

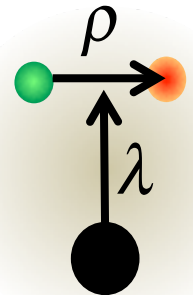
Charmed baryons

- Structure, productions and decays -

Atsushi Hosaka, RCNP, Osaka
HHIQCD - symposium, YITP, Kyoto, 2/15-3/21

With Noumi, Shirotori, Kim, Sadato, Yoshida, Oka,
Hiyama, Yasui, Nagahiro

Prog. Theor. Exp. Phys. (2014) 103D01



Contents

1. Introduction -- Exotics with heavy quarks
2. Charmed baryons: *to probe $\rho\lambda$ modes* → **diquark**
3. Productions
4. Decays

1. Introduction

Quark model and EXOTICS: Now 51 years old

A SCHEMATIC MODEL OF BARYONS AND MESONS

M. GELL-MANN

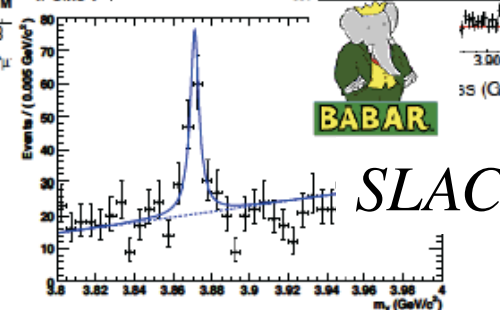
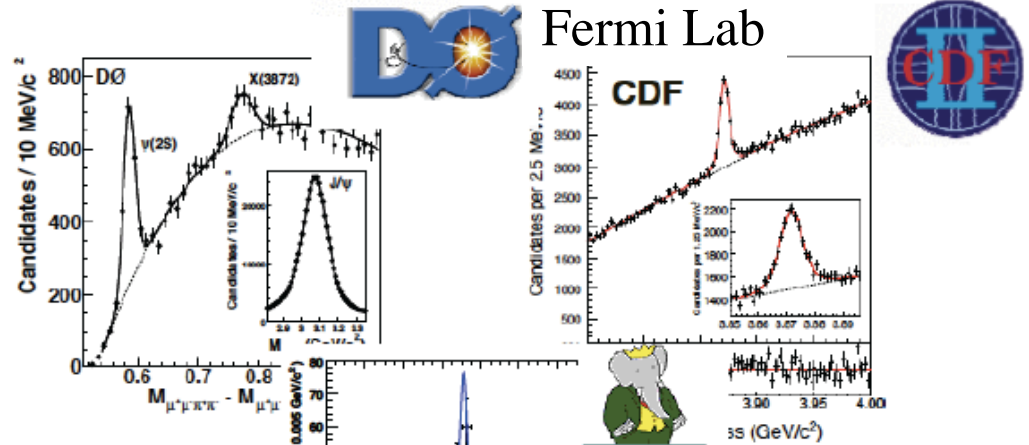
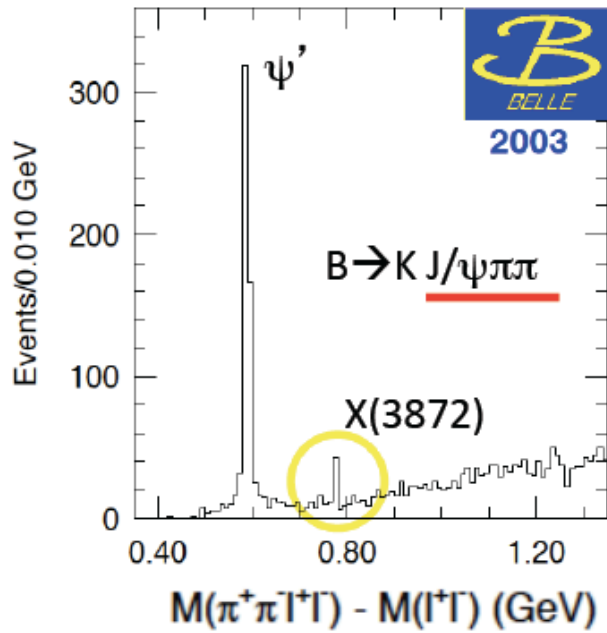
California Institute of Technology, Pasadena, California

Received 4 January 1964

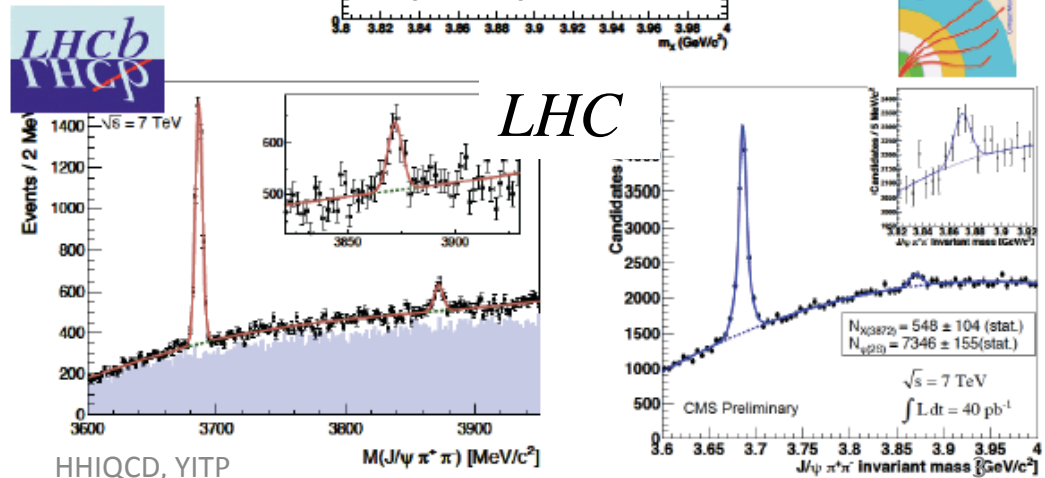
anti-triplet as anti-quarks \bar{q} . Baryons can now be constructed from quarks by using the combinations (qqq) , $(qqq\bar{q})$, etc., while mesons are made out of $(q\bar{q})$, $(qq\bar{q}\bar{q})$, etc. It is assuming that the lowest baryon configuration (qqq) gives just the representations **1**, **8**, and **10** that have been observed, while the lowest meson configuration $(q\bar{q})$ similarly gives just **1** and **8**.

X (3872)

Discovery by Belle in 2003, followed by D0, CDF, BaBar.



And more recently also by LHCb, CMS



Quarks bonding differently at LHCb $Z^+(4430)$

<http://www.theguardian.com/science/life-and-physics/2014/apr/13/quarks-bonding-differently-at-lhcb>



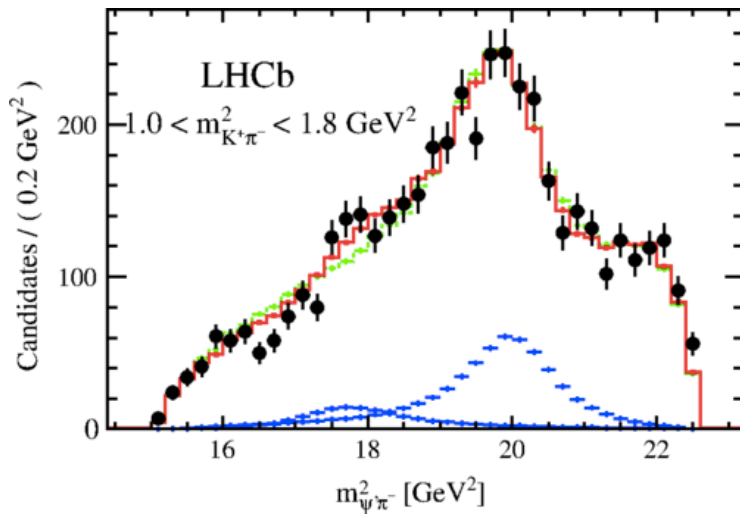
So until last week there were two known types of hadron.

.....

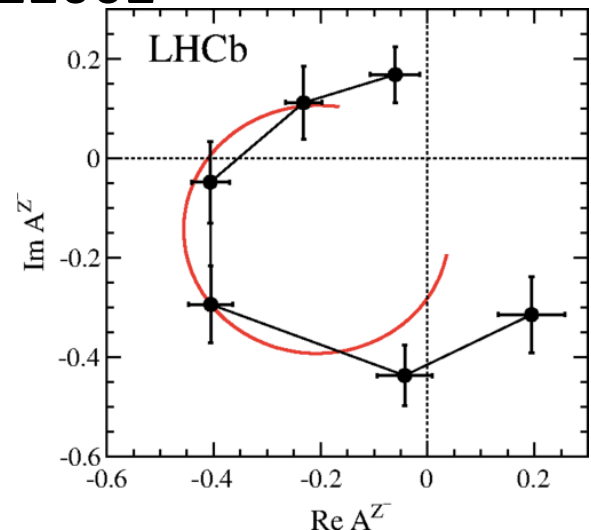
LHCb has just confirmed what data from other experiments had already led us to suspect.

There is a third way.

Phys. Rev. Lett. **112**, 222002

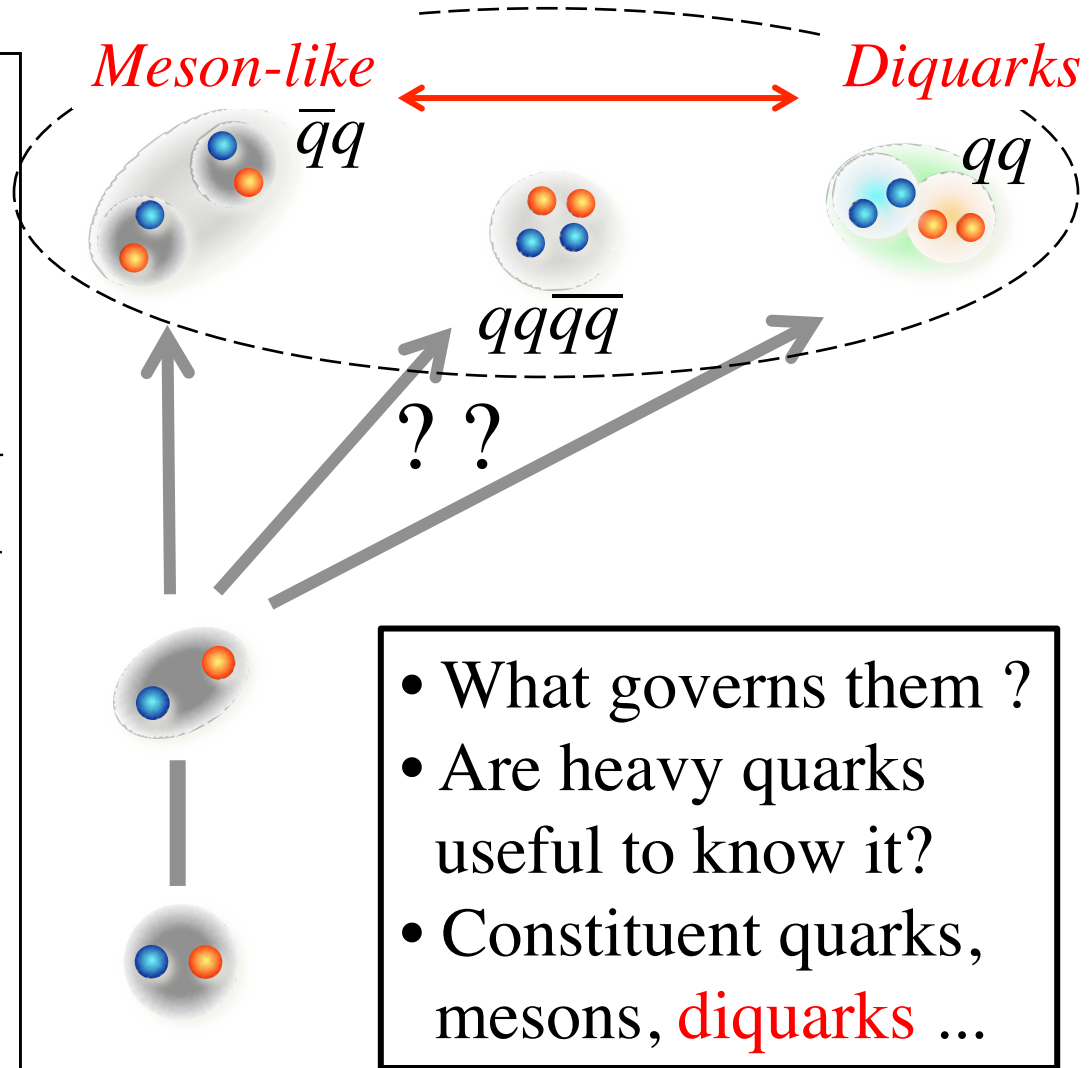
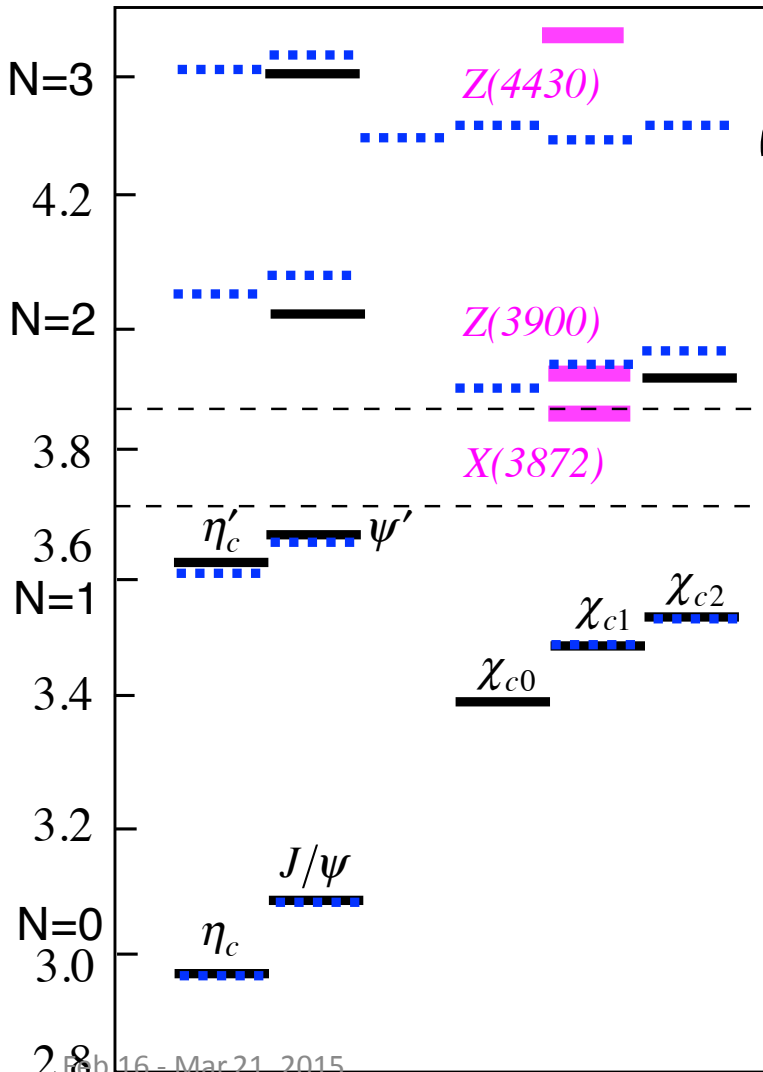


HHIQCD, YITP



EXOTIC threshold (and above)

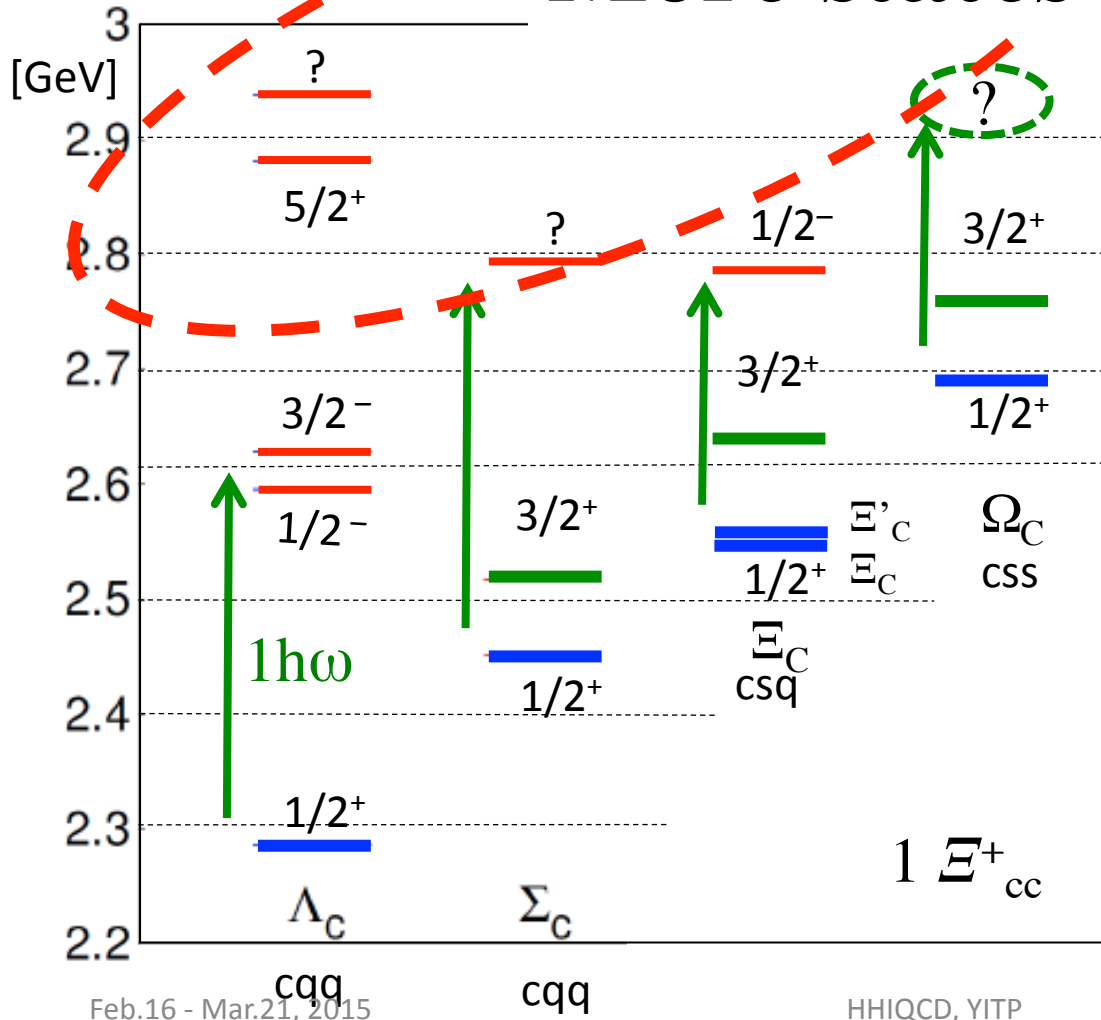
$q\bar{q}$ creation and rearrangement of multiquarks



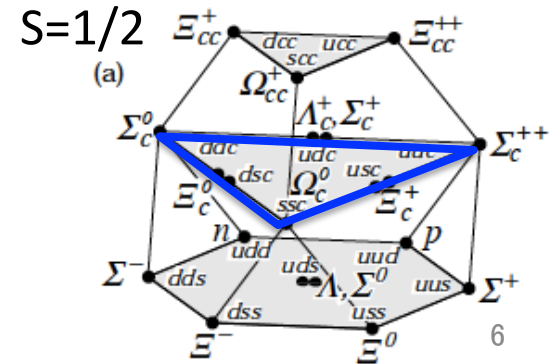
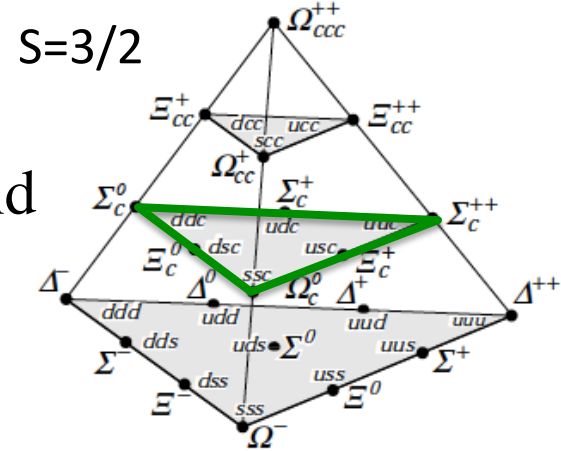
- What governs them ?
- Are heavy quarks useful to know it?
- Constituent quarks, mesons, **diquarks** ...

2. Charmed baryons

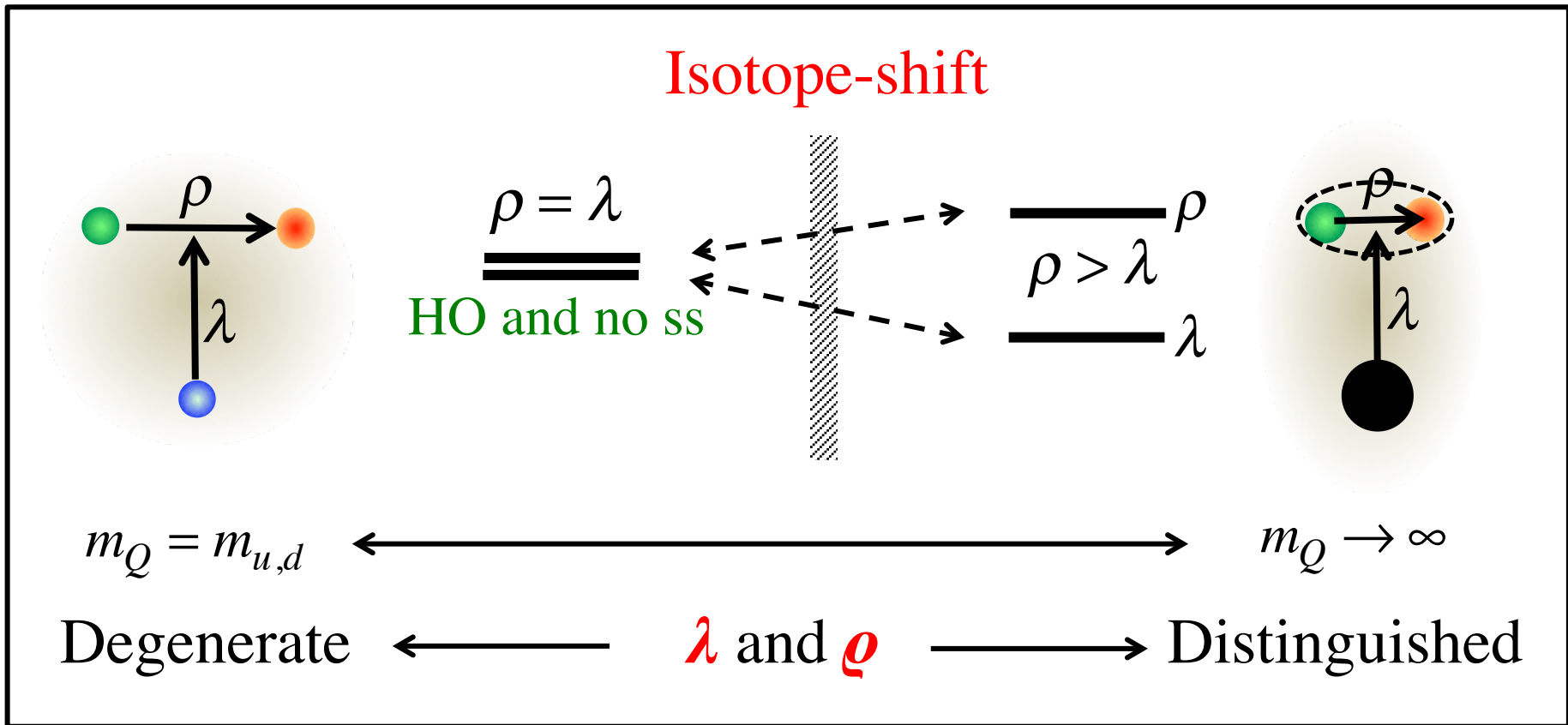
More states at J-PARC



Ground states



Qqq distinguishes the fundamental modes
 λ and $\rho \rightarrow$ place to look at diquarks



Quark model calculation

Yoshida, Sadato, Hiyama, Oka, Hosaka

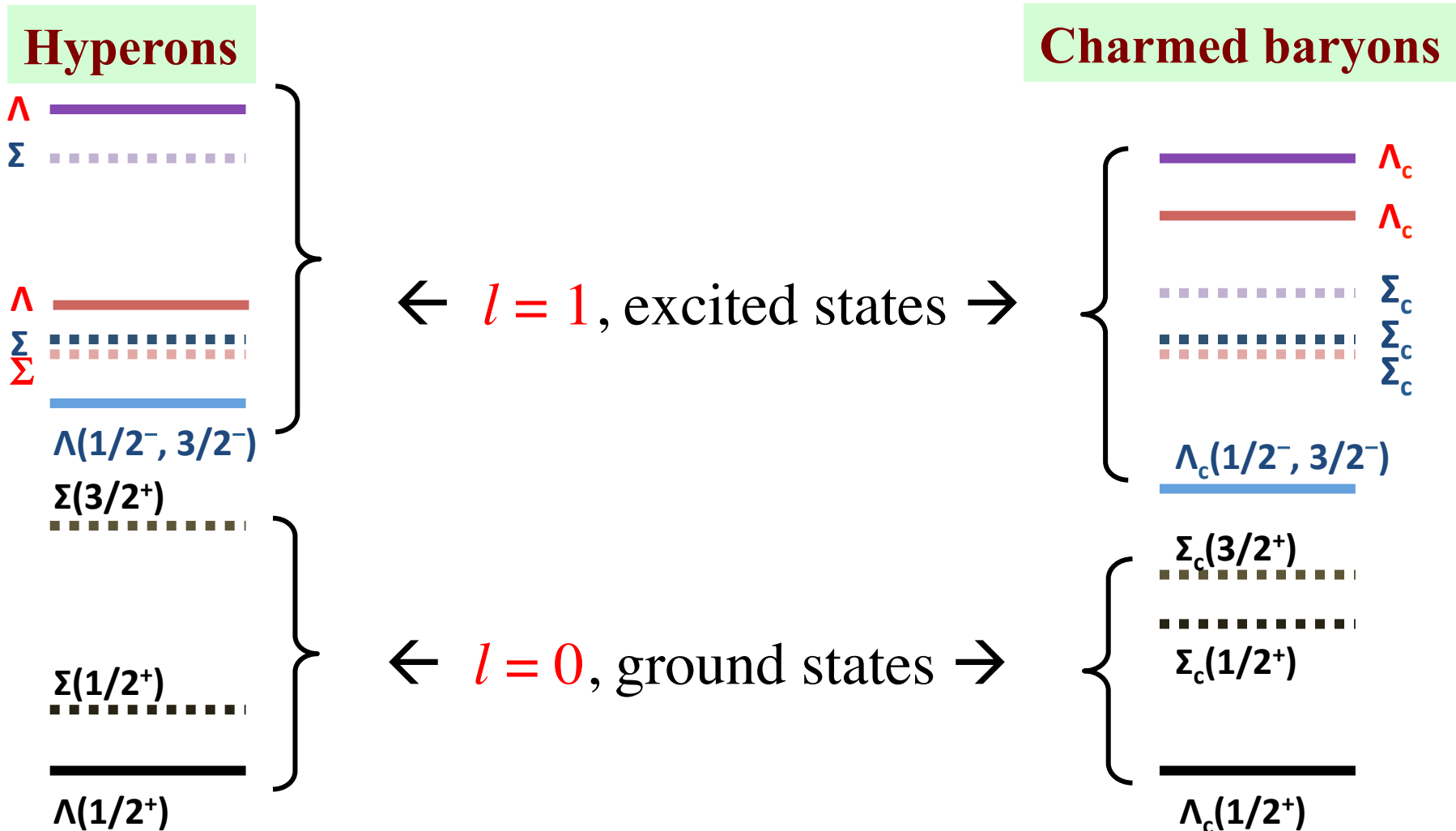
- Model Hamiltonian

$$H = \frac{p_1^2}{2m_q} + \frac{p_2^2}{2m_q} + \frac{p_3^2}{2M_Q} - \frac{P^2}{2M_{tot}} \\ + V_{conf}(HO) + V_{spin-spin}(Color - magnetic) + \dots$$

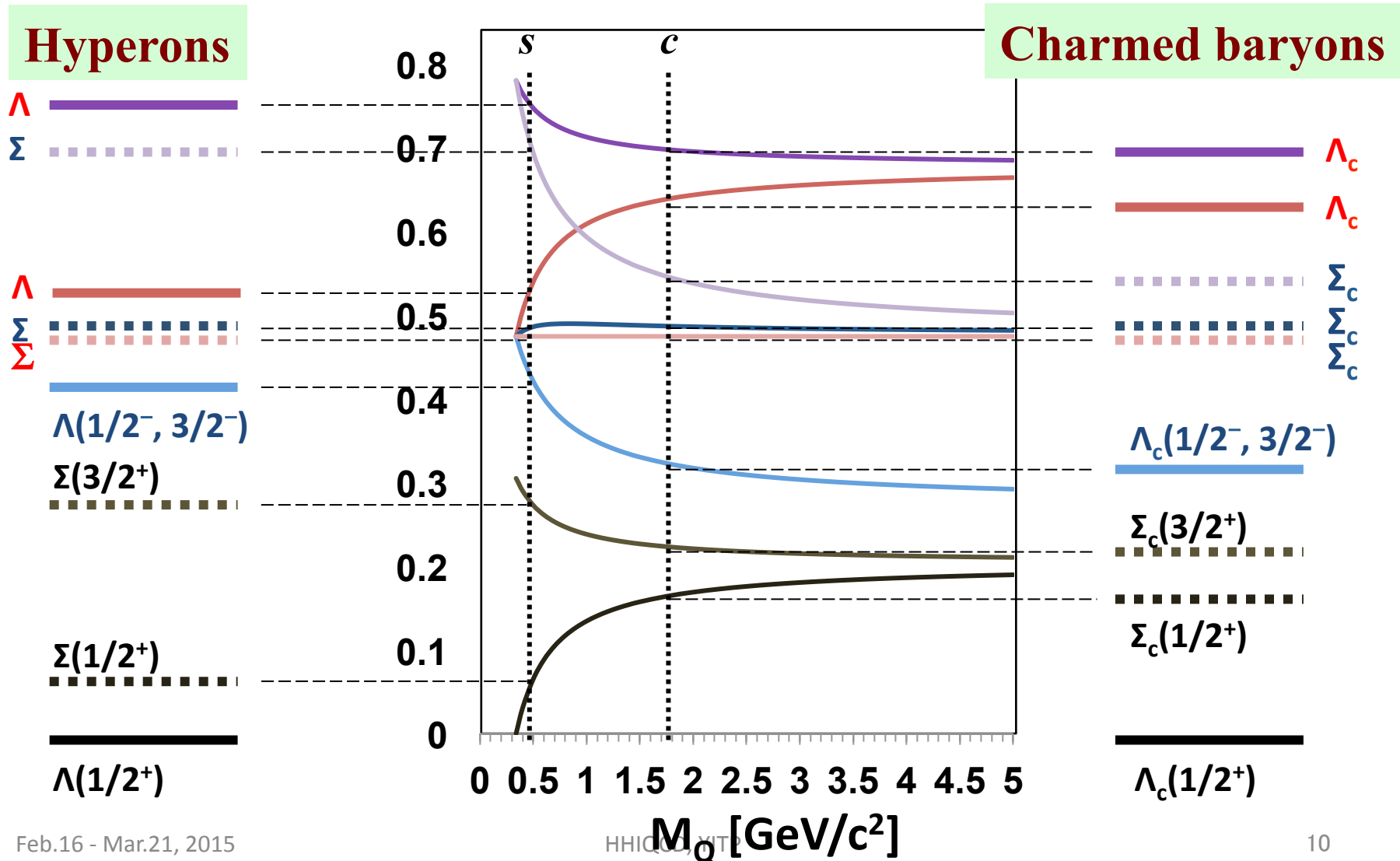
- Solved by the Gaussian expansion method

See how systems change as M_Q is varied

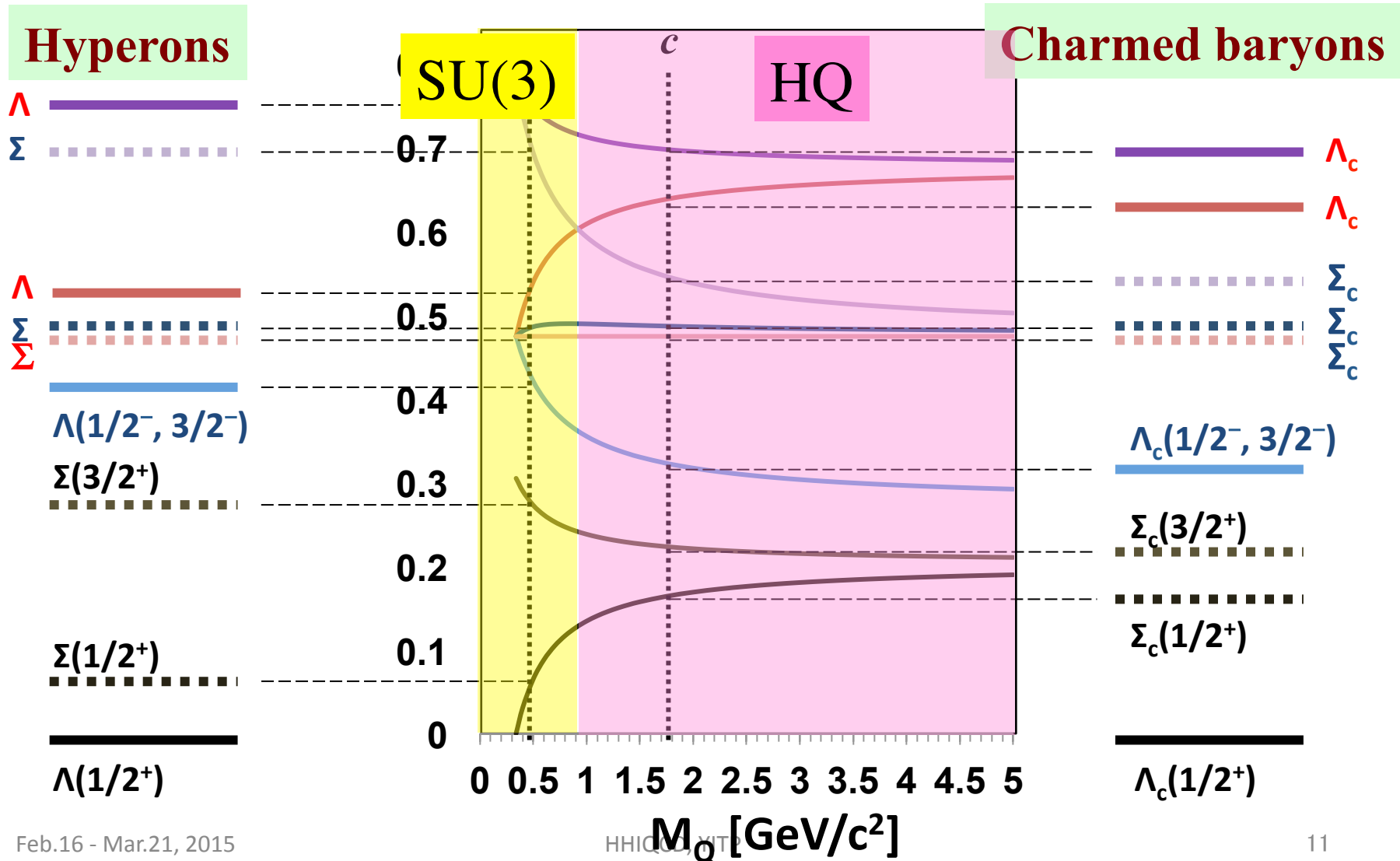
Excitation spectrum



Excitation spectrum



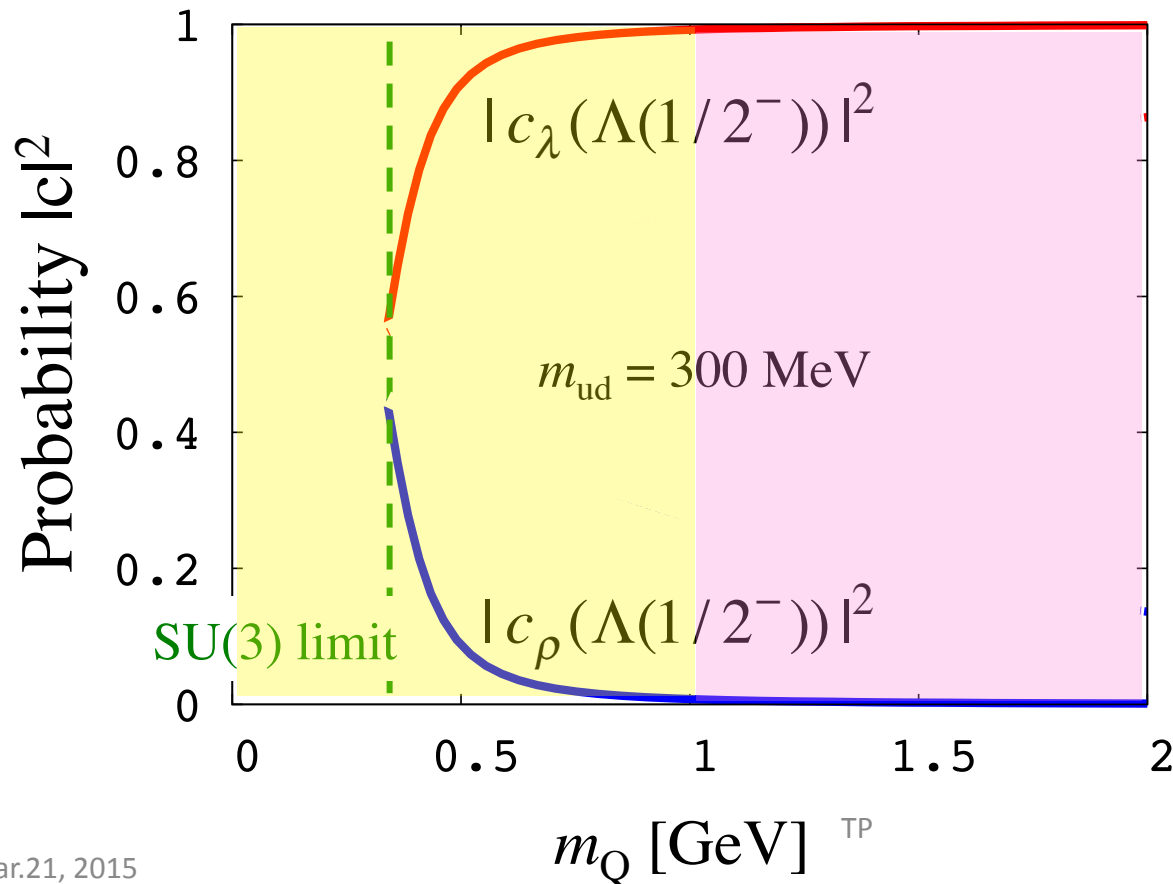
Excitation spectrum



Wave function

Mixing of $\Lambda(\text{phys}) = c_\lambda \Lambda(^2\lambda) + c_\rho \Lambda(^2\rho)$

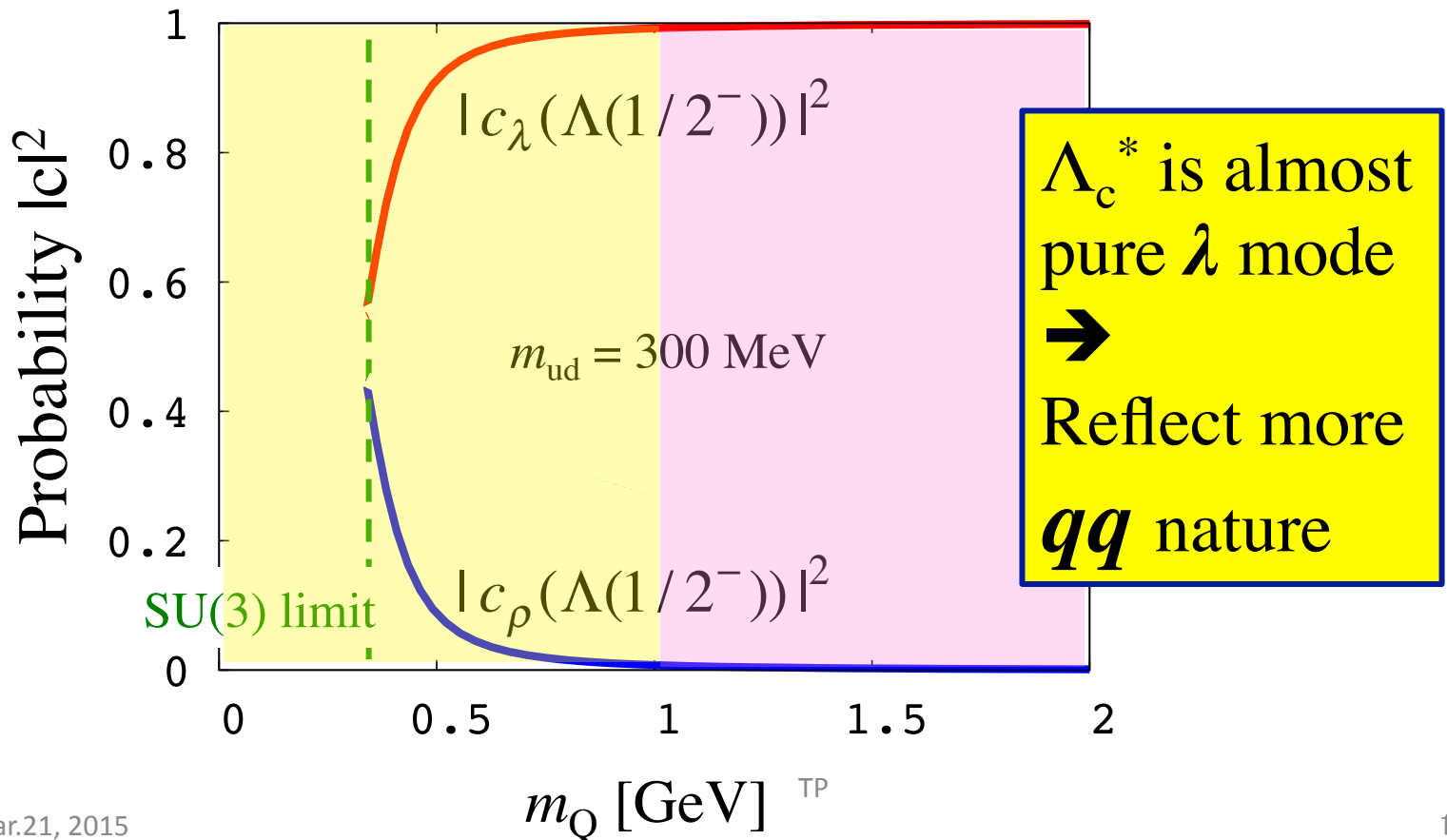
e.g. λ -mode dominant state: How much the other mode mixes?



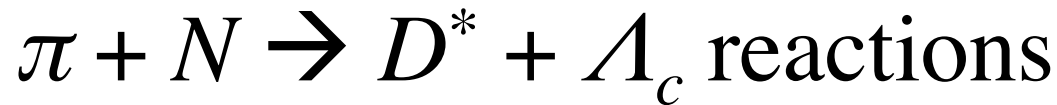
Wave function

Mixing of $\Lambda(\text{phys}) = c_\lambda \Lambda(^2\lambda) + c_\rho \Lambda(^2\rho)$

e.g. λ -mode dominant state: How much the other mode mixes?



3. Charmed baryon Productions



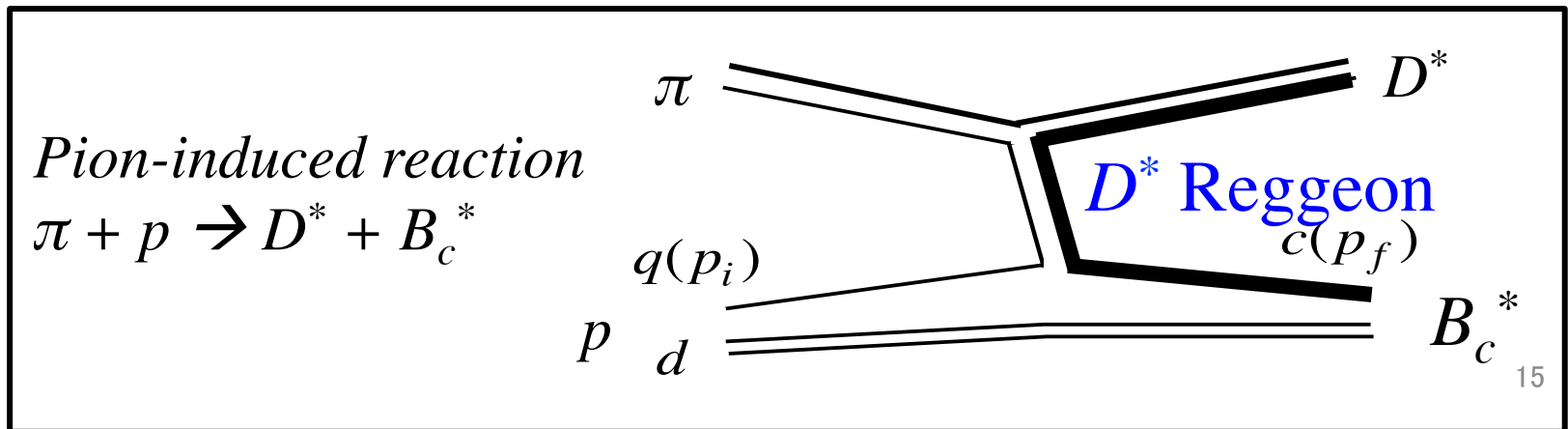
Production rate (Λ_c/Λ) and Ratios (B_c^*/B_c)

Strategy:

Forward peak (high energy) \rightarrow t-channel dominant

We look at:

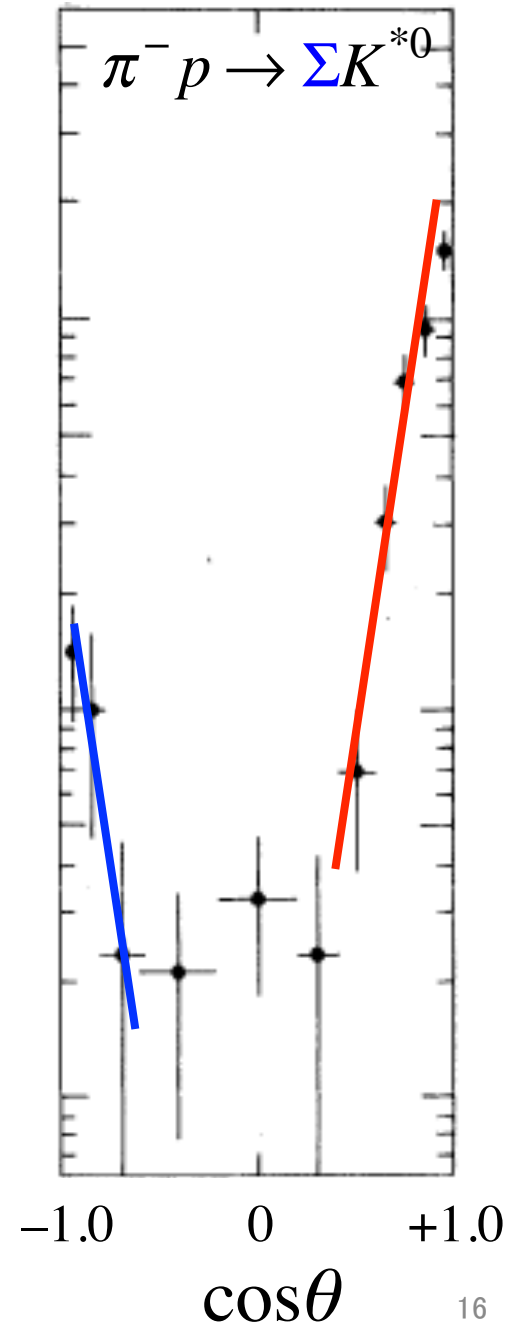
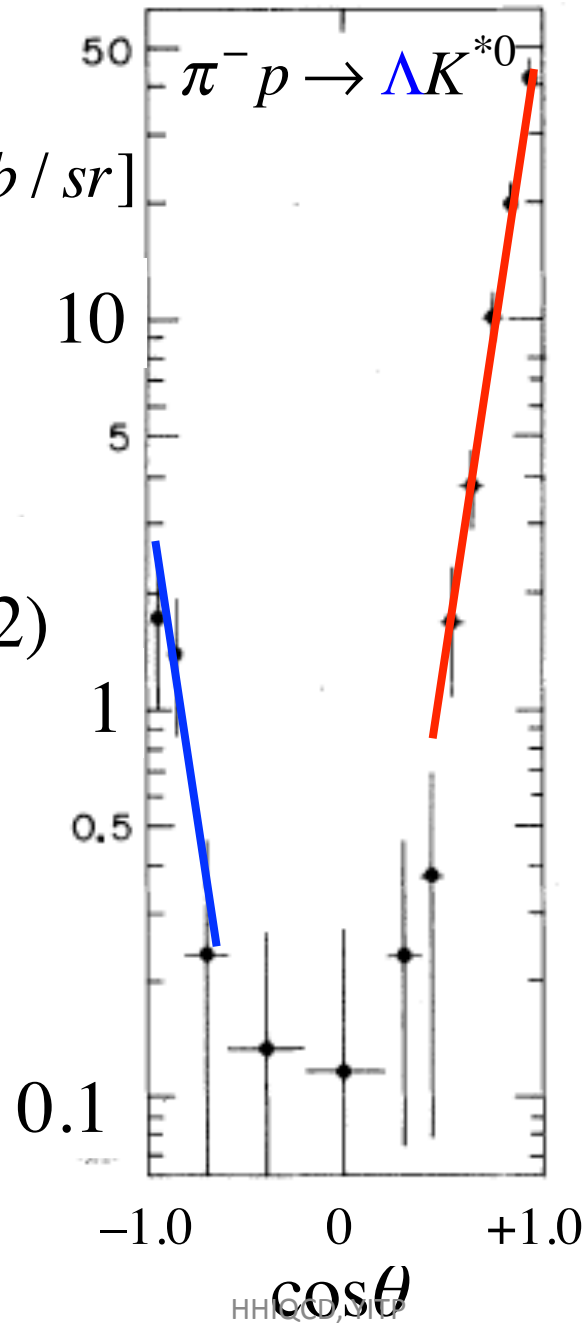
- Rates (Λ_c/Λ) by the Regge model, D^* Reggeon
- Relative ratios of B_c^* (λ modes) / B_c
by a one step process of Qd picture for λ -mode



$p_{\pi, \text{Lab}} = 4.5 \text{ GeV}$

D.J. Krennel et al
PRD6, 1220 (1972)

$$\frac{d\sigma}{d\Omega} [\mu\text{b} / \text{sr}]$$

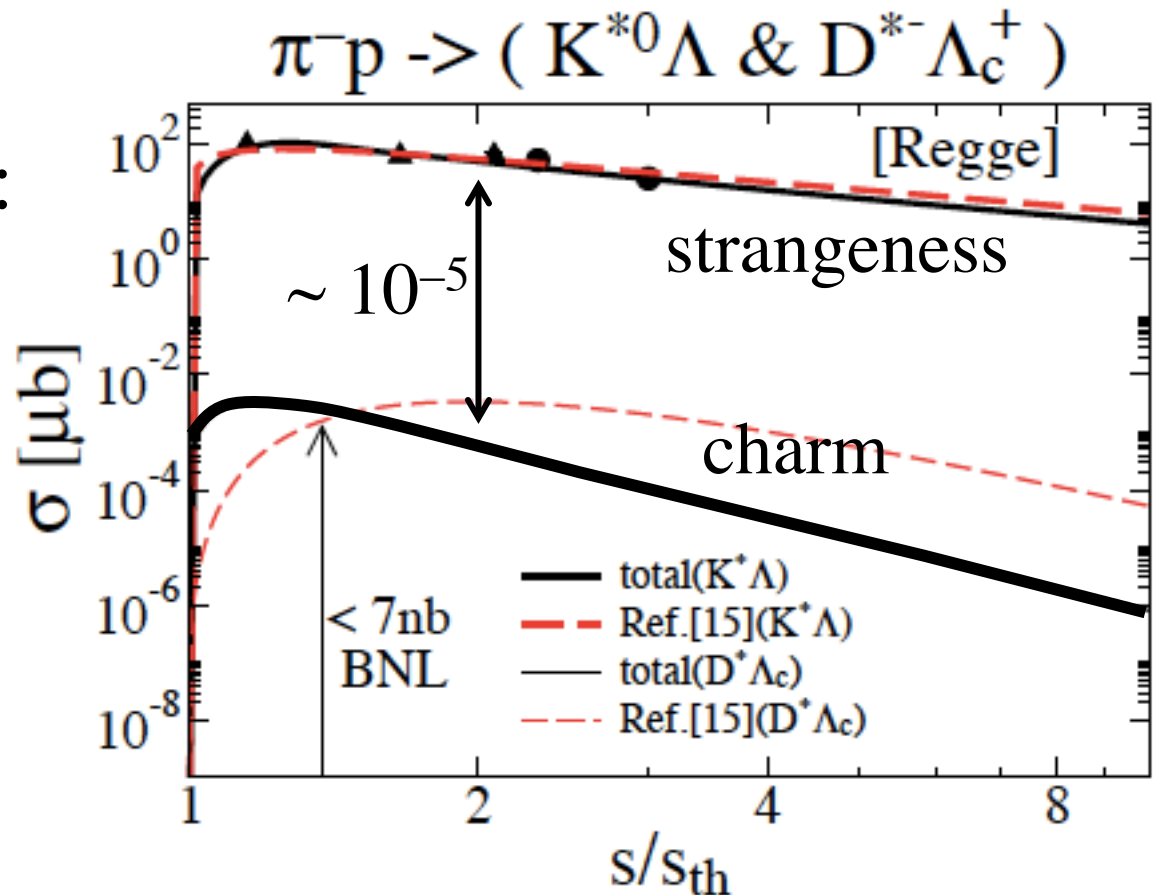


Absolute values

Regge model (Sang-Ho Kim, in preparation)

We have examined:

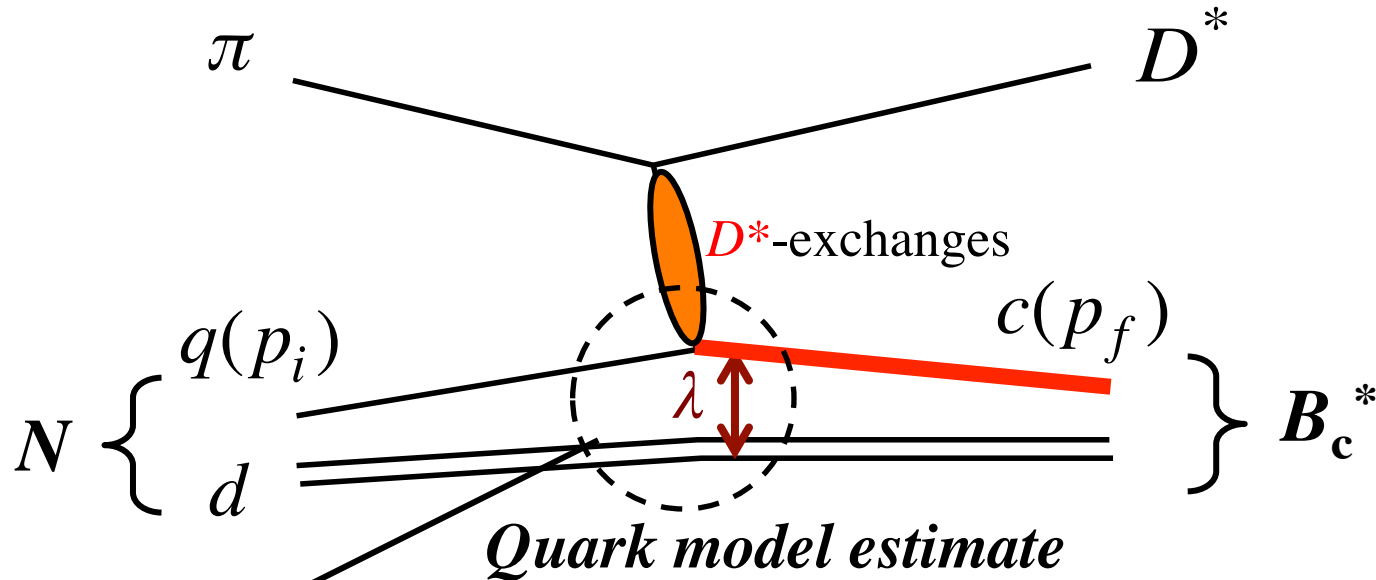
- K^* (strange) productions
- D and D^* Reggeon
- Angular dependence
- u -channels
 - ~ Baryon Regge
- Normalizations



Charm/strangeness ratio: $10^{-4} \sim 10^{-6}$
→ several nb near the threshold

Relative rates of (B_c^*/B_c)

One step process for Qd λ -mode



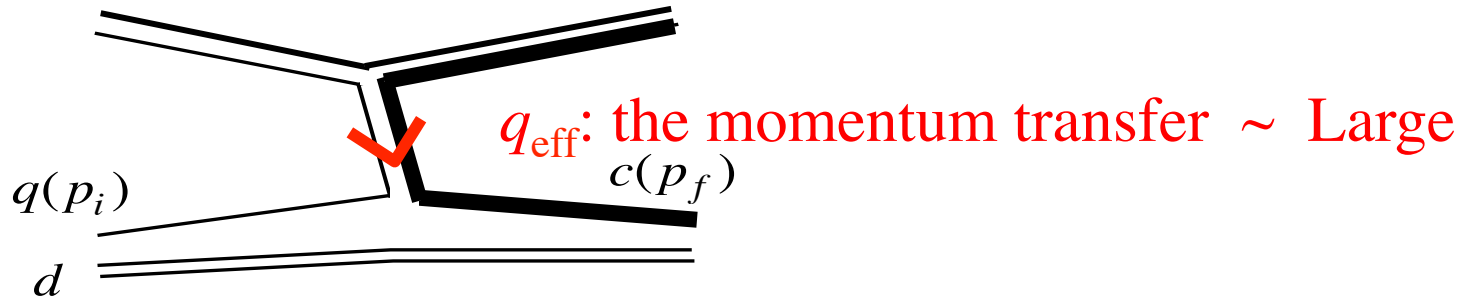
$$t_{fi} \sim \vec{k}_\pi \times \vec{e} \cdot \vec{J}_{fi}$$

$$\sim \langle B_c^* | \vec{e}_\perp \cdot \vec{\sigma} e^{i\vec{q}_{eff} \cdot \vec{x}} | N \rangle = (\text{Geometric}) \times (\text{Dynamic})$$

$D^* \sim \text{Transverse}$

CG coefficients

Dynamical part \sim radial integral



$$\text{GS} \quad \langle B_c(\text{S-wave}) | \vec{e}_\perp \cdot \vec{\sigma} e^{i\vec{q}_{\text{eff}} \cdot \vec{x}} | N(\text{S-wave}) \rangle_{\text{radial}} \sim 1 \times \exp\left(-\frac{q_{\text{eff}}^2}{4A^2}\right)$$

Excited states

$$\langle B_c(\text{P-wave}) | \vec{e}_\perp \cdot \vec{\sigma} e^{i\vec{q}_{\text{eff}} \cdot \vec{x}} | N(\text{S-wave}) \rangle_{\text{radial}} \sim \left(\frac{q_{\text{eff}}}{A}\right)^1 \times \exp\left(-\frac{q_{\text{eff}}^2}{4A^2}\right)$$

$$\langle B_c(\text{D-wave}) | \vec{e}_\perp \cdot \vec{\sigma} e^{i\vec{q}_{\text{eff}} \cdot \vec{x}} | N(\text{S-wave}) \rangle_{\text{radial}} \sim \left(\frac{q_{\text{eff}}}{A}\right)^2 \times \exp\left(-\frac{q_{\text{eff}}^2}{4A^2}\right)$$

Transitions to excited states are not suppressed

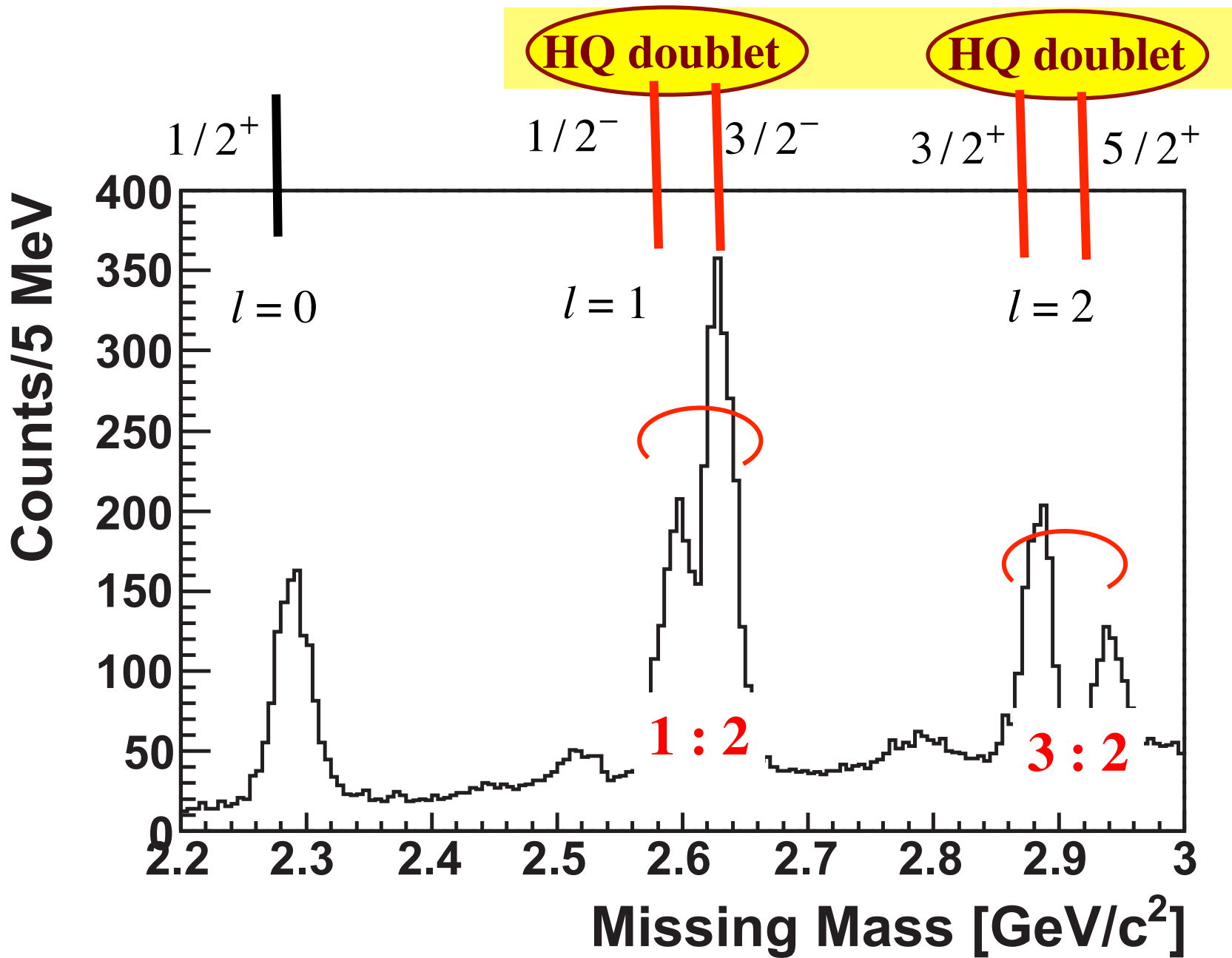
Results

Charm $k_{\pi}^{CM} = 2.71$ [GeV], $k_{\pi}^{Lab} = 16$ [GeV]

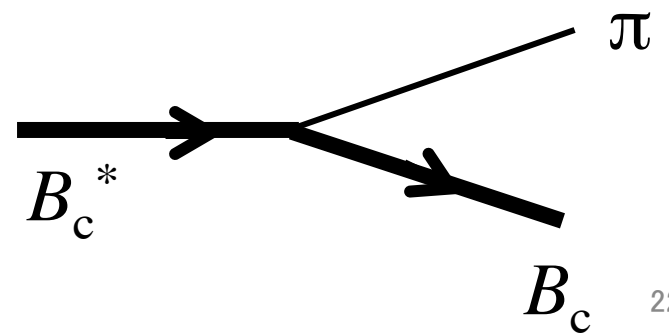
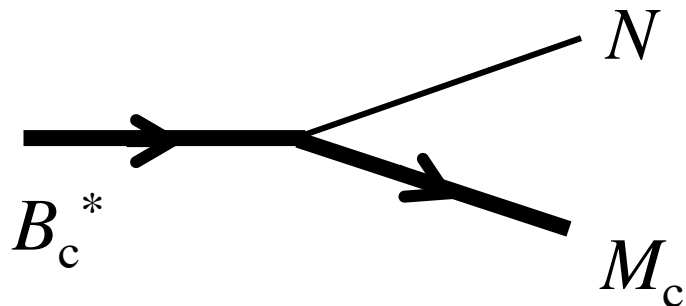
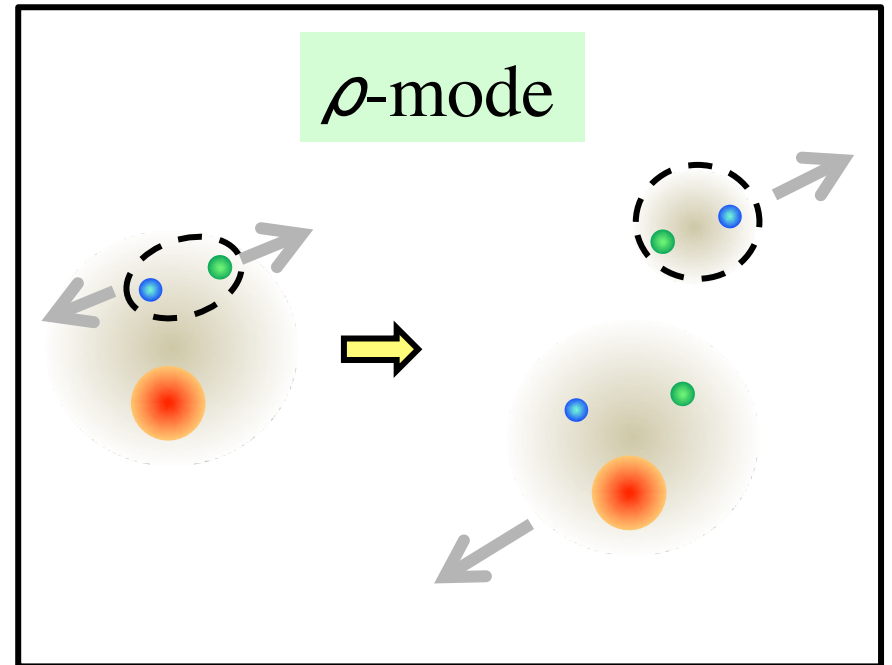
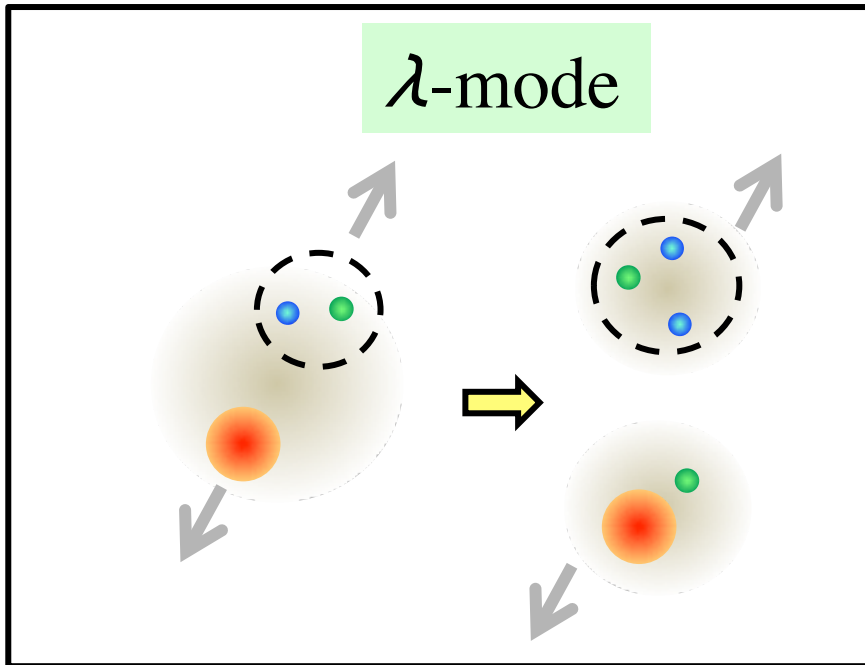
$l = 0$	$\Lambda_c(\frac{1}{2}^+)$ 1.00	$\Sigma_c(\frac{1}{2}^+)$ 0.02	$\Sigma_c(\frac{3}{2}^+)$ 0.16					
<u>$l = 1$</u>	$\Lambda_c(\frac{1}{2}^-)$ 0.90	$\Lambda_c(\frac{3}{2}^-)$ 1.70	$\Sigma_c(\frac{1}{2}^-)$ 0.02	$\Sigma_c(\frac{3}{2}^-)$ 0.03	$\Sigma'_c(\frac{1}{2}^-)$ 0.04	$\Sigma'_c(\frac{3}{2}^-)$ 0.19	$\Sigma'_c(\frac{5}{2}^-)$ 0.18	
<u>$l = 2$</u>	$\Lambda_c(\frac{3}{2}^+)$ 0.50	$\Lambda_c(\frac{5}{2}^+ -)$ 0.88	$\Sigma_c(\frac{3}{2}^+)$ 0.02	$\Sigma_c(\frac{5}{2}^+)$ 0.02	$\Sigma'_c(\frac{1}{2}^+)$ 0.01	$\Sigma'_c(\frac{3}{2}^+)$ 0.03	$\Sigma'_c(\frac{5}{2}^+)$ 0.07	$\Sigma'_c(\frac{5}{2}^+)$ 0.07

Strange $k_{\pi}^{CM} = 1.59$ [GeV], $k_{\pi}^{Lab} = 5.8$ [GeV]

$l = 0$	$\Lambda_-(\frac{1}{2}^+)$ 1.00	$\Sigma_-(\frac{1}{2}^+)$ 0.067	$\Sigma_-(\frac{3}{2}^+)$ 0.44					
$l = 1$	$\Lambda_-(\frac{1}{2}^-)$ 0.11	$\Lambda_-(\frac{3}{2}^-)$ 0.23	$\Sigma_-(\frac{1}{2}^-)$ 0.007	$\Sigma_-(\frac{3}{2}^-)$ 0.01	$\Sigma'_-(\frac{1}{2}^-)$ 0.01	$\Sigma'_-(\frac{3}{2}^-)$ 0.07	$\Sigma'_-(\frac{5}{2}^-)$ 0.067	
$l = 2$	$\Lambda_-(\frac{3}{2}^+)$ 0.13	$\Lambda_-(\frac{5}{2}^+ -)$ 0.20	$\Sigma_-(\frac{3}{2}^+)$ 0.007	$\Sigma_-(\frac{5}{2}^+)$ 0.01	$\Sigma'_-(\frac{1}{2}^+)$ 0.004	$\Sigma'_-(\frac{3}{2}^+)$ 0.02	$\Sigma'_-(\frac{5}{2}^+)$ 0.038	$\Sigma'_-(\frac{5}{2}^+)$ 0.04



4. Decays



Pion emission – quark model --on going

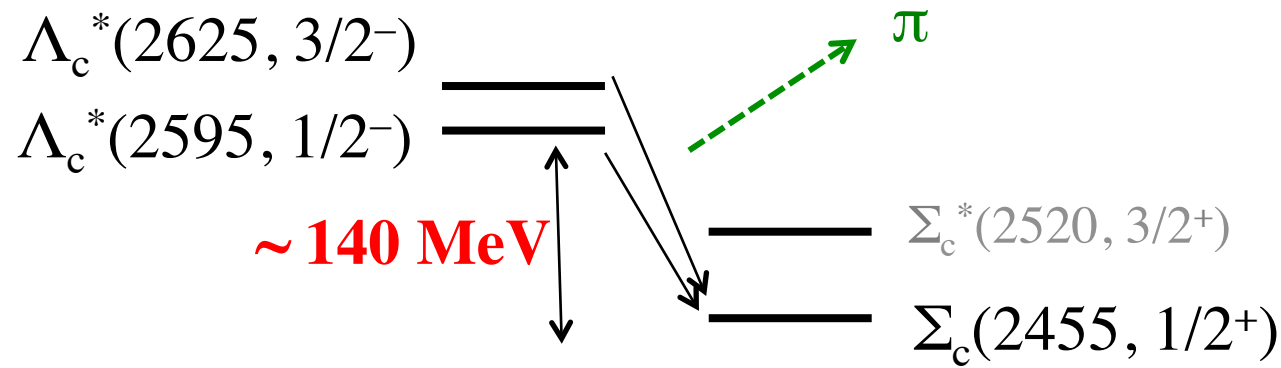
Things to be looked at:

- Pion emission \sim **very near the threshold**
- D-meson emission \sim The coupling of D -meson to cq

Pion emission – quark model --on going

Things to be looked at:

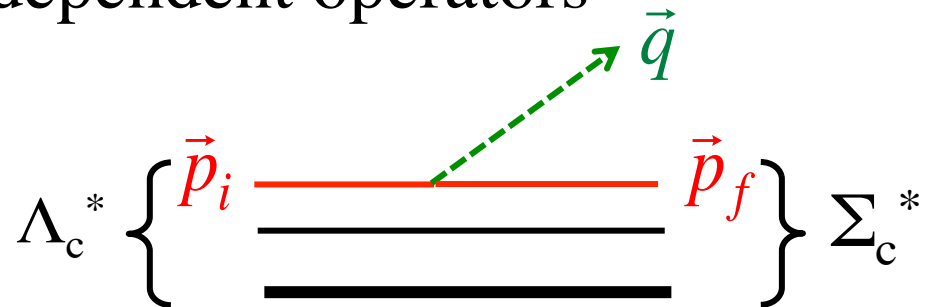
- Pion emission \sim **very near the threshold**



Place to look at the two independent operators

$$\bar{q}\gamma_5 q \phi_\pi, \bar{q}\gamma^\mu \gamma_5 q \partial_\mu \phi_\pi$$

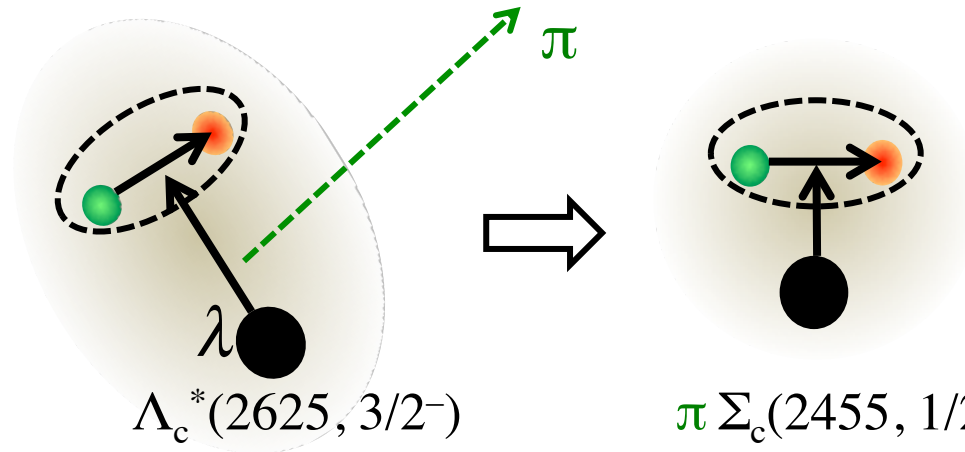
$$\vec{\sigma} \cdot \vec{p}_i, \vec{\sigma} \cdot \vec{p}_f \quad (\vec{\sigma} \cdot \vec{q})$$



Pion emission – quark model --on going

Diquark correlation may suppress pion emission

λ -mode decay



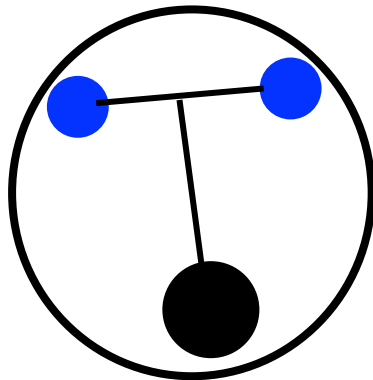
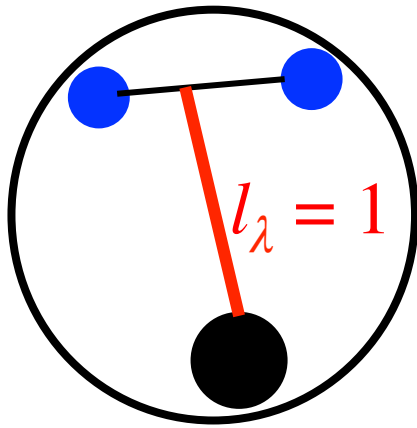
overlap vanishes
with strong correlation

$$\begin{aligned}
 A_h &= \frac{q}{m} \int d^3 X e^{-i(P-q-k)X} \times \int d^3 \rho e^{-i\vec{q}_\rho \cdot \vec{\rho}} \langle \psi'_0(\vec{\rho}) | \psi_0(\vec{\rho}) \rangle \\
 &\times \int d^3 \lambda \langle [[[\psi_0(\vec{\lambda}), d^1]^1, \chi^c]_h^{J_f} | e^{-i\vec{q}\lambda \cdot \vec{\lambda}} \sigma_3 | [[[\psi_1(\vec{\lambda}), d^0]^1, \chi^c]_h^{J_i} \rangle \\
 &\quad \Sigma_c(2455, 1/2^+) + \pi \leftarrow \Lambda_c^*(2625, 3/2^-) \\
 &\quad \text{Pion emission by the } \lambda \text{ motion}
 \end{aligned}$$

Radiative decay: $1/2^- \rightarrow 1/2^+$ E1

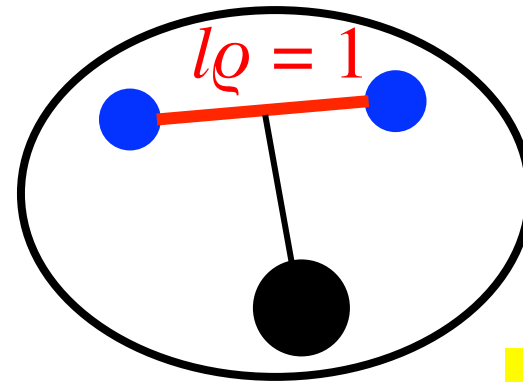
λ mode

Good diquark 0^+

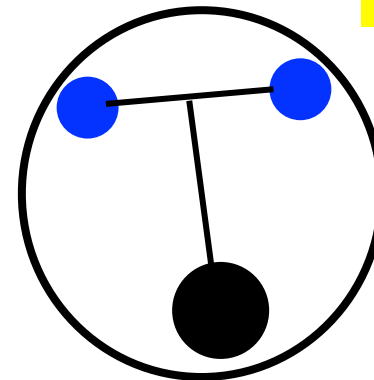


ρ mode

3P0 diquark 0^-



Good diquark 0^+

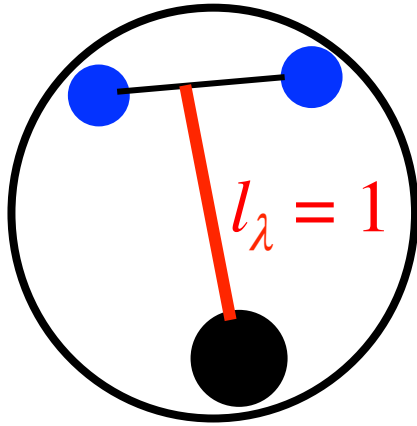


$0^- \rightarrow 0^+$ is
forbidden

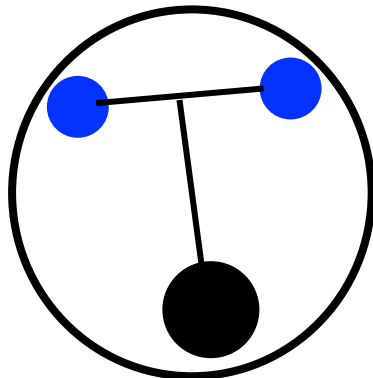
Radiative decay: $5/2^- \rightarrow 1/2^+$ M2, E3

λ mode

3S_1 diquark 1^+



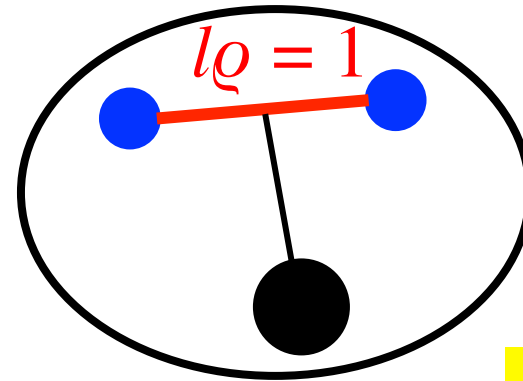
Both M2 E3



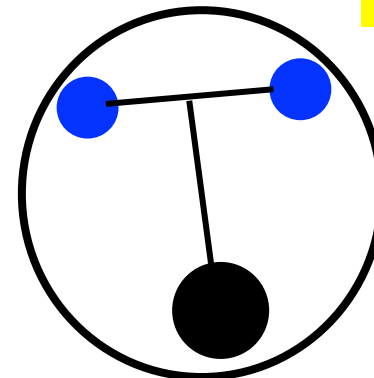
Good diquark 0^+

ρ mode

3P_2 diquark 2^-



$2^- \rightarrow 0^+$ is
only M2



Summary

- Charmed baryons: yet unexplored states
- J-PARC plans to observe them
- ρ and λ modes are distinct (Isotope shift)
better chance to look at in Λ^* than in Σ^* \rightarrow diquarks
- Systematic study in strangeness is important
- Higher excited (Λ) states may be produced a lot
- Pion decay provides unique opportunity
- Radiative decay selection rules are also helpful


The 10th International Workshop on the Physics of Excited Nucleons

NSTAR 2015



May 25(Mon)-28(Thu), 2015

Ichō Kaikan Suita Campus,
Osaka University
Osaka, Japan

- 
- Complete experiments of meson-photoproduction reactions for N^*
 - Reaction models and PWA for resonance parameters
 - Dynamical models and coupled channel analysis
 - Baryon (hadron) structure, form factors, and transition amplitudes
 - Chiral symmetry and baryon resonances
 - Baryons (hadrons) from lattice QCD and new approaches
 - Facilities and future projects
 - Other topics related to N^* physics