

Light nuclei from 2+1 flavor lattice QCD

Takeshi Yamazaki



University of Tsukuba

K.-I. Ishikawa, Y. Kuramashi, and A. Ukawa

Refs: PRD81:111504(R)(2010); PRD84:054506(2011); PRD86:074514(2012)

[arXiv:1502.04182\[hep-lat\]](https://arxiv.org/abs/1502.04182)

Hadrons and Hadron Interactions in QCD 2015 – *Effective Theories and Lattice* –

© Yukawa Institute for Theoretical Physics, Feb. 15–Mar. 21 2015

Outline

1. Introduction
2. Calculation method of nuclei in lattice QCD
3. Simulation parameters
4. Results
 - ${}^4\text{He}$ and ${}^3\text{He}$ channels
 - NN channels
5. Summary and future work

Introduction

Binding force $\left\{ \begin{array}{l} \text{protons and neutrons} \rightarrow \text{nuclei} \\ \text{quarks and gluons} \rightarrow \text{protons and neutrons} \end{array} \right.$
both from fundamental strong interaction of quark and gluon
well known, but hard to prove

Spectrum of nuclei: Shell model

degrees of freedom of protons and neutrons

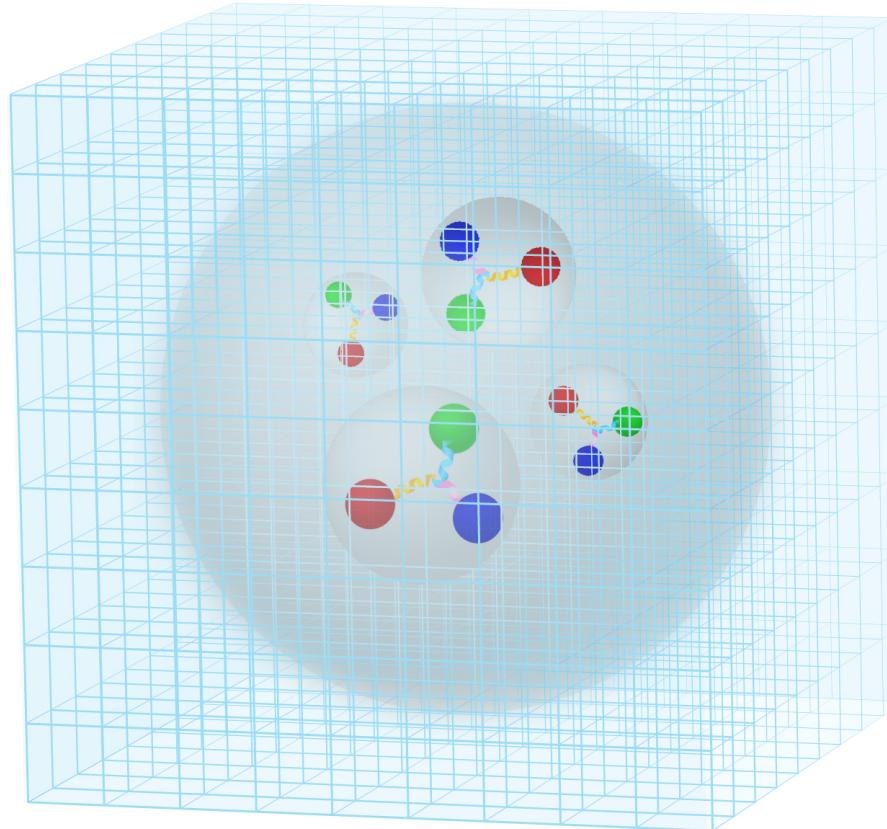
Spectrum of proton and neutron (nucleons)

success of non-perturbative calculation of QCD
degrees of freedom of quarks and gluons

goal: quantitatively understand property of nuclei from QCD

quarks and gluons $\xrightarrow{\text{Shell model}}$ protons and neutrons \rightarrow nuclei
 $\xrightarrow{\text{lattice QCD}}$

Ultimate goal of lattice QCD



<http://www.jicfus.jp/jp/promotion/pr/mj/2014-1/>

quantitatively understand property of nuclei from QCD

Introduction

Motivation :

Understand property of nuclei from QCD

If we can study nuclei from QCD, we may be able to

1. reproduce spectrum of nuclei
2. predict property of nuclei hard to calculate or observe
such as neutron rich nuclei

So far not so many studies for multi-baryon bound states

Before studying such difficult problems, we should study

→ Can we reproduce known binding energy in light nuclei?

Multi-baryon bound state from lattice QCD

Not observed before '09 (except H-dibaryon '88 Iwasaki *et al.*)

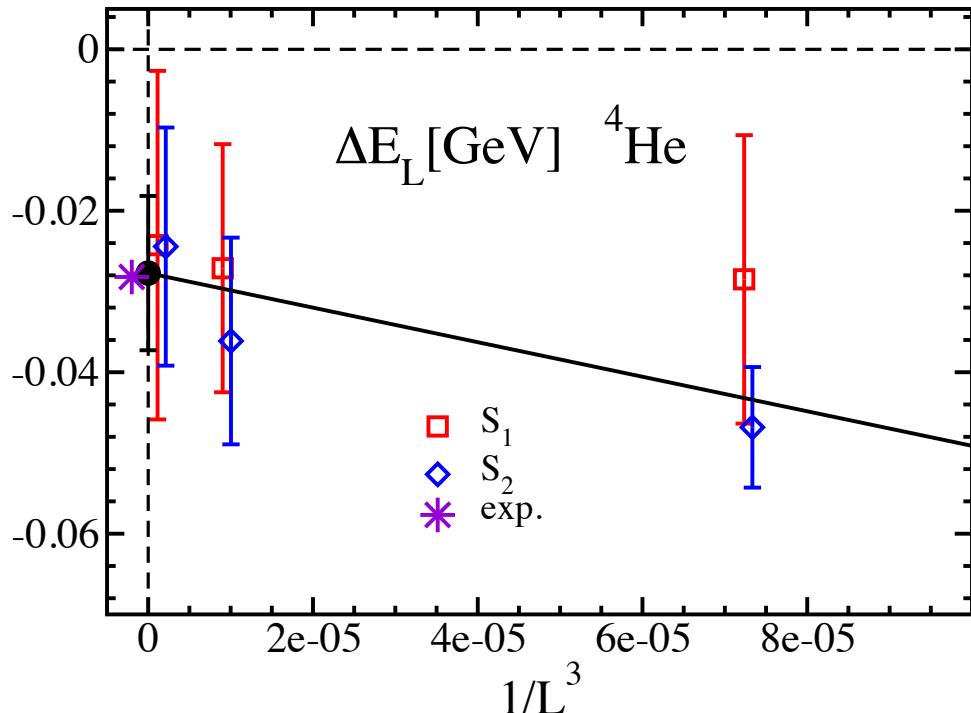
1. ^4He and ^3He

'10 PACS-CS $N_f = 0$ $m_\pi = 0.8$ GeV PRD81:111504(R)(2010)

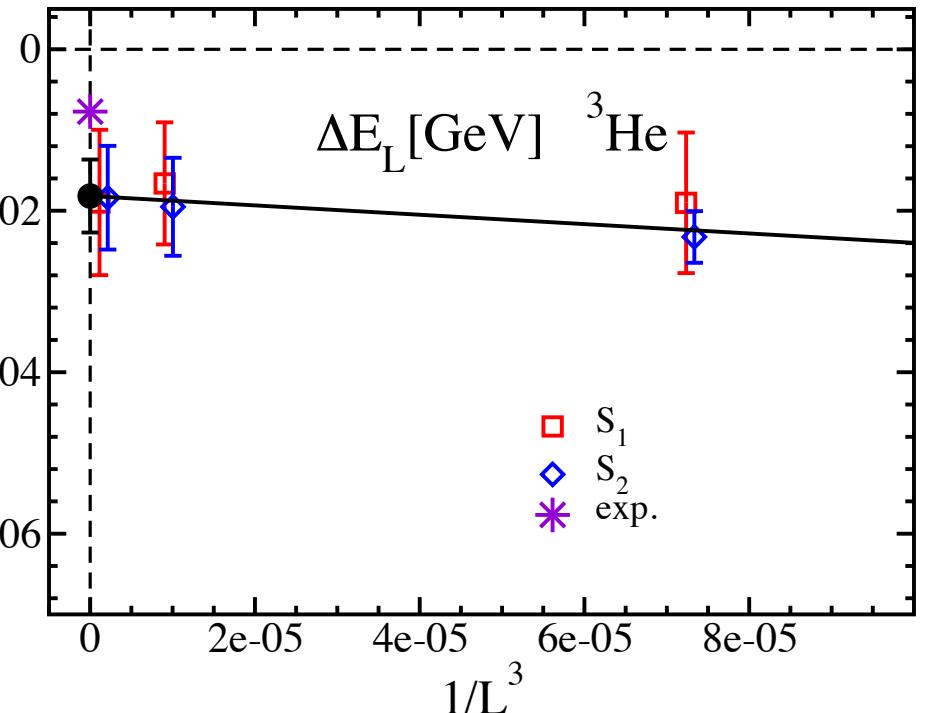
Exploratory study of three- and four-nucleon systems

PACS-CS Collaboration, PRD81:111504(R)(2010)

Identification of bound state from volume dependence of ΔE



$$\Delta E_{4\text{He}} = 27.7(7.8)(5.5) \text{ MeV}$$



$$\Delta E_{3\text{He}} = 18.2(3.5)(2.9) \text{ MeV}$$

1. Observe bound state in both channels
2. Same order of ΔE to experiment

Several systematic errors included, e.g., $N_f = 0$, $m_\pi = 0.8$ GeV

Multi-baryon bound state from lattice QCD

Extend our exploratory study to $N_f = 2 + 1$ calculation

1. ${}^4\text{He}$ and ${}^3\text{He}$

'10 PACS-CS $N_f = 0$ $m_\pi = 0.8$ GeV PRD81:111504(R)(2010)

'12 HALQCD $N_f = 3$ $m_\pi = 0.47$ GeV, $m_\pi > 1$ GeV ${}^4\text{He}$

'12 NPLQCD $N_f = 3$ $m_\pi = 0.81$ GeV

'12 TY et al. $N_f = 2 + 1$ $m_\pi = 0.51$ GeV PRD86:074514(2012)

'15 TY et al. $N_f = 2 + 1$ $m_\pi = 0.30$ GeV arXiv:1502.04182

2. H dibaryon in $\Lambda\Lambda$ channel ($S=-2$, $I=0$)

'11, '12 NPLQCD $N_f = 2 + 1$ $m_\pi = 0.39$ GeV, $N_f = 3$ $m_\pi = 0.81$ GeV

'11, '12 HALQCD $N_f = 3$ $m_\pi = 0.47\text{--}1.02$ GeV

'11 Luo et al. $N_f = 0$ $m_\pi = 0.5\text{--}1.3$ GeV

3. NN

'11 PACS-CS $N_f = 0$ $m_\pi = 0.8$ GeV PRD84:054506(2011)

'12 NPLQCD $N_f = 2 + 1$ $m_\pi = 0.39$ GeV (Possibility)

'12 NPLQCD $N_f = 3$ $m_\pi = 0.81$ GeV

'12 TY et al. $N_f = 2 + 1$ $m_\pi = 0.51$ GeV PRD86:074514(2012)

'15 TY et al. $N_f = 2 + 1$ $m_\pi = 0.30$ GeV arXiv:1502.04182

Other states: $\Xi\Xi$, '12 NPLQCD; spin-2 $N\Omega$, ${}^{16}\text{O}$ and ${}^{40}\text{Ca}$, '14 HALQCD

Calculation method of multi-nucleon bound state

Traditional method for example ${}^4\text{He}$ channel

$$\langle 0 | O_{{}^4\text{He}}(t) O_{{}^4\text{He}}^\dagger(0) | 0 \rangle = \sum_n \langle 0 | O_{{}^4\text{He}} | n \rangle \langle n | O_{{}^4\text{He}}^\dagger | 0 \rangle e^{-E_n t} \xrightarrow[t \gg 1]{} A_0 e^{-E_0 t}$$

Difficulties for multi-nucleon calculation

1. Statistical error

$$\text{Statistical error} \propto \exp \left(N_N \left[m_N - \frac{3}{2} m_\pi \right] t \right)$$

2. Calculation cost

$$\begin{aligned} \text{Wick contraction for } {}^4\text{He} &= p^2 n^2 = (udu)^2 (dud)^2: 518400 \\ \text{proton} &= p = (udu): 2 \end{aligned}$$

3. Identification of bound state on finite volume

Finite volume effect of attractive scattering state

$$\Delta E_L = E_0 - N_N m_N = O(L^{-3}) < 0 \leftrightarrow \text{binding energy}$$

Calculation method of multi-nucleon bound state

Traditional method for example ${}^4\text{He}$ channel

$$\langle 0 | O_{{}^4\text{He}}(t) O_{{}^4\text{He}}^\dagger(0) | 0 \rangle = \sum_n \langle 0 | O_{{}^4\text{He}} | n \rangle \langle n | O_{{}^4\text{He}}^\dagger | 0 \rangle e^{-E_n t} \xrightarrow[t \gg 1]{} A_0 e^{-E_0 t}$$

Difficulties for multi-nucleon calculation

1. Statistical error

$$\text{Statistical error} \propto \exp \left(N_N \left[m_N - \frac{3}{2} m_\pi \right] t \right)$$

2. Calculation cost

$$\begin{aligned} \text{Wick contraction for } {}^4\text{He} &= p^2 n^2 = (udu)^2 (dud)^2: 518400 \\ \text{proton} &= p = (udu): 2 \end{aligned}$$

Most severe problem before '09: (every t) $\times N_{\text{meas}} \sim O(10^6)$

3. Identification of bound state on finite volume

Finite volume effect of attractive scattering state

$$\Delta E_L = E_0 - N_N m_N = O(L^{-3}) < 0 \leftrightarrow \text{binding energy}$$

Calculation method of multi-nucleon bound state

Traditional method for example ${}^4\text{He}$ channel

$$\langle 0 | O_{{}^4\text{He}}(t) O_{{}^4\text{He}}^\dagger(0) | 0 \rangle = \sum_n \langle 0 | O_{{}^4\text{He}} | n \rangle \langle n | O_{{}^4\text{He}}^\dagger | 0 \rangle e^{-E_n t} \xrightarrow[t \gg 1]{} A_0 e^{-E_0 t}$$

Difficulties for multi-nucleon calculation

1. Statistical error

$$\text{Statistical error} \propto \exp \left(N_N \left[m_N - \frac{3}{2} m_\pi \right] t \right)$$

→ heavy quark $m_\pi = 0.8\text{--}0.3 \text{ GeV}$ + large # of measurements

2. Calculation cost PACS-CS PRD81:111504(R)(2010)

Wick contraction for ${}^4\text{He} = p^2 n^2 = (udu)^2 (dud)^2$: 518400 → 1107

→ reduction using $p(n) \leftrightarrow p(n)$, $p \leftrightarrow n$, $u(d) \leftrightarrow u(d)$ in $p(n)$
+ block of 3 quark props(parallel) and contraction(workstation)

Multi-baryon: '12 Doi and Endres; Detmold and Orginos; '13 Günther et al.

3. Identification of bound state on finite volume

attractive scattering state $\Delta E_L = E_0 - N_N m_N = O(L^{-3}) < 0$

'86, '91 Lüscher, '07 Beane et al. [Sharpe's talk]

→ Volume dependence of $\Delta E_L \rightarrow \Delta E_\infty \neq 0 \rightarrow$ bound state

Spectral weight: '04 Mathur et al., Anti-PBC '05 Ishii et al.

Calculation method of multi-nucleon bound state

Traditional method for example ${}^4\text{He}$ channel

$$\langle 0 | O_{{}^4\text{He}}(t) O_{{}^4\text{He}}^\dagger(0) | 0 \rangle = \sum_n \langle 0 | O_{{}^4\text{He}} | n \rangle \langle n | O_{{}^4\text{He}}^\dagger | 0 \rangle e^{-E_n t} \xrightarrow[t \gg 1]{} A_0 e^{-E_0 t}$$

Difficulties for multi-nucleon calculation

1. Statistical error

$$\text{Statistical error} \propto \exp \left(N_N \left[m_N - \frac{3}{2} m_\pi \right] t \right)$$

Most severe problem at present

2. Calculation cost PACS-CS PRD81:111504(R)(2010)

$$\begin{aligned} \text{Wick contraction for } {}^4\text{He} &= p^2 n^2 = (udu)^2 (dud)^2: 518400 \\ \text{proton} &= p = (udu): 2 \end{aligned}$$

Used to be most severe problem

3. Identification of bound state on finite volume

Finite volume effect of attractive scattering state

$$\Delta E_L = E_0 - N_N m_N = O(L^{-3}) < 0 \leftrightarrow \text{binding energy}$$

Simulation parameters

$N_f = 2 + 1$ QCD

Iwasaki gauge action at $\beta = 1.90$

$a^{-1} = 2.194$ GeV with $m_\Omega = 1.6725$ GeV '10 PACS-CS

non-perturbative $O(a)$ -improved Wilson fermion action

$m_\pi = 0.51$ GeV and $m_N = 1.32$ GeV PRD86:074514(2012)

$m_\pi = 0.30$ GeV and $m_N = 1.05$ GeV arXiv:1502.04182

$m_s \sim$ physical strange quark mass

${}^4\text{He}$, ${}^3\text{He}$, NN(${}^3\text{S}_1$ and ${}^1\text{S}_0$)

		$m_\pi = 0.5$ GeV		$m_\pi = 0.3$ GeV		R
L	L [fm]	N_{conf}	N_{meas}	N_{conf}	N_{meas}	
32	2.9	200	192			
40	3.6	200	192			
48	4.3	200	192	400	1152	12
64	5.8	190	256	160	1536	5

$$R = (N_{\text{conf}} \cdot N_{\text{meas}})_{0.3\text{GeV}} / (N_{\text{conf}} \cdot N_{\text{meas}})_{0.5\text{GeV}}$$

Computational resources

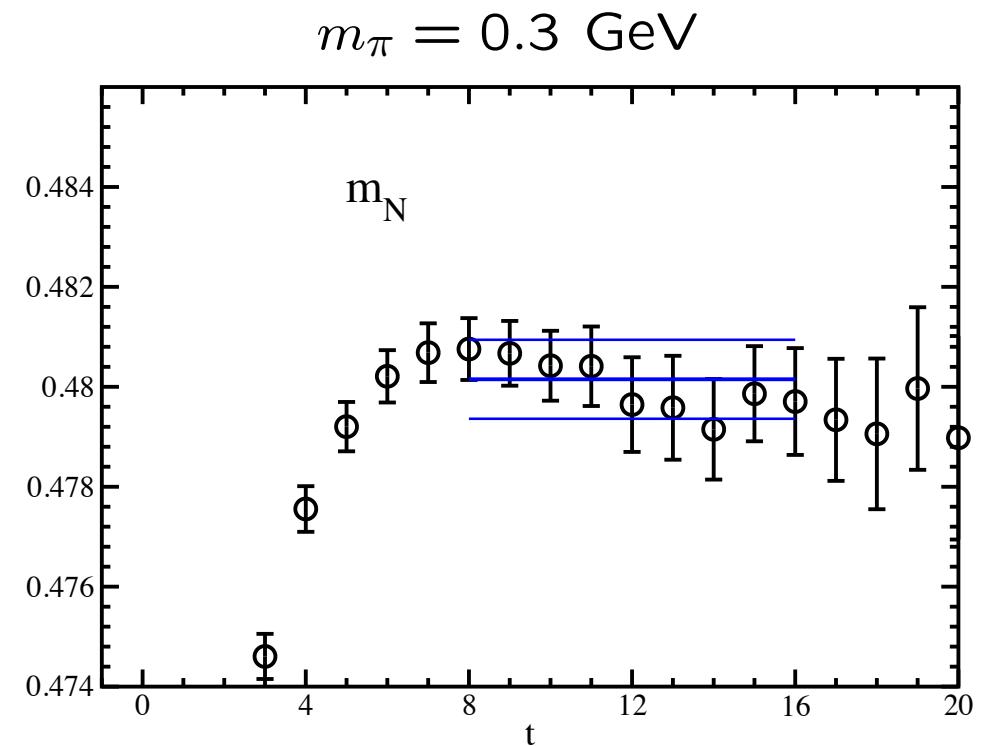
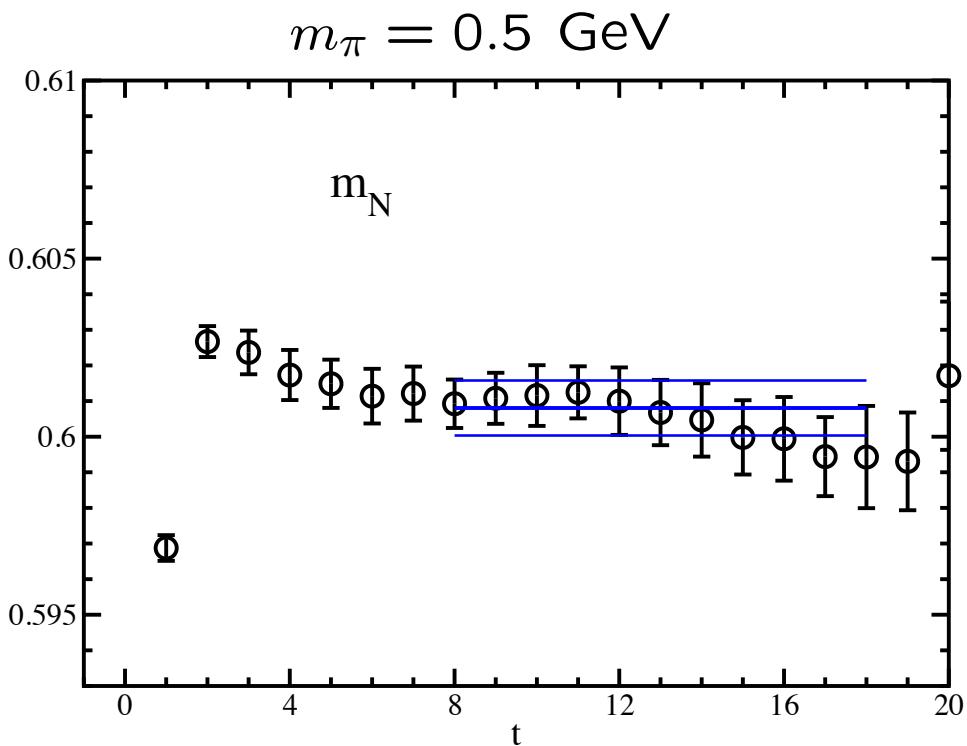
PACS-CS, T2K-Tsukuba, HA-PACS, COMA at Univ. of Tsukuba

T2K-Tokyo and FX10 at Univ. of Tokyo, and K at AICS

Results

Effective mass of nucleon on $L = 5.8$ fm

$$\text{Effective } m_N = \log \left(\frac{C_N(t)}{C_N(t+1)} \right)$$



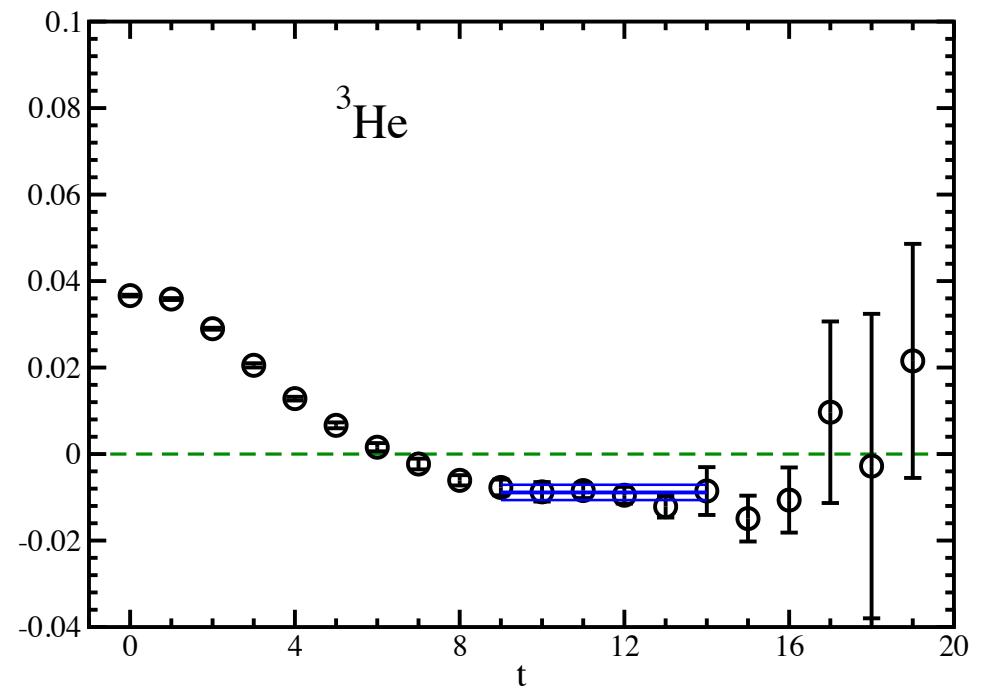
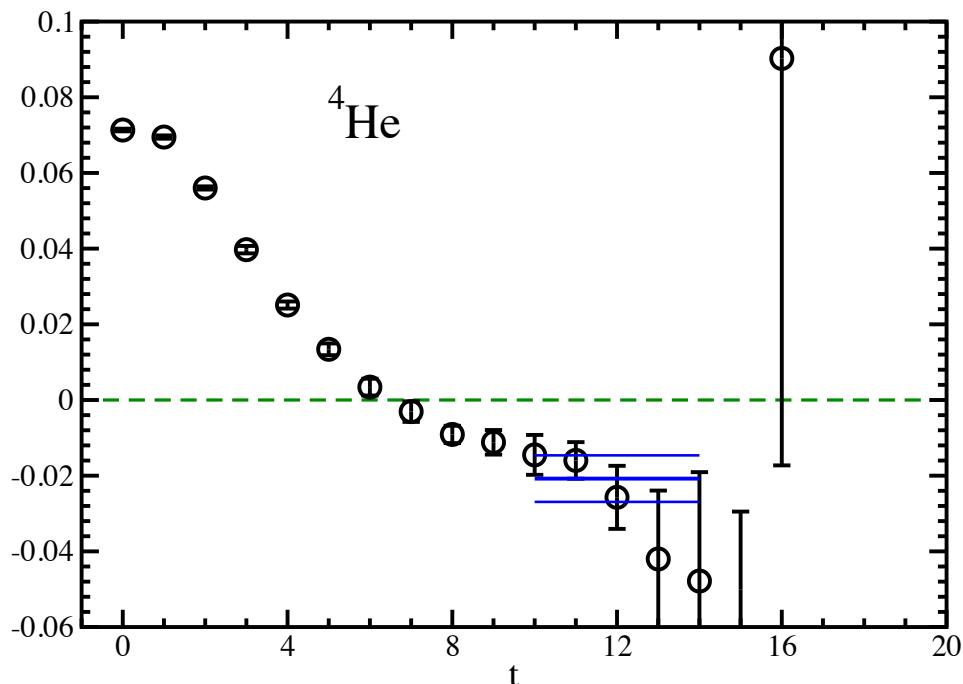
- Good plateau $t \gtrsim 7$
- Statistical error $< 0.2\%$

$\Delta E_L = E_0 - N_N m_N$ in ^4He and ^3He channels

at $m_\pi = 0.5$ GeV on $L = 5.8$ fm

TY et al., PRD86:074514(2012)

$$\Delta E_L = \log \left(\frac{R_{^4\text{He}}(t)}{R_{^4\text{He}}(t+1)} \right) \text{ with } R_{^4\text{He}}(t) = \frac{C_{^4\text{He}}(t)}{(C_N(t))^4}$$

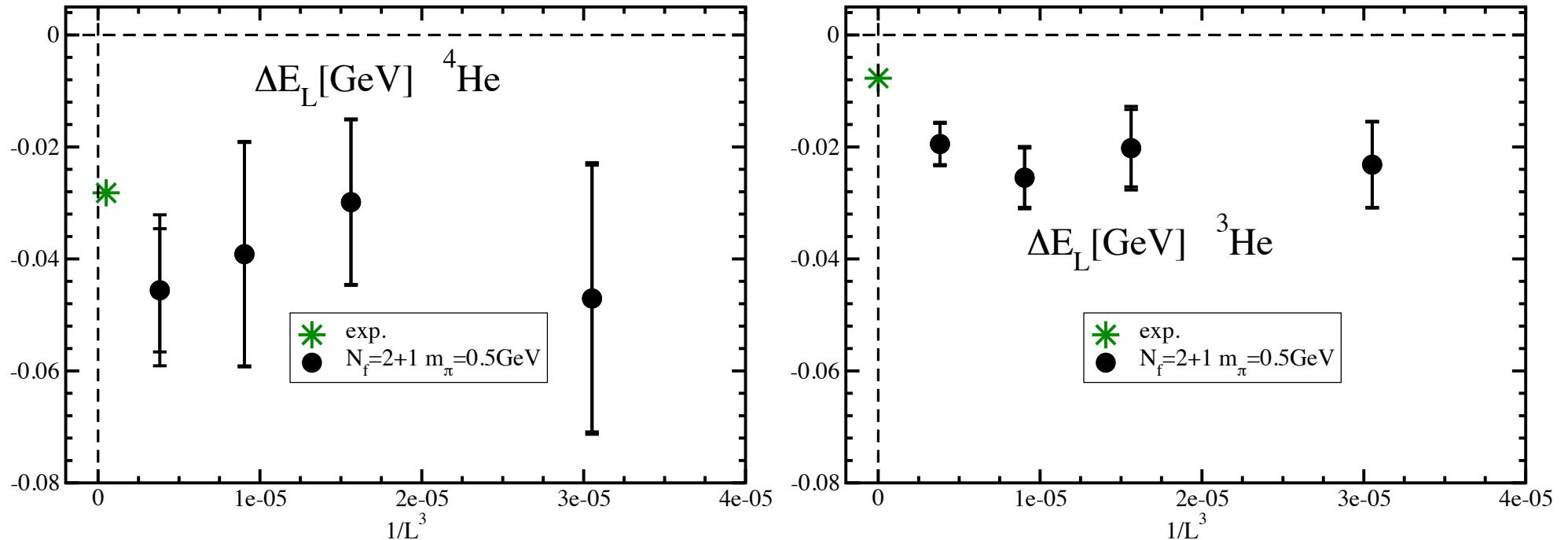


- Larger error in ${}^4\text{He}$ channel
- Statistical error under control in $t < 12$
- Negative ΔE_L in both channels

${}^4\text{He}$ and ${}^3\text{He}$ channels $\Delta E_L = E_0 - N_N m_N$ at $m_\pi = 0.5$ GeV

TY et al., PRD86:074514(2012)

Identification of bound state from volume dependence of ΔE

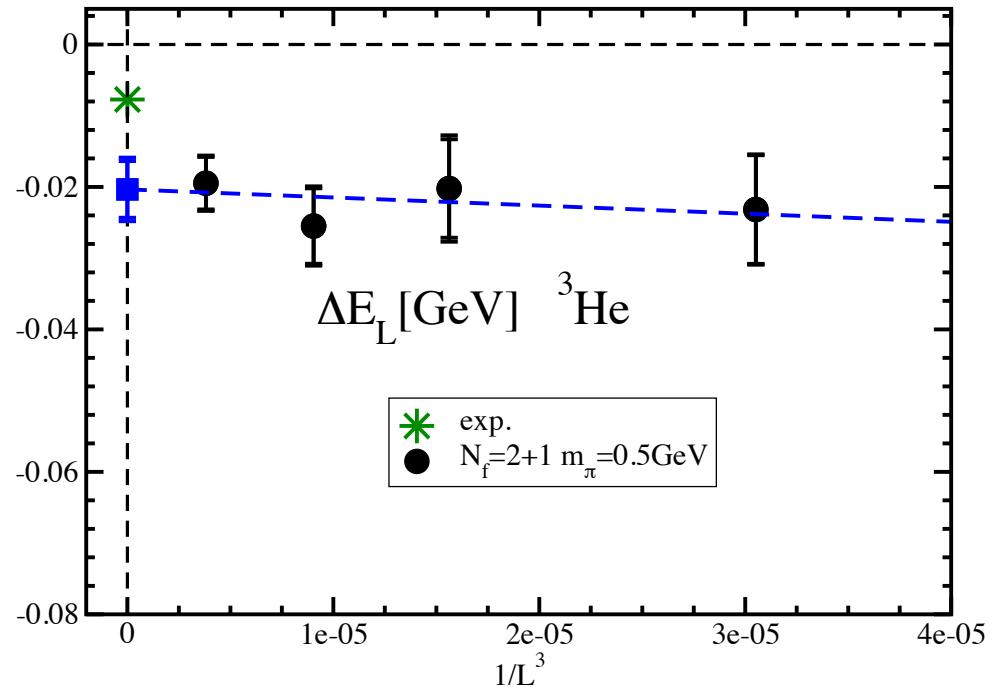
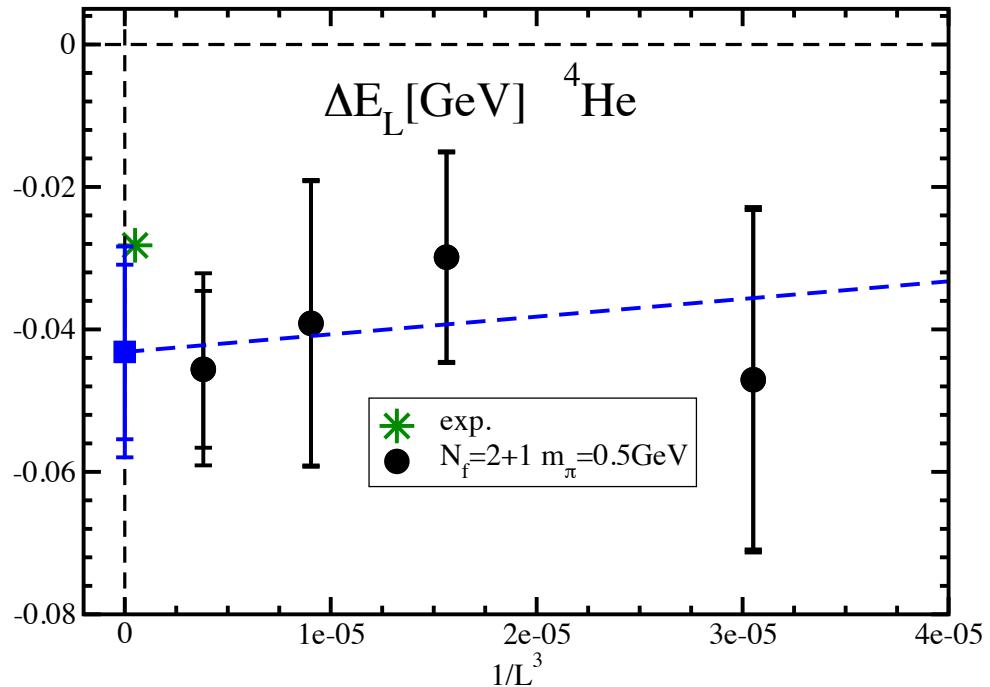


- $\Delta E_L < 0$ and mild volume dependence
- Infinite volume extrapolation with $\Delta E_L = -\Delta E_{\text{bind}} + C/L^3$
small difference with $\exp(-cL)$ fit due to large error

${}^4\text{He}$ and ${}^3\text{He}$ channels $\Delta E_L = E_0 - N_N m_N$ at $m_\pi = 0.5$ GeV

TY et al., PRD86:074514(2012)

Identification of bound state from volume dependence of ΔE

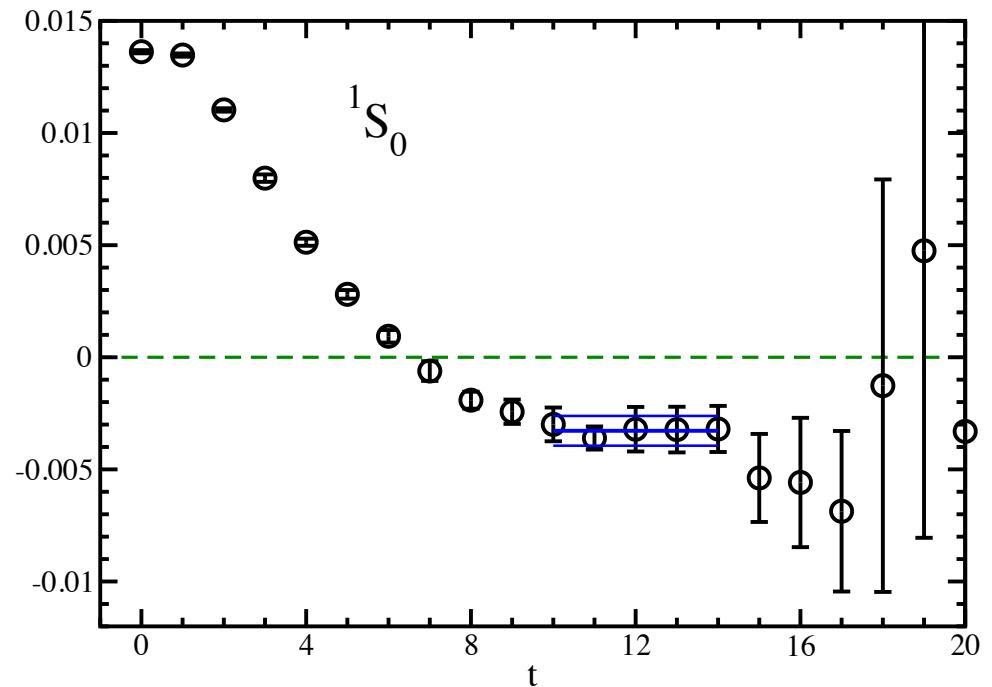
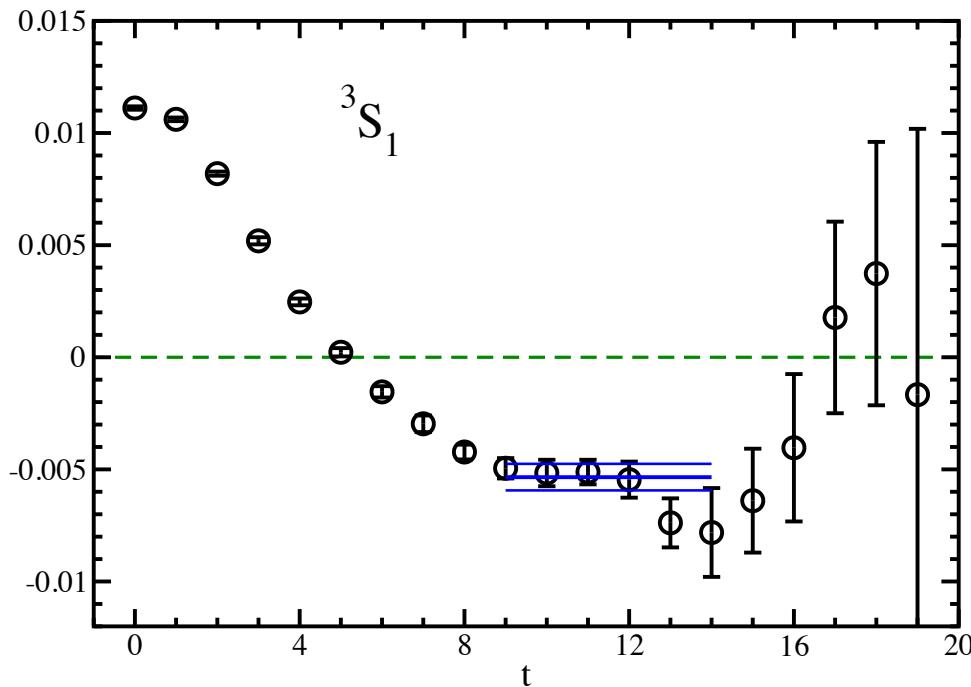


Observe bound state in both channels

ΔE_L in 2-nucleon channels at $m_\pi = 0.5$ GeV on $L = 5.8$ fm

TY et al., PRD86:074514(2012)

$$\Delta E_L = \log \left(\frac{R_{\text{NN}}(t)}{R_{\text{NN}}(t+1)} \right) \text{ with } R_{\text{NN}}(t) = \frac{C_{\text{NN}}(t)}{(C_N(t))^2}$$

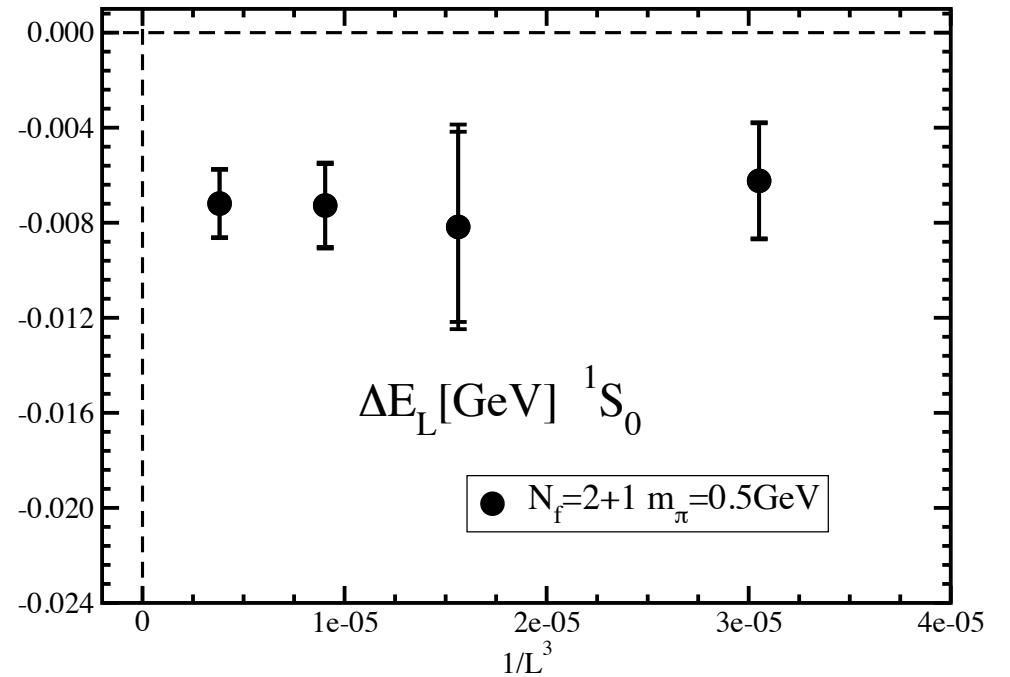
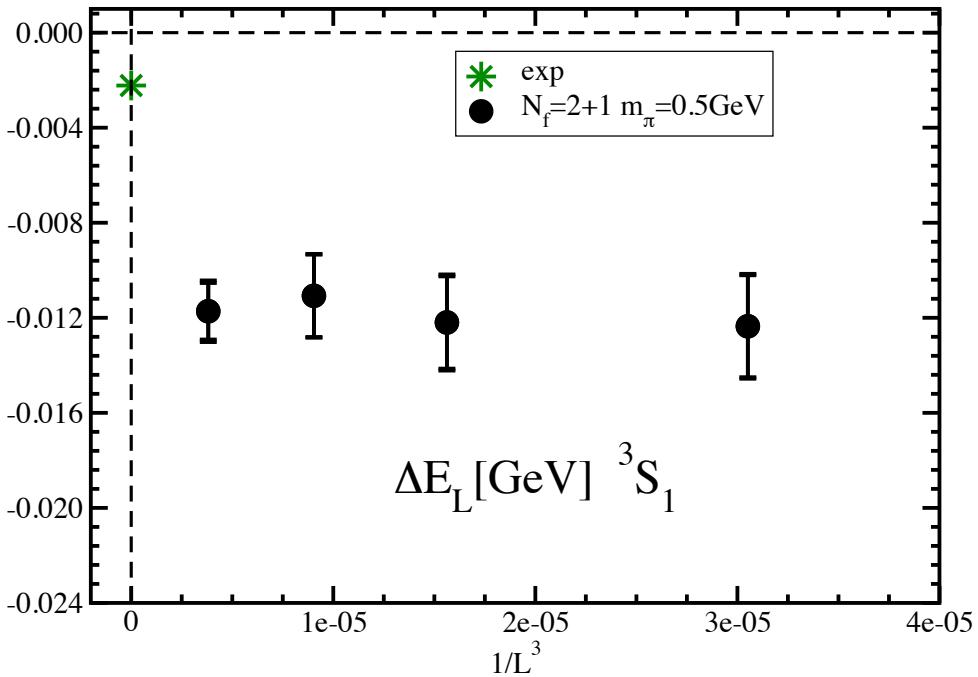


- Statistical error under control in $t \leq 12$
- Smaller error than ${}^4\text{He}$ and ${}^3\text{He}$ channels
- Negative ΔE_L in both channels

NN (3S_1 and 1S_0) channels $\Delta E_L = E_0 - 2m_N$ at $m_\pi = 0.5$ GeV

TY *et al.*, PRD86:074514(2012)

Identification of bound state from volume dependence of ΔE



- Negative ΔE_L
- Infinite volume extrapolation of ΔE_L

'04 Beane *et al.*, '06 Sasaki & TY

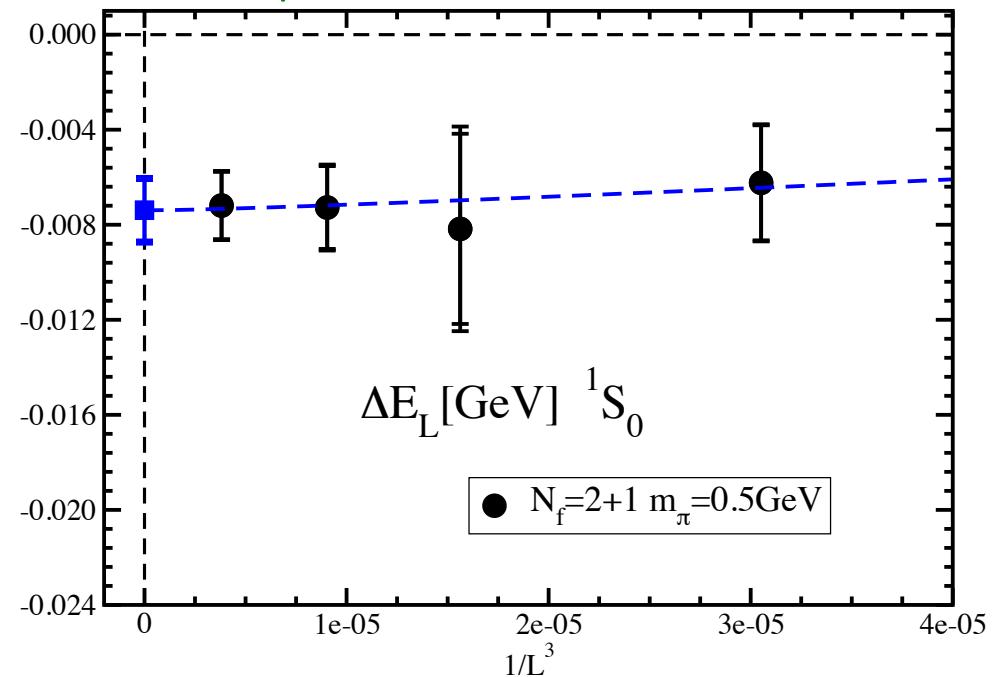
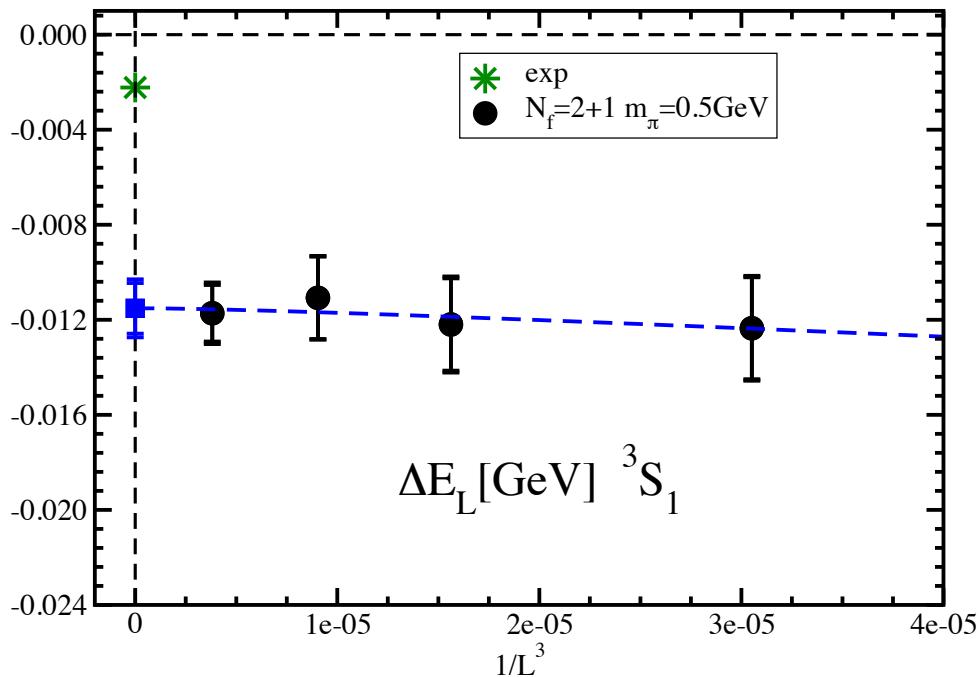
$$\Delta E_L = -\frac{\gamma^2}{m_N} \left\{ 1 + \frac{C_\gamma}{\gamma L} \sum'_{\vec{n}} \frac{\exp(-\gamma L \sqrt{\vec{n}^2})}{\sqrt{\vec{n}^2}} \right\}, \quad \Delta E_{\text{bind}} = \frac{\gamma^2}{m_N}$$

based on Lüscher's finite volume formula

NN (3S_1 and 1S_0) channels $\Delta E_L = E_0 - 2m_N$ at $m_\pi = 0.5$ GeV

TY et al., PRD86:074514(2012)

Identification of bound state from volume dependence of ΔE



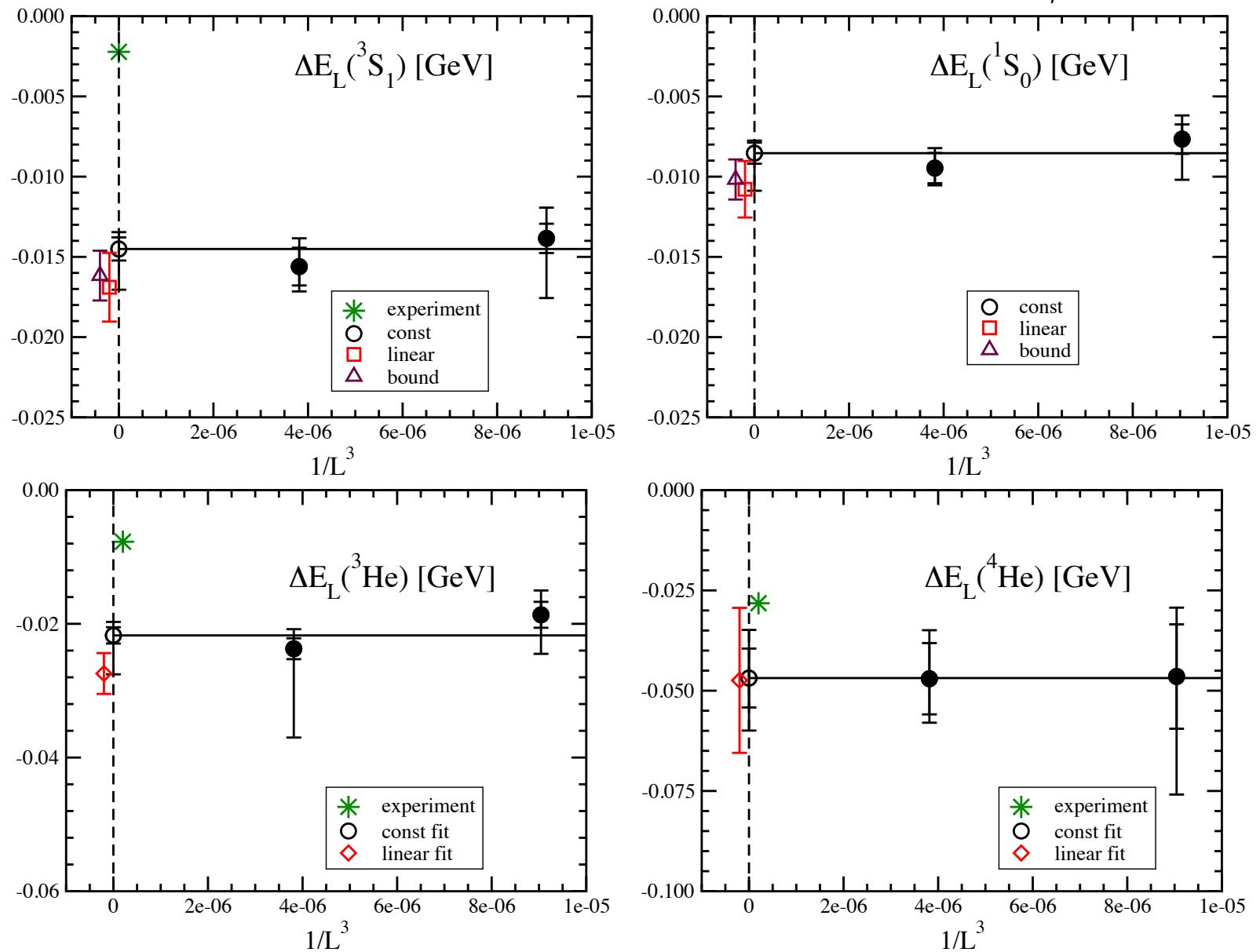
Bound state in both channels \leftarrow different from experiment

$$\Delta E_{3S_1} = 11.5(1.1)(0.6) \text{ MeV}$$

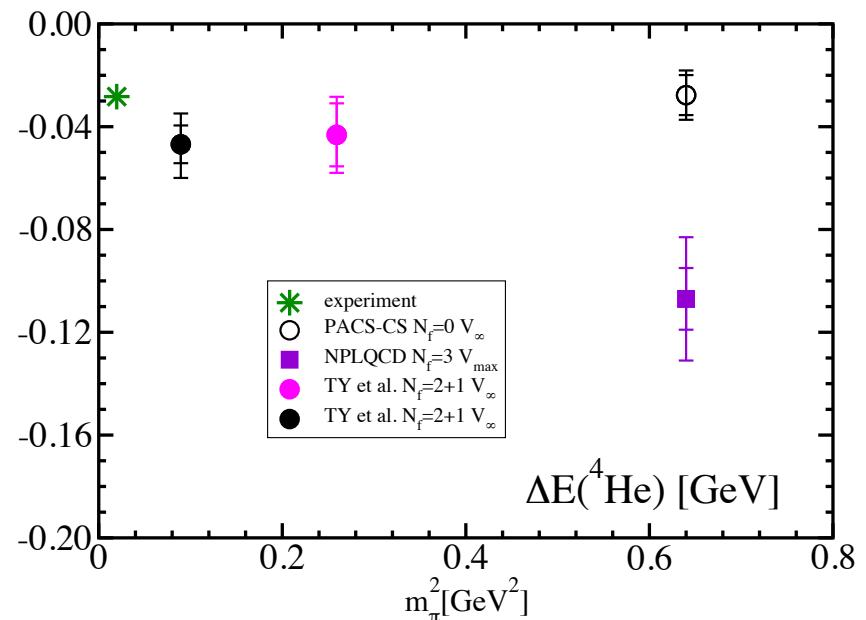
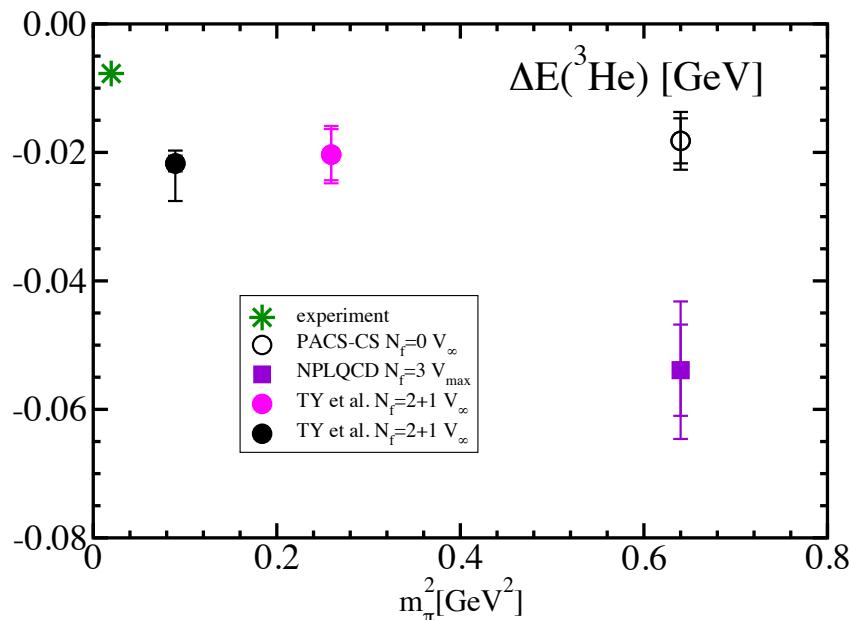
$$\Delta E_{1S_0} = 7.4(1.3)(0.6) \text{ MeV}$$

Results at $m_\pi = 0.3$ GeV with two volumes

TY et al., arXiv:1502.04182



Comparison of ${}^3\text{He}$ and ${}^4\text{He}$ nuclei

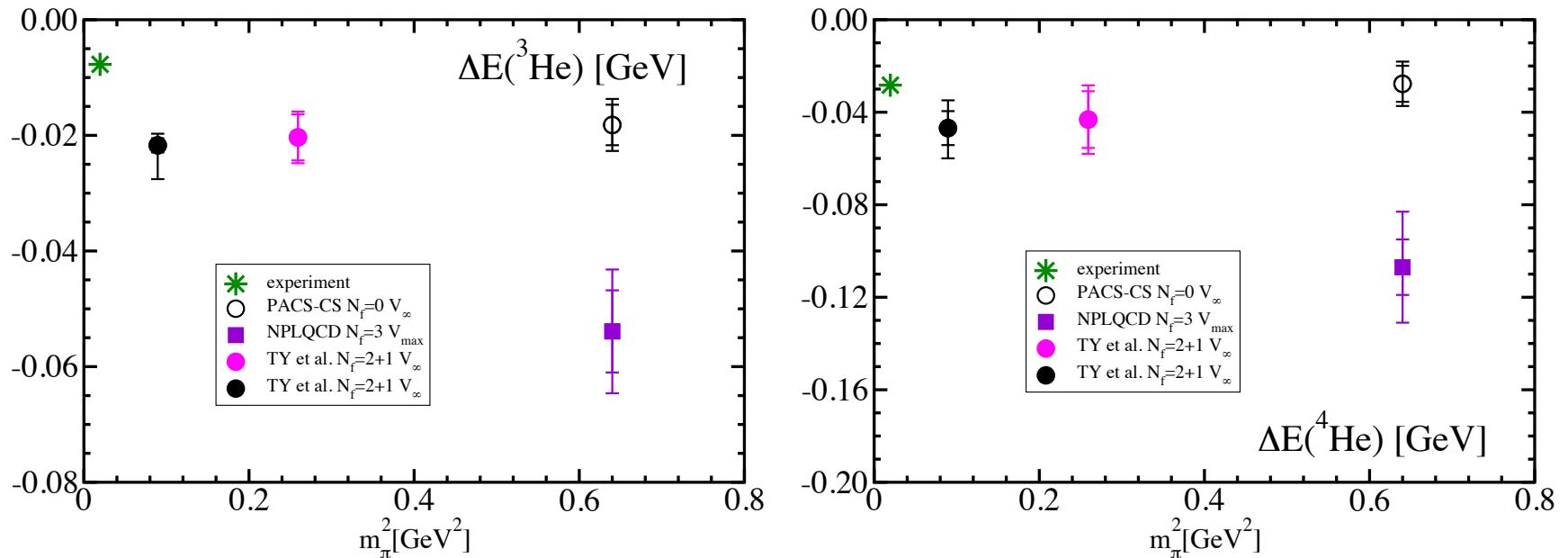


$L^3 \rightarrow \infty$ results only

Light nuclei likely formed in $0.3 \text{ GeV} \leq m_\pi \leq 0.8 \text{ GeV}$

Same order of ΔE to experiments

Comparison of ${}^3\text{He}$ and ${}^4\text{He}$ nuclei



$L^3 \rightarrow \infty$ results only

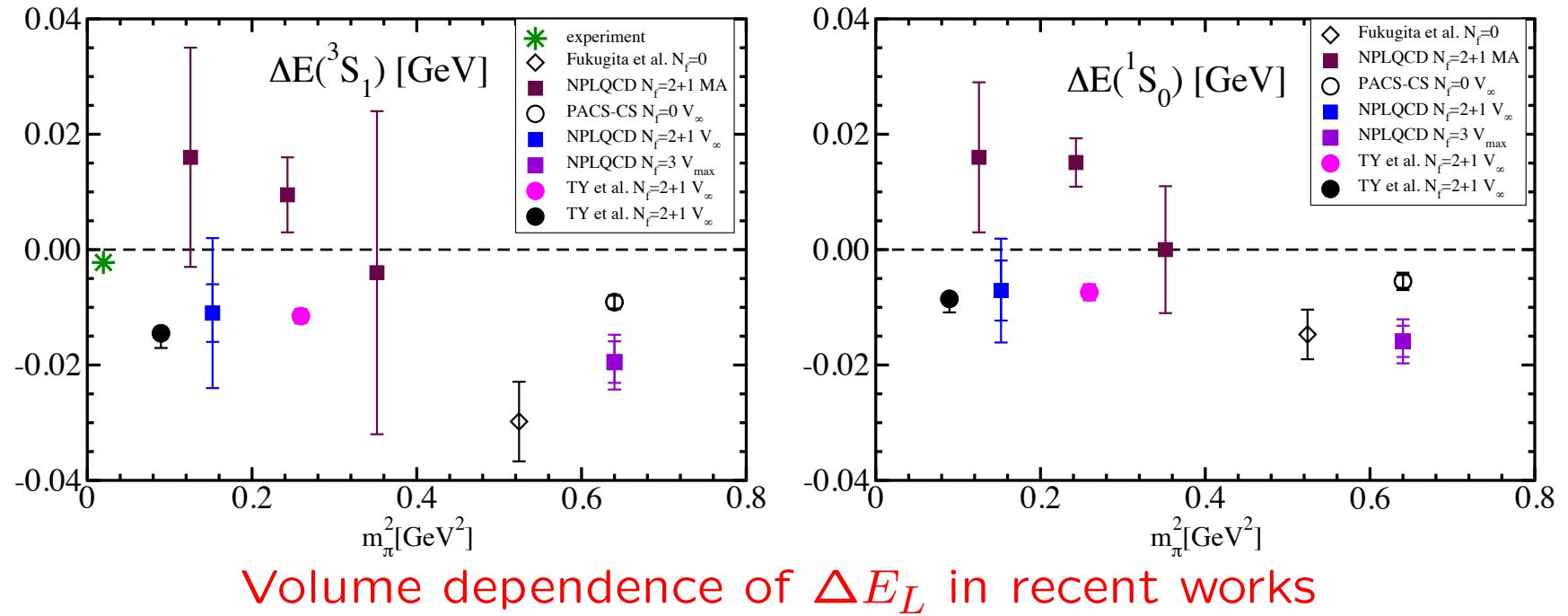
Light nuclei likely formed in $0.3 \text{ GeV} \leq m_\pi \leq 0.8 \text{ GeV}$

Same order of ΔE to experiments \rightarrow relatively easier than NN
 large $|\Delta E|$ makes less V dependence at physical m_π

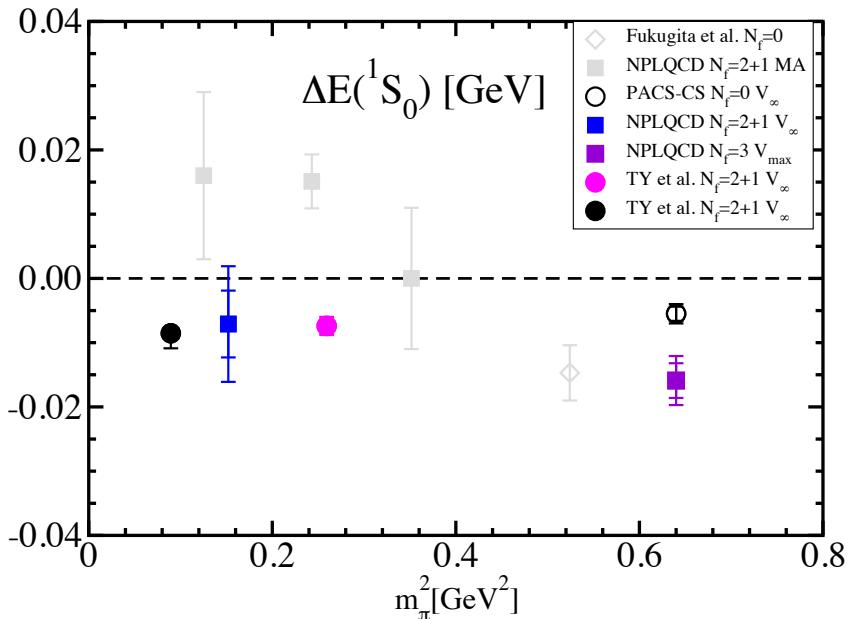
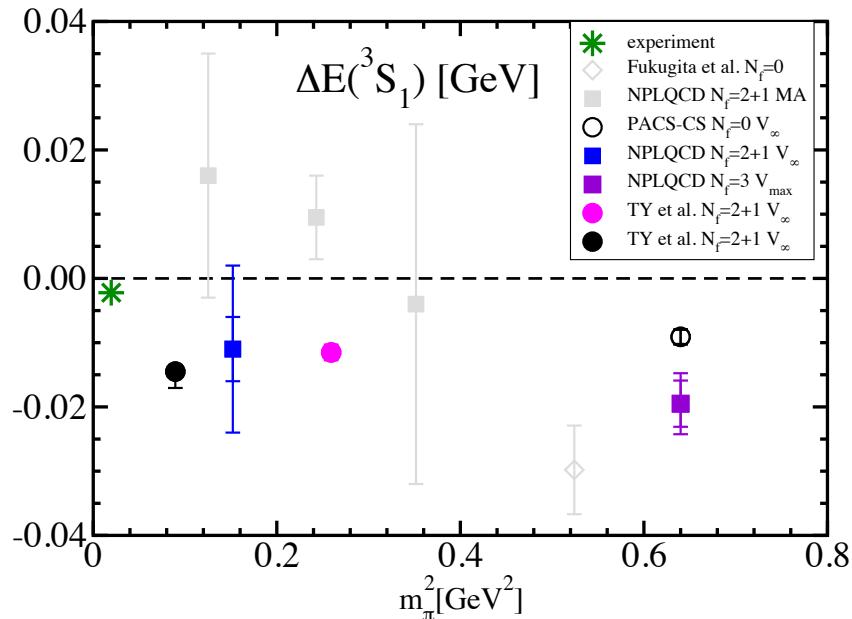
touchstone of quantitative understanding of nuclei from lattice QCD

Investigations of m_π dependence $\rightarrow m_\pi \sim 0.14 \text{ GeV}$ on $L \sim 8 \text{ fm}$

Comparison of NN channels



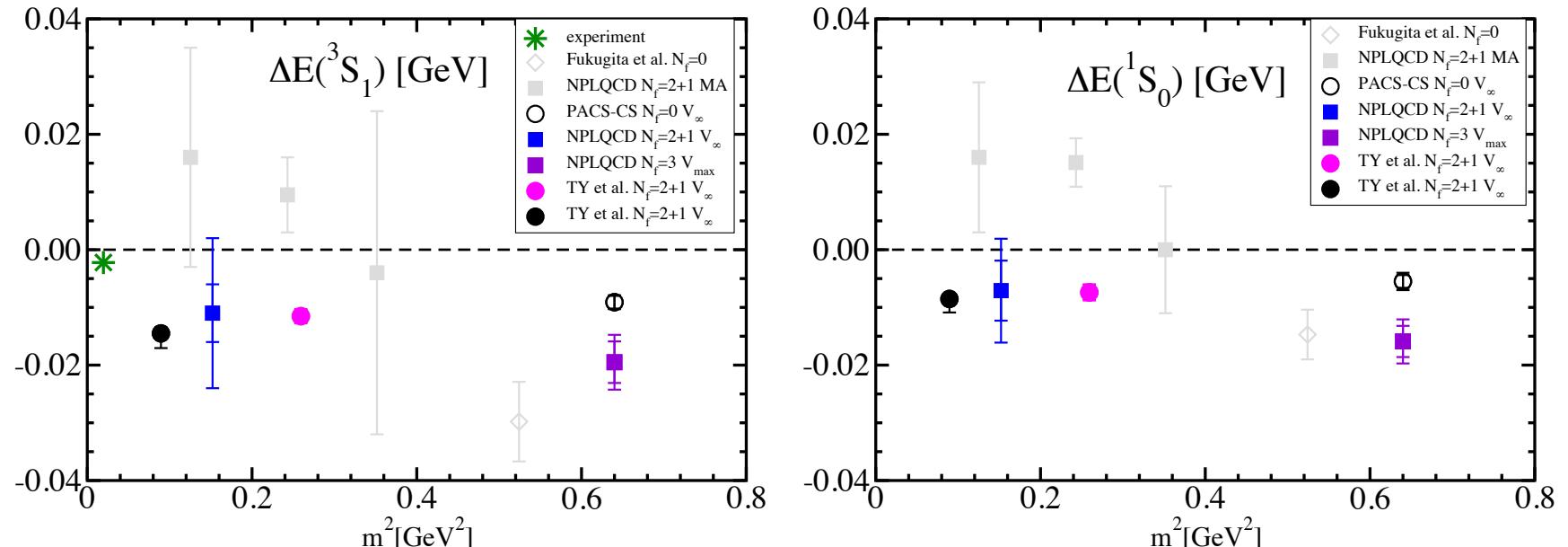
Comparison of NN channels



$L^3 \rightarrow \infty$: existence of bound states in 3S_1 and 1S_0
inconsistent with experiment due to larger m_π (?)

Investigations of m_π dependence $\rightarrow m_\pi \sim 0.14$ GeV on $L \sim 8$ fm

Comparison of NN channels with preliminary results



$L^3 \rightarrow \infty$: existence of bound states in 3S_1 and 1S_0
inconsistent with experiment due to larger m_π (?)

Investigations of m_π dependence $\rightarrow m_\pi \sim 0.14$ GeV on $L \sim 8$ fm

Large finite volume effect expected even on $L \sim 8$ fm '86 Lüscher, '04 Beane

$$^3S_1: \Delta E_{\text{exp}} = 2.2 \text{ MeV}$$

$$\Delta E_L = -(\Delta E_{\text{exp}} + \mathcal{O}(\exp(-L\sqrt{m_N \Delta E_{\text{exp}}})) \sim -4 \text{ MeV}$$

$$^1S_0: a_0^{\text{exp}} = 23.7 \text{ fm}$$

$$\Delta E_L = -\frac{4\pi a_0^{\text{exp}}}{m_N L^3} + \mathcal{O}(1/L^4) \sim -2 \text{ MeV}$$

Summary

$N_f = 2 + 1$ lattice QCD at $m_\pi = 0.5$ and 0.3 GeV

- Volume dependence of ΔE

$\Delta E \neq 0$ of 0th state in infinite volume limit

→ bound state in ${}^4\text{He}$, ${}^3\text{He}$, 3S_1 and 1S_0
at $m_\pi = 0.5$ and 0.3 GeV

- ΔE larger than experiment and small m_π dependence
- Bound state in 1S_0 not observed in experiment
Deep bound state in $N_f = 3$ at $m_\pi = 0.8$ GeV ('12 NPLQCD)
- No bound state in HALQCD method

Need further investigations

e.g. systematic error from large m_π and finite lattice spacing

$N_f = 2 + 1$ $m_\pi \sim 0.14$ GeV on $L \sim 8$ fm calculation is ongoing.

Very preliminary results of

$$\Delta E = E_0 - N_N m_N$$

at $m_\pi \sim 0.14$ GeV on $L \sim 8$ fm

