Viscosity divergence and dynamical slowing down at the jamming transition

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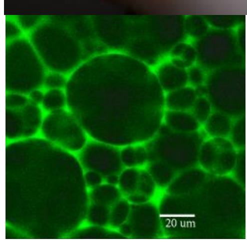
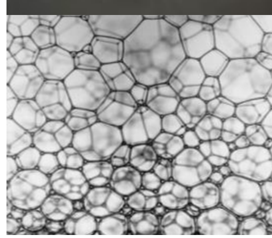
Jamming transition

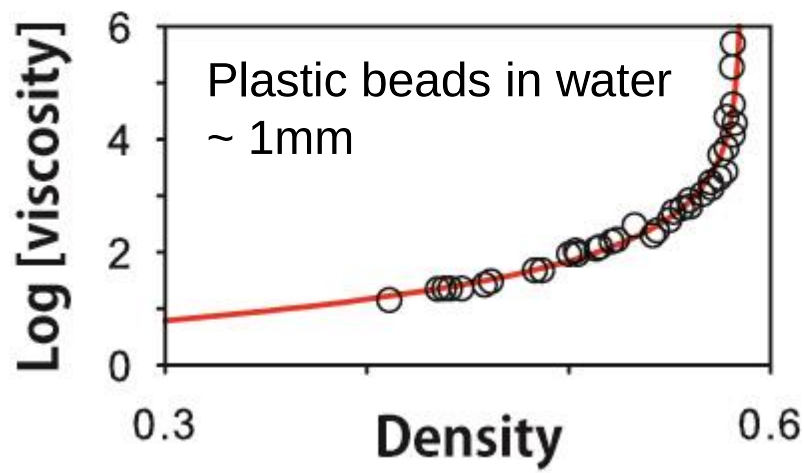
◆ **Athermal particles**: particles that are large enough to be free from Brownian motions

- ◆ Large colloidal particles, emulsions, foams, granular materials etc

◆ **Jamming transition = Viscosity divergence of athermal particles at the critical densitiy ~ random close packing density.**







$$\eta \propto (\varphi_{\mathrm{J}} - \varphi)^{-\gamma}, \quad \gamma \approx 2$$

[Boyer et al. 2011]

Glass vs Jamming

- **Glass transition = Viscosity divergence of thermal particles**
  - Small colloidal particles, atoms, molecules etc
  - Viscosity increases following **Vogel-Fulcher law**
  - **Relaxation time increases** as the viscosity increases (Green-Kubo formula)

- **Jamming transition = Viscosity divergence of athermal particles**
  - Large colloidal particles, foams, grains etc
  - Viscosity diverges following **the power-law**
  - **Relaxation time do NOT increase** near the transition

→ I will explain these points.

Jamming: Simple model

- **Athermal soft particles in solvent**
  - **Inter-particle interaction**

$$v(r) = \epsilon \delta^2 \qquad \delta = \begin{cases} 1 - \frac{r}{\sigma} & (r < \sigma) \\ 0 & (r > \sigma) \end{cases}$$

Overlap length between particles

* Energy penalty when overlapping

* Finite range, repulsive

contacts between particles are well-defined

  - **Overdamp equation of motion**

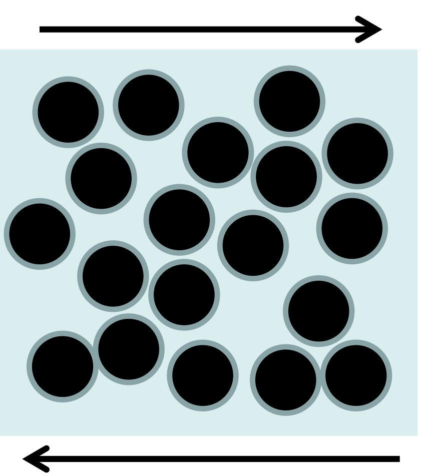
$$\xi \left(\frac{\partial \vec{r}_i}{\partial t} - \dot{\gamma} y_i \vec{e}_x \right) = - \sum_{j \neq i} \frac{\partial v(|\vec{r}_i - \vec{r}_j|)}{\partial \vec{r}_i}$$

Viscous dissipation

ξ = damping coefficient

$\dot{\gamma}$ = shear rate

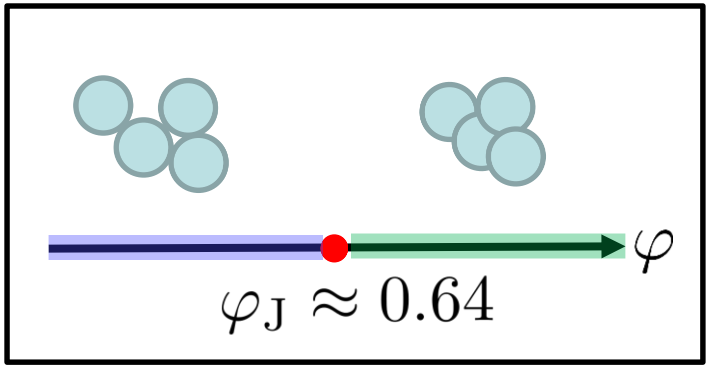
Inter-particle interaction



[Olsson-Teitel 2007, Ikeda et al. 2012 etc.]

Jamming: Phase diagram

- **Control parameter**
  - Packing density φ (or pressure)
  - Shear rate $\dot{\gamma}$ (or shear stress)



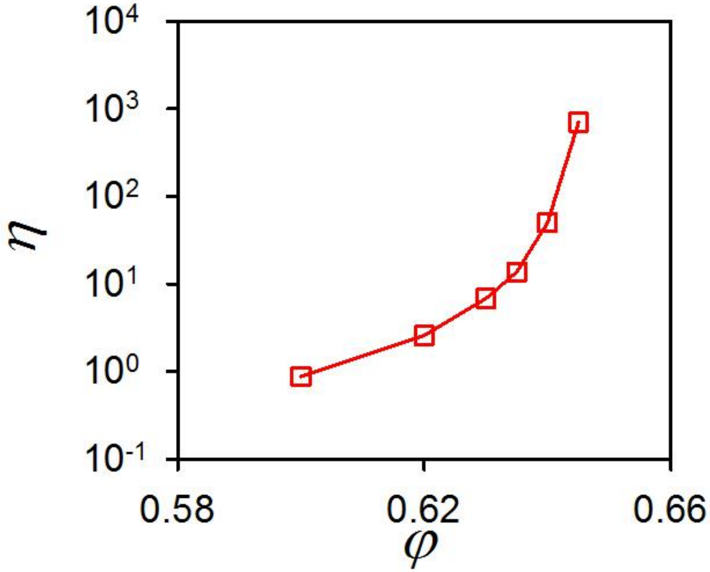
- **Phase diagram at low shear rate limit**
  - Low density: $\varphi < \varphi_{\mathrm{J}}$
    - Newtonian flow
    - Particles are just touching each other
  - High density: $\varphi > \varphi_{\mathrm{J}}$
    - Yielding of solid
    - Particles are overlapping
  - Critical density: $\varphi = \varphi_{\mathrm{J}}$
    - Marginally stable solid
    - Number of contacts per particle becomes isostatic

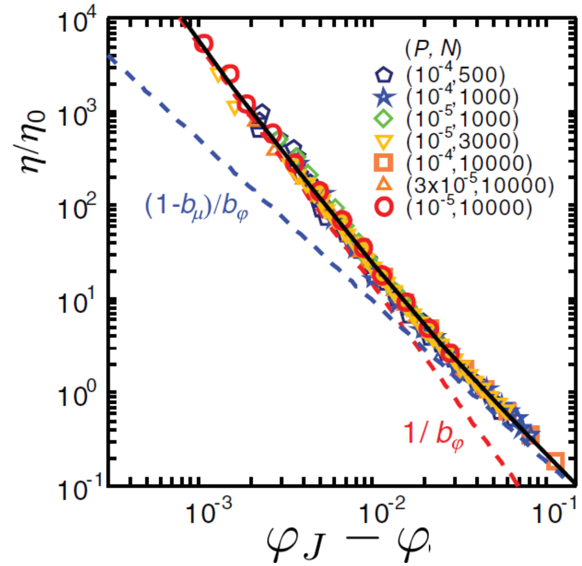
$$z = 2d \; (= 6)$$

[O'Hern 2003, Olsson-Teitel 2007]

Jamming: Viscosity

- **Newtonian viscosity** (at low shear rate)





- **Power-law divergence**

$$\boxed{\eta \propto (\varphi_{\mathrm{J}} - \varphi)^{-2.6} \quad \varphi_{\mathrm{J}} \approx 0.64}$$

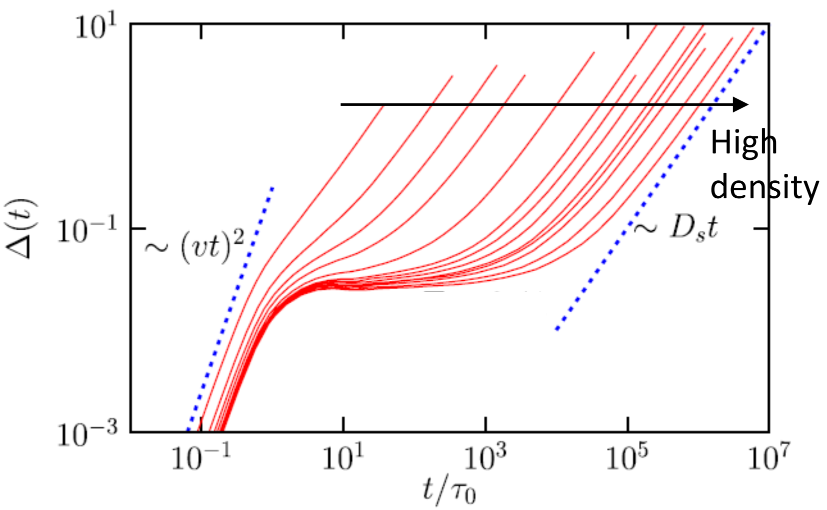
- Exponent changes near the transition: 1.7 → 2.6

[Lerner et al. 2012, Kawasaki et al. 2015]

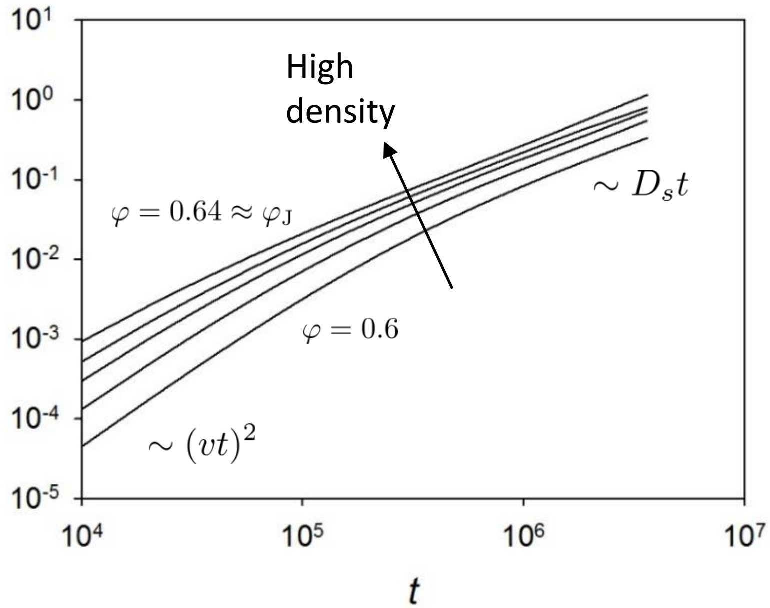
Jamming: Dynamics

- **Mean-square displacement** (at low shear rate, Newtonian regime)

Thermal particles (glass transition)



Athermal particles (jamming transition)



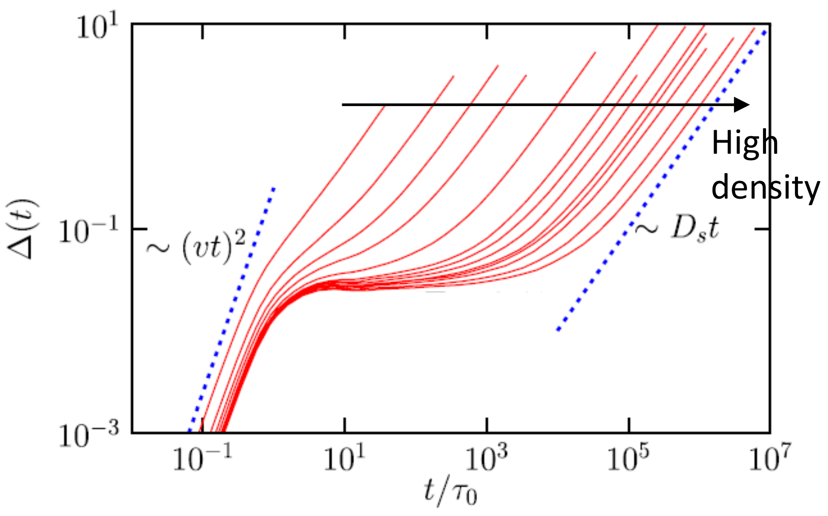
- **Jamming transition do not slow down the dynamics.**

[Heussinger et al. 2010, Ikeda et al. 2012]

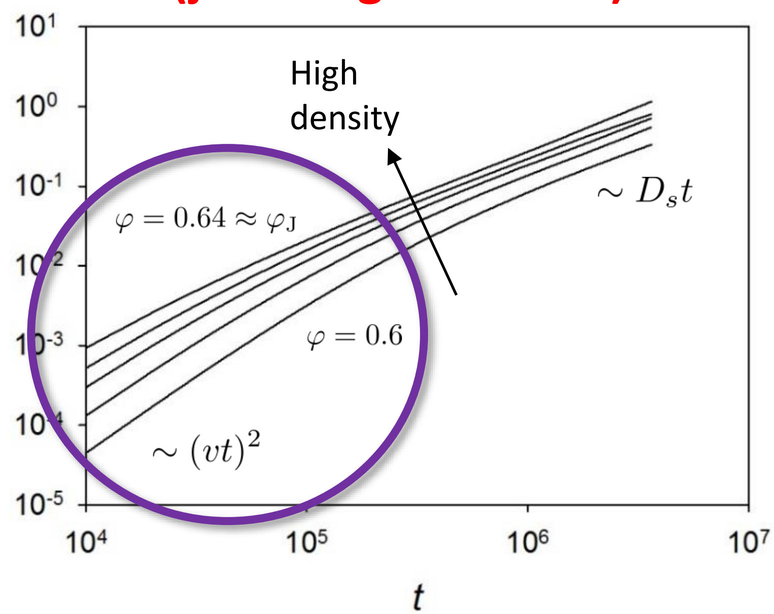
Jamming: Dynamics

- **Mean-square displacement** (at low shear rate, Newtonian regime)

Thermal particles (glass transition)



Athermal particles (jamming transition)

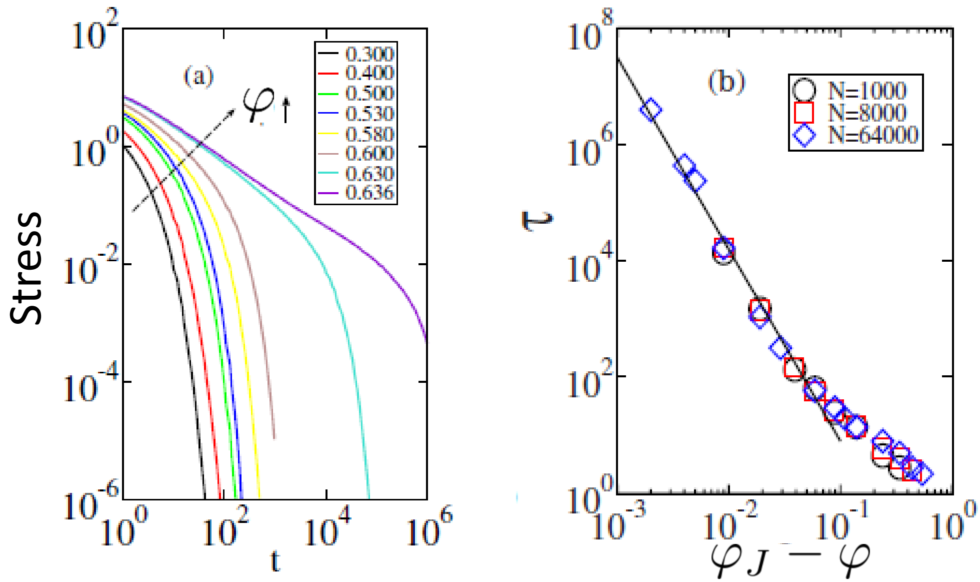


- **Jamming transition do not slow down the dynamics.**
- **Jamming even "speeds up" the short-time ballistic dynamics!**

[Heussinger et al. 2010, Ikeda et al. 2012]

Slowing down (1)

- **However in several settings, slowing down near the jamming have been observed**
  - Prepare particles configuration by the steepest descent without shear
  - Apply an infinitesimally small step strain.
  - Then, the system is relaxed by overdamped dynamics without shear and the relaxation of the stress is studied



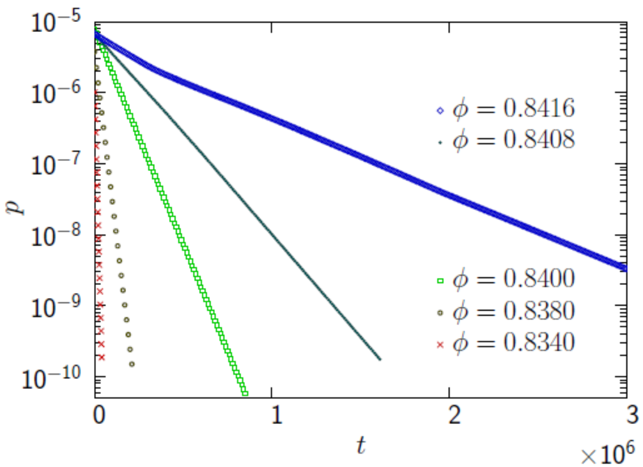
Relaxation time diverges at the jamming transition

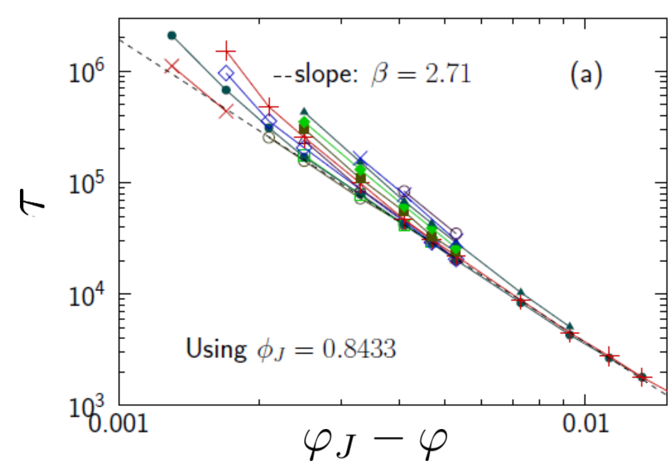
$$\tau \propto (\varphi_{\mathrm{J}} - \varphi)^{-3.3}$$

[Hatano 2010]

Slowing down (2)

- **However in several settings, slowing down near the jamming have been observed**
    - Perform overdamped dynamics with shear
    - Stop the shear. Then, the system is relaxed by overdamped dynamics without shear, and the relaxation of the pressure is studied





- **Viscosity and relaxation time diverges with the same power-law**

$$\eta \propto \tau \propto (\varphi_\mathrm{J} - \varphi)^{-2.7}$$

[Olsson 2014]

This work

?? Speed up near the jamming in steady-state, while slowing down in some cases ??

This work: Comprehensive study of the "dynamics" of the system near the jamming transition. And connect the "dynamics" to "the viscosity divergence".

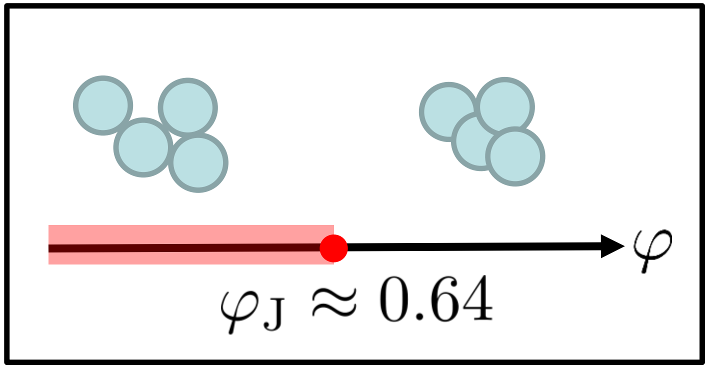
- **Relaxation dynamics near the jamming in the simplest setting**

Setting

◆ **Put particles randomly in a box.**

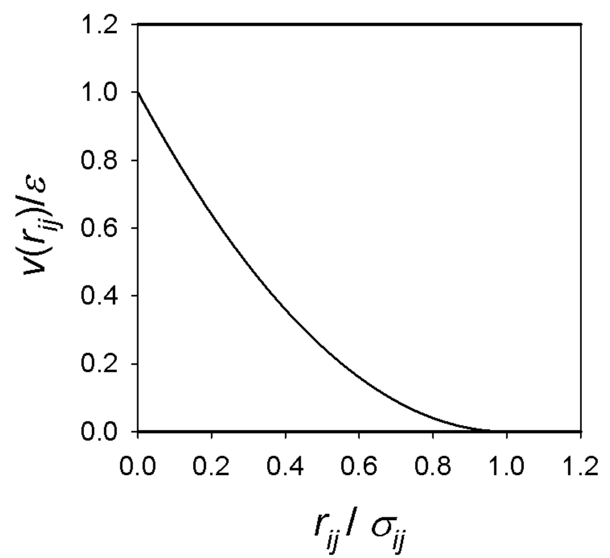
Focus only on the unjammed phase:

$$\varphi < \varphi_{\mathrm{J}}$$



◆ **Inter-particles interaction is**

$$v(r) = \begin{cases} \epsilon\left(1 - \frac{r}{\sigma}\right)^2 & (r \leq \sigma) \\ 0 & (r > \sigma) \end{cases}$$

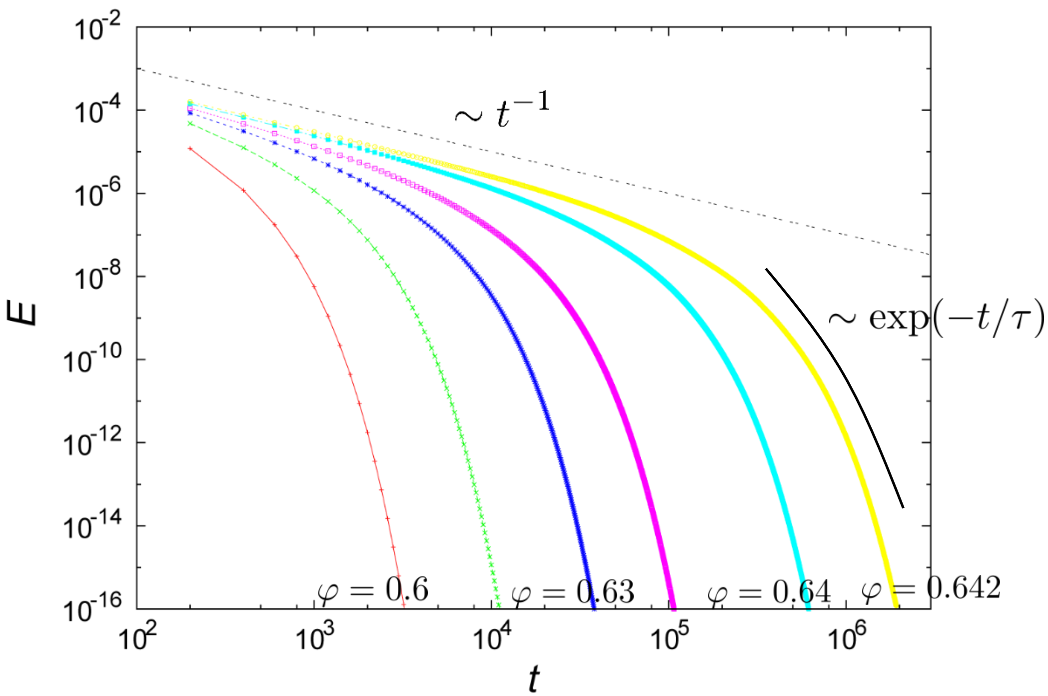


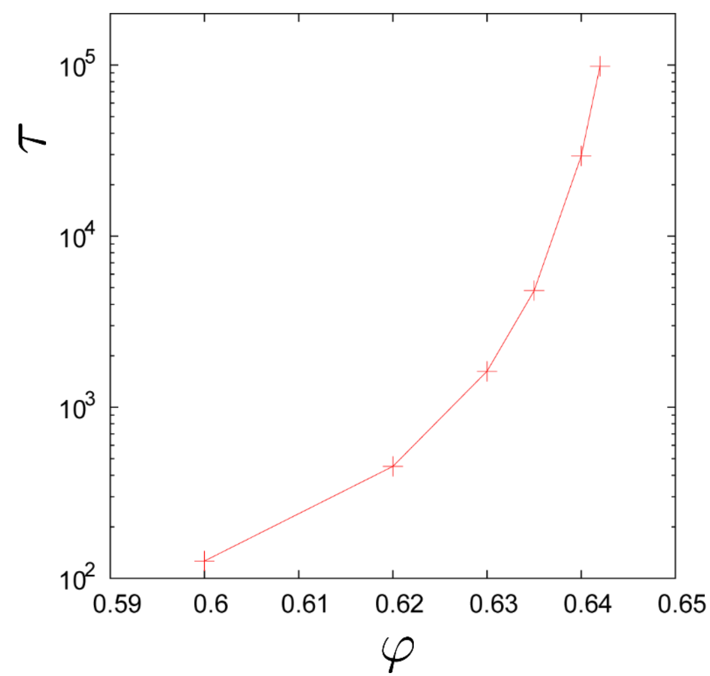
◆ **Then study the relaxation dynamics (No shear)**

$$\xi \frac{\partial \vec{r}_i}{\partial t} = -\sum_{j \neq i} \frac{\partial v(|\vec{r}_i - \vec{r}_j|)}{\partial \vec{r}_i}$$

Relaxation dymamics

- **Relaxation dynamics of the potential energy**

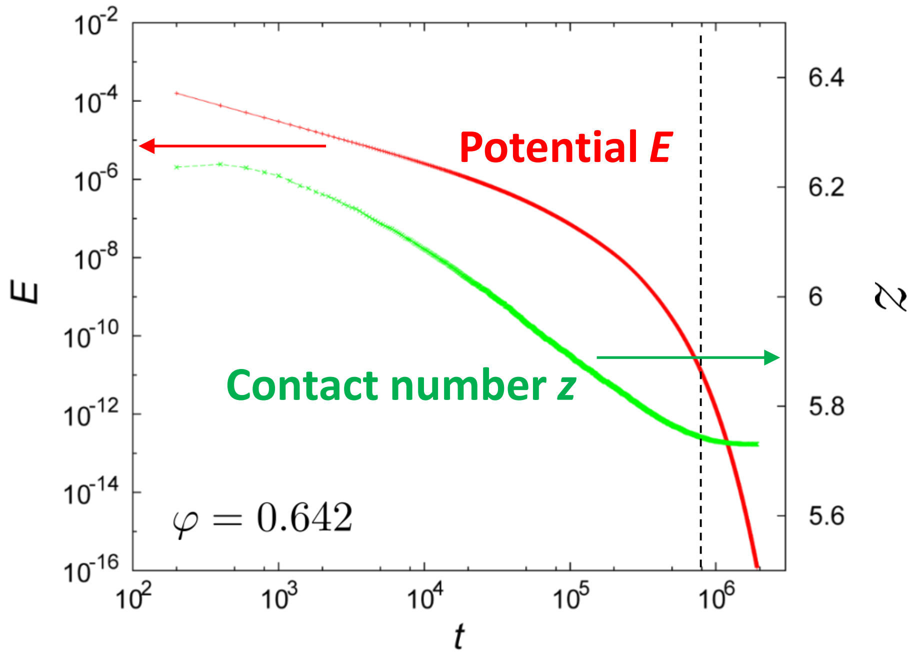




- Power-law*exponential: $E(t) \sim t^{-1} \exp(-t/\tau)$
- Relaxation time τ diverges at $\varphi \to \varphi_{\mathrm{J}}$

Contact number

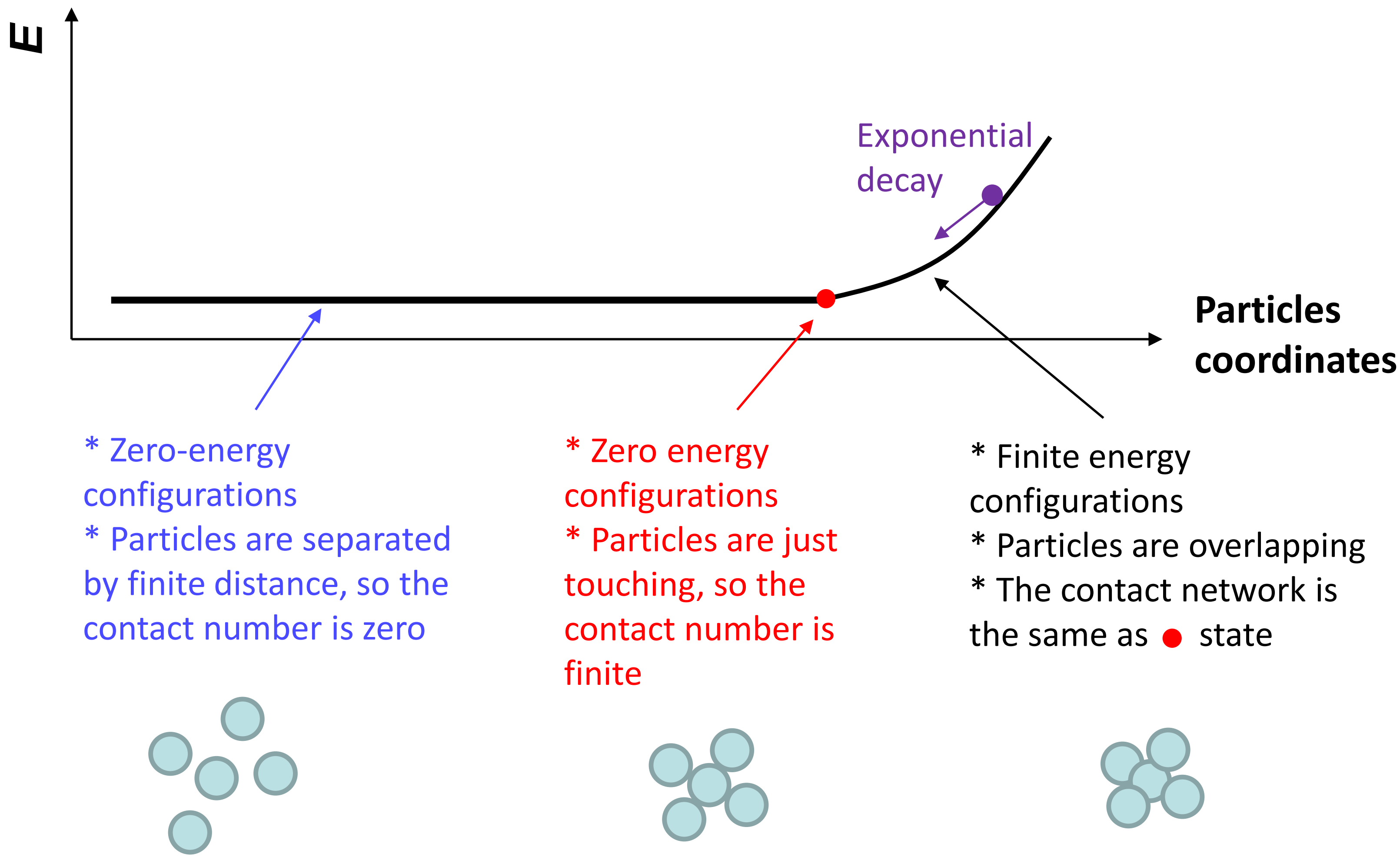
- Relaxation dynamics of contact number z
  - Contact number = average number of overlapping particles per particle



- Power-law region：*z* decreases (contacts are broken)
- Exponential region：*z* converges into a constant
  - **Relaxation without changing the contact network of particles**

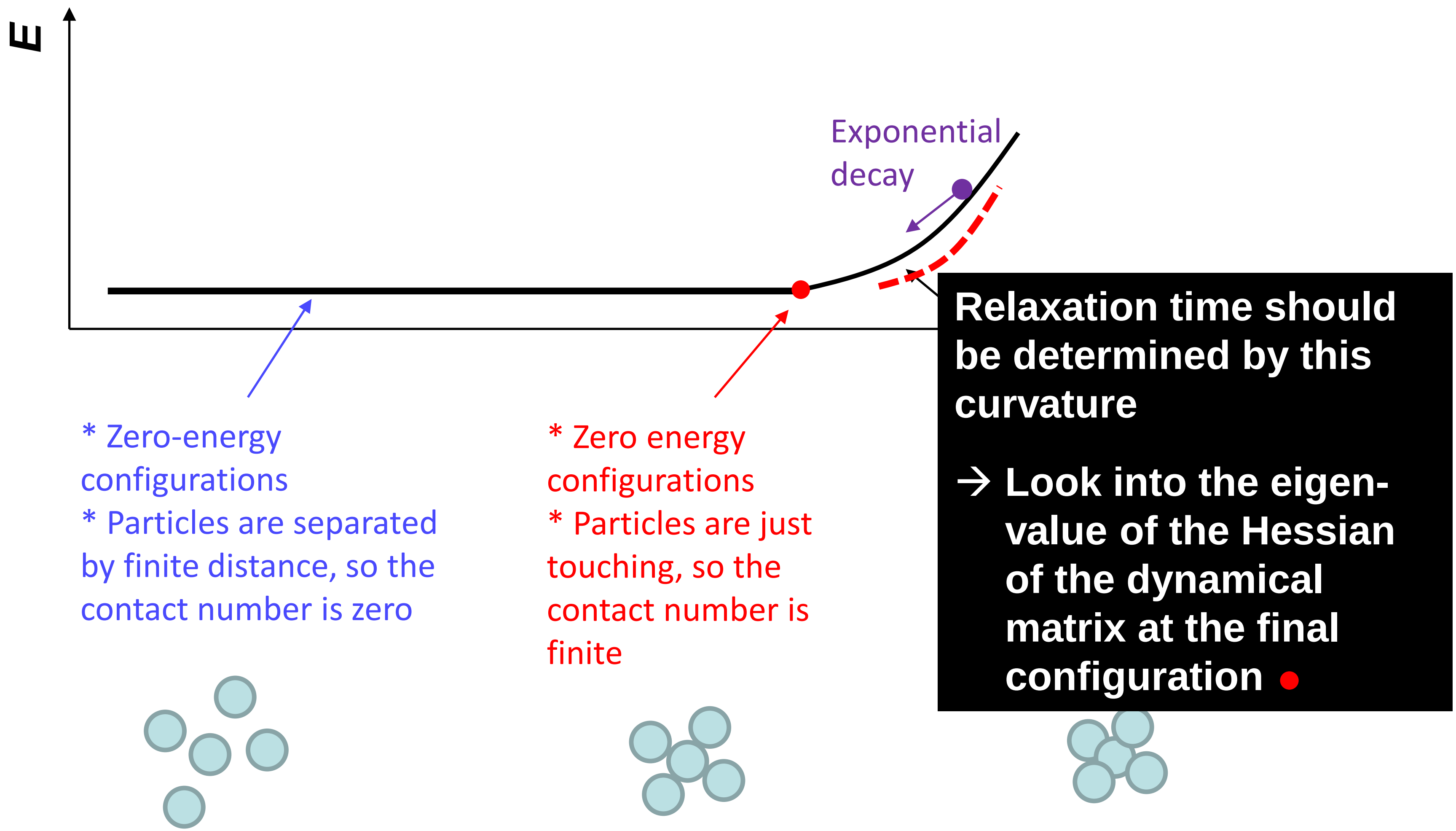
What's going on?

- Potential energy landscape at $\varphi < \varphi_{\mathrm{J}}$



What's going on?

- Potential energy landscape at $\varphi < \varphi_{\mathrm{J}}$



Eigenvalues

- For each φ , many final configurations are obtained from many initial configurations
- Obtain the Hessian of each final configuration

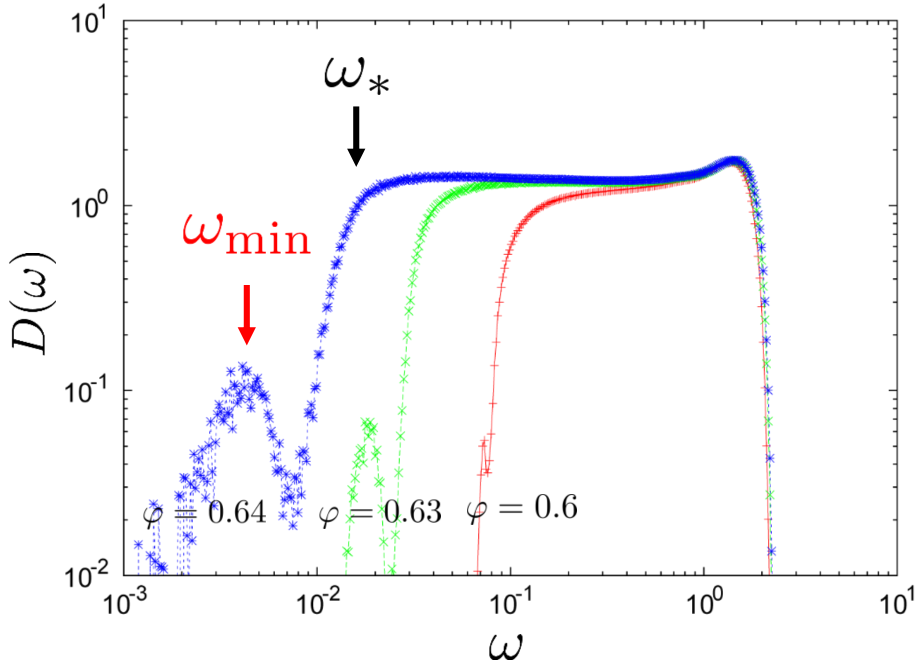
$$\mathbf{M} \equiv \frac{\partial^2 V}{\partial \vec{r}_i \partial \vec{r}_j} \qquad V = \sum_{i,j} v(|\vec{r}_i - \vec{r}_j|)$$

- Diagonalize the Hessian, obtain the eigenvalues $\{\lambda_\alpha\}$, calculate the vibrational density of states:

$$D(\omega) \equiv \frac{1}{N} \sum_\alpha \delta(\omega - \omega_\alpha) \qquad \omega_\alpha \equiv \sqrt{\lambda_\alpha}$$

  - We ignore the 3N*(6-z) zero modes, because we consider the relaxation from finite energy configurations
- Average $D(\omega)$ over obtained final configurations

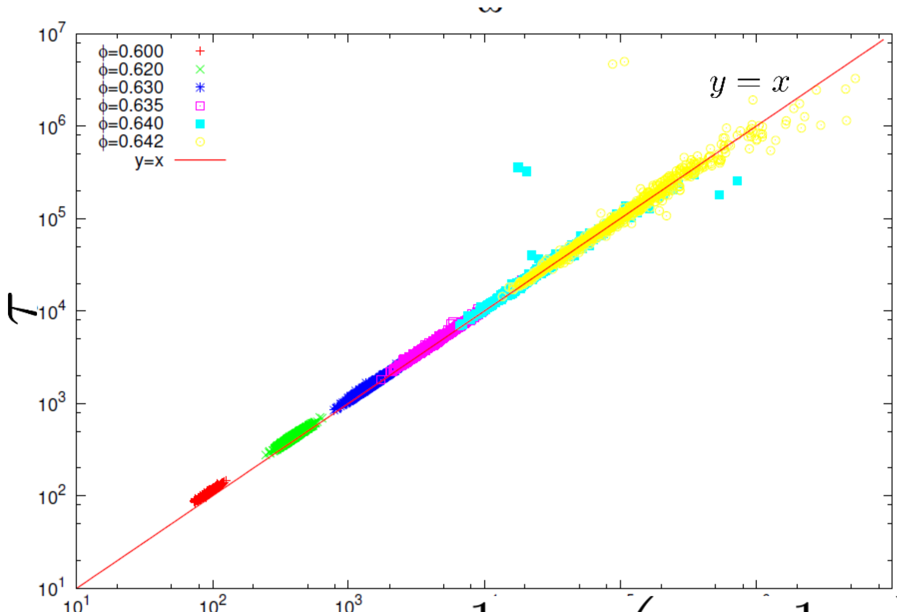
Eigenvalues



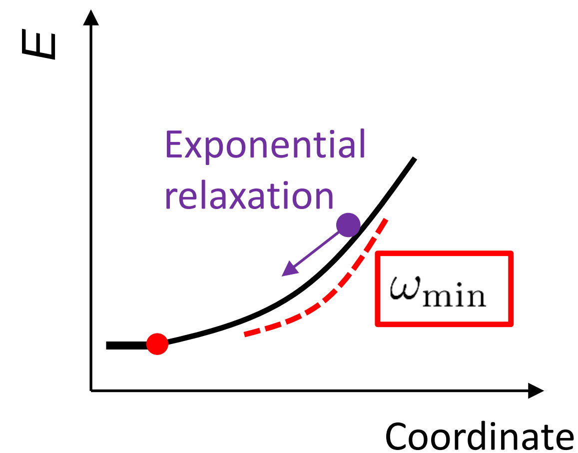
- $\omega > \omega_*$: Flat density of states
- $\omega = \omega_{\min}$: **There is one anomalously soft mode**
  - This mode is one isolated mode for one configuration

Relaxation time vs $\omega_{\min}$

- Plot the relaxation time against $\omega_{\min}$ for each configuration



$$\frac{1}{2\omega_{\min}^2} \quad \left(= \frac{1}{2\lambda \min} \right)$$

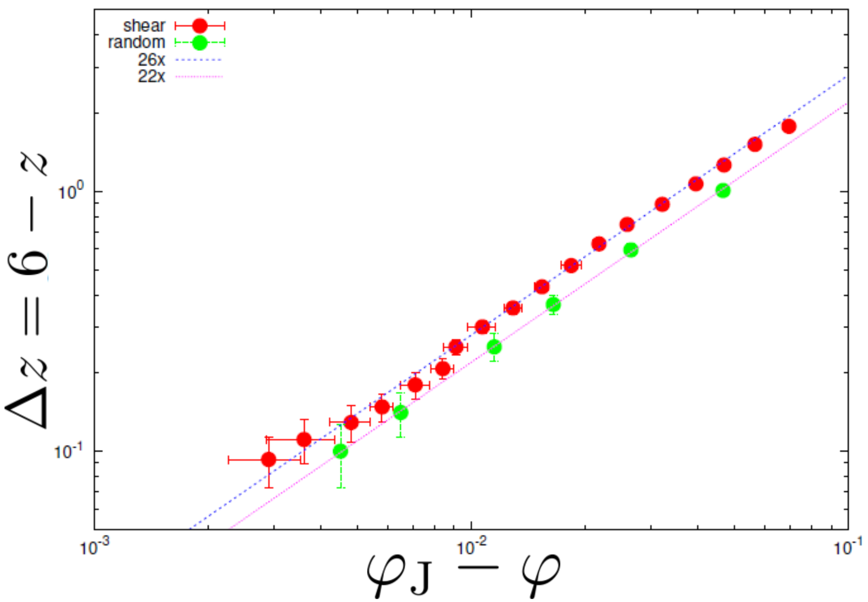


One configuration has only one extremely soft mode.

Exponential relaxation is the relaxation along this mode.

Critical behavior

- The terminal contact number vs density

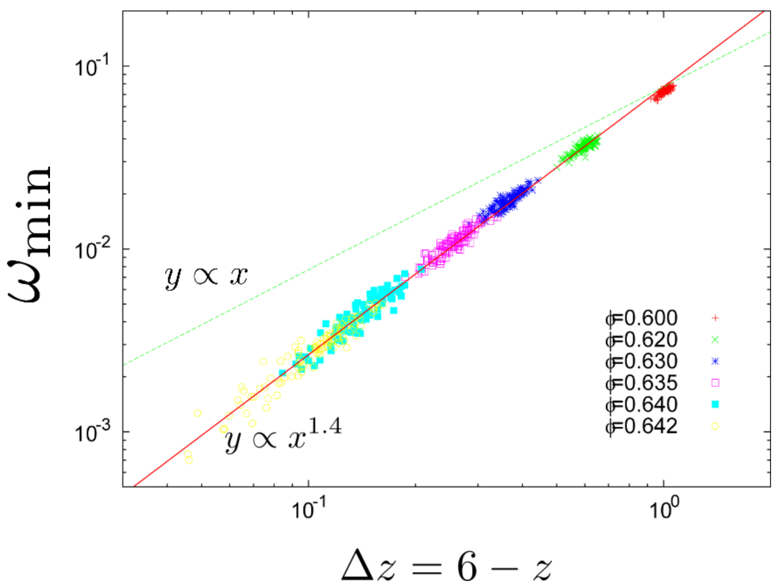


- **The terminal contact number linearly approaches the isostatic value, 6.**

$$\Delta z \propto (\varphi_J - \varphi)$$

Critical behavior

- The terminal contact number vs $\omega_{\min}$

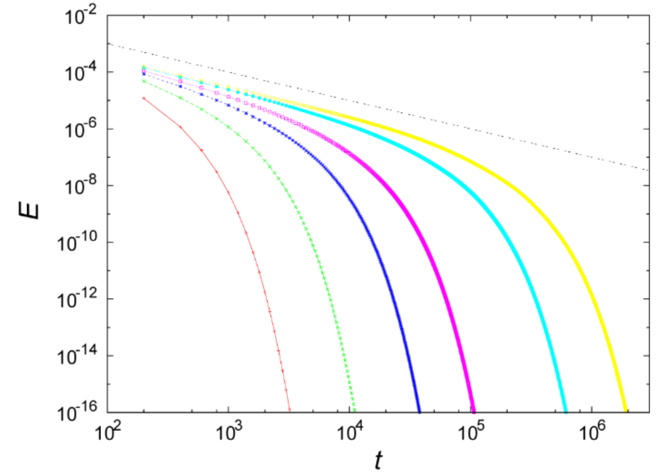


- Lowest frequency: $\omega_{\min} \propto \Delta z^{1.4} \propto (\varphi_{\rm J} - \varphi)^{1.4}$

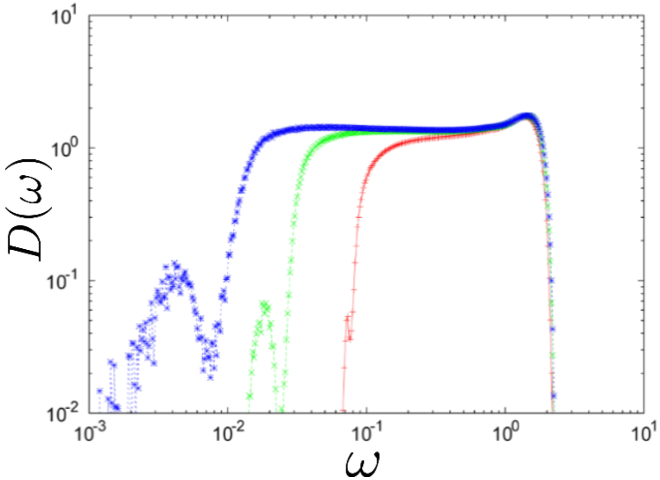
- Relaxation time： $\tau \sim (\varphi_{\rm J} - \varphi)^{-2.8}$

Summary

- The terminal relaxation is exponential, and takes place without changing the contact network.



- Each unjammed configuration has one anomalously soft mode. The terminal relaxation is the relaxation along this mode.



- The relaxation time (= inverse of the eigenvalue of the anomalously soft mode) diverges as

$$\tau \sim (\varphi_{\mathrm{J}} - \varphi)^{-2.8}$$