

*Charged Quantum Entanglement  
and  
Symmetry Protected Topological phases*

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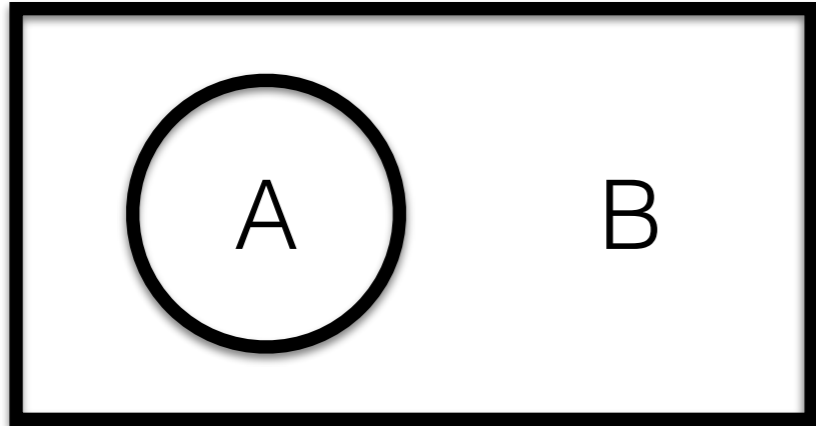
with/ Alexandre Belin, Janet Hung, Alexander Maloney  
Robert Myers, Todd Sierens

# QUANTUM ENTANGLEMENT

Useful probe in quantum systems

- Quantum Information
  - efficiency of information transmission
- High energy physics / Quantum Gravity
  - emergent space time, black hole information
- Condensed Matter Physics
  - topological order, quantum phase transition

# QUANTUM ENTANGLEMENT



Reduced density matrix

$$\rho_A = \text{Tr}_B \rho$$

Entanglement Hamiltonian

$$\rho_A = \frac{e^{-H_E/T_0}}{Z(T_0)}$$

Entanglement Entropy

$$S_{EE} = -\text{Tr}_A \rho_A \log \rho_A$$

Number of states in B that are consistent with all measurements in A

# Entanglement vs Thermal

## Entanglement Entropy

$$S_{EE} = -\text{Tr}_A \rho_A \log \rho_A$$

$$\rho_A = \frac{e^{-H_E/T_0}}{Z(T_0)}$$

- finite at zero temperature
- one state
- in and out of equilibrium
- grand canonical ensemble

[Belin-Hung-Myers-SM-Sierens]

## Thermal Entropy

$$S = -\text{Tr} \rho_{therm} \log \rho_{therm}$$

$$\rho_{therm} = \frac{e^{-H/T}}{Z(T)}$$

- zero at zero temperature
- all states
- in equilibrium
- Laws of thermodynamics

$$dF = -S_{therm}dT + \mu dQ$$

# Entanglement Chemical potential

Grand Canonical ensemble

$$\tilde{\rho}_{therm} = \frac{e^{-(H - \mu Q)/T}}{Z(T, \mu)}$$

# Entanglement Chemical potential

Grand Canonical ensemble in Entanglement

$$\tilde{\rho}_A = \frac{e^{-(H_E - \mu_E Q)/T_0}}{Z(T_0, \mu_E)}$$

$\mu_E$  Entanglement chemical potential

Total system has a fixed charge (say  $Q_{\text{total}}=0$ )  
while subsystems can have any charges ( $Q_A = -Q_B$ )

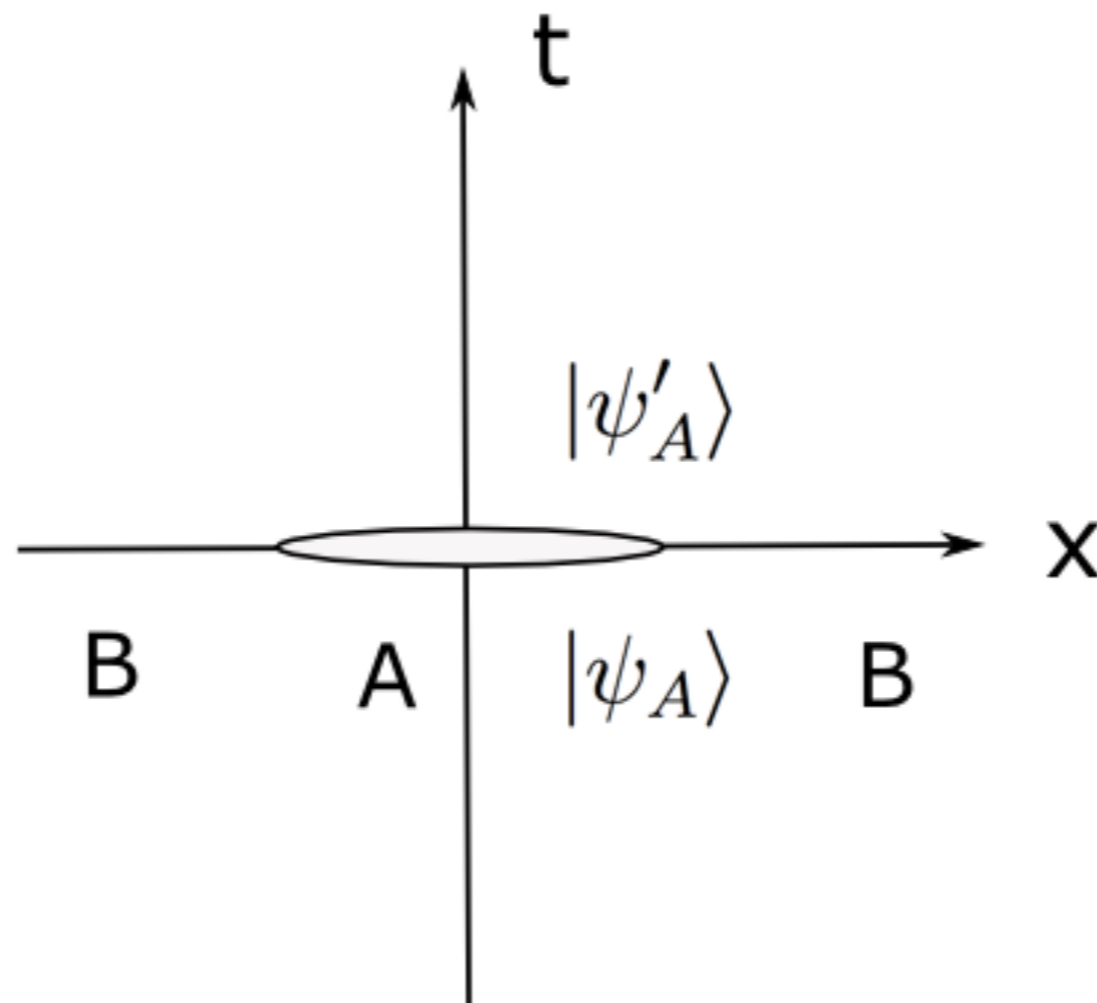
$$|\Psi\rangle = \sum_i a_i |A, Q_i\rangle \otimes |B, -Q_i\rangle$$

$$S_n(\mu, A) = S_n(-\mu, B)$$

# Path Integral

Reduced density matrix of ground state

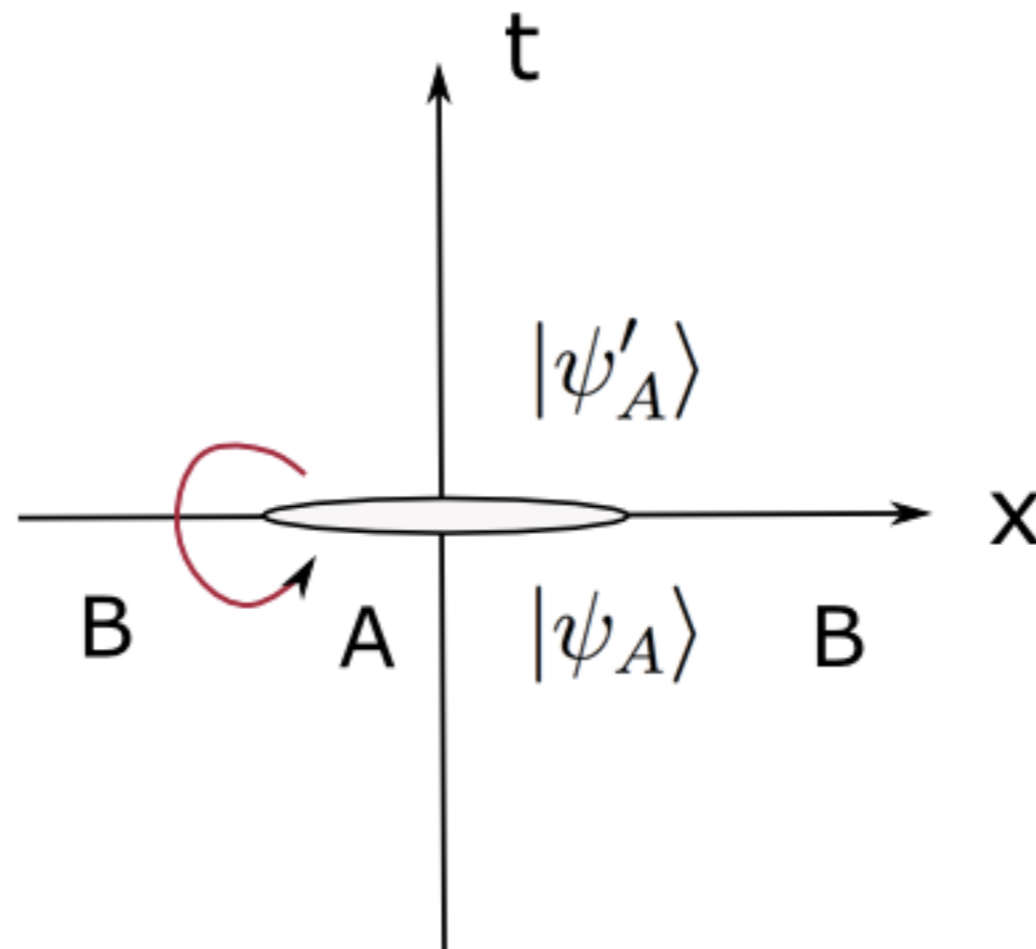
$$\langle \psi_A | \rho | \psi'_A \rangle = \frac{1}{Z} \int_{\substack{\psi(t=0_-, x \in A) = \psi_A \\ \psi(t=0_+, x \in A) = \psi'_A}} \mathcal{D}\psi e^{-S_E}$$



# Path Integral

Reduced density matrix of ground state

$$\tilde{\rho}_A = \frac{1}{Z(T_0)} \langle \psi_A | e^{-H_E/T_0} | \psi'_A \rangle$$

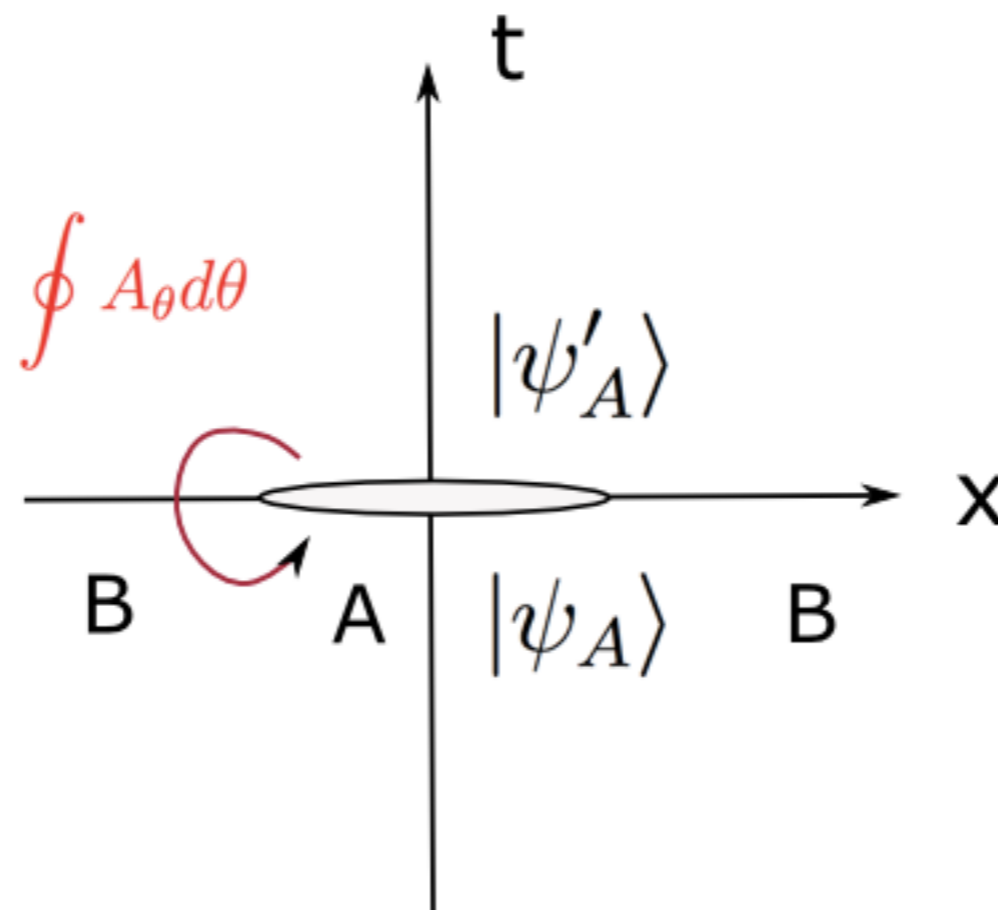




# Path Integral

Reduced density matrix of ground state

$$\tilde{\rho}_A = \frac{1}{Z(T_0, \mu_E)} \langle \psi_A | e^{-(H_E - i\mu_E Q)/T_0} | \psi'_A \rangle$$



Aharonov-Bohm effect around the entangling surface

# Conformal Field Theories

Entanglement Hamiltonian of a spherical region

$$\begin{aligned} ds_{R^d}^2 &= dt_E^2 + dr^2 + r^2 d\Omega_{d-2}^2 \quad (r < R) \\ &= \Omega^{-2} (d\tau_E^2 + R^2 (du^2 + \sinh^2 u d\Omega_{d-2}^2)) \end{aligned}$$

$$\Omega = \left| 1 + \cosh \left( u + i \frac{\tau_E}{R} \right) \right| \quad \exp \left[ - \left( u + i \frac{\tau_E}{R} \right) \right] = \frac{R - (r + it_E)}{R + (r + it_E)}$$

$$\begin{aligned} H_E &= 2\pi \int_A d^{d-1}x \frac{R^2 - r^2}{2R} T^{00}(x) \\ &= \partial_{\tau_E} \end{aligned}$$

Reduced density matrix on a spherical region  
= Thermal density matrix on a hyperbolic space

# Conformal Field Theories

Entanglement chemical potential of a spherical region

$$\begin{aligned} ds_{R^d}^2 &= dt_E^2 + dr^2 + r^2 d\Omega_{d-2}^2 \quad (r < R) \\ &= \Omega^{-2} (d\tau_E^2 + R^2 (du^2 + \sinh^2 u d\Omega_{d-2}^2)) \end{aligned}$$

Wilson loop along  $\tau_E$  direction  $\int_A A_{\tau_E} J^{\tau_E}$

Finite density system on the hyperbolic space

# Conformal Field Theories

Renyi entropy

$$S_n = \frac{1}{1-n} \log \frac{Z(T_0/n)}{Z(T_0)^n}$$

By using the following relations

$$F(T) = -T \log Z(T), \quad S_{therm} = -\frac{\partial F(T)}{\partial T}$$

we obtain

$$\begin{aligned} S_n &= \frac{n}{1-n} \frac{1}{T_0} (F(T_0) - F(T_0/n)) \\ &= \frac{n}{1-n} \frac{1}{T_0} \int_{T_0/n}^{T_0} S_{therm}(T) dT \end{aligned}$$

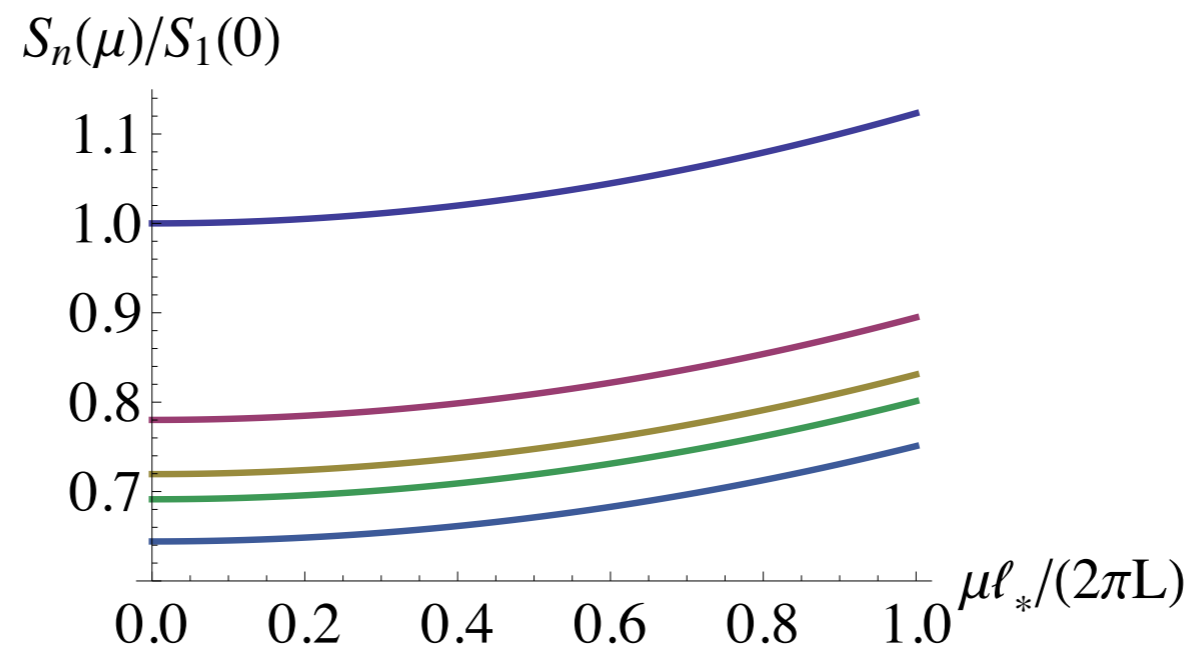
# AdS/CFT

Charged hyperbolic black hole

$$I_{E-M} = \frac{1}{2\ell_{\text{P}}^{d-1}} \int d^{d+1}x \sqrt{-g} \left( \frac{d(d-1)}{L^2} + \mathcal{R} - \frac{\ell_*^2}{4} F_{\mu\nu} F^{\mu\nu} \right)$$

$$ds^2 = -f(r) \frac{L^2}{R^2} d\tau_E^2 + \frac{dr^2}{f(r)} + r^2 (du^2 + \sinh^2 u d\Omega_{d-2}^2)$$

$$f(r) = \frac{r^2}{L^2} - 1 - \frac{m}{r^{d-2}} + \frac{q^2}{r^{2d-4}}$$



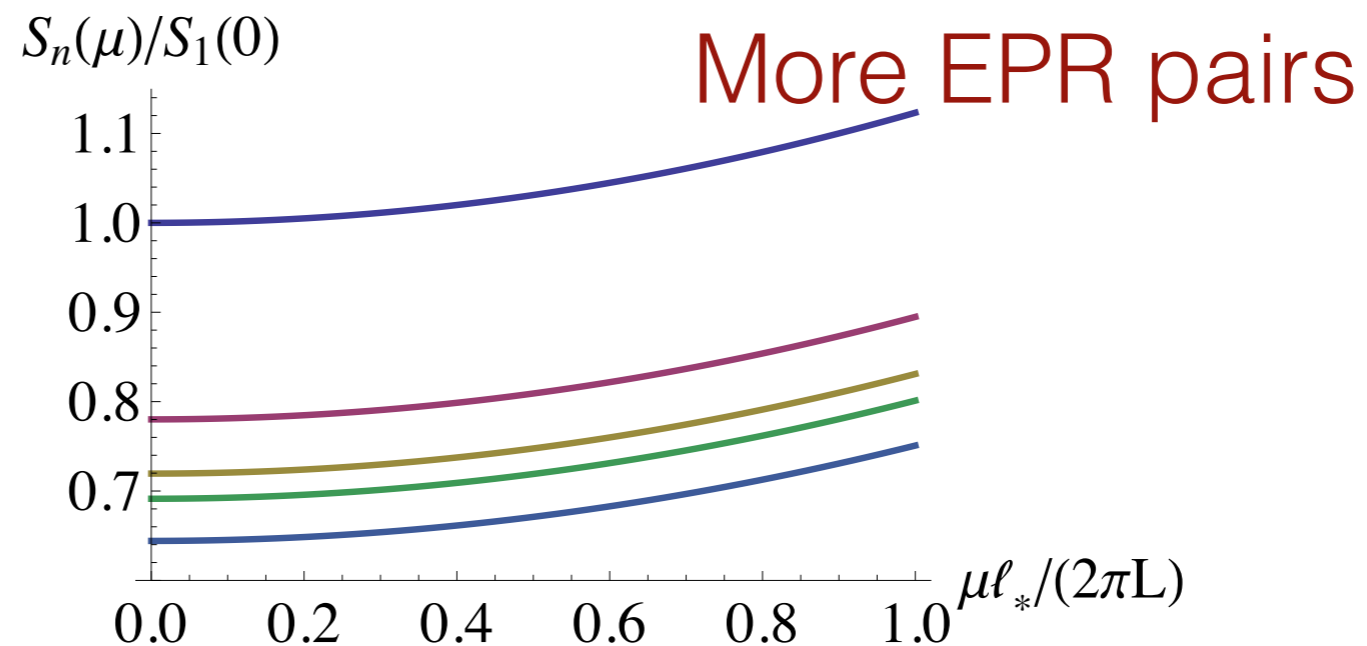
# AdS/CFT

Charged hyperbolic black hole

$$I_{E-M} = \frac{1}{2\ell_{\text{P}}^{d-1}} \int d^{d+1}x \sqrt{-g} \left( \frac{d(d-1)}{L^2} + \mathcal{R} - \frac{\ell_*^2}{4} F_{\mu\nu} F^{\mu\nu} \right)$$

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$$f(r) = \frac{r^2}{L^2} - 1 - \frac{m}{r^{d-2}} + \frac{q^2}{r^{2d-4}}$$



# Phase diagram

[Dias, Monteiro, Real, Santos/Belin, Hung, Maloney, SM]

## Einstein-Maxwell-scalar

$$I_{E-M} = \frac{1}{2\ell_P^2} \int d^4x \sqrt{-g} \left( \frac{6}{L^2} + \mathcal{R} - \frac{\ell_*^2}{4} F_{\mu\nu} F^{\mu\nu} - V(|\Phi|) - \frac{1}{2} |\nabla\Phi - iqA\Phi|^2 \right)$$

$$ds^2 = -f(r) \frac{L^2}{R^2} d\tau_E^2 + \frac{dr^2}{f(r)} + r^2 (du^2 + \sinh^2 u d\Omega_{d-2}^2)$$

Instability in the scalar field

$$\left( (\nabla^\mu - iqA^\mu)(\nabla_\mu - iqA_\mu) - m_\Phi^2 \right) \Phi = 0$$

Time dependent solution

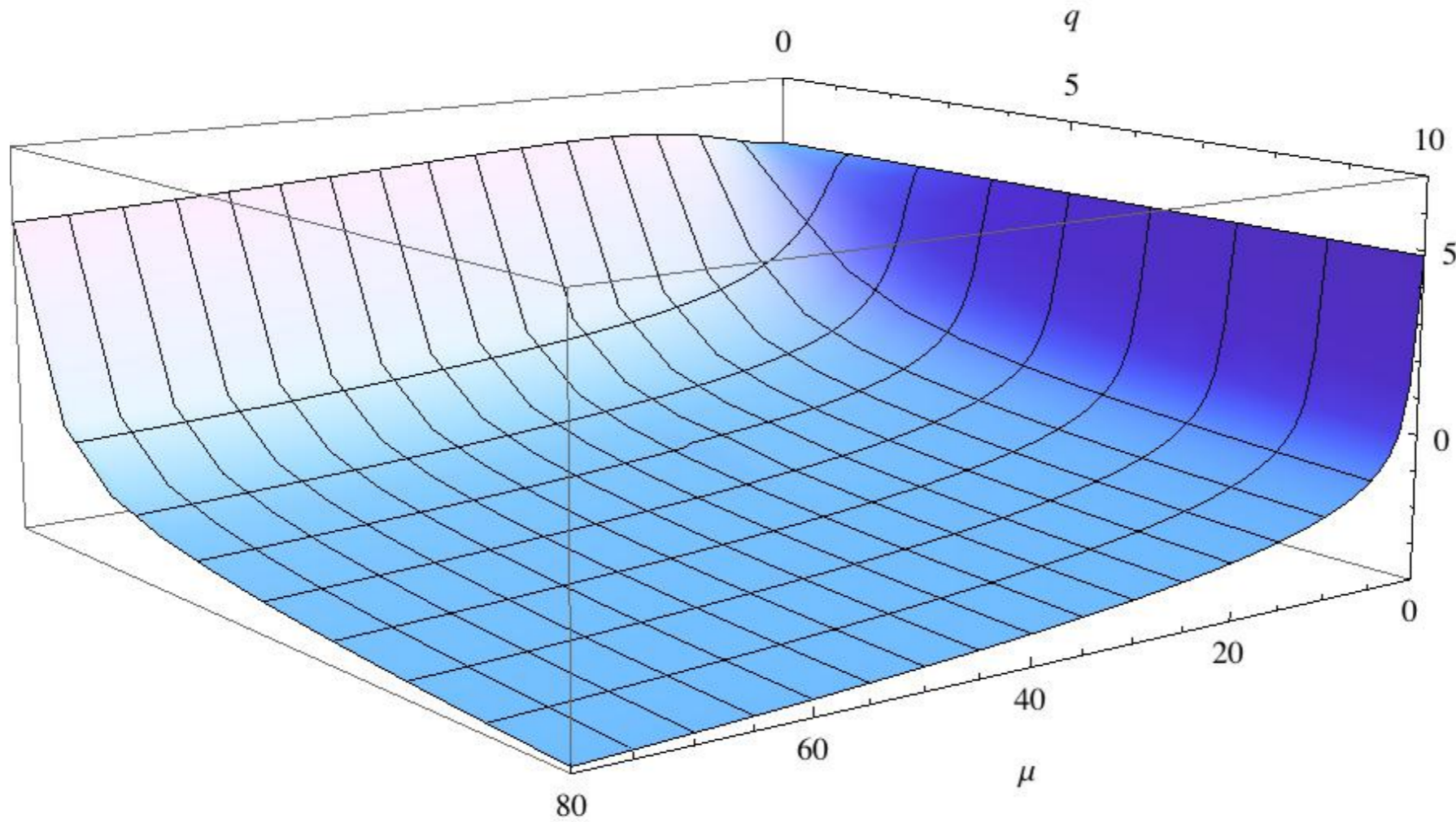
$$\Phi = \frac{\phi(r) e^{\omega t} Y(\sigma)}{r} \quad \nabla_{H_2}^2 Y = -\lambda Y \quad \lambda > 1/4$$

Effective mass/BF bound estimation

$$-\frac{d^2}{4} < m_\Phi^2 L^2 < -\frac{f''(r_H)}{8} - \frac{(d-2)^2}{4r_H^2} + \frac{2q^2\mu^2}{4\pi R^2 r_H^2 f''(r_H)}$$



Scalar charge



Critical Renyi  
parameter  
 $\log(n_c)$

Chemical potential

$$\Delta = 2$$

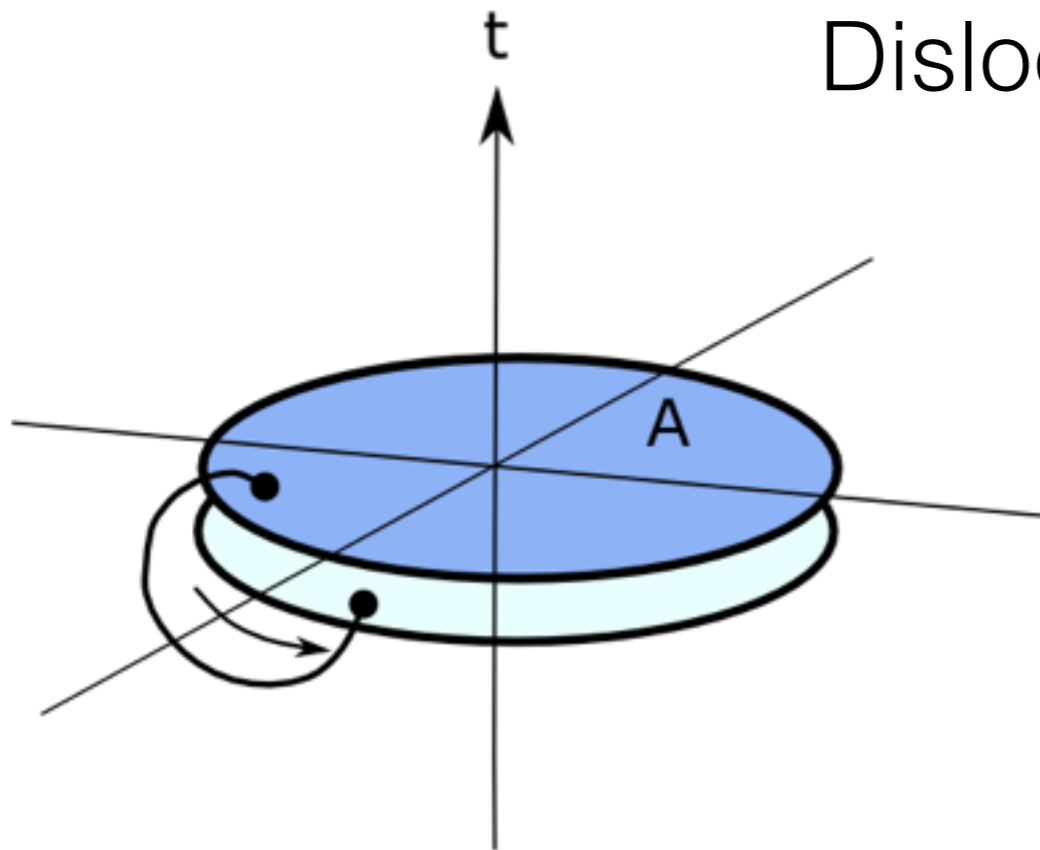


# Rotating Entanglement entropy

$$\tilde{\rho}_A = \frac{e^{-(H_E - \Omega_E J)/T_0}}{Z(T_0, \Omega_E)}$$

$\Omega_E$  : Entanglement angular potential

Dislocation along the entangling surface



Finite temperature system on hyperbolic space with

$$\phi(0, \theta) = \phi(2\pi, \theta + 2\pi\Omega_E)$$

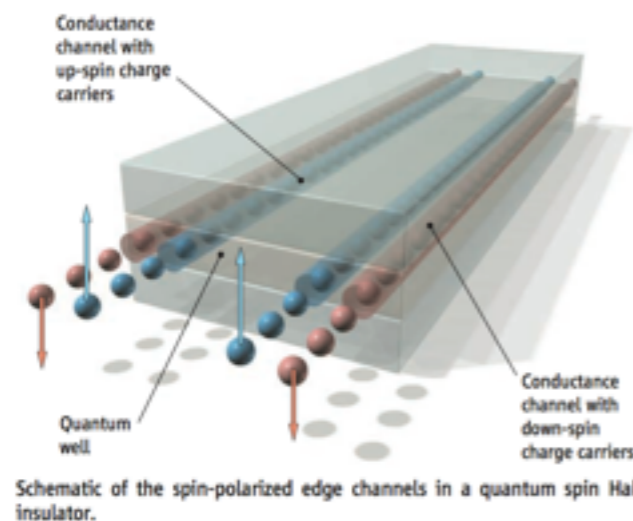
Rotating hyperbolic black holes

# Application to Topological phases

## Fractional Quantum Hall:

- Topological Order
- Ground State degeneracy
- Fractional charge
- Anyon statistics
- Long range entanglement

## Quantum Spin Hall :



[C. Bickel/Science]

- No Topological Order
- Unique ground state
- robust gapless edge states
- Short range entanglement

**Topological Insulators/Superconductors**  
**Symmetry Protected Topological phases**

# Application to SPT phases

w/Janet Hung

## Chern Simons theory

(Interacting bosons and fermions)

$$S = \frac{1}{4\pi} \int K_{IJ} a^I \wedge da^J$$

SPT phase:  $|\det K| = 1$       No topological order

Simplest bosonic case:  $K = \sigma_x$

Unstable without symmetry: Assume global symmetry  $Z_N$

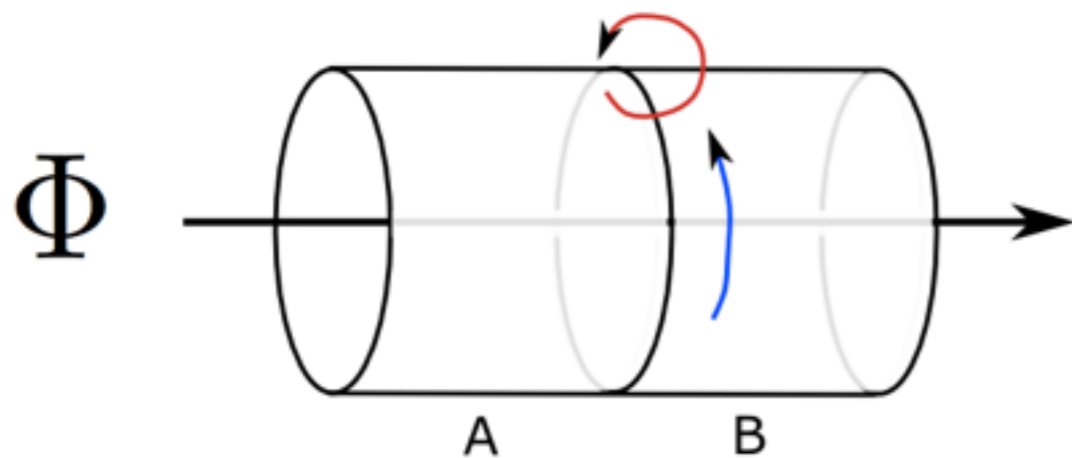
edge modes  $\phi^I \rightarrow \phi^I + 2\pi \frac{t^I}{N}$        $t^T = (1, q), \quad q \in [1, N]$

Entanglement chemical potential = External gauge field

$$S = \frac{1}{4\pi} \int K_{IJ} a^I \wedge da^J + \frac{1}{2\pi} \int K_{IJ} t^I A da^J$$

$$A_0 = \mu \quad \mu K_{IJ} t^I \int_A da^J = \mu K_{IJ} t^I \int_{\partial A} dx^i a_i^J$$

Entanglement chemical potential generates  
Wilson loop **along** entangling surface



Flux threading (Laughlin)

$$A_y = \Phi$$

Flux threading generates  
Wilson loop **around** entangling surface

# CFT story

Partition function  $Z(\tau) = \text{tr} \exp(-2\pi(\tau H - i\mu Q))$

Hamiltonian and charge operator

$$H = \frac{1}{2\pi} \int_0^{2\pi} dx \frac{1}{2} ((\partial_x \phi_L)^2 + (\partial_x \phi_R)^2)$$

$$Q = (\partial_x \phi_2 + q \partial_x \phi_1)$$

$$\phi_1(x + 2\pi, t) = \phi_1(x, t) + 2\pi(n + \frac{k_1}{N})$$

$$\phi_2(x + 2\pi, t) = \phi_2(x, t) + 2\pi(m + \frac{l_1}{N})$$

$$k_1, l_1 \in [1, N], \quad l_1 = qk_1$$

Thermodynamical limit

$$\log Z(L/\tau \rightarrow \infty) = \frac{c}{24} \frac{L}{\tau} - \gamma(= 0) + \text{Link number} \left( 2\pi i \mu \frac{l_1 + qk_1}{N} \right)$$

# Conclusion

**Generalize quantum entanglement  
by twisting boundary conditions**

**CFT**: central charge, current algebra,

**SPT**: generalized topological entanglement

Thermodynamical laws  
Einstein/Maxwell equations  
Phase diagrams  
Quantum corrections  
Discrete symmetries