

Chiral Thermodynamics with Charm

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- [1] C.S., Phys. Rev. D 90, no. 11, 114007 (2014).
- [2] C.S. and K. Redlich, arXiv:1412.7365 [hep-ph].

Introduction: quark flavors vs. conserved charges

- chiral: melting quark condensates, $T_{\text{u,d}} \sim 155 \text{ MeV} < T_{\text{s}} \sim 200 \text{ MeV}$
- deconfinement: conserved charges
 - B : baryon number, S : strangeness, Q : electric charge
 - $\mu = B\mu_B + S\mu_S + Q\mu_Q$
 - kurtosis of baryon number fluctuations

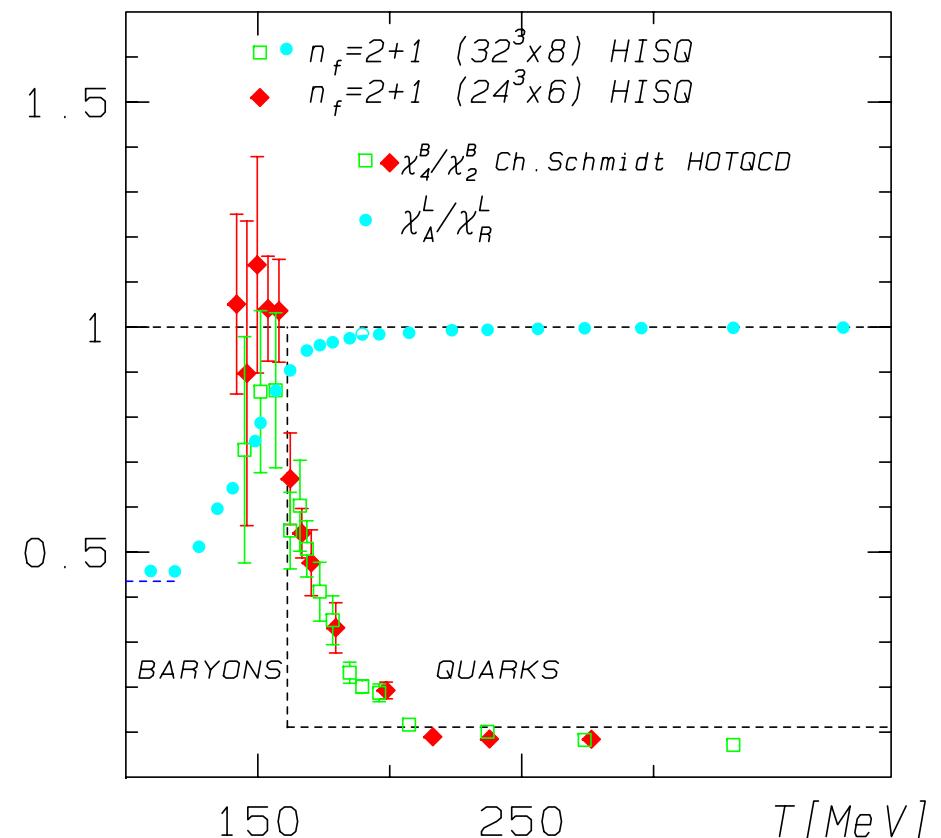
$$\kappa = \chi_4^B / \chi_2^B ,$$

$$\chi_n^B = \partial^n (P/T^4) / \partial (\mu_B/T)^n$$

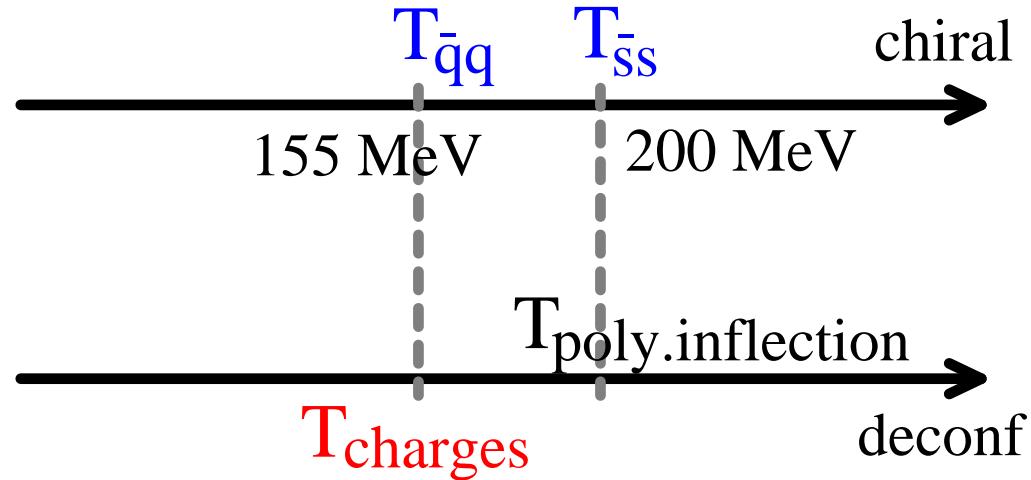
asymptotic values:

$$\kappa = B^2 = \begin{cases} 1, & T \ll T_c \\ 1/9, & T \gg T_c \end{cases}$$

- coincidence: $T_{\text{ch}} \simeq T_{\text{dec}}$



- crossover temperatures: *not unique!*



- flavor basis vs. conserved charge basis: strange mesons deconfined at T_{ch} !

$$\mu_u = \frac{1}{3}\mu_B + \frac{2}{3}\mu_Q, \quad \mu_d = \frac{1}{3}\mu_B - \frac{1}{3}\mu_Q, \quad \mu_s = \frac{1}{3}\mu_B - \frac{1}{3}\mu_Q - \mu_S.$$

- charm? . . . lessons from lattice QCD:

- EoS not affected by dynamical c quark around T_{ch} [Borsanyi et al. ('11)]
- charmed mesons deconfined together with light mesons [Basavov et al. ('14)]

- correlations between light and heavy-flavor physics

\Rightarrow how are heavy-light hadrons modified toward chiral crossover?

$D_s \sim c\bar{s}$ is like $K \sim q\bar{s}$? . . . NO!

I. Chiral Structure of Heavy-light Mesons

Symmetries of QCD in the heavy quark mass limit

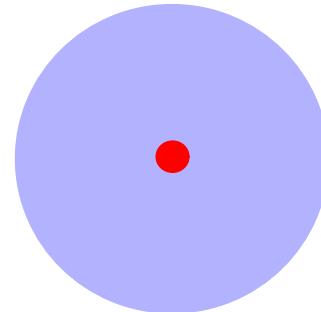
- flavor symmetries

chiral symmetry : $m_{u,d}/\Lambda_{\text{QCD}} \ll 1$, $m_s/\Lambda_{\text{QCD}} < 1$.

heavy quark symmetry : $\Lambda_{\text{QCD}}/m_{c,b} \ll 1$.

- heavy-light ($Q\bar{q}$) mesons Q : heavy quark and q : light quark

e.g. D mesons: $Q = c$, $q = u, d, s$

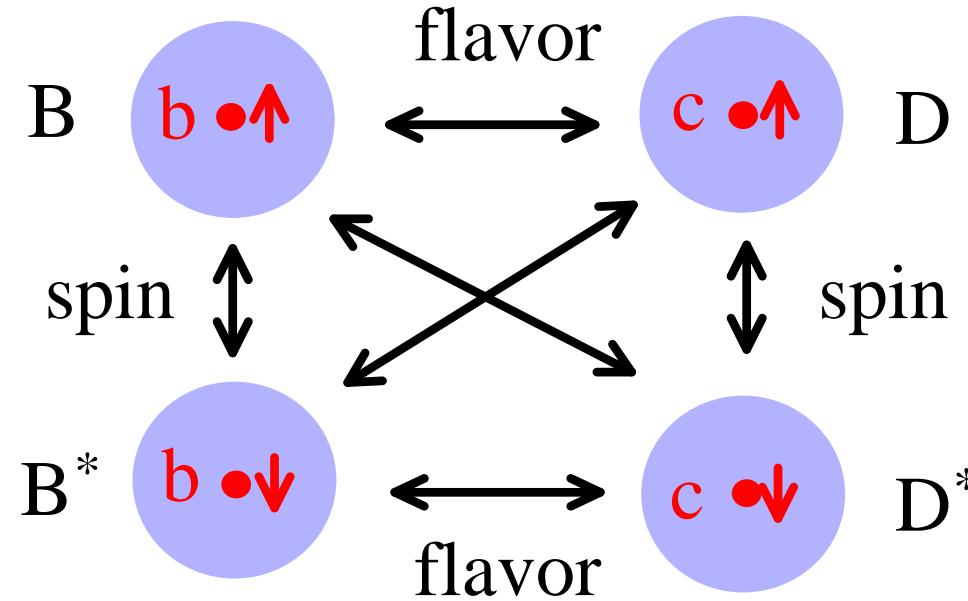


- physical picture ($m_Q \rightarrow \infty$)

- **flavor symmetry** ($c \leftrightarrow b$): cloud does not feel the flavor of Q .
 - **spin symmetry**: cloud does not feel the spin of Q .

Spin and flavor symmetries of heavy quarks are entangled!

- $SU(2N_{Q_f})$ **spin-flavor symmetry**: [Shuryak ('81), Isgur-Wise ('89)]
light d.o.f. (q) do not feel the flavor and spin of the heavy quark (Q).



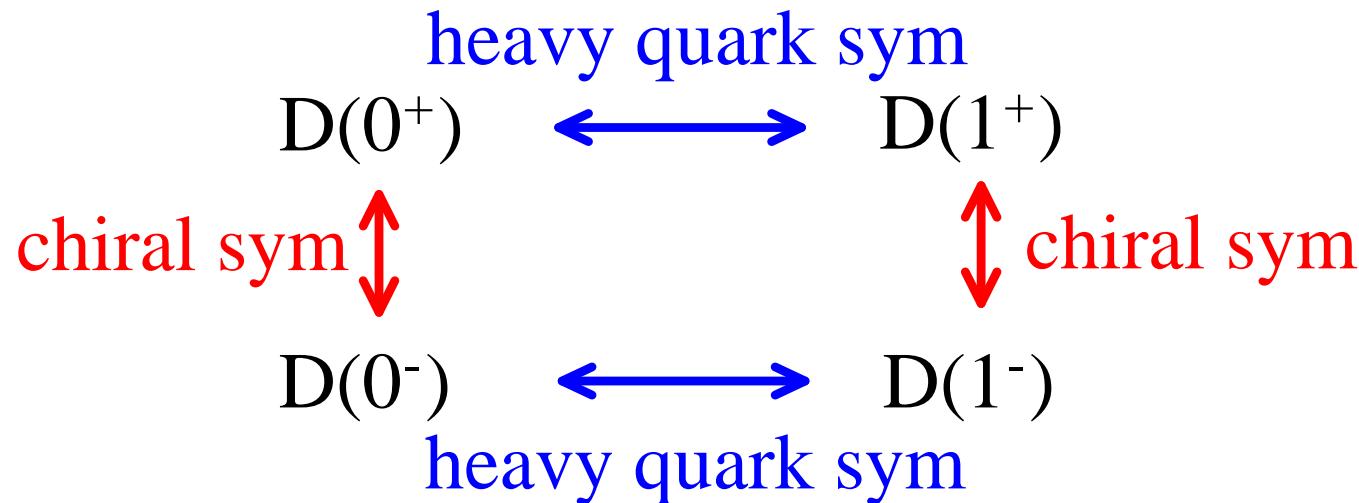
- spin partners: $D(0^-)$ and $D(1^-)$, $B(0^-)$ and $B(1^-)$
- **real world**:
 - $m_{D^*} - m_D = 142 \text{ MeV}$, $m_{B^*} - m_B = 46 \text{ MeV} \ll \Lambda_{\text{QCD}}$
 - ... $1/m_Q$ corrections
 - $m_{D_s} - m_{Dd} = 100 \text{ MeV}$, $m_{B_s} - m_{Bd} = 90 \text{ MeV} \ll \Lambda_{\text{QCD}}$
 - ... m_q corrections

Role of light flavor (chiral) symmetry

- **observation:** 2nd lowest spin doublets

$$\begin{array}{lll} D_{u,d}(0^+) : 2308 \text{ MeV} & [\text{Belle (03)}] & D_{u,d}(1^+) : 2427 \text{ MeV} & [\text{Belle (03)}] \\ D_s(0^+) : 2317 \text{ MeV} & [\text{Babar (03)}] & D_s(1^+) : 2460 \text{ MeV} & [\text{CLEO (03)}] \end{array}$$

- mass difference of parity doublets: $\delta m = 300 - 400 \text{ MeV} \sim \Lambda_{\text{QCD}}$
NOTE: potential model for D mesons (cf. hydrogen atom) does not work!
- chiral doubling [Nowak-Rho-Zahed (92); Bardeen-Hill (93)]



effective theory for heavy-light system based on the two relevant symmetries

II. Thermodynamics

Embedding D, D_s in a linear sigma model

- chiral fields $\Sigma = \sigma + i\pi$, heavy-light meson fields $H(0^-, 1^-), G(0^+, 1^+)$

$$\Sigma \rightarrow g_L \Sigma g_R^\dagger, \quad \mathcal{H}_{L,R} \rightarrow S \mathcal{H}_{L,R} g_{L,R}^\dagger.$$

- Lagrangian

$$\mathcal{L} = \mathcal{L}_L(\Sigma) + \mathcal{L}_{HL}(\mathcal{H}, \Sigma), \quad V_{HL} = V_{HL}(\mathcal{H}^2, \mathcal{H}^4; \Sigma) + V_{HL}^{(\text{exp})}.$$

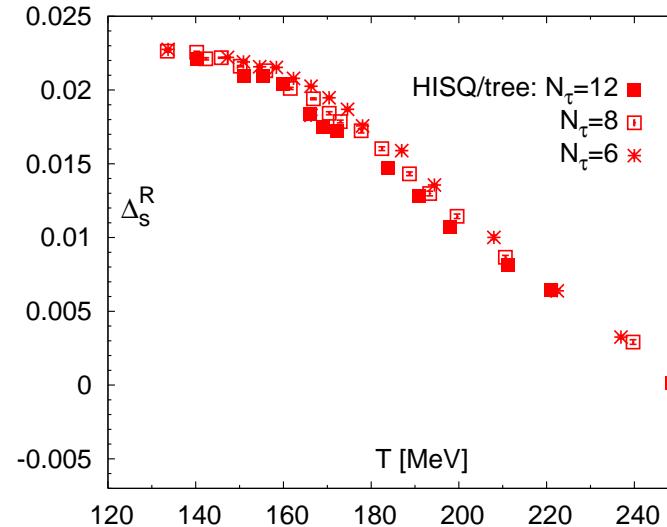
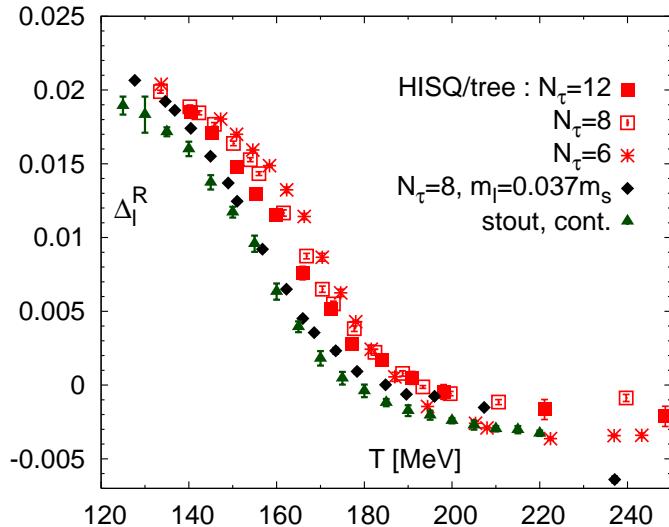
- 6 parameters fixed with $T = 0$ physics

$$V_{HL}^{(2)} : m_0, \underbrace{g_\pi^q, g_\pi^s}_{\Sigma \leftrightarrow \mathcal{H}^2}, \quad V_{HL}^{(4)} : k_0, \underbrace{k_q, k_s}_{\Sigma \leftrightarrow \mathcal{H}^4}$$

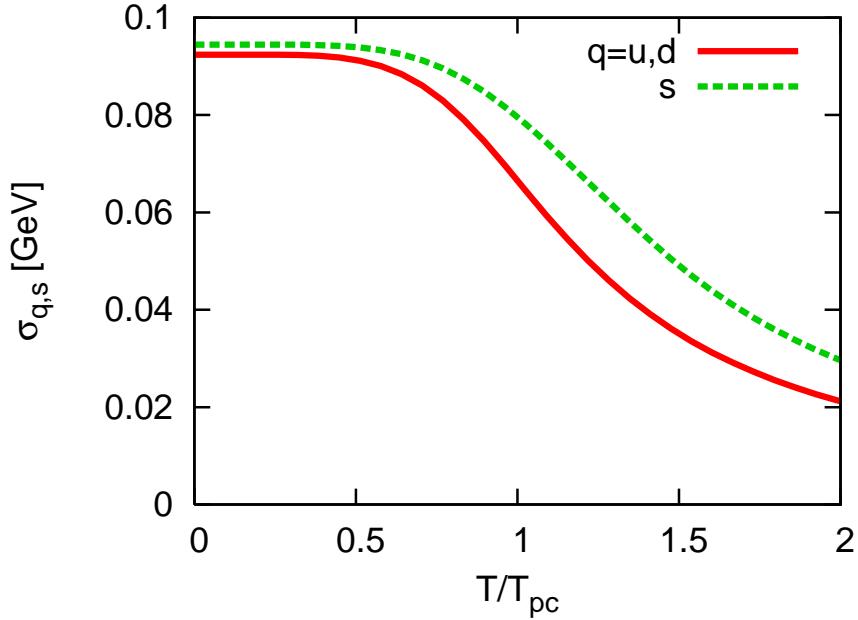
- isospin sym & mean field approximation: $\langle \sigma_q \rangle, \langle \sigma_s \rangle, \langle D_q \rangle, \langle D_s \rangle$

conventional approach ... then?

Chiral condensates: role of charmed-meson MF



[HotQCD Collaboration ('12)]

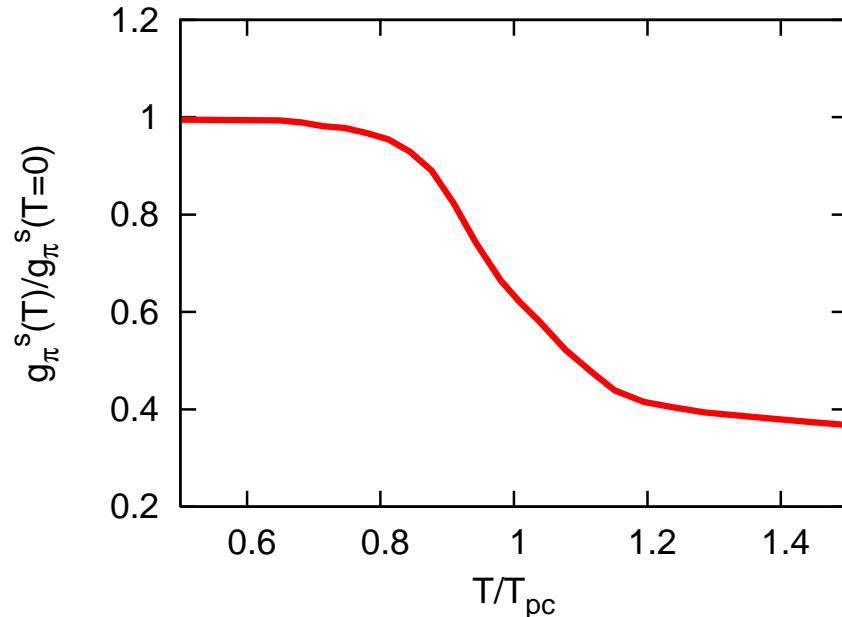
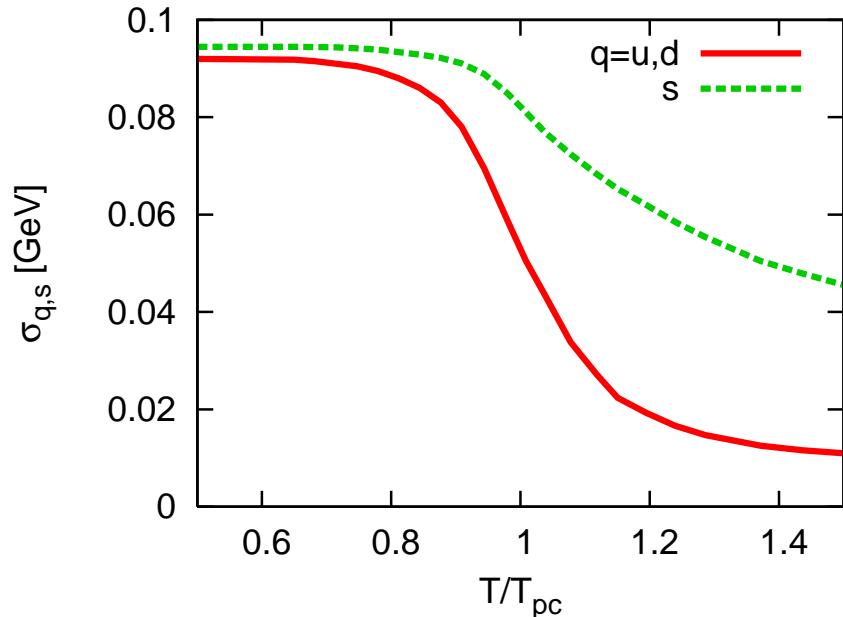


- lattice: qualitative diff. between $\langle \bar{q}q \rangle$ and $\langle \bar{s}s \rangle \dots$ $SU(2+1)$: $T_c^{(u,d)} < T_c^{(s)}$
- chiral model: $\sigma_{q,s}$ — approx. $SU(3)!$?
- induced chiral sym. breaking:

$$h_q^* = h_q - D_q^2 \left(\frac{1}{2} g_\pi^q + 2k_q D_q^2 \right),$$

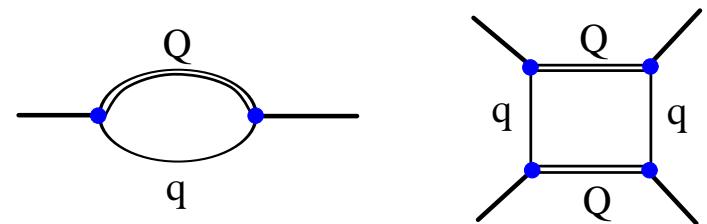
$$h_s^* = h_s - \frac{1}{\sqrt{2}} D_s^2 \left(\frac{1}{2} g_\pi^s + 2k_s D_s^2 \right).$$

Intrinsic thermal effects



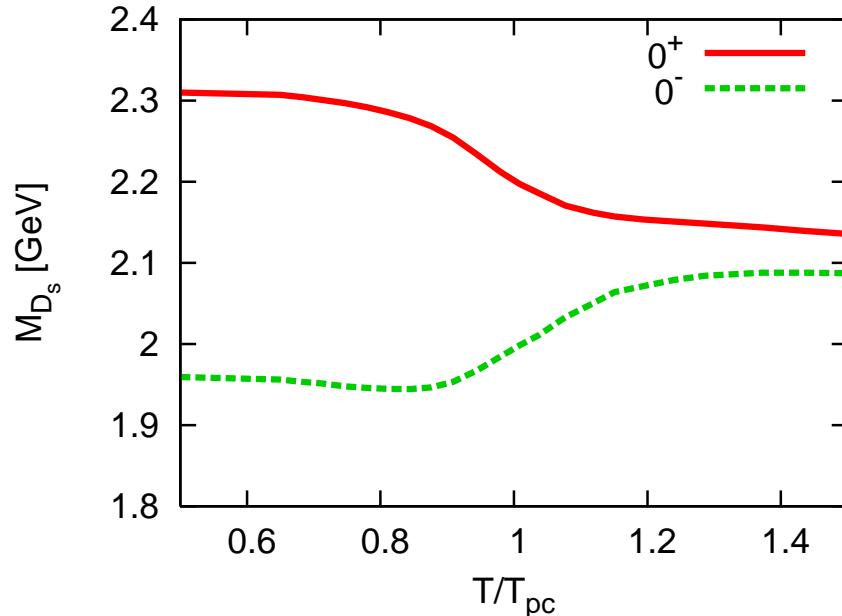
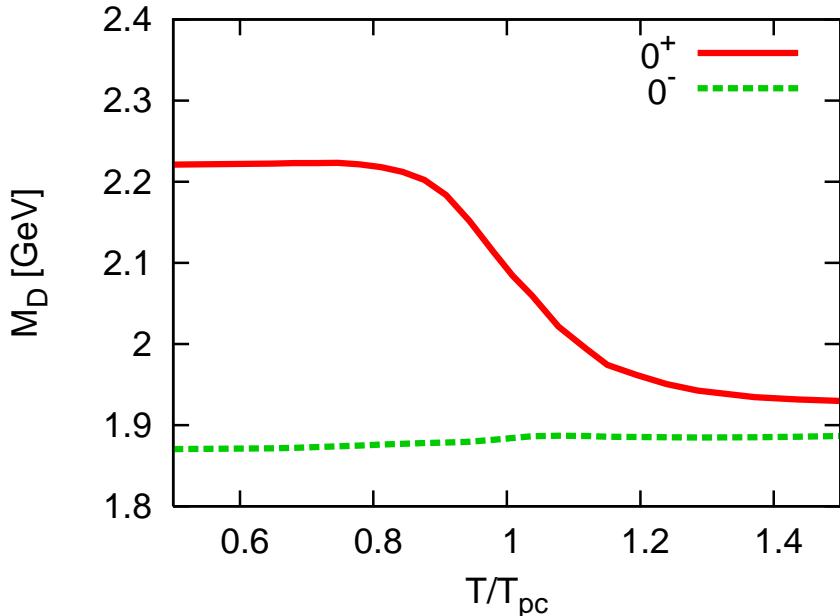
- concept of EFT: generating functional, Green's functions

$$Z = \int \mathcal{D}q \mathcal{D}g e^{S_{\text{QCD}}[q,g]} \equiv \int \mathcal{D}U e^{S_{\text{eff}}[U]}$$

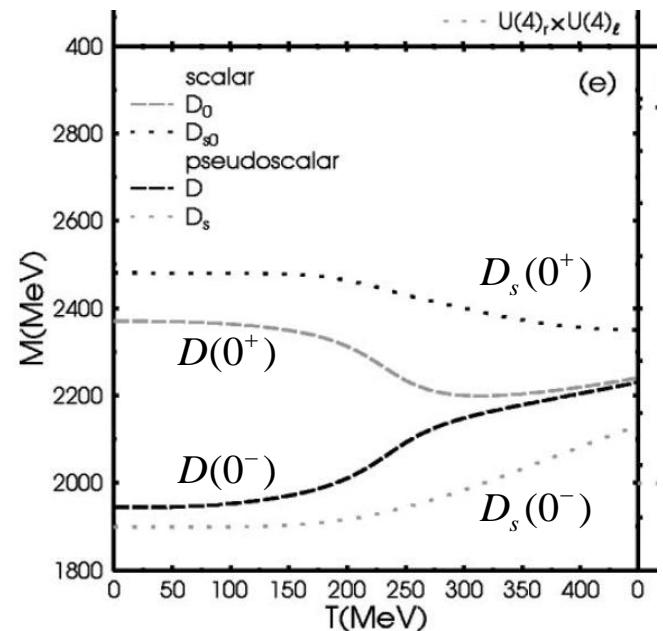


- low-energy constants: high-frequency modes integrated out
 \Rightarrow in a hot/dense medium: effective couplings dep. on T/n
- L: $T_{pc}^{\text{lat}} = 154$ MeV $\Rightarrow m_\sigma = 400$ MeV
 HL: $\sigma_{q,s}$ profiles from lattice QCD $\Rightarrow g_\pi^{q,s}(T)$ etc.

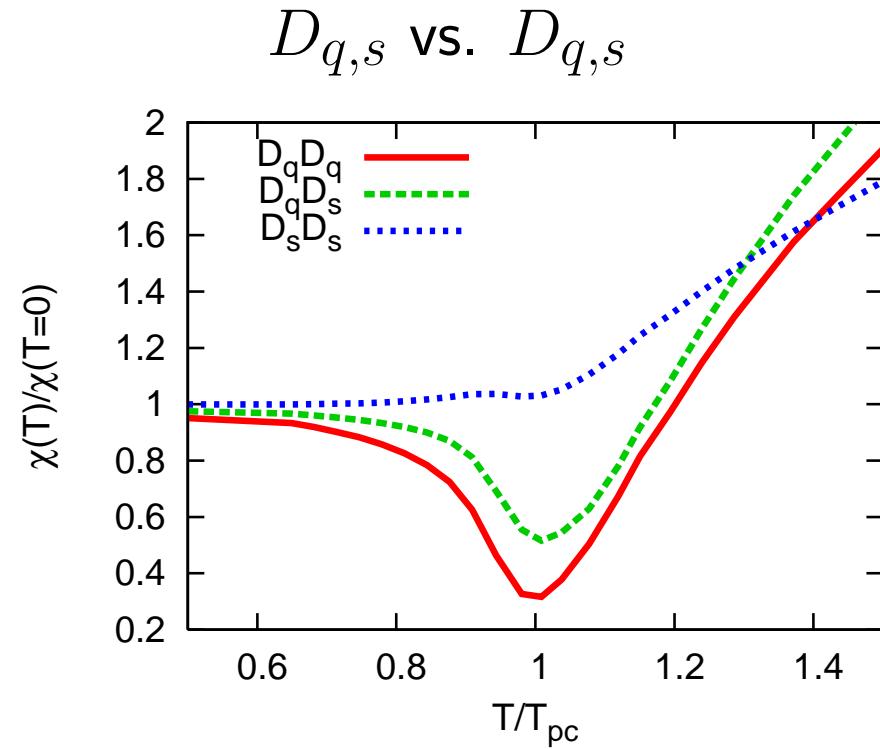
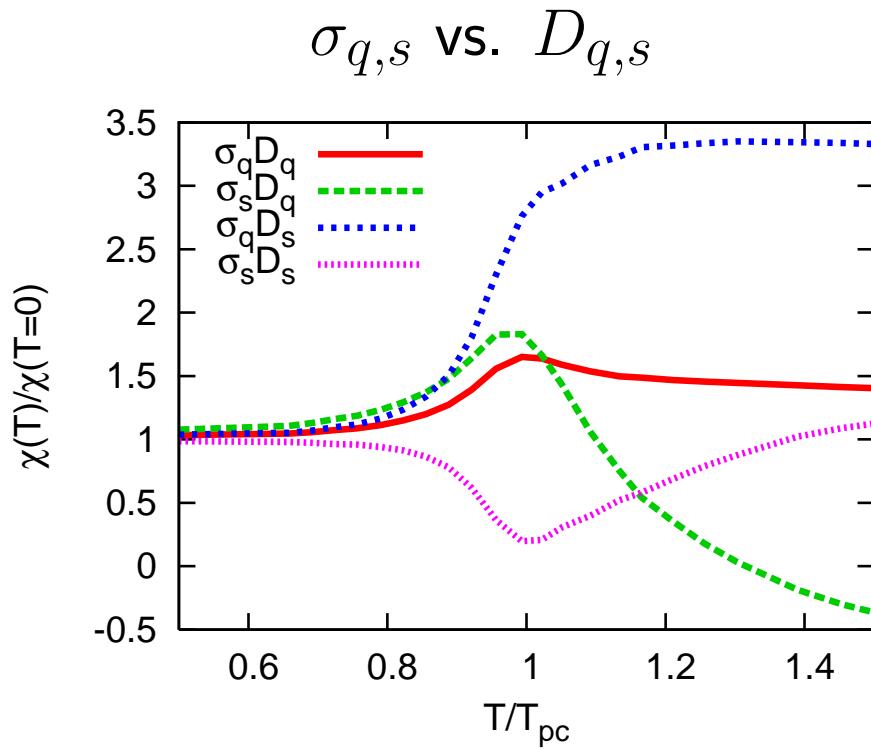
In-medium charmed-meson masses



- chiral splitting at T_{pc} : $\delta M_D \simeq \delta M_{D_s}$
 ... *insensitive to light flavors!*
 ⇒ heavy quark symmetry
- light mesons at T_{pc} : $\delta M_{\pi-\sigma} \ll \delta M_{K-\kappa}$
 ... $SU(2+1) \neq SU(3)$
- cf. chiral $SU(4)$: [Roder-Ruppert-Rischke ('03)]
 $\delta M_D \ll \delta M_{D_s}$



Correlations between light and heavy-light mesons [CS-Redlich ('14)]



qualitative changes set in at $T \sim T_{pc}$: (NOTE: $\chi_{ch} \sim \partial \sigma_{q,s} / \partial m_{q,s}$)

$$\begin{aligned}\hat{\chi}_{\sigma D} &= -\hat{\chi}_{ch} \hat{\mathcal{C}}_{HL} \hat{\chi}_D, & \hat{\chi}_{D\sigma} &= -\hat{\chi}_D \hat{\mathcal{C}}_{HL} \hat{\chi}_{ch}, \\ \hat{\chi}_{DD} &= \hat{\mathcal{C}}_D - \hat{\mathcal{C}}_{HL} \hat{\chi}_{ch} \hat{\mathcal{C}}_{HL} \equiv \hat{\chi}_D.\end{aligned}$$

in-medium D_s as a probe of O(4)!

Chiral restoration vs. deconfinement in QCD

- small μ : only *crossover*, $T_{\text{ch}} \sim T_{\text{dec}}$ from lattice studies

To which extent is the hadronic picture reliable?

- **lesson 1:** [Rapp, van Hees]
in-medium ρ meson vs. SPS dilepton data
- **lesson 2:** [Kitazawa, Kunihiro, Nemoto]
 $N_f = 2$ NJL, mesonic collective modes above chiral crossover
 \Rightarrow talk by T. Kunihiro (Thu)
- **lesson 3:** [Gattringer; Bruckmann, Gattringer, Hagen; Synatschke, Wipf, Langfeld; Gongyo, Doi, Iritani, Suganuma; Glozman, Lang, Schrock]
LL Dirac eigenmodes, confined phase with unbroken chiral symmetry
 \Rightarrow talk by T. Doi (Thu)
- large μ : ???
quark-meson-nucleon hybrid model at $T = 0$ and large μ_B
 $\Rightarrow \rho_{\text{ch}}$ separated from ρ_{dec} [Benic- Mishustin-CS, ('15)]

Summary and Remarks

- **Synthesis of light and heavy quark dynamics**

$$\frac{m_q}{m_c} \sim \mathcal{O}(10^{-3}), \frac{m_s}{m_c} \sim \mathcal{O}(10^{-2}) \ll 1$$

- at T_{pc} : chiral mass splittings of HL mesons insensitive to light flavors.

$$\delta M_{D,B} \simeq \delta M_{D_s,B_s} \quad \text{vs.} \quad \delta M_{\pi-\sigma} \ll \delta M_{K-\kappa}$$

- remnant of O(4) in HL mixed fluctuations.
- anomalous suppression of D_s decay widths as a sign of CSR

in-medium D_s as a probe of O(4)!

- **Issues**

- dissociation vs. local field approximation
 - · · chiral restored phase with confinement? — topology, lattice!
- application to a dense system *under* strange and charm number conservation