

Λ - Λ Interaction from Relativistic Heavy Ion Collisions

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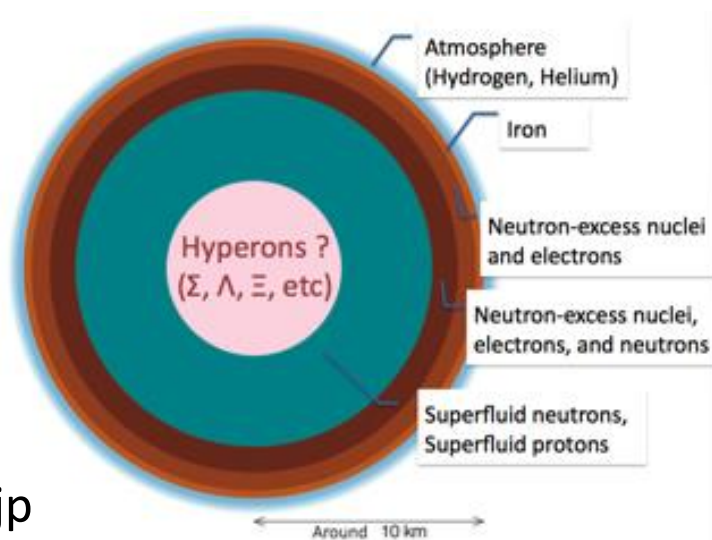
Outline

Goal : $\Lambda\Lambda$ interaction from HIC

- 1. Introduction**
- 2. Effect of $\Lambda\Lambda$ interaction (w/ simplest source)**
- 3. +Collective Expansion**
 - 1. Spherical source**
 - 2. Bjorken and transverse expansion**
- 4. Feed-down Contribution**
- 5. Residual Correlation**
- 6. Summary**

Role of $\Lambda\Lambda$ Interaction

Possible Emergence of Hyperons in NS core



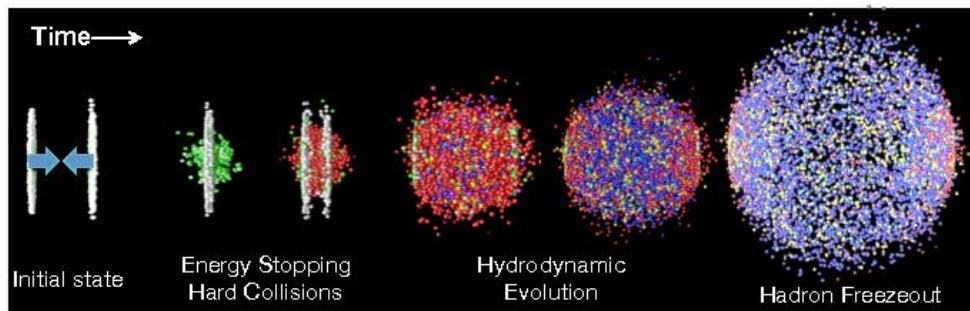
from KEK.jp

To understand EoS,
Information on
Hyperon-Hyperon
Interaction is
indispensable

H-dibaryon ($uuddss$)?

- Λ - Λ bound state due to strong attraction?
- Resonance?

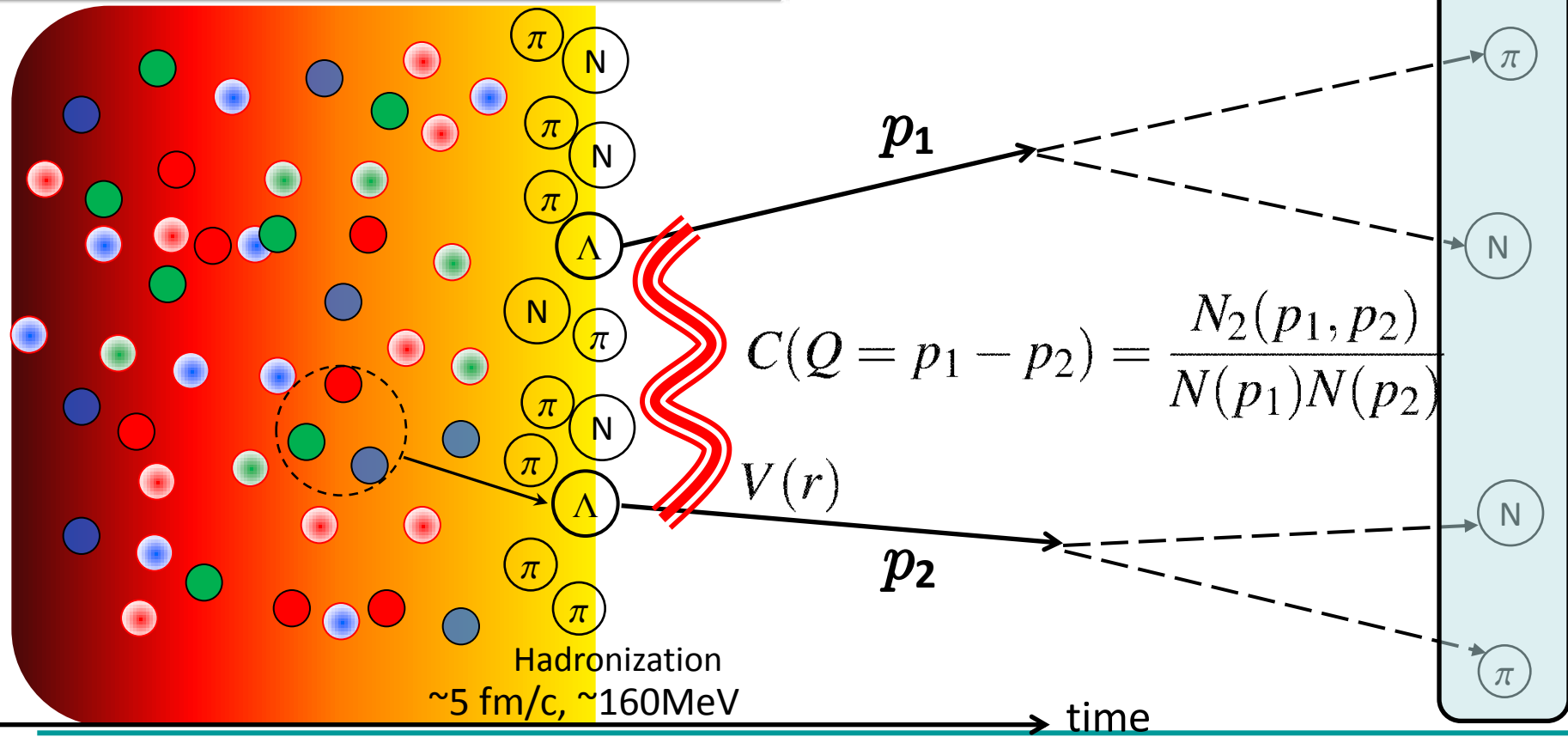
RHIC Can Tell Us about $\Lambda\Lambda$ interaction



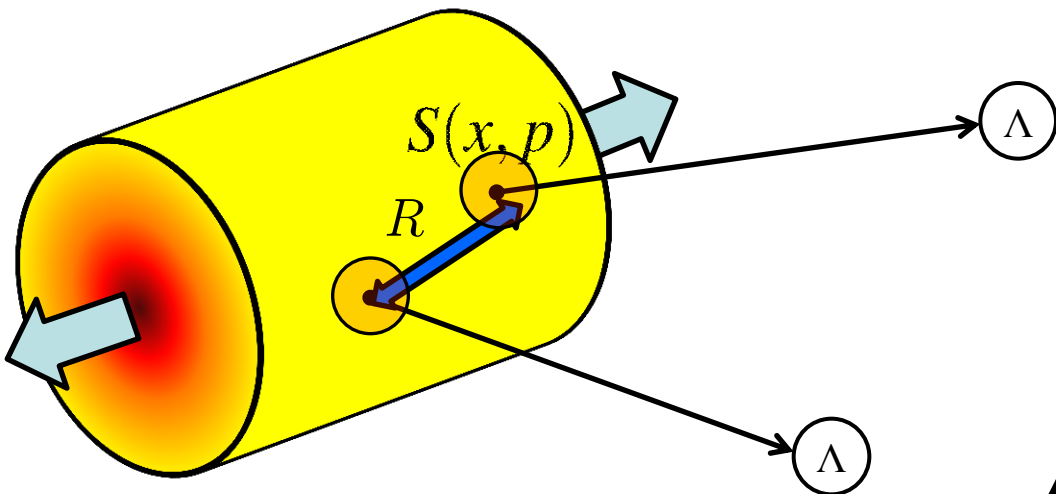
Au+Au 200A GeV

$$\frac{dN_\Lambda}{dy} \simeq 0.6 - 13 \quad \times 10^8 \text{ events}$$

(60-80%) (0-5%)



$\Lambda\Lambda$ Correlation in HIC



Independent (Chaotic) emission
(\leftarrow Thermal Source)

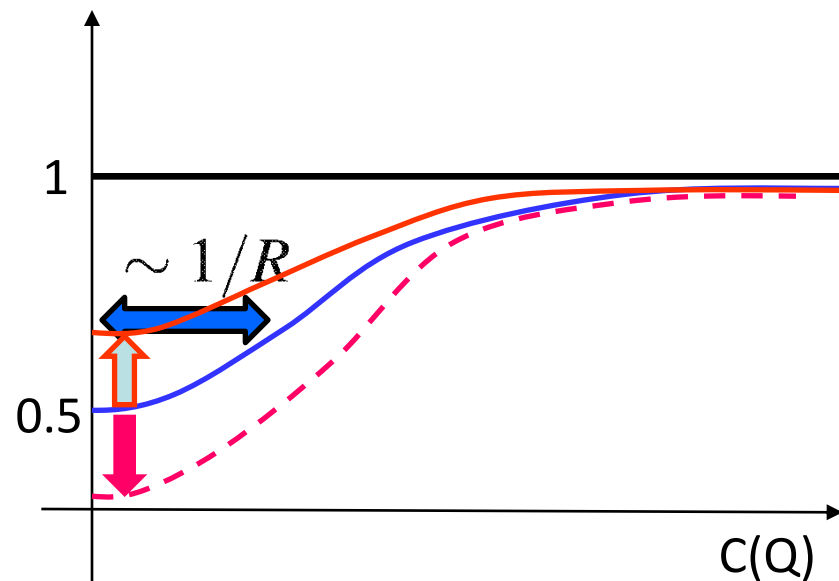
Identical particle correlation from
quantum statistics (HBT effect)

$C(Q)$: effective source size

$\Lambda\Lambda$ Interaction : No Coulomb!

Affects $C(Q)$ when effective range
 r_{eff} is comparable with R

Different results for repulsive and
attractive interaction



Approach : Thermal Source + $V_{\Lambda\Lambda}$

Formula (Gong et al., '91)

$$C_2(Q, K) = \frac{\int d^4x_1 d^4x_2 S(x_1, K) S(x_2, K) |\Psi_{12}(Q, x_1 - x_2 - (t_2 - t_1)K/m)|^2}{\int d^4x_1 d^4x_2 S(x_1, k_1) S(x_1, k_2)}$$

Emission source func.



Thermal source model
(Mimic hydro)

$\Lambda\Lambda$ relative S-wave func.



Various potentials (via 2 or 3 range
Gaussian Fit)

- ◆ Static Spherically Symmetric
- ◆ Spherically Symmetric + Hubble Flow
- ◆ Cylindrically Symmetric + Boost-invariance + Transverse flow

Meson exchange models (Nijmegen model D, F, Soft Core89/97, ESC08)

Phenomenological (Ehime)

Quark model (fss2)

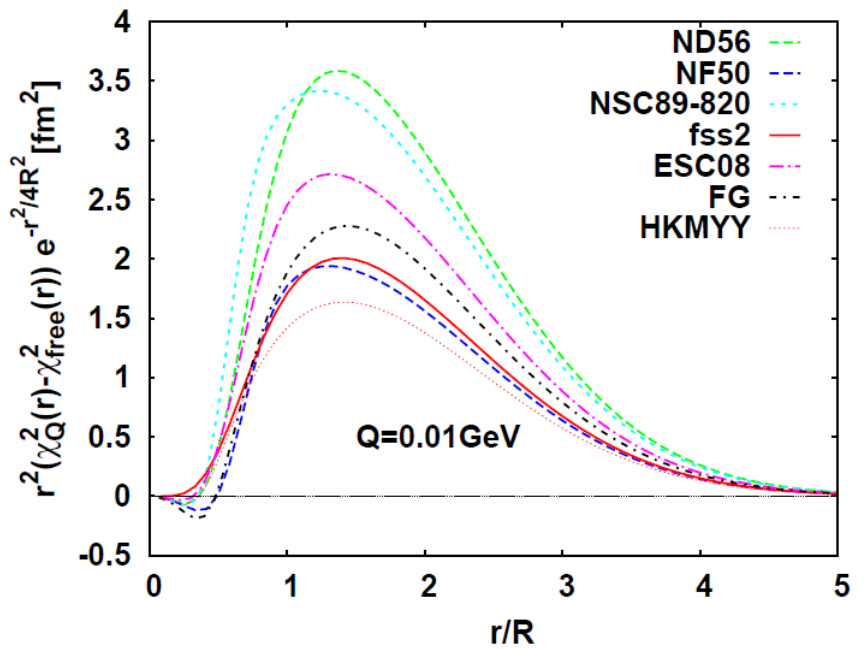
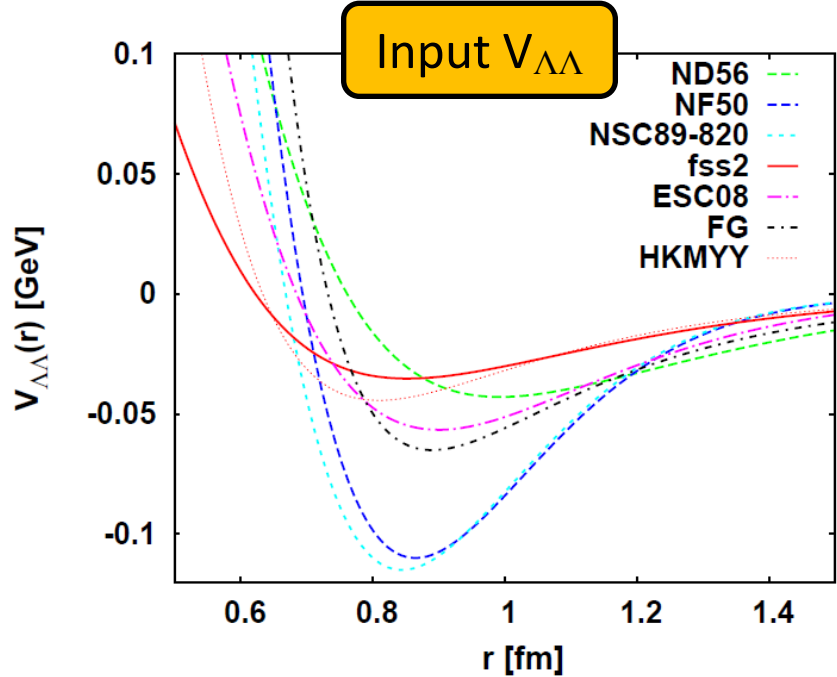
Fit to $_{\Lambda\Lambda}^6\text{He}$ (Nagara) Filikhin-Gal (FG)
Hiyama et al. (HKMY)

Potential, Wave func., / Correlation

Correlation Function for the Static Source

$$C_{\text{stat}}(Q) = 1 - \frac{1}{2} e^{-Q^2 R^2} + \frac{1}{4\sqrt{\pi}R^3} \int_0^\infty dr r^2 e^{-\frac{r^2}{4R^2}} \underbrace{[\chi_Q(r)]^2 - [j_0(Qr/2)]^2}_{\text{Interaction : deviation from free w.f.}}$$

HBT (size R) S-wave w.f.

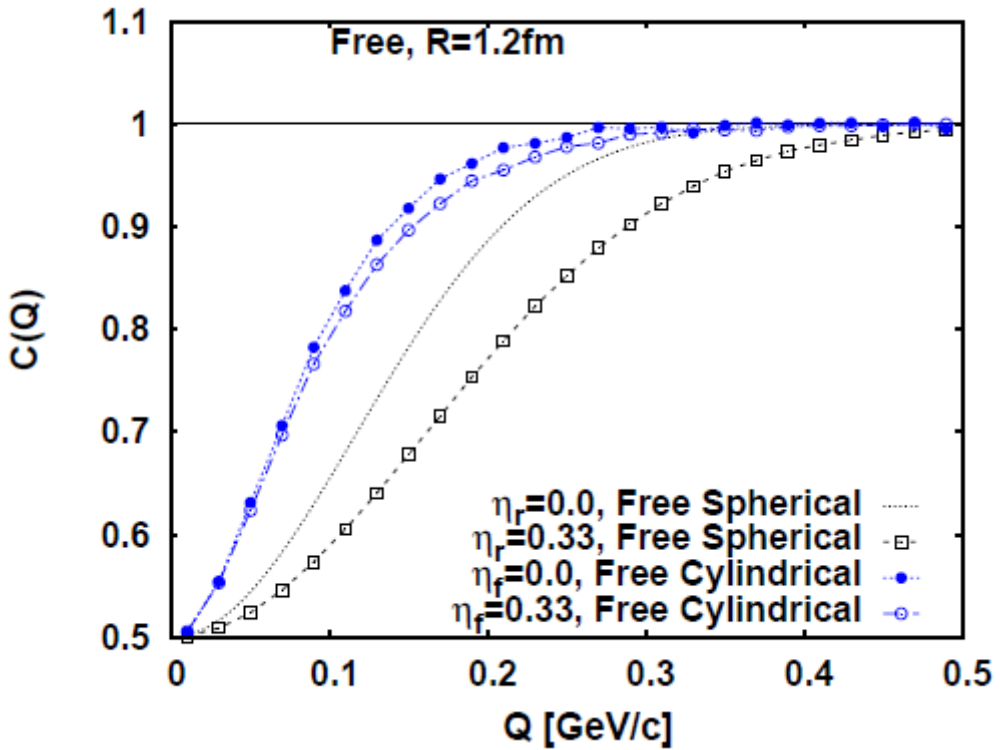


Expanding Source Model

$$S(x, k) \propto m_T \cosh(y - Y_L) n_F(u \cdot k / T) \exp \left[-\frac{(\tau - \tau_0)}{2(\Delta\tau)^2} - \frac{x^2 + y^2}{2R^2} \right]$$

$$\sim \exp \left[-\frac{\gamma_T M_T}{T} \cosh(y - Y_L) \right] \exp \left[\frac{\gamma v_T k_T \cos \phi}{T} \right]$$

Give a finite longitudinal extent $R_L \sim \tau_0 \sqrt{\frac{T}{M_T}}$

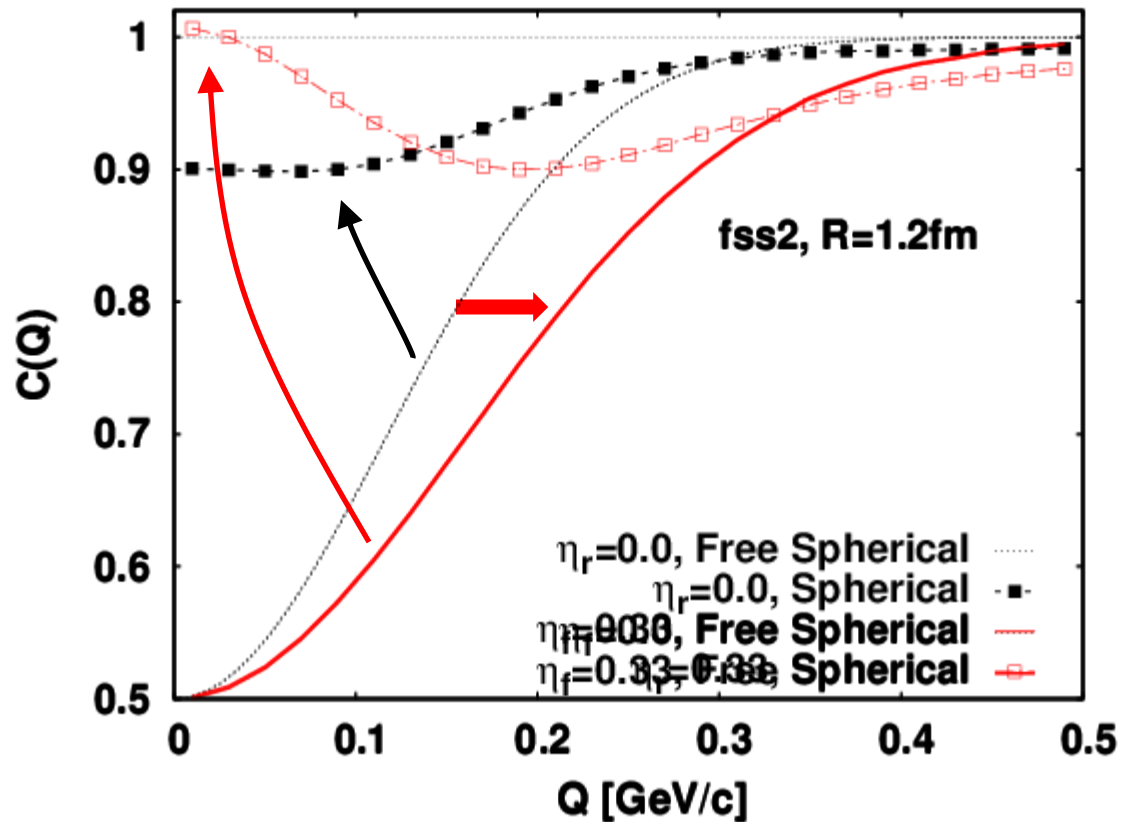


Width of C(Q) ($Q = |p_1 - p_2|$): effective 3d size

Large R_L : narrow width

Transverse flow: fit to p_t distribution of Λ (STAR '12)

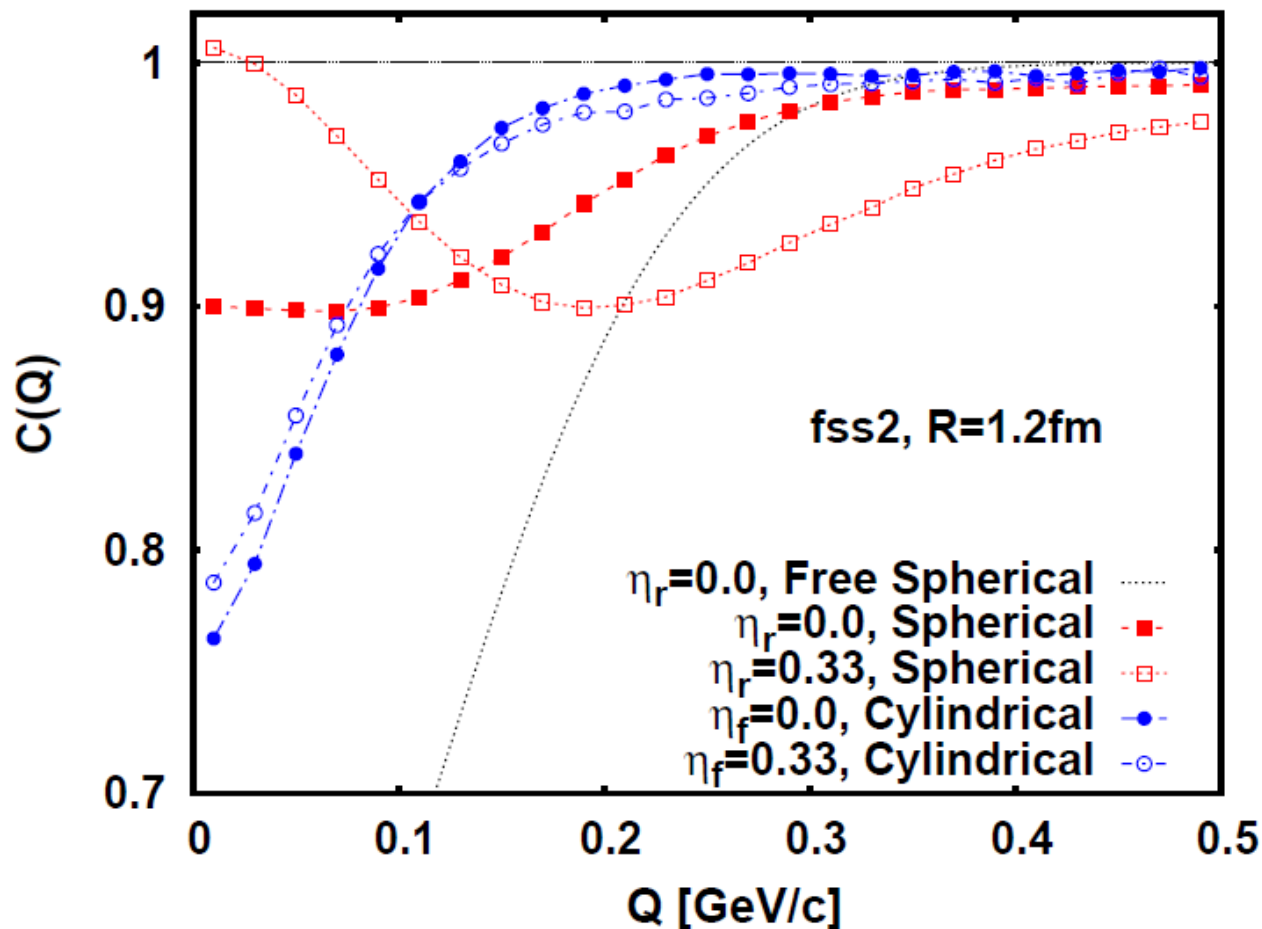
Illustration : Spherical Source



Expansion makes the apparent source size smaller

Attraction in C(Q) becomes stronger

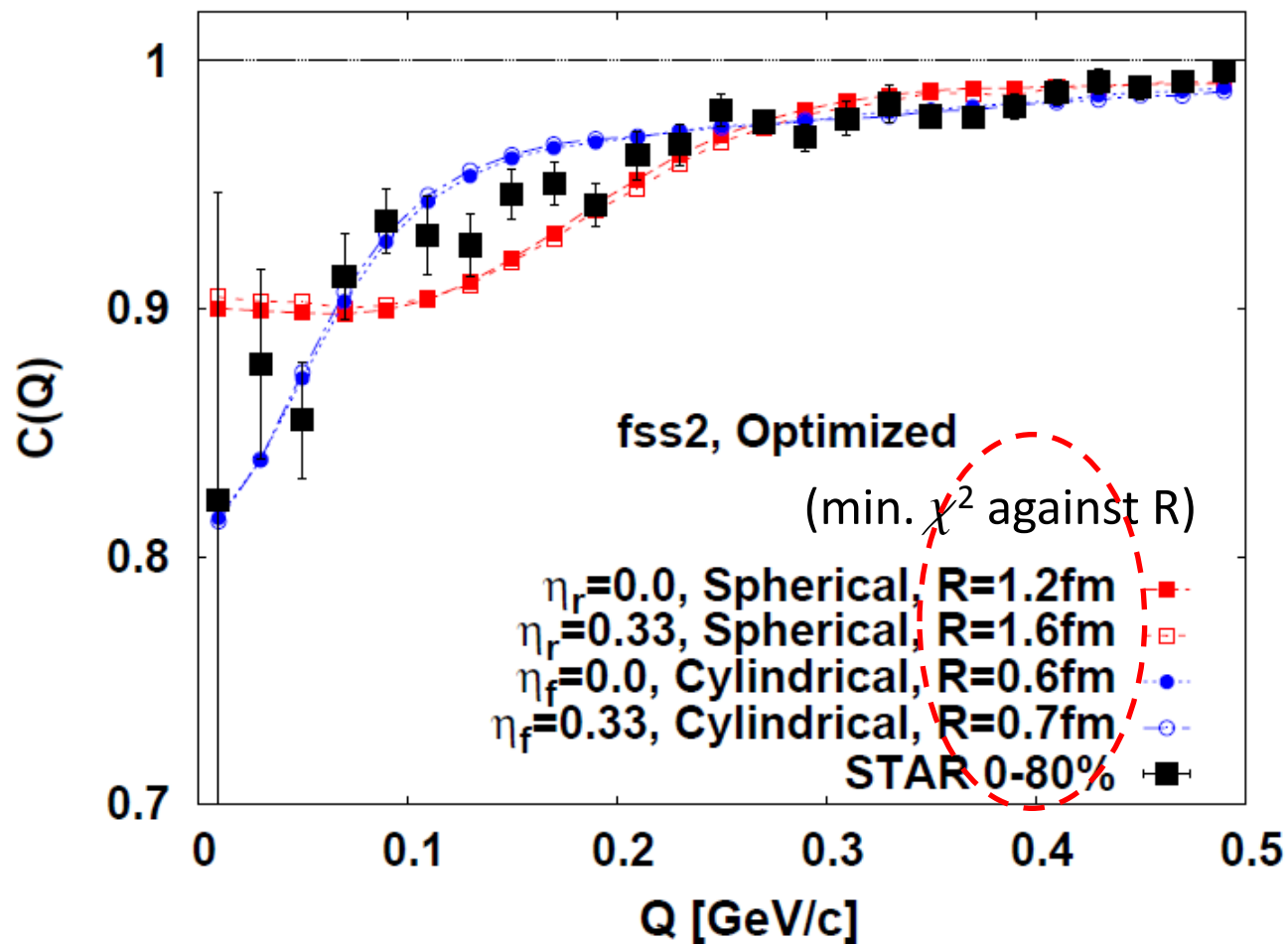
Results from Boost-invariant Source



Effect of η_f is rather small : longitudinal expansion dominates $C(Q)$

Behavior at small Q is different from the static source !

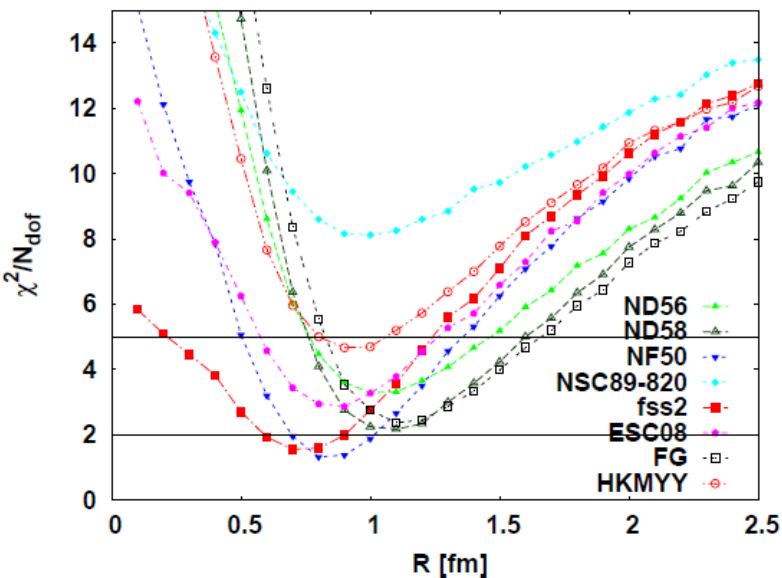
Combined effects from flow and $V_{\Lambda\Lambda}$



Effect of η_f is absorbed into change of R_{opt}

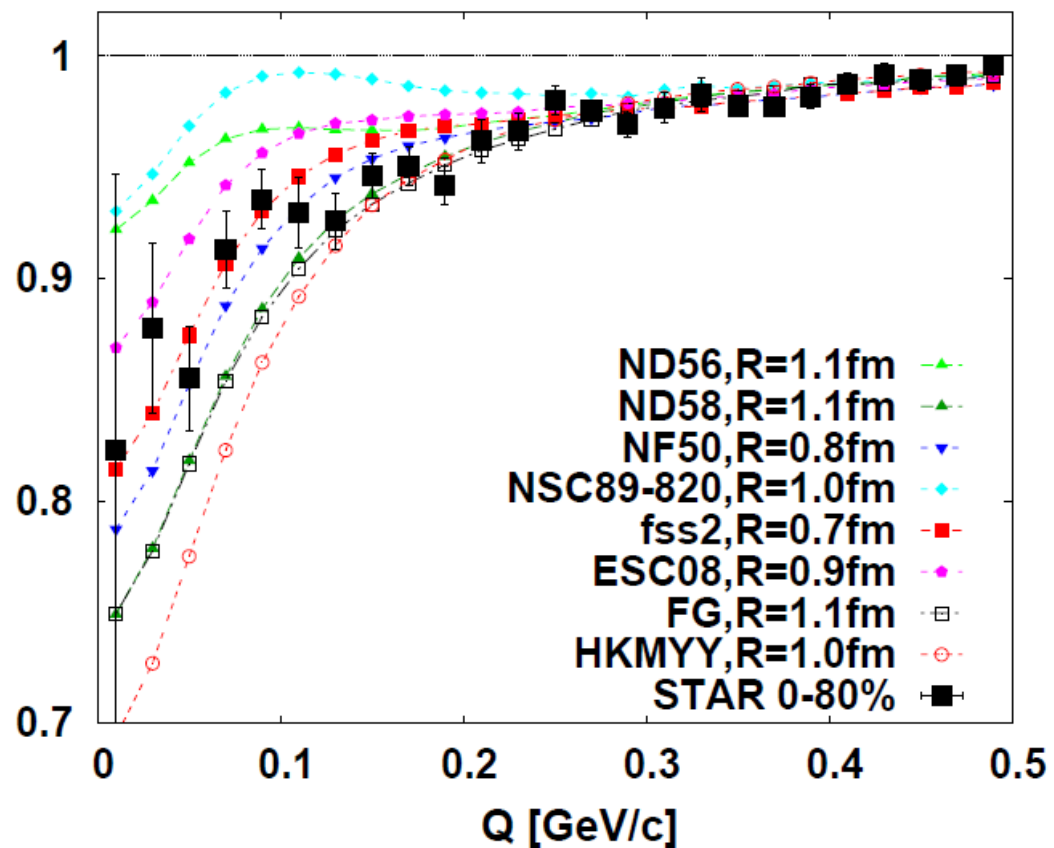
Longitudinal expansion gives another type of “best fit”

$V_{\Lambda\Lambda}$ from Expanding Source Model



Same $V_{\Lambda\Lambda}$ are chosen with *smaller R!*

$C(Q)$



Feed-Down Contribution

Assume Ξ is not included

DCA < 0.4cm from the primary vertex

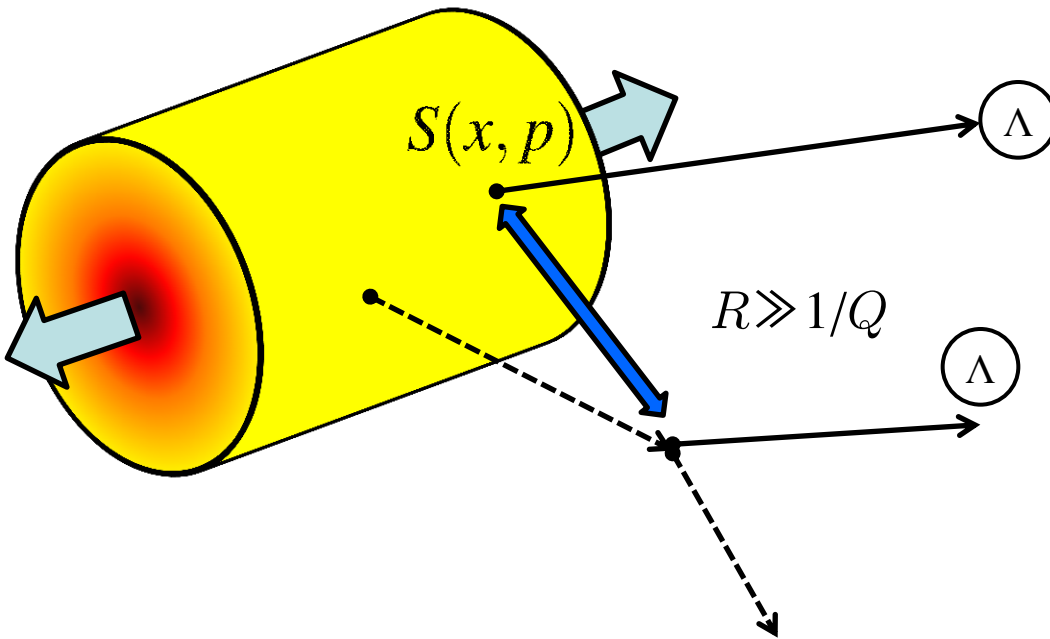
$\Sigma^0 \rightarrow \Lambda + \gamma$

$$C(Q) \rightarrow 1 + \left(\frac{\Lambda^{\text{dir}}}{\Lambda^{\text{tot}}} \right)^2 (C(Q) - 1)$$



$$(0.67)^2$$

consistent w/ thermal model



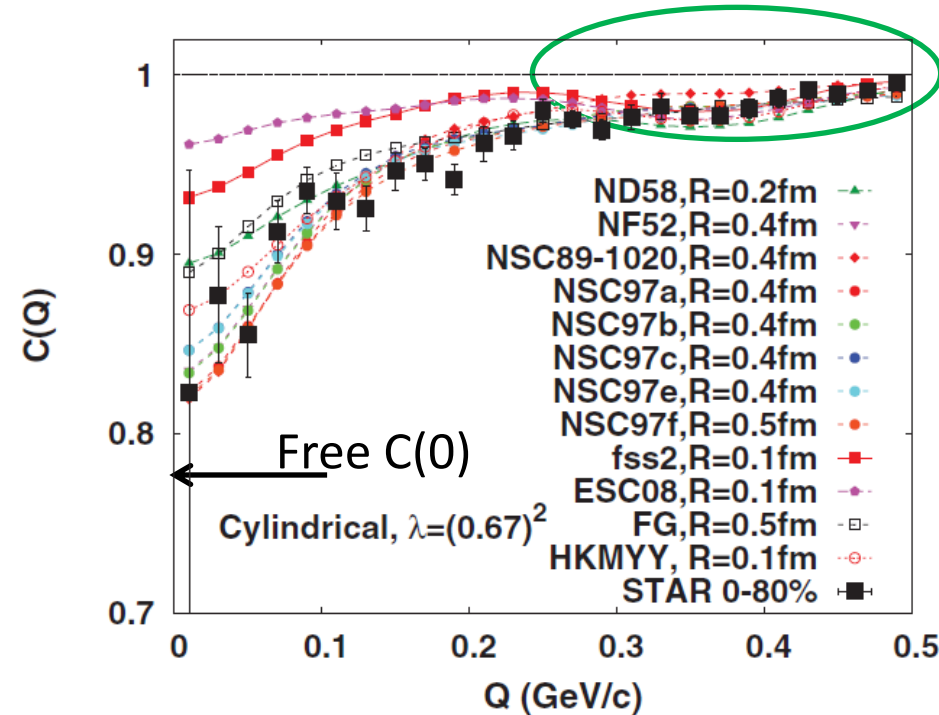
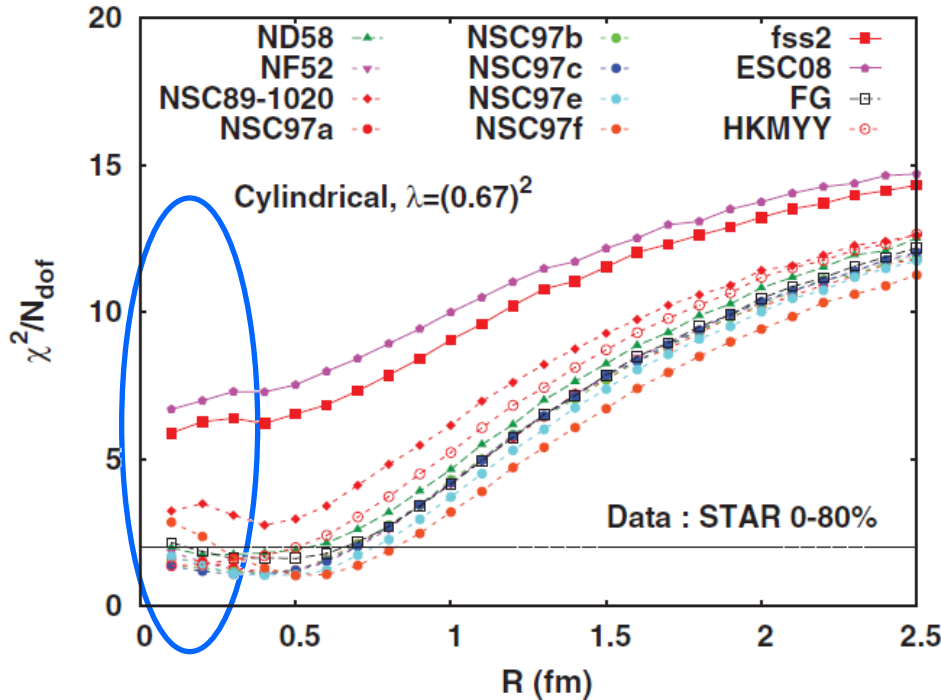
Using $\Sigma^0/\Lambda = 0.278$ (p+Be data) and $\Xi/\Lambda = 0.85$ (RHIC), $\Lambda^{\text{dir}}/\Lambda^{\text{tot}} = 0.67$

(0.52 if including Ξ)

Feed-Down Effect

$$C(Q) \rightarrow 1 + (0.67)^2 (C(Q) - 1)$$

 Sensitivity at low Q is reduced



 **Unphysically small size** is preferred;

due to the **long tail** in $C(Q)$ which cannot be included in the present framework

Long Tail : Residual Correlation?

Assuming a short-range correlation

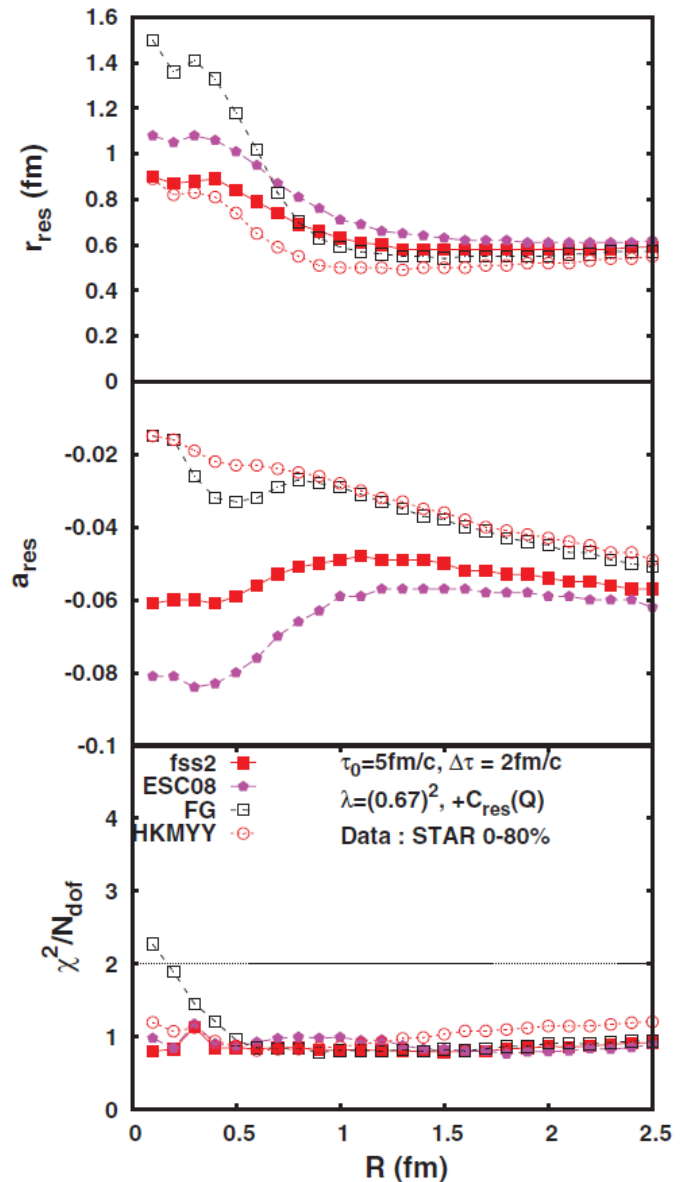
$$C(Q) \rightarrow C(Q) + a_{res} e^{-r_{res}^2 Q^2}$$

Search for min χ^2 in (a_{res}, r_{res}) for each R

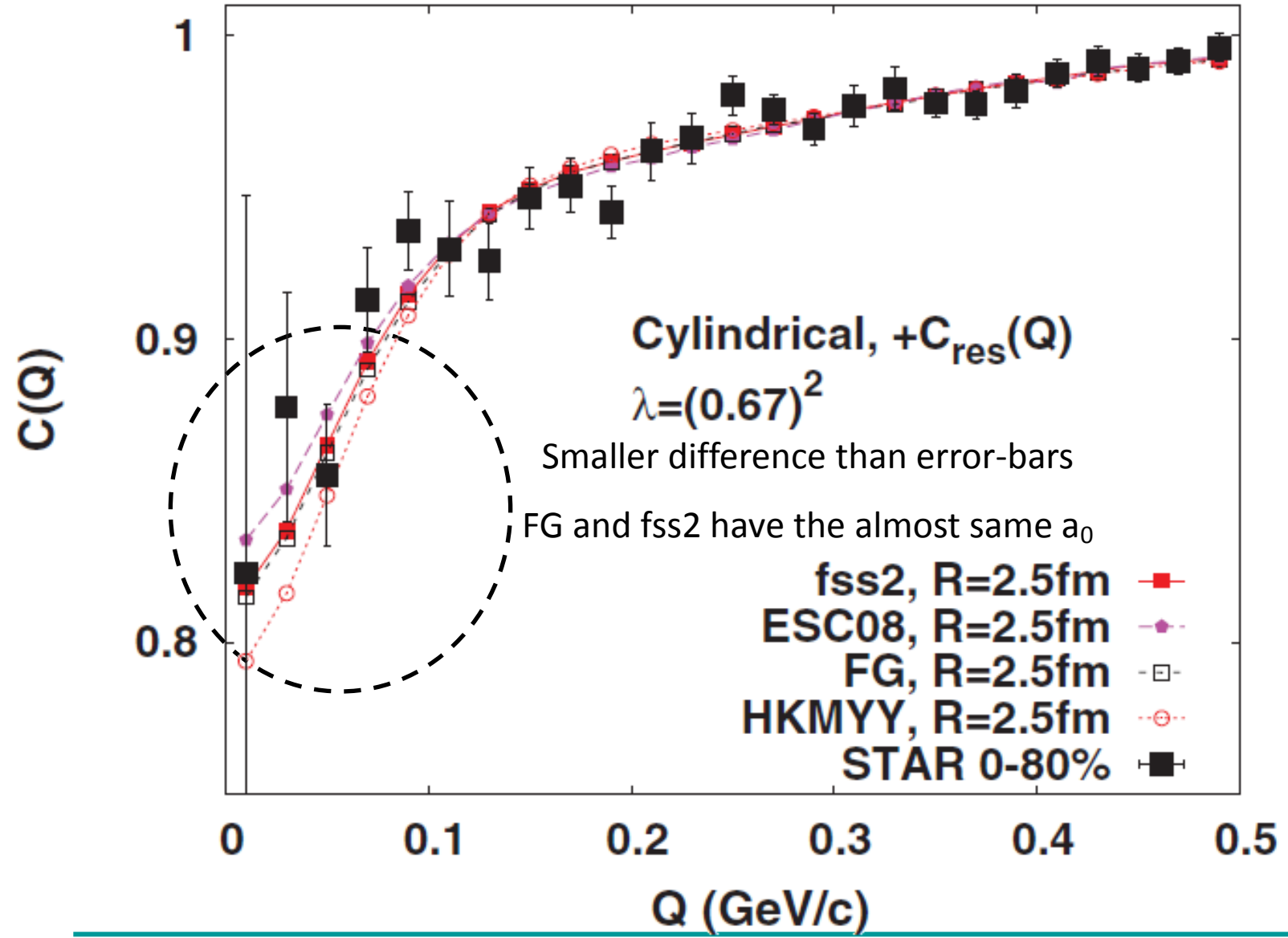
$\chi^2/N_{dof} \sim 1$ for shown R range

Small R : Large $r_{res} > R$

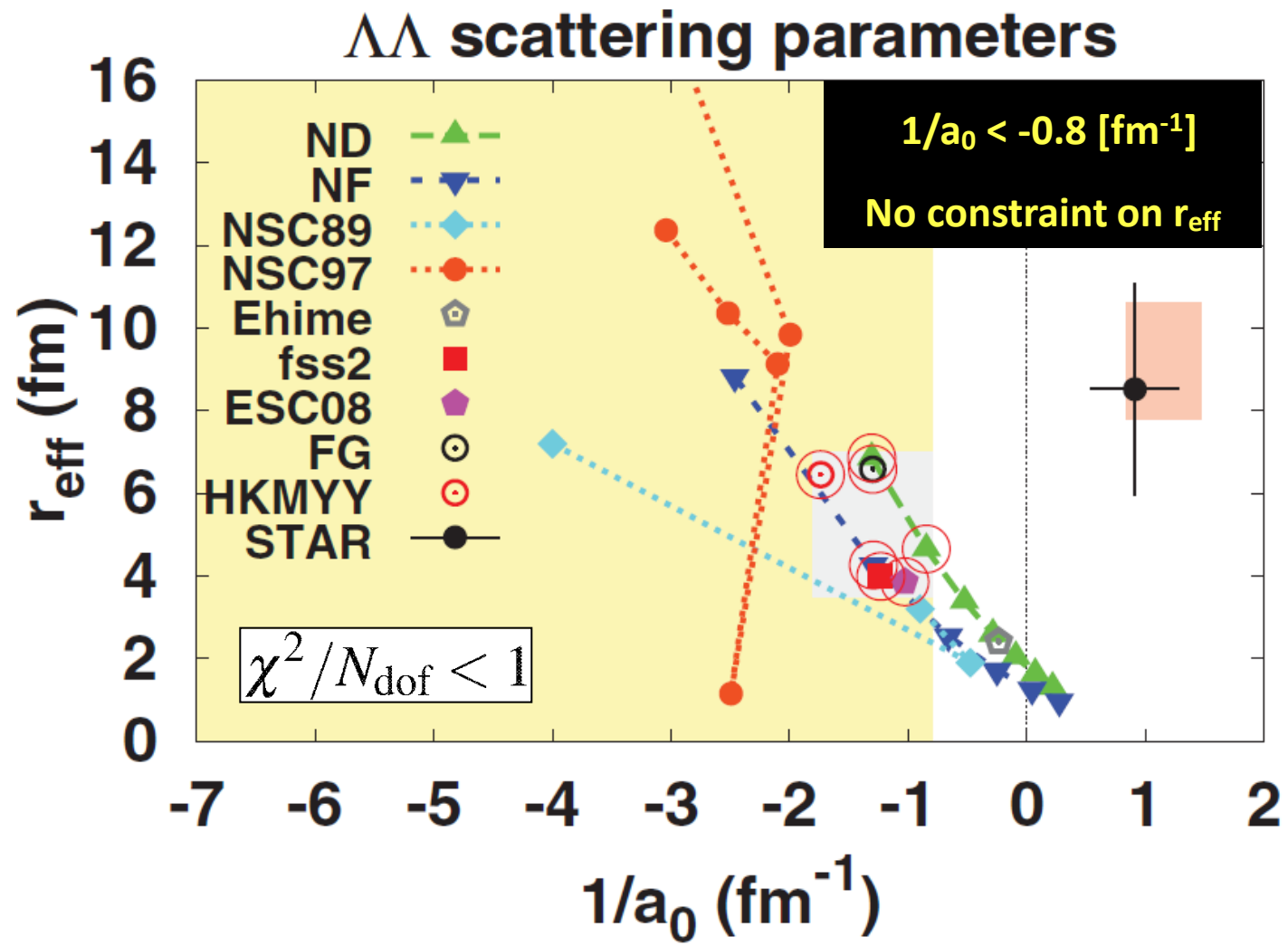
Low Q behavior \longleftrightarrow $\Lambda\Lambda$ Interaction



An Example for Agreement



Constraints on a_0 and r_{eff}



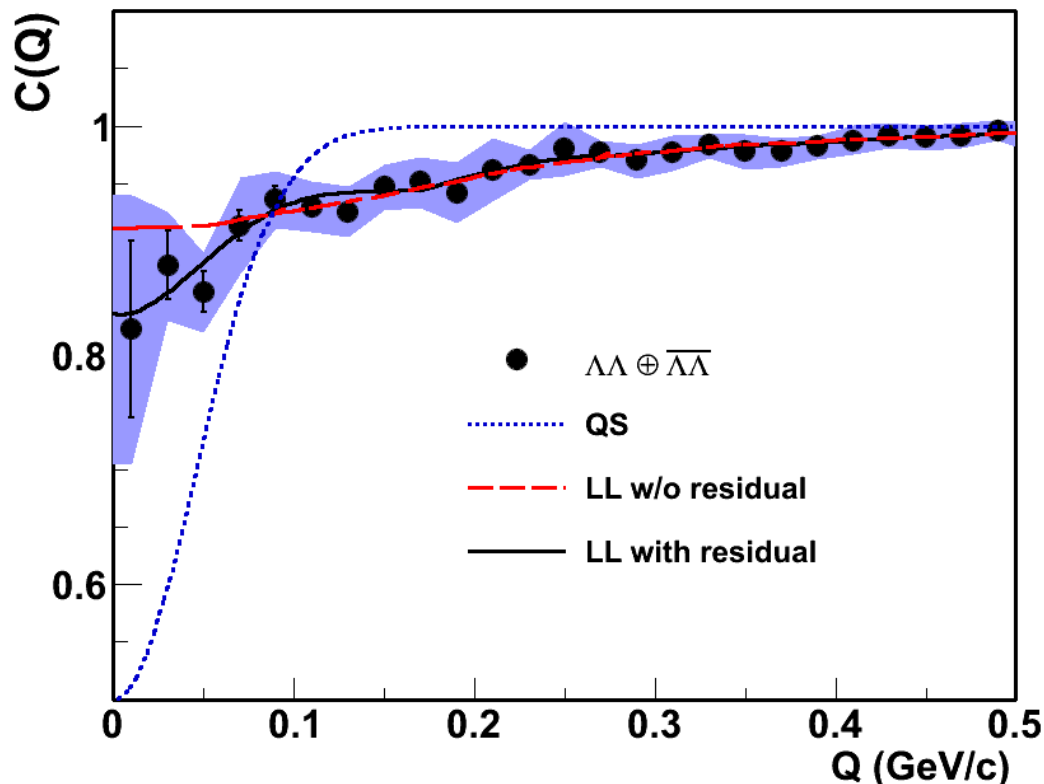
Summary and Outlook

- RHIC has potential to determine $\Lambda\Lambda$ interaction
- Expansion effect modifies low Q behavior
- Ideal measurements (i.e., decay contribution is subtracted) will give strong constraints
- Feed-down effect reduces resolving power
- Long-tail in STAR data needs to be subtracted and its origin needs to be understood
- Scattering length $1/a_0 < -0.8 \text{ fm}^{-1}$

Backup

Analysis by STAR Coll.

arXiv:1408.4360



Fit w/ Lednicky-Lyuboshitz model
 Data: long tail (\rightarrow Small source size)
 Introduce “residual correlation” for
 a better fit to data

$$a_0 = -1.10 \pm 0.37_{-0.08}^{+0.68} \text{ fm}$$

$$r_{\text{eff}} = 8.52 \pm 2.56_{-0.74}^{+2.09} \text{ fm}$$

$$\chi^2/N_{\text{dof}} = 0.56$$

$$C_{\text{fit}}(Q) = N \left[1 + \lambda \left\{ -\frac{1}{2} e^{-r_0^2 Q^2} + \chi(a_0, r_{\text{eff}}) \right\} + a_{\text{res}} \exp(-r_{\text{res}}^2 Q^2) \right]$$

	HBT (size r_0)	Interaction	Residual correlation
1.006	0.18	2.96 fm	-0.044
			0.43 fm

TABLE I: $\Lambda\Lambda$ potentials. The scattering length (a_0) and effective range (r_{eff}) are fitted using a two-range gaussian potential, $V_{\Lambda\Lambda}(r) = V_1 \exp(-r^2/\mu_1^2) + V_2 \exp(-r^2/\mu_2^2)$.

Model	a_0 (fm)	r_{eff} (fm)	μ_1 (fm)	V_1 (MeV)	μ_2 (fm)	V_2 (MeV)	Ref.
ND46	4.621	1.300	1.0	-144.89	0.45	127.87	[18] $r_c = 0.46$ fm
ND48	14.394	1.633	1.0	-150.83	0.45	355.09	[18] $r_c = 0.48$ fm
ND50	-10.629	2.042	1.0	-151.54	0.45	587.21	[18] $r_c = 0.50$ fm
ND52	-3.483	2.592	1.0	-150.29	0.45	840.55	[18] $r_c = 0.52$ fm
ND54	-1.893	3.389	1.0	-147.65	0.45	1114.72	[18] $r_c = 0.54$ fm
ND56	-1.179	4.656	1.0	-144.26	0.45	1413.75	[18] $r_c = 0.56$ fm
ND58	-0.764	6.863	1.0	-137.74	0.45	1666.78	[18] $r_c = 0.58$ fm
NF42	3.659	0.975	0.6	-878.97	0.45	1048.58	[19] $r_c = 0.42$ fm
NF44	23.956	1.258	0.6	-1066.98	0.45	1646.65	[19] $r_c = 0.44$ fm
NF46	-3.960	1.721	0.6	-1327.26	0.45	2561.56	[19] $r_c = 0.46$ fm
NF48	-1.511	2.549	0.6	-1647.40	0.45	3888.96	[19] $r_c = 0.48$ fm
NF50	-0.772	4.271	0.6	-2007.35	0.45	5678.97	[19] $r_c = 0.50$ fm
NF52	-0.406	8.828	0.6	-2276.73	0.45	7415.56	[19] $r_c = 0.52$ fm
NSC89-1020	-0.250	7.200	1.0	-22.89	0.45	67.45	[20] $m_{\text{cut}} = 1020$ MeV
NSC89-920	-2.100	1.900	0.6	-1080.35	0.45	2039.54	[20] $m_{\text{cut}} = 920$ MeV
NSC89-820	-1.110	3.200	0.6	-1904.41	0.45	4996.93	[20] $m_{\text{cut}} = 820$ MeV
NSC97a	-0.329	12.370	1.0	-69.45	0.45	653.86	[21]
NSC97b	-0.397	10.360	1.0	-78.42	0.45	741.76	[21]
NSC97c	-0.476	9.130	1.0	-91.80	0.45	914.67	[21]
NSC97d	-0.401	1.150	0.4	-445.77	0.30	373.64	[21]
NSC97e	-0.501	9.840	1.0	-110.45	0.45	1309.55	[21]
NSC97f	-0.350	16.330	1.0	-106.53	0.45	1469.33	[21]
Ehime	-4.21	2.41	1.0	-146.6	0.45	720.9	[23]
fss2	-0.81	3.99	0.92	-103.9	0.41	658.2	[25]
ESC08	-0.97	3.86	0.80	-293.66	0.45	1429.27	[22]

TABLE II: $\Lambda\Lambda$ potentials from Nagara event. The scattering length (a_0) and effective range (r_{eff}) are fitted using a three-range gaussian potential, $V_{\Lambda\Lambda}(r) = V_1 \exp(-r^2/\mu_1^2) + V_2 \exp(-r^2/\mu_2^2) + V_3 \exp(-r^2/\mu_3^2)$.

Model	a_0 (fm)	r_{eff} (fm)	μ_1 (fm)	V_1 (MeV)	μ_2 (fm)	V_2 (MeV)	μ_3 (fm)	V_3 (MeV)	Ref.
HKMY	-0.575	6.45	1.342	-10.96	0.777	-141.75	0.35	2136.6	[3]
FG	-0.77	6.59	1.342	-21.49	0.777	-250.13	0.35	9324.0	[2]

$$N = 1.006 \pm 0.001$$

$$\lambda = 0.18 \pm 0.05^{+0.12}_{-0.06}$$

$$a_0 = -1.10 \pm 0.37^{+0.68}_{-0.08}$$

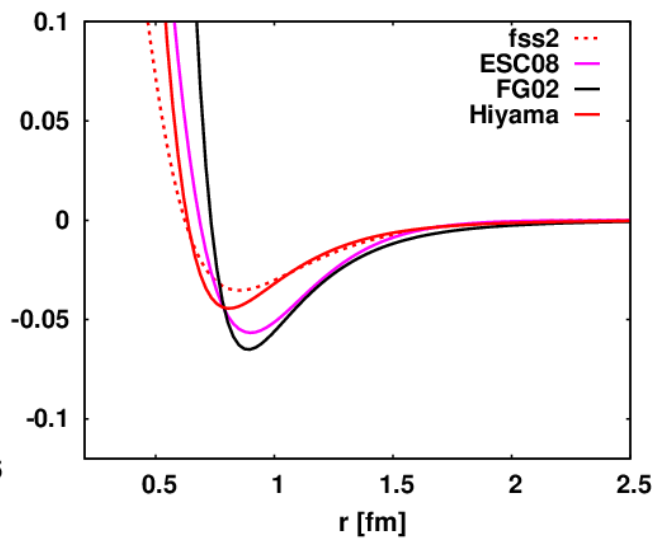
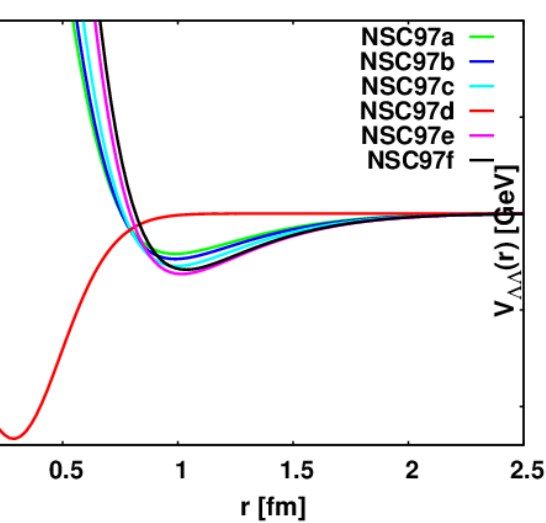
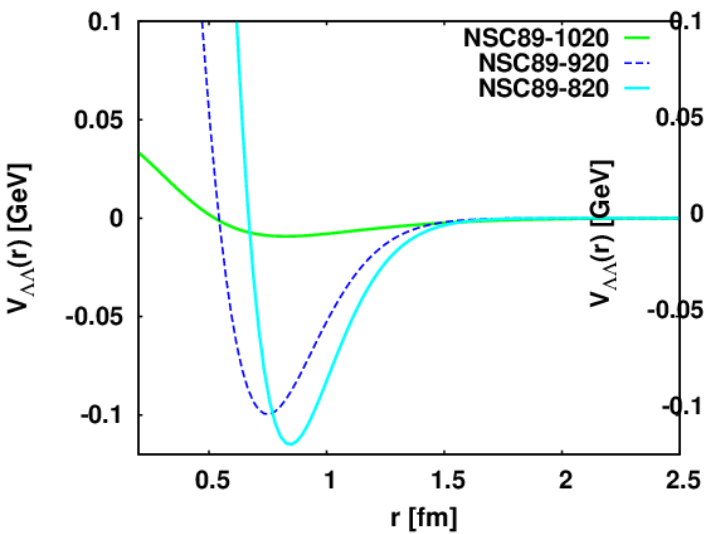
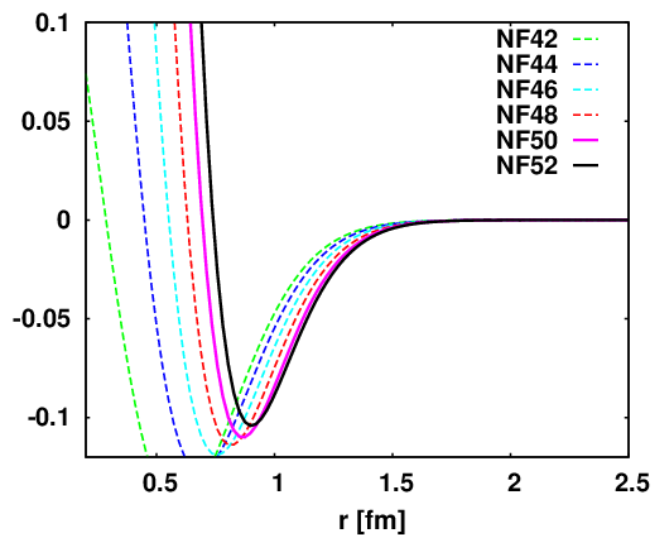
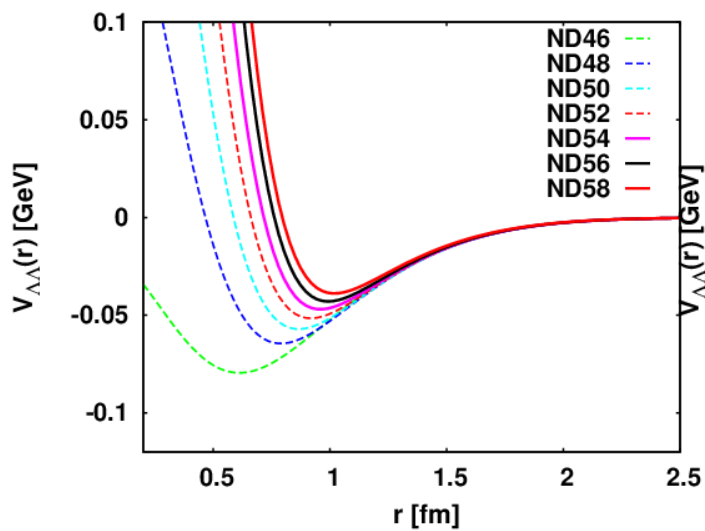
$$r_{\text{eff}} = 8.52 \pm 2.56^{+2.09}_{-0.74}$$

$$r_0 = 2.96 \pm 0.38^{+0.96}_{-0.02}$$

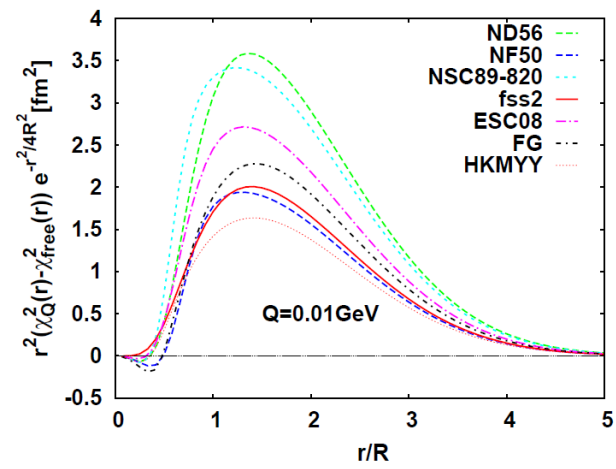
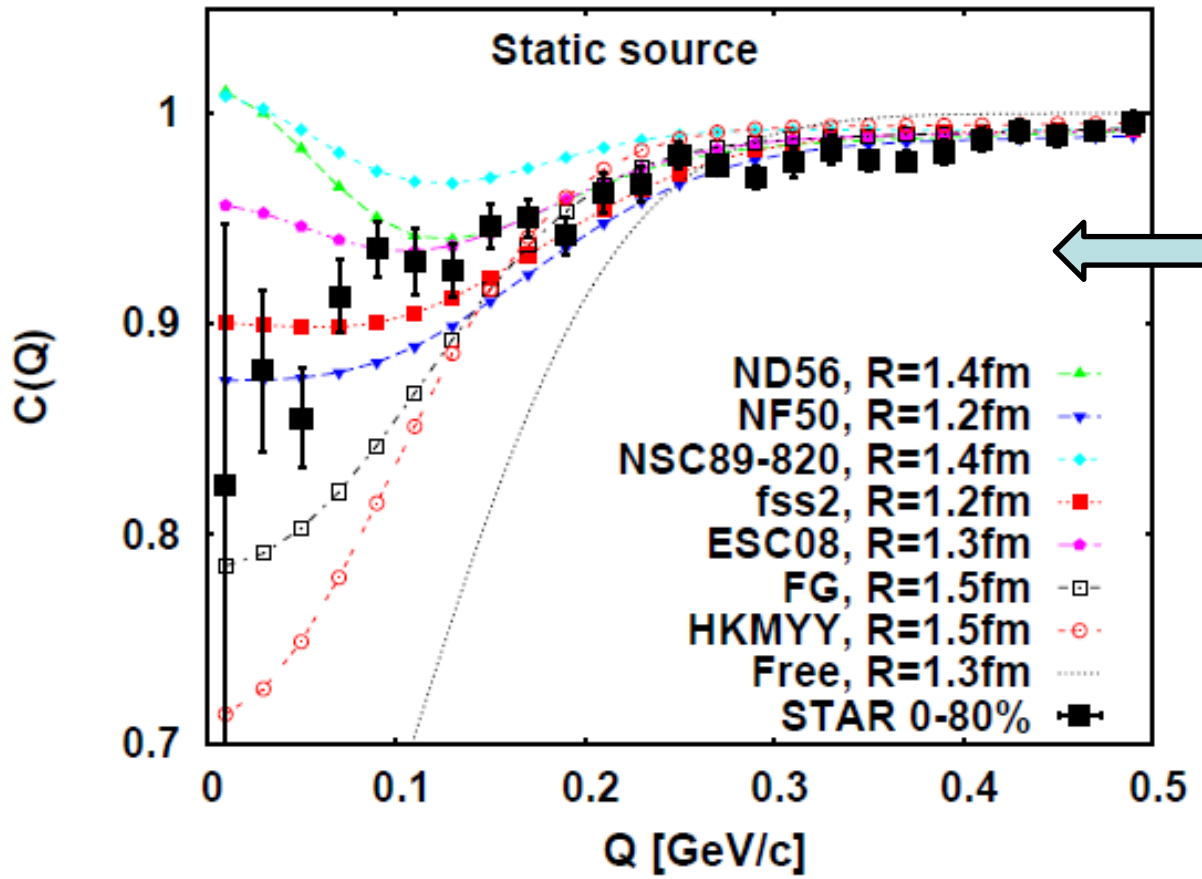
$$a_{\text{res}} = -0.044 \pm 0.004^{+0.048}_{-0.009}$$

$$r_{\text{res}} = 0.43 \pm 0.04^{+0.43}_{-0.03}$$

$$\chi^2/N_{\text{dof}} = 0.56$$



Results from the Static Source



Strong attraction is reflected onto $C(Q)$

Larger variation among potentials than data error-bars

Size : determined from $\min. \chi^2$

Collectivity Deforms Source Function

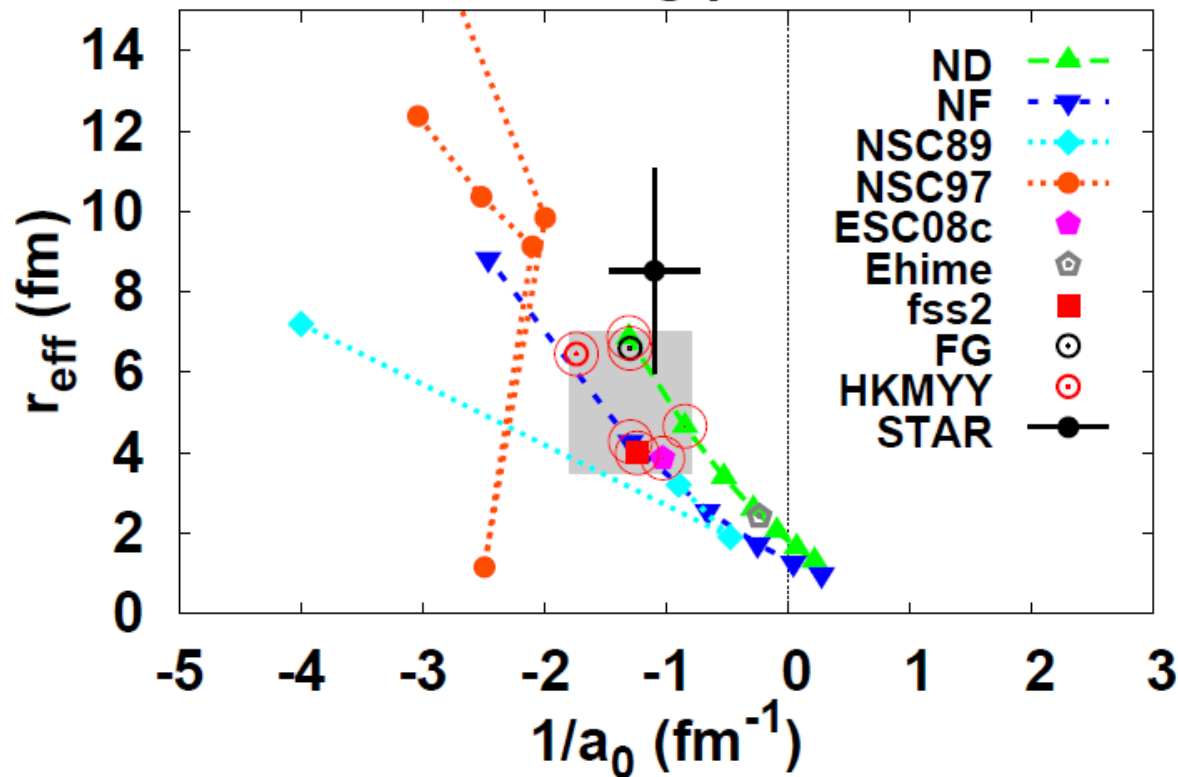
$$C_2(Q, K) = \frac{\int d^4x_1 d^4x_2 \mathcal{S}(x_1, K) \mathcal{S}(x_2, K) |\Psi_{12}(Q, x_1 - x_2 - (t_2 - t_1)K/m)|^2}{\int d^4x_1 d^4x_2 \mathcal{S}(x_1, k_1) \mathcal{S}(x_1, k_2)}$$

Influence on the best-fit potentials?

$C(Q)$ is fairly sensitive
to interaction

Scattering Length and Effective Range

$\Lambda\Lambda$ scattering parameters



$$-1.8 < 1/a_0 < -0.8 \text{ [fm}^{-1}\text{]}$$

$$3.5 < r_{\text{eff}} < 7 \text{ [fm]}$$

 Σ^0 feed-down affects this ?