

Early Emission in Shock Breakout of Binary Neutron Star Merger

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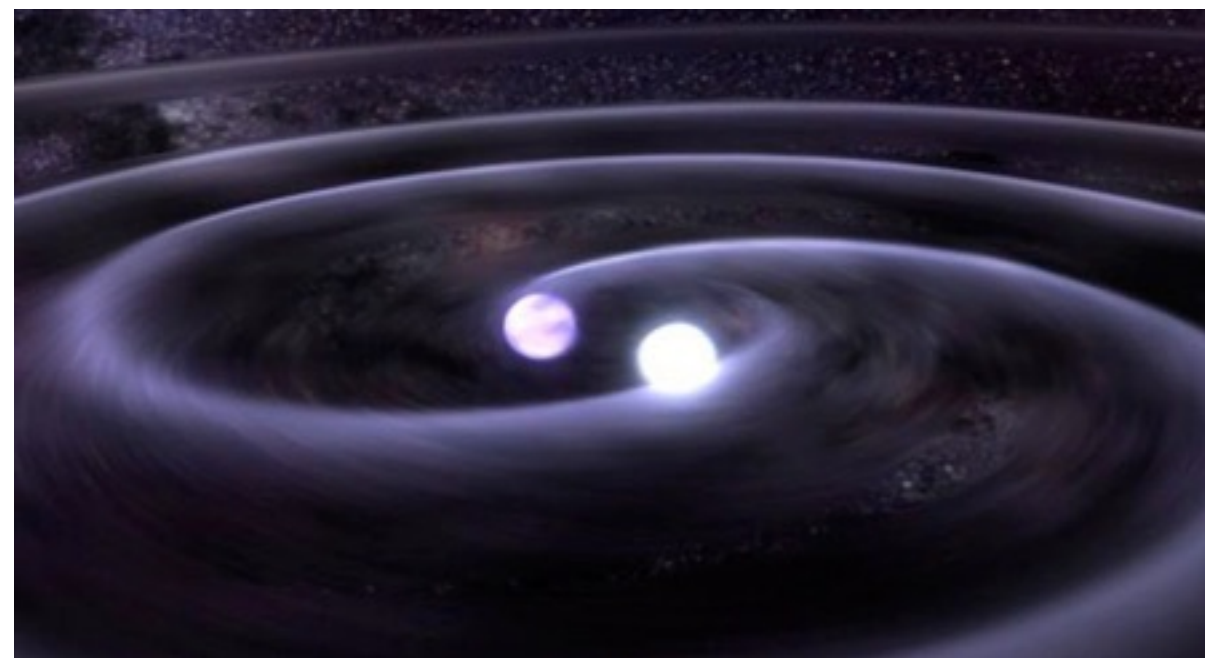


Emission from Binary Neutron Star Merger

(Li & Paczynski 1998, B. D. Metzger et al. 2010, ...)

GW170817

Electromagnetic emission
was detected over wide
wavelength range



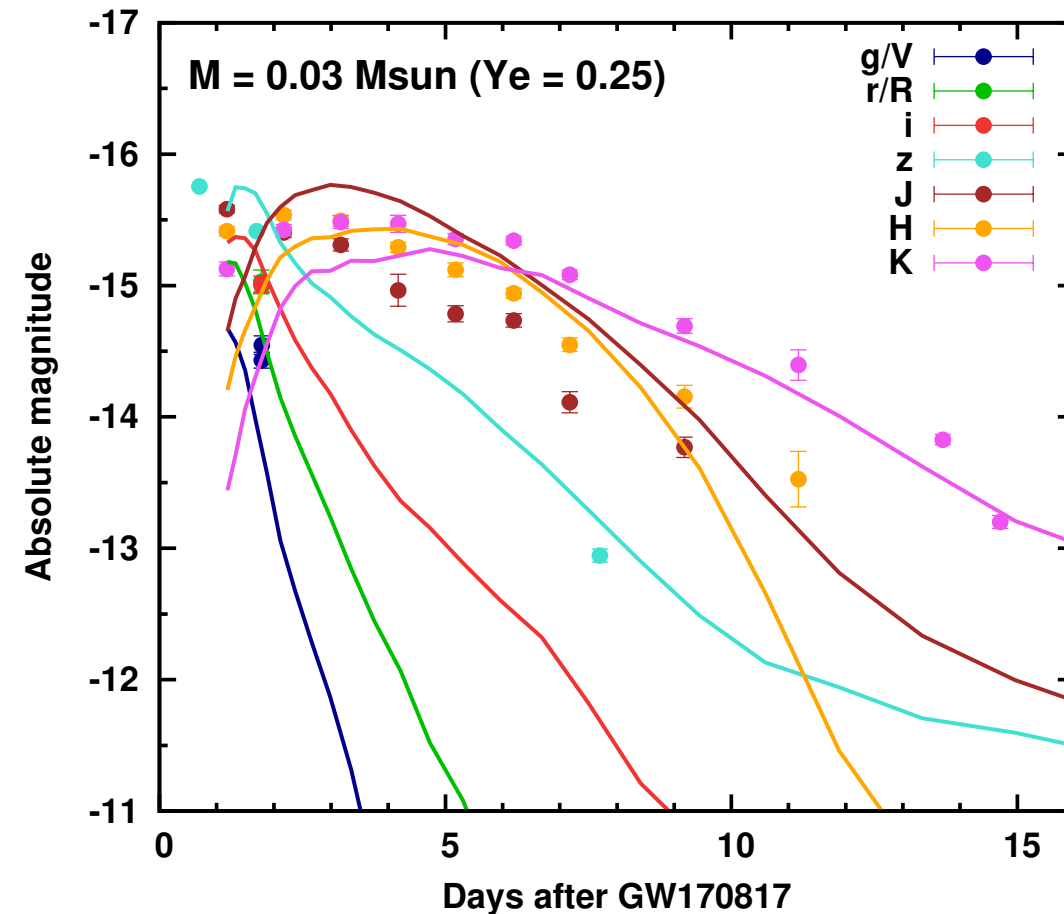
<http://aasnova.org/2015/10/28/what-do-you-get-when-two-neutron-stars-merge/>

Observed emission could be almost explained by kilonova model

- Neutron star matter is ejected at merging
(Neutron-rich ejecta for r-process nucleosynthesis)
- Emission by radioactive decay is detected

Early Emission from Neutron Star Merger

(M. Tanaka et al. 2017, Y. Utsumi et al. 2017, I. Arcavi 2018, ...)



- Observations started in ~ 11 h after merging event
- Observed early emission (~ 1 day) is more luminous and bluer than model computation

Early emission can provide us with rich information

→ It relates to the shock breakout

Emission in Shock Breakout

(C. D. Matzner and C. F. McKee 1999)

	NS merger	Supernova	long GRB
timescale (R/v)	$\sim 10^{-4}$ s	$10 \sim 10^4$ s ? (depends on progenitor)	a few 10^{-2} s ? ($\sim R/c\Gamma$)
nuclear reaction	occur	not occur	not occur
velocity	mildly relativistic	non relativistic	highly relativistic
opacity source	heavy elements H or He in outer region	depends on progenitor (H, He, O, C, Si, Fe, Ti, Ca...)	fully ionized gas e^{\pm}

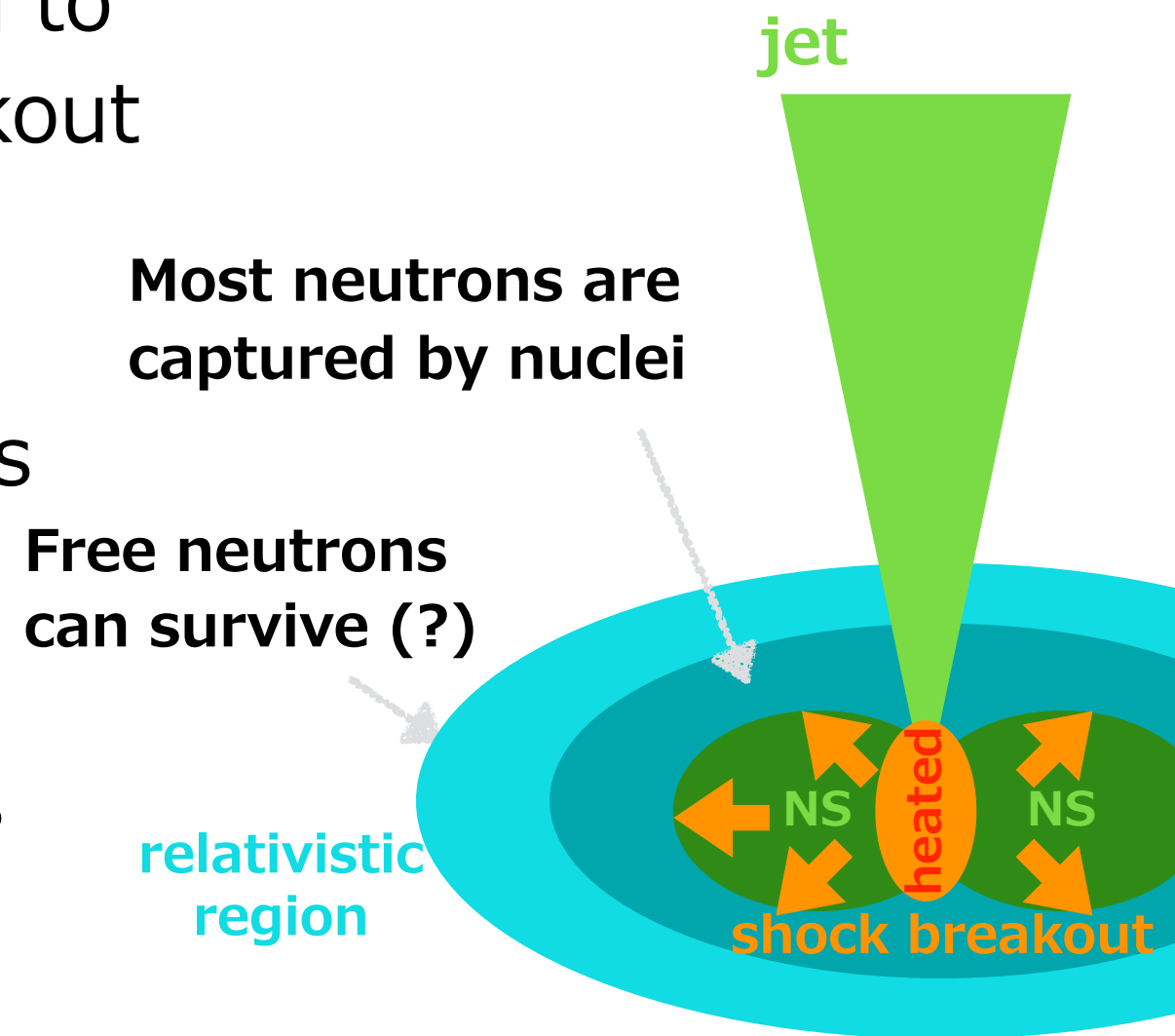
- The timescale may relate to the observability
- Radiation mediated shock should be considered
(\rightarrow it is neglected in subsequent topic for simplicity)

Free Neutron Precursor

(B. D. Metzger et al. 2015, Metzger 2017)

- Outermost ejecta is accelerated to relativistic speed in shock breakout (K. Kyutoku et al. 2014)
- Outermost ejecta expands sufficiently rapidly that neutrons avoid capture ($M \sim 10^{-4} M_{\text{sun}}$) (Goriely et al. 2014, Just et al. 2014)
- β -decay of free neutron powers "precursor" to kilonova

peaks at ~ few hours



Smoothed Particle Hydrodynamics (SPH) simulation (Just et al. 2015)

→ Can the similar result be obtained in grid-based simulations?

Objectives

Examining early emission by free-neutron-powered precursor in shock breakout of binary neutron star merger

Step 1

- Developing relativistic Lagrangian hydrodynamics code and reproducing shock breakout of neutron star merger

Step 2

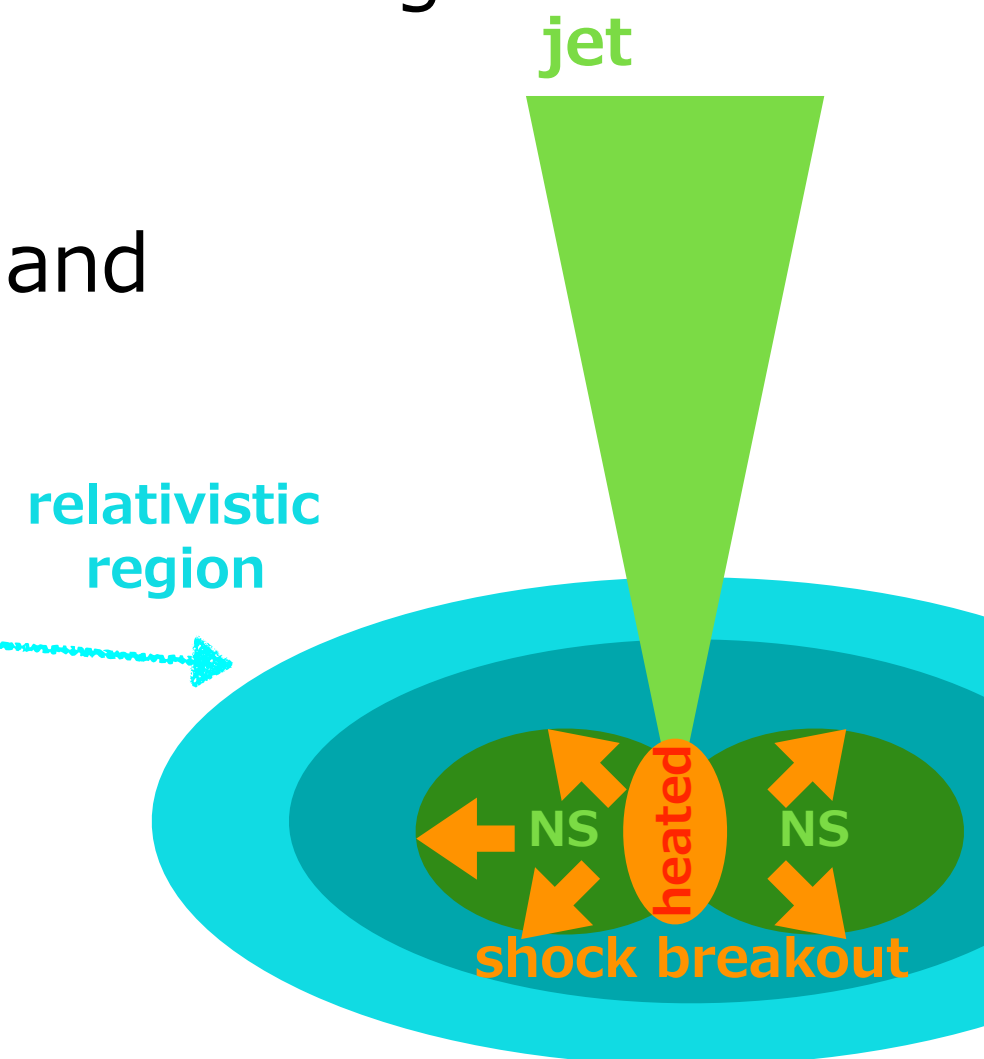
- Estimating surviving free neutron mass fraction with e^+ and e^- captures and some nuclear reactions

Step 3

- Calculating mass of region where reasonable amount of free neutron is surviving

Step 4

- Calculating emission from β -decay of free neutron



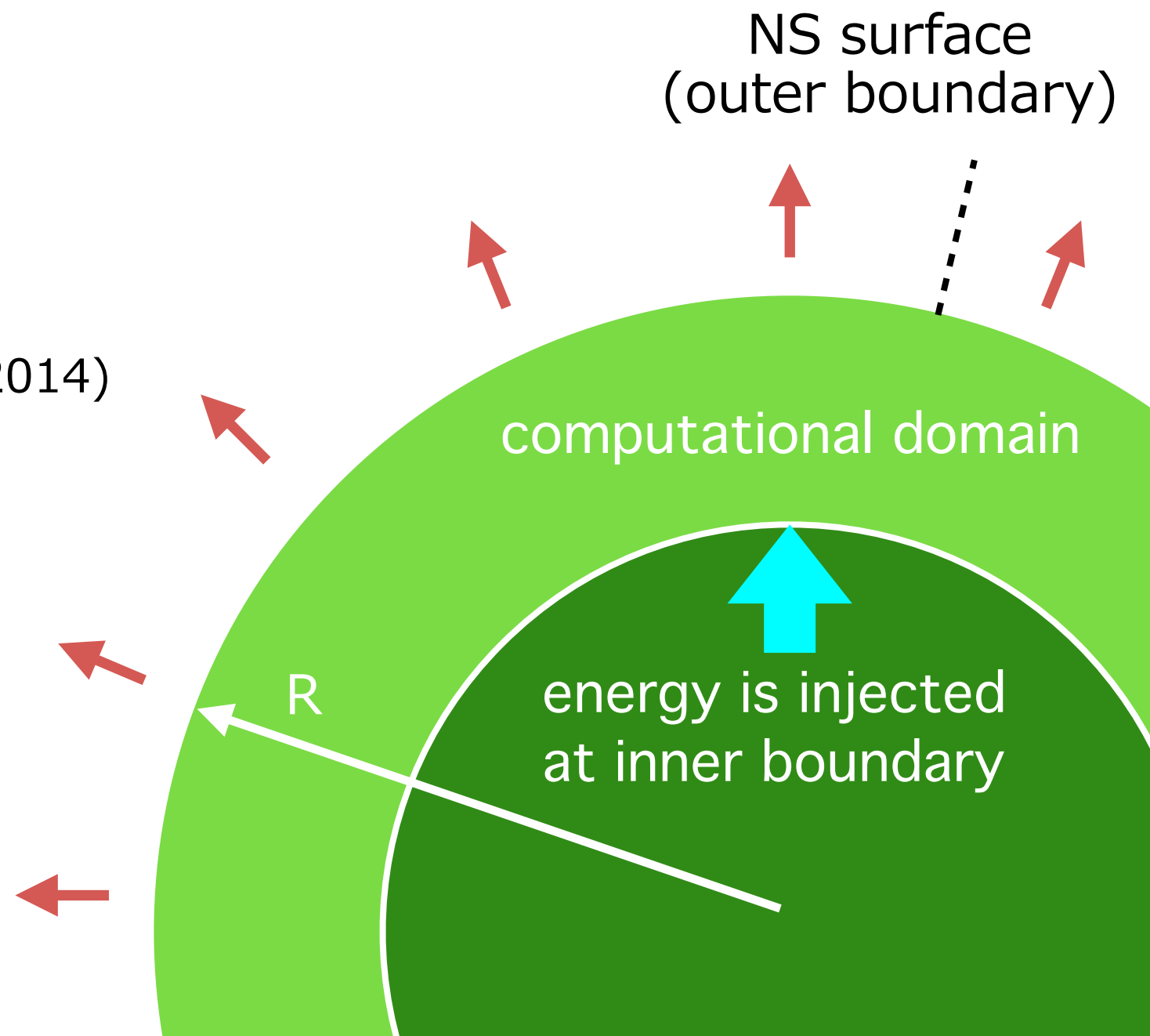
Simulation Condition

- Relativistic Lagrangian hydro simulation
- 1D spherical symmetric coordinate
- 500 computational cells in radial direction
- $E_{\text{final}} = 10^{47} - 10^{50}$ erg
- $R = 15, 20, 25, 30$ km
- $M_{\text{shell}} = 10^{-3} M_{\text{sun}}$
- $\rho \propto (R - r)^3$ (K. Kyutoku et al. 2014)

Shock wave propagates
through merging NS



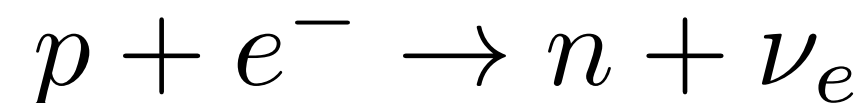
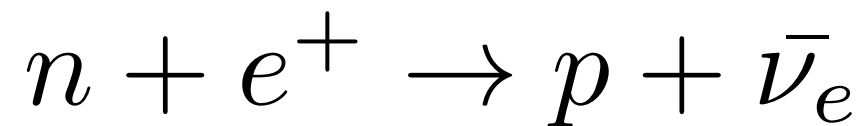
Shock breakout occurs
when it reaches NS surface



Estimation of free neutron rate

- Free neutron rate X_n is set to be 0.9 initially ($Y_e=0.1$) (beta equilibrium of cold dense matter)

- e^\pm is generated by shock heating



- Time scale of positron and electron capture processes are obtained depending on temperature ($\tau_+(T)$, $\tau_-(T)$) (L. Kawano 1992, B. D. Metzger et al. 2015)

- Time evolution of X_n is calculated by

$$\frac{dX_n}{dt} = -\frac{X_n}{\Gamma\tau_+(T)} + \frac{(1-X_n)}{\Gamma\tau_-(T)}$$

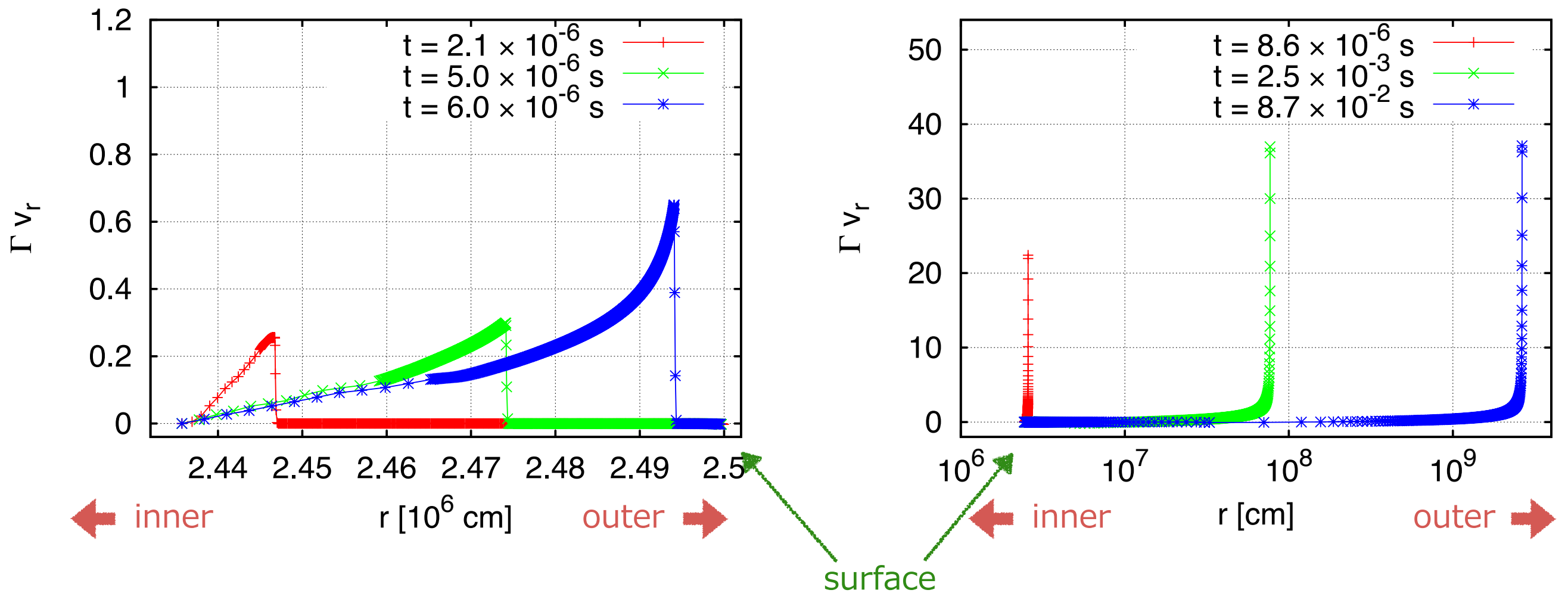
- Nuclear reaction network calculations are performed after temperature decreases down to 10^{10} K (Shigeyama et al. 2010)

Results in Shock Breakout

Before shock breakout

$R = 25 \text{ km}, E_{\text{final}} = 10^{49} \text{ erg}$

After shock breakout



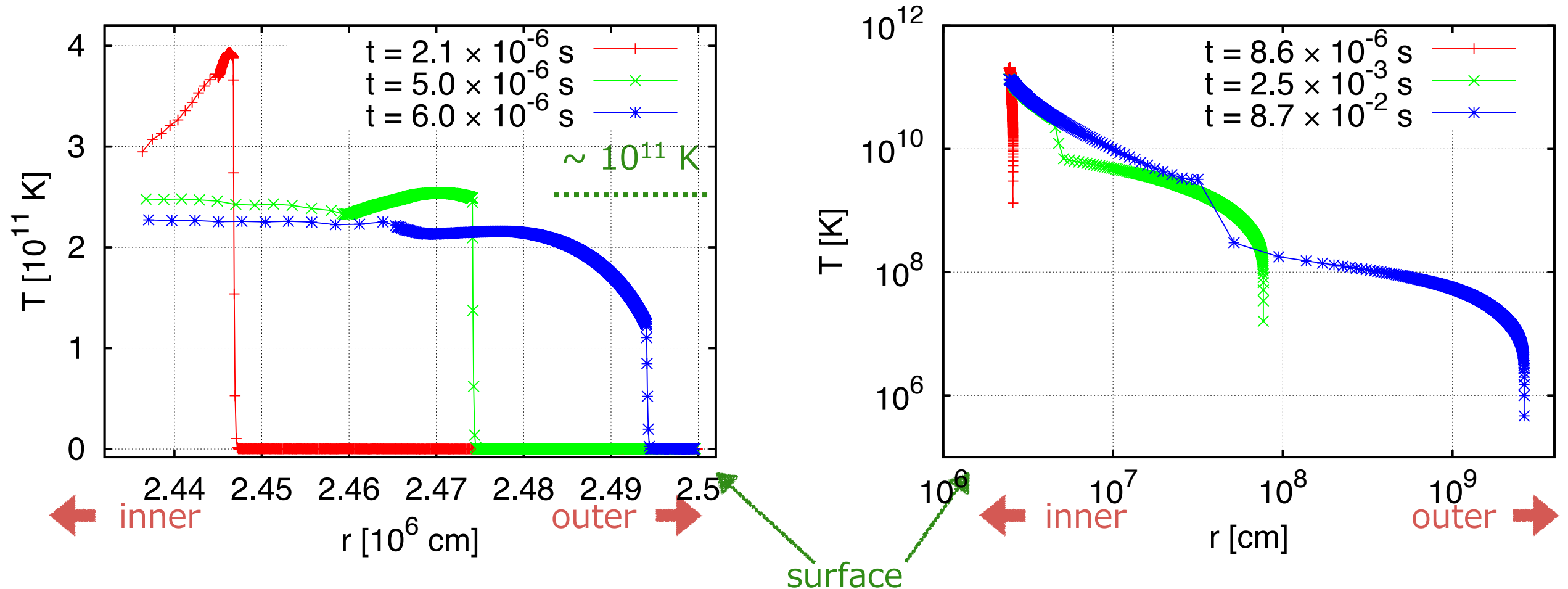
- Accelerated shock wave in breakout can be reproduced
- Ejecta in outermost region has relativistic speed

Results in Shock Breakout

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After shock breakout

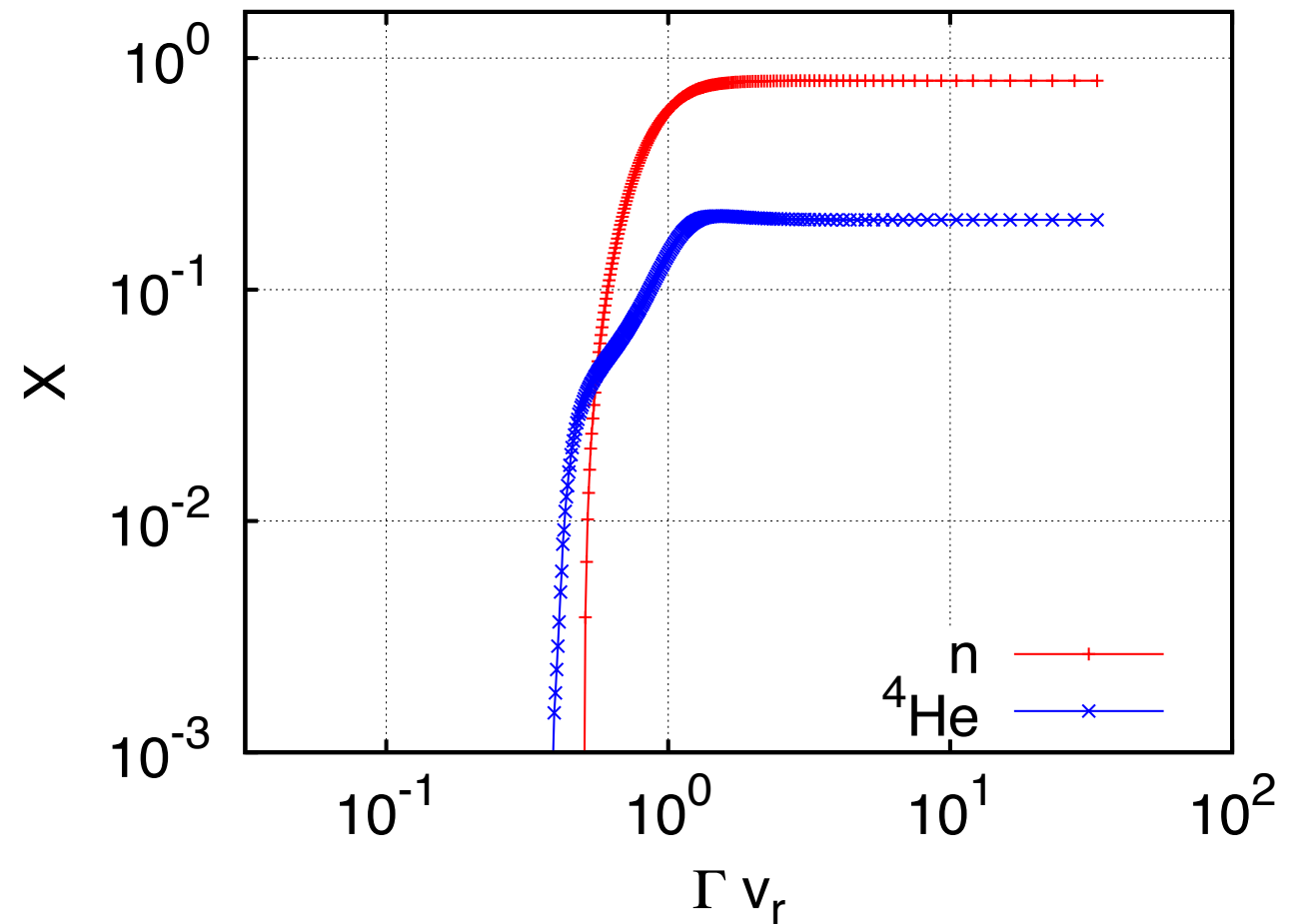


- Radiation pressure is assumed to be dominant ($P = \frac{aT^4}{3}$)
- Temperature decreases rapidly after shock breakout

$$\tau_+ \simeq 2.1 (T/\text{MeV})^{-5} \text{ s} \xrightarrow{10^{11} \text{ K}} \sim 4.41 \times 10^{-5} \text{ s}$$

Distribution of Mass Fraction

$R = 25 \text{ km}$
 $E_{\text{final}} = 10^{49} \text{ erg}$
 $t = 9 \times 10^{-2} \text{ s}$
 $T < 10^8 \text{ K}$



Inner region

- Neutrons are captured by nuclei to produce heavy elements

Middle region

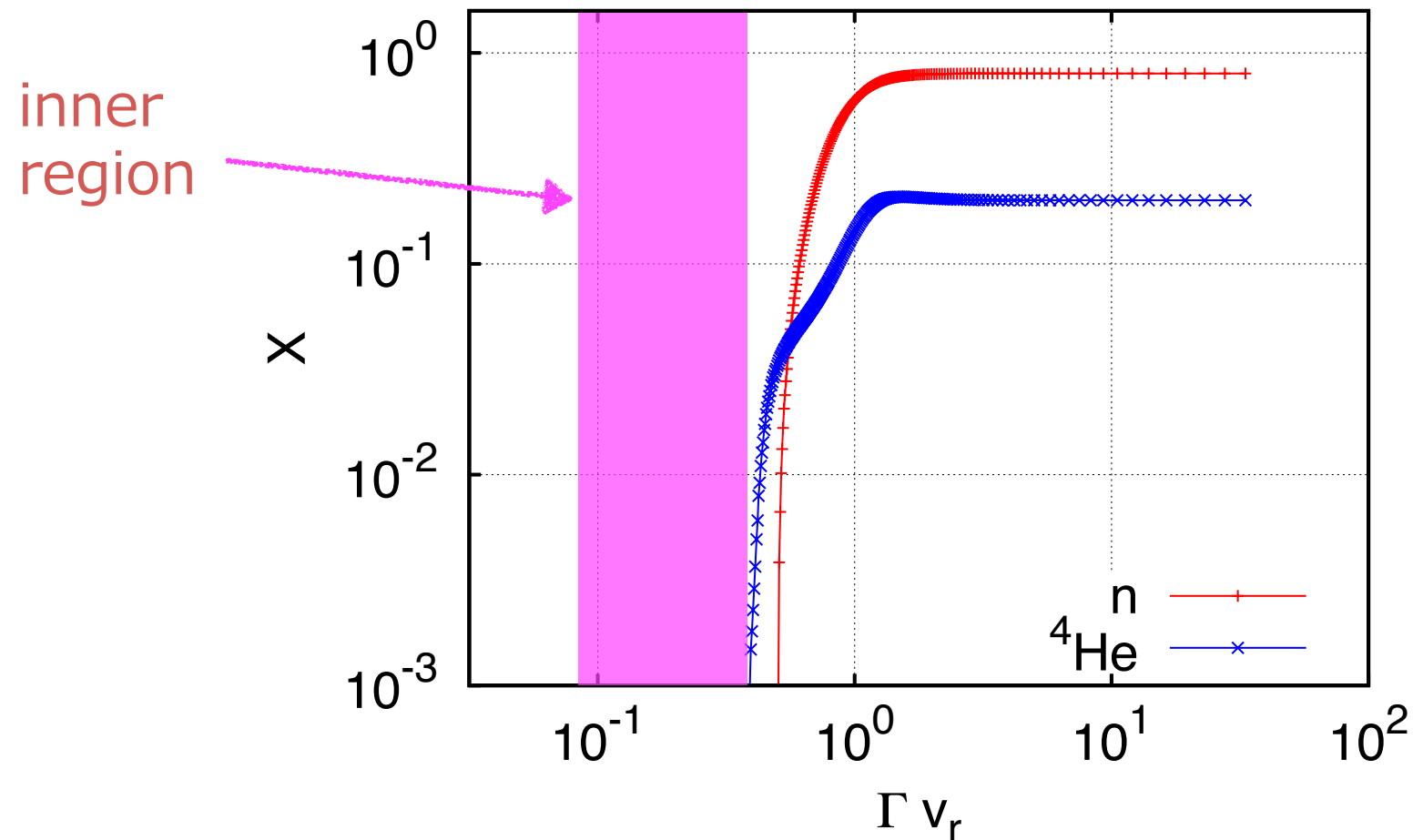
- $p(n,\gamma)d$ reactions consume all neutrons to produce ${}^4\text{He}$

Outermost region

- Free neutron layer is produced due to low density and temperature

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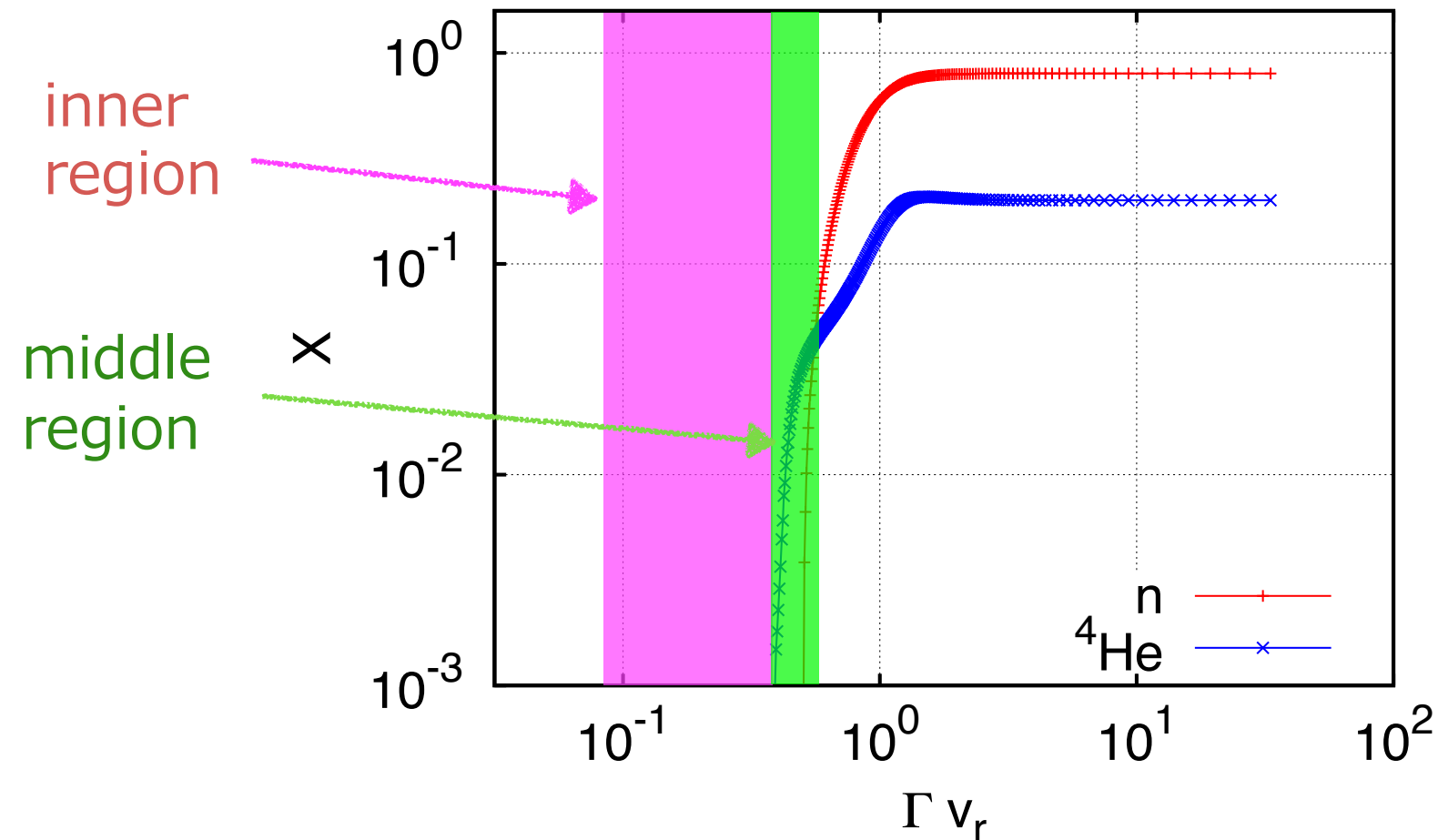
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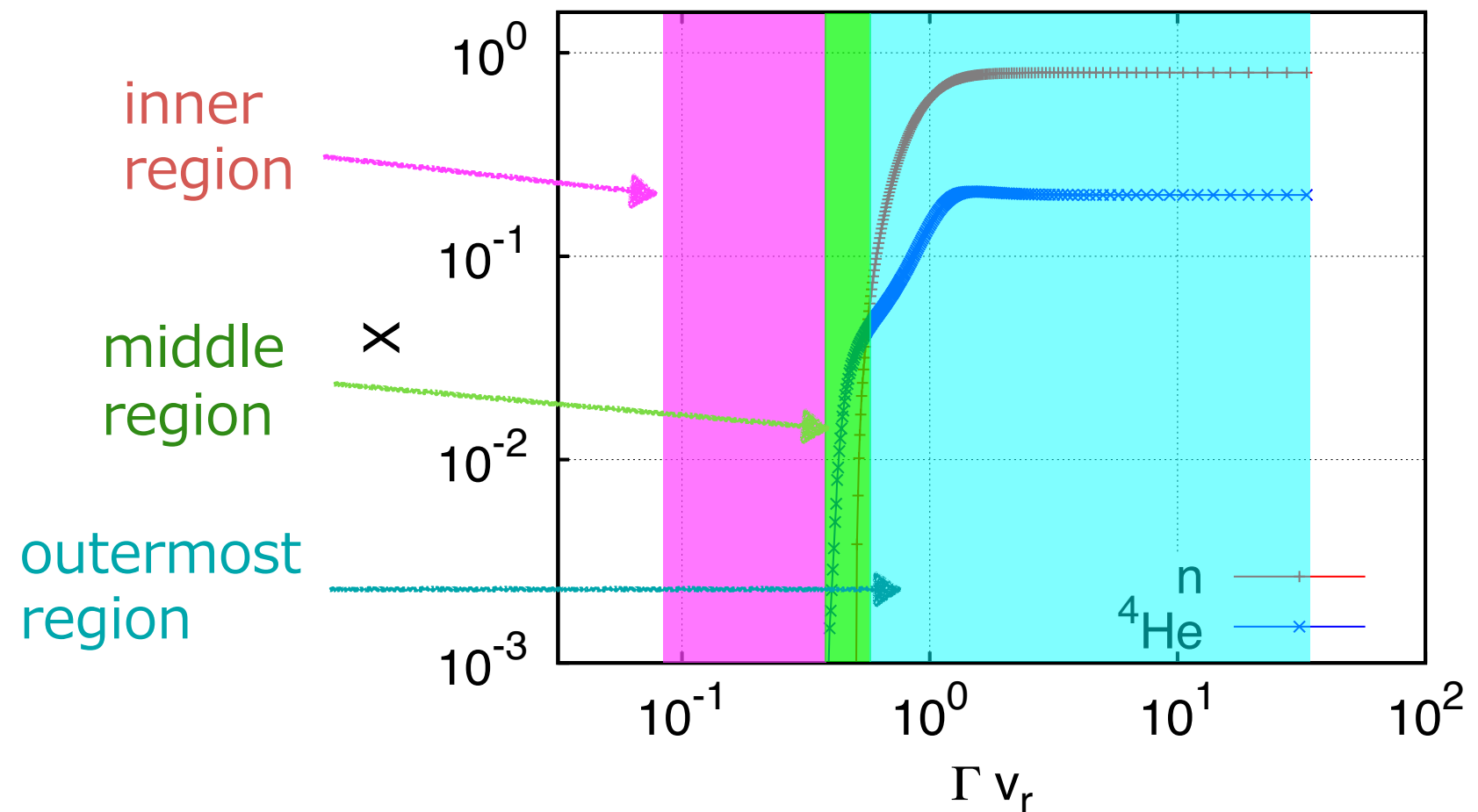
- $p(n,\gamma)d$ reactions produce ${}^4\text{He}$ with $X_n \sim 0.5$

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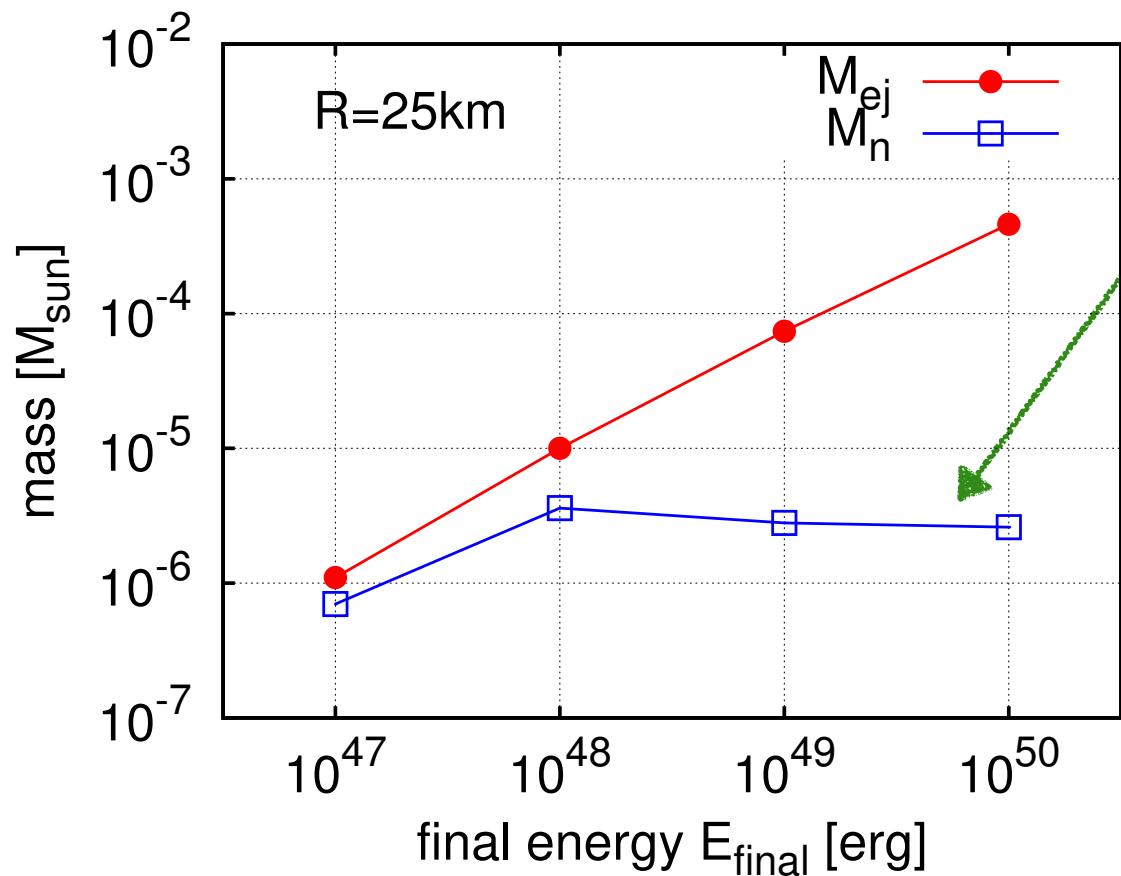
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Total Mass of Free Neutron

R = 25 km



$$M_n = \sum_i (X_{n,i} \times m_i)$$

(Total mass of free neutron layer)

- Preferred energy that yields maximum amount of free neutrons is 10^{48} erg
- M_n value is smaller than previous SPH work by more than 2 orders ($\sim 10^{-4} M_{\text{sun}}$)

R [km] ($E_f = 10^{48}$ erg)	15	20	25	30
M_n [M_{sun}]	9.2×10^{-7}	2.1×10^{-6}	3.6×10^{-6}	5.2×10^{-6}

Emission from free neutron layer

$$M_{\text{ej}} = 10^{-5} M_{\text{sun}}, E_{\text{final}} = 10^{48} \text{ erg, ejecta velocity } \sim c/3, \text{ opacity } \sim 0.4 \text{ cm}^2 \text{ g}^{-1}$$

- Photon diffusion velocity becomes comparable to expansion velocity ($c/3$) at $\sim 1,500$ s
- Energy density at neutron decay time (~ 800 s) is

$$\epsilon_0 \sim 10^6 \text{ erg cm}^{-3} \left(M_n / 3.6 \times 10^{-6} M_{\text{sun}} \right)$$

- Considering subsequent adiabatic expansion up to 1,500 s, luminosity L is estimated by

$$L \sim 7.6 \times 10^{41} \text{ erg s}^{-1} \left(\frac{t}{1,500 \text{ s}} \right)^{-2} \left(\frac{M_n}{3.6 \times 10^{-6} M_{\text{sun}}} \right)$$

(Ultraviolet, timescale of ~ 30 min)

- This is detectable with Swift if observations start immediately after the merger

Summary

- Shock breakout in neutron star merger was reproduced by relativistic Lagrangian hydrodynamics code
- Free neutrons can survive especially with $E_f \sim 10^{48}$ erg
- Total mass of neutron surviving region is $\sim 10^{-6} M_{\text{sun}}$ (two orders smaller than previous SPH work)
→ due to the different ejecta component or low resolution with a small number of SPH particles
- Luminosity of free neutron emission is $\sim 7 \times 10^{41}$ erg s^{-1} in optical band at ~ 30 min after merger event

Future work

- Monte Carlo Radiative transfer computation (A. Ishii et al. 2017) with thermal photons from free neutron decays

Reaction timescales

timescale for positron capture

$$\tau_+ \simeq 2.1 \left(\frac{kT}{\text{MeV}} \right)^{-5} \text{ S} \quad (\text{B. D. Metzger et al. 2015})$$

timescale for electron capture

$$\tau_- \simeq \frac{\tau e^{qz}}{\left(\frac{5.252}{z} - \frac{16.229}{z^2} + \frac{18.059}{z^3} + \frac{34.181}{z^4} + \frac{27.617}{z^5} \right)}$$

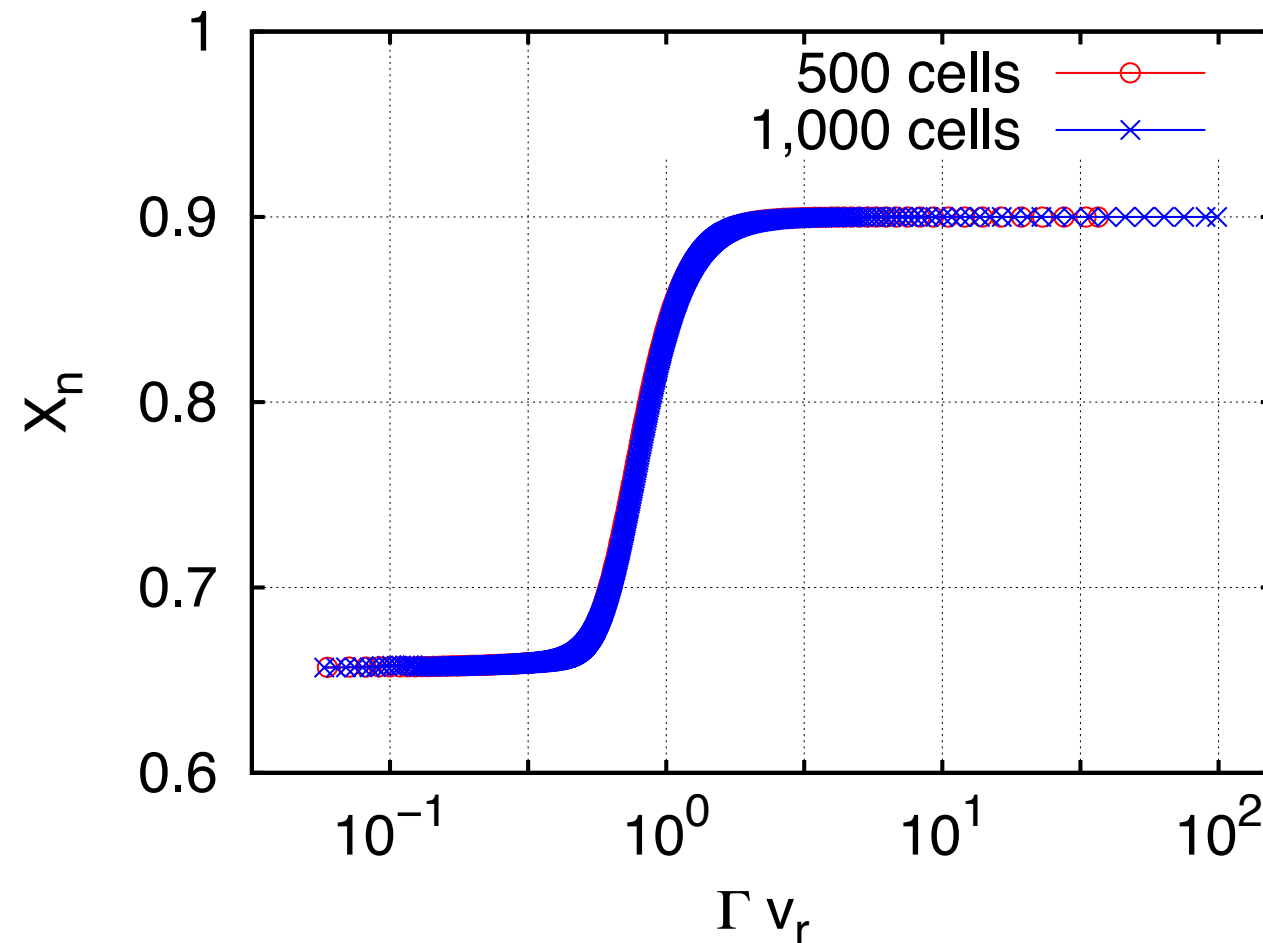
τ : neutron lifetime

q : $(m_n - m_p)/m_e$

z : $m_e c^2/kT$

(L. Kawano 1992)

Grid Convergence Test



- Computation was performed with 1,000 cells
- Distribution extends to larger Γv_r region than 500 cells
- M_n values are almost equivalent
($1.4 \times 10^{-6} M_{\text{sun}}$ and $1.3 \times 10^{-6} M_{\text{sun}}$)
→ Computation with 500 cells are converged

Endothermic reaction

- Disintegration for the heavy element nuclei is endothermic reaction
- The energy density: $a_r T^4 / \rho \sim 3 \times 10^{19} \text{ erg g}^{-1}$ ($T \sim 2.5 \times 10^{11} \text{ K}$)
- The energy of the endothermic reaction for Fe: $\sim 9 \text{ MeV}$ per nucleon
→ Energy density for the reaction: $9 \text{ MeV}/\mu \sim 8.6 \times 10^{18} \text{ erg g}^{-1}$
- Thus the energy is brought out by the endothermic reaction by a few tens of a percent
- The expelled energy might be decreased with the realistic composition of the merging neutron stars