

# Axial symmetry at finite temperature and Dirac operator eigenmodes

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YITP HHIQCD2015, Kyoto  
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# Outline

Literature

Finite  
temperature  
Axial symm.  
**Introduction**

Methods

&

Results

**Work in progress**

Final  
thoughts

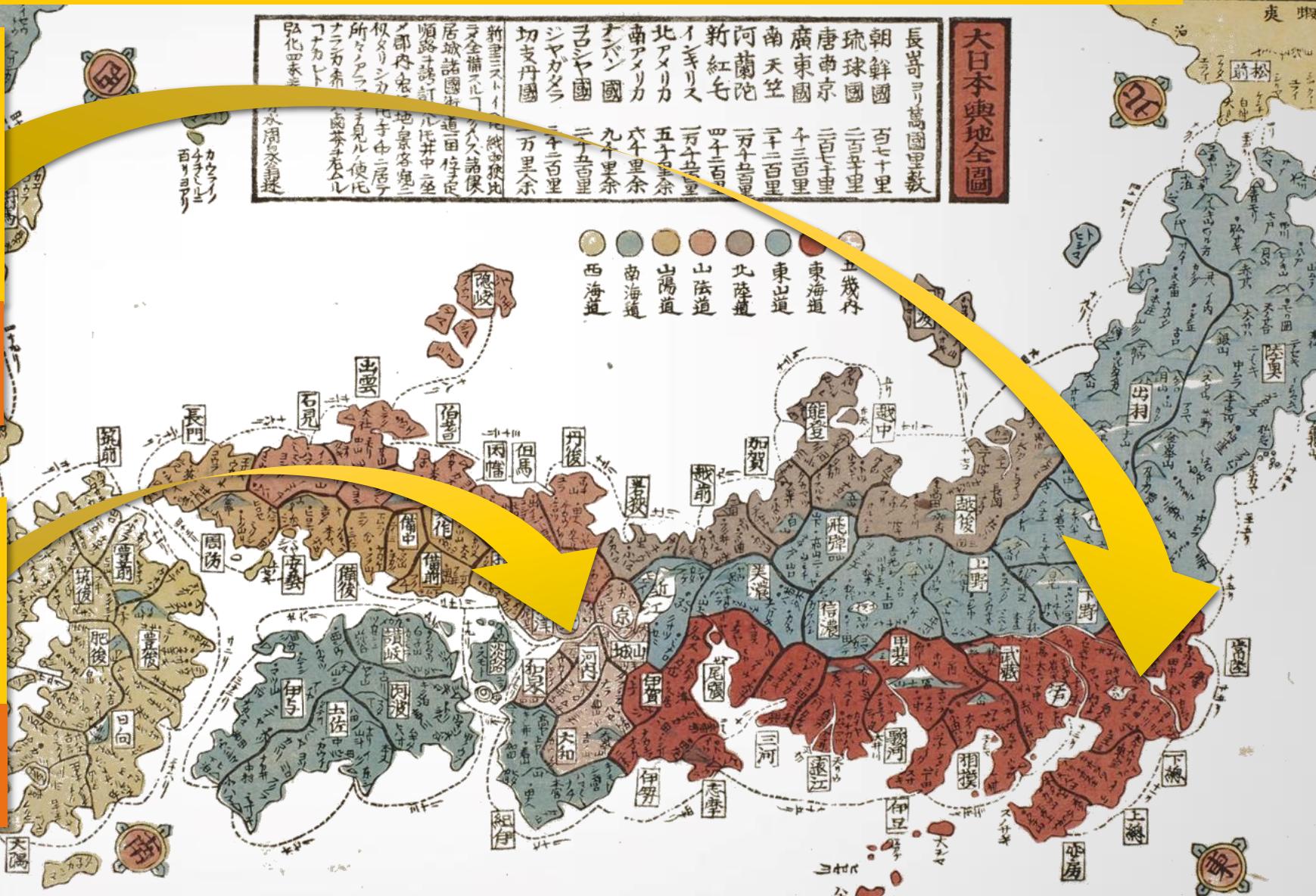
# JLQCD collaborators

# Shoji Hashimoto Jun-ichi Noaki Takashi Kaneko

**KEK**

# Hidenori Fukaya Akio Tomiya

# Osaka Un.



Finite  
temperature  
Axial symm.  
**Introduction**

Literature

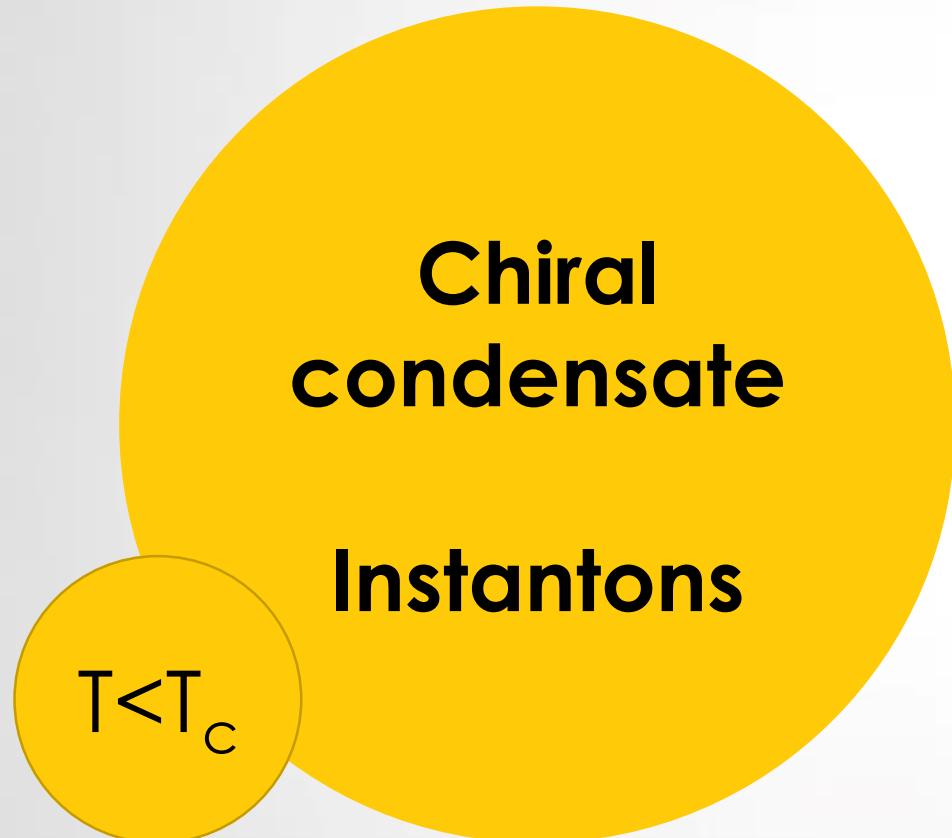
Methods  
&  
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# Low temperature – symmetries

$$SU(2)_L \times SU(2)_R \times U(1)_V \times U(1)_A$$



**Dirac operator eigenmodes**

Near zero modes density

$$\Sigma = \pi \rho(0)$$

Zero modes

$$\int \partial_\mu J_{\mu 5} \propto Q$$

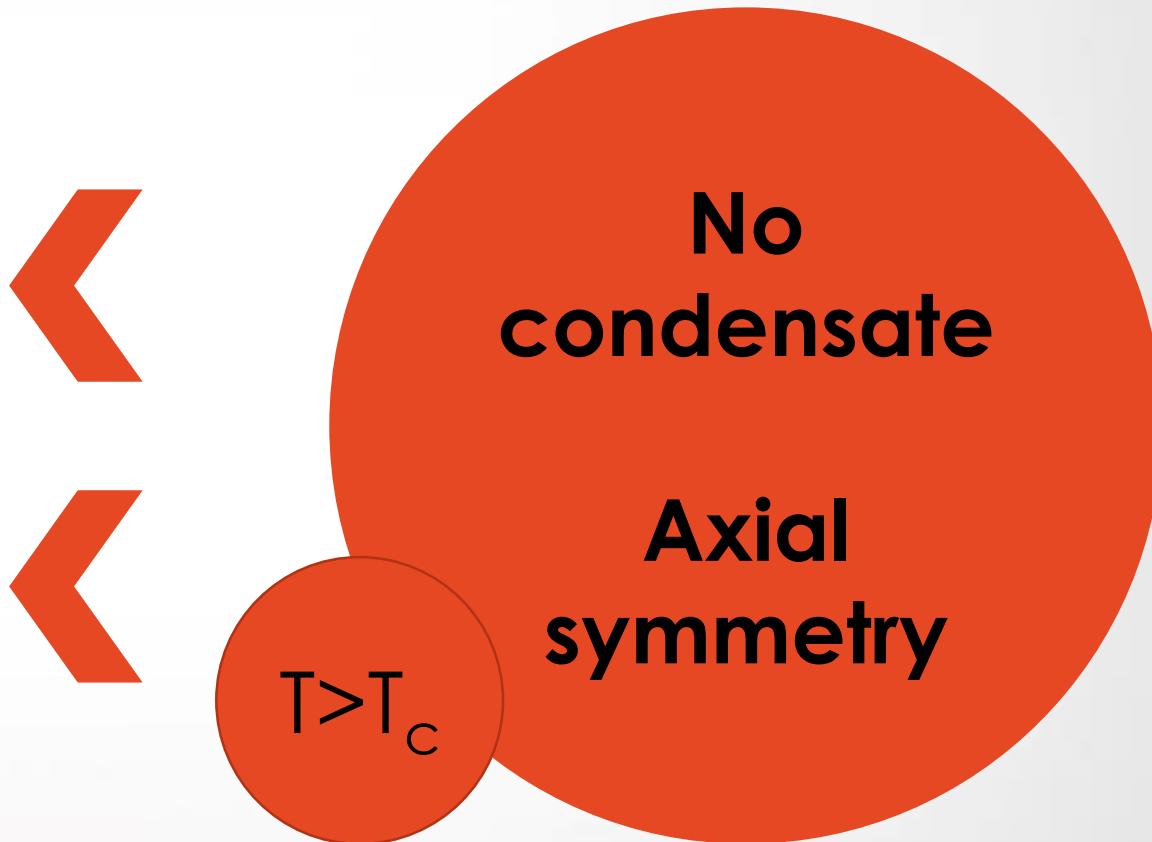
# High temperature – symmetries

$$SU(2)_L \times SU(2)_R \times U(1)_V \times U(1)_A$$

**Current knowledge**

**Restoration** of chiral symmetry at  $T_c$

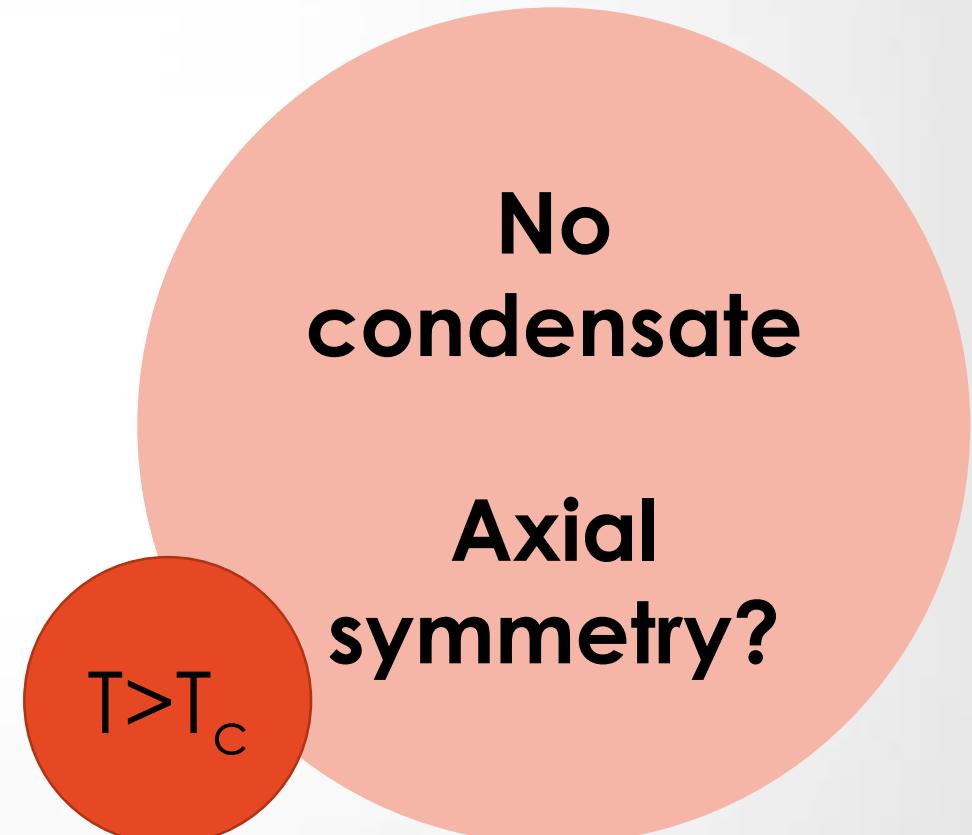
**Restoration** at  $T \rightarrow \infty$



# High temperature – **symmetries**

$$SU(2)_L \times SU(2)_R \times U(1)_V \times U(1)_A$$

$T \gtrsim T_c ?$



# Methods & Results

**Work in progress**

Finite  
temperature  
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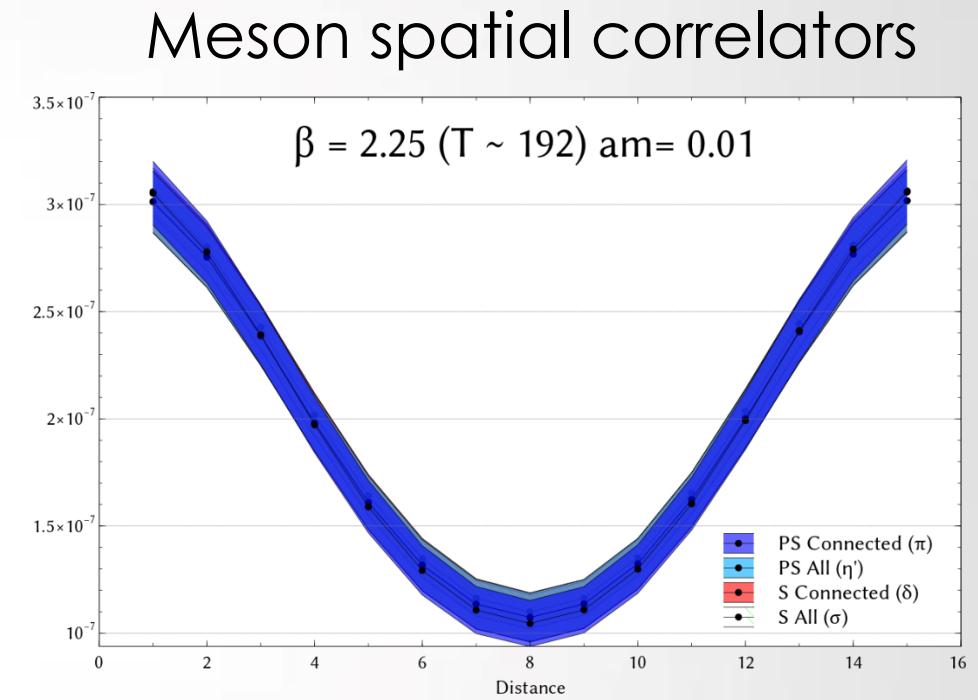
Final  
thoughts

# Recent literature - I

G. Cossu et al. (2013) for JLQCD  
Disconnected meson diagrams  
**vanish** at temperatures above  $T_c$

Related: **Gap** in the Dirac spectrum

Aoki, Fukaya, Taniguchi (2012)  
Analytic calculation (Overlap)  
Dirac spectrum  $\rho(\lambda) \sim c\lambda^3$   
Implies **U(1)<sub>A</sub> anomaly invisible**



$$\pi = \delta = \rho = \sigma$$

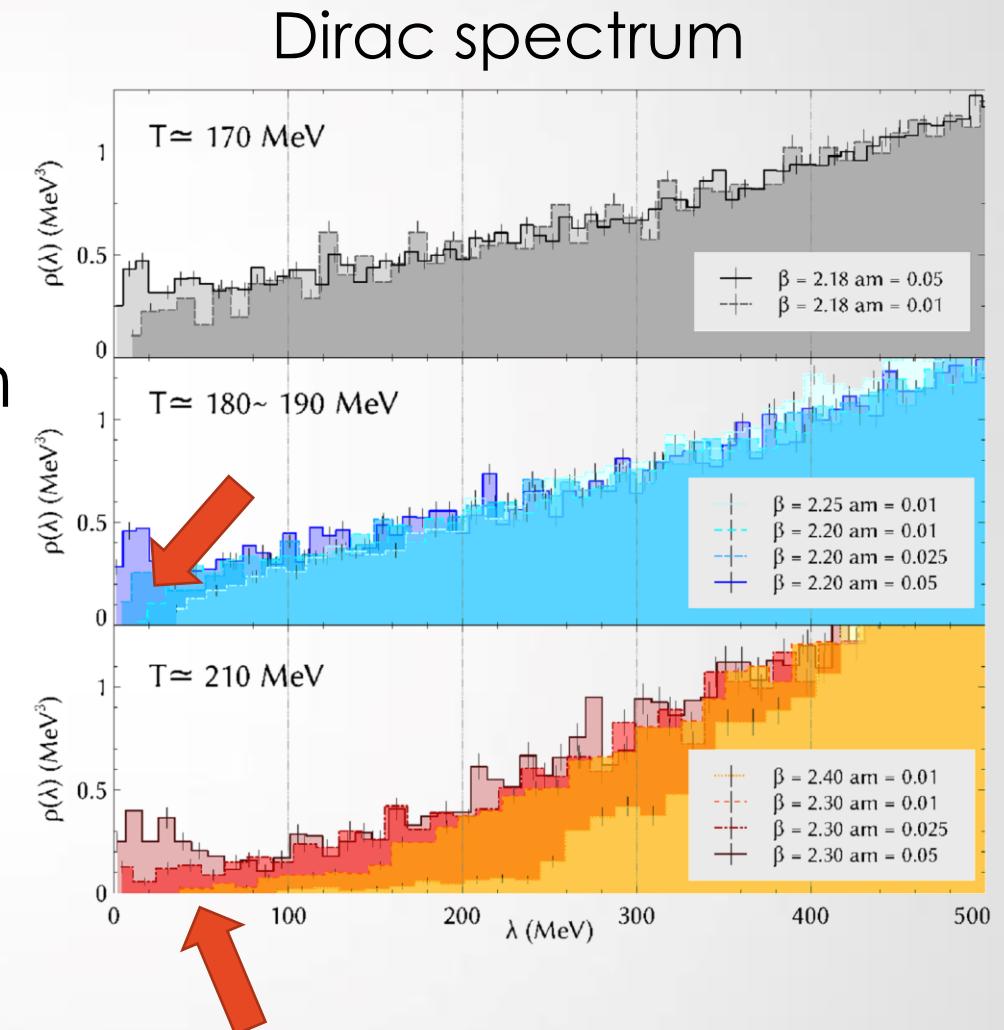
**Restored**

# Recent literature - II

G. Cossu et al. (2013) for JLQCD  
Disconnected meson diagrams  
**vanish** at temperatures above  $T_c$

Related: **Gap** in the Dirac spectrum

Aoki, Fukaya, Taniguchi (2012)  
Analytic calculation (Overlap)  
Dirac spectrum  $\rho(\lambda) \sim c\lambda^3$   
Implies **U(1)<sub>A</sub> anomaly invisible**



# Recent literature - III

Bazavov *et al.* (2012-13)

Domain wall, several volumes

Dirac spectrum, susceptibilities

**NOT restored**

Ohno *et al.*, Sharma *et al.* (2012-15)

Overlap on HISQ configurations

Dirac spectrum

**NOT restored**

Brandt *et al.* (2013)

Wilson improved fermions

Screening masses

**NOT restored**

Our previous study  
**Exact** chiral symmetry (Overlap)  
**topology fixed**  
**Only**  $16^3 \times 8$  volume  
**Mass dependence**  
**No** continuum limit

# Methods & Results

**Work in progress**

Finite  
temperature  
Axial symm.  
**Introduction**

Literature

Final  
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# Chiral symmetry on the lattice

$$\{D, \gamma_5\} = 0$$

Nielsen-Ninomiya no-go theorem:  
chiral symmetry implies unwanted doublers

The Ginsparg-Wilson relation (1982)

$$\{D, \gamma_5\} = a D \gamma_5 D$$

# Generalized Domain Wall

$$D^4(m) = \frac{1+m}{2} + \frac{1-m}{2} \gamma_5 \text{sgn}(H)$$

Play with the **sign function**

Möbius Kernel

$$H_M = \gamma_5 \frac{b D_W}{2 + c D_W}$$

Function approximation

Transfer matrix in 5D

- Hyperbolic tangent
- Rational approximation

**Reduced residual mass**

**b=2 c=1 Scaled Shamir, m<sub>res</sub> ~ 10<sup>-4</sup>**

# Status of simulations



**Symanzik + smeared 2-flavors DWF**

**Multipurpose code**, HMC & measurements

Available on request, soon online

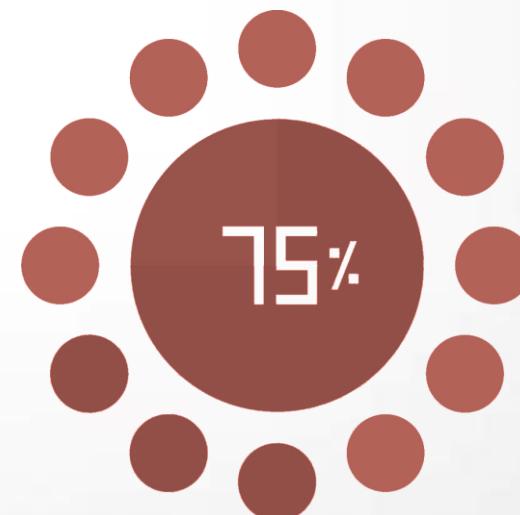
**Optimized for BlueGene/Q**

Webpage: [http://suchix.kek.jp/guido\\_cossu/](http://suchix.kek.jp/guido_cossu/)

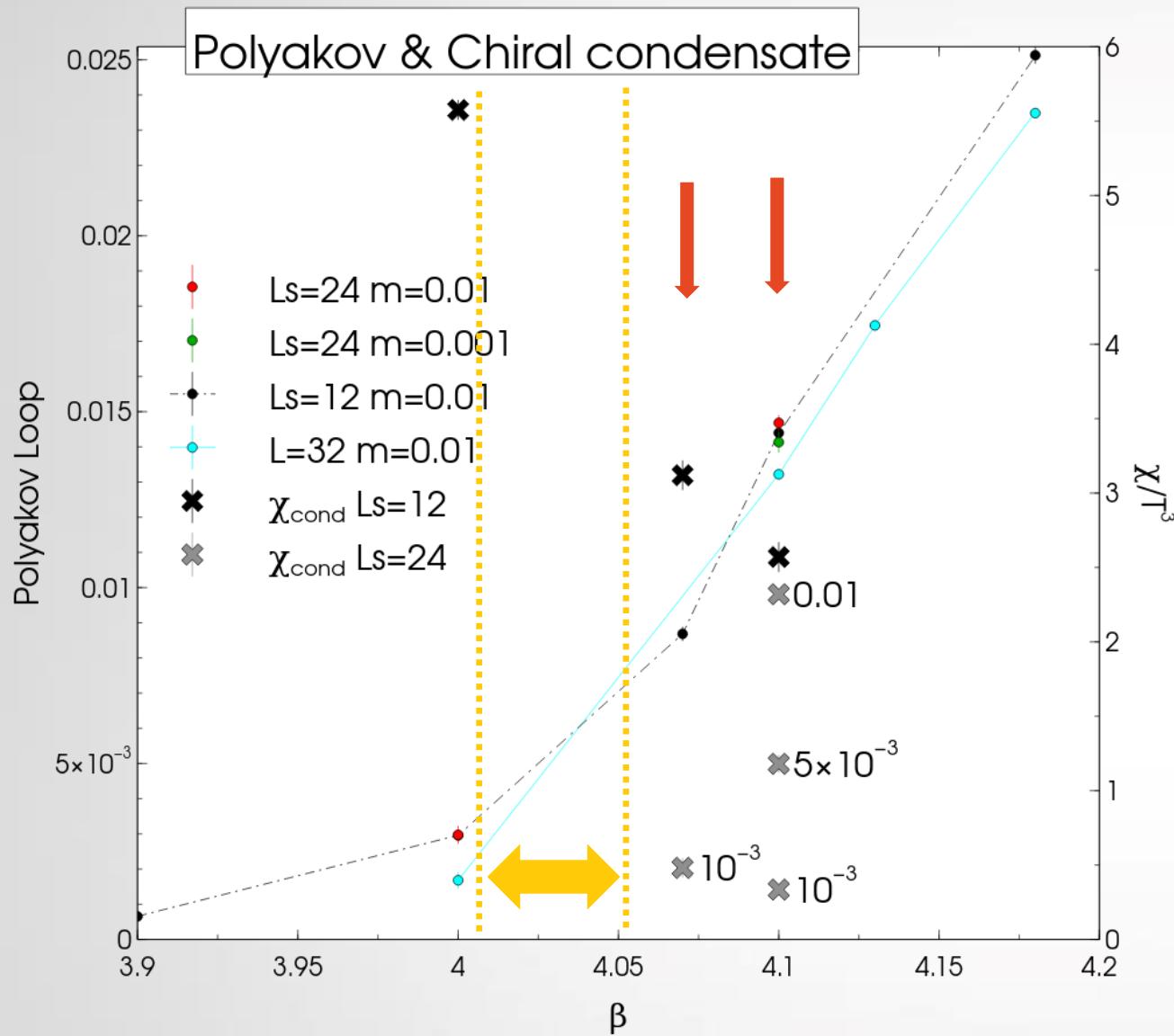
## Collected data

- 2 lattice volumes
- 3 masses
- 5 temperatures
- Topology changes
- $N_t=8, N_f=12$  (finer lattice)

**Full analysis in progress**



# Phase transition



Today:  
**T ~ 184, 200 MeV (red arrows)**

Phase transition at ~180 MeV

**2 volumes**  
**Mass dependence**

$N_f=12$  running now.  
Analysis almost done

$$\text{Delta} \quad \Delta = \chi_\pi - \chi_\delta \quad \chi_x = \int \langle X(0)X(x) \rangle$$

Theoretically clean: **zero** if axial symmetry is restored

**Veeery delicate measurement**

**Talk: breakdown of the signal sources (physics/artifacts)**

First: integrating correlators is **bad**, so

**Stochastic** measurements of Dirac traces

**Eigenmode** decomposition

$$\Delta = \int \frac{2m^2 \rho(\lambda, m)}{(\lambda^2 + m^2)^2}$$

# Source of the signal

Discrete spectral sum

**Zero modes**

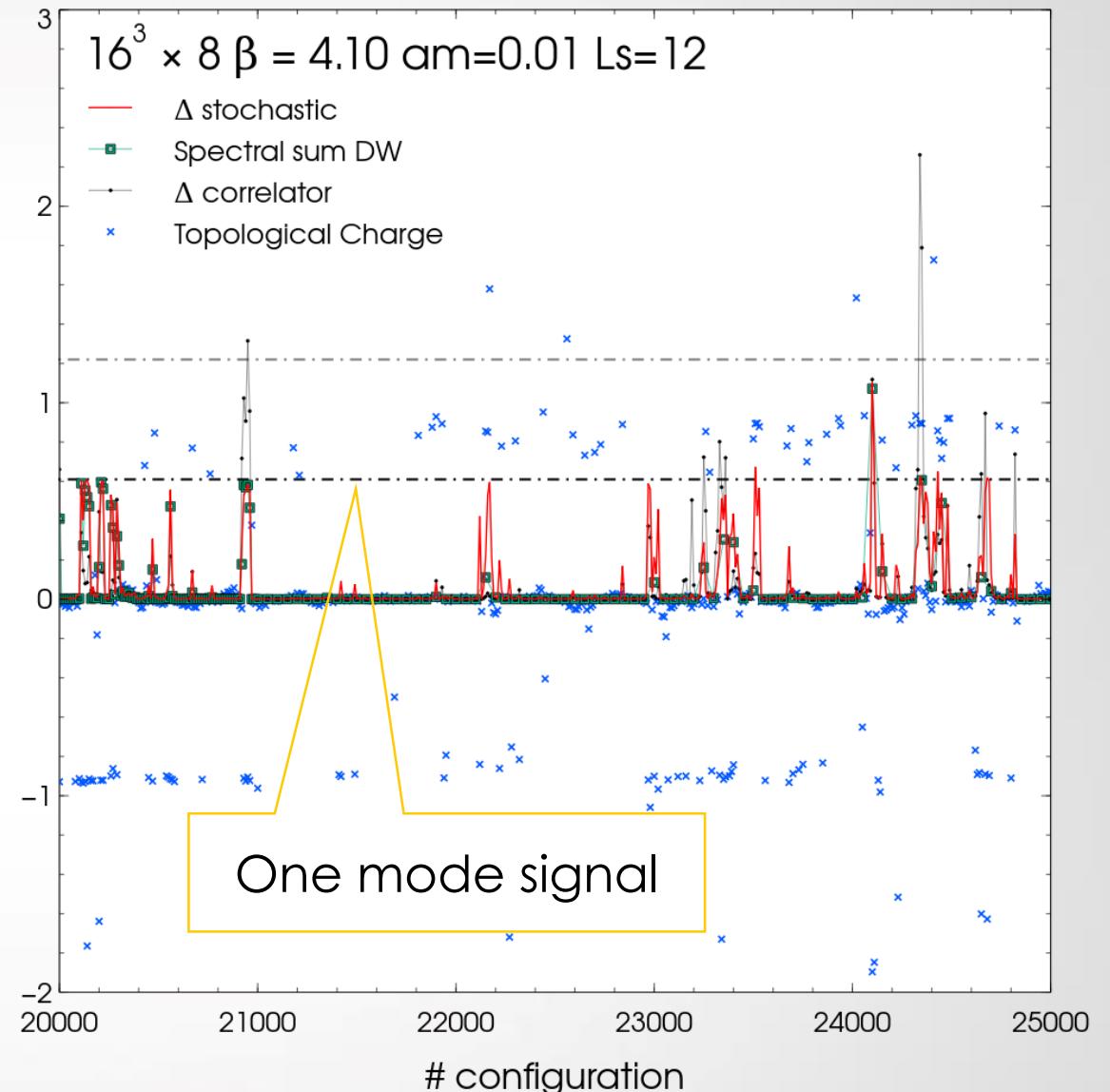
$$\Delta = \frac{2N_0}{Vm^2} + \sum_{\lambda \neq 0} \frac{2m^2}{V(\lambda^2 + m^2)^2}$$

**Bulk**

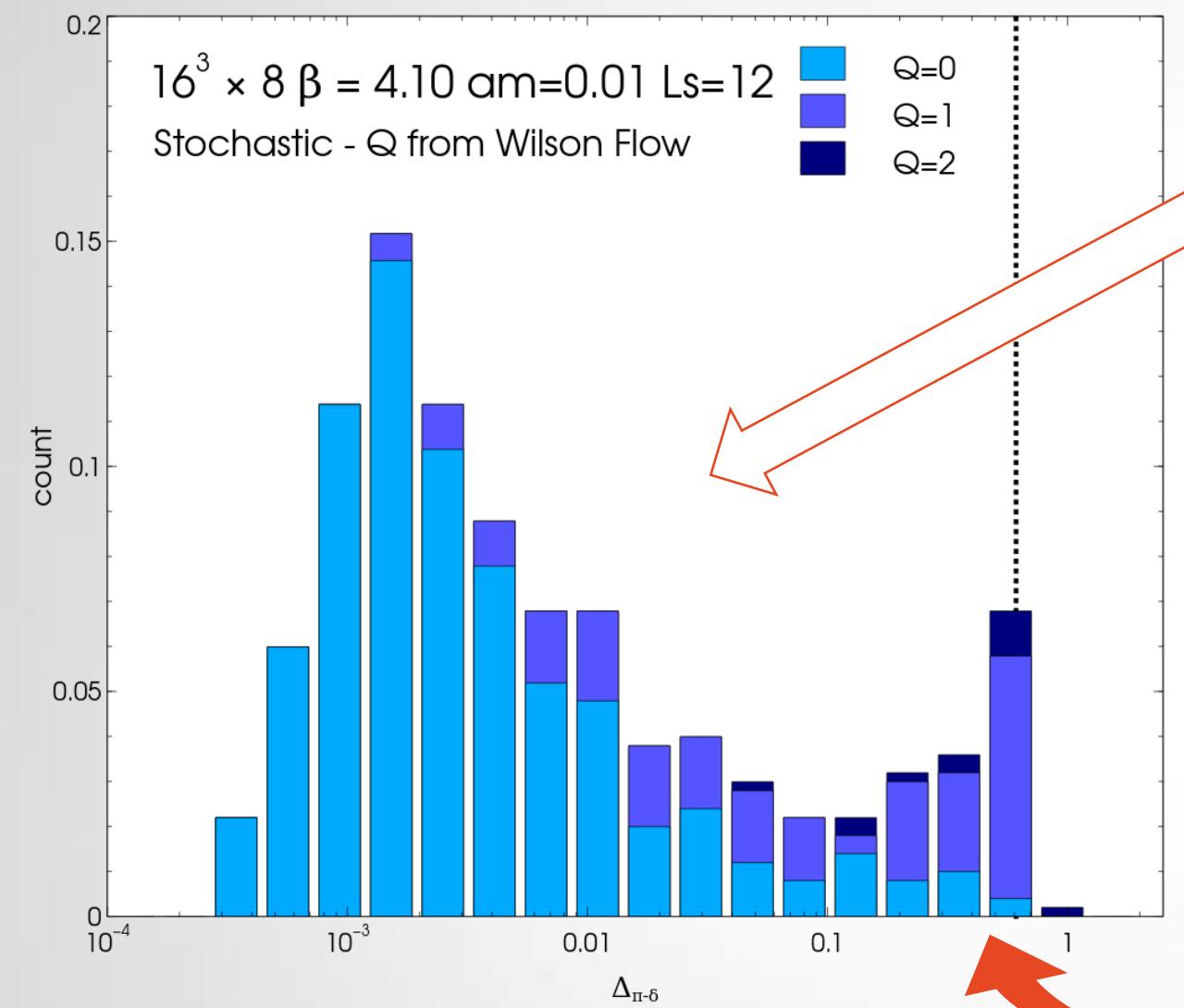
Peaks dominate the signal



Fluctuations of  
3 orders of magnitude



# $\Delta$ - Topology correlation



**Mild correlation**

**Tension** with spectral sum expectations

Two sources

- GW violations
- $F\tilde{F}$  estimate

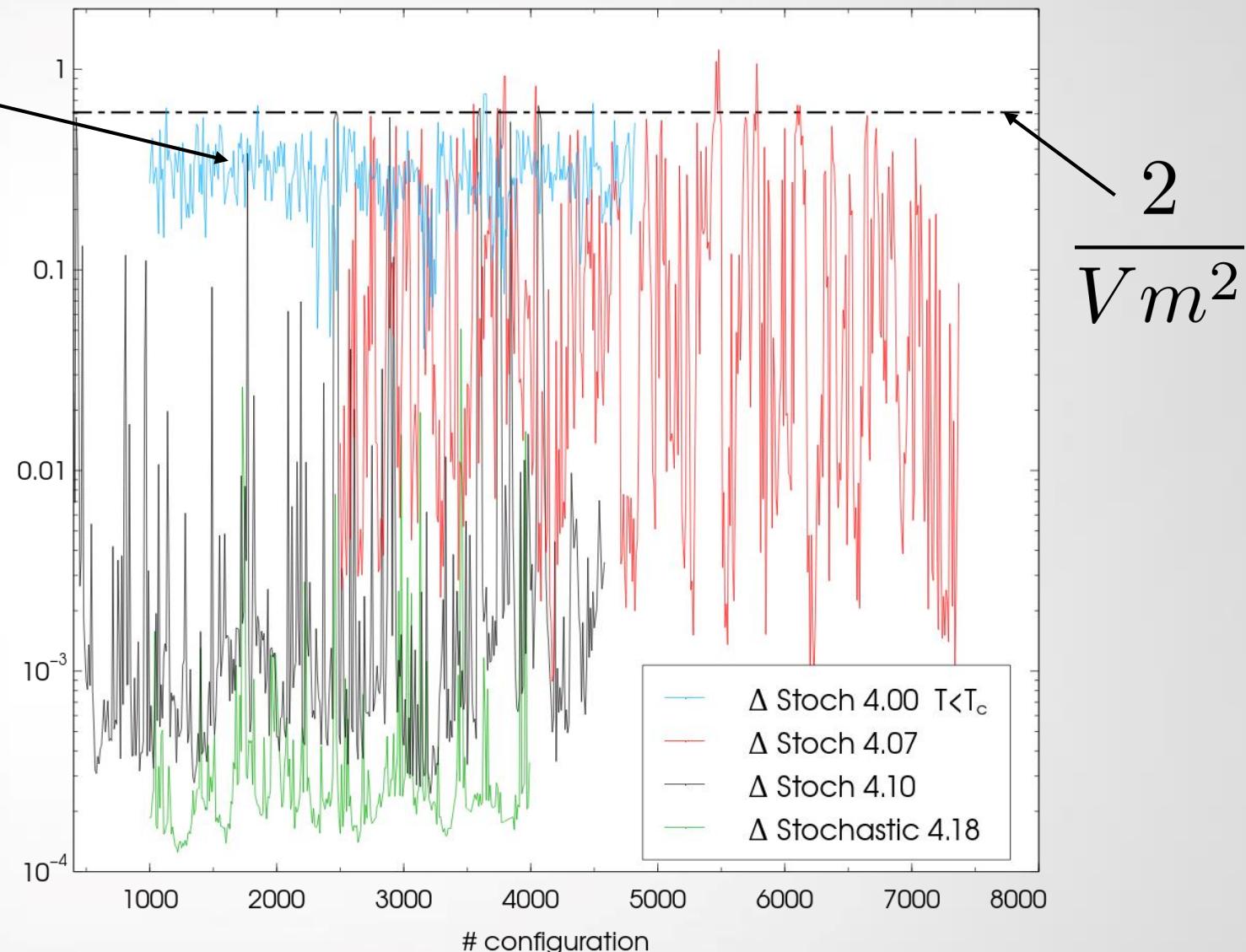
$Q=0$  near zero modes

# Temperature dependence

Broad picture arising  
at *this stage*:

- **Just above** the phase transition zero modes **dominate**
- **Then** they are **strongly suppressed** and the signal goes down

$$T < T_c$$



# Let's increase volume – $m=0.01$

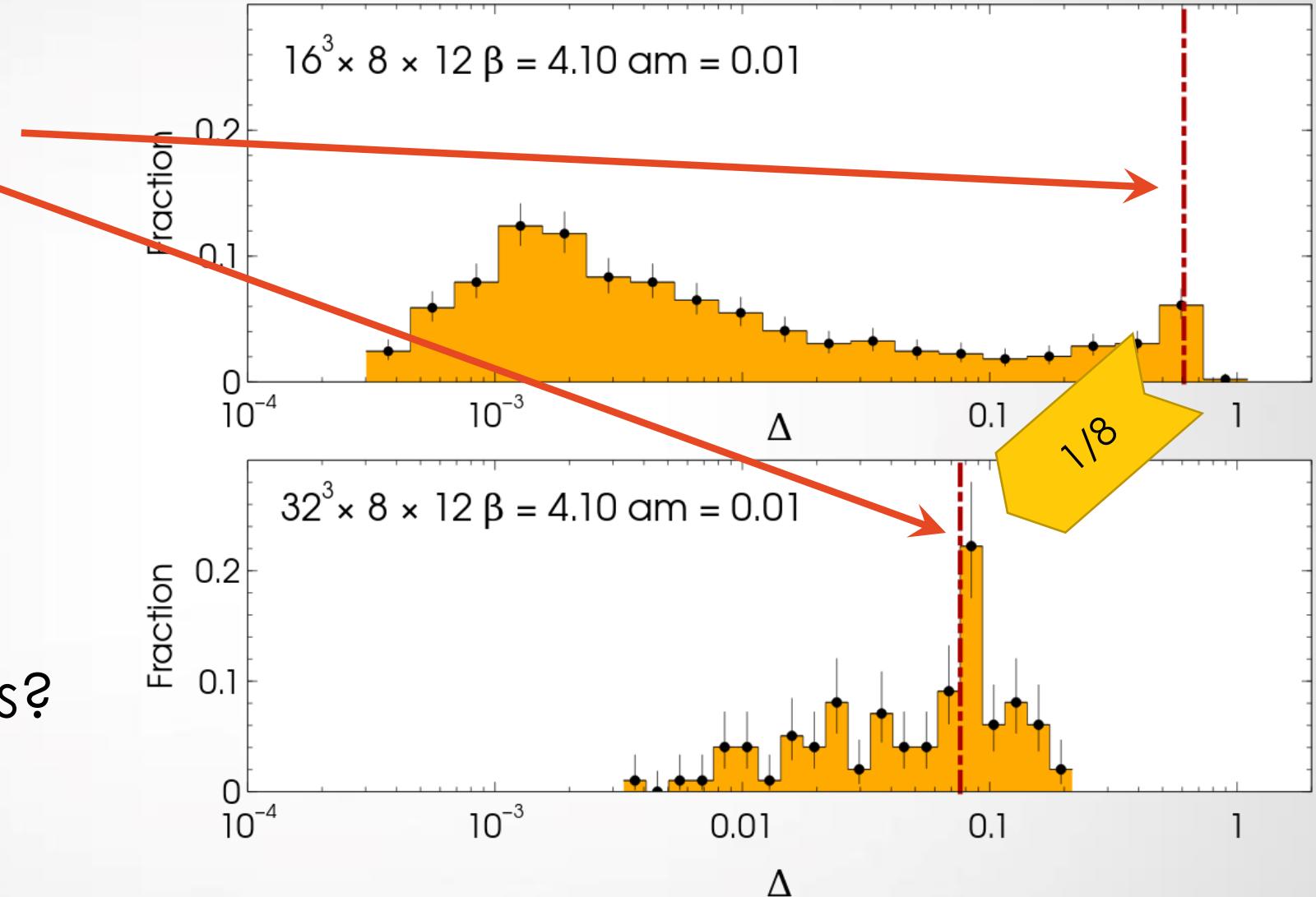
**Zero mode  
contribution**

**suppressed  $\sim 1/V$**

As expected from  
spectral sum

**Bulk contribution  
increases**

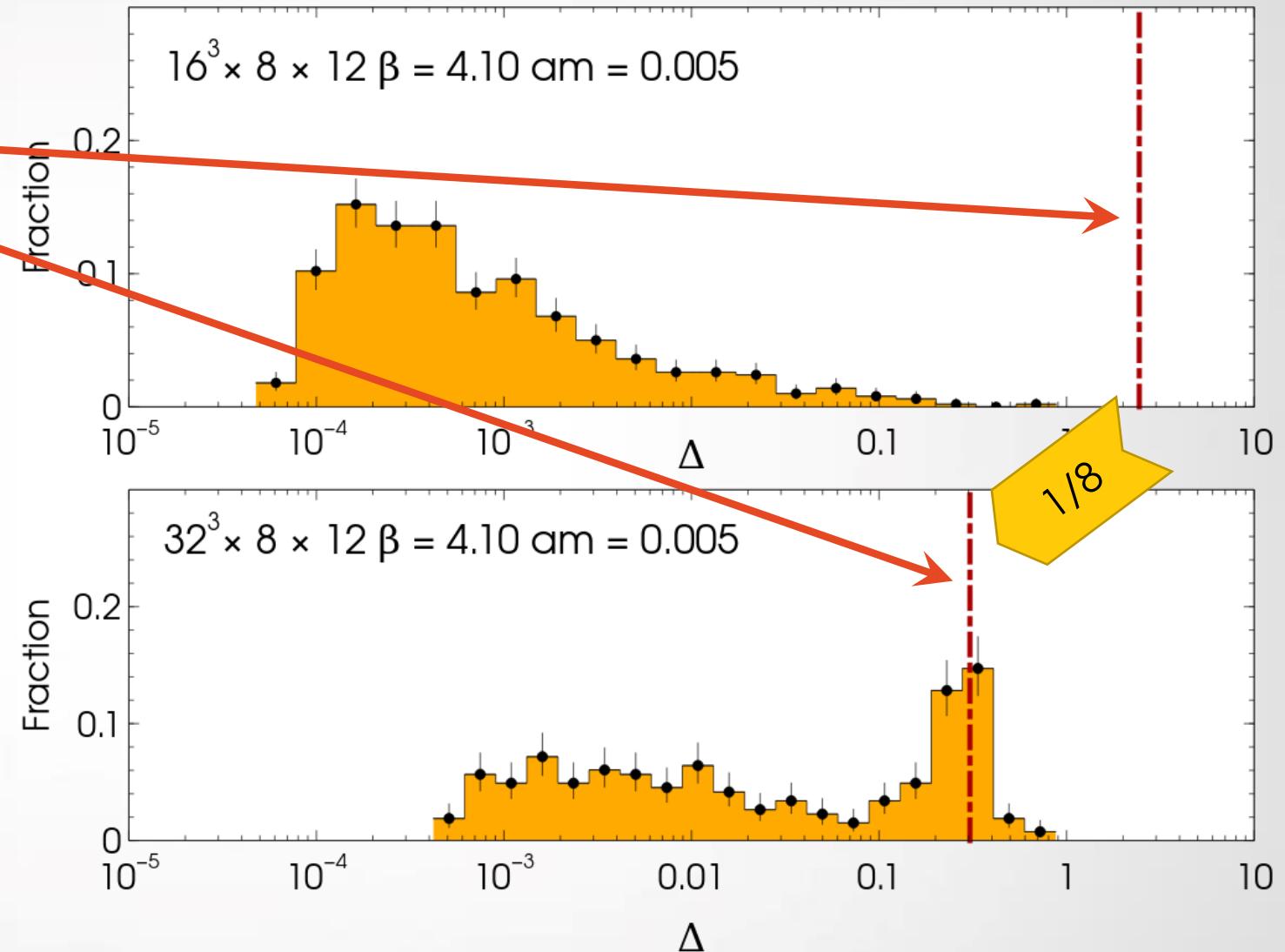
Decrease the mass?



# Let's increase volume – $m=0.005$

**Zero mode contribution suppressed  $\sim 1/V$**   
As expected from spectral sum

**Bulk contribution increases**

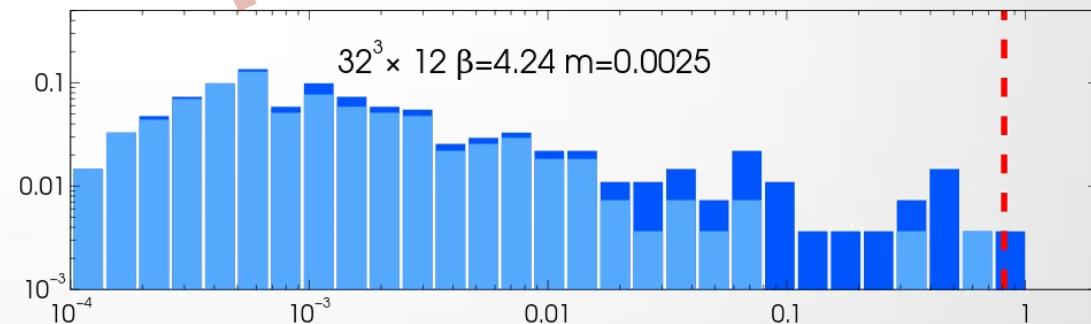
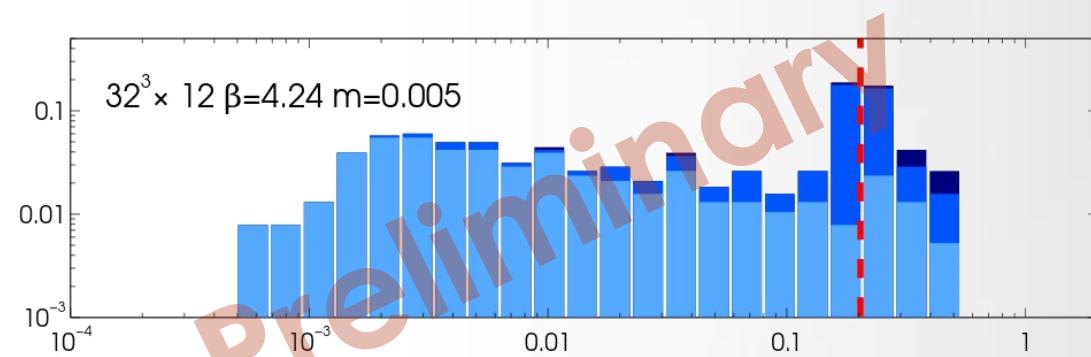
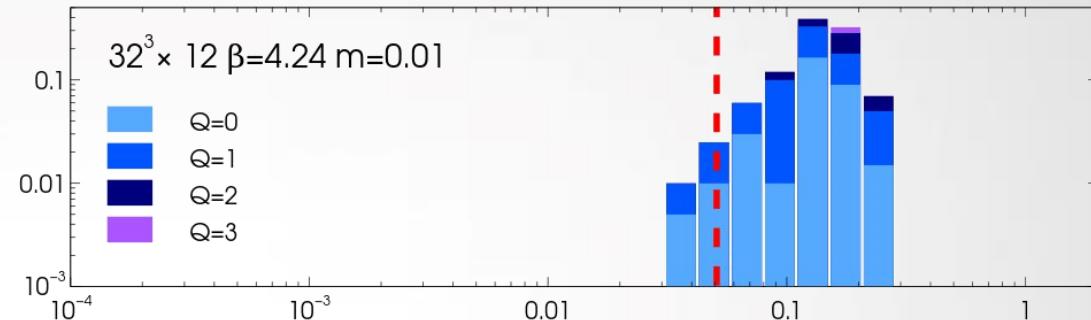


# Continuum limit

**N<sub>t</sub> from 8 to 12**

No big news, violations reduced

T ~ 200 MeV



# Volume&mass dependence

$$\chi_t = \lim_{V \rightarrow \infty} \frac{\langle Q^2 \rangle}{V} = \text{const.} \rightarrow \frac{N_0}{V} \rightarrow 0$$

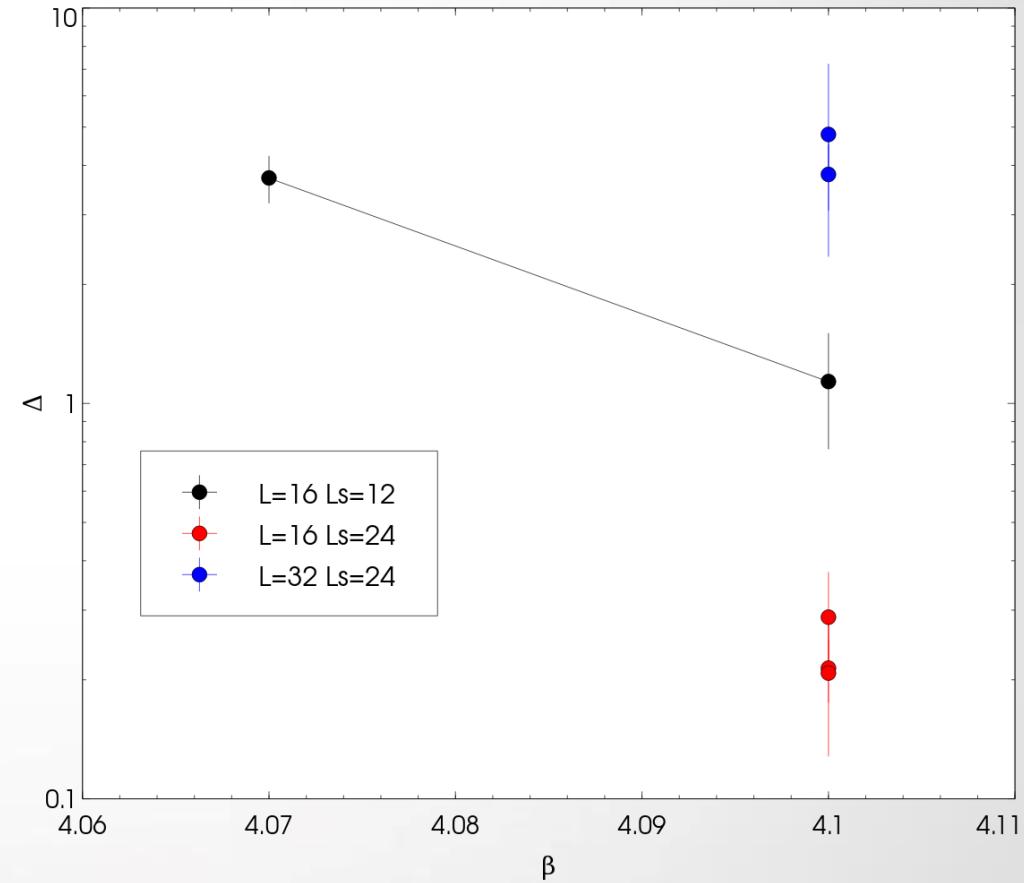
Zero modes  
contribution  
**vanishes**

Conclusion: **signal from the bulk part, near zero modes**

Let's cut all configurations with  $Q > 0$  (naïve cut)



Signal **constant with the mass**



# Is everything all right? – I

From the Ginsparg-Wilson relation we can measure the amount of violation **for each mode**,  $\textcolor{red}{g_{nn}}$

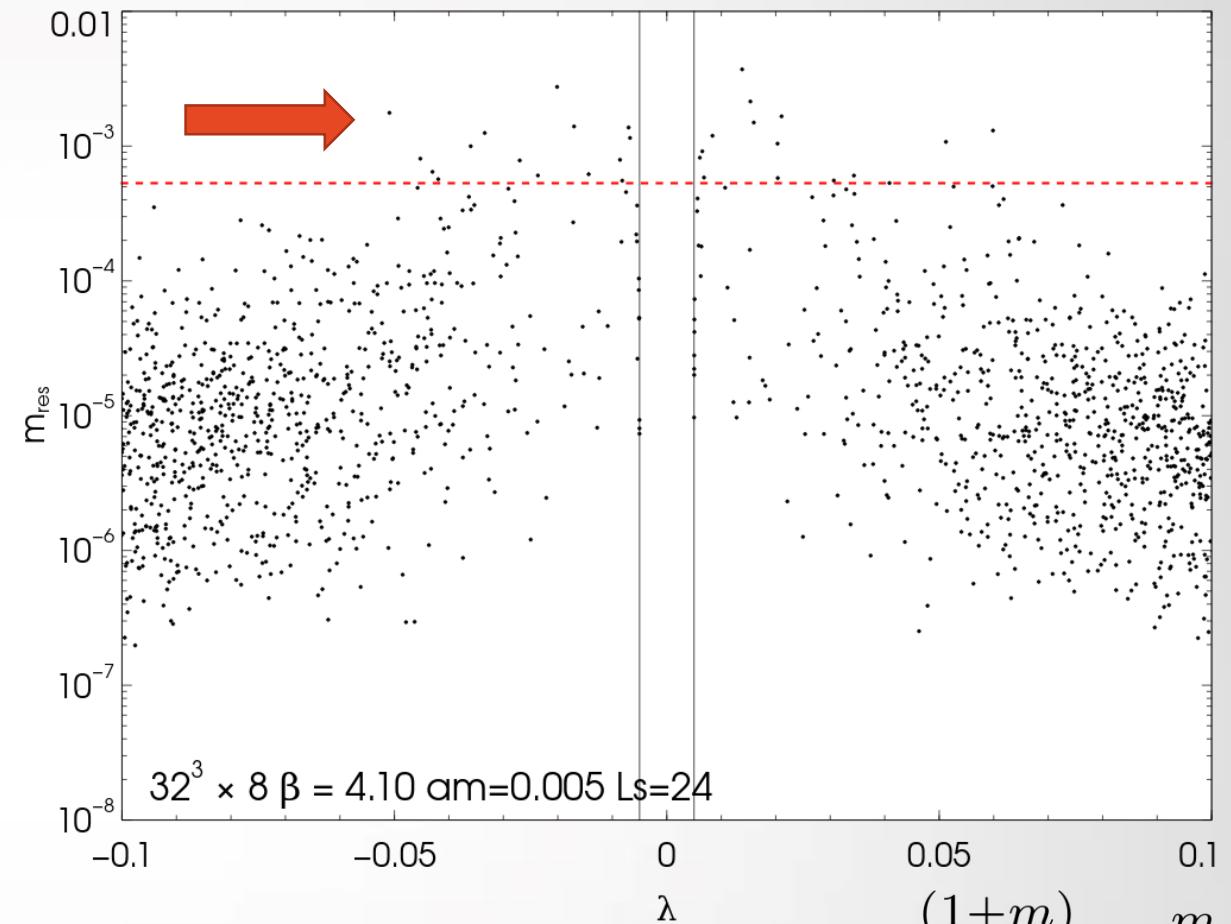
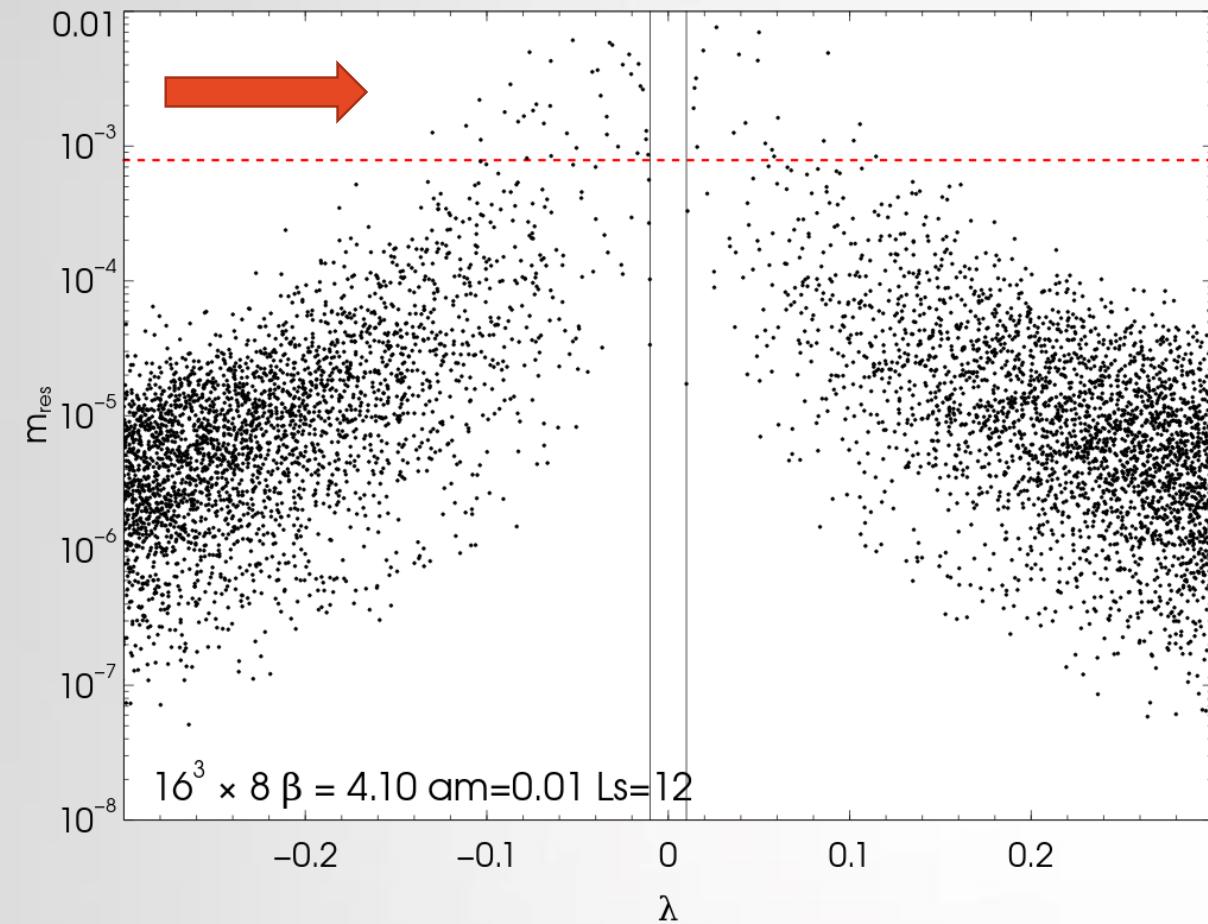
$$\{\hat{\gamma}_5, H_0\} = \Delta_{\text{viol}} \quad \hat{\gamma}_5 = \gamma_5 - H_m$$

$$\langle \psi_n | \gamma_5 | \psi_n \rangle = \frac{\lambda_n^2 + m}{\lambda_n(1+m)} + g_{nn}$$

$$m_{\text{res}} \sim \frac{\sum_n \frac{(1+m)}{(1-m)^2 \lambda_m^n} g_{nn}^m}{\sum_n \frac{1}{(\lambda_m^n)^2}}$$



# Is everything all right? – II



Lowest modes show violations of GW by 1  
**order of magnitude bigger than the average**

$$m_{\text{res},k} = \frac{\frac{(1+m)}{(1-m)^2 \lambda_m^k} g_{kk}^m}{\sum_k \frac{1}{(\lambda_m^k)^2}}$$

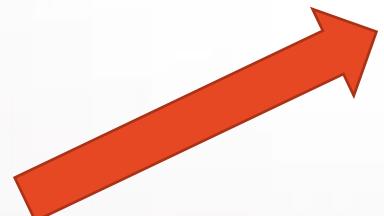
# Is everything all right? – III

**Exact result** for susceptibility

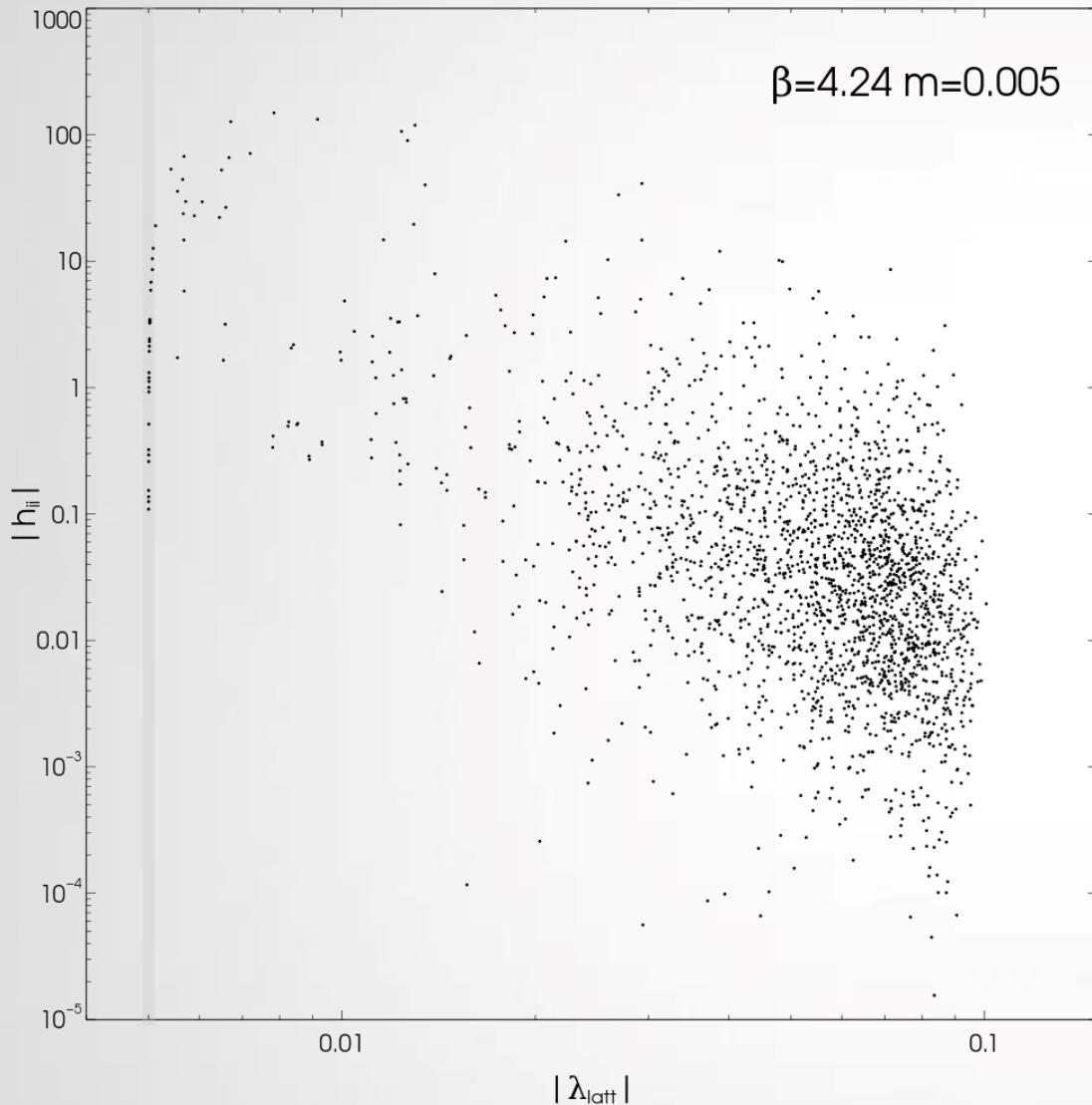
$$\begin{aligned}\chi_\pi - \chi_\delta = & \frac{1}{V(1-m^2)^2} \sum \frac{2m^2(1-\lambda_n^2)^2}{\lambda_n^4} + \\ & + \frac{1}{V(1-m)^2} \sum \left[ \frac{h_{nn}}{\lambda_n} - 4 \frac{g_{nn}}{\lambda_n} \right]\end{aligned}$$

GW violation terms

$$\{\hat{\gamma}_5, H_0\} = \Delta_{\text{viol}}$$



# Is everything all right? – IV

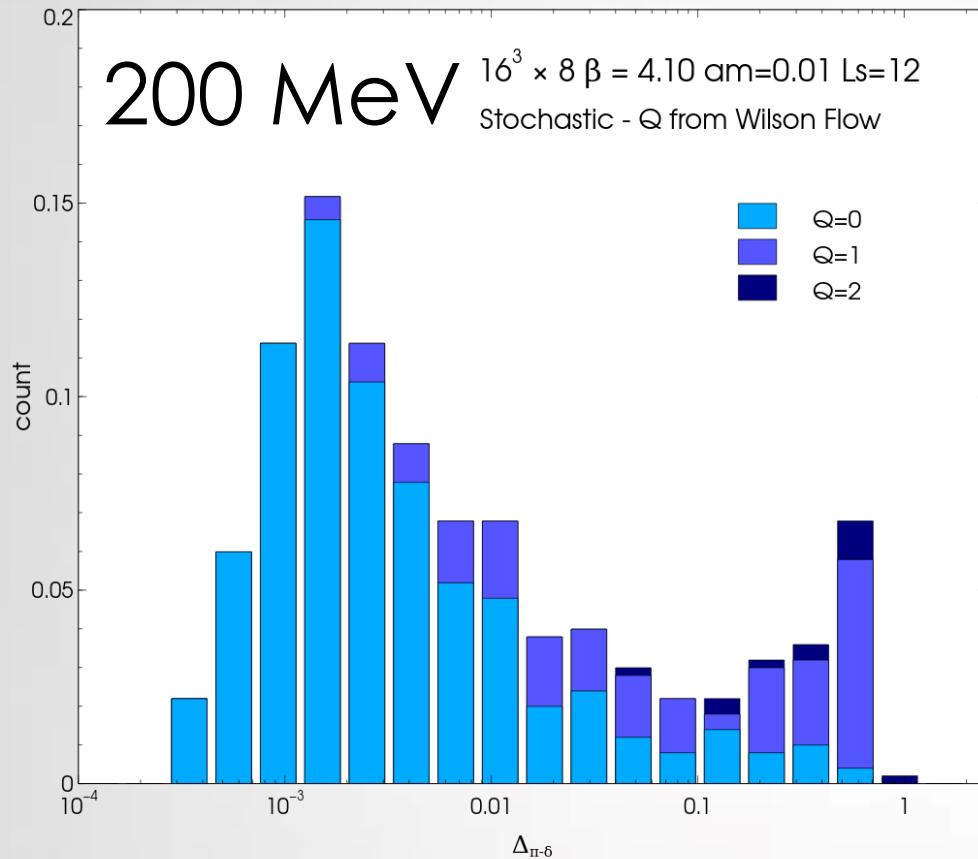


$$\chi_\pi - \chi_\delta = \frac{1}{V(1-m^2)^2} \sum \frac{2m^2(1-\lambda_n^2)^2}{\lambda_n^4} + \\ + \frac{1}{V(1-m)^2} \sum \left[ \frac{h_{nn}}{\lambda_n} - 4 \frac{g_{nn}}{\lambda_n} \right]$$

4.07	0.001	32	8	24	18	25	$0.967 \pm 0.006$	$0.90837 \pm 0.02868$
4.10	0.001	16	8	24	124	84	$0.690 \pm 0.030$	$1.05913 \pm 0.05022$
4.10	0.005	32	8	24	94	28	$0.463 \pm 0.031$	$0.94124 \pm 0.00959$
4.10	0.001	32	8	24	43	36	$0.928 \pm 0.022$	$0.88578 \pm 0.02801$
4.18	0.01	32	12	16	54	17	$0.080 \pm 0.005$	$0.98895 \pm 0.00402$
4.22	0.01	32	12	16	50	38	$0.056 \pm 0.007$	$1.01573 \pm 0.01243$
4.23	0.01	32	12	16	55	27	$0.036 \pm 0.005$	$0.99317 \pm 0.00963$
4.23	0.005	32	12	16	55	22	$0.143 \pm 0.022$	$0.97322 \pm 0.01076$
4.24	0.01	32	12	16	249	38	$0.045 \pm 0.003$	$0.99284 \pm 0.00774$
4.24	0.005	32	12	16	69	33	$0.115 \pm 0.022$	$0.99036 \pm 0.00675$

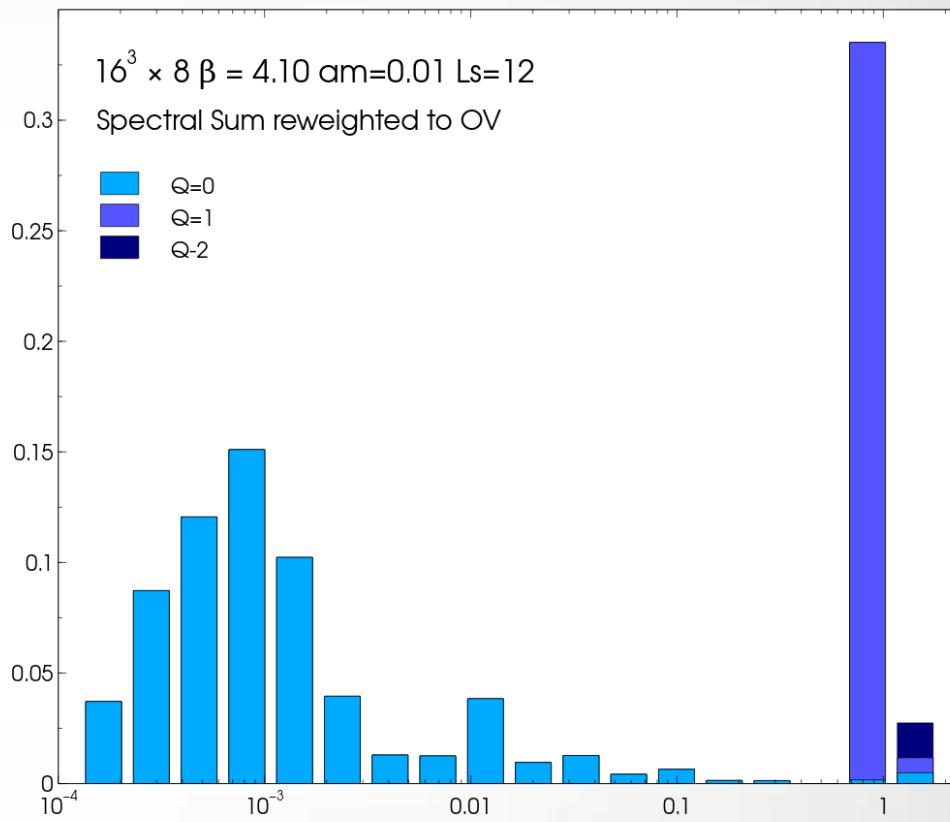
# Reweighting (DWF to Overlap)

# Before



76%

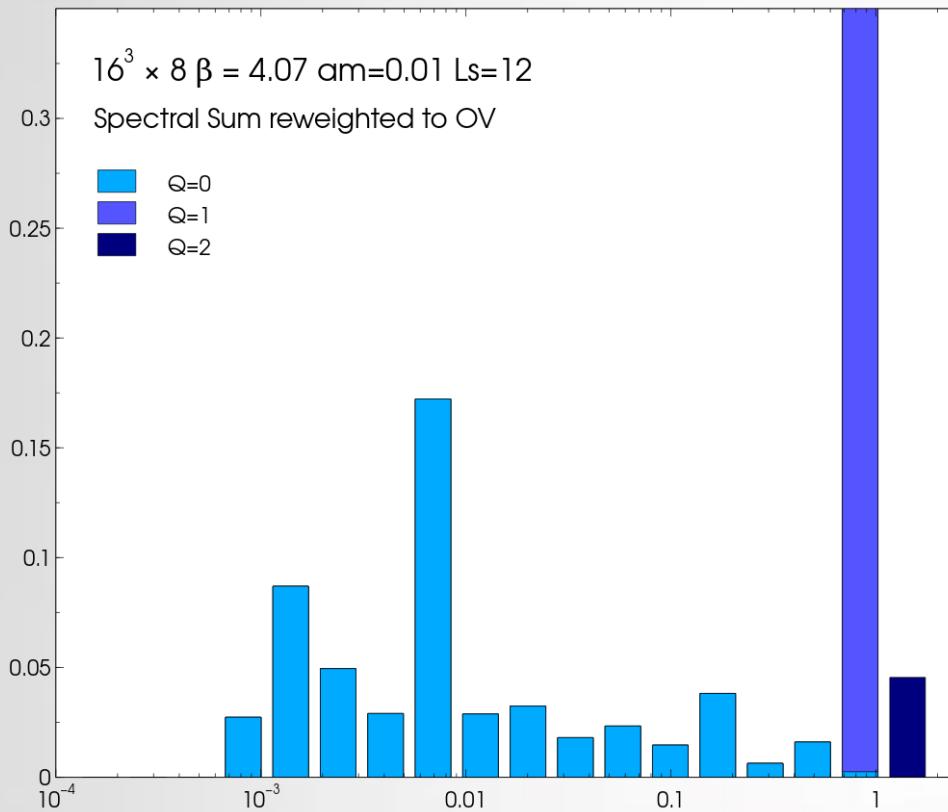
# After



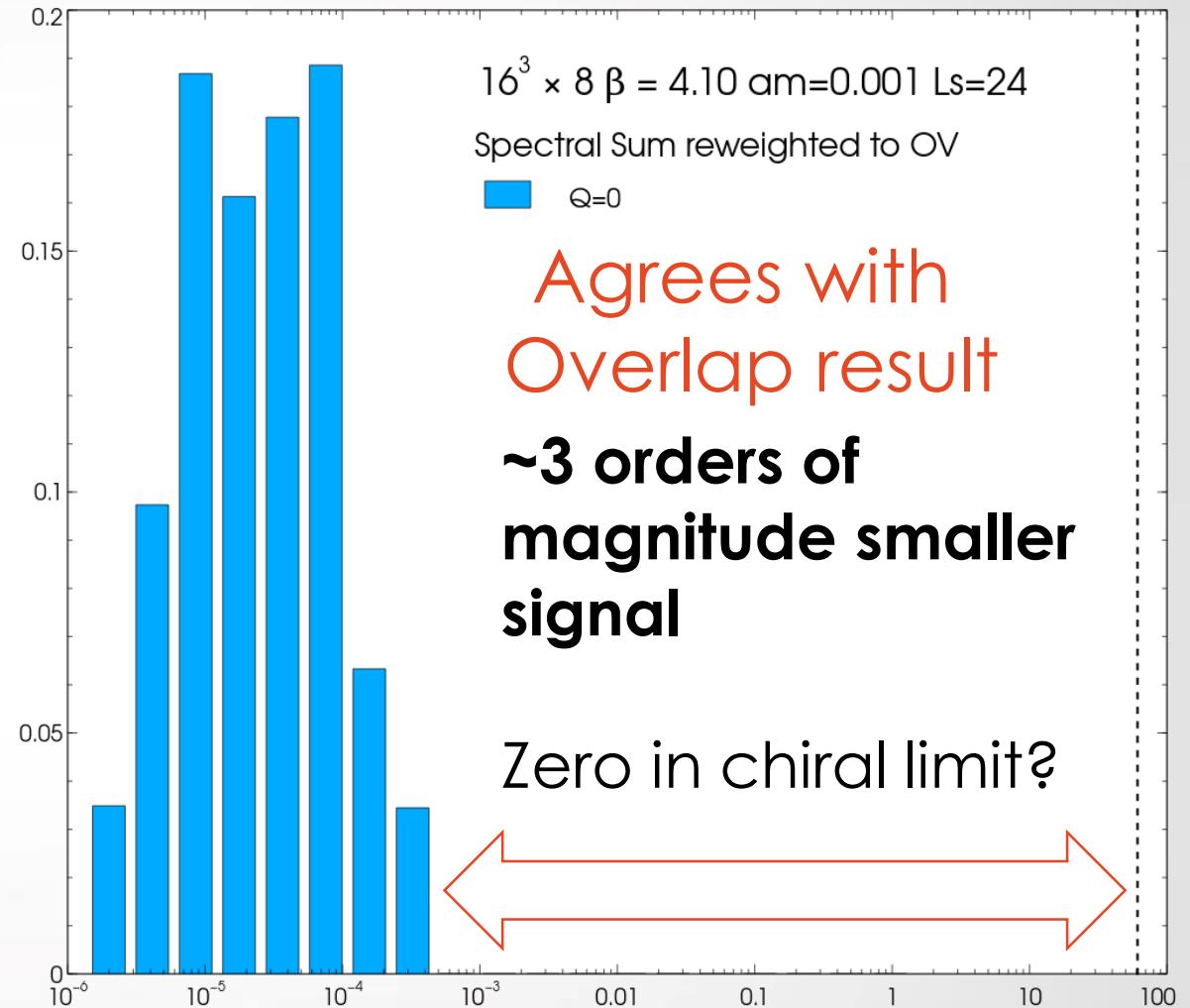
85%

# Temperature and mass dependence

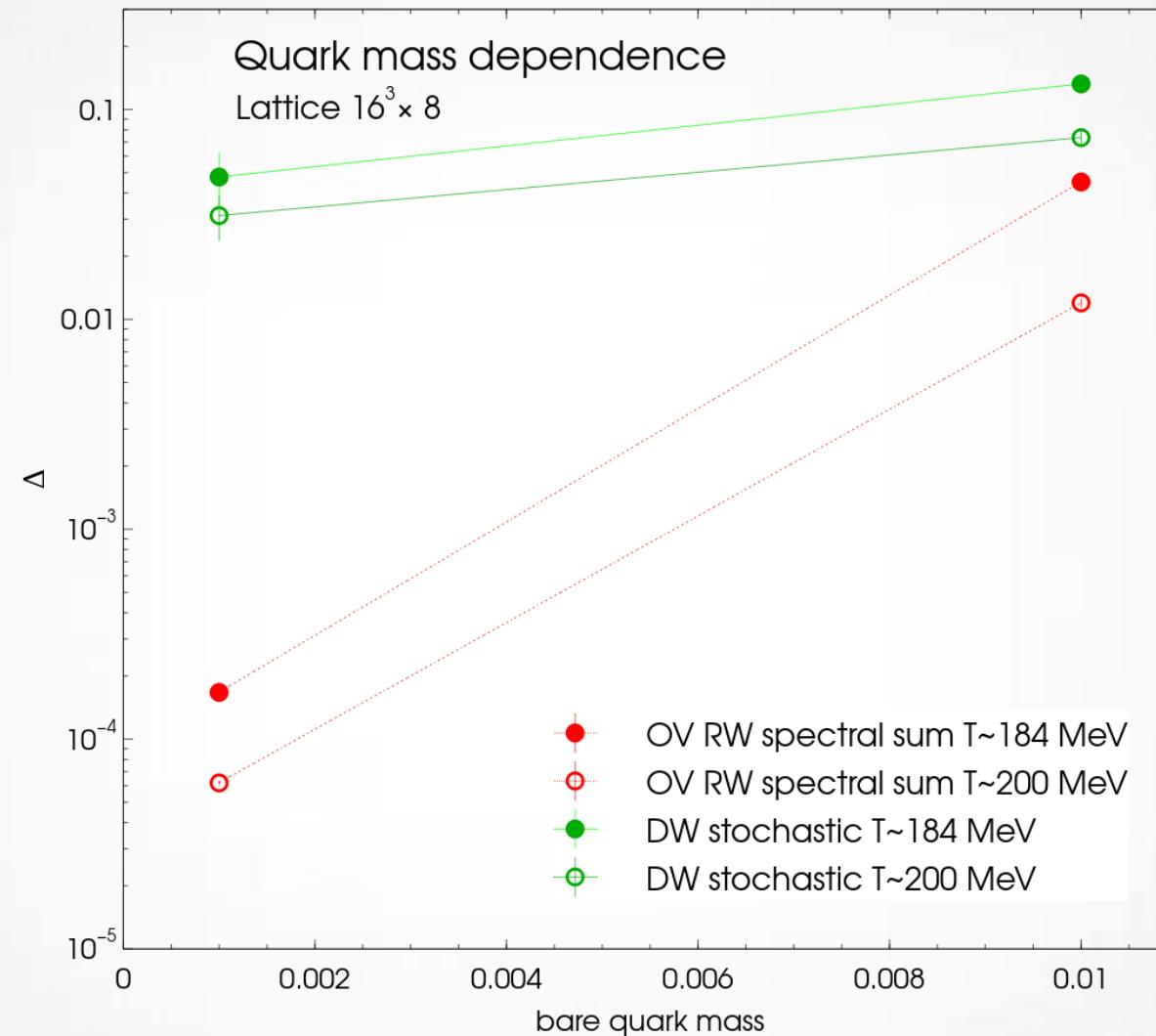
$T \sim 184$  MeV Just above  $T_c$



Quark mass 10 times smaller



# Quark mass dependence



# Methods & Results

**Work in progress**

$U(1)_A$   
symmetry  
Finite  
temperature

Literature

Final  
thoughts

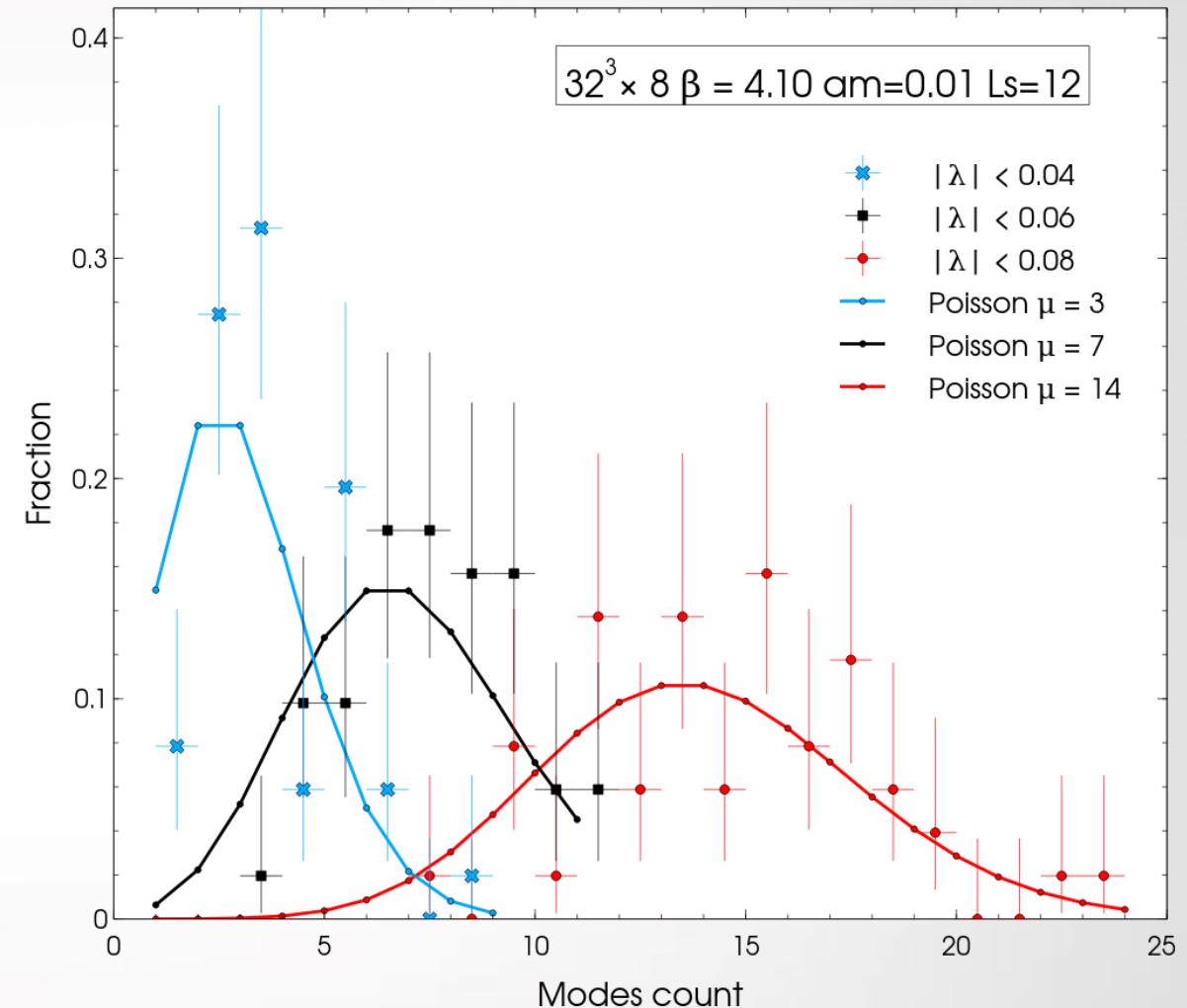
# Instanton gas – hints?

**Results not yet conclusive  
(analysis running right now)**

**If the large volume signal is  
not coming from lattice  
artifacts**

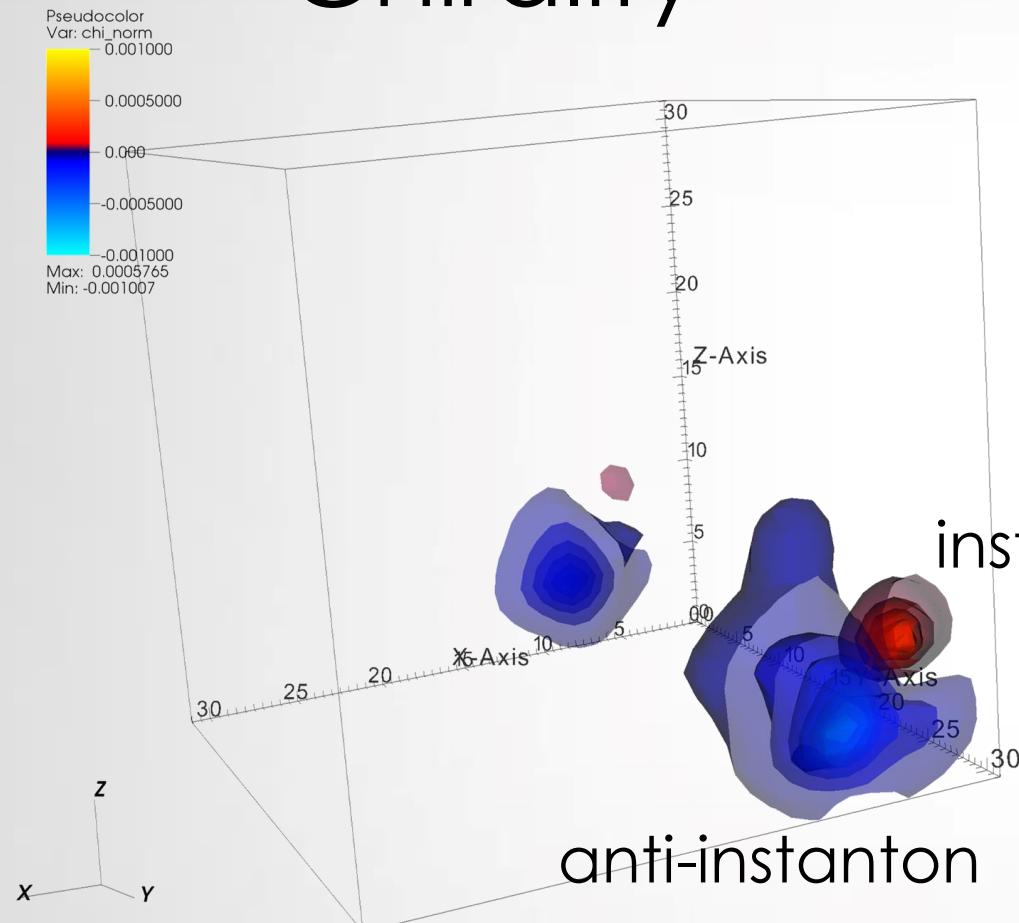
Near zero modes are  
responsible for breaking U(1)

*What are they?  
**Poisson distributed?***

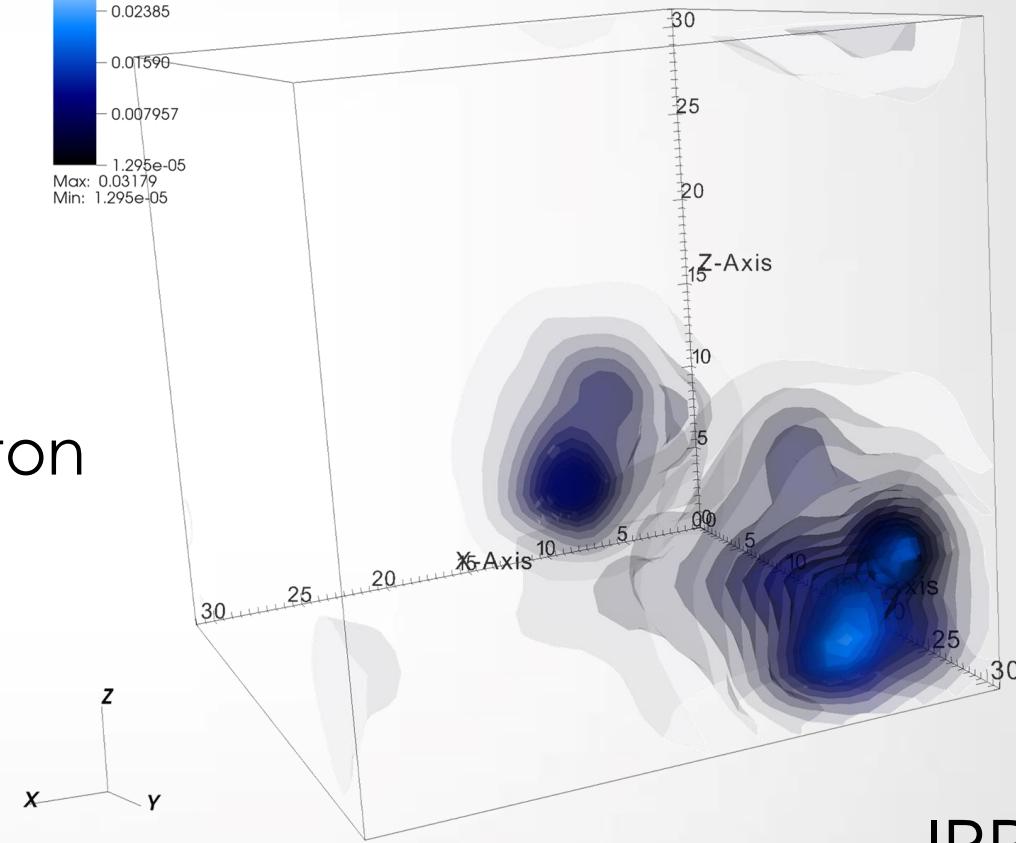


# Fun with 3D – put your glasses on

## Chirality



## Norm



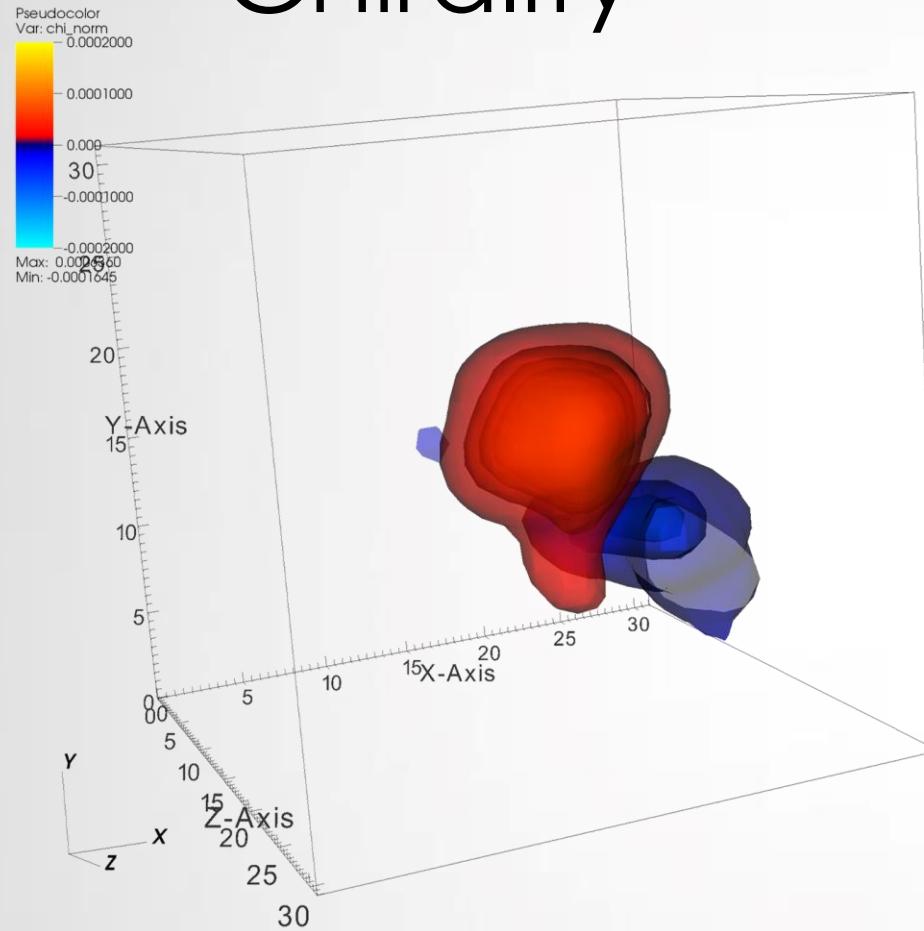
Eigenvalue



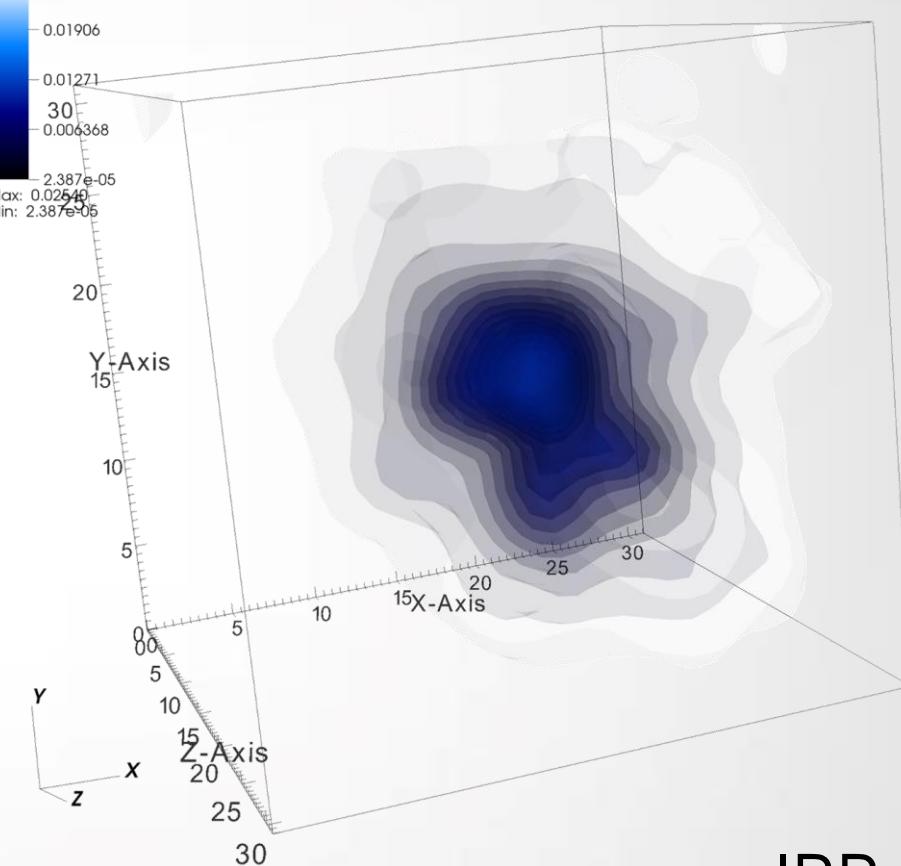
IPR ~34

# Fun with 3D – put your glasses on

## Chirality



## Norm



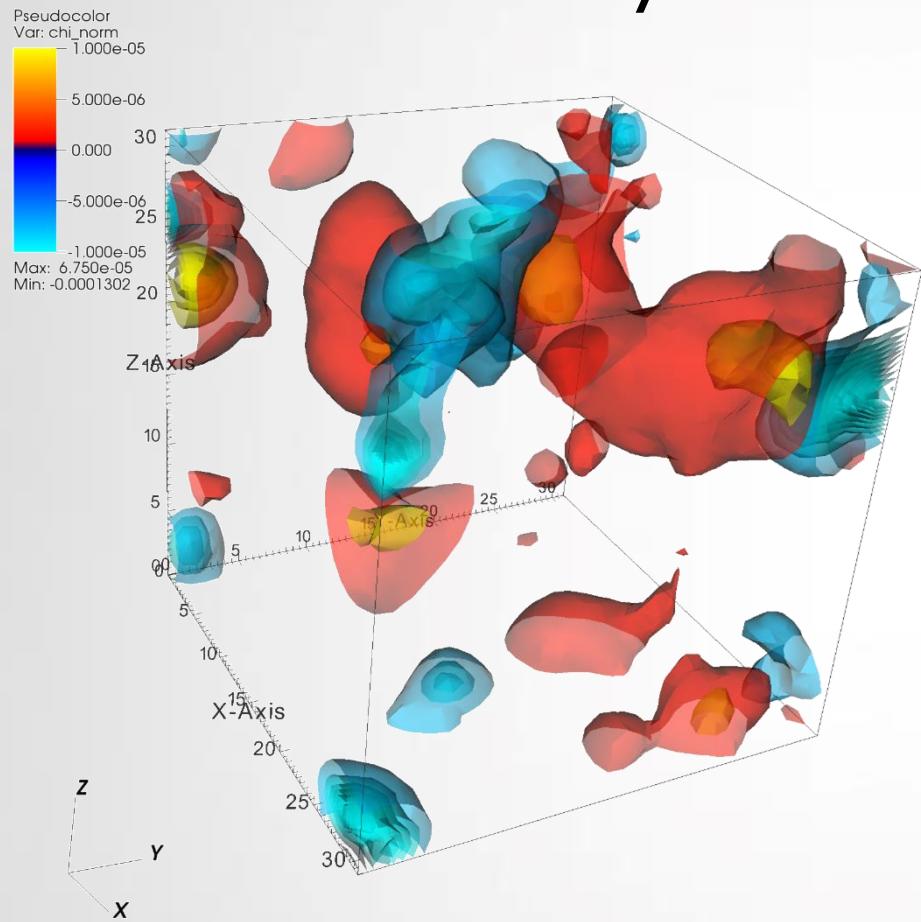
Eigenvalue



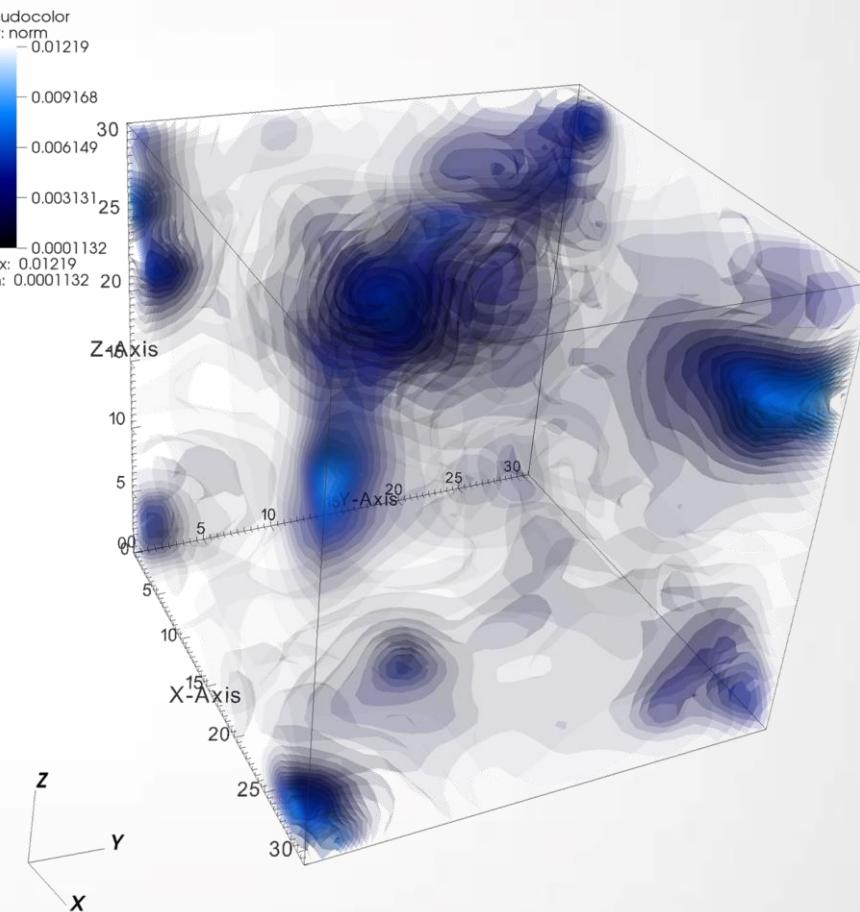
IPR ~34

# Fun with 3D – put your glasses on

## Chirality



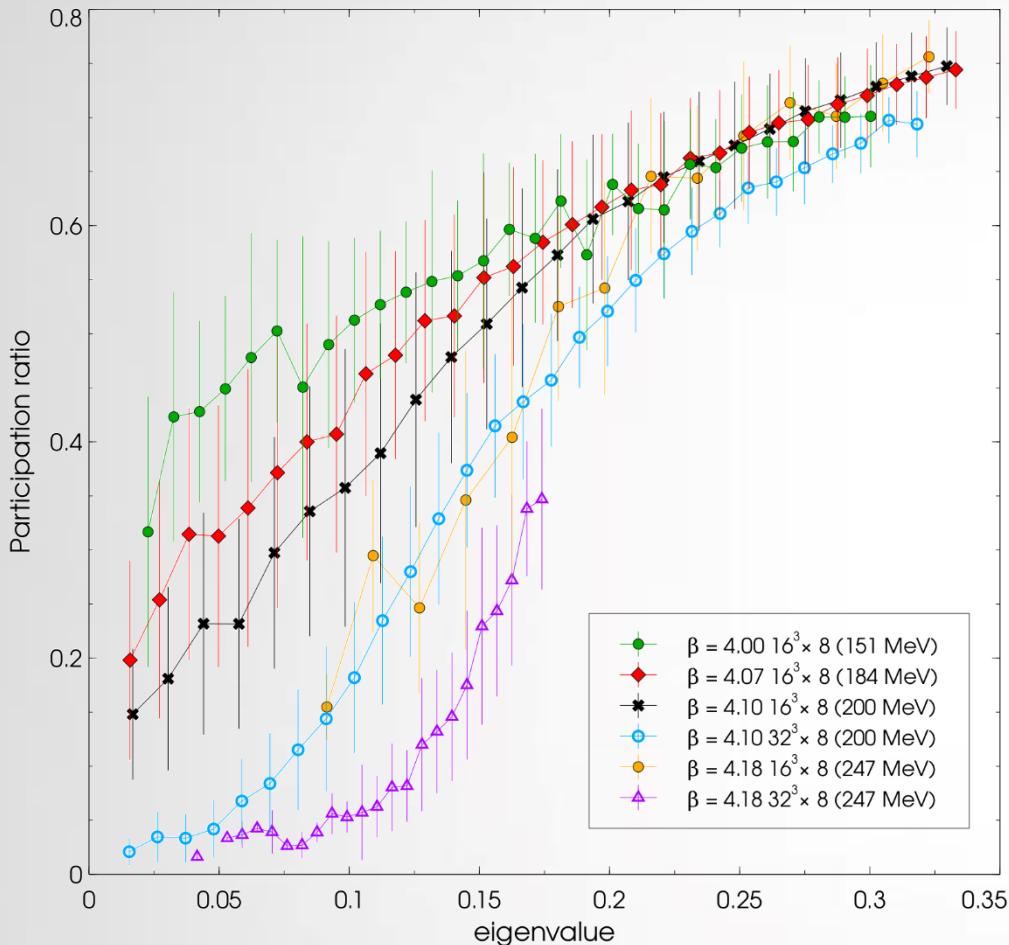
## Norm



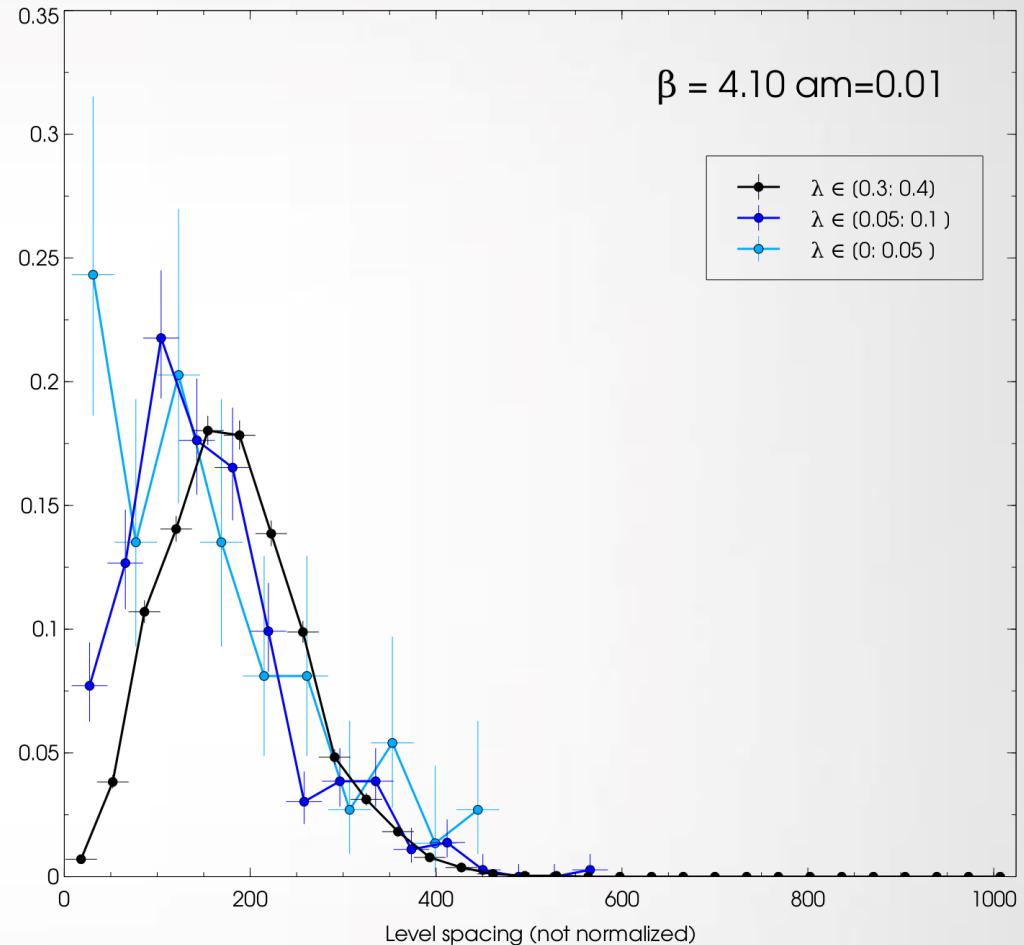
Eigenvalue

IPR ~5

# Localized modes



$$\langle PR \rangle = (V \sum (\psi^\dagger(x)\psi(x))^2)^{-1}$$



# Summary – one more slide...

DWF volume & mass dependence suggests that **near zero modes are the source of U(1) breaking**

Lattice artifacts can spoil the signal

Exact chiral symmetry results differ from DWF

DWF lowest modes look like an instanton weakly interacting gas

# Are we finished?

The talk is over the work is not! (but almost there)

Some collected data yet to analyze

- Reweighting
- Continuum limit
- Chiral limit

Lattice artifacts?

Gas of instanton pairs, dyons?  
Correlation with Polyakov loop?  
U(1) restoration above critical  
temperature  
**is still an open question.**



**Thanks!**

# Axial symmetry at finite temperature and Dirac operator eigenmodes

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**Guido Cossu**

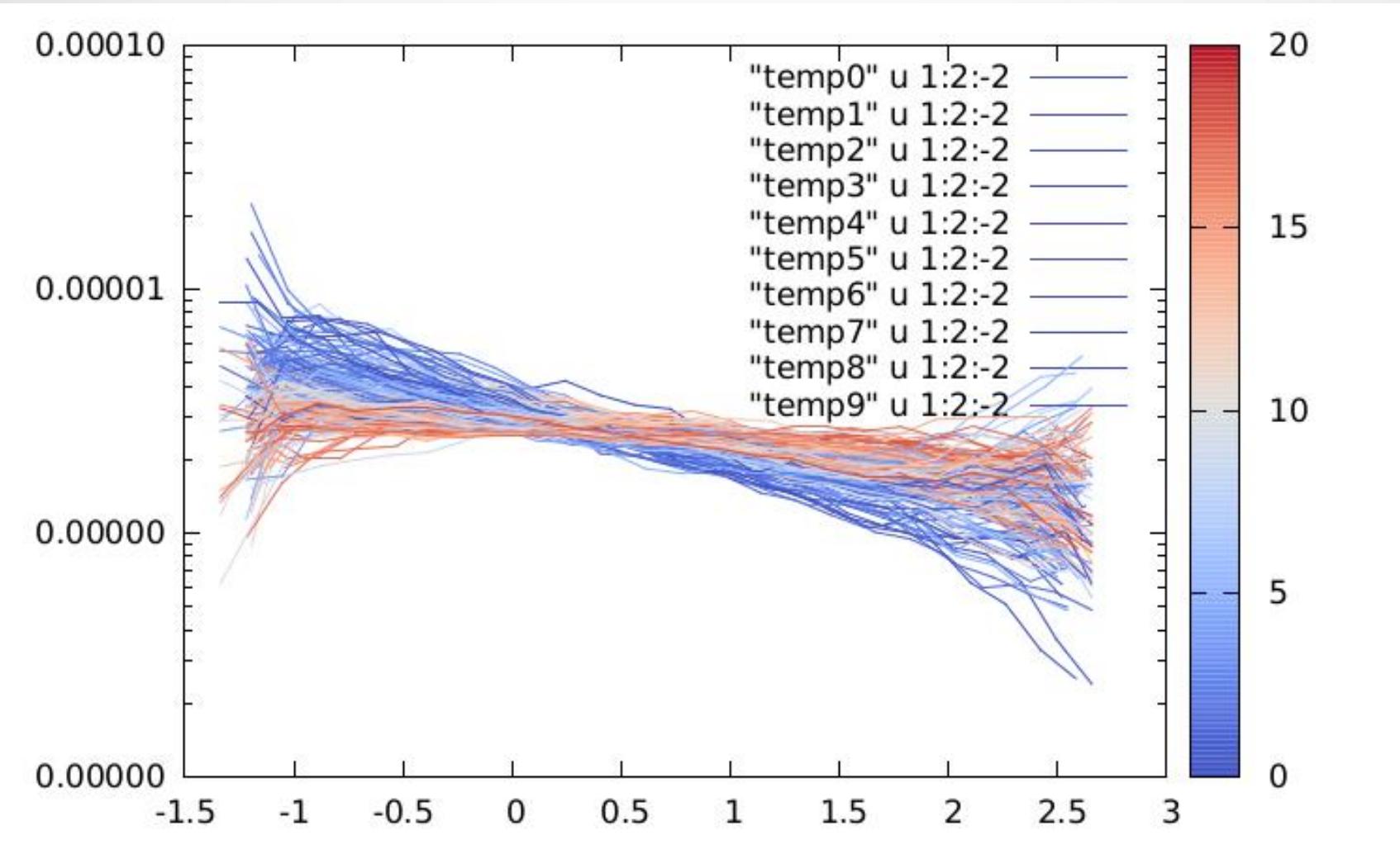
High Energy Accelerator Organization (KEK)

Joint Institute for Computational Fundamental Science (JICFuS)

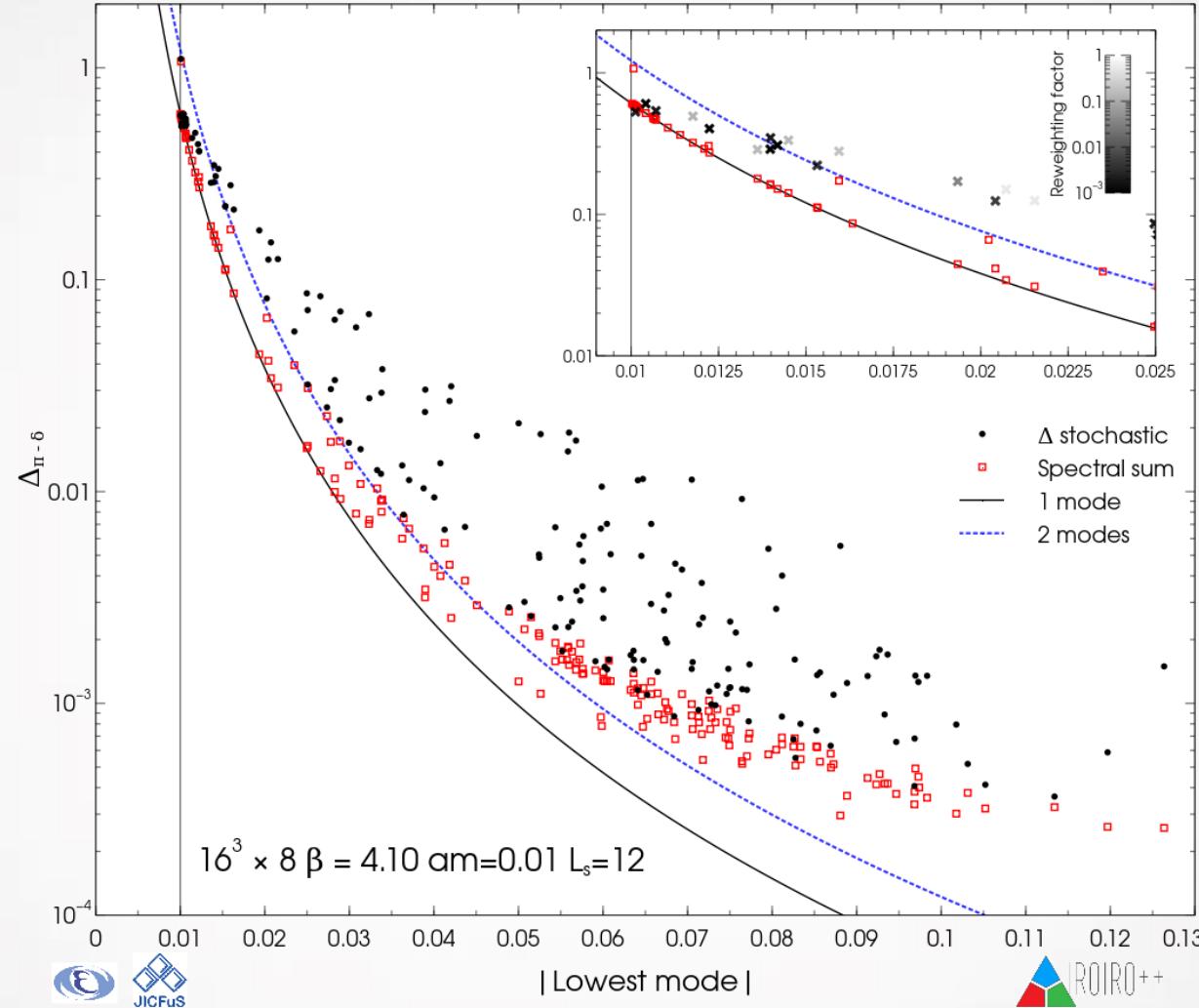
YITP HHIQCD2015, Kyoto  
March 5th 2015



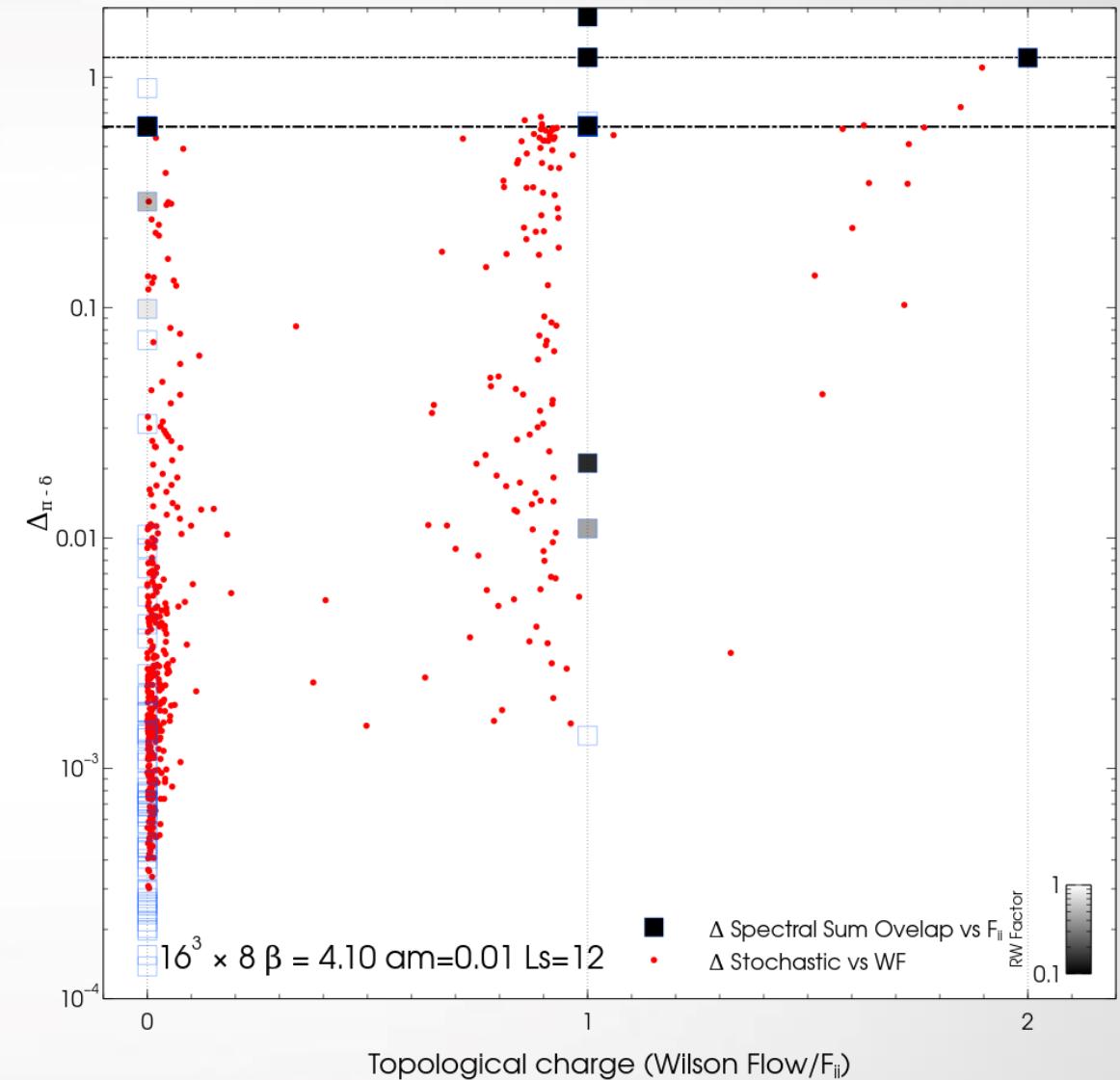
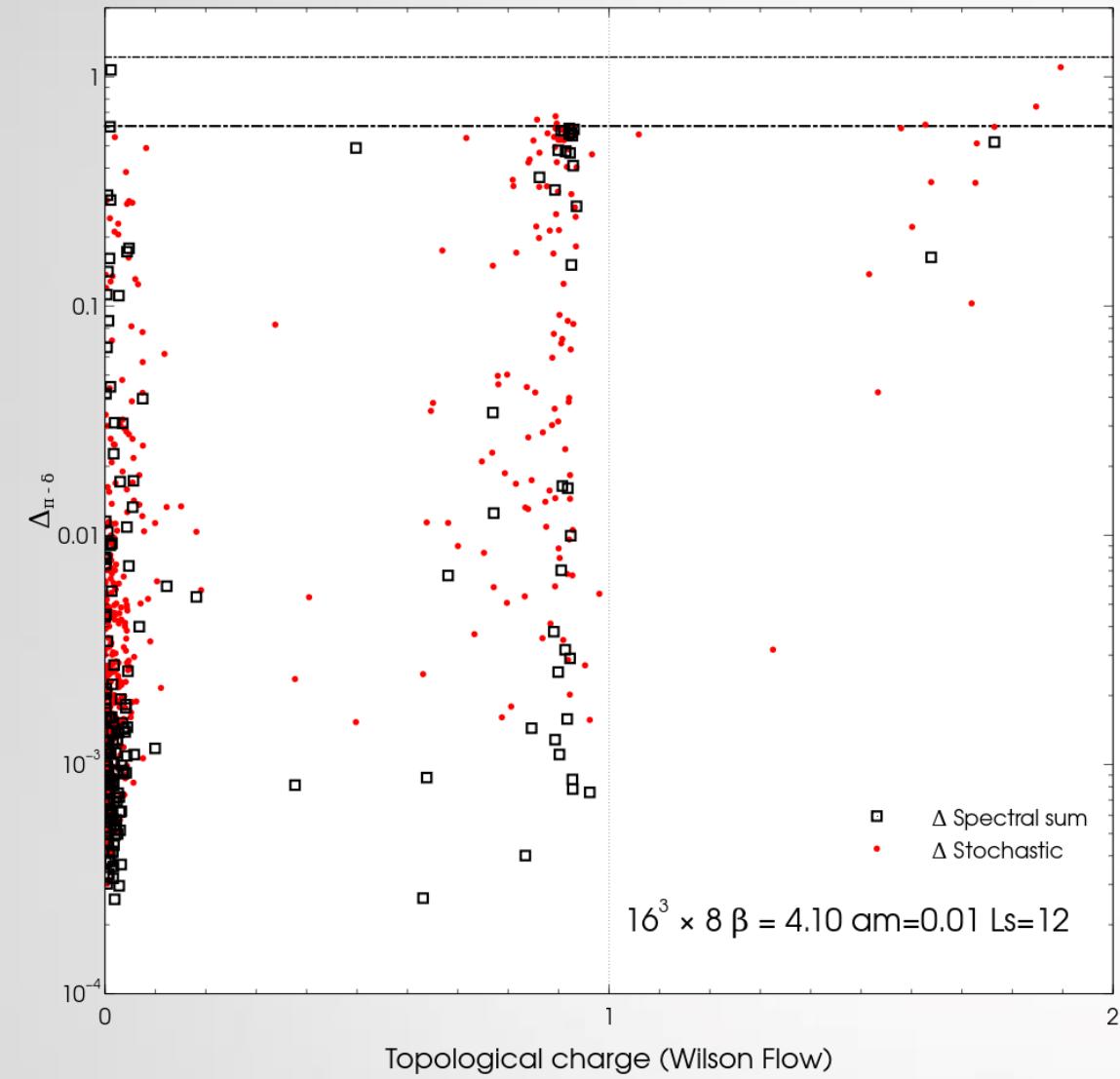
Backup  
slides



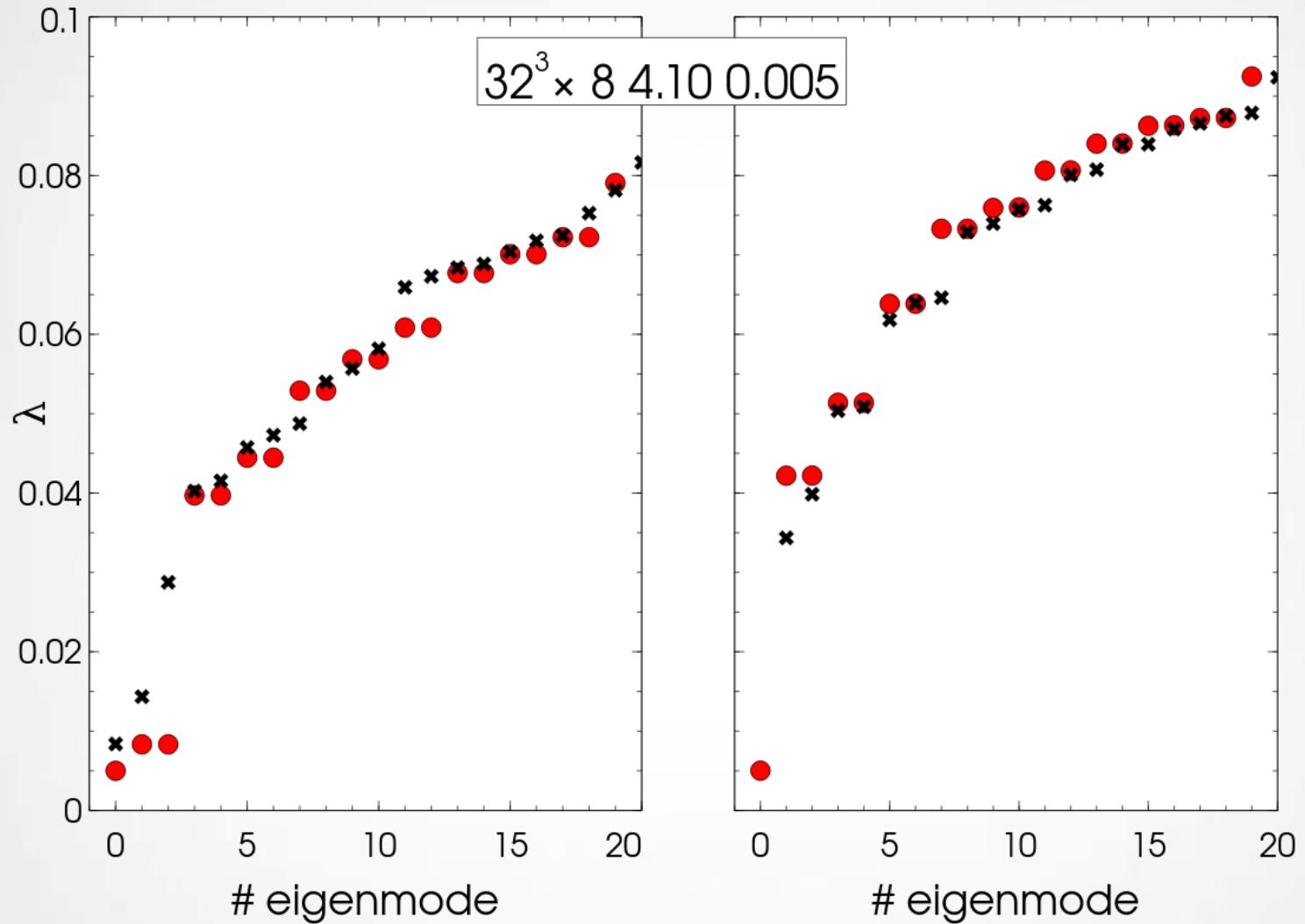
# GW violations



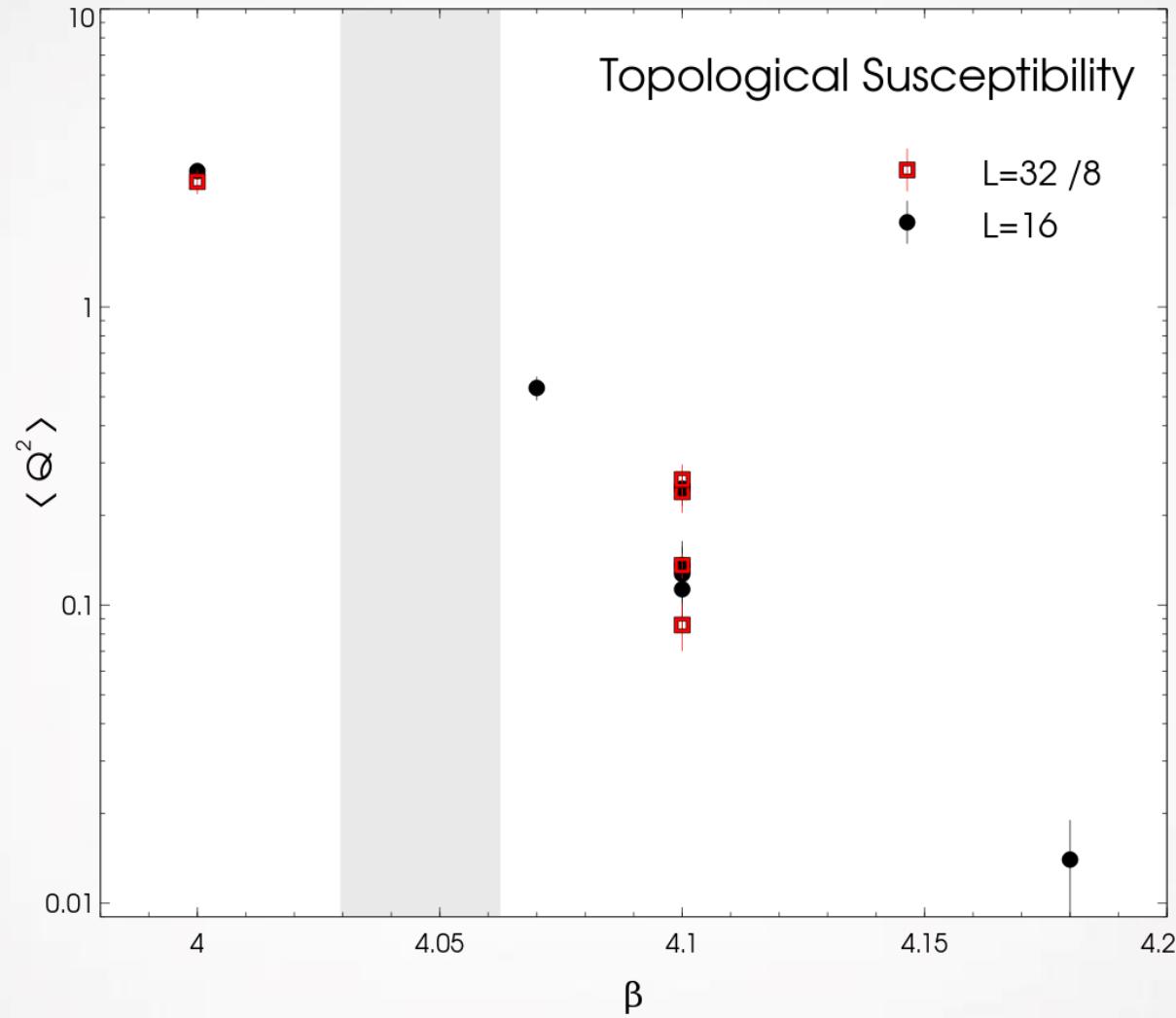
# Many configurations violate GW



# DW – OV eigenvalue mismatch



# Susceptibility scales with volume



# Let's increase volume

Zero mode contribution  $\sim 1/V$  - Bulk contribution increases

