-Sarcevic-Tan (2010)
- Watanabe-Suzuki (2012, 2014)
- Watanabe-Li (2015) and so on...

# Preceding studies

- (2003-2009) Important studies on DIS from holography have been done almost completely...
  - e.g., in 2003, Polchinski and Strassler have demonstrated string calculations for structure functions and shown that the Callan-Gross relation holds.
- (2010) Brower-Djuric-Sarcevic-Tan have studied the nucleon structure function at small  $x$  by considering the single Pomeron exchange, and well reproduced the HERA data for  $F_2^p$ .



# Holographic description of structure functions

- Polchinski-Strassler (2003)
- Brower-Polchinski-Strassler-Tan (2007)
- **Brower-Djuric-Sarcevic-Tan (2010)**

derived Pomeron exchange kernel

studied nucleon structure functions

$$\mathcal{A}(s,t) = 2is \int d^2b e^{iq \cdot b} \int dz dz' P_{13}(z) P_{24}(z') \{1 - e^{i\chi(s,b,z,z')}\}$$

$$F_2(x, Q^2) = \frac{Q^2}{2\pi^2} \int dz dz' P_{13}(z, Q^2) P_{24}(z') \text{Im}[\chi(s, z, z')]$$

$z$  and  $z'$  : 5th coordinate

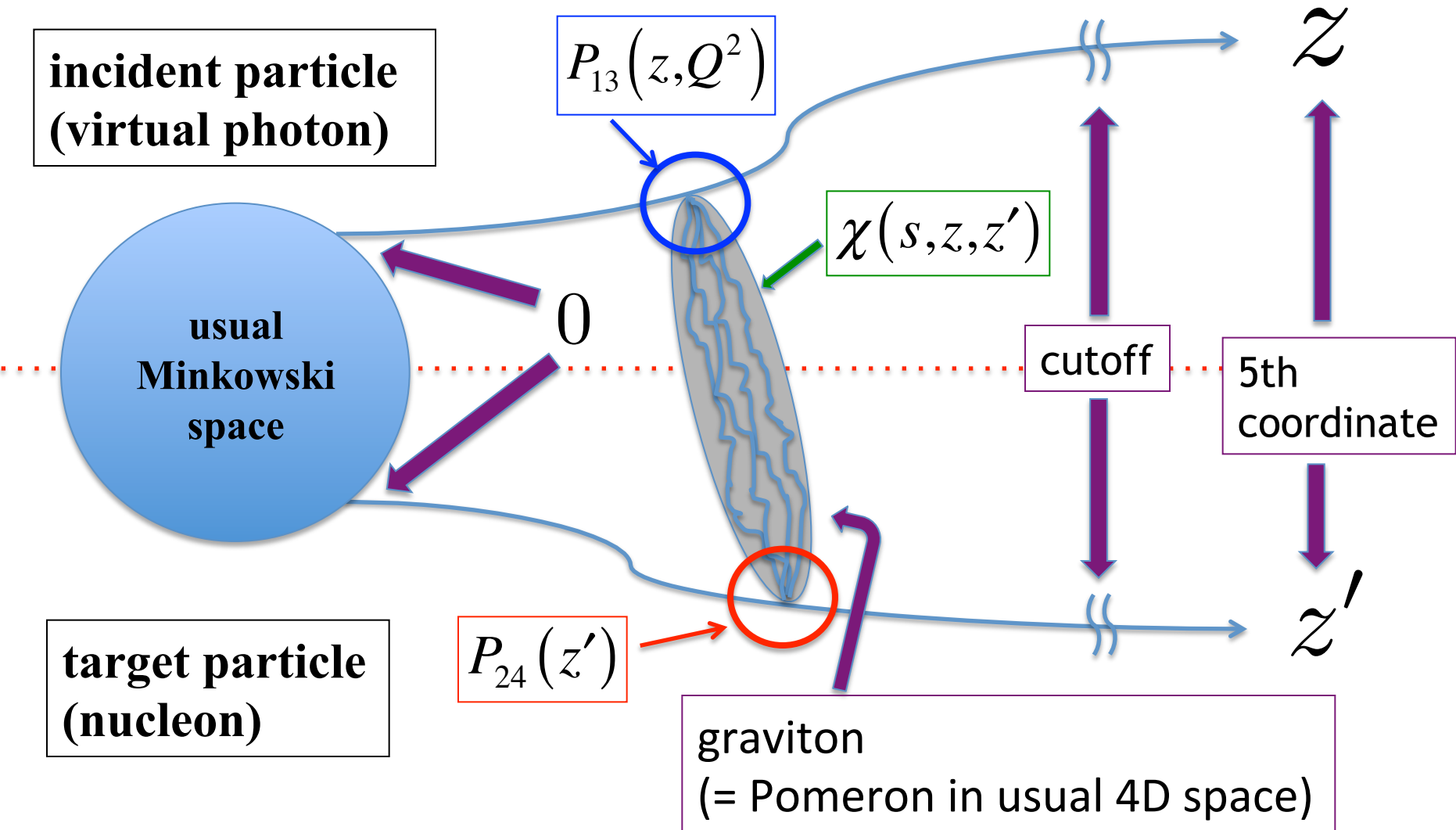
$\chi$  : Pomeron exchange kernel in the AdS space

$P_{13}(z, Q^2)$  : incident particle  
(virtual photon, 4-momentum  $Q$ )

$P_{24}(z')$  : target particle

overlap functions  
(density distributions in the AdS space)

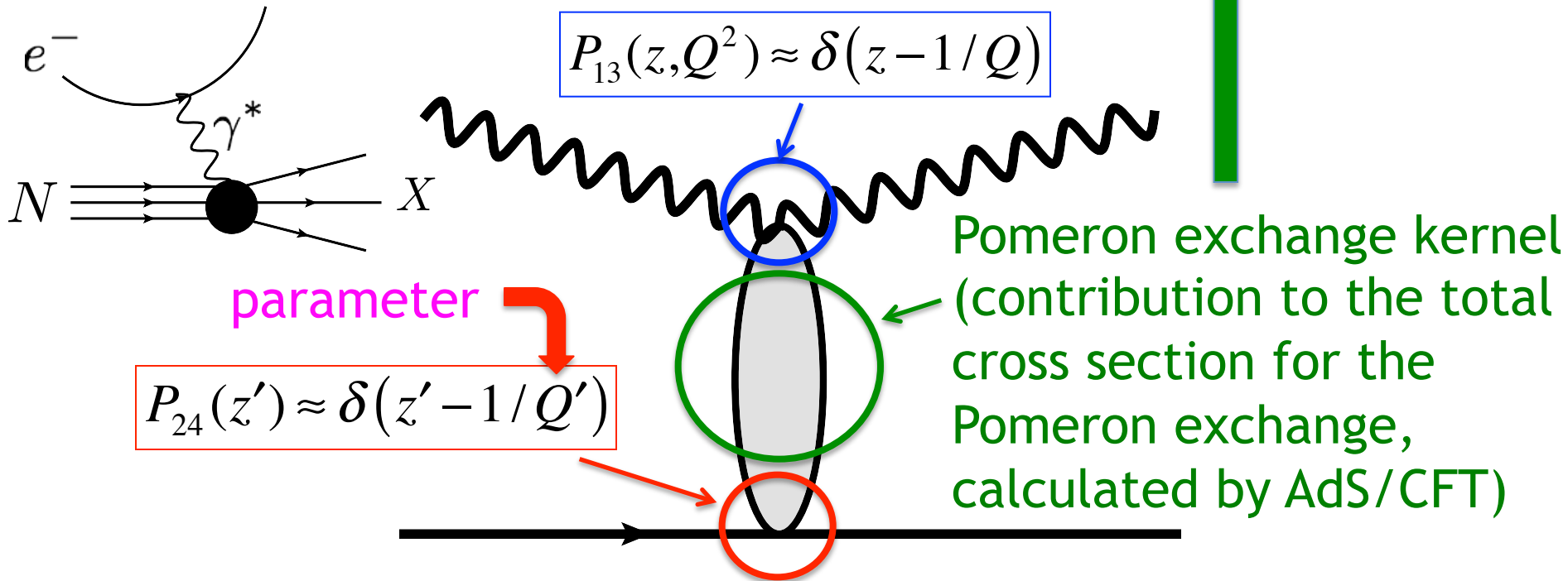
# 5D background spacetime ( $AdS_5$ )



# A preceding study

Brower-Djuric-Sarcevic-Tan (2010)

$$F_2(x, Q^2) = \frac{Q^2}{2\pi^2} \int dz dz' P_{13}(z, Q^2) P_{24}(z') \text{Im}[\chi(s, z, z')]$$



- Incident and target particles are simply replaced with delta functions.

# Pomeron exchange kernel

**Brower-Polchinski-Strassler-Tan (2007)**

**Brower-Strassler-Tan (2009)**

$$F_i(x, Q^2) = \frac{g_0^2 \rho^{3/2} Q^2}{32\pi^{5/2}} \int dz dz' P_{13}^{(i)}(z, Q^2) P_{24}(z') (zz') \text{Im}[\chi(s, z, z')]$$

$i = 2 \text{ or } L$

$$\text{Im}[\chi_c(s, z, z')] \equiv e^{(1-\rho)\tau} e^{-\frac{\log^2 z/z'}{\rho\tau}} / \tau^{1/2}$$

$$\tau = \log(\rho z z' s / 2)$$

$$\text{Im}[\chi_{\text{mod}}(s, z, z')] \equiv \text{Im}[\chi_c(s, z, z')] + \mathcal{F}(z, z', \tau) \text{Im}[\chi_c(s, z, z_0^2 / z')]$$

$$\mathcal{F}(z, z', \tau) = 1 - 2\sqrt{\rho\pi\tau} e^{\eta^2} \text{erfc}(\eta)$$

$$\eta = \left( -\log \frac{zz'}{z_0^2} + \rho\tau \right) / \sqrt{\rho\tau}$$

↓  
confinement  
effect

3 adjustable parameters :

$$\rho, g_0^2, z_0$$

# Our model setup

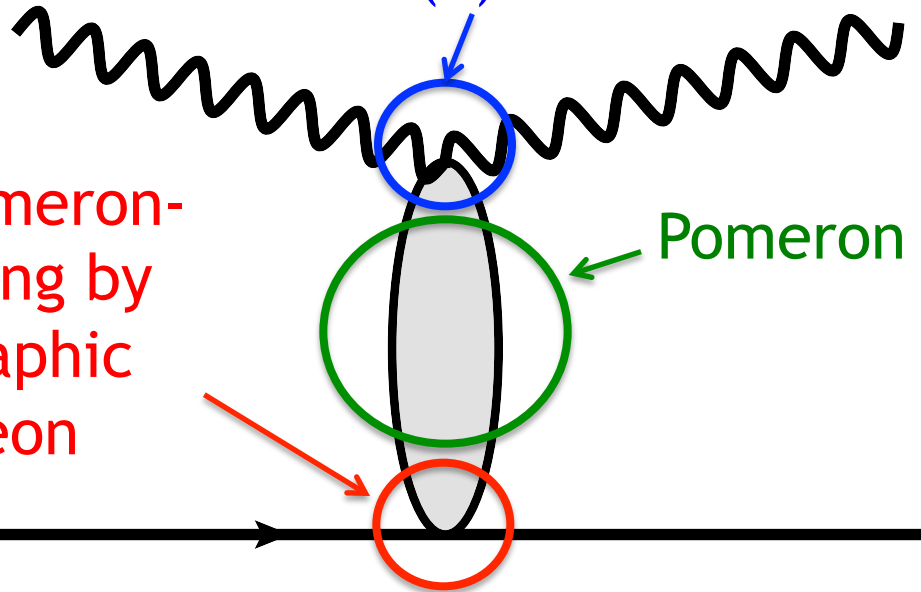
using wave function of  
the 5D U(1) vector field

virtual photon

calculating Pomeron-  
nucleon coupling by  
using a holographic  
model of nucleon

Pomeron exchange kernel

nucleon

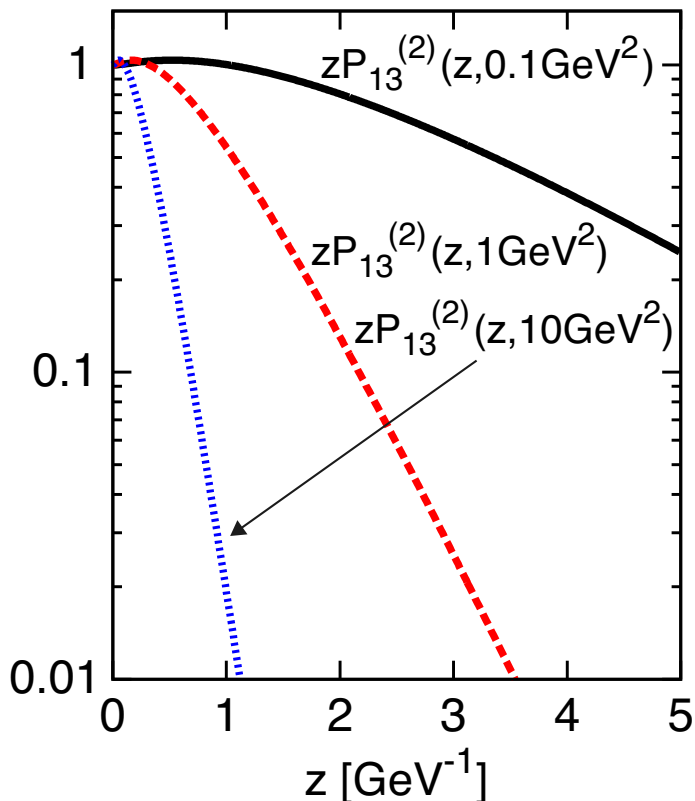


- This is a more consistent description of structure functions.
- In this model, we can consider structure functions of various hadrons and longitudinal structure functions, which can not be considered in the preceding model setup.

# Density distribution of the probe photon

Polchinski-Strassler (2003)

- As a density distribution of the virtual photon, we use wave function of the 5D U(1) vector field



$$P_{13}^{(2)}(z, Q^2) = Q^2 z \left( K_0^2(Qz) + K_1^2(Qz) \right)$$

(to calculate  $F_2$ )

$$P_{13}^{(L)}(z, Q^2) = Q^2 z K_0^2(Qz)$$

(to calculate  $F_L$ )

$P_{13}^{(2)}$  are localized at the origin  
with  $Q^2$  increasing  
(the behavior of  $P_{13}^{(L)}$  is similar)

# Density distribution for the target particle (1)

Abidin-Carlson (2009)

The matrix element of the energy momentum tensor in respect to spin 1/2 particle:

$$\langle p_2, s_2 | T^{\mu\nu}(0) | p_1, s_1 \rangle = u(p_2, s_2) \left( A(t) \gamma^{(\mu} p^{\nu)} + B(t) \frac{i p^{(\mu} \sigma^{\nu)\alpha} q_\alpha}{2m} + C(t) \frac{q^\mu q^\nu - q^2 \eta^{\mu\nu}}{m} \right) u(p_1, s_1)$$

To calculate A(t) etc. with the holographic model of nucleons,

$$S_F = \int d^5x \sqrt{g} e^{-\kappa^2 z^2} \left( \frac{i}{2} \bar{\Psi} e_A^N \Gamma^A D_N \Psi - \frac{i}{2} (D_N \Psi)^\dagger \Gamma^0 e_A^N \Gamma^A \Psi - (M + \kappa^2 z^2) \bar{\Psi} \Psi \right)$$

we introduce the metric perturbation,  $\eta_{\mu\nu} \rightarrow \eta_{\mu\nu} + h_{\mu\nu}$ , in the 5D classical action, and pick up the  $h\psi\psi$  terms.

By comparing the Lorentz structure of them, we can obtain the form factors. (in this case, only A(t) remains)

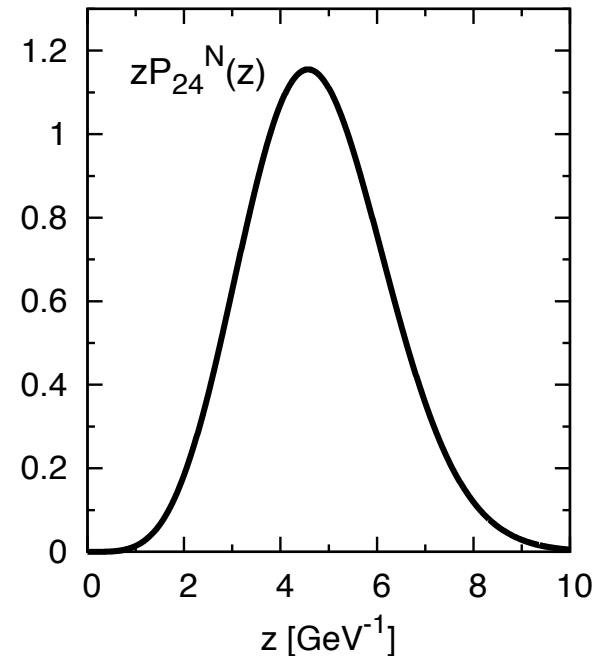
# Density distribution for the target particle (2)

Finally, one can obtain the density distribution for the target nucleon:

$$P_{24}(z') = \frac{e^{-\kappa^2 z'^2}}{2z'^3} (\psi_L^2(z') + \psi_R^2(z')) \quad \int dz' P_{24}(z') = 1$$

where  $\psi_{L,R}$  are 5D wave functions describing a nucleon as a 5D Dirac fermion with chiral symmetry breaking.

- The peak position of  $P_{24}$  is in the large  $z$  region, which is obviously different from  $P_{13}$ .





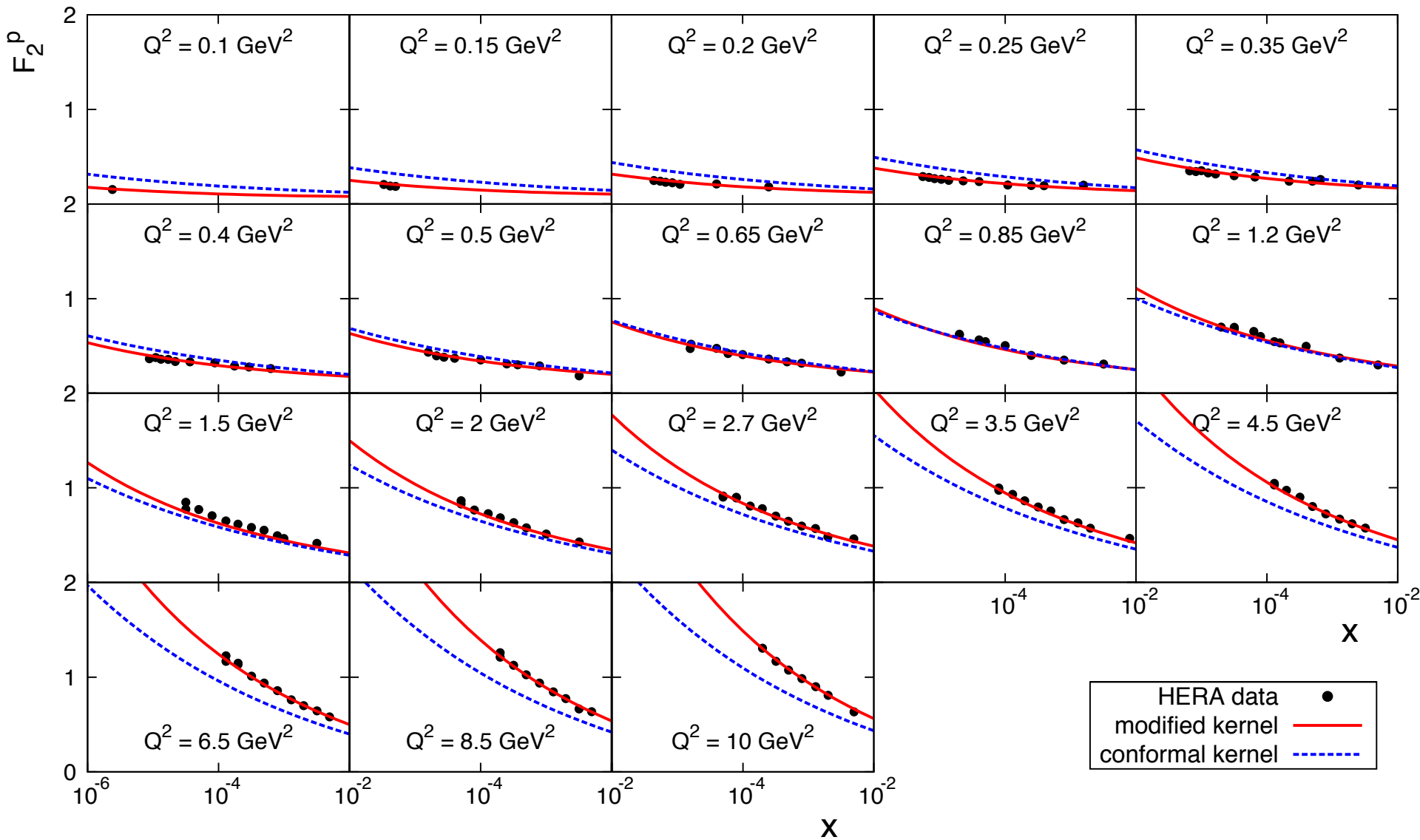
# Fixing parameters

- There are 3 adjustable parameters ( $\rho$ ,  $g_0$ , and  $z_0$ )
- They are fixed with the experimental data for  $F_2^p$  measured at HERA

ZEUS collaboration (1999)

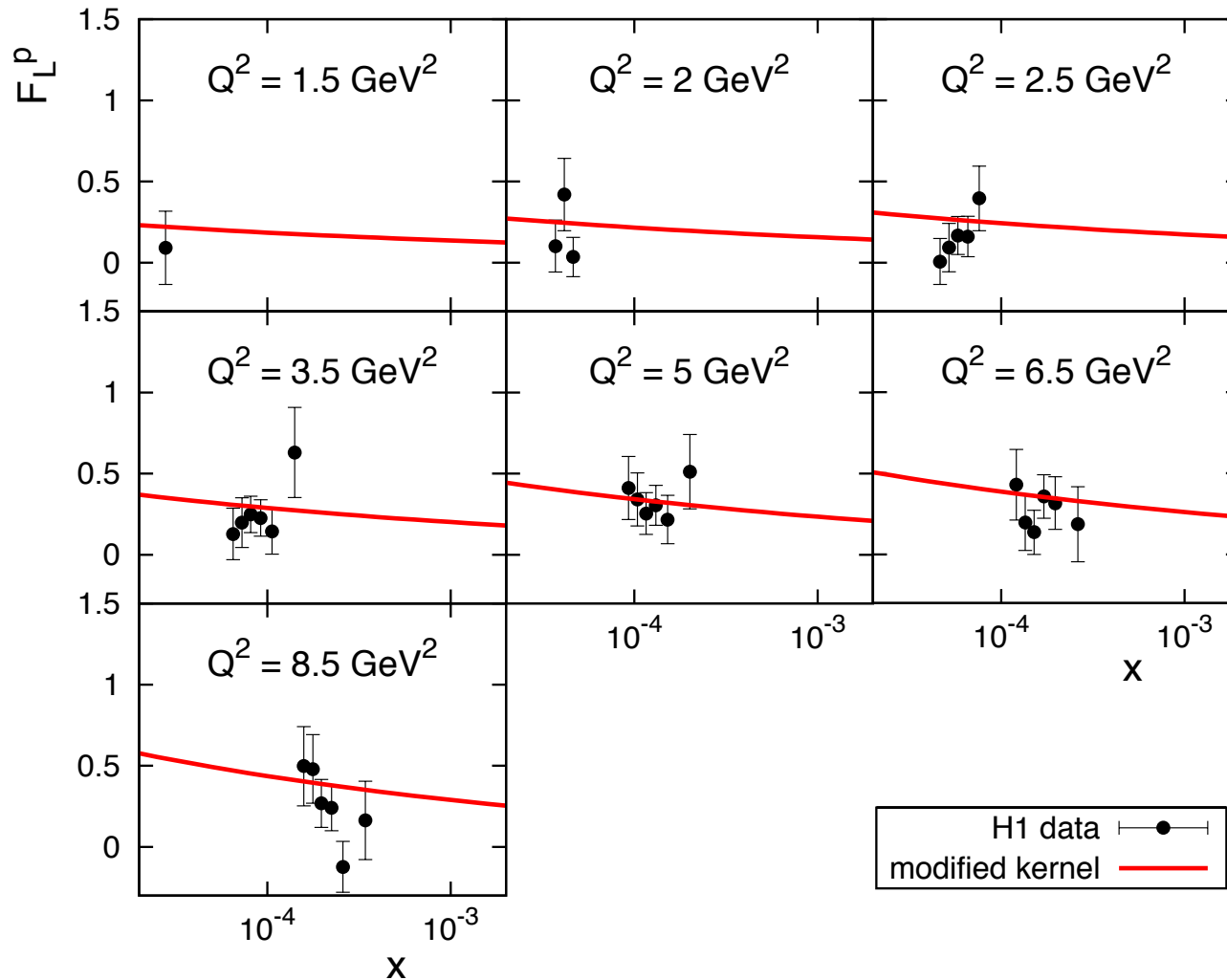
H1 and ZEUS collaborations (2010)

# Proton structure function



# Proton longitudinal structure function

- Replacing the density distribution of the probe photon with the longitudinal component



Photon structure functions at small  $x$

# Photon's "internal structure"?

- In high energy scattering (here it is the electron-photon DIS), it can fluctuate into  $q\bar{q}$  pairs or vector mesons.  
-> "dressed photon"
- One can investigate the parton structure inside (or inside the cloud of) the photon, which has both pointlike and hadronic components.

# Motivation

- A solid understanding of properties of the elementary particles is basically important.
- Hadronic contribution to cross sections for electron-photon DIS becomes dominant at small  $x$ . (→ Pomeron exchange may work)
- Two of the three adjustable parameters of the model have already been fixed via the study on the nucleon. We shall test the predictive power.
- Electron-photon DIS is a cleaner process.
- Preceding studies at small  $x$  are quite limited.
- The predictions can be tested at ILC in the future.

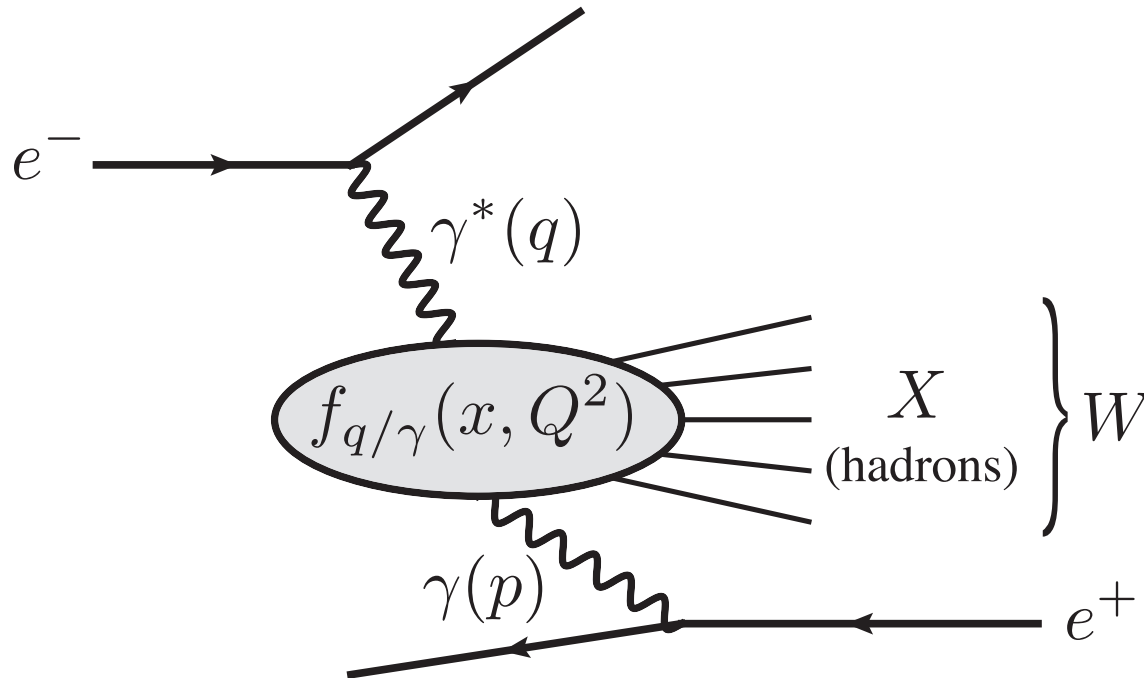
# History

- 1971 :  $e\gamma$  scattering cross section in the VMD model (Brodsky-Kinoshita-Terazawa, Walsh)
- 1973 : Calculating the structure function in the QM (Walsh-Zerwas)
- 1977 : Calculating the LO QCD corrections (Witten)
- 1981 : First experimental test (PLUTO Collaboration)
- 1979, 1980 : QCD calculations at NLO (Bardeen-Buras, Duke-Owens) -> negative values of  $F_2^\gamma$  at small  $x$
- 1992 : An algebraic error in *Bardeen-Buras (1979)* was corrected (Fontanaz-Pilon, Gluck-Reya-Vogt)
- 2002 : QCD calculations at NNLO (Moch-Vermaseren-Vogt)

# Deeply inelastic electron-photon scattering

$$\frac{d^2 \sigma_{e\gamma \rightarrow eX}}{dx dQ^2} = \frac{2\pi\alpha^2}{xQ^4} \left[ \left\{ 1 + (1-y)^2 \right\} F_2^\gamma(x, Q^2) - y^2 F_L^\gamma(x, Q^2) \right]$$

$$x = \frac{Q^2}{Q^2 + W^2 + P^2} \quad \text{when} \quad W^2 \gg Q^2 \gg P^2 \quad \rightarrow \quad x \approx \frac{Q^2}{W^2}$$





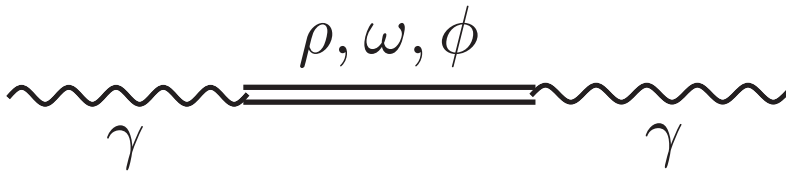
# Two components in electron-photon DIS

$$\frac{d^2\sigma}{dQ^2 dx} = \frac{2\pi\alpha^2}{\pi Q^4} \left[ \left\{ 1 + (1-y)^2 \right\} F_2^\gamma - y^2 F_L^\gamma \right]$$

$$x = \frac{Q^2}{Q^2 + W^2}, \quad y = \frac{Q^2}{xs}$$

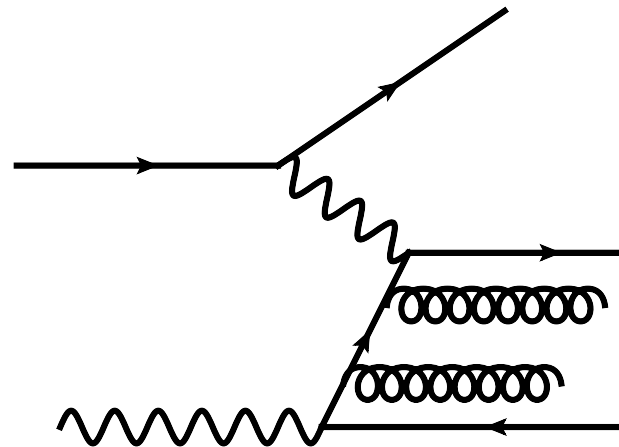
in quark model  $F_2^\gamma(x, Q^2) = 2x \sum_q e_q^2 q^\gamma(x, Q^2), \quad F_L^\gamma = 0$

“hadronic” (small x)



- A photon behaves like  $\rho$  meson (and other vector mesons)
- Utilizing the vector meson dominance model is the only way to calculate cross sections

“pointlike” (large x)



- One can predict cross sections by pQCD

# Virtual photon – quasi-real photon scattering

$$F_i^\gamma(x, Q^2) = \frac{g_0^2 \rho^{3/2} Q^2}{32\pi^{5/2}} \int dz dz' P_{13}^{(i)}(z, Q^2) P_{24}(z', P^2 \ll 1 \text{ GeV}^2) \text{Im}[\chi(s, z, z')]$$

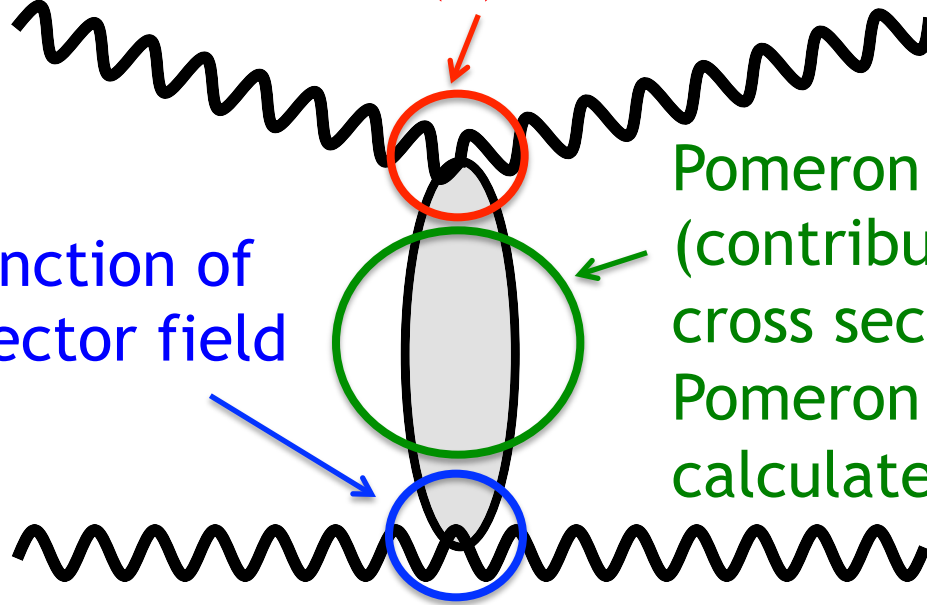
using wave function of  
the 5D U(1) vector field

virtual photon

using wave function of  
the 5D U(1) vector field

quasi-real  
photon

Pomeron exchange kernel  
(contribution to the total  
cross section for the  
Pomeron exchange,  
calculated by AdS/CFT)



We reuse the parameters  $\rho$  and  $z_0$  which were fixed in the calculations for the nucleon  $F_2$  structure function.

# Virtual photon – $\rho$ meson scattering

$$F_i^\gamma(x, Q^2) = \frac{g_0^2 \rho^{3/2} Q^2}{32\pi^{5/2}} \int dz dz' P_{13}^{(i)}(z, Q^2) P_{24}(z') \text{Im}[\chi(s, z, z')]$$

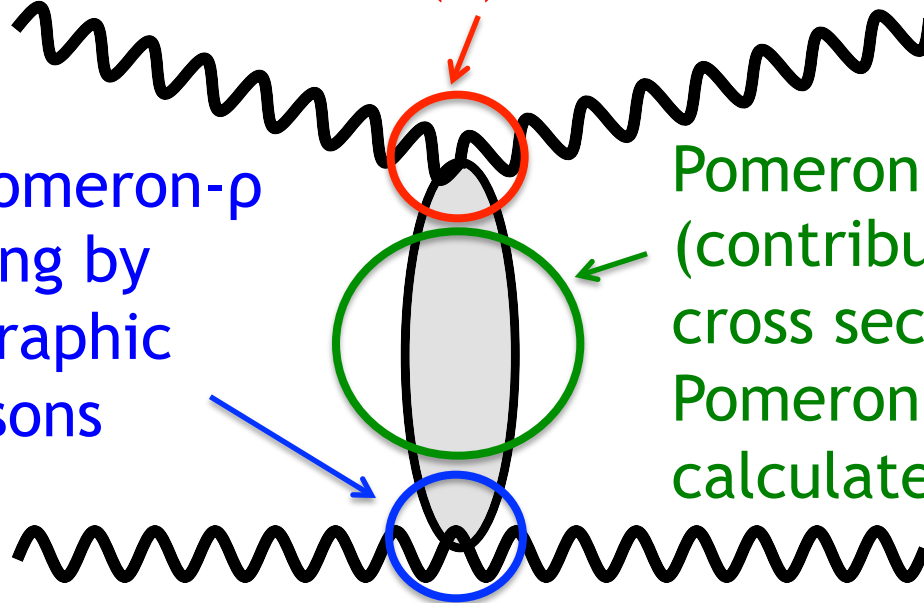
using wave function of  
the 5D U(1) vector field

virtual photon

calculating Pomeron- $\rho$   
meson coupling by  
using a holographic  
model of mesons

real photon  
( $\rho$  meson)

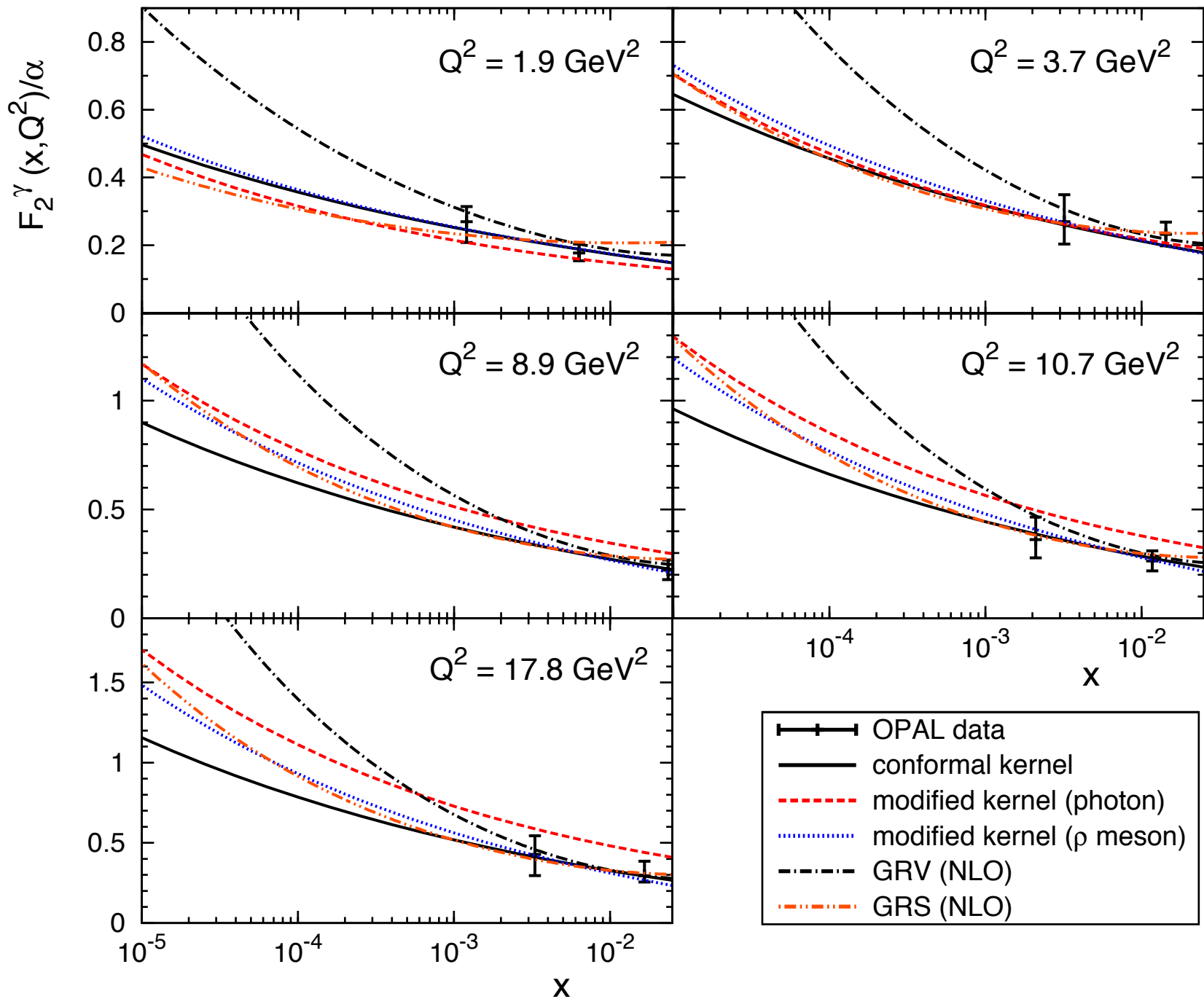
Pomeron exchange kernel  
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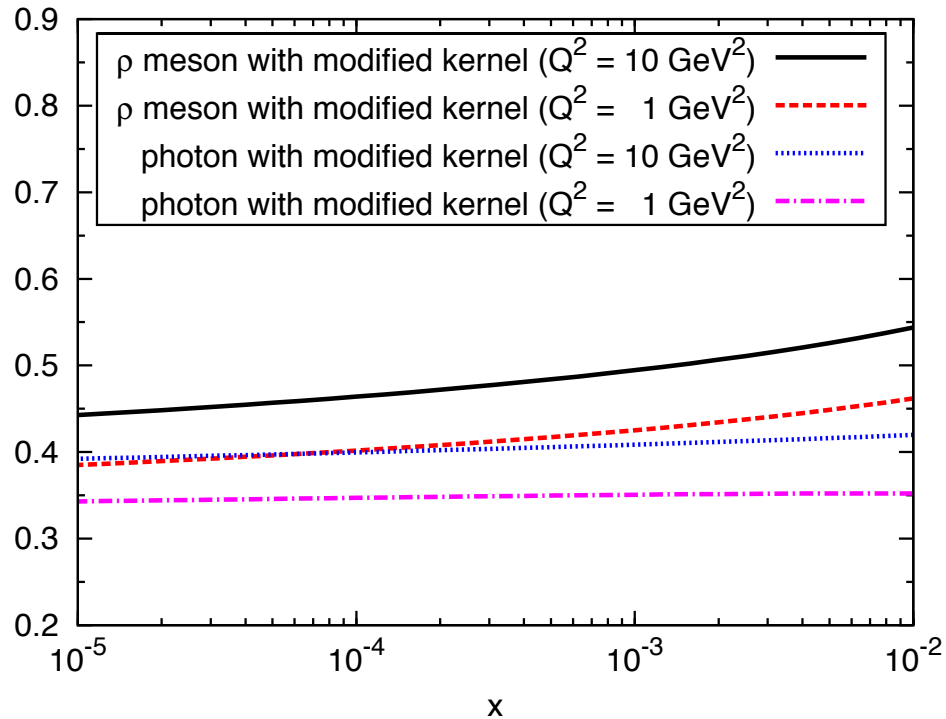
# Directions

- Since the target is a non-normalizable mode this time, we need to newly fix the overall factor by experimental data.
- We adopt the data measured by the OPAL collaboration at LEP.
- Since the available experimental data for the small  $x$  region are quite limited, we also compare our calculations (predictions) with those calculated from the well-known PDF sets, GRV (Gluck-Reya-Vogt, 1992) and GRS (Gluck-Reya-Schienbein, 1999).

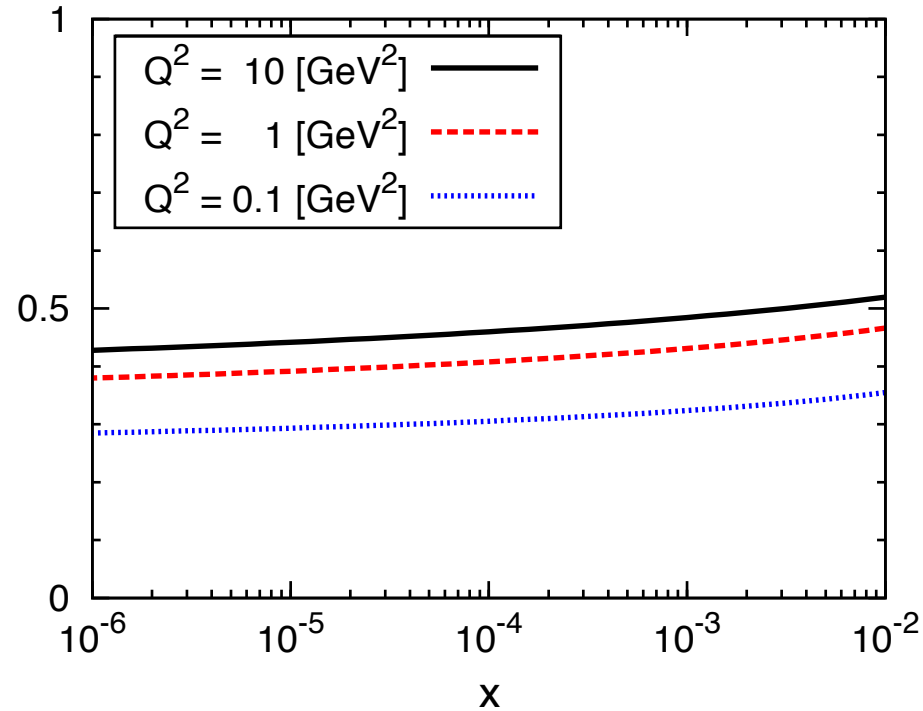


# Longitudinal-to-transverse ratio $R = F_L/F_T$

Photon



Nucleon



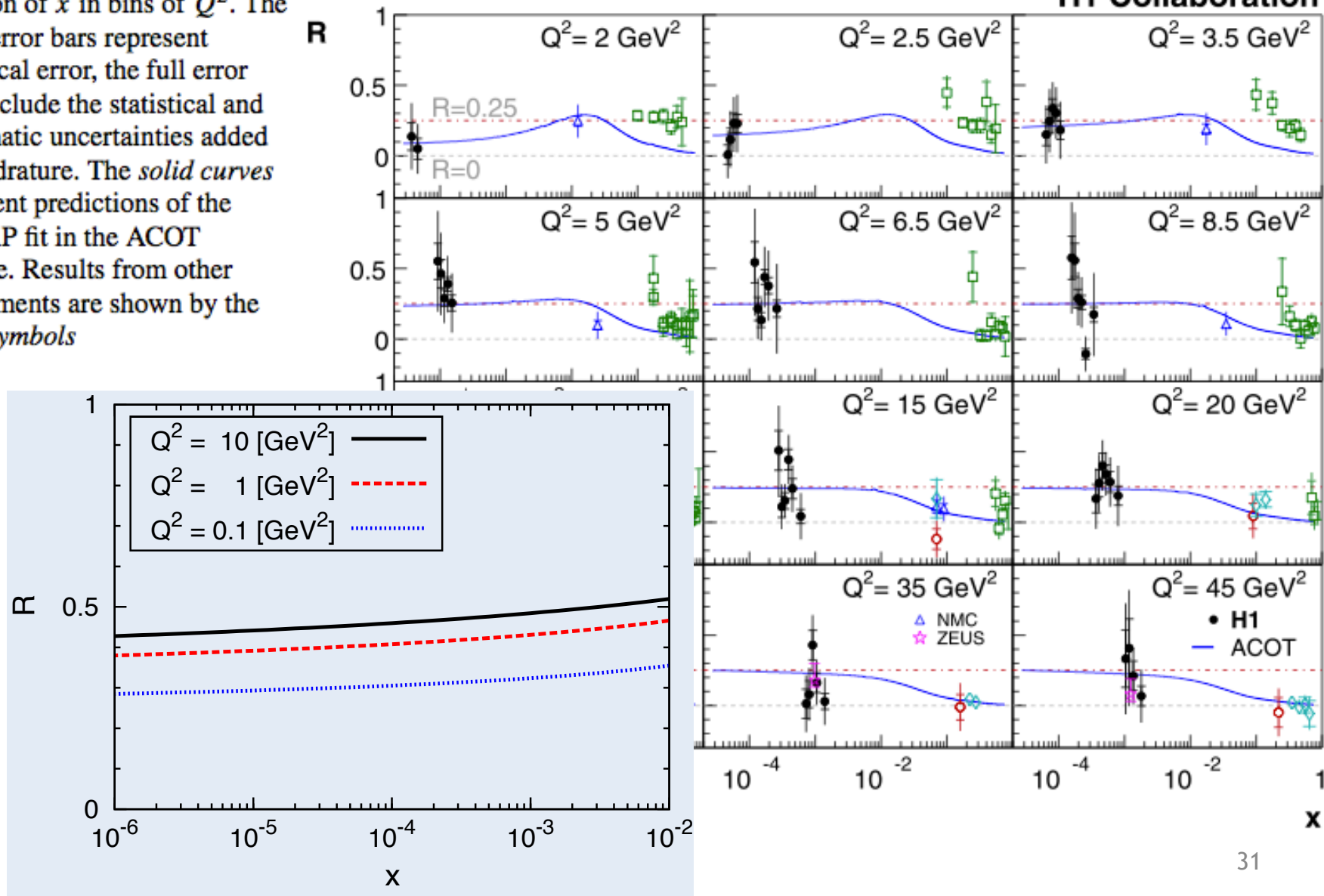
The longitudinal structure function is not negligible in this kinematical region.

# $x$ dependence of $R$

H1 Collaboration (2011)

**Fig. 14** The ratio  $R$  as a function of  $x$  in bins of  $Q^2$ . The inner error bars represent statistical error, the full error bars include the statistical and systematic uncertainties added in quadrature. The *solid curves* represent predictions of the DGLAP fit in the ACOT scheme. Results from other experiments are shown by the *open symbols*

H1 Collaboration



# Summary

- We have studied the  $F_2$  and  $F_L$  nucleon structure functions at small  $x$  in the framework of holographic QCD. The results are consistent with the data.
- We have also investigated the photon structure functions at small  $x$  in a similar model setup with a single adjustable parameter.
- Our calculations are in agreement with the experimental data and the GRS predictions.
  - Single Pomeron exchange works.
  - Vector meson dominance model works.
- The results can be tested at future linear colliders, such as the planned ILC.