Entangled Operators in AdS/CFT

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based on arXiv:1405.5946 [hep-th] with M.Nozaki and T.Takayanagi

Outline

- Motivation
- Renyi entropies for excited states
- Large N and large c
- Holographic analysis
- Summary/Future directions

AdS/CFT ("early")



Physics of local operators in CFT



Physics of Strings on AdS

$$\langle \mathcal{O}(x_1)\cdots\mathcal{O}(x_n)\rangle = \frac{\delta}{\delta\phi_0(x_1)}\cdots\frac{\delta}{\delta\phi_0(x_1)}Z_{\text{Gravity}}[\phi_0(x)]$$

Lots of developments: 2 & 3-pt functions, Integrability, Operators with large R-charge and LLM geometries

Seems well understood and explored...

AdS/CFT ("modern")

Entanglement



Spacetime Geometry

 $S_A = \frac{\operatorname{Area}(\gamma_A^d)}{4G_N^{d+2}}$

[Ryu,Takayanagi'06] [Hubeny,Rangamani, Takayanagi'07]

A common platform for holography, quantum many-body physics, QFT, information theory....

Time dependence and thermalization, global and local quench, Universal Laws of entanglement

[Cardy&Calabrese;Solodukhin; Myers Sinha; Sinha et al.; Maldacena,Harman; Iizuka,Ogawa;Takayanagi,Bhattacharya;Ugajin; Das,Galante,Myers;Lopez et al.; Nozaki,Numasawa,Takayanagi;.....] Universal physics of excitations created by a local operator acting on the ground state (milder version of a local quench)

$$\left|\psi\right\rangle = O(z,\bar{z})\left|0\right\rangle$$

Connection with AdS/CFT: Large N or large c

Program: Modern look on early holography

(Can we learn anything new about local operators from the perspective of entanglement ?)

Our Results:

- von-Neuman (n=1) and Renyi entropies behave very differently
- Naive large N limit breaks down for the Entanglement Entropy
- Universal scaling of the Renyi entropy for excited states at large c
- Holographic analysis consistent with large c and large N analysis...

Main Tool:

[M.Nozaki's talk]

Relative Renyi entropies for states excited by local operators

$$\rho(t) = \mathcal{N} \cdot e^{-iHt} e^{-\epsilon H} O(0, x^{i}) |0\rangle \langle 0| O^{\dagger}(0, x^{i}) e^{-\epsilon H} e^{iHt}$$
$$= \mathcal{N} \cdot O(w_{2}, \bar{w}_{2}, \mathbf{x}) |0\rangle \langle 0| O^{\dagger}(w_{1}, \bar{w}_{1}, \mathbf{x}),$$

Reduced $\rho_A = \operatorname{Tr}_B(\rho)$

$$\Delta S_A^{(n)} = \frac{1}{1-n} \log \left(\frac{\text{Tr}(\rho_A^n)}{\text{Tr}(\rho_A^{(0)})^n} \right) = \frac{1}{1-n} \log \left[\frac{\langle O(w_1, \bar{w}_1) O^{\dagger}(w_2, \bar{w}_2) \dots O(w_{2n}, \bar{w}_{2n}) \rangle_{\Sigma_n}}{(\langle O(w_1, \bar{w}_1) O^{\dagger}(w_2, \bar{w}_2) \rangle_{\Sigma_1})^n} \right]$$

Connection with AdS/CFT: Large N or large c

CFTs and large N limit

Operator:

 $n \ge 2$

$$Tr(\mathcal{Z}^J) = Tr(\phi_1 + i\phi_2)^J$$

e.g. J=2 (Free Field or L-R decomposition)

$$\begin{split} \Delta S_R^{(n)} &= \frac{1}{1-n} \log \left(2^{1-2n} + \frac{1}{2^n N^{2(n-1)}} \right) \\ \Delta S_A^{(n)} &= \frac{Jn-1}{n-1} \log 2 \end{split}$$

(n=1) von-Neuman entropy

$$\Delta S_R^{(1)} = \log\left(2\sqrt{2}N\right)$$

 $N^2 \,\, {
m d.o.f}$

<u>Ground State</u> Both scale as c~N^2 2d CFTs

(n=2)

$$\frac{\langle O(w_1, \bar{w}_1) O(w_2, \bar{w}_2) O(w_3, \bar{w}_3) O(w_4, \bar{w}_4) \rangle_{\Sigma_2}}{\left(\langle O(w_1, \bar{w}_1) O(w_2, \bar{w}_2) \rangle_{\Sigma_1} \right)^2} = |z|^{2\Delta_O} |1 - z|^{2\Delta_O} G_O(z, \bar{z})$$

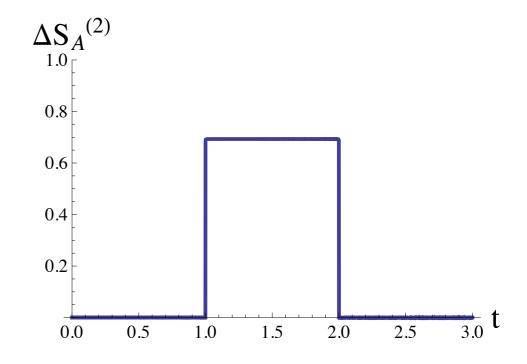
At late time $(z, \overline{z}) \rightarrow (1, 0)$

In rational CFTs

$$G(z,\bar{z}) \simeq F_{00}[O] \cdot (1-z)^{-2\Delta_O} \bar{z}^{-2\Delta_O}$$

$$\Delta S_A^{(2)} = -\log F_{00}[O] = \log d_O$$
$$d_\Delta = \frac{S_{0\Delta}}{S_{00}}$$

quantum dimension



2d CFTs at large c

[PC,M.Nozaki,T.Takayanagi14]

Conformal block expansion

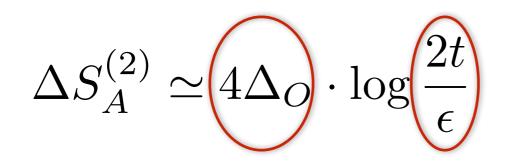
$$G(z,\bar{z}) = \sum_{b} (C_{OO^{\dagger}}^{b})^2 F_O(b|z) \bar{F}_O(b|\bar{z})$$

at large central charge c

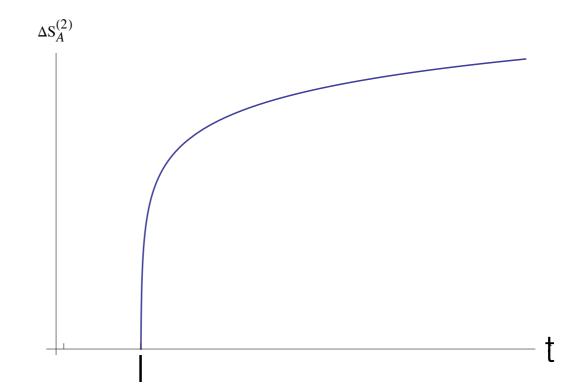
[Fateev,Ribault'11]

$$F_O(b|z) \simeq z^{\Delta_b - 2\Delta_O} \cdot {}_2F_1(\Delta_b, \Delta_b, 2\Delta_b, z)$$

at late time



similar to a local quench



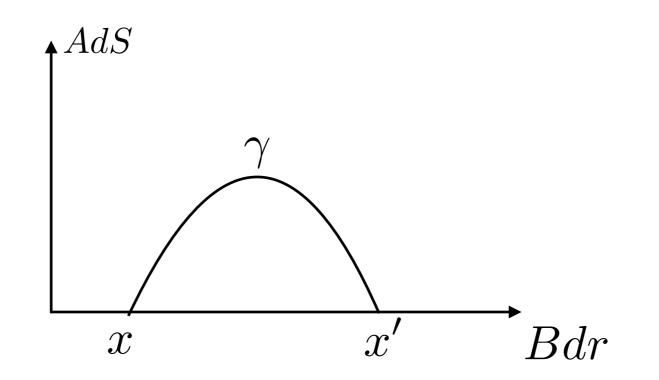
Holographic Checks

Geodesics and propagators

$$\langle \mathcal{O}(x_1)\cdots\mathcal{O}(x_n)\rangle = \frac{\delta}{\delta\phi_0(x_1)}\cdots\frac{\delta}{\delta\phi_0(x_1)}Z_{\text{Gravity}}[\phi_0(x)]$$

for operators with "large" conformal dimension (semiclassical)

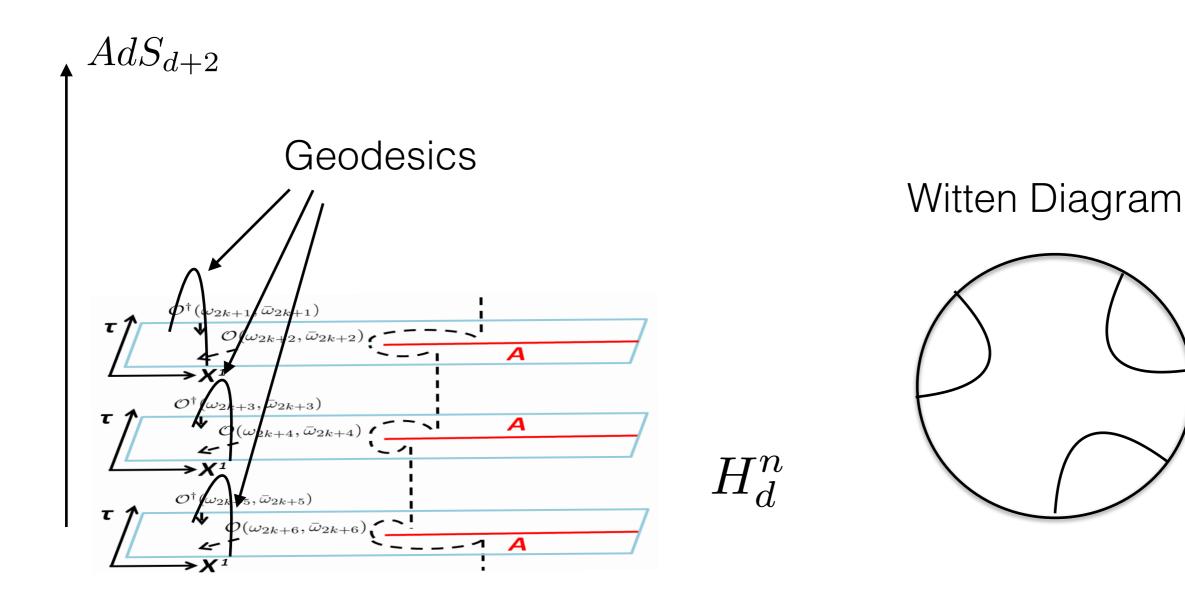
$$\langle \Phi_{\Delta}(x)\Phi_{\Delta}(x')\rangle \sim e^{-\frac{\Delta}{R}L(\gamma)}$$
 [Balasubramanian,Ross'99]



Geodesics in topological Black Holes



horizon H_d^n



Geodesics in topological Black Holes

[Emparan'99] [Cassini,Huerta,Myers' 11]

$$ds^{2} = f(r)d\tau^{2} + \frac{dr^{2}}{f(r)} + r^{2}d\phi^{2} + r^{2}e^{-2\phi}dx_{i}^{2}$$
$$f(r) = -1 - \frac{\mu}{r^{d-2}} + \frac{r^{2}}{R^{2}}$$

Temperature and "n"

$$\beta = 2\pi nR$$

holographic ratio of the propagators (eg d=2)

$$e^{-\frac{2\Delta_O}{R}\left(L^{(n)}-L^{(1)}\right)} = \left(\frac{\cosh\left(\Delta\phi\right) - \cos\left(\frac{\Delta\tau}{R}\right)}{n^2\left(\cosh\left(\frac{\Delta\phi}{n}\right) - \cos\left(\frac{\Delta\tau}{nR}\right)\right)}\right)^{2\Delta_O}$$

perfectly matches the 2d CFT result

Results $n \ge 2$

 AdS_3

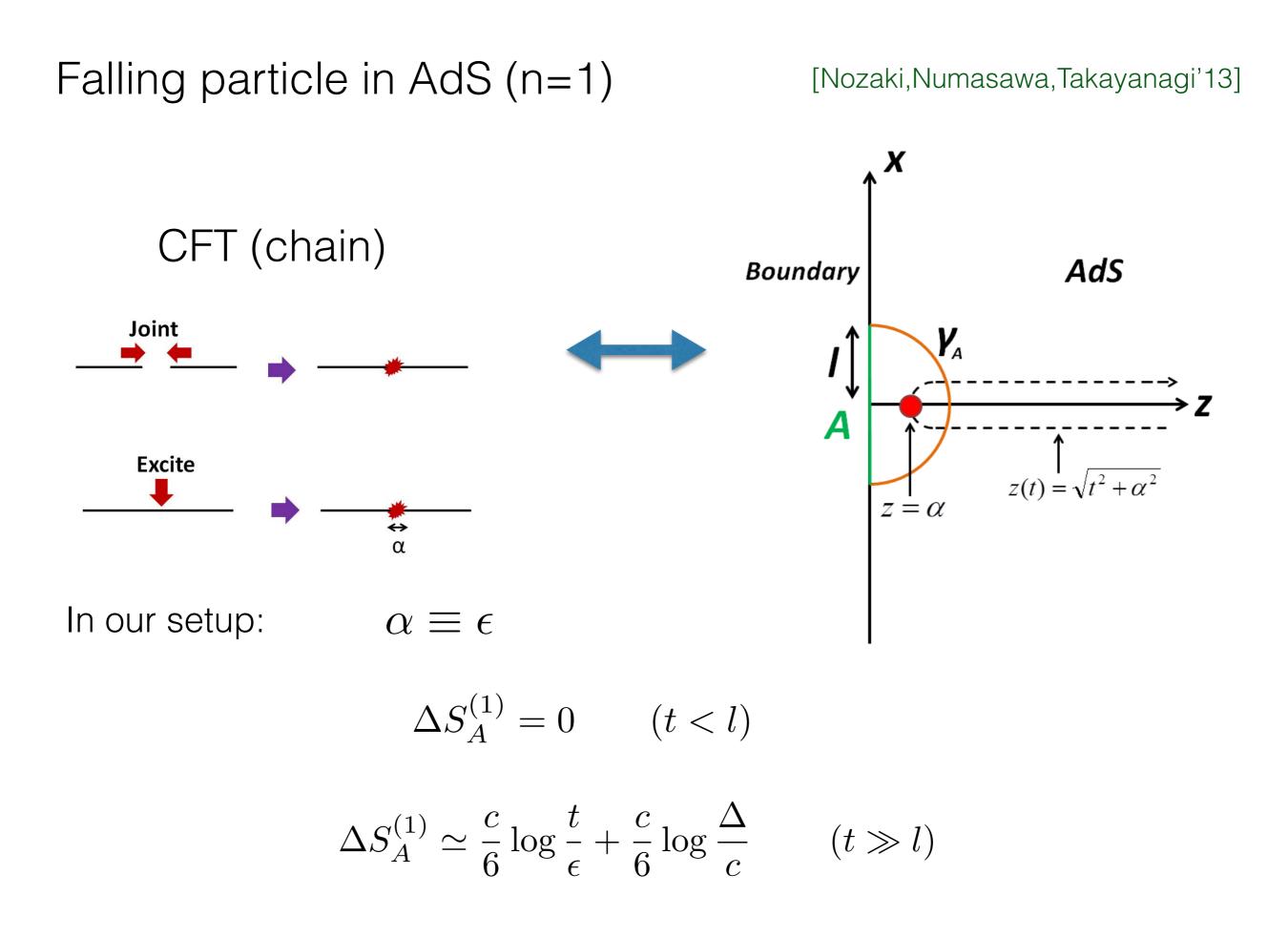
$$\Delta S_A^{(n)} \simeq \frac{2n\Delta_O}{n-1}\log\left(\frac{n\sin\left(\frac{\pi}{n}\right)t}{\epsilon}\right) - \frac{1}{n-1}\log 2$$

higher d at late time

$$\Delta S_A^{(n)} \simeq \underbrace{\frac{4n\Delta}{d(n-1)}} \log\left(\frac{F_{(d,n)}t}{2\epsilon}\right) + C_{(n,d)} - \frac{1}{n-1}\log 2$$

n->1

n->1 from holography?



General picture (summary)

rational CFT
$$D_n = (d_O)^{n-1}$$

free field (large N)
$$D_n = 2^{Jn-1} + O(N^{-2})$$

$$\Delta S_A^{(n)} = -\frac{1}{n-1} \log \frac{\langle O^{\dagger} O \cdots O^{\dagger} O \rangle_{\Sigma_n}}{(\langle O^{\dagger} O \rangle_{\Sigma_1})^n}$$
$$\simeq -\frac{1}{n-1} \log \left[\frac{1}{D_n} + \mu_n \cdot \left(\frac{\epsilon}{t}\right)^{\nu_n} \right]$$

Large c
(strongly interacting)
$$\Delta S_A^{(n)} \simeq \frac{\nu_n}{n-1} \log \frac{t}{\epsilon}$$
$$\nu_n \simeq \frac{4n\Delta}{d} + O\left(\frac{\Delta^2}{c}\right)$$

$$D_n \sim e^{b_n \cdot N^{a_n}}$$

$$\lim_{n \to 1} \frac{\nu_n}{n-1} \simeq c$$
?

Thank You