

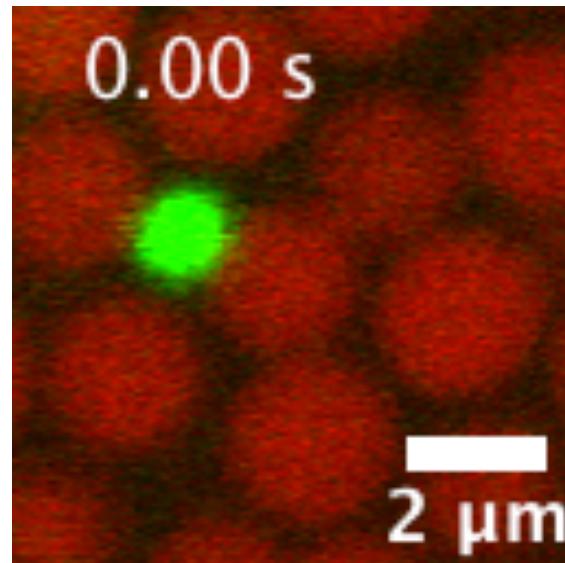
Role of surface roughness and friction in the rheology of dense colloidal suspensions

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Colloids vs non-Brownian particles

Brownian motion

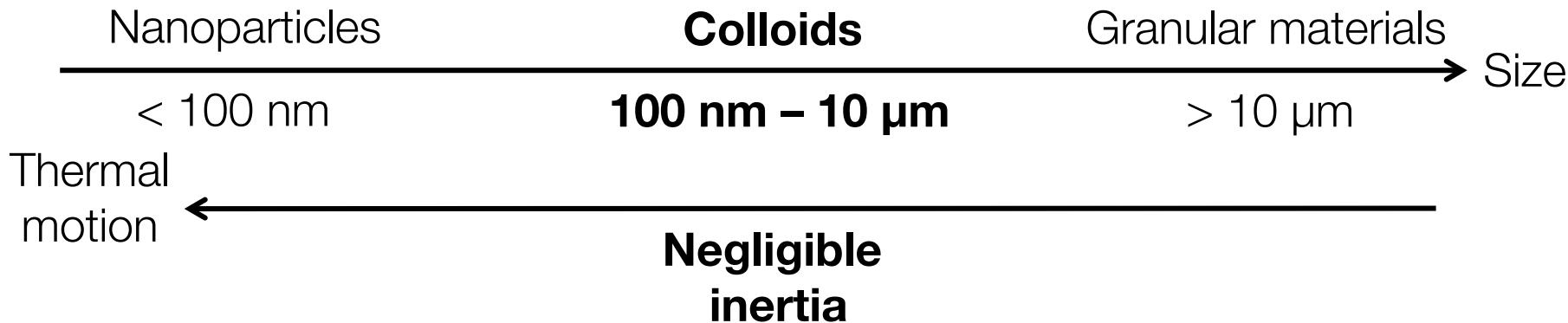


Pair potential interactions

$$\lim_{\Delta t \rightarrow \infty} \langle r^2(t) \rangle = 6D_s^\infty(\phi)\Delta t$$

$$\lim_{\Delta t \rightarrow \infty} \langle \varphi^2(t) \rangle = 4D_R^\infty(\phi)\Delta t$$

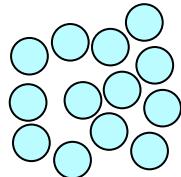
$$U(r) = 4\varepsilon \left(\left(\frac{\sigma}{r} \right)^{2\alpha} - \left(\frac{\sigma}{r} \right)^\alpha \right)$$



Flow properties and connections to microstructure

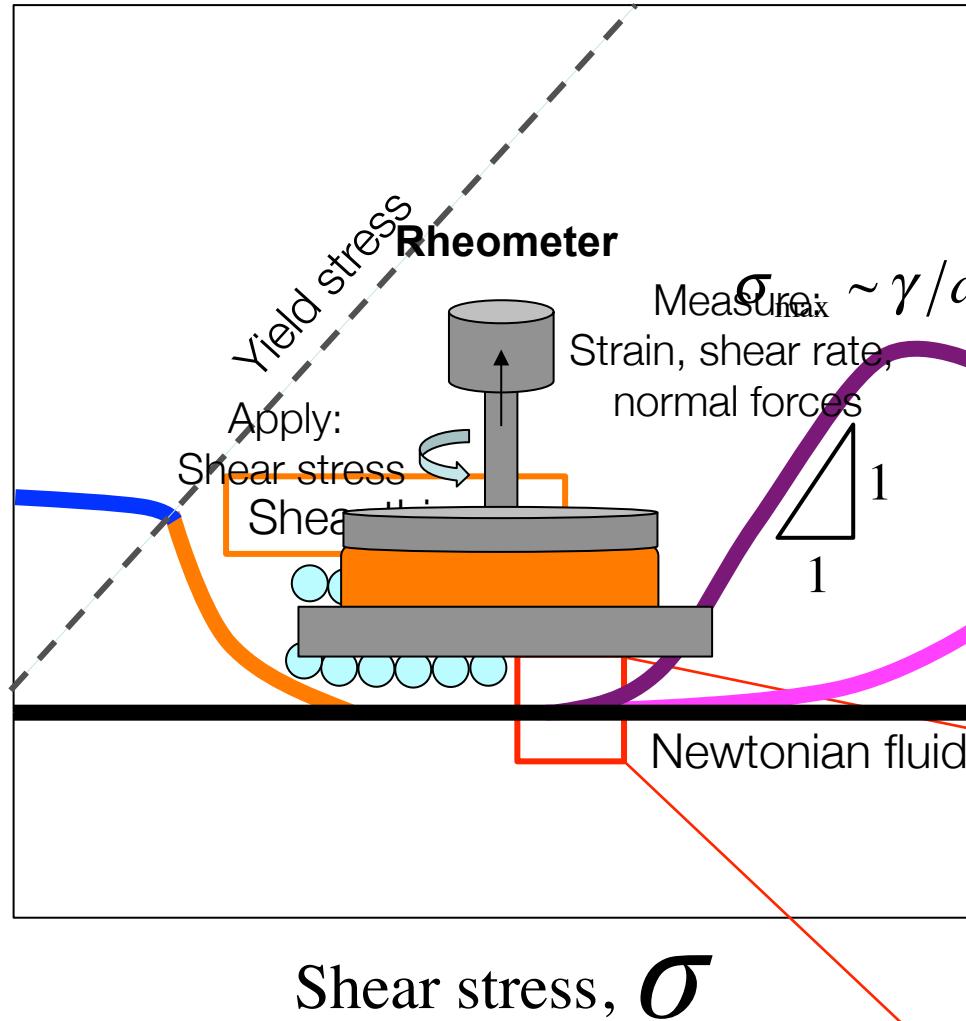
$$\eta = \frac{\sigma}{\dot{\gamma}}$$

Zero shear viscosity

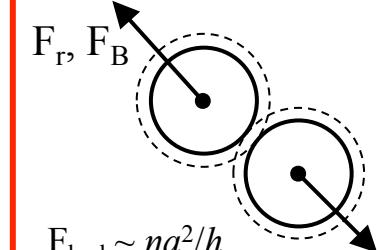


Suspensions at equilibrium

Viscosity, η



Force balance at σ_c



Force chains

"DST"

Hydroclusters

"CST"

Confocal rheometer for visualization and rheology

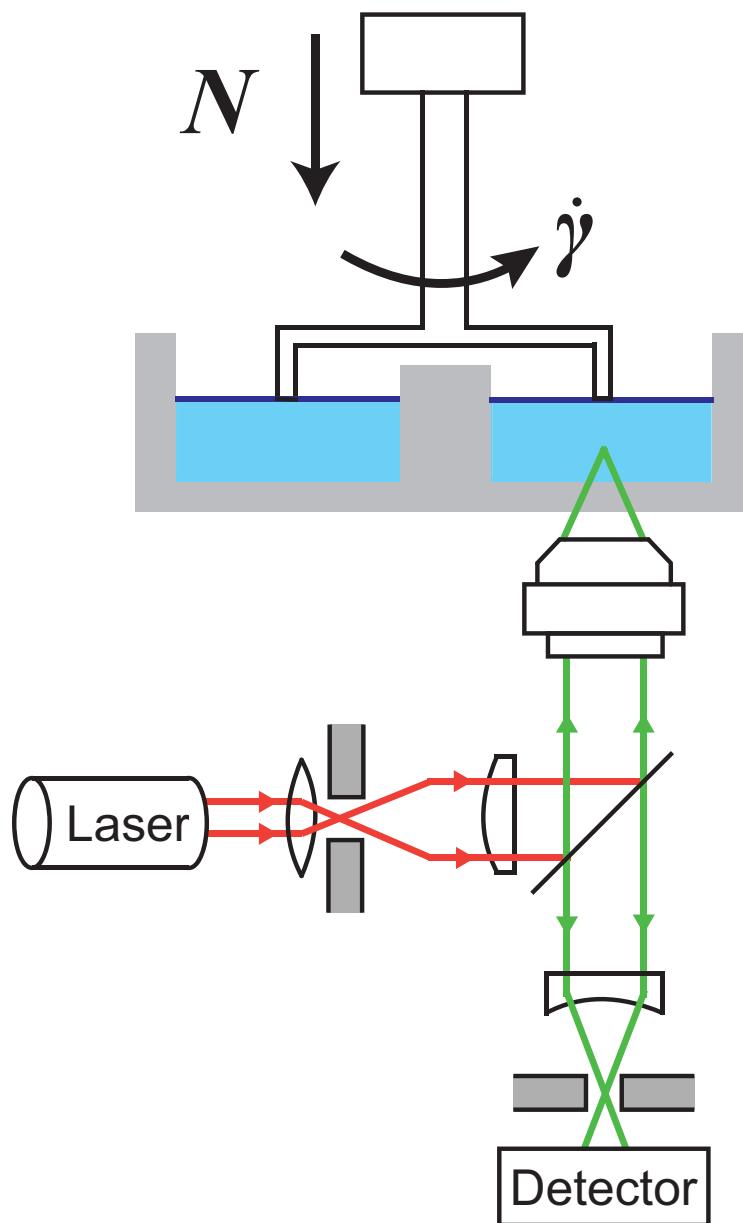
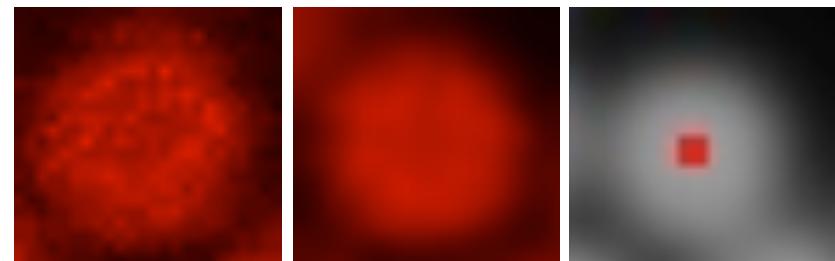
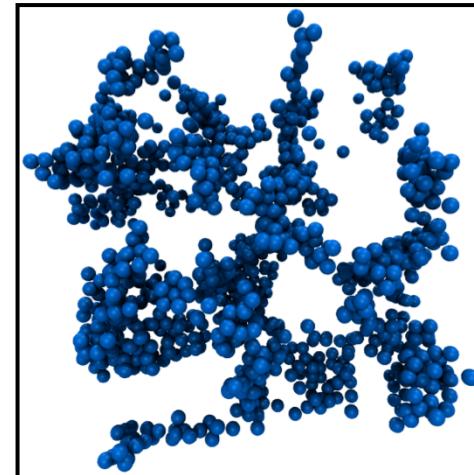


Image acquisition from confocal microscopy



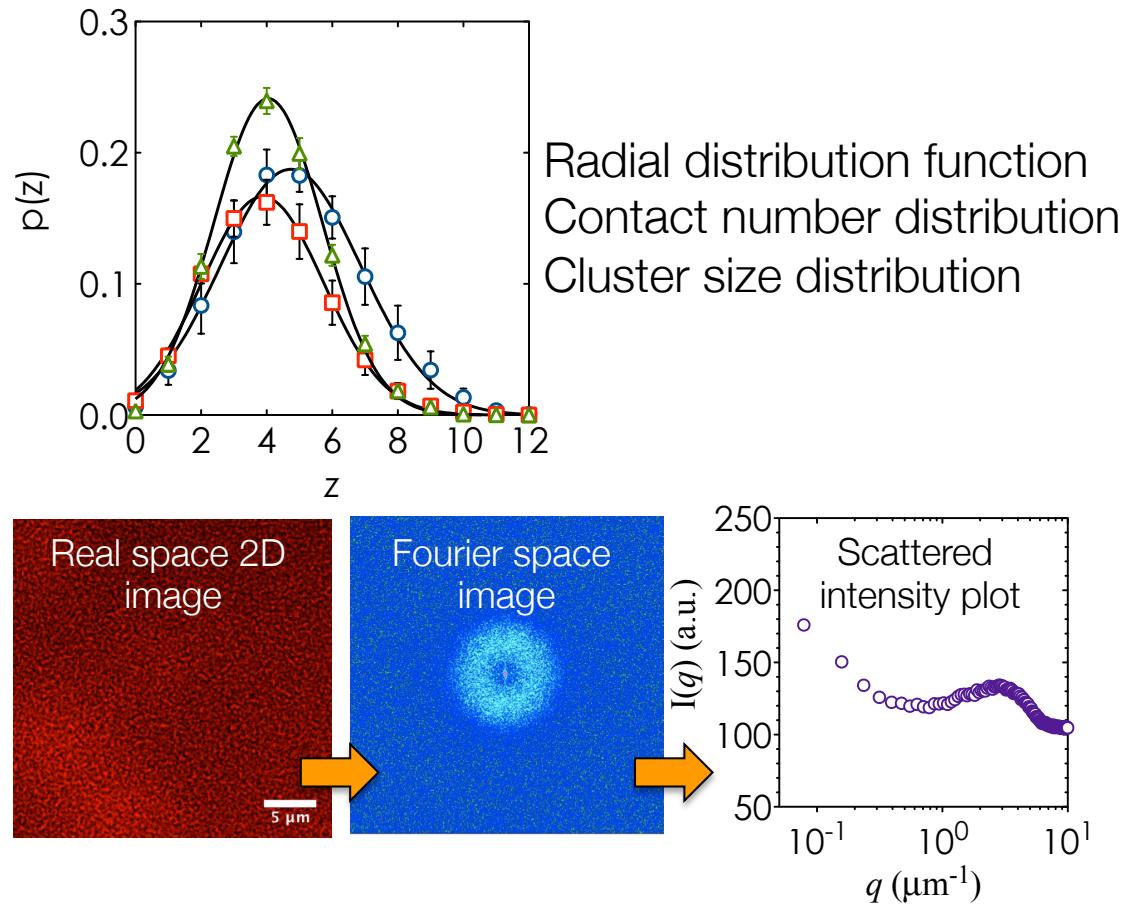
Post-processing

Space-time information

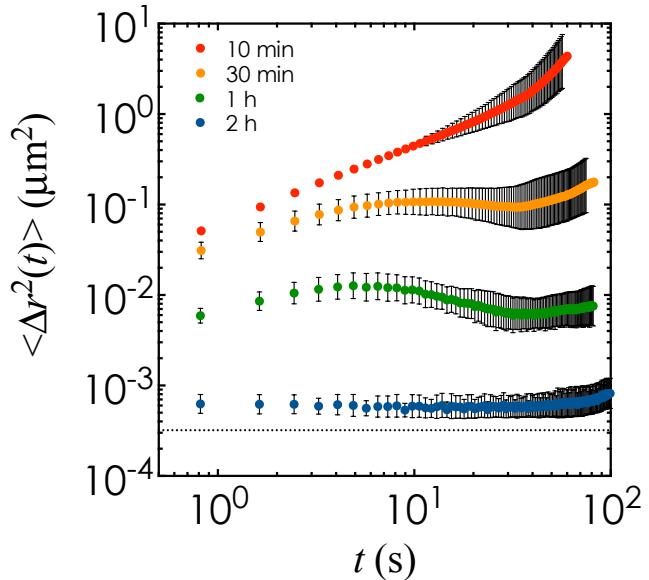
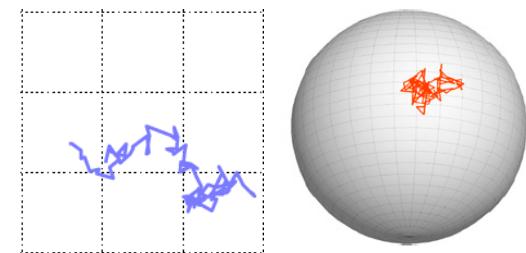


Characterizing structure and dynamics from images

Microstructure

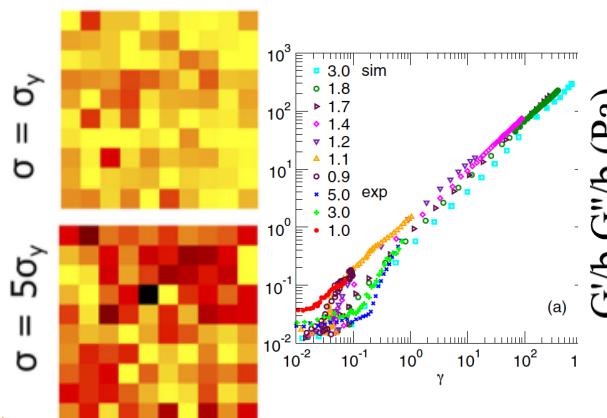
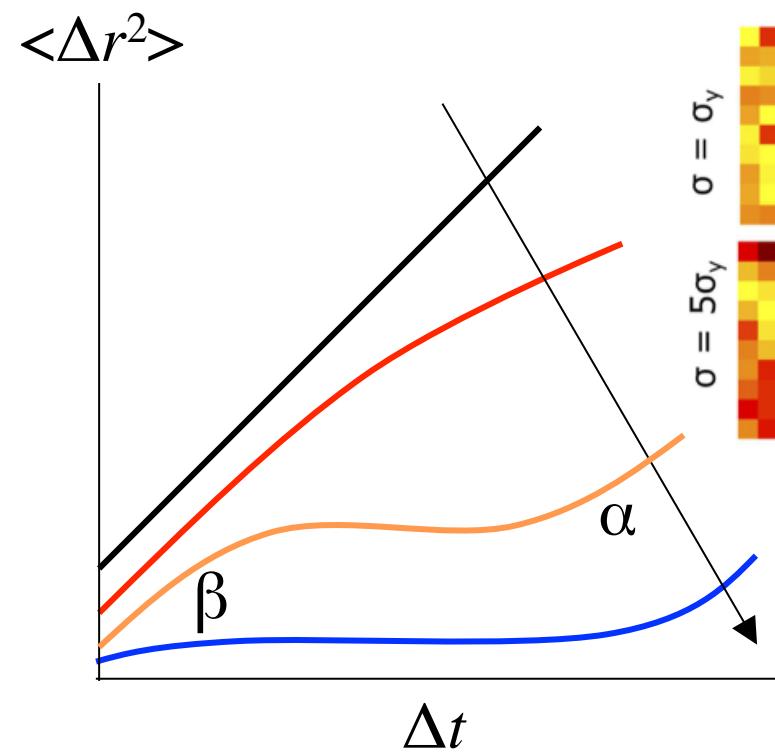
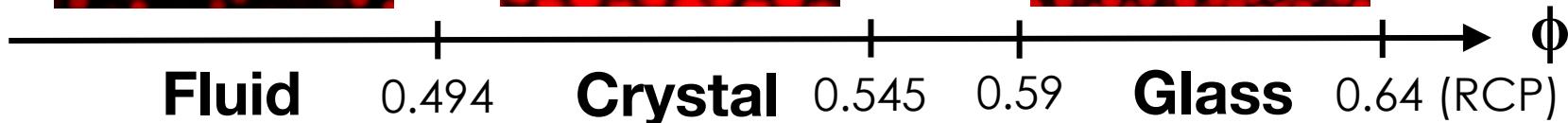
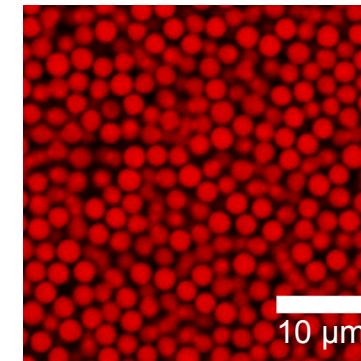
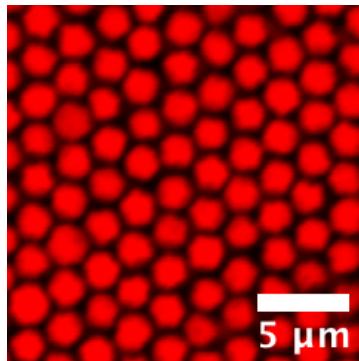
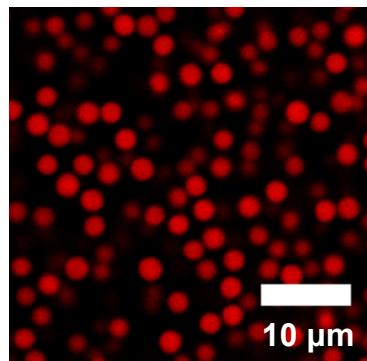


Translational and rotational dynamics

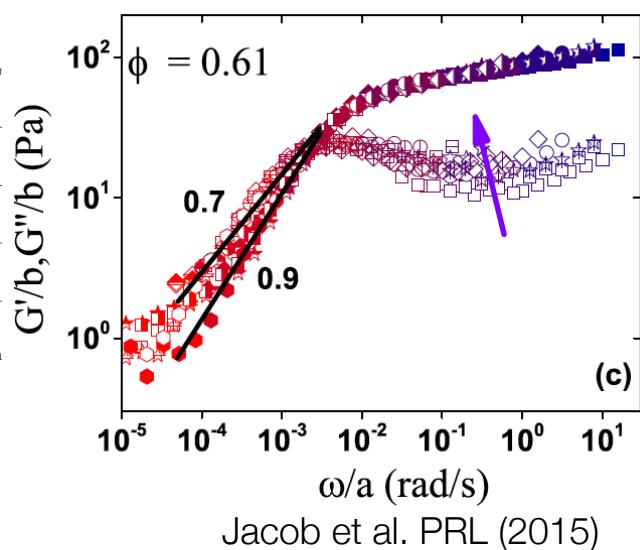


Mean squared displacement
Single particle correlations

Rheology and dynamics of colloidal glasses



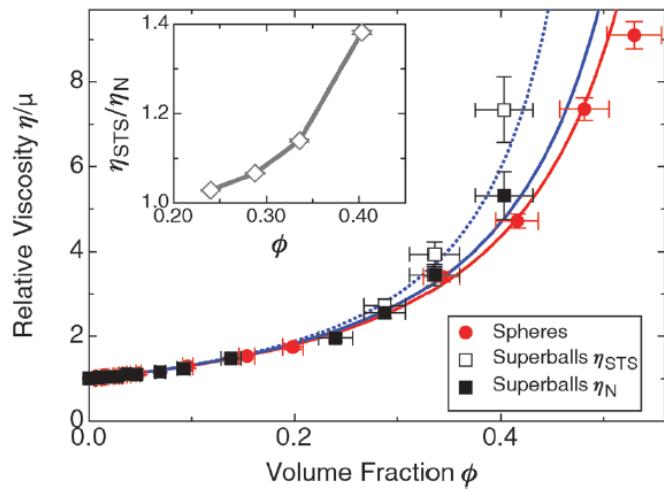
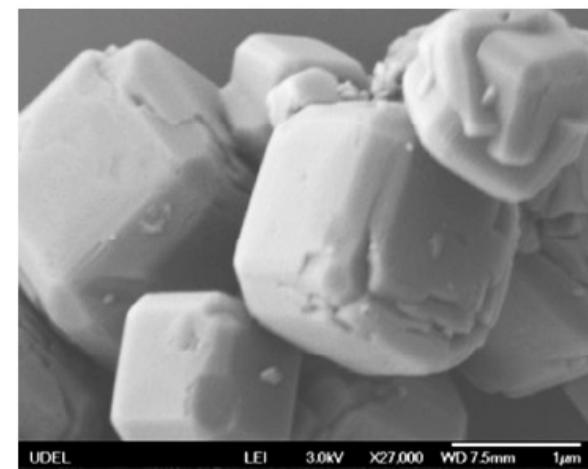
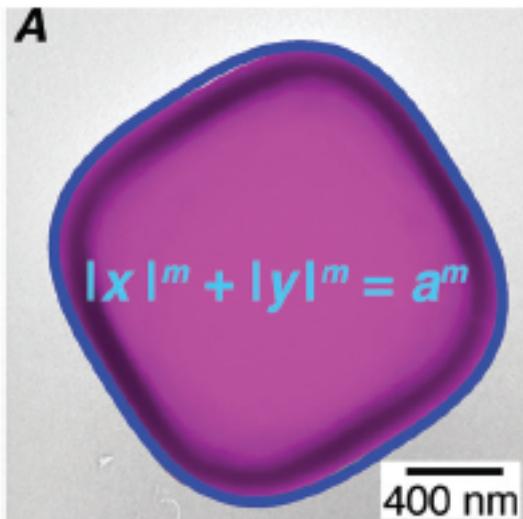
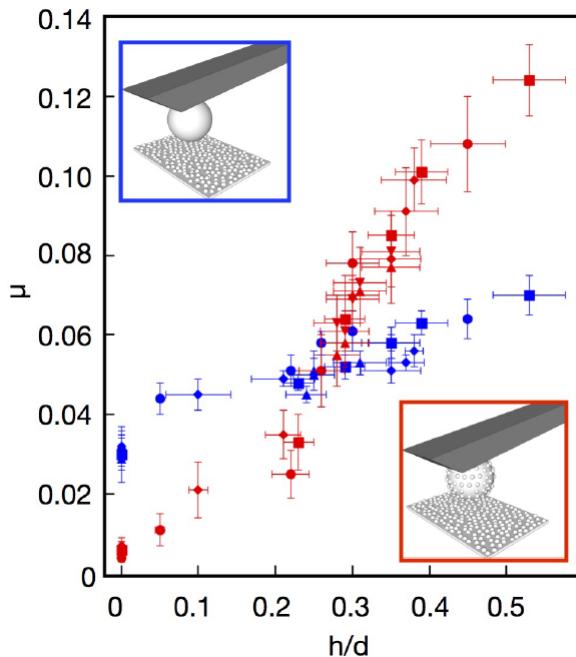
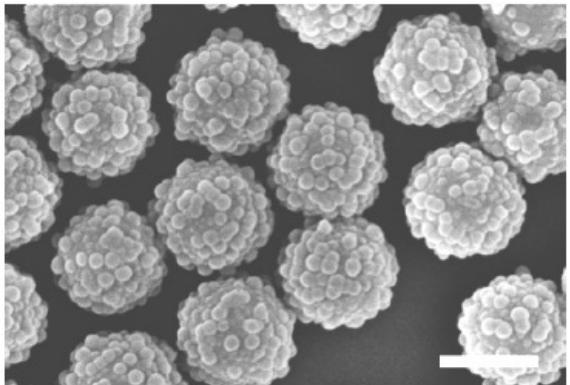
Sentjabrskaja et al. Sci Rep
(2015).



Jacob et al. PRL (2015)

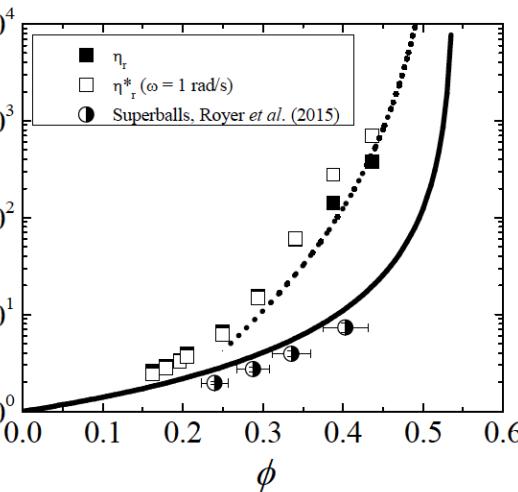
Two dynamic relaxation modes responsible
for cage + bond release under shear

Effect of surface anisotropy on dense suspensions



Hsu & Isa et al. PNAS (2018).

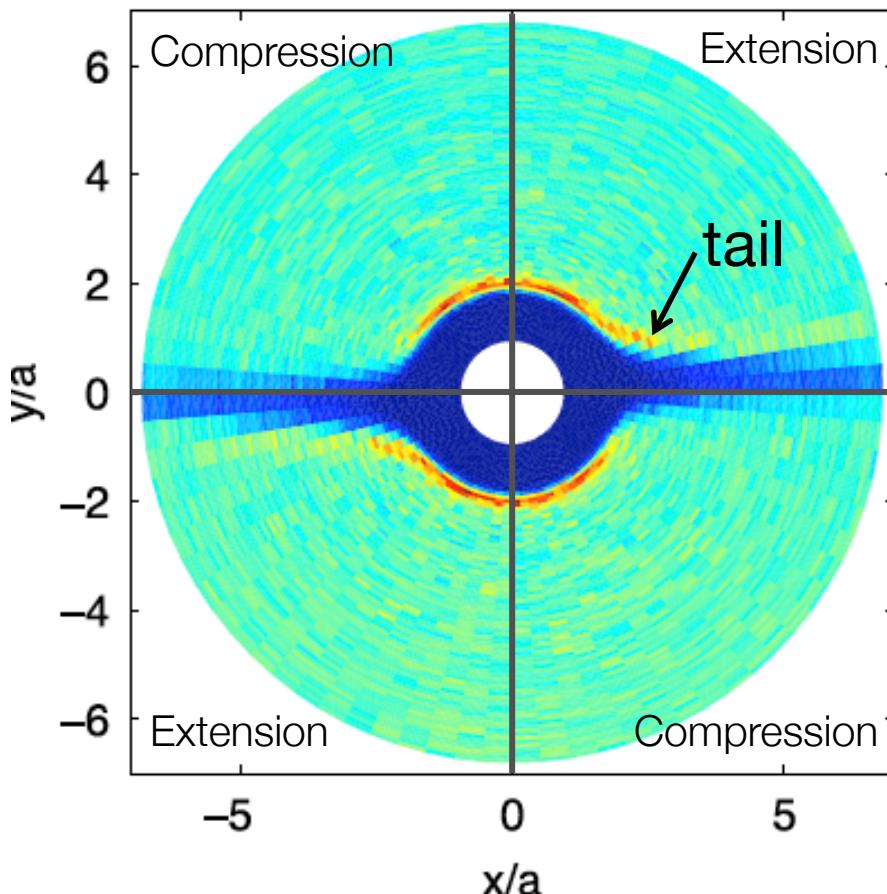
Royer, Blair, Hudson et al. Soft Matter (2015).



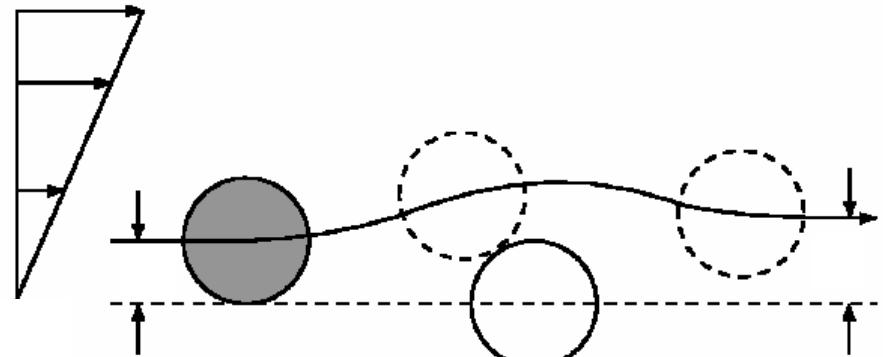
Cwalina & Wagner. Soft Matter (2016).

Roughness breaks fore-aft symmetry in simple shear

Experiments with non-Brownian PMMA ($\varphi = 0.05$) in Couette cell



Blanc et al. Phys Rev Lett (2011)

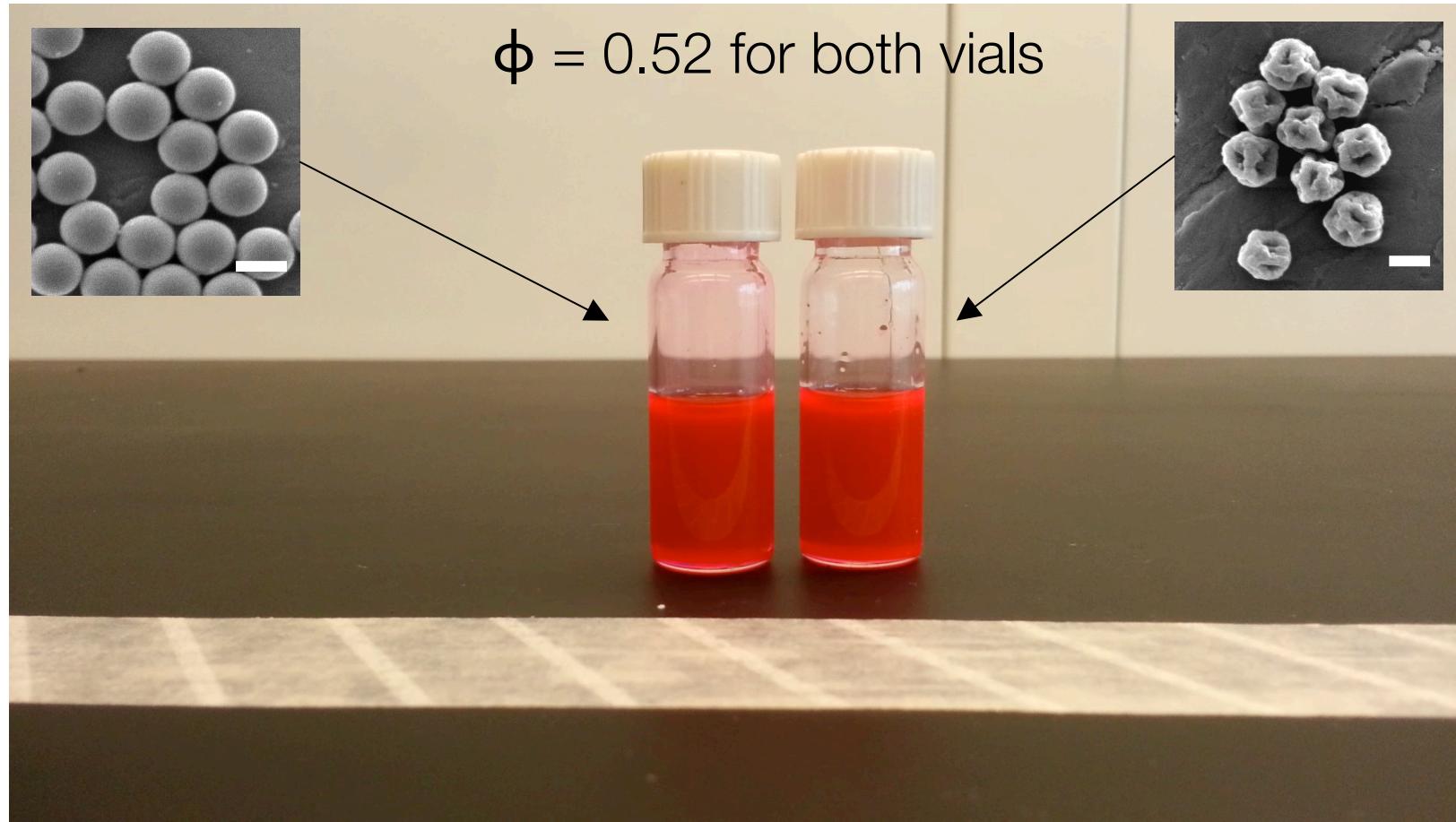


Davis et al. Phil Trans R Soc Lond A (2003)

Theoretical developments:

- Asperity prevents particle contact
- 2 limiting cases: slide (frictionless) and interlocking rigid bodies (frictional)
- Roughness brings Stokes flow irreversibility and symmetry breaking

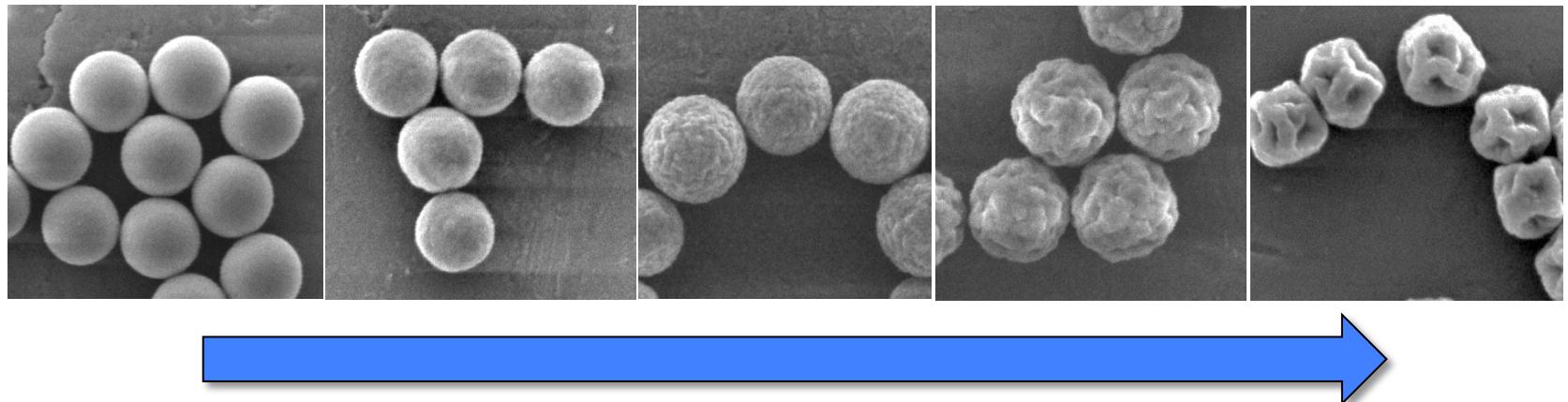
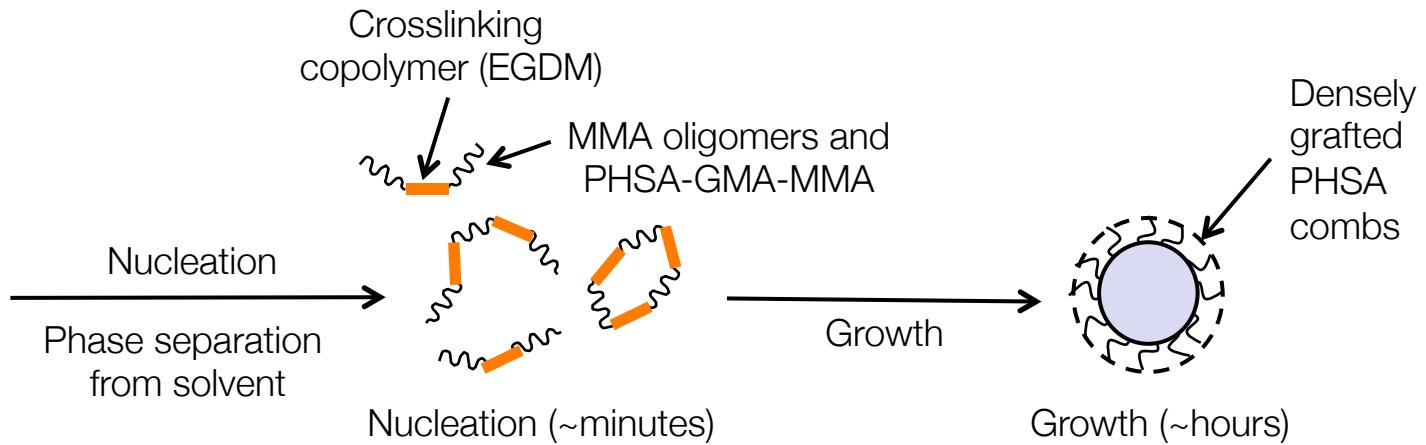
Particle roughness and shear thickening



Synthesis of PHSA-PMMA rough colloids

Free-radical dispersion polymerization

Initiator (AIBN)
 Monomer (MMA)
 Crosslinker (EGDM)
 Stabilizer (PHSA-GMA-MMA)

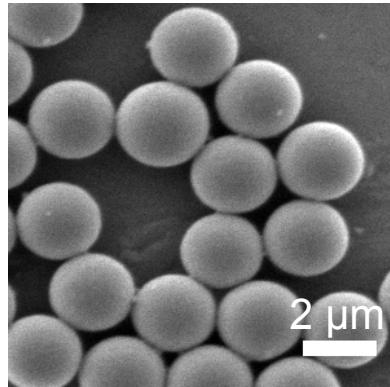


Increasing crosslinker concentration, Increasing roughness

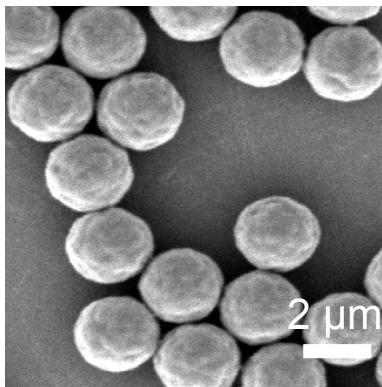
Colloidal particles of varying roughness

PHSA-stabilized PMMA colloids, $2a_{\text{eff}} = 1.9$ to $2.5 \mu\text{m}$, Size polydispersity = 3-4%

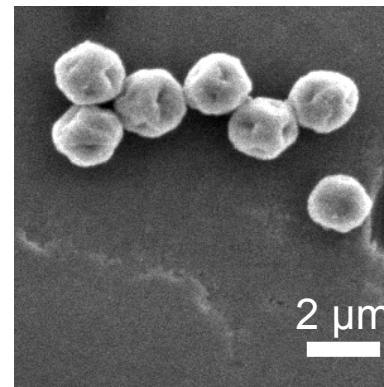
Smooth



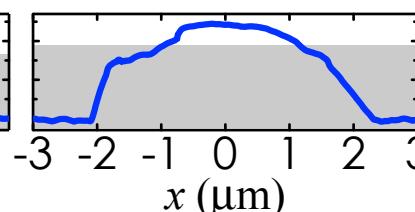
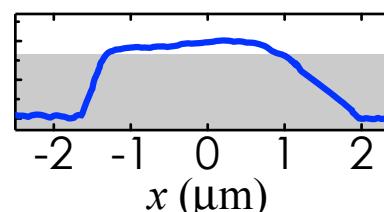
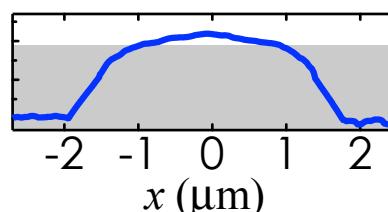
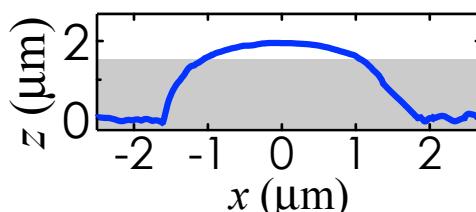
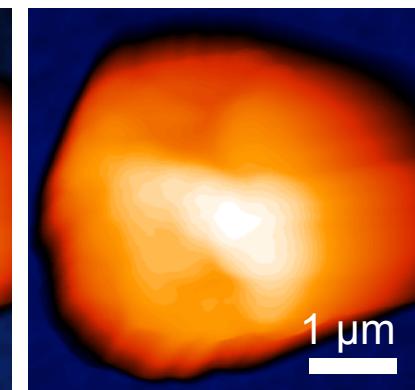
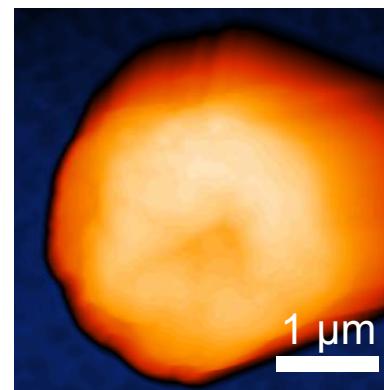
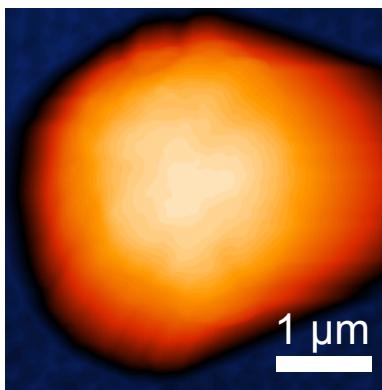
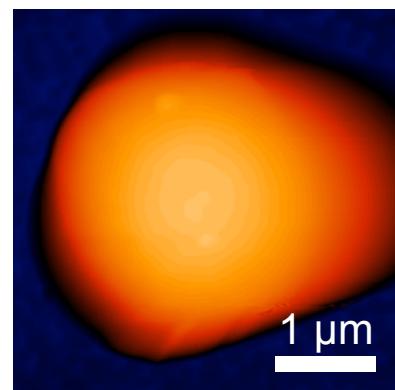
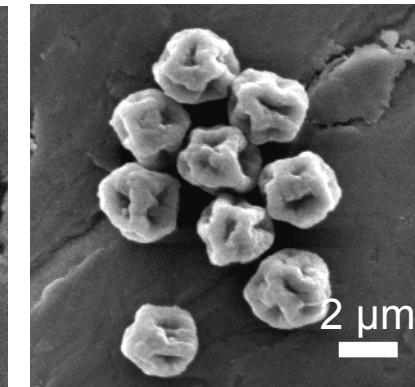
Slightly rough



Medium rough



Very rough

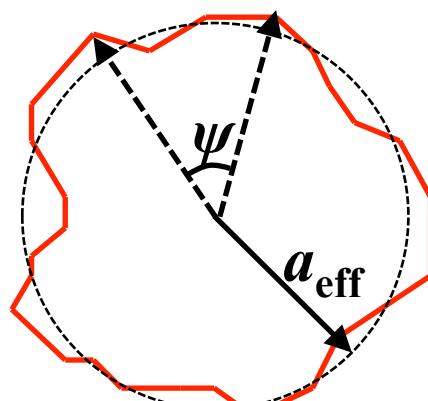
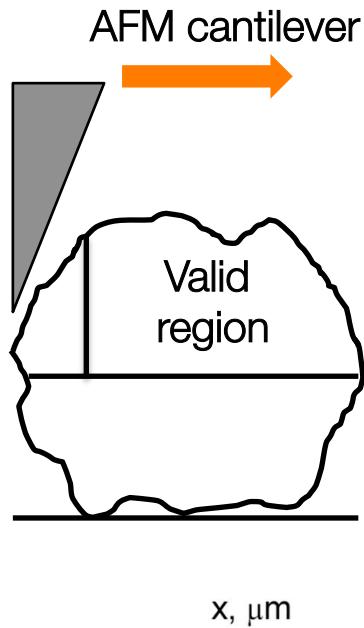


0.00

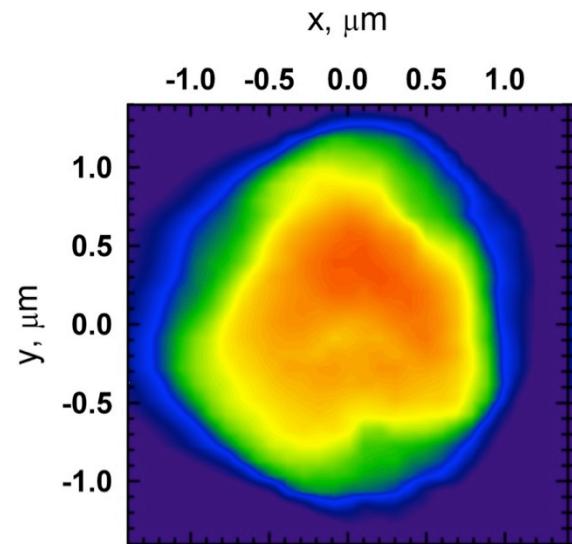
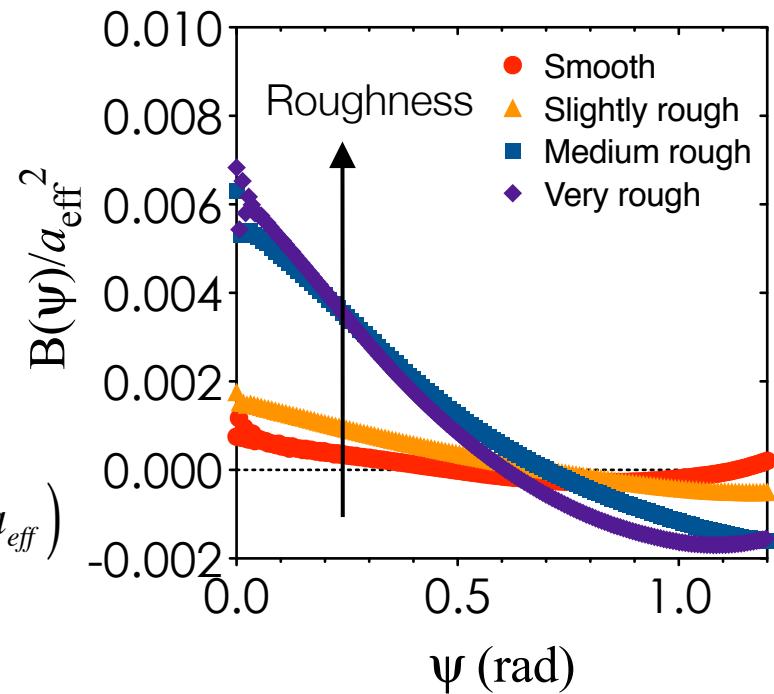
2.70

$z (\mu\text{m})$

Extracting RMS roughness from AFM measurements



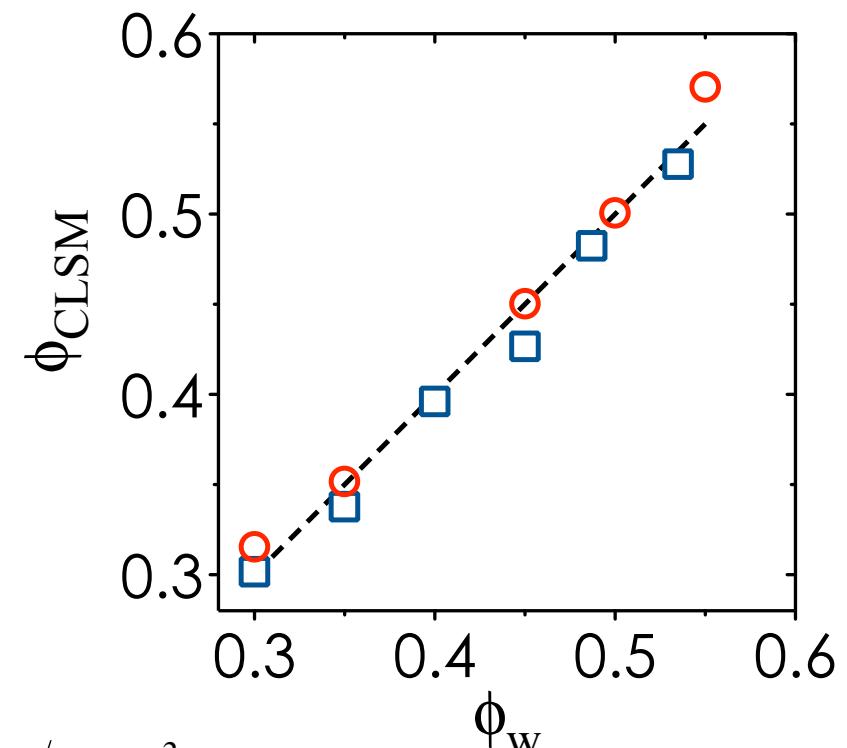
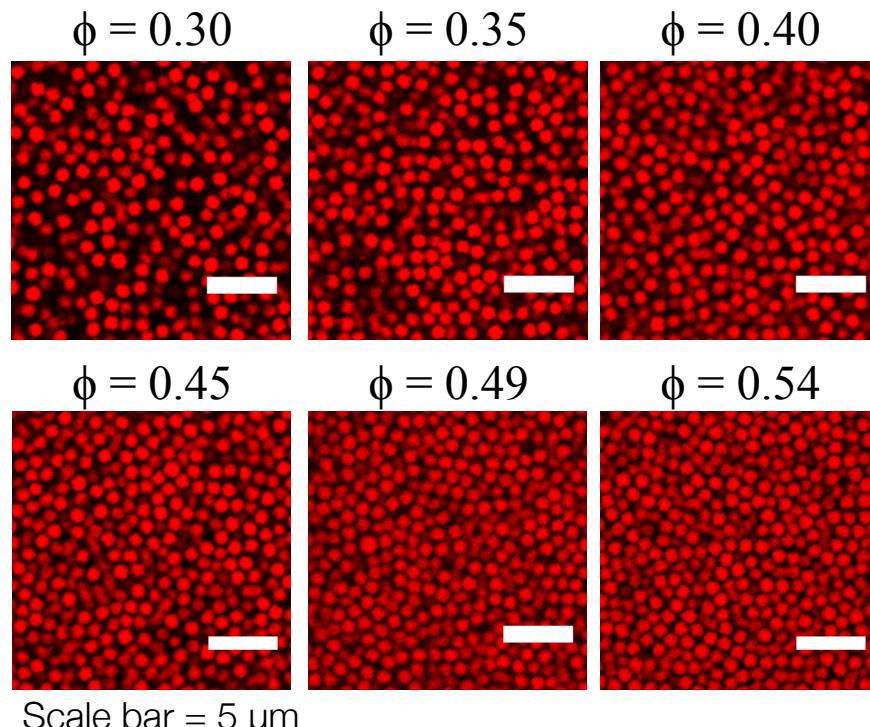
$$B(\psi) = \frac{1}{N} \sum_{\substack{i,j=1 \\ i \neq j}}^{\psi} (|r_i| - a_{\text{eff}})(|r_j| - a_{\text{eff}})$$



Geometry	$2a_{\text{eff}}$ (μm)	$(B(\psi = 0)/a_{\text{eff}}^2)^{1/2}$	Range of ϕ
Smooth	$2.27 \mu\text{m} \pm 5\%$ $1.60 \mu\text{m} \pm 4\%$	0.026 ± 0.003	0.30 to 0.55
Slightly rough	$2.55 \mu\text{m} \pm 2\%$	0.040 ± 0.002	0.45 to 0.55
Rough	$1.95 \mu\text{m} \pm 5\%$ $2.06 \mu\text{m} \pm 4\%$ $2.47 \mu\text{m} \pm 5\%$ $1.91 \mu\text{m} \pm 5\%$	0.075 ± 0.005	0.30 to 0.55
Very rough	$2.78 \mu\text{m} \pm 6\%$	0.082 ± 0.003	0.45 to 0.50

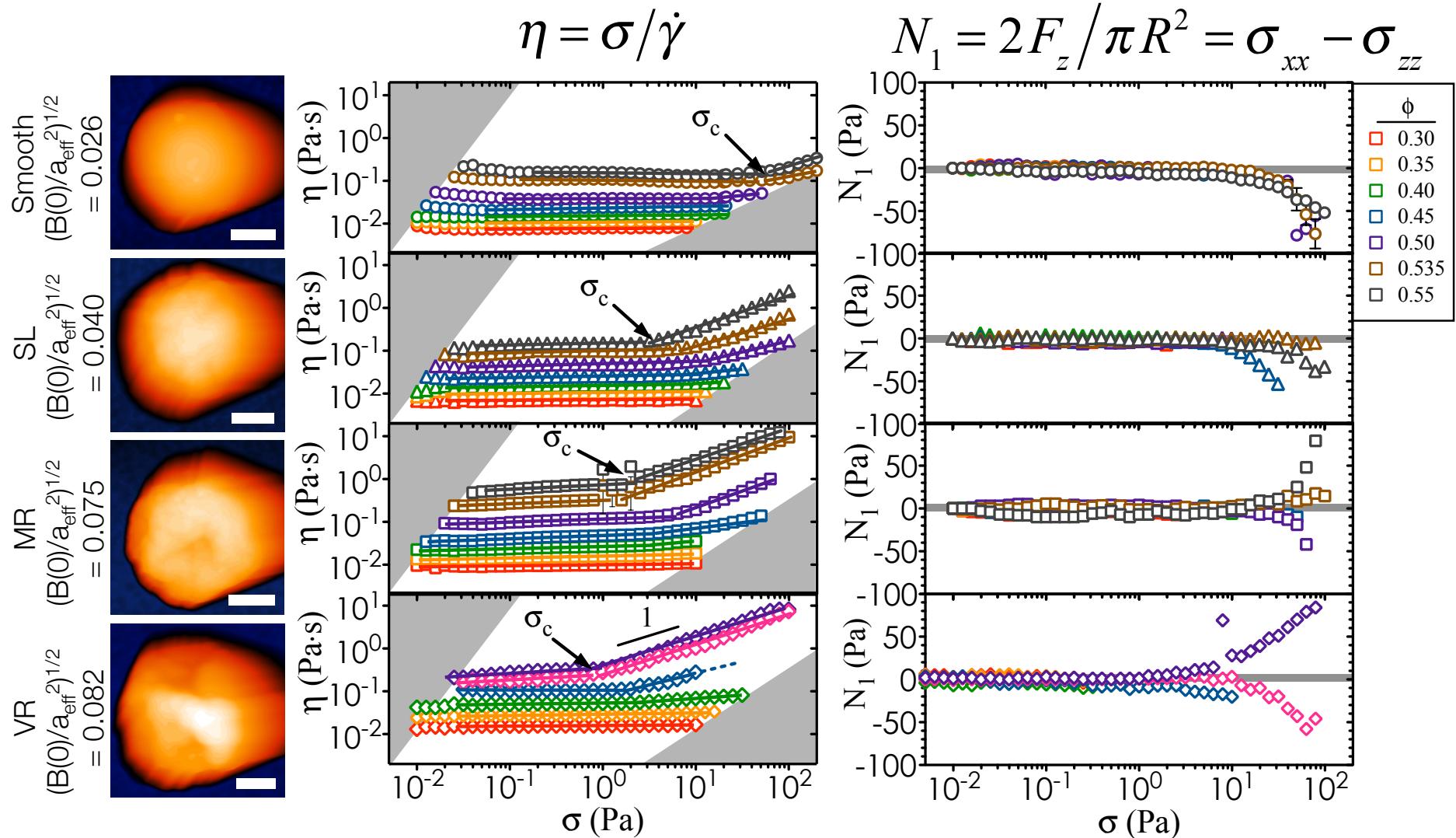
Calculation of the volume fraction

Effective size from AFM/SEM, particle counting from 3D confocal microscopy



$$\phi_{\text{CLSM}} = \frac{4/3\pi a_{\text{eff}}^3 N_p}{V_{\text{box}}}$$

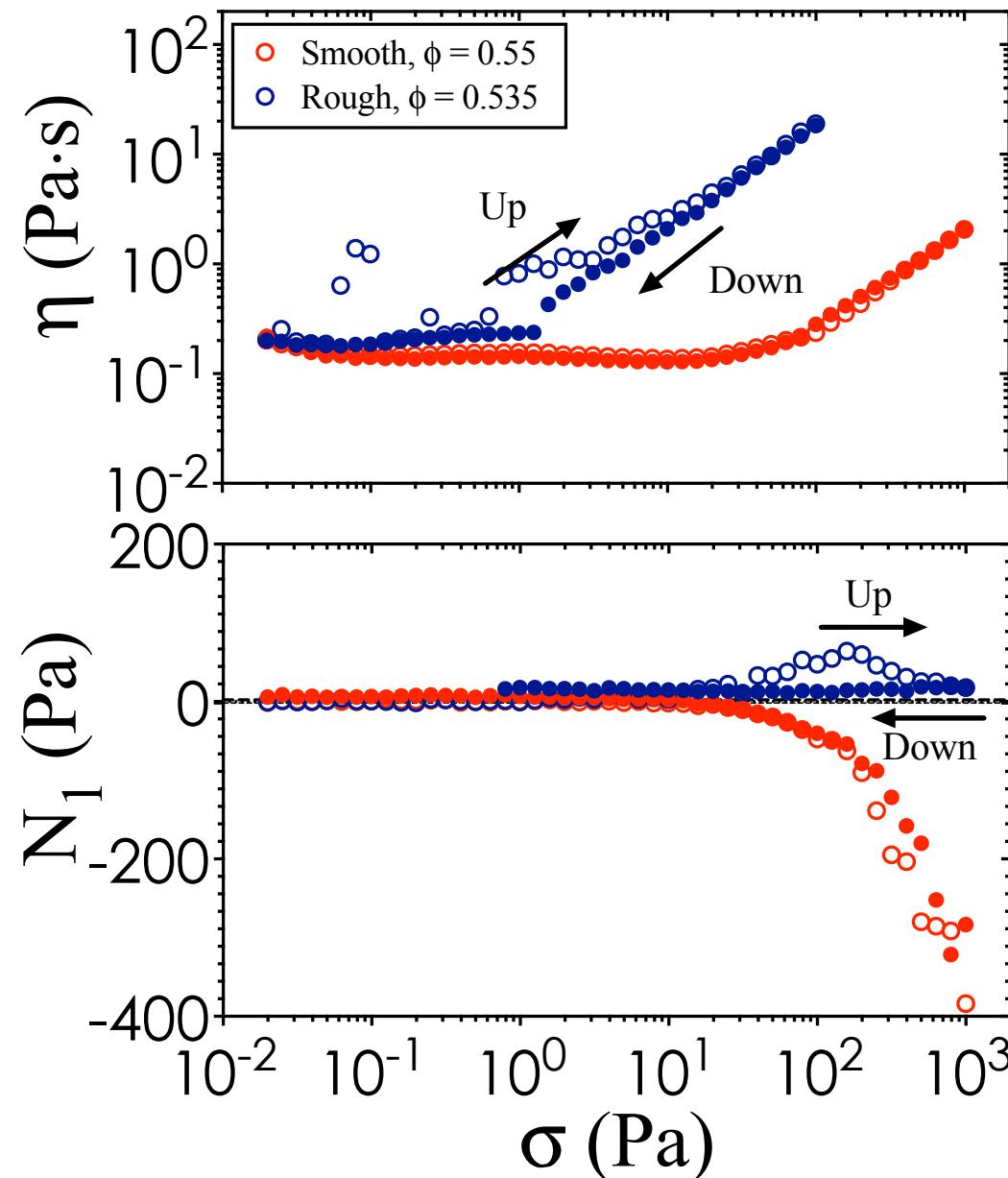
Steady state viscosity and first normal stresses



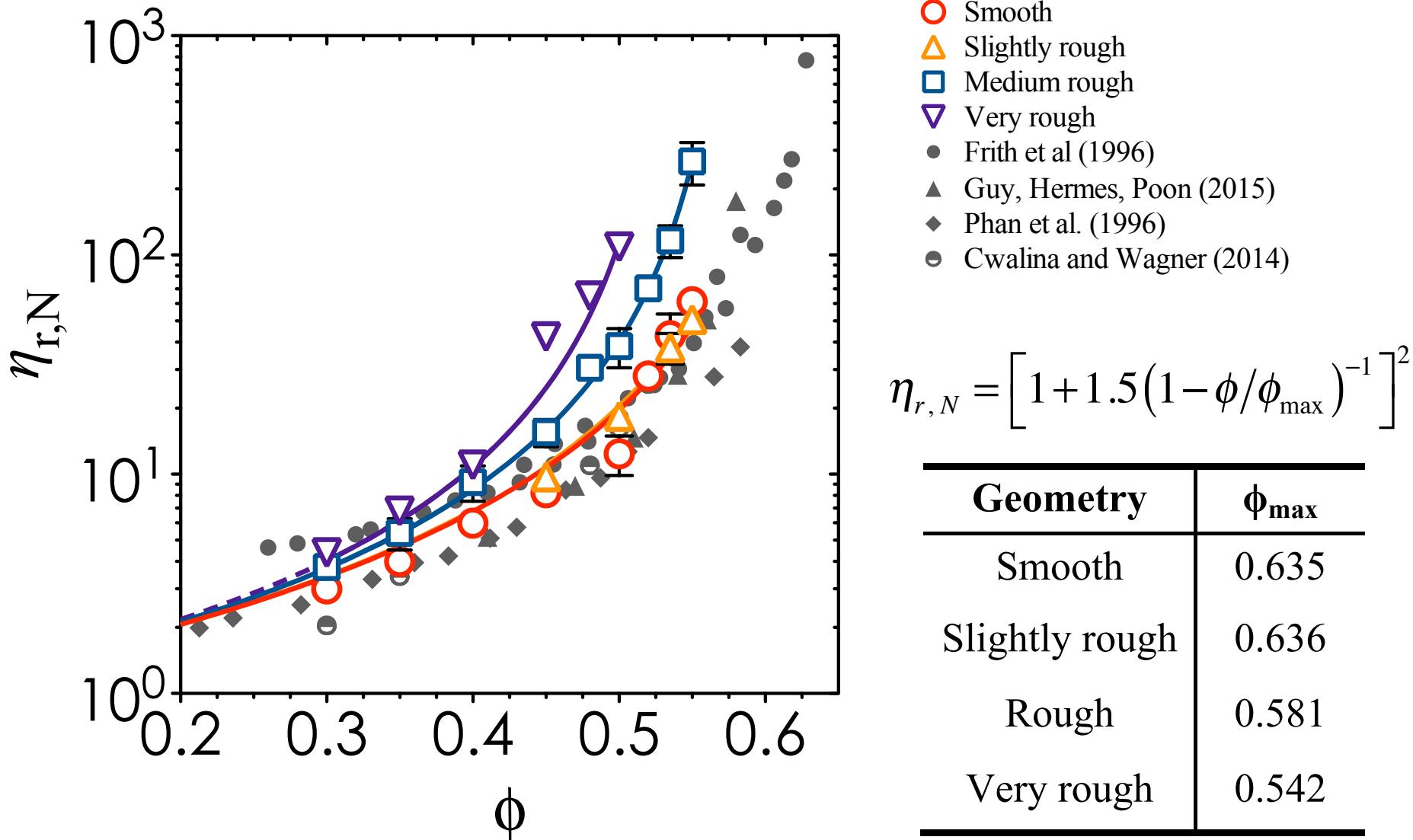
$$St = \frac{\rho_P a^2 \dot{\gamma}}{\eta_f}$$

$$10^{-7} \leq St \leq 10^{-3}$$

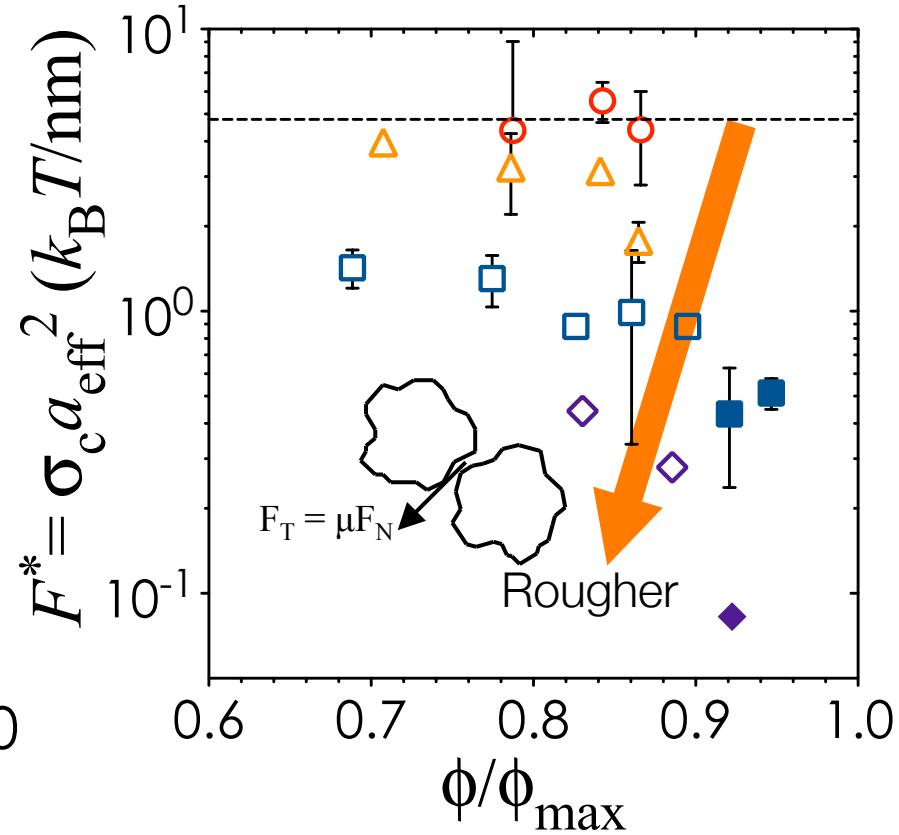
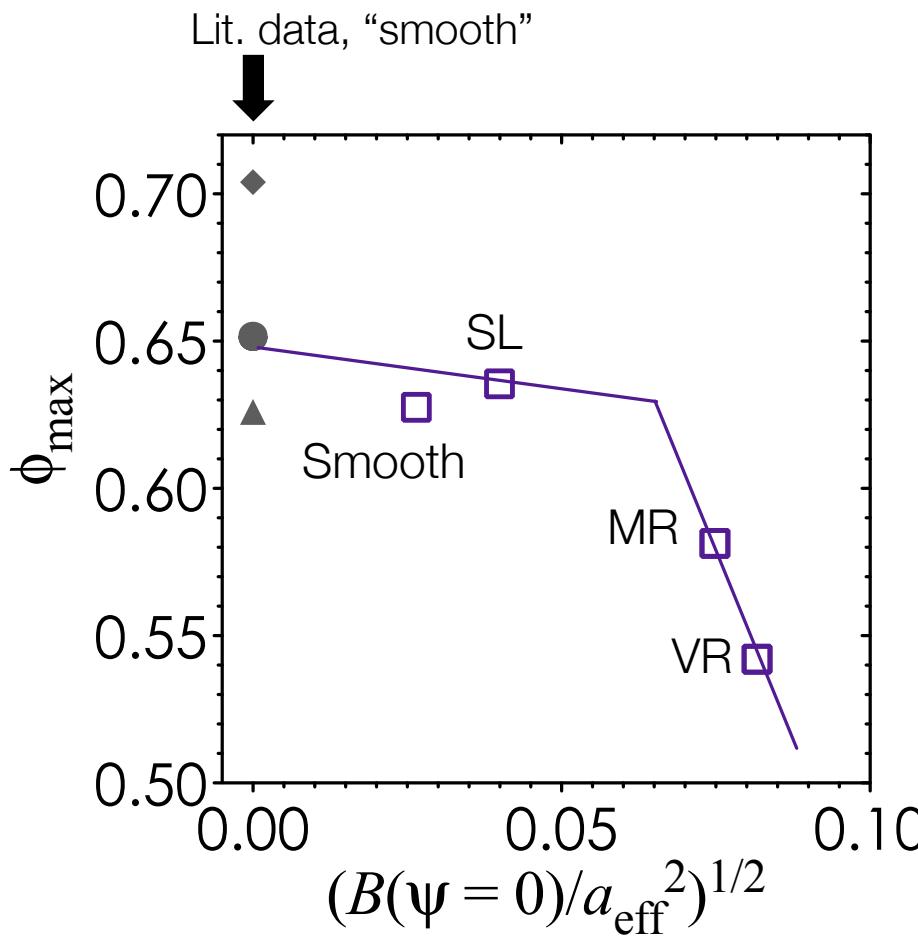
Hysteretic flow curves for frictional particles



Fitting the high-shear viscosity to the Eilers model



Effect of roughness on excluded volume

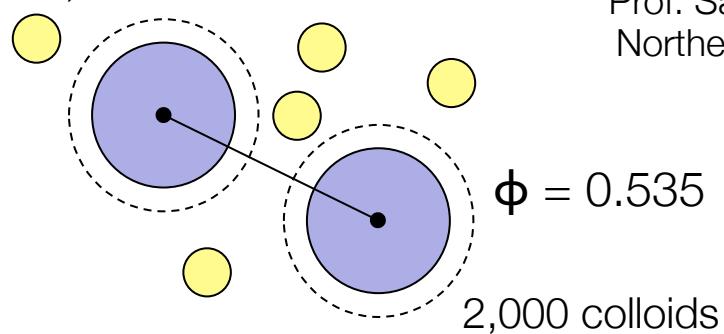


- When RMS roughness increases, lubrication becomes less effective in keeping particles apart
- It takes less force to push them into contact

Comparison with DPD simulations

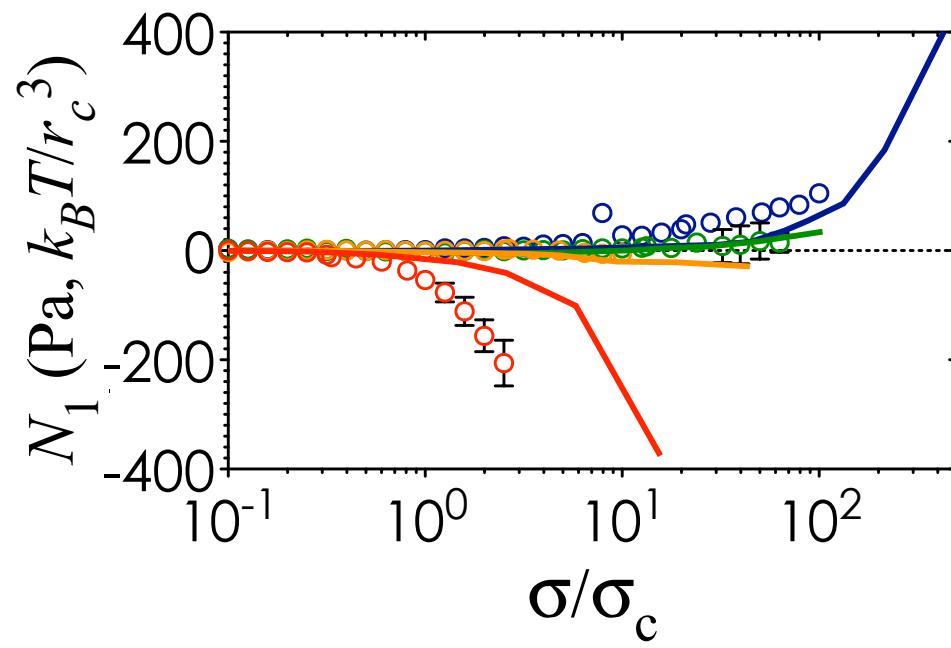
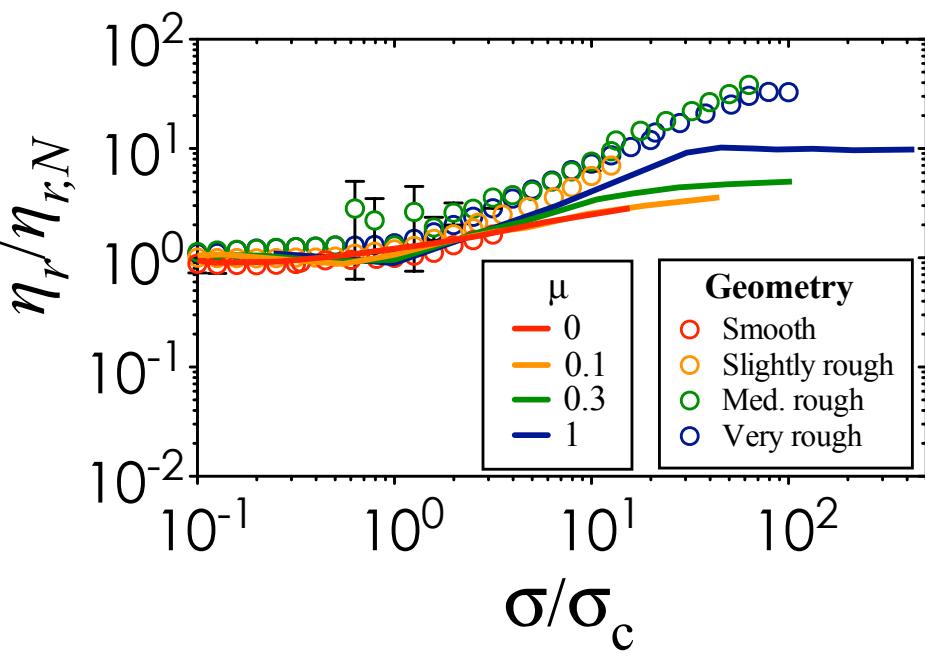
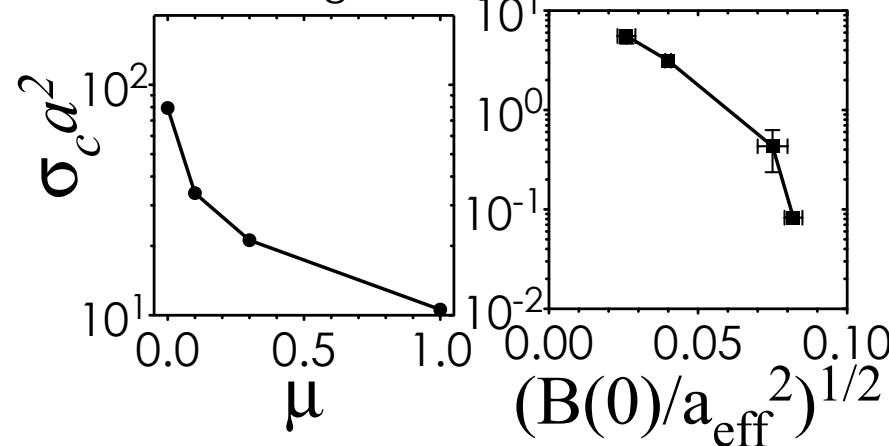
$$m \frac{dv}{dt} = f(F_T, F_N, F_h, F_{solvent})$$

50,000 solvent molecules

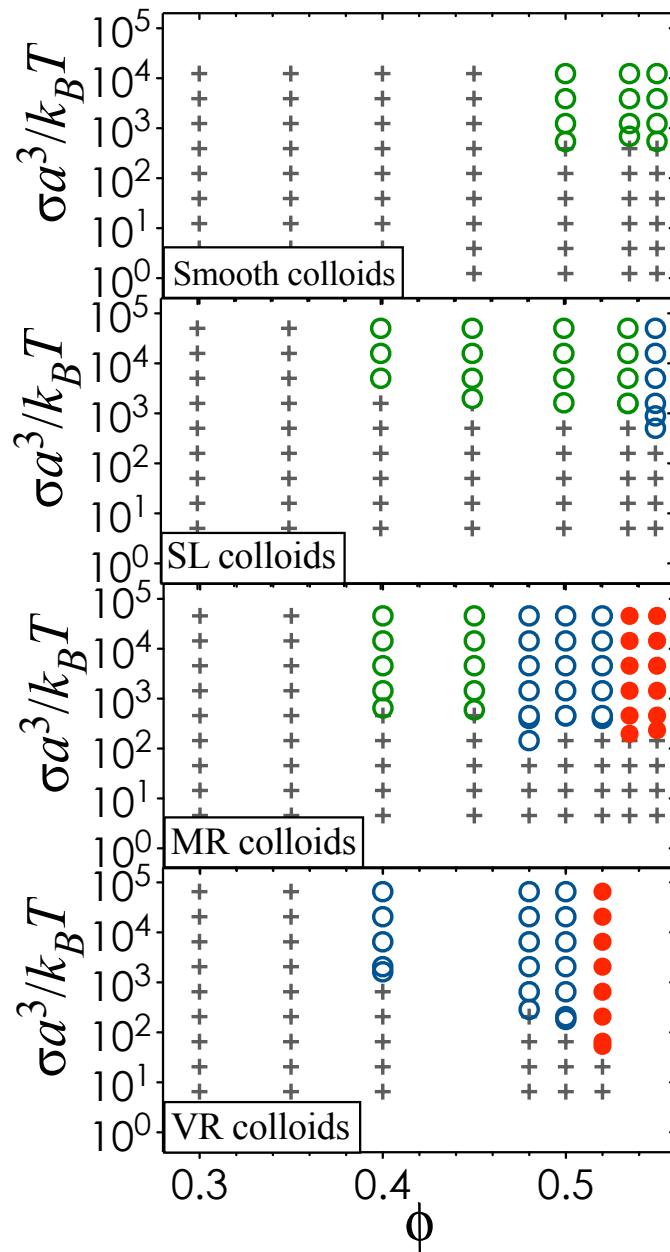


Prof. Safa Jamali
Northeastern U.

Onset stress depends on roughness and friction



State diagrams for rough colloids in shear flow

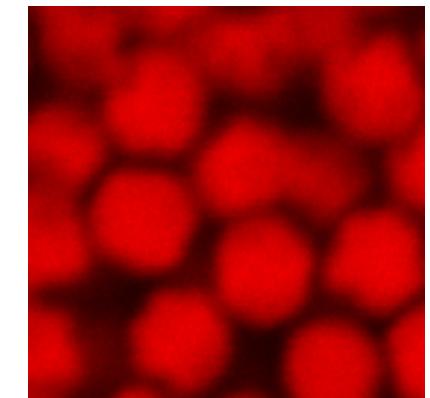


Engineering approach: Assume particle is a sphere.

But !

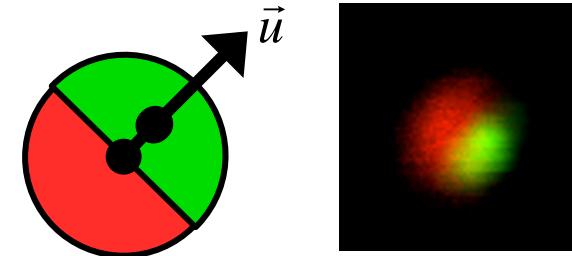
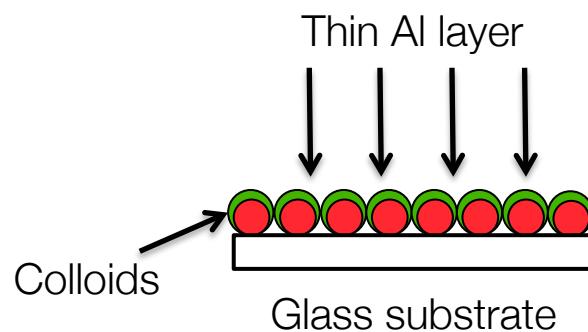
What is the microscopic reason?

- ⊕ = Newtonian flow
- = weak thickening
- = strong thickening
- = dilatant ($N_1 > 0$)



Measuring the 3D rotational dynamics of colloids

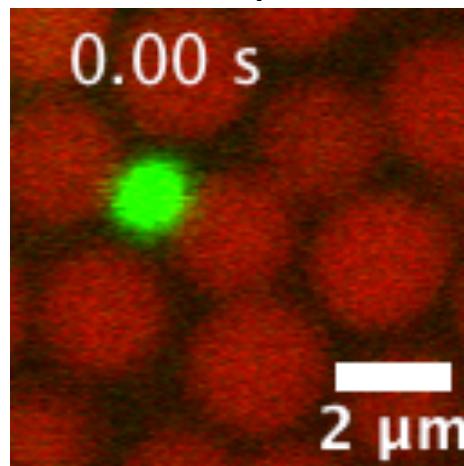
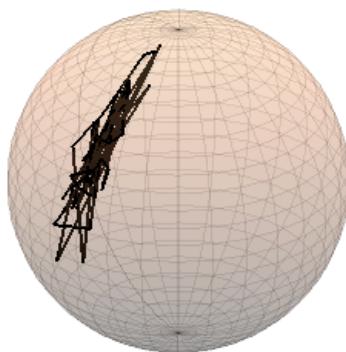
Inert Janus tracers to track rotational dynamics



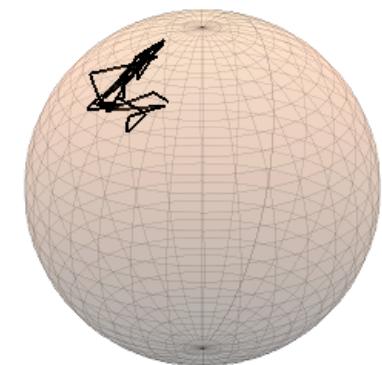
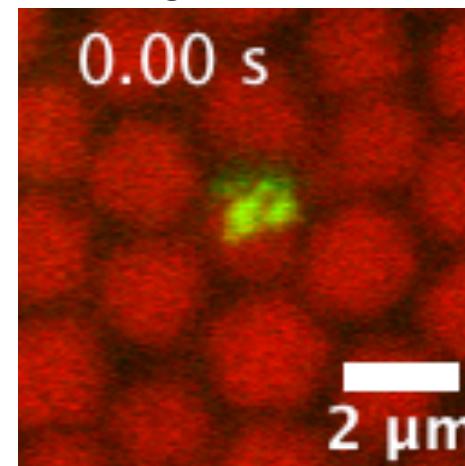
Mean-squared angular displacement

$$\langle \Delta \vec{\phi}^2(\Delta t) \rangle = \left\langle [\vec{\phi}(t + \Delta t) - \vec{\phi}(t)]^2 \right\rangle$$

Smooth, $\phi = 0.50$

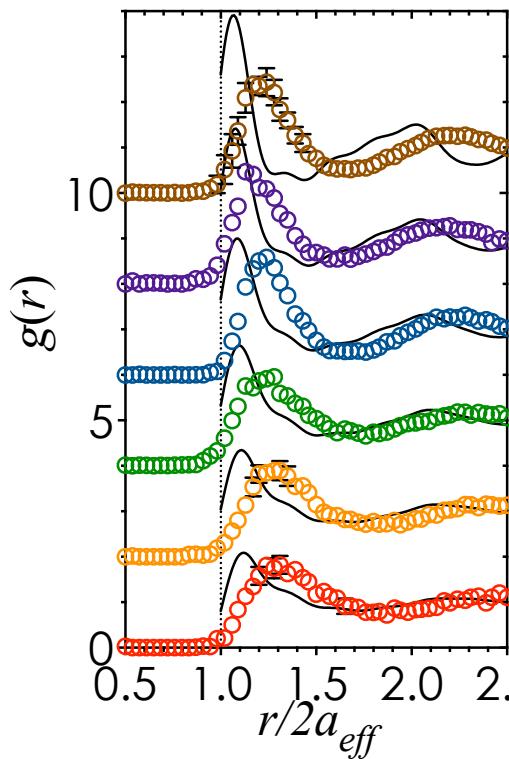
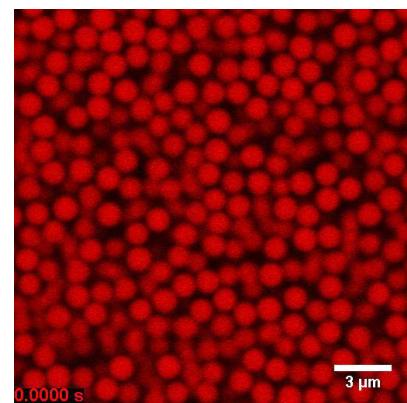


Rough, $\phi = 0.50$



Accounting for repulsion in colloidal suspensions

Smooth



Yukawa electrostatics

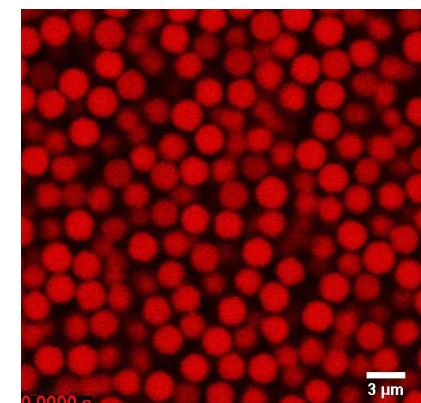
$$U(r) = \pi \epsilon_0 \epsilon \psi^2 (2a)^2 \exp[-\kappa(r - 2a)] / r$$

$$c(x) = \begin{cases} A + Bx + \frac{1}{2} \eta Ax^3 + \frac{C \sinh kx}{x} + \frac{F(\cosh kx - 1)}{x}, & x = \frac{r}{2a} < 1 \\ -\gamma \exp(-kx) / x, & x > 1 \end{cases}$$

$$g(x) = 1 + \frac{1}{12\pi\eta x} \int_0^\infty [S(K) - 1] K \sin Kx dK$$

$\phi = 0.55$

Rough

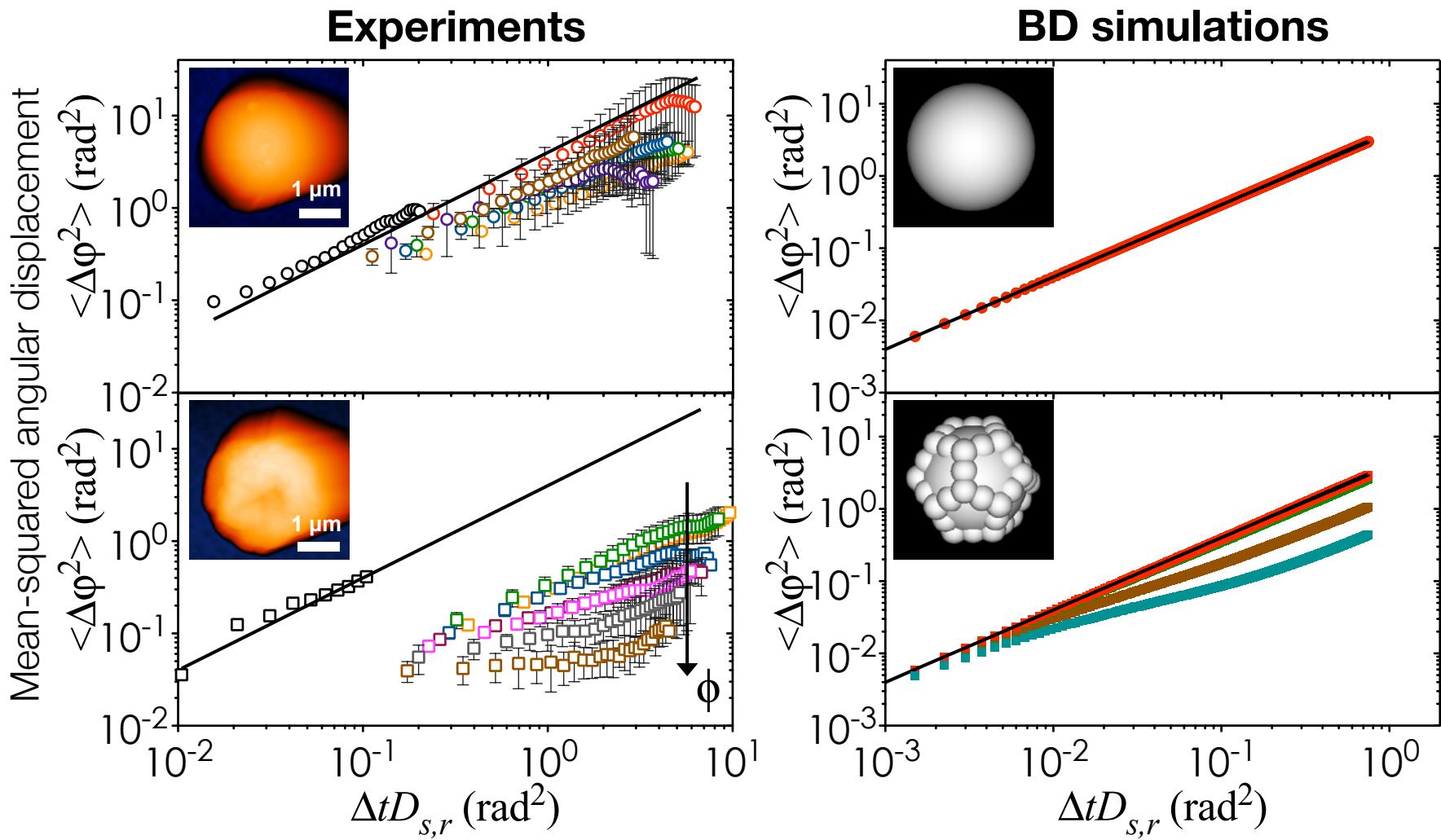


$\phi = 0.30$

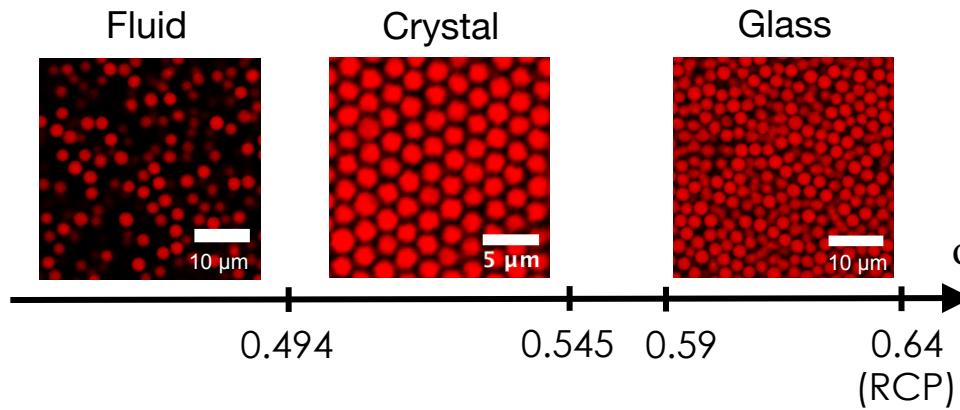
Mean spherical approximation closure
to the Ornstein-Zernicke equation

Hayter and Penfold, Mol Phys (1981)

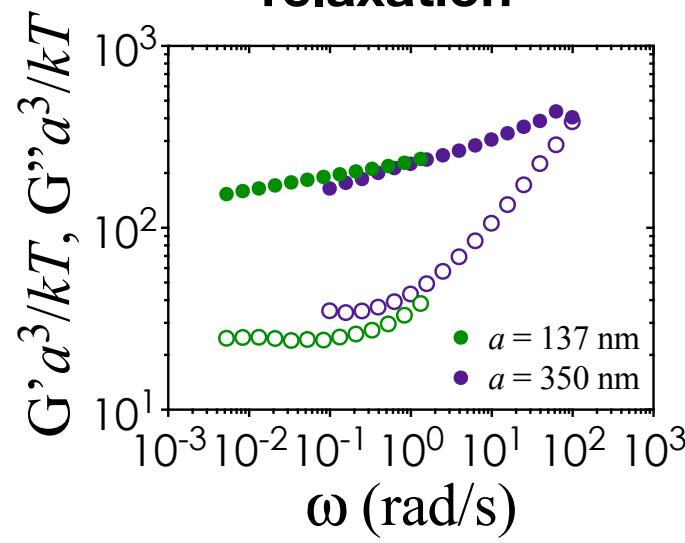
Hindered rotational dynamics from surface geometry



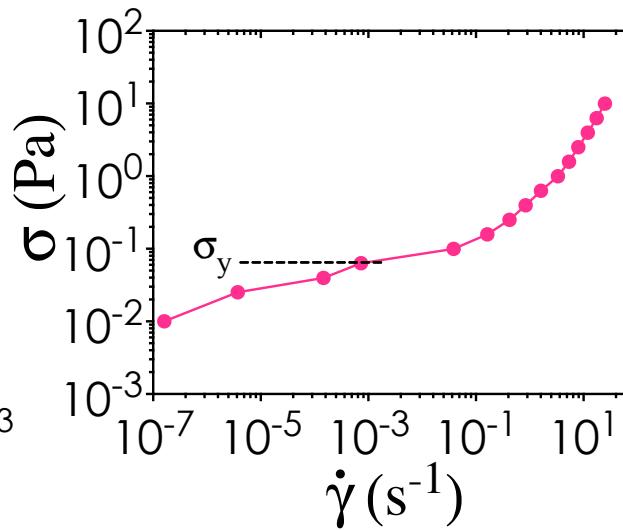
Rheological properties of colloidal glasses



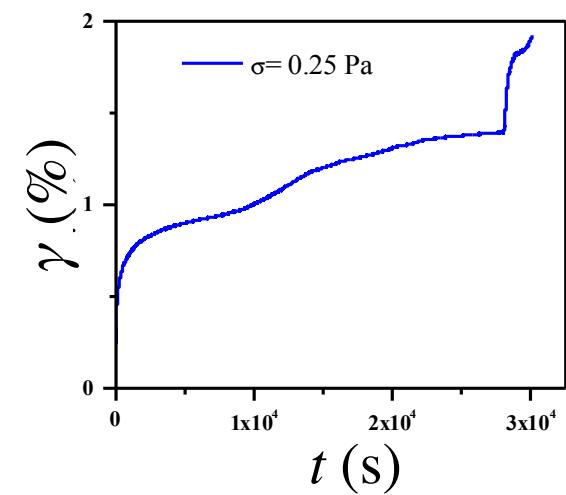
Viscoelastic 2-step relaxation



Yield stress behavior



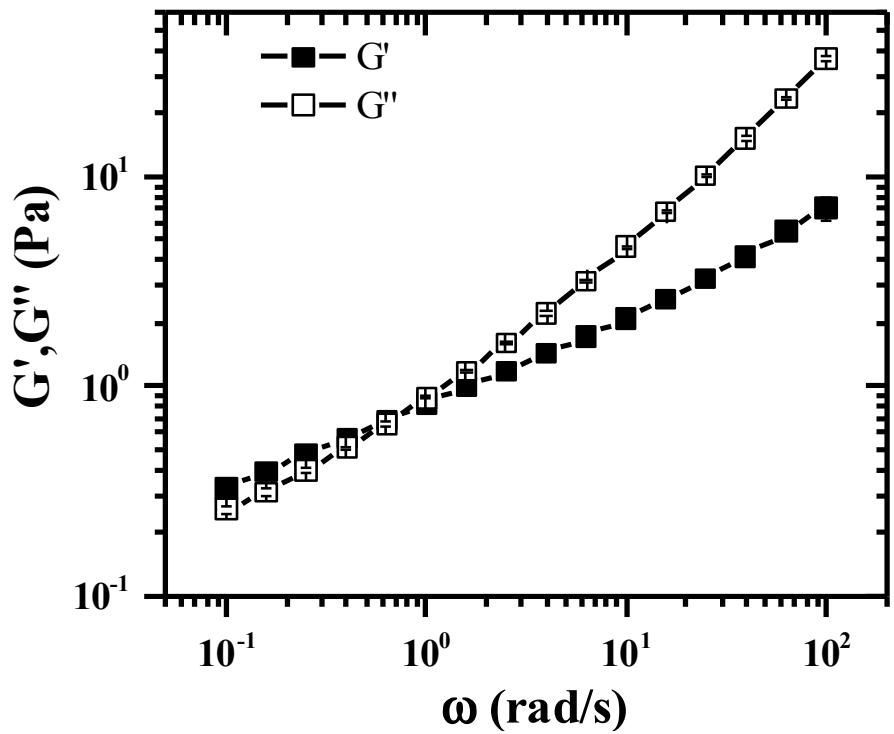
Creep near σ_s



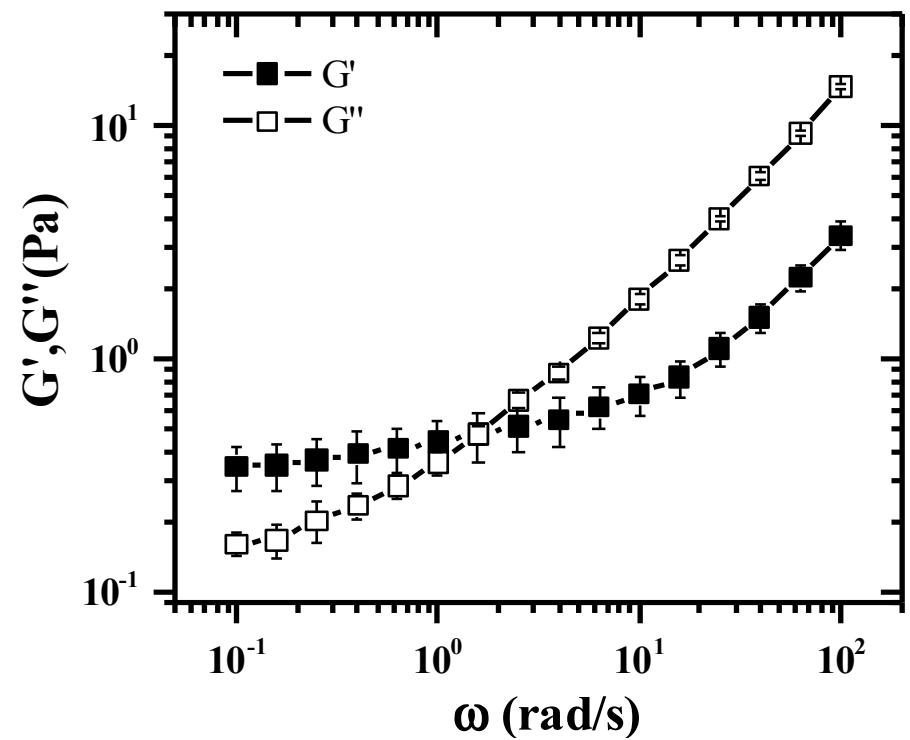
Courtesy: Alan Jacob

Viscoelastic comparison

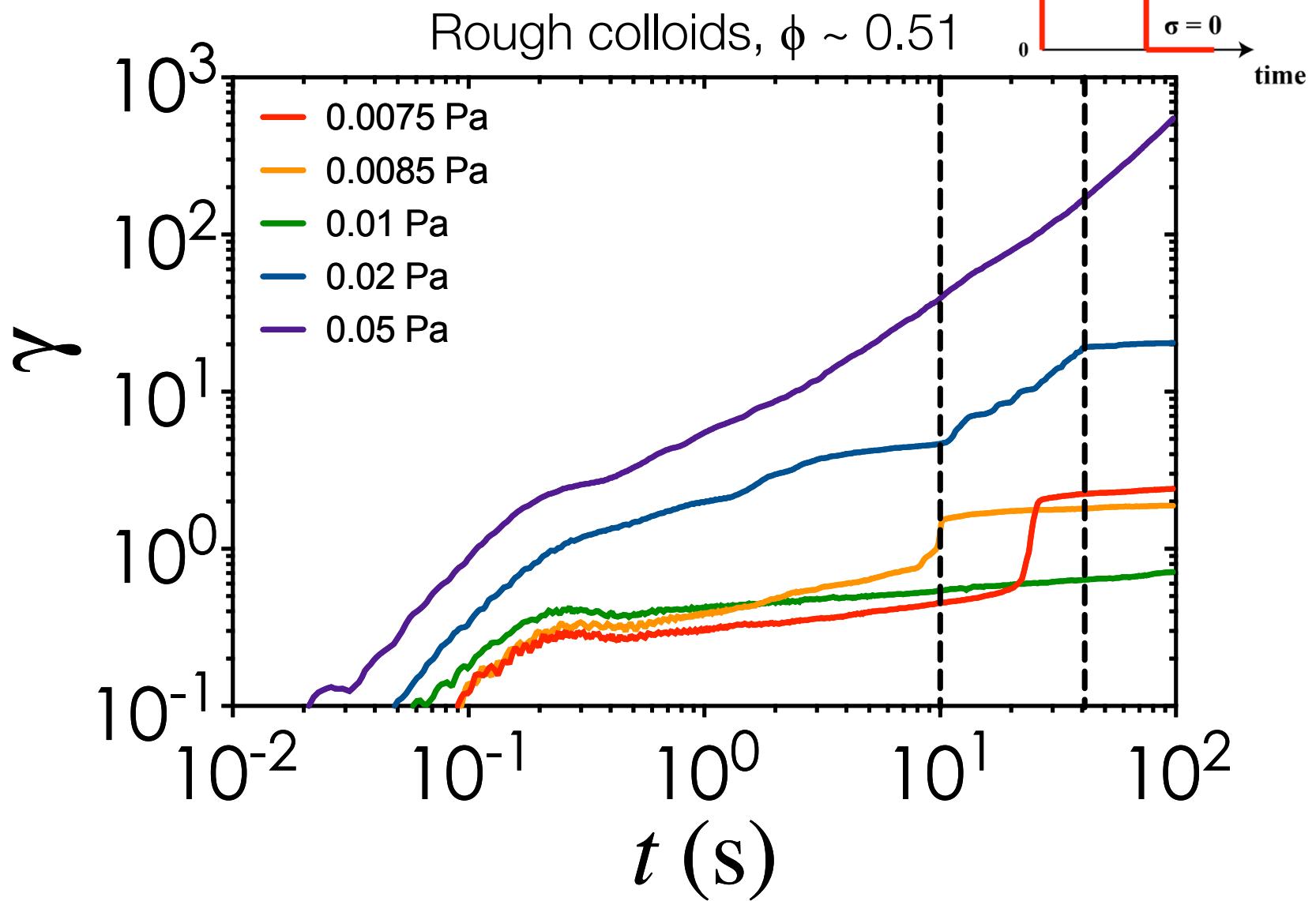
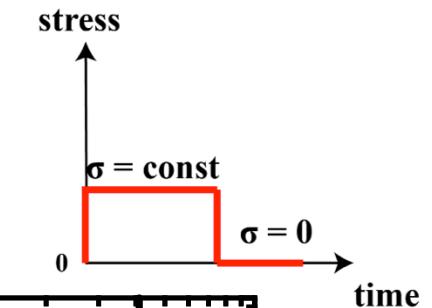
Smooth, $\phi = 0.64$



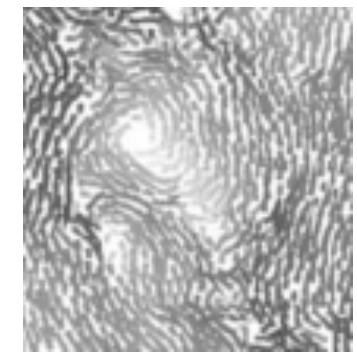
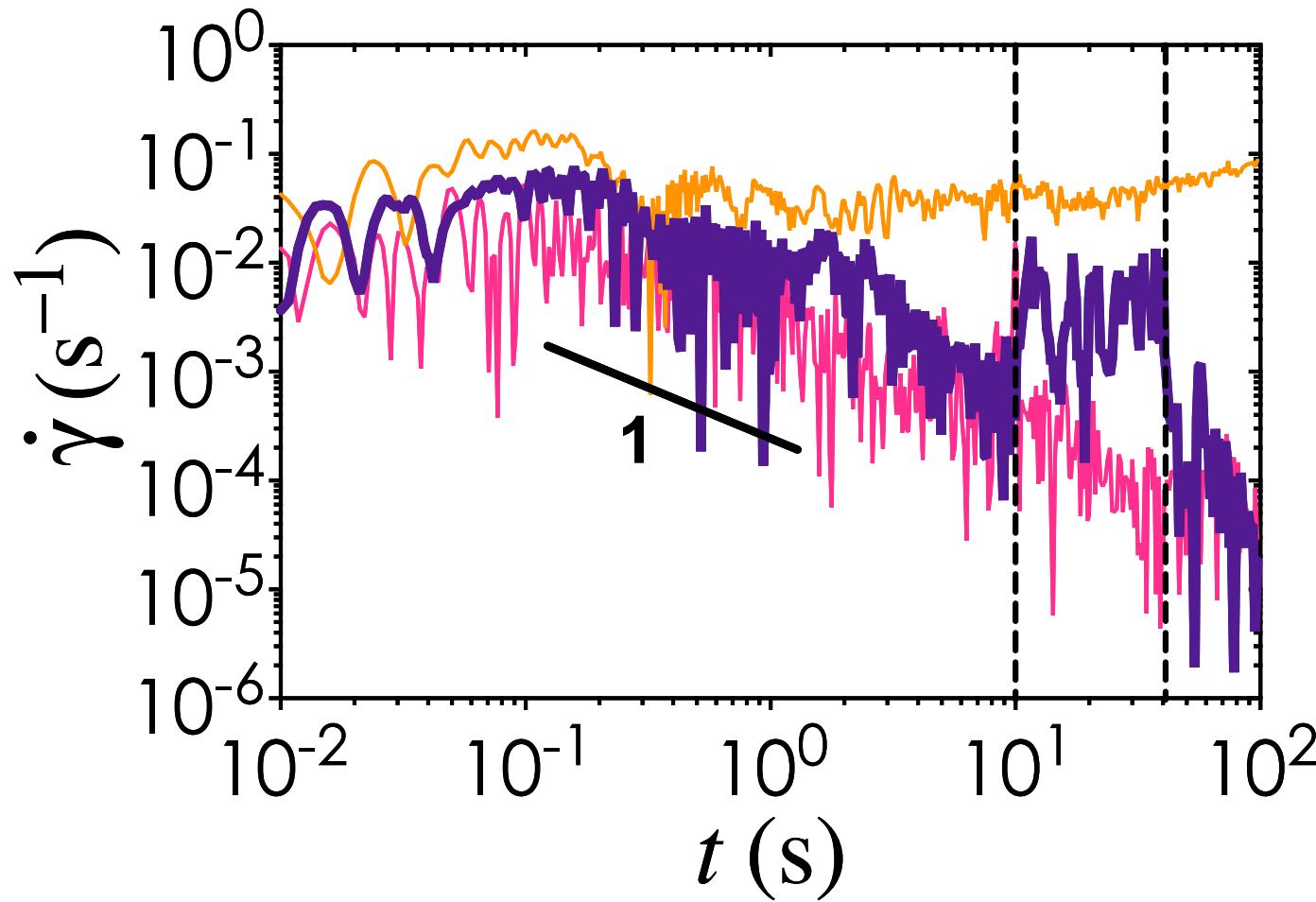
Rough, $\phi = 0.55$



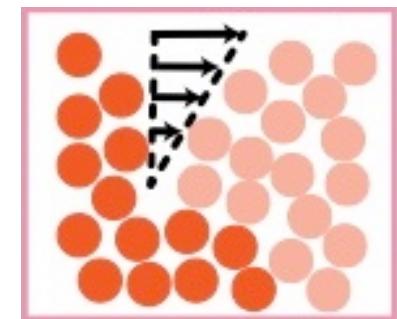
Creep under step stress



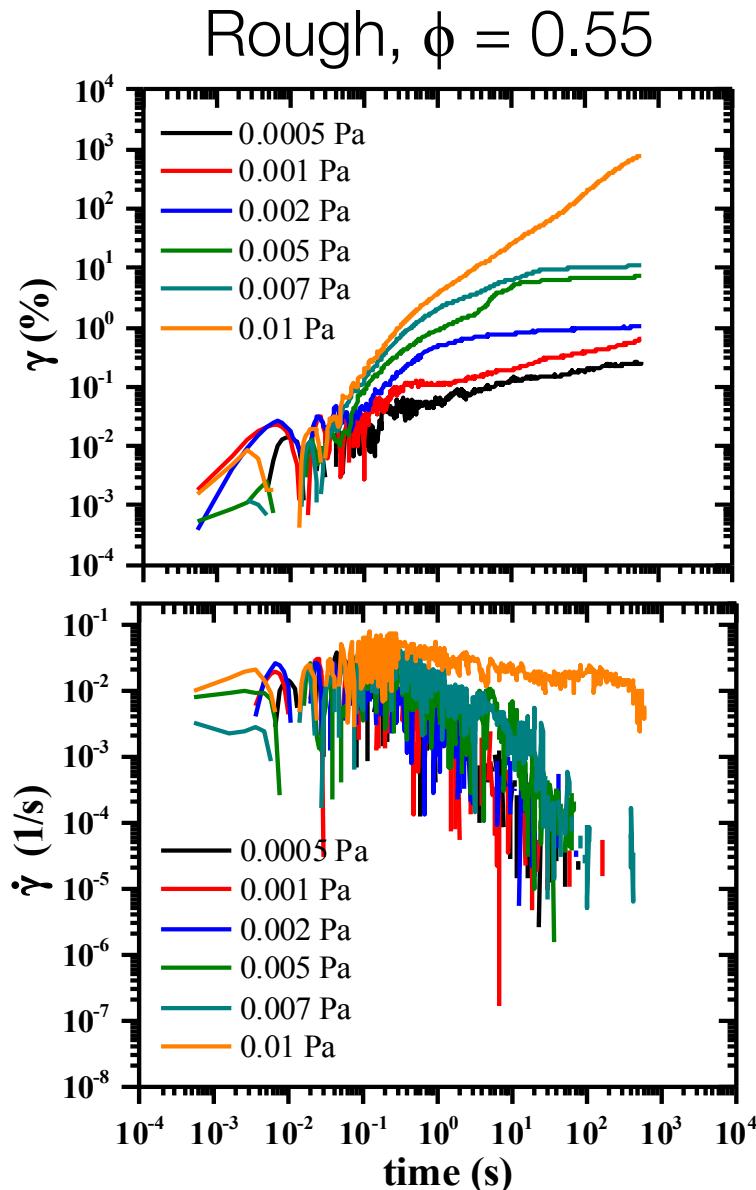
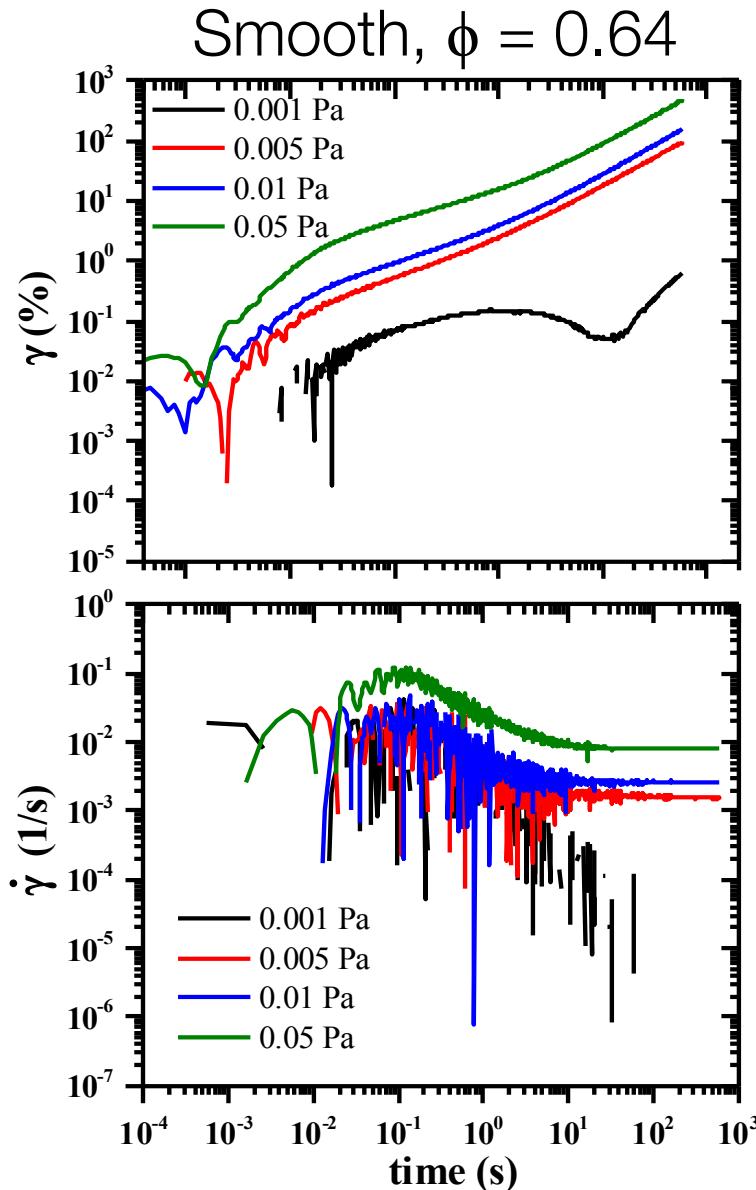
Creep shear rate showing transient fluidization



Maloney & Lemaitre.
PRE (2006).

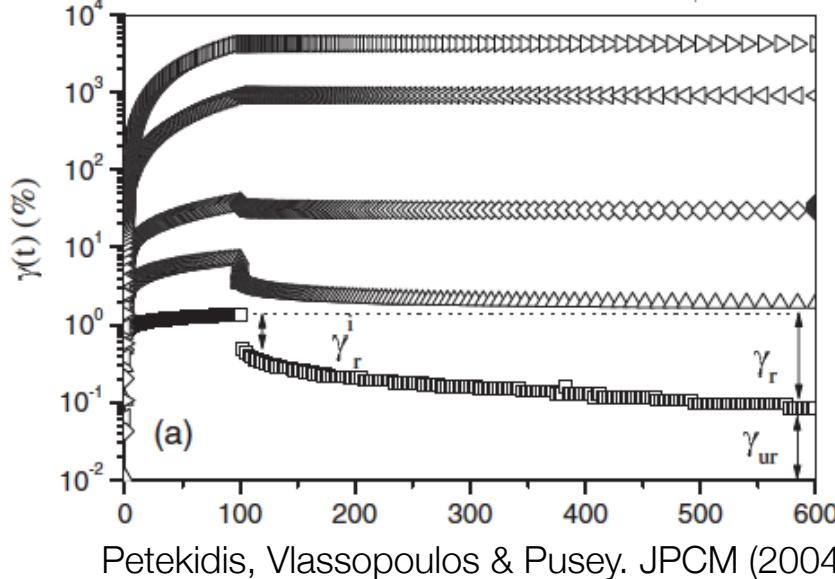


Comparisons between smooth and rough colloids



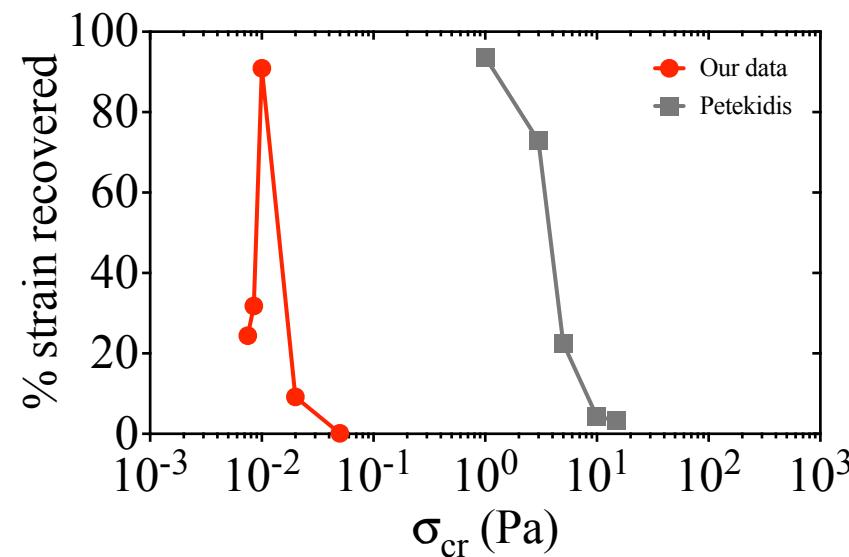
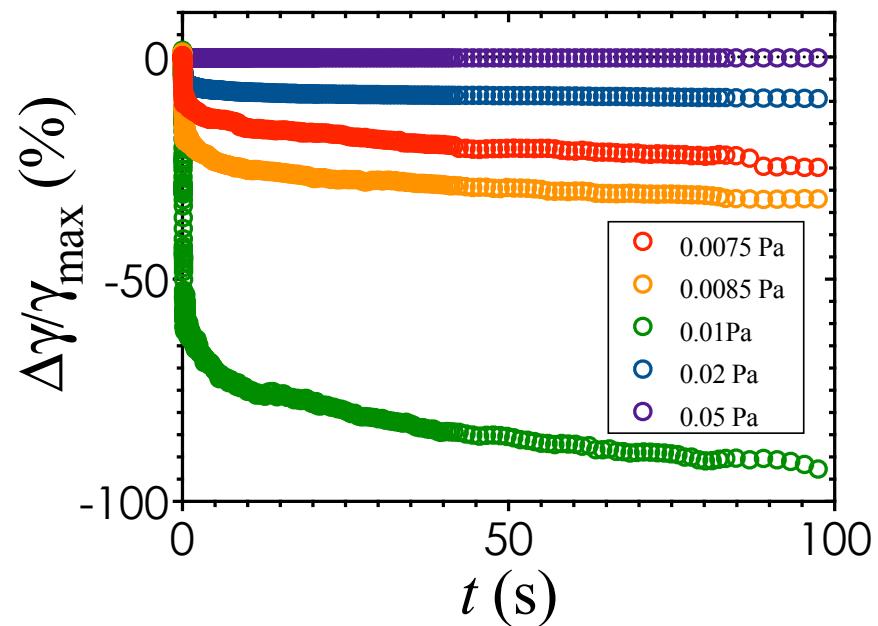
Strain recovery after creep cessation

PHSA-PMMA smooth colloids
 $(a = 183 \text{ nm} \pm 12\%, \phi = 0.62)$



Petekidis, Vlassopoulos & Pusey. JPCM (2004).

PHSA-PMMA rough colloids
 $(a = 1.7 \mu\text{m} \pm 5\%, \phi = 0.49)$



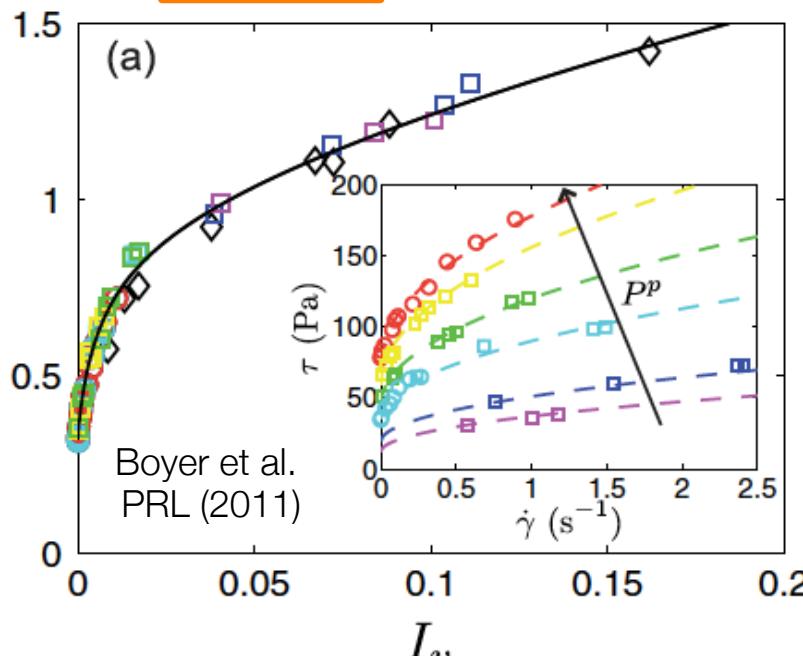
Testing nonlocal theories in dense suspensions

$$\mu(I, J, \phi) = \frac{\sigma_{xy}(J, \phi)}{P(J, \phi)}, I = \frac{2a\dot{\gamma}}{\sqrt{P/\rho_p}}, J = \frac{\eta_f \dot{\gamma}}{P}$$

Henann & Kamrin nonlocal theory

$$\boxed{\sigma_{ij}} = -\boxed{P\delta_{ij}} + \frac{2P}{g}\boxed{\dot{\gamma}_{ij}}$$

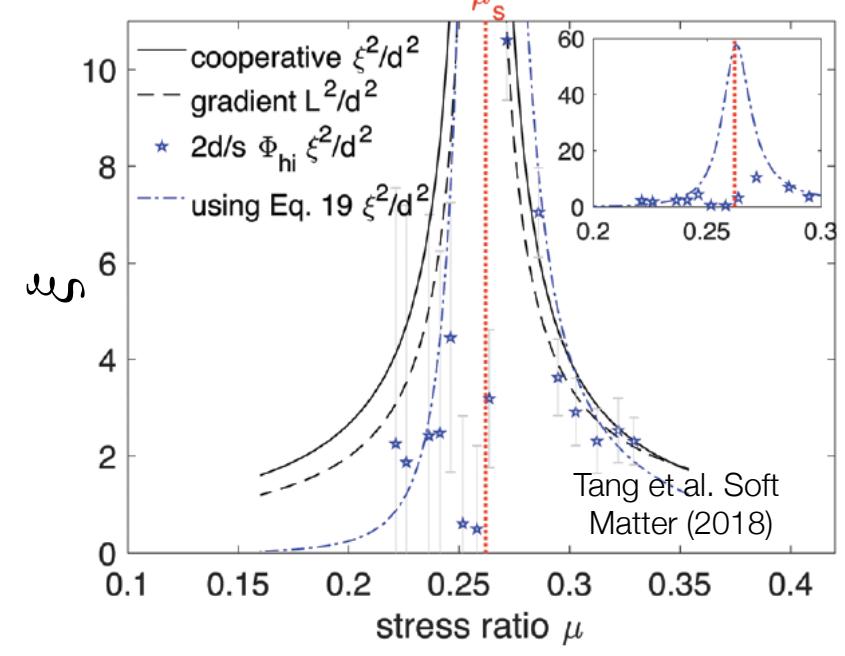
$$\nabla^2 g = \frac{1}{[\xi(\mu)]^2} (g - g_{loc}), g = \frac{\dot{\gamma}}{\mu}$$



Bouzid et al. nonlocal theory

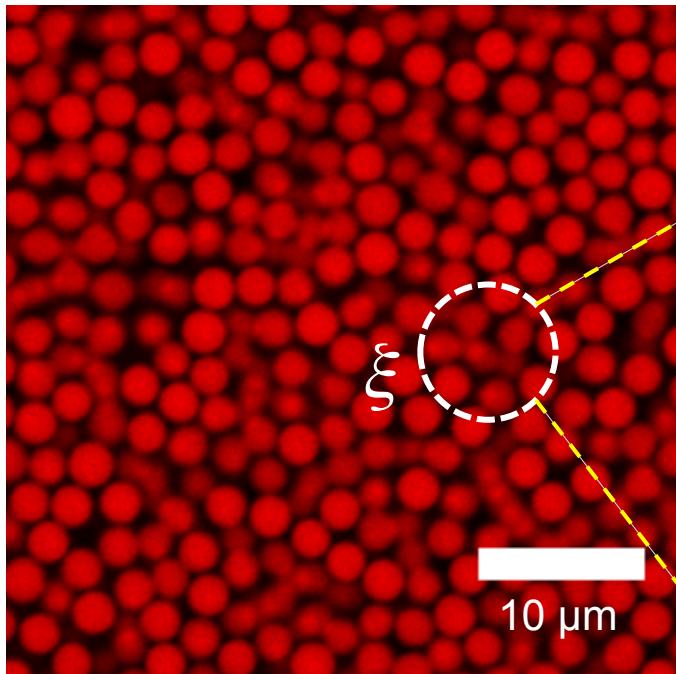
$$Y = \frac{\boxed{\sigma_{ij}}}{\boxed{\mu_s P}} = \frac{\mu}{\mu_s} \left(1 - A \frac{d^2 (\nabla^2 I)}{I} \right)$$

$$\boxed{\dot{\gamma}} = \frac{I_{loc}(f)}{t} - l^2 \nabla^2 f, f = \frac{\dot{\gamma}}{Y}$$

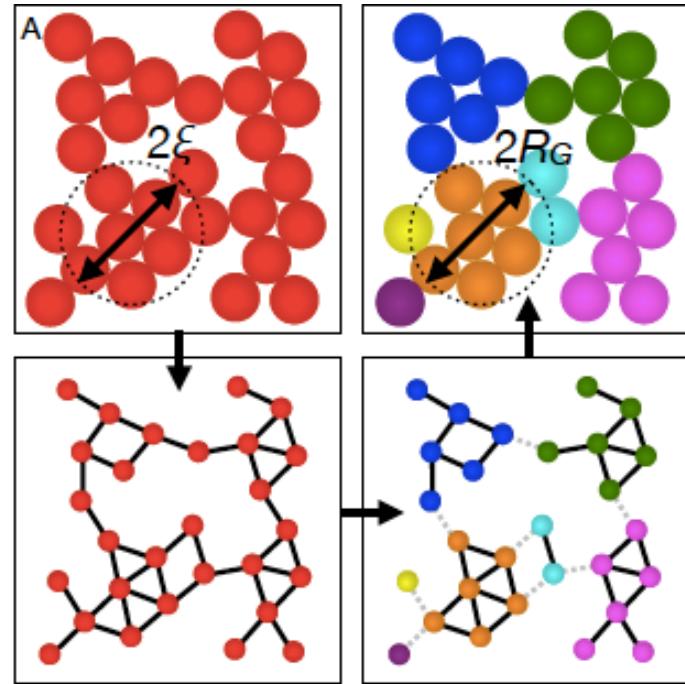
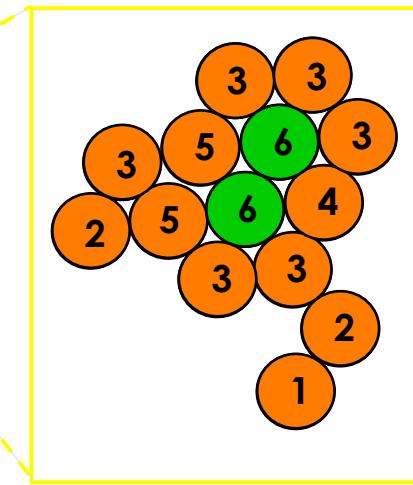


Measuring a correlation length scale

$$\xi(\phi, \mu, \dot{\gamma})$$



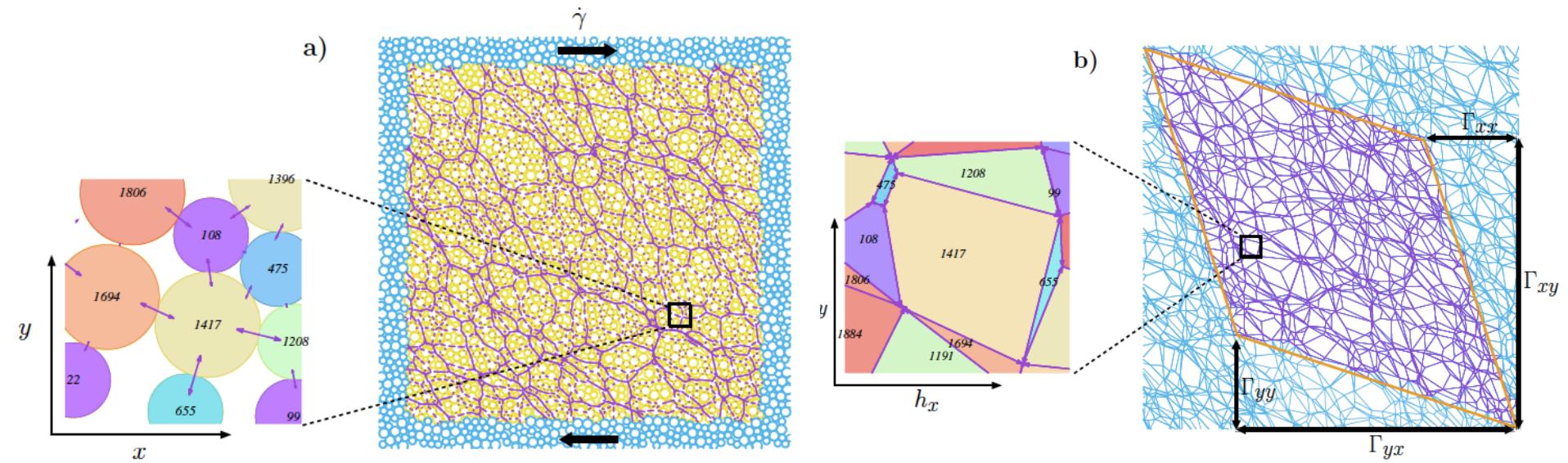
At $\phi = 0.60$,
 $\langle z \rangle \sim 8$ for smooth
 $\langle z \rangle \sim 6$ for rough



Whitaker & Hsiao et al. Submitted (2018).

- Are local rigid clusters responsible for rheology?
- Can we determine the length scale ξ as a function of ϕ ?
- Can we measure long-range velocity correlations?

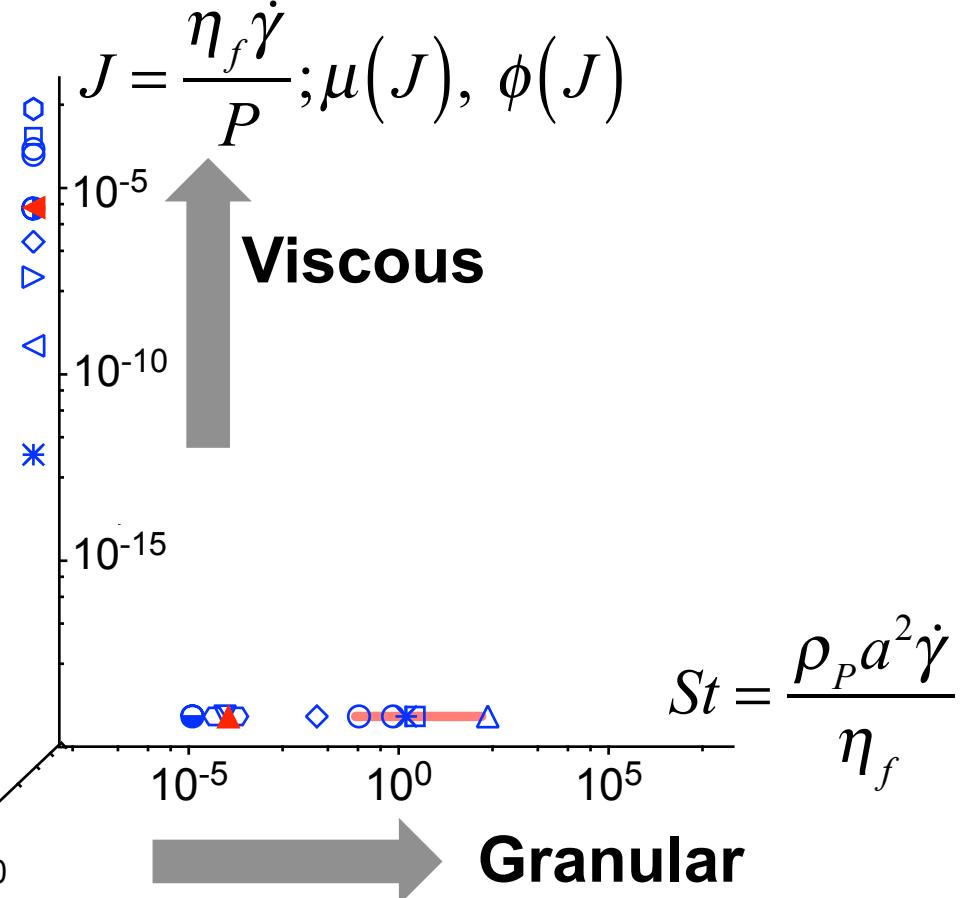
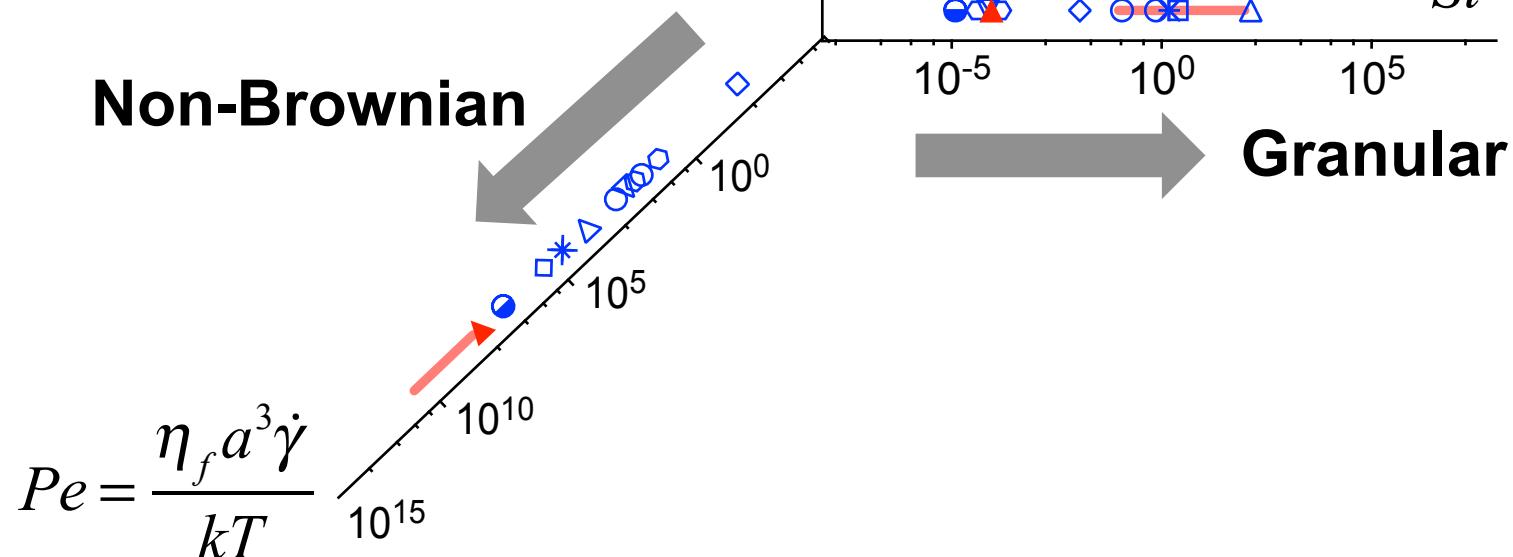
Testing force tiling concepts in suspensions



$$\mu = \frac{\sigma_{xy}}{P} \sim \frac{\sqrt{N_1^2 + \Sigma_{xy}^2}}{\Sigma_{xx} + \Sigma_{yy}}, \Sigma_{ij} = \sum_{i \neq j} L_{ij} F_{ij}$$

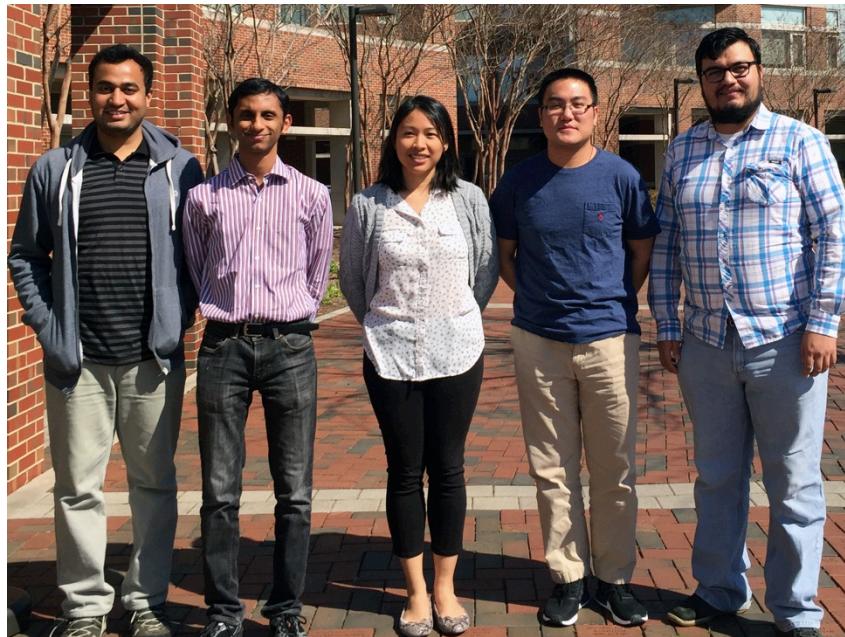
A state diagram for suspensions and granular flows

- Jerolmack et al., Nat Comm (2015)
- △ Fenistein et al., Phys Rev Lett (2004)
- ▽ Siavoshi et al., Phys Rev E (2006)
- ◇ Cagny et al., J Rheol (2015)
- Boyer et al., Phys Rev Lett (2004)
- Tanner & Dai, J Rheol (2016)
- × Behringer et al., Phys Rev Lett (2007)
- * Daniels et al., Soft Matter (2018)
- Jaeger et al., Nature (2016)
- ▲ Hsiao et al., Phys Rev Lett (2017)
- Wang & Brady, Phys Rev Lett (2015)



Summary

- Roughness introduces friction because it is difficult for particles to undergo full rotation
- Shifts maximum packing (based on viscosity divergence) to lower ϕ
- Causes glassy behavior at values of ϕ that is normally that of a fluid
- We may soon be able to test theories of $\mu(l, J)$ rheology in colloidal suspensions



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