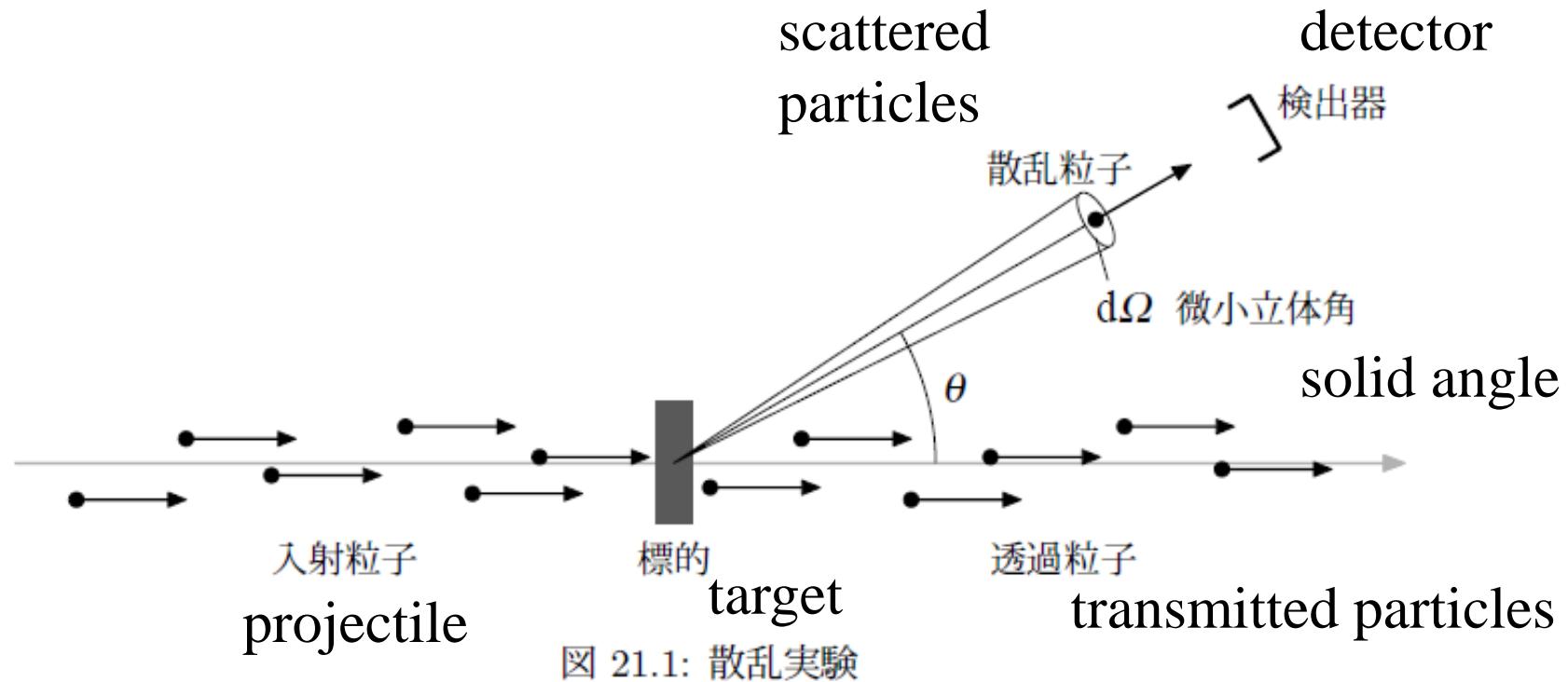


Nuclear Reactions

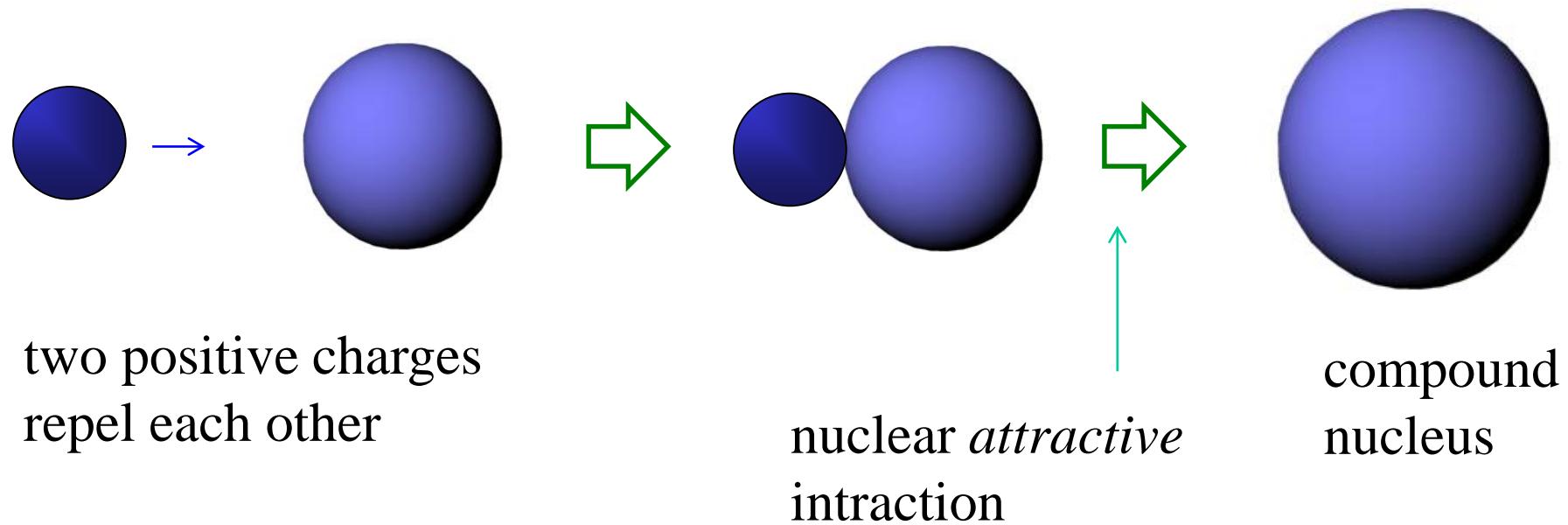
Shape, interaction, and excitation structures of nuclei ← scattering expt.



http://www.th.phys.titech.ac.jp/~muto/lectures/QMII11/QMII11_chap21.pdf

K. Muto (TIT)

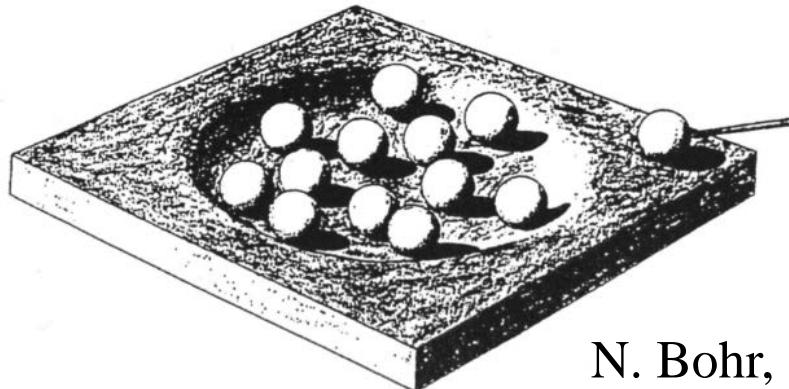
Nuclear fusion reactions



Fusion reactions: compound nucleus formation

Niels Bohr (1936)

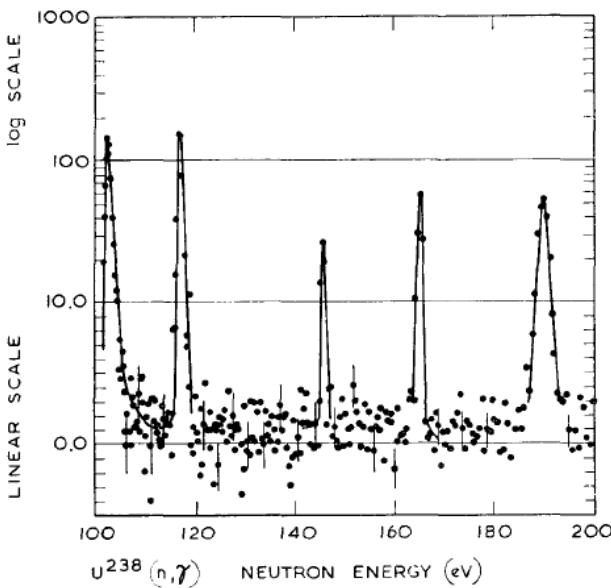
Neutron capture of nuclei → compound nucleus



N. Bohr,
Nature 137 ('36) 351



Wikipedia



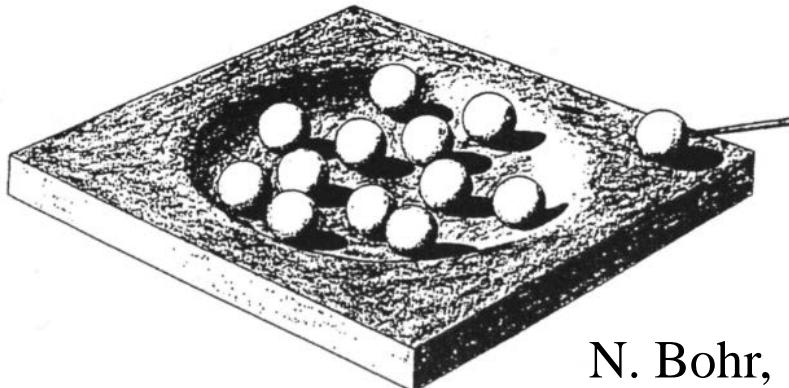
cf. Experiment of Enrico Fermi (1935)
many very narrow (=long life-time)
resonances (width ~ eV)

M. Asghar et al., Nucl. Phys. 85 ('66) 305

Fusion reactions: compound nucleus formation

Niels Bohr (1936)

Neutron capture of nuclei → compound nucleus

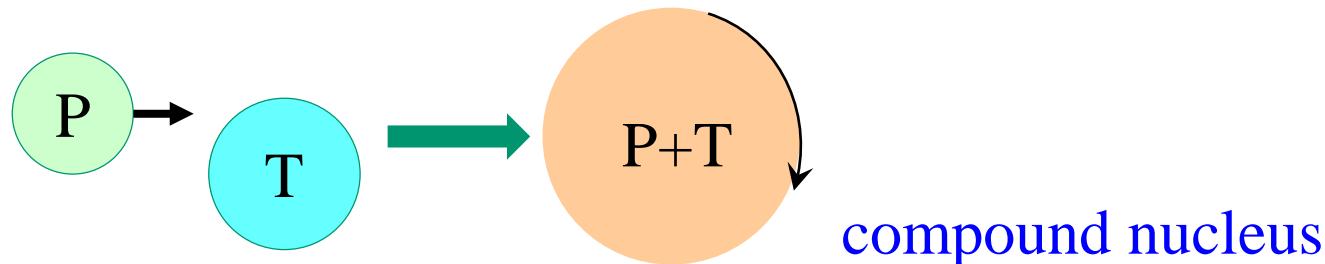


N. Bohr,
Nature 137 ('36) 351

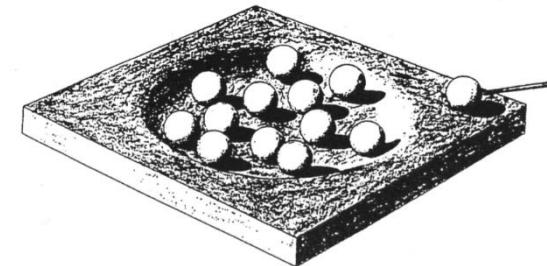
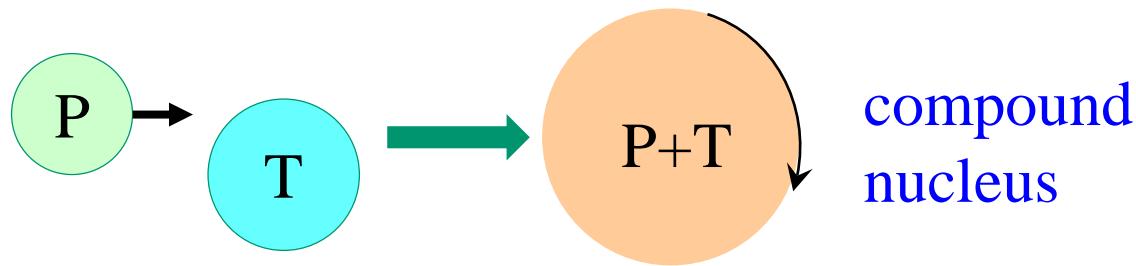


Wikipedia

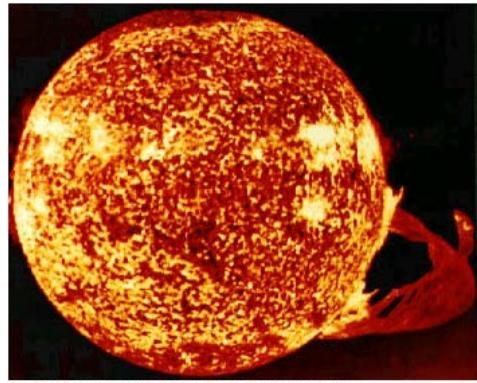
forming a compound nucleus with heavy-ion reactions = H.I. fusion



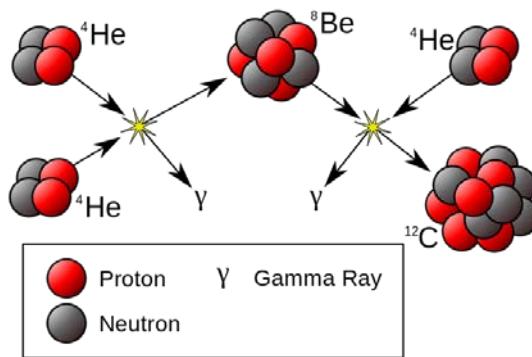
Fusion reactions: compound nucleus formation



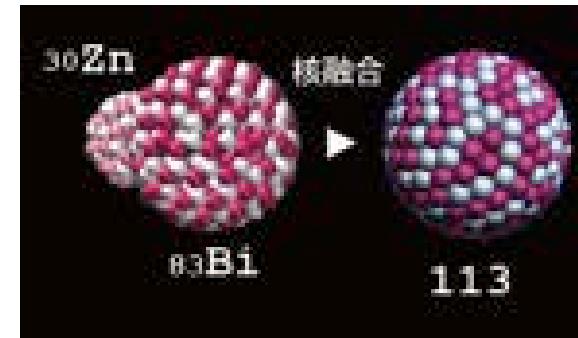
cf. Bohr '36



energy production
in stars (Bethe '39)



nucleosynthesis

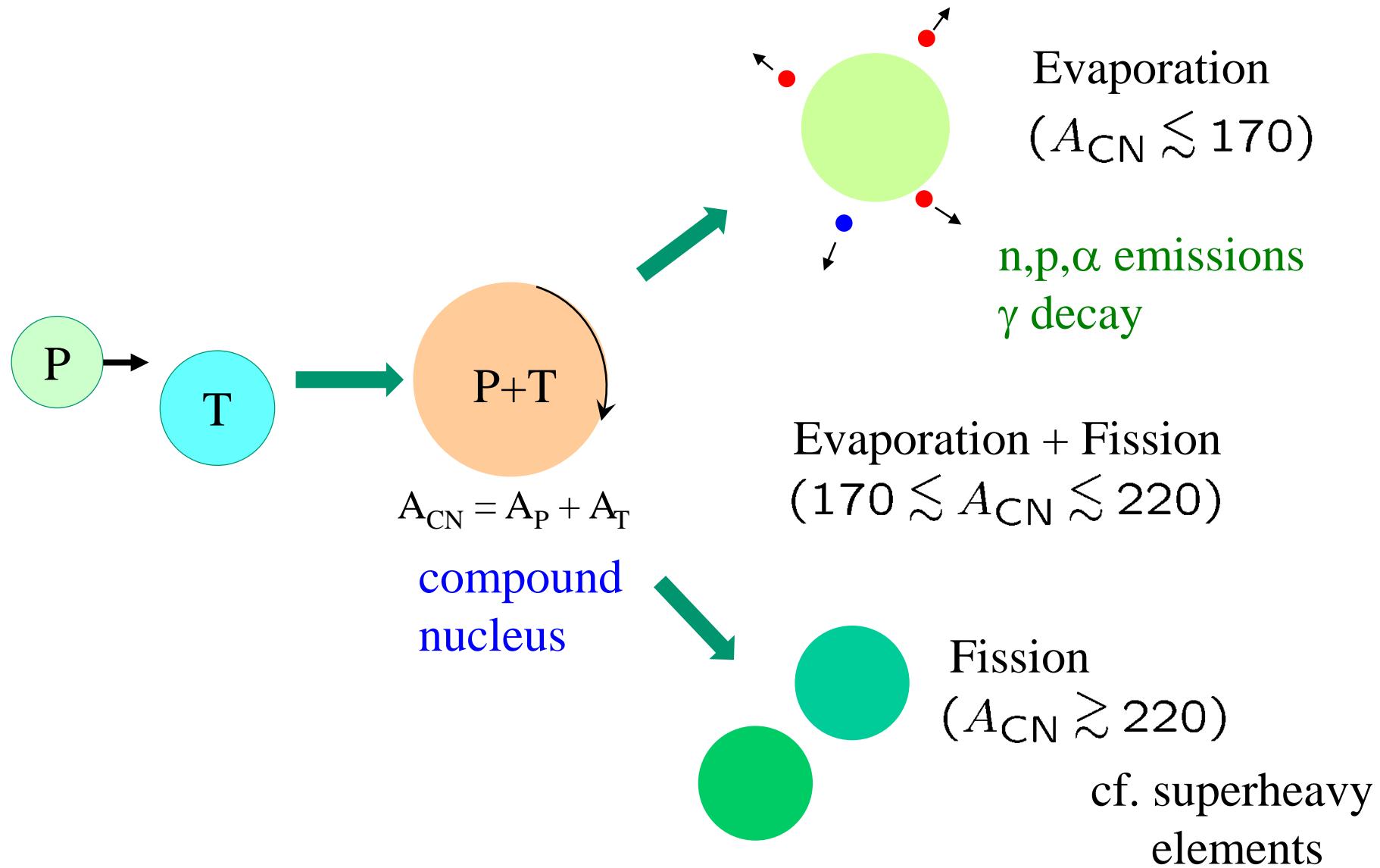


superheavy elements

Fusion and fission: large amplitude motions of quantum many-body systems with strong interaction

← microscopic understanding: an ultimate goal of nuclear physics

Fusion reactions: compound nucleus formation



Inter-nucleus potential

Two interactions:

1. Coulomb force

long range repulsion

2. Nuclear force

short range attraction



potential barrier

due to a cancellation

between the two

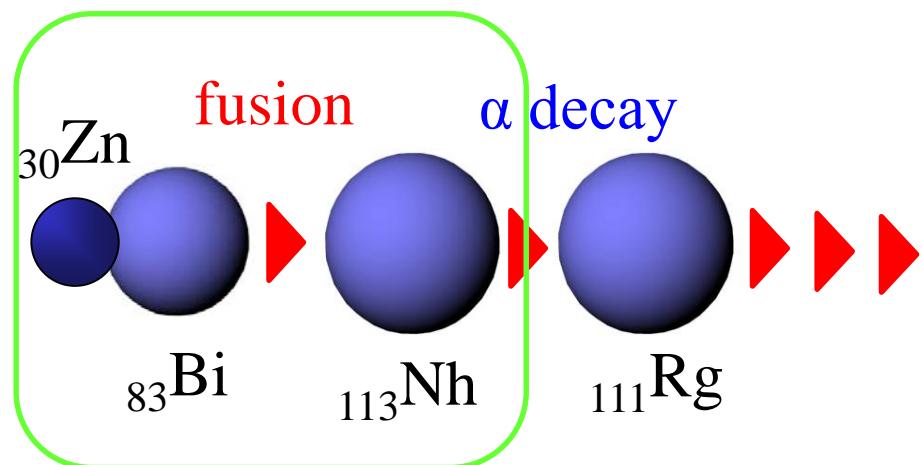
(Coulomb barrier)

- Above-barrier energies
- • Sub-barrier energies
(energies around the Coulomb barrier)
- Deep sub-barrier energies

Why sub-barrier fusion?

two obvious reasons:

113 Nh nihonium	115 Mc moscovium
117 Ts tennessine	118 Og oganesson



superheavy elements

cf. $^{209}\text{Bi} ({}^{70}\text{Zn}, \text{n}) {}^{278}\text{Nh}$

$$V_B \sim 260 \text{ MeV}$$

$$E_{\text{cm}}^{\text{(exp)}} \sim 262 \text{ MeV}$$

Why sub-barrier fusion?

two obvious reasons:

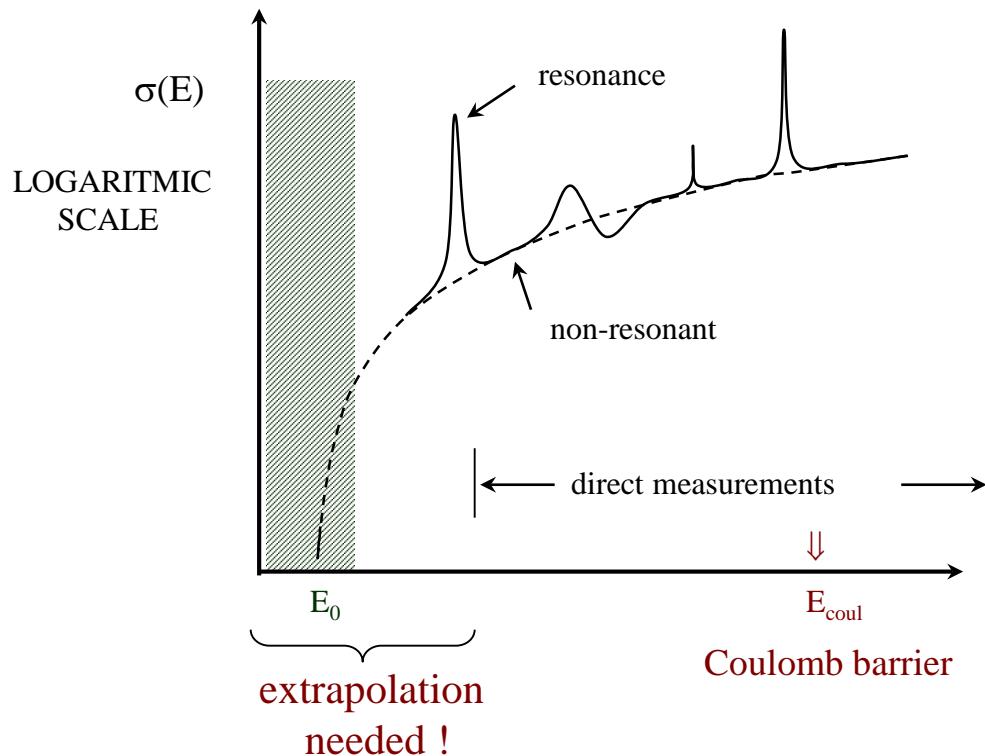
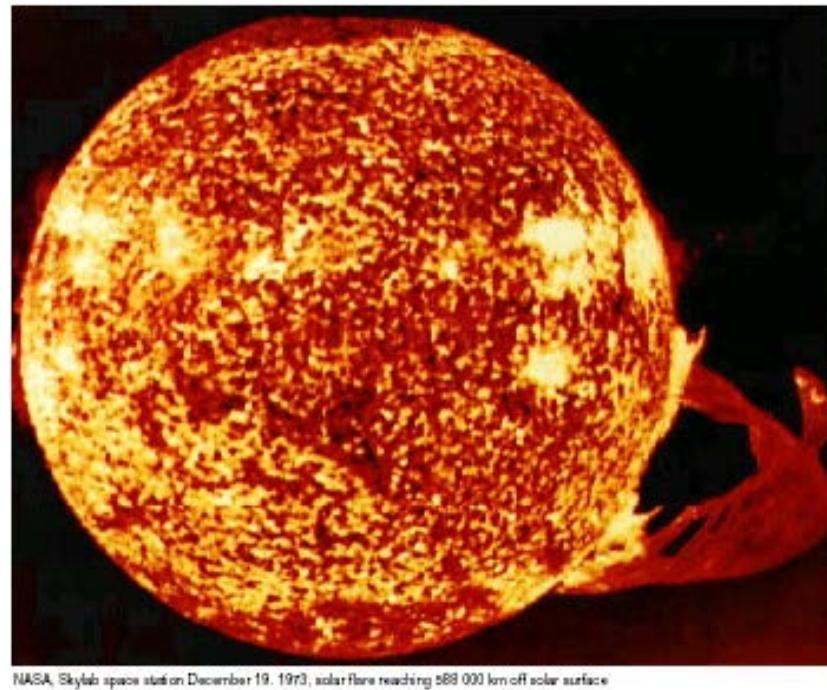


figure: M. Aliotta



nuclear astrophysics
(nuclear fusion in stars)

cf. extrapolation of data

Why sub-barrier fusion?

two obvious reasons:

- ✓ superheavy elements
- ✓ nuclear astrophysics

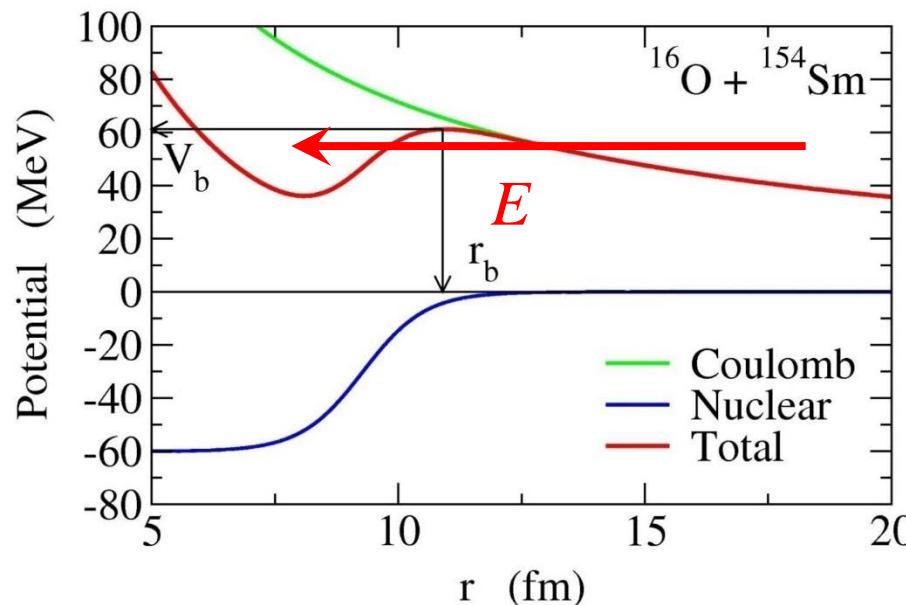
other reasons:

- ✓ reaction dynamics

strong interplay between reaction and structure

cf. high E reactions: much simpler reaction mechanisms

- ✓ many-particle tunneling



Why sub-barrier fusion?

two obvious reasons:

- ✓ superheavy elements
- ✓ nuclear astrophysics

other reasons:

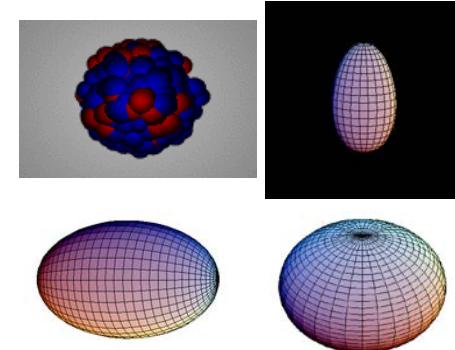
- ✓ reaction dynamics

strong interplay between reaction and structure

cf. high E reactions: much simpler reaction mechanisms

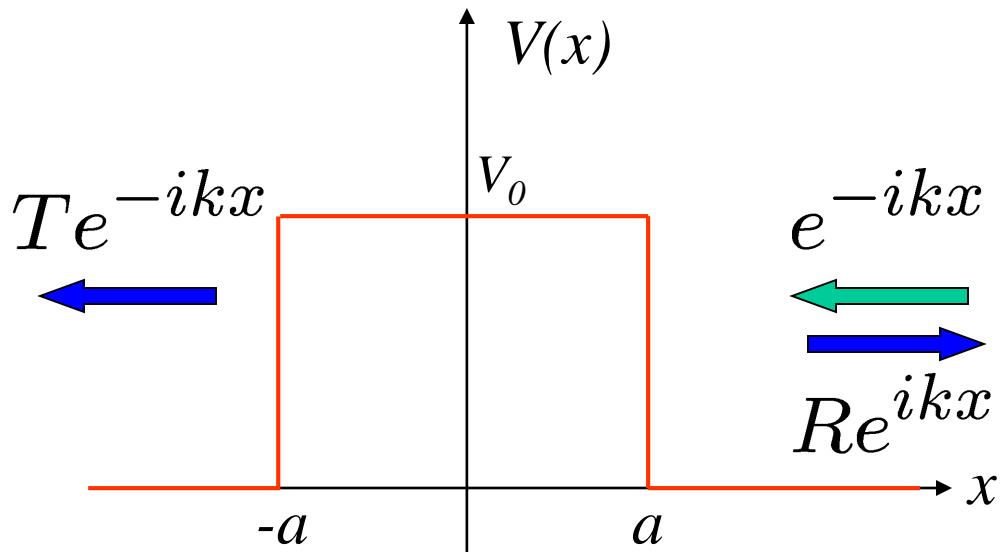
- ✓ many-particle tunneling

- many types of intrinsic degrees of freedom
(several types of collective vibrations,
deformation with several multipolarities)
- energy dependence of tunneling probability
cf. alpha decay: fixed energy



H.I. fusion reaction = an ideal playground to study quantum
tunneling with many degrees of freedom

Quantum Tunneling Phenomena

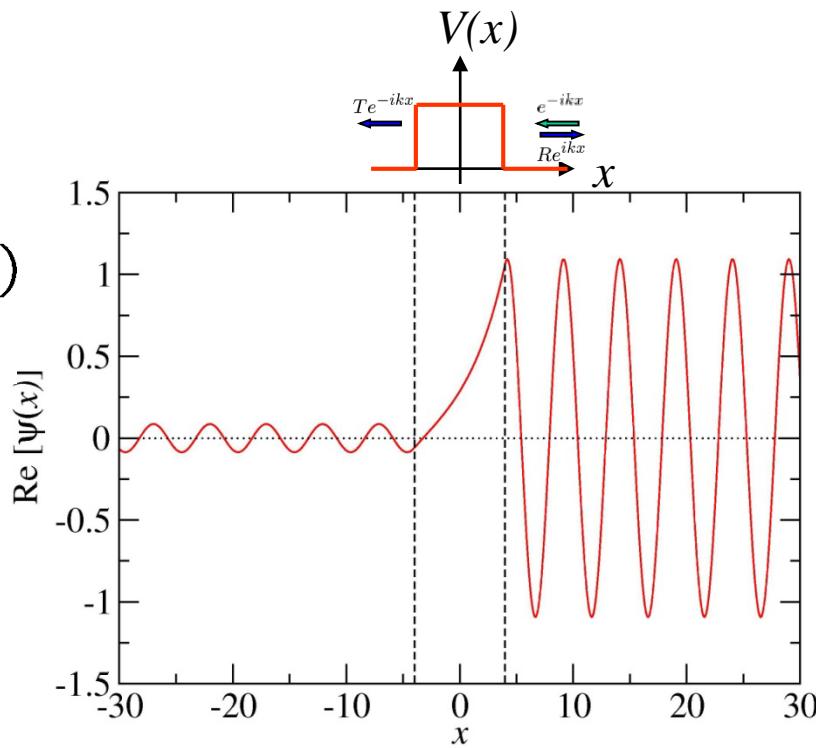


$$\begin{aligned}\psi(x) &= T e^{-ikx} & (x \leq -a) \\ &= A e^{-\kappa x} + B e^{\kappa x} & (-a < x < a) \\ &= e^{-ikx} + R e^{ikx} & (x \geq a)\end{aligned}$$

$$k = \sqrt{2mE/\hbar^2}$$

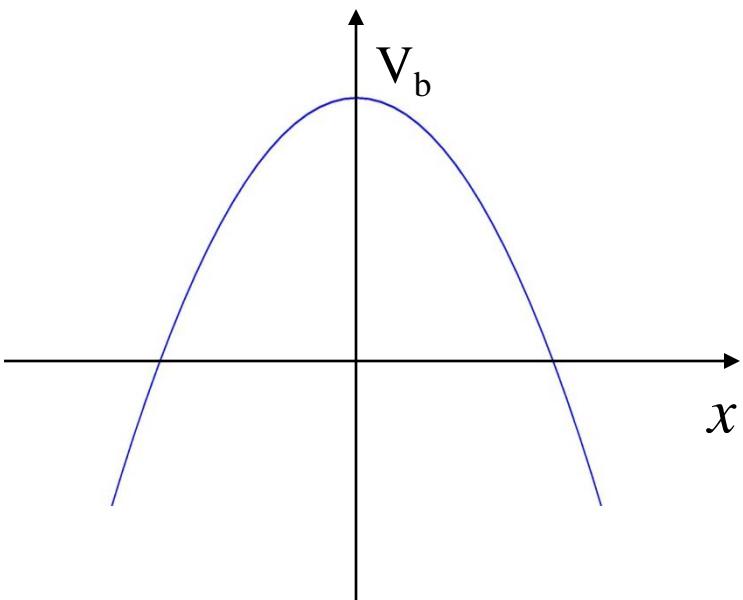
$$\kappa = \sqrt{2m(V_0 - E)/\hbar^2}$$

Tunnel probability: $P(E) = |T|^2$

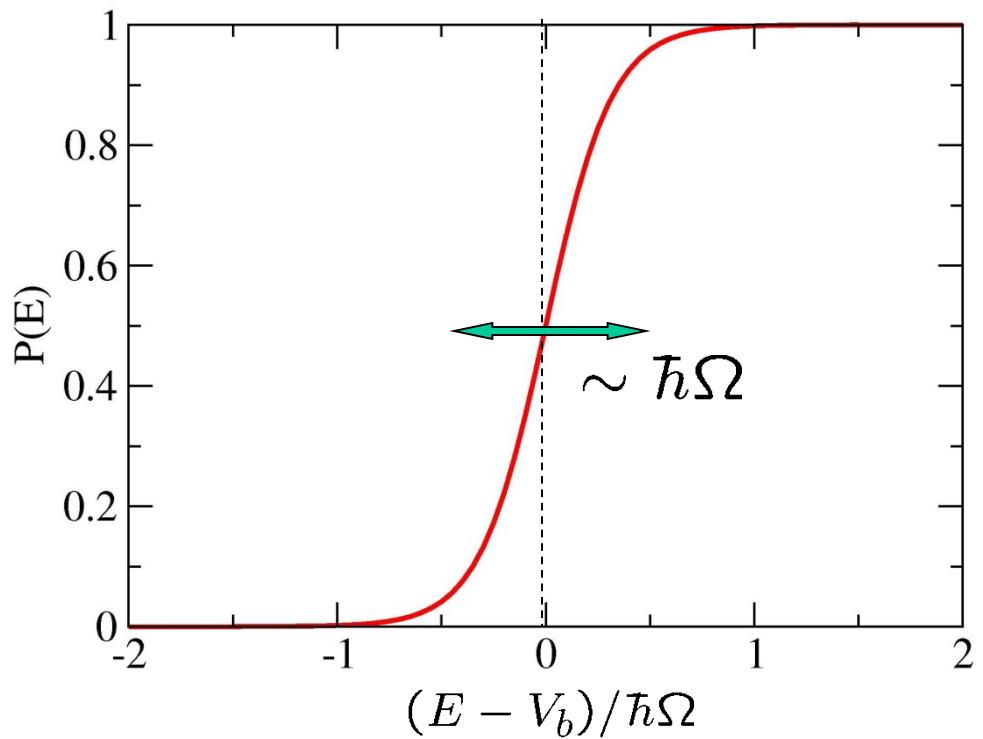


For a parabolic barrier.....

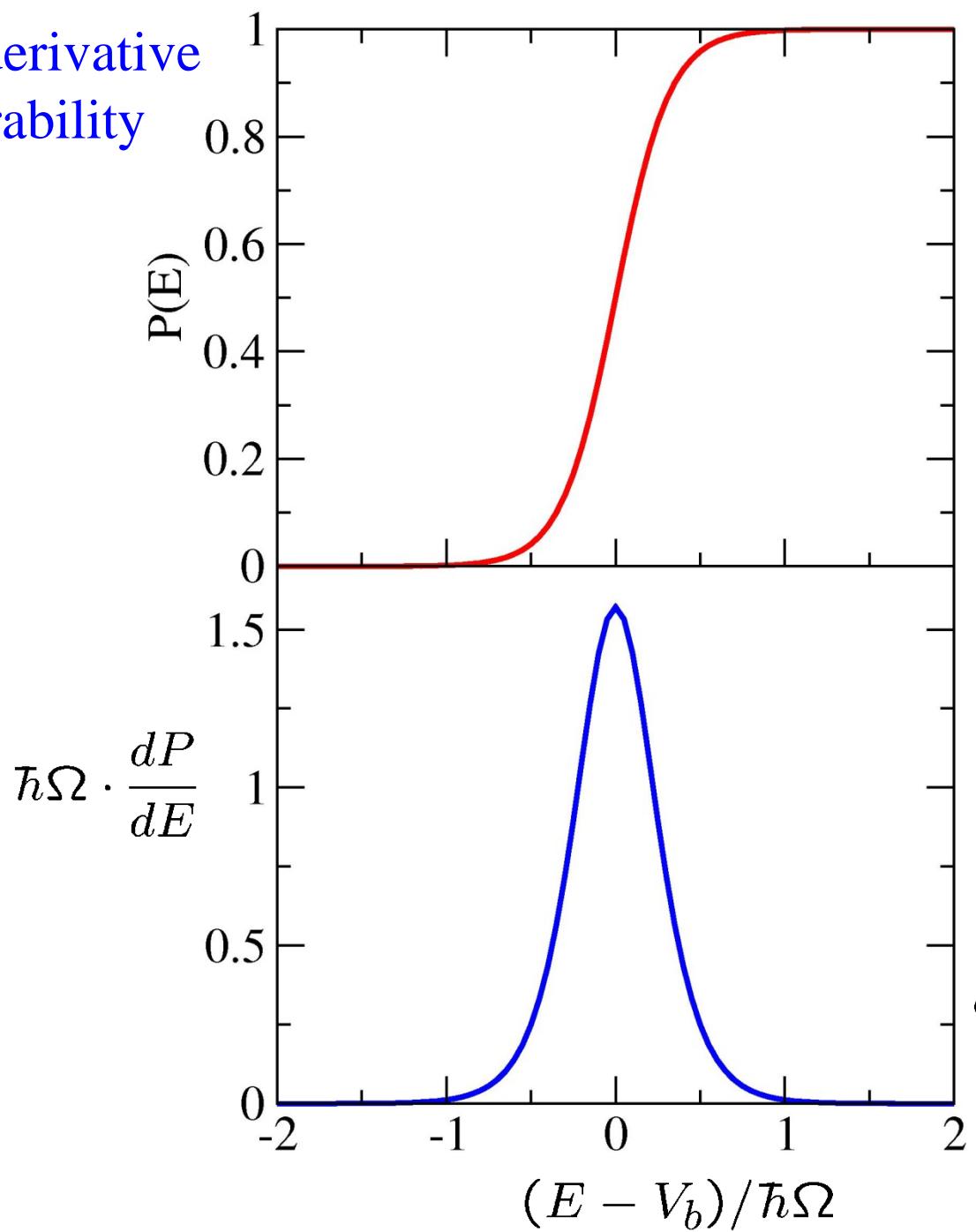
$$V(x) = V_b - \frac{1}{2}m\Omega^2x^2$$



$$P(E) = \frac{1}{1 + \exp \left[\frac{2\pi}{\hbar\Omega} (V_b - E) \right]}$$



Energy derivative of penetrability



(note) Classical limit

$$P(E) = \theta(E - V_b)$$

$$dP/dE = \delta(E - V_b)$$

cf. WKB approximation

One dimensional Schrodinger equation:

$$-\frac{\hbar^2}{2m} \frac{d^2}{dx^2} \psi(x) + (V(x) - E)\psi(x) = 0$$

→ $\frac{d^2}{dx^2} \psi(x) + \frac{p^2(x)}{\hbar^2} \psi(x) = 0, \quad p(x) \equiv \sqrt{2m(E - V(x))}$

Assume

$$\psi(x) = e^{iS(x)/\hbar}$$

$$\rightarrow \psi' = \frac{i}{\hbar} S' \psi$$

$$\psi'' = \frac{i}{\hbar} S'' \psi - \frac{1}{\hbar^2} (S')^2 \psi$$

$$\rightarrow \frac{i}{\hbar} S'' - \frac{1}{\hbar^2} (S')^2 + \frac{p(x)^2}{\hbar^2} = 0$$

cf. WKB approximation

$$i\hbar S'' - (S')^2 + p(x)^2 = 0$$

Expand S as: $S(x) = S_0(x) + \hbar S_1(x) + \hbar^2 S_2(x) + \dots$

$$S_0(x) = \pm \int^x p(x') dx'$$

$$S_1(x) = \frac{i}{2} \ln p(x) + const.$$



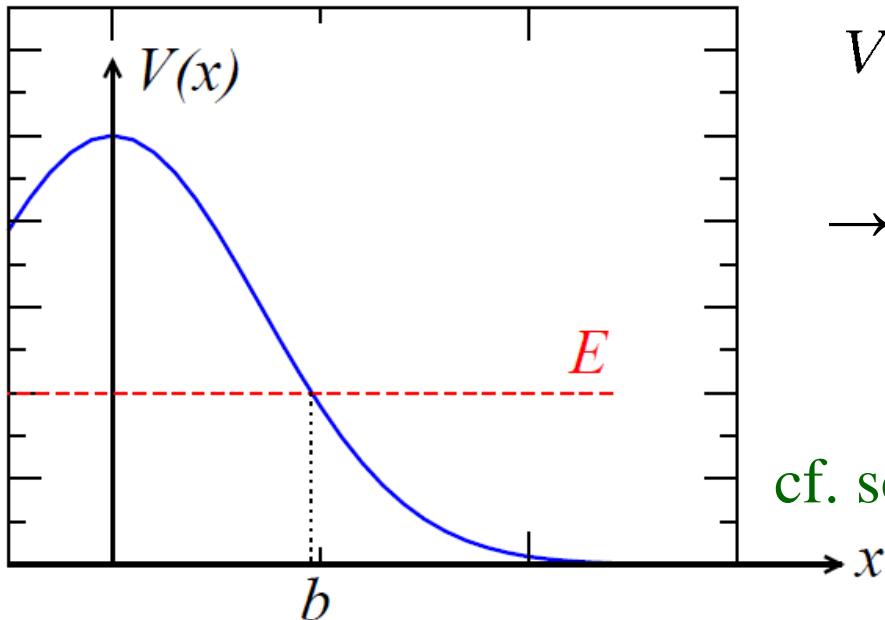
$$\psi(x) = e^{iS(x)/\hbar} \sim \frac{1}{\sqrt{p(x)}} e^{\pm \frac{i}{\hbar} \int^x p(x') dx'}$$

cf. WKB approximation

$$\psi(x) = e^{iS(x)/\hbar} = \frac{1}{\sqrt{p(x)}} e^{\pm \frac{i}{\hbar} \int^x p(x') dx'}$$

this wave function breaks down at x which satisfies $p(x) = 0$.

$$p(x) = \sqrt{2m(E - V(x))} \rightarrow V(x) = E \quad (\text{a turning point})$$

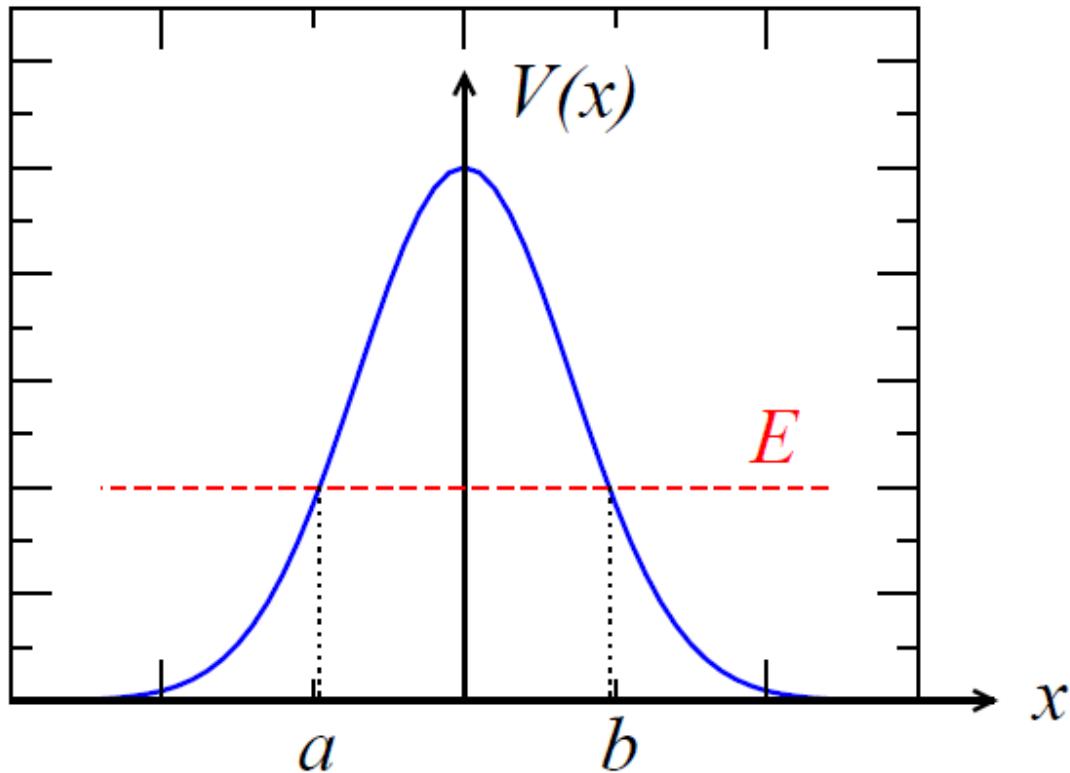


$V(x) \sim V(b) + V'(b)(x - b)$
around $x \sim b$
→ connect wave function from $x > b$
to $x < b$
(WKB connection formula)

cf. solution for a linear potential
: Airy function

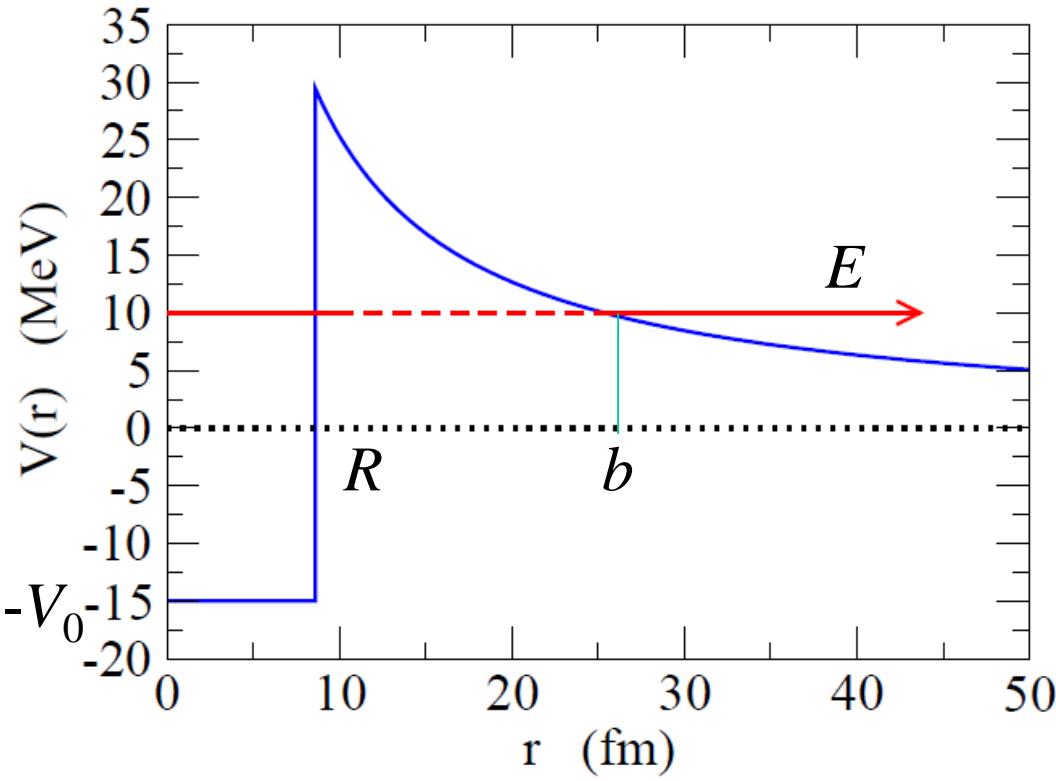
cf. WKB approximation

if applied to a tunneling problem:



$$P(E) = \exp \left[-2 \int_a^b dx \sqrt{\frac{2m}{\hbar^2} (V(x) - E)} \right]$$

for a Coulomb potential



$$\begin{aligned} V(r) &= -V_0 & (r < R) \\ &= \frac{Z_1 Z_2 e^2}{r} & (r \geq R) \end{aligned}$$

$$P = \exp \left[-2 \int_R^b dr \sqrt{\frac{2\mu}{\hbar^2} \left(\frac{Z_1 Z_2 e^2}{r} - E \right)} \right]$$

$$P = \exp \left[-2 \int_R^b dr \sqrt{\frac{2\mu}{\hbar^2} \left(\frac{Z_1 Z_2 e^2}{r} - E \right)} \right]$$

$$P = e^{-G} \quad G \sim \sqrt{\frac{2\mu R}{\hbar^2} Z_1 Z_2 e^2} \left(\pi \sqrt{\frac{Z_1 Z_2 e^2}{R E}} - 4 \right)$$

(note) for $R \rightarrow 0$

$$G \sim \sqrt{\frac{2\mu R}{\hbar^2} Z_1 Z_2 e^2} \left(\pi \sqrt{\frac{Z_1 Z_2 e^2}{R E}} \right) = \pi Z_1 Z_2 e^2 \sqrt{\frac{2\mu}{\hbar^2 E}} = 2\pi \frac{Z_1 Z_2 e^2}{\hbar v}$$

$\equiv \eta(E)$

$$P(E) \sim e^{-2\pi\eta(E)}$$

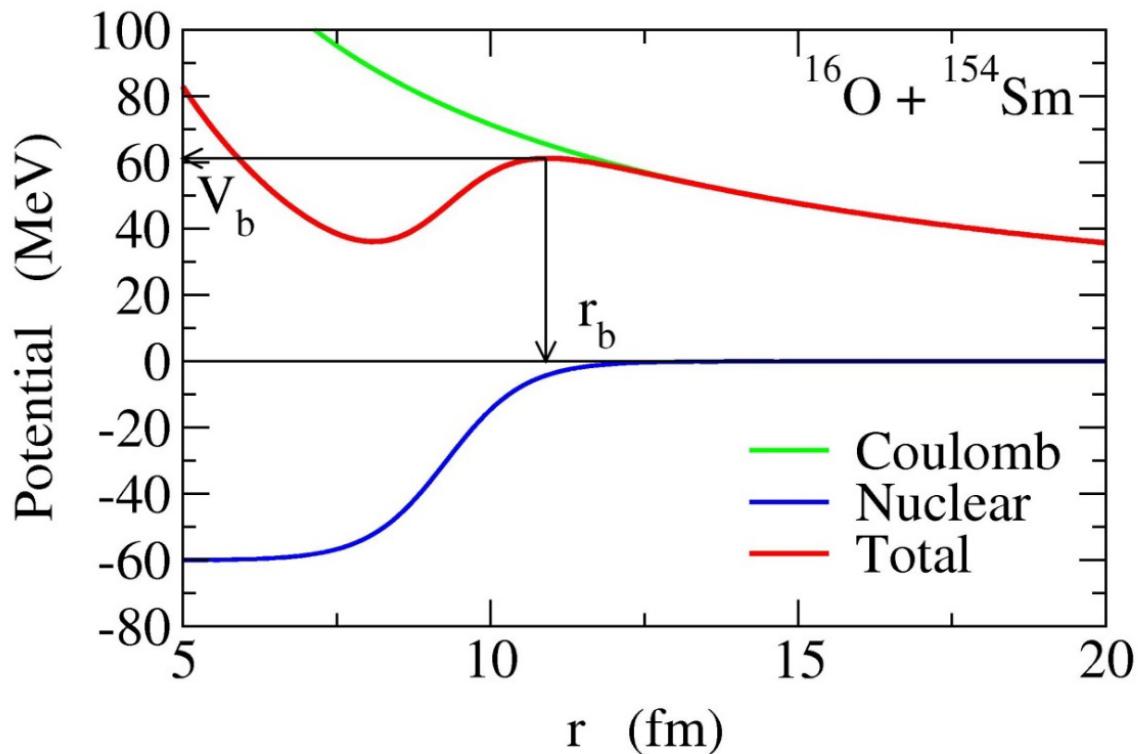
Sommerfeld
parameter

The simplest approach to fusion: potential model

Potential model: $V(r)$ + absorption

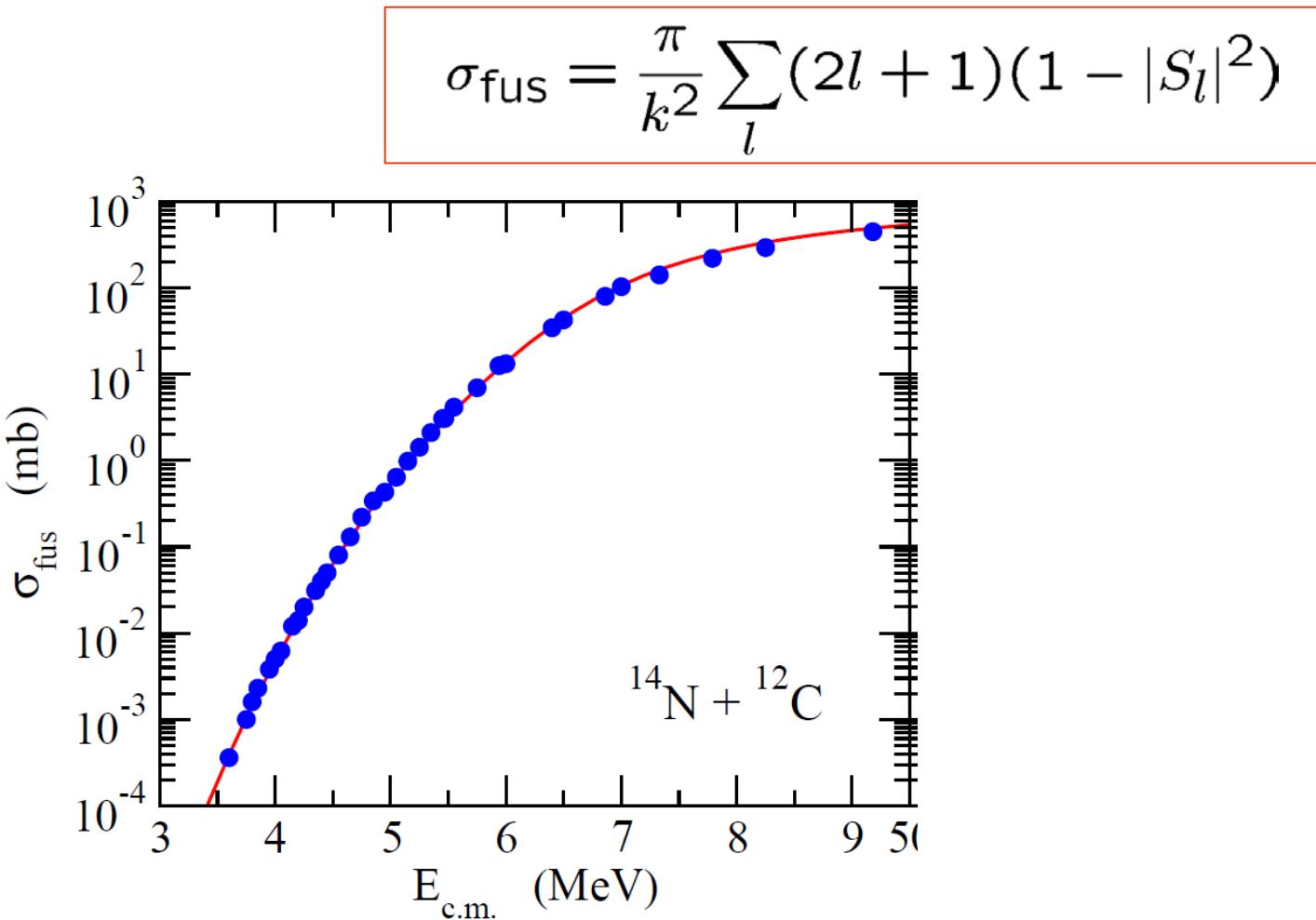
$$\sigma_{\text{fus}}(E) = \frac{\pi}{k^2} \sum_l (2l + 1) P_l(E)$$

$P_l(E)$: barrier penetrability



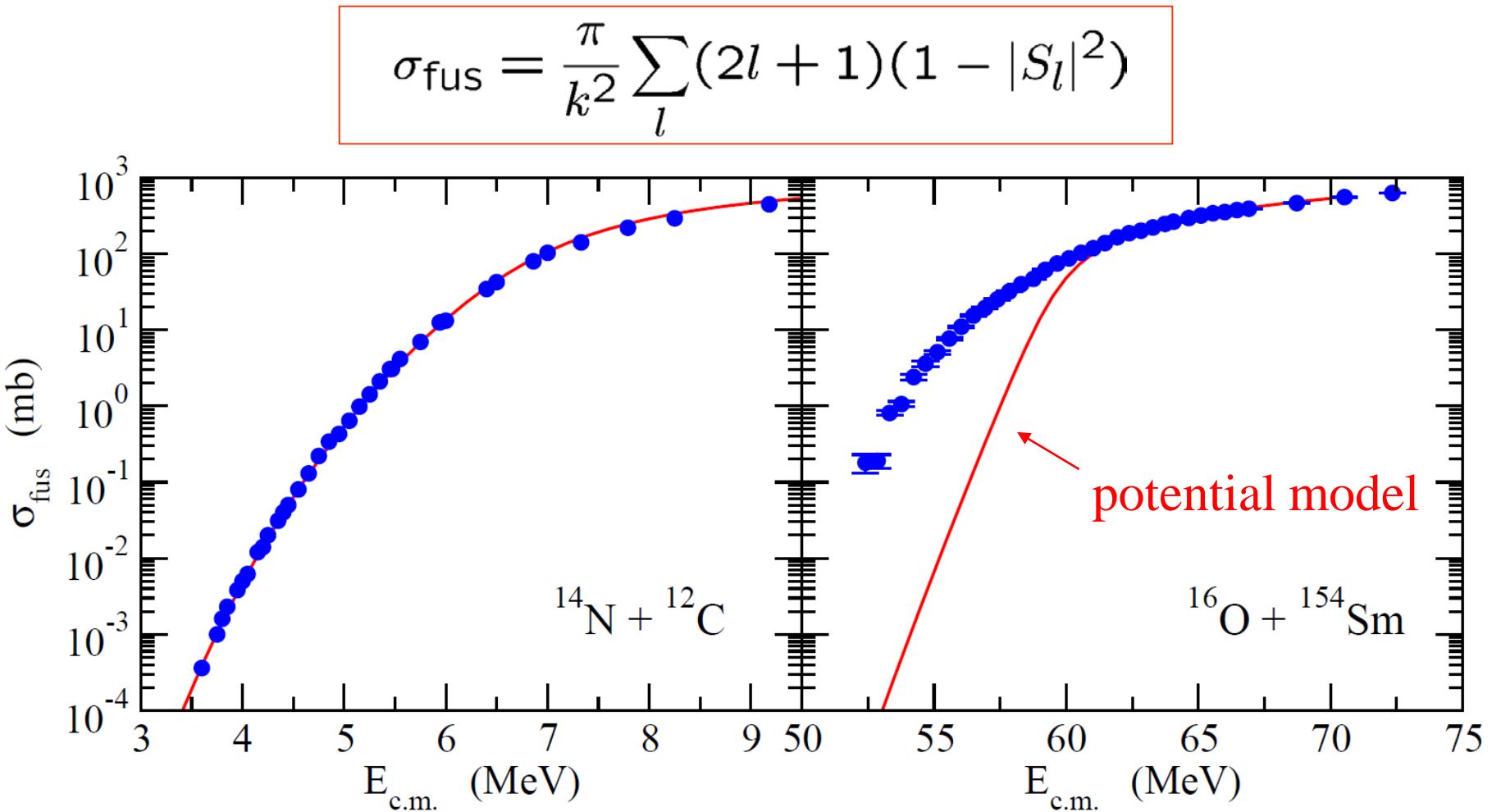
Comparison with experimental data: large enhancement of σ_{fus}

Potential model: $V(r)$ + absorption



Comparison with experimental data: large enhancement of σ_{fus}

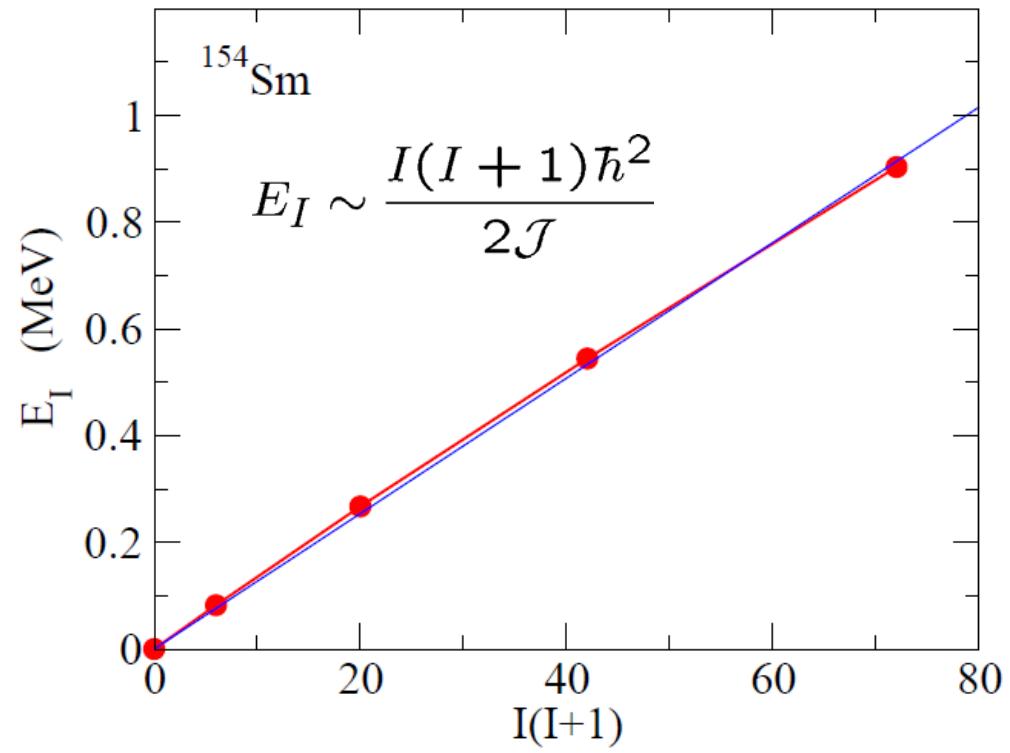
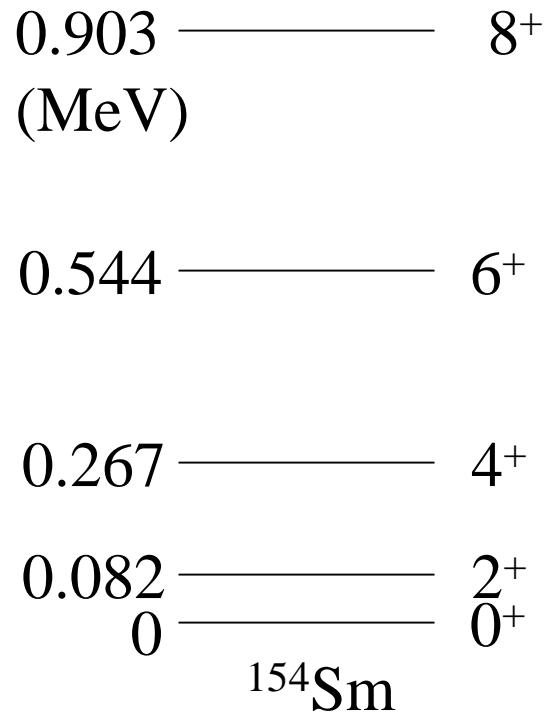
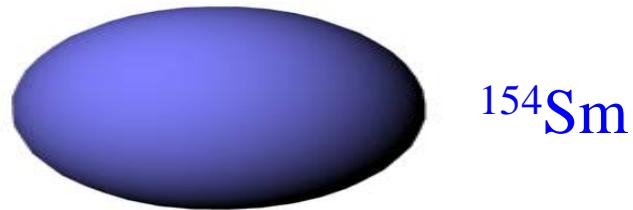
Potential model: $V(r)$ + absorption



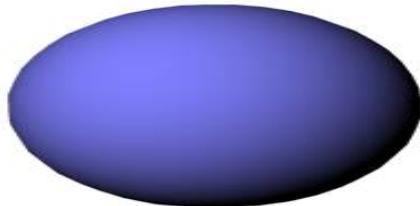
cf. seminal work:

R.G. Stokstad et al., PRL41('78) 465

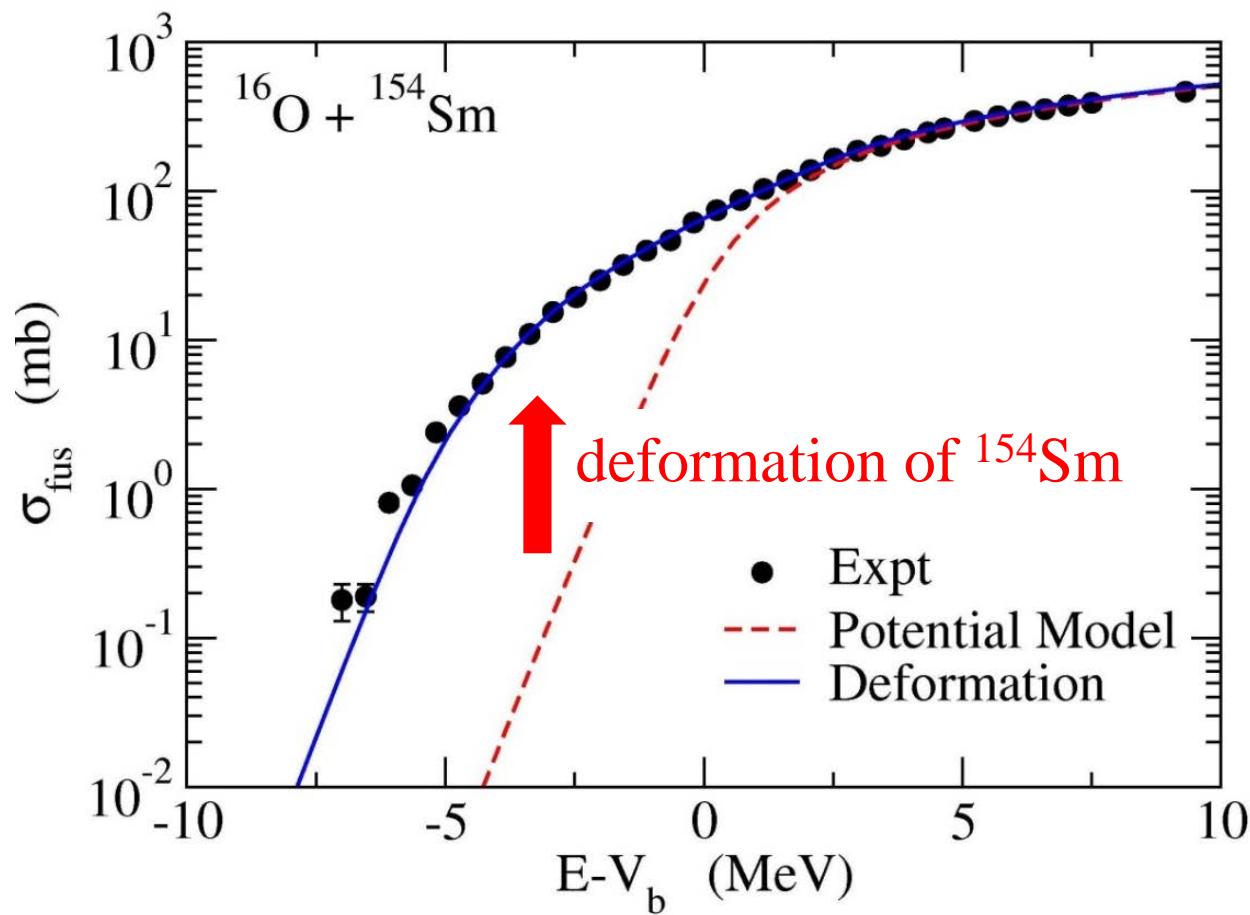
^{154}Sm : a typical deformed nucleus



^{154}Sm : a typical deformed nucleus

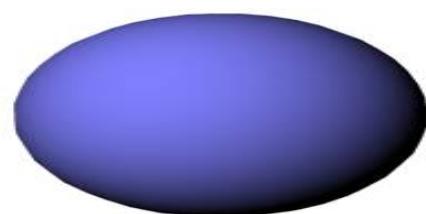


^{154}Sm

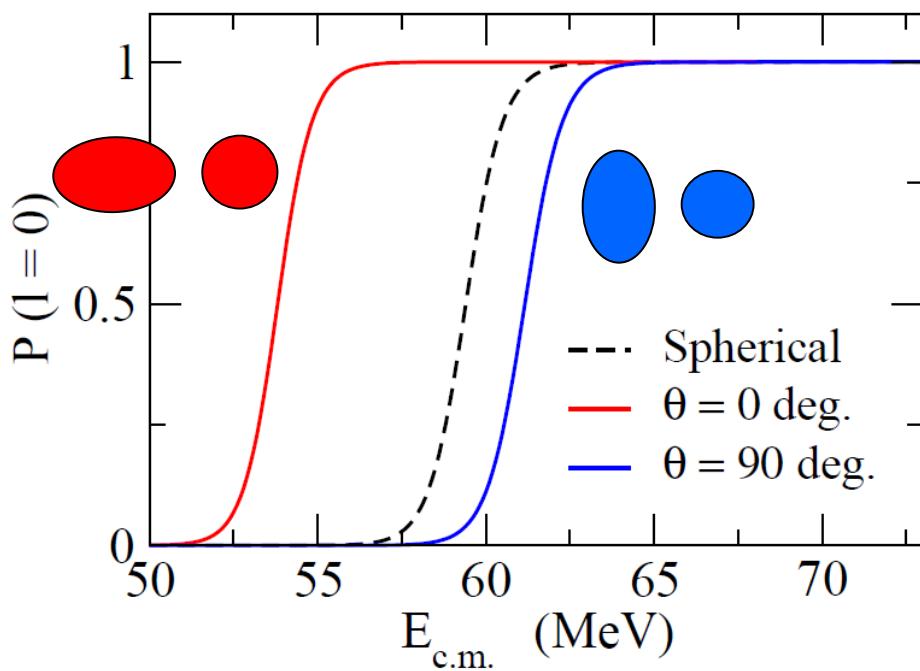
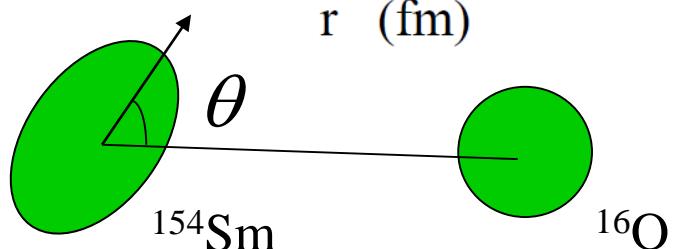
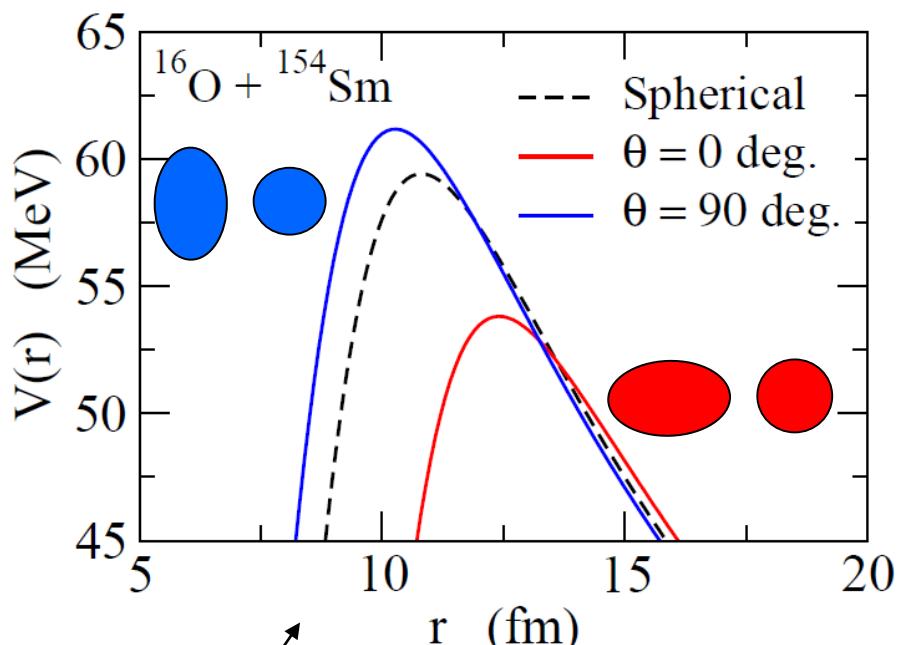


Effects of nuclear deformation

^{154}Sm : a typical deformed nucleus

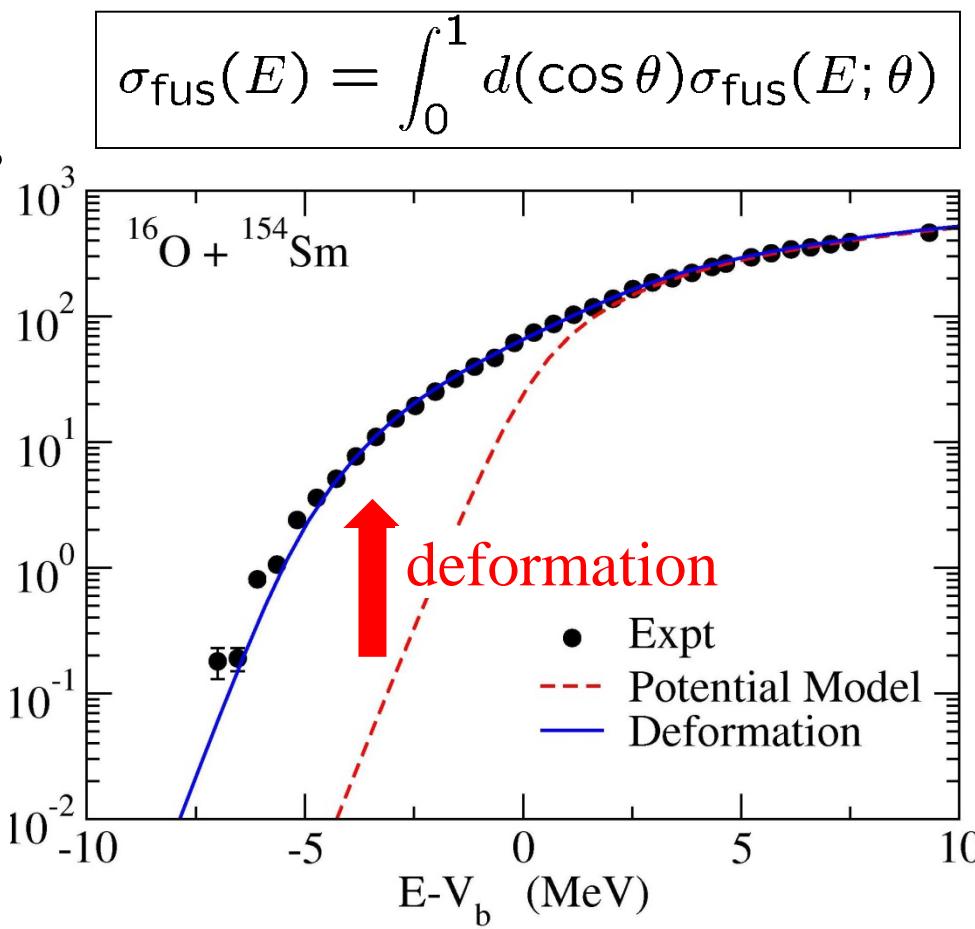
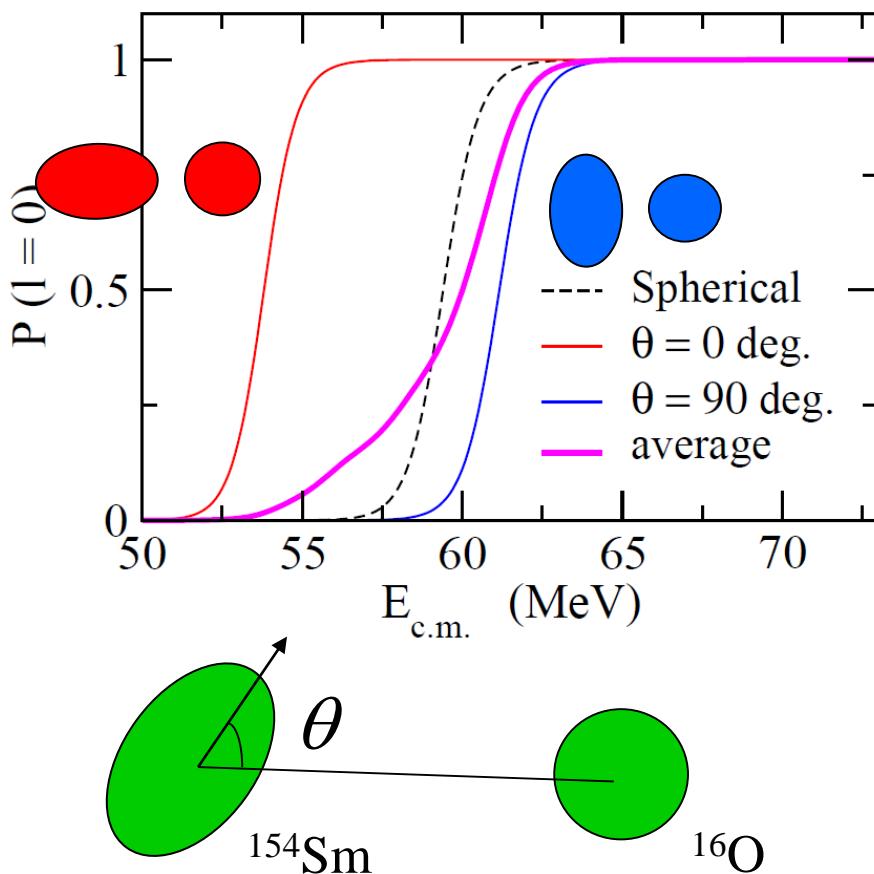


^{154}Sm



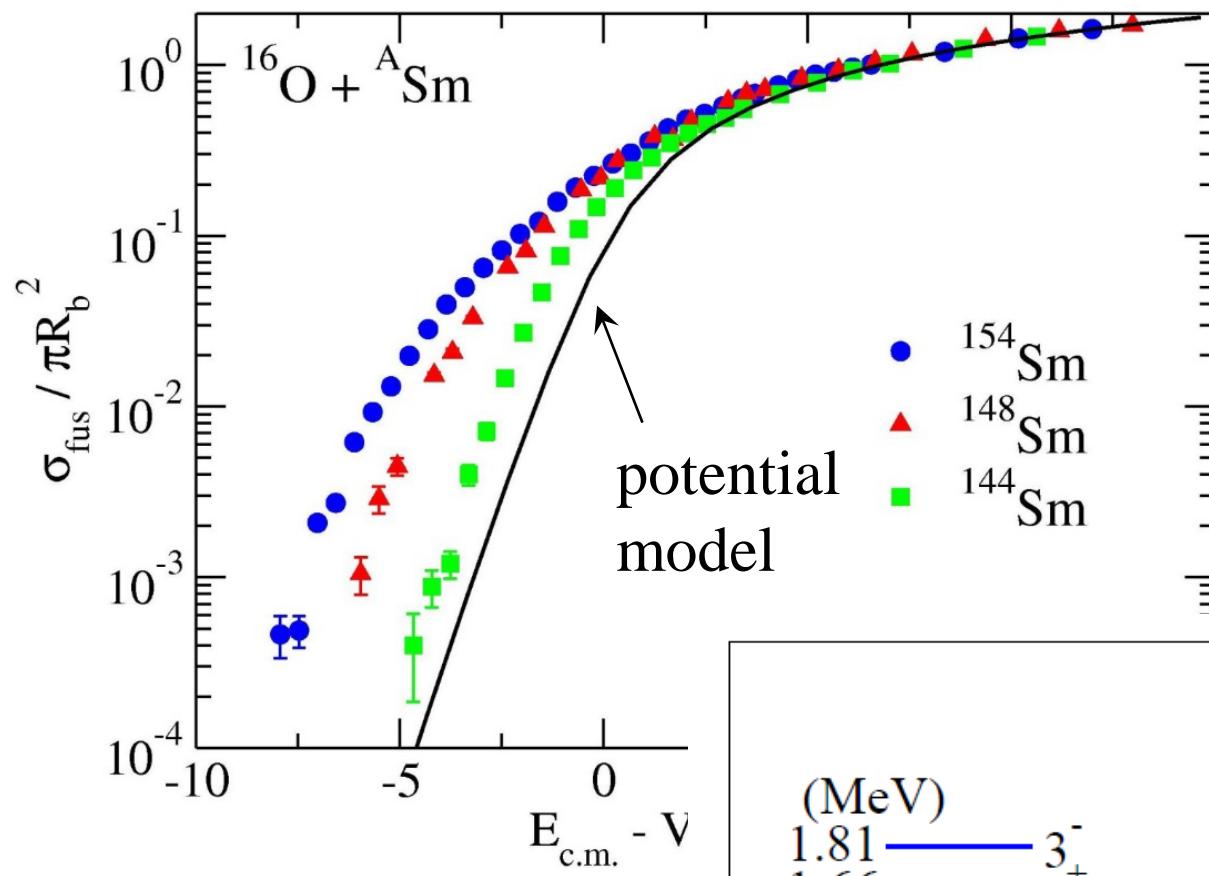
Effects of nuclear deformation

^{154}Sm : a typical deformed nucleus



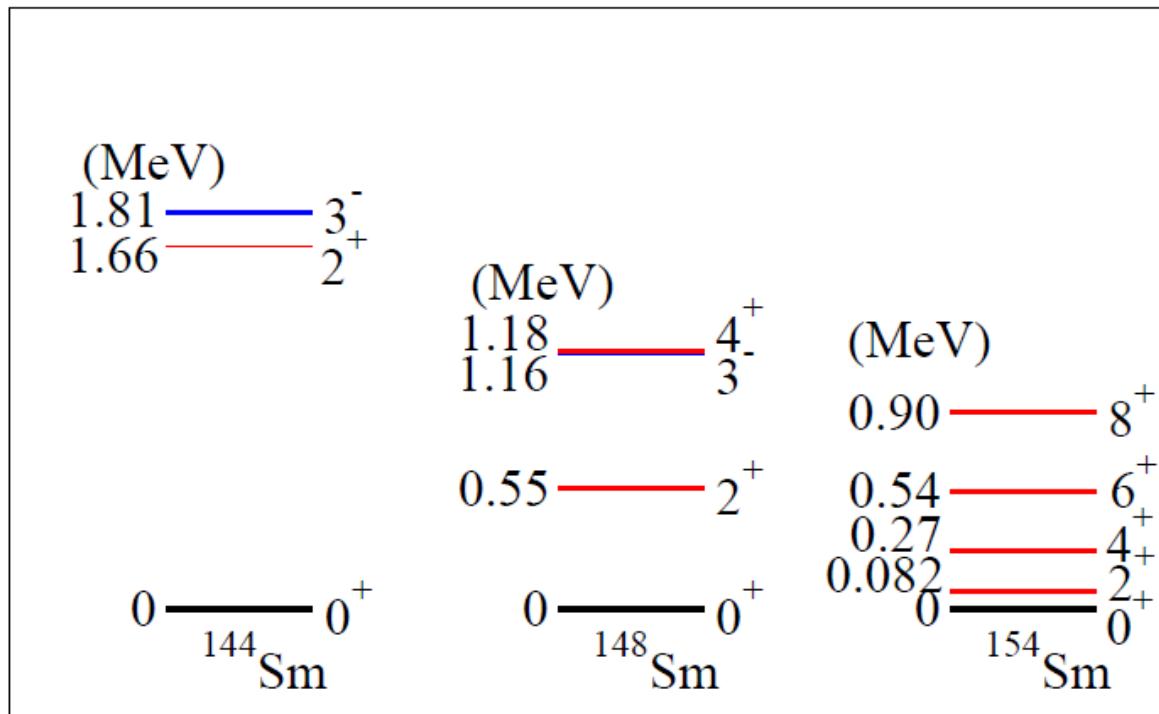
Fusion: strong interplay between nuclear structure and reaction

* Sub-barrier enhancement also for non-deformed targets:
couplings to low-lying collective excitations → coupling assisted tunneling



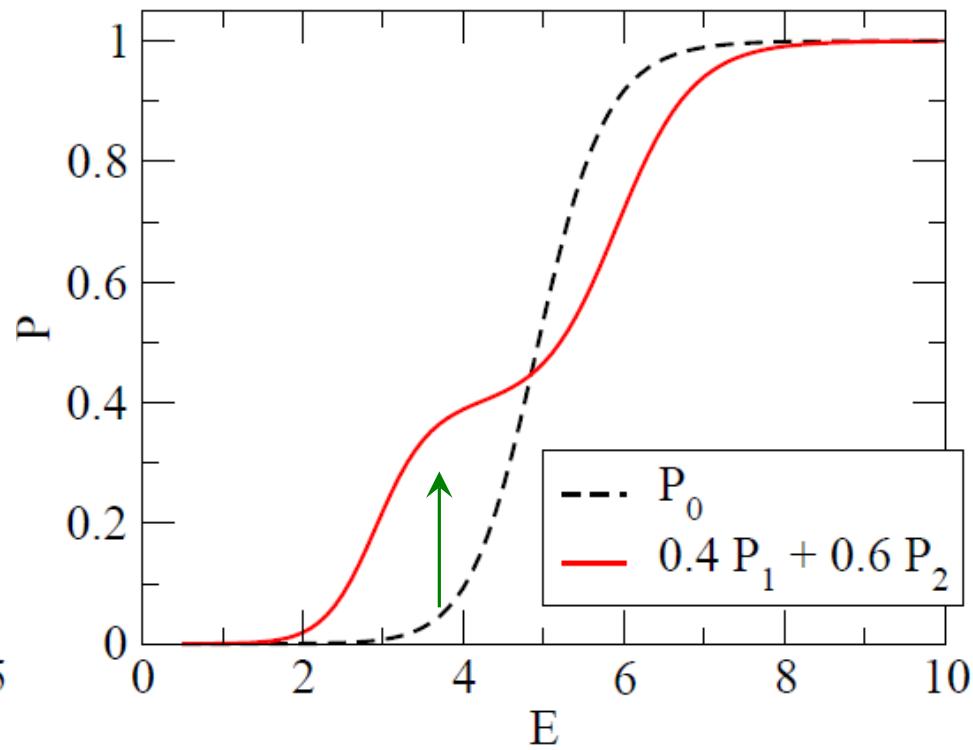
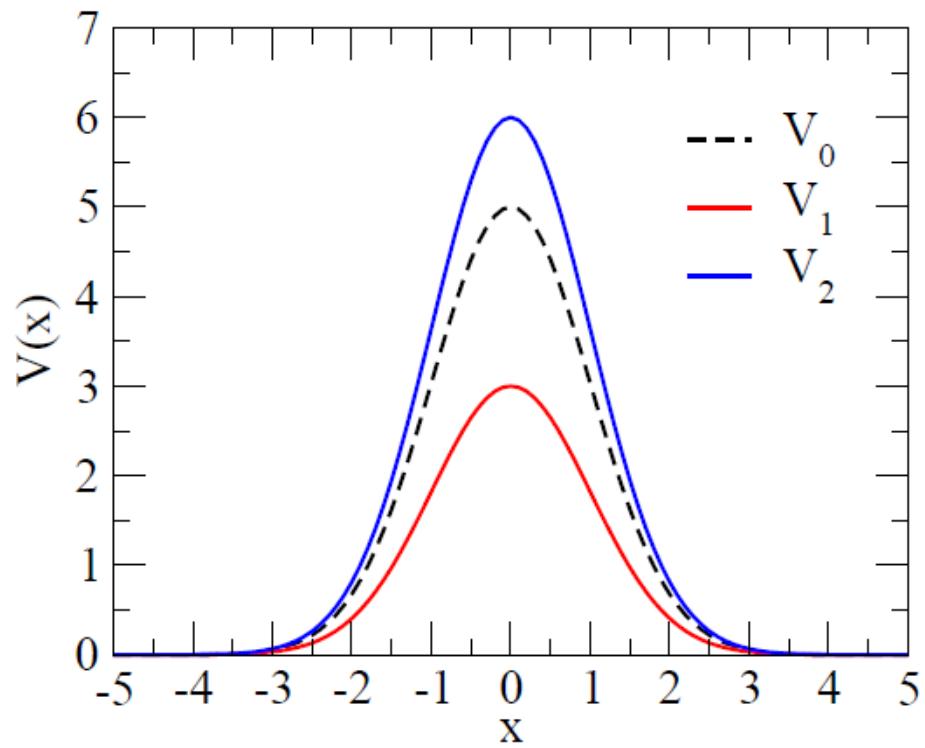
enhancement of fusion
cross sections
: a general phenomenon

strong correlation with
nuclear spectrum
→ coupling assisted
tunneling



Enhancement of tunneling probability : a problem of two potential barriers

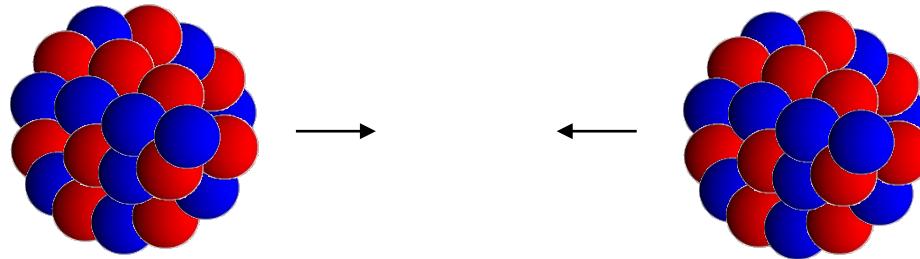
$$P(E) = P(E; V_0) \rightarrow w_1 P(E; V_1) + w_2 P(E; V_2)$$



“barrier distribution” due to couplings to excited states
in projectile/target nuclei

Coupled-channels method: a quantal scattering theory with excitations

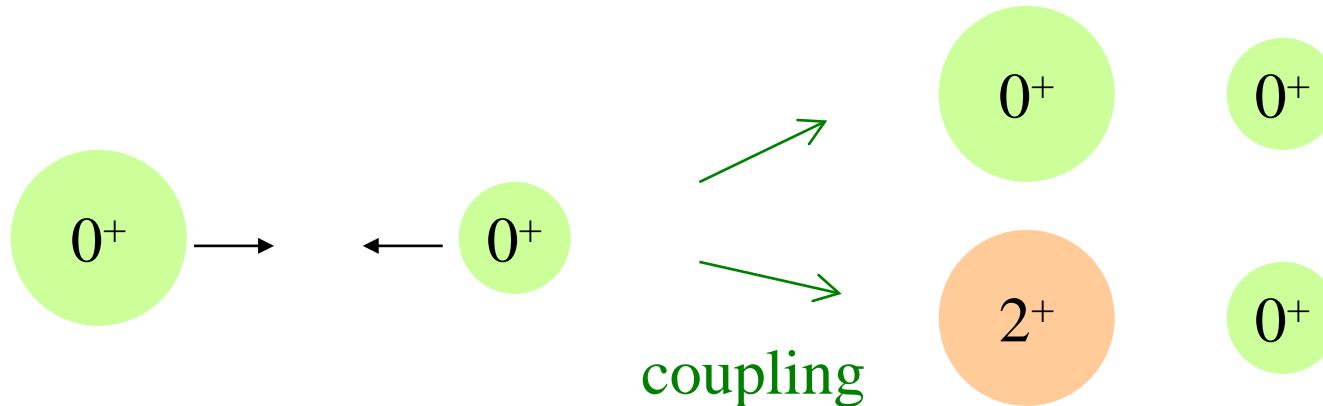
many-body problem



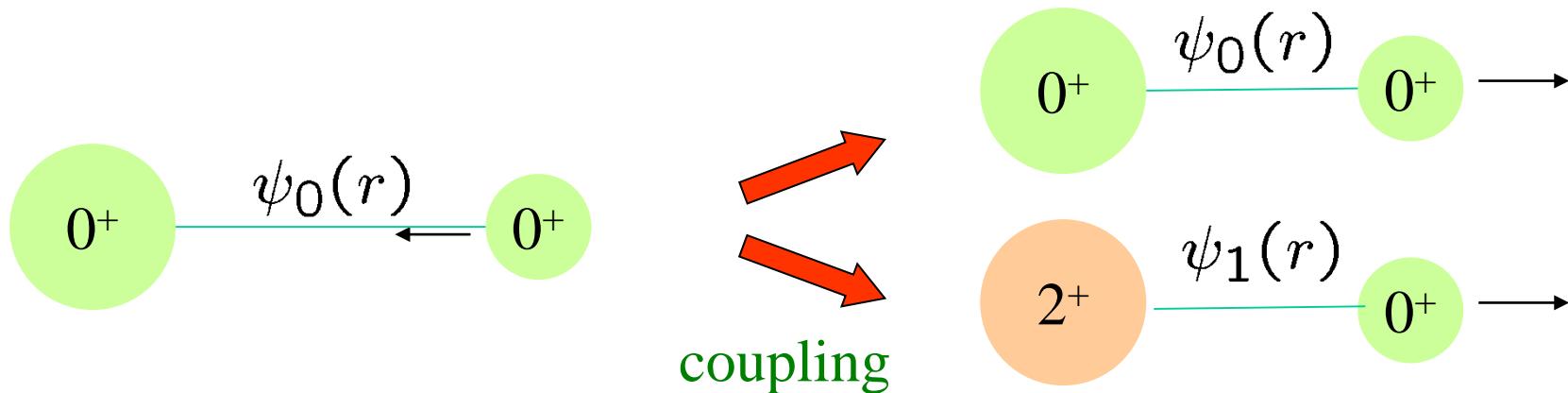
still very challenging



two-body problem, but with excitations
(coupled-channels approach)



Coupled-channels method: a quantal scattering theory with excitations



$$\left[-\frac{\hbar^2}{2\mu} \nabla^2 + \stackrel{\leftrightarrow}{V}(r) - \stackrel{\leftrightarrow}{E} \right] \vec{\psi}(r) = 0$$

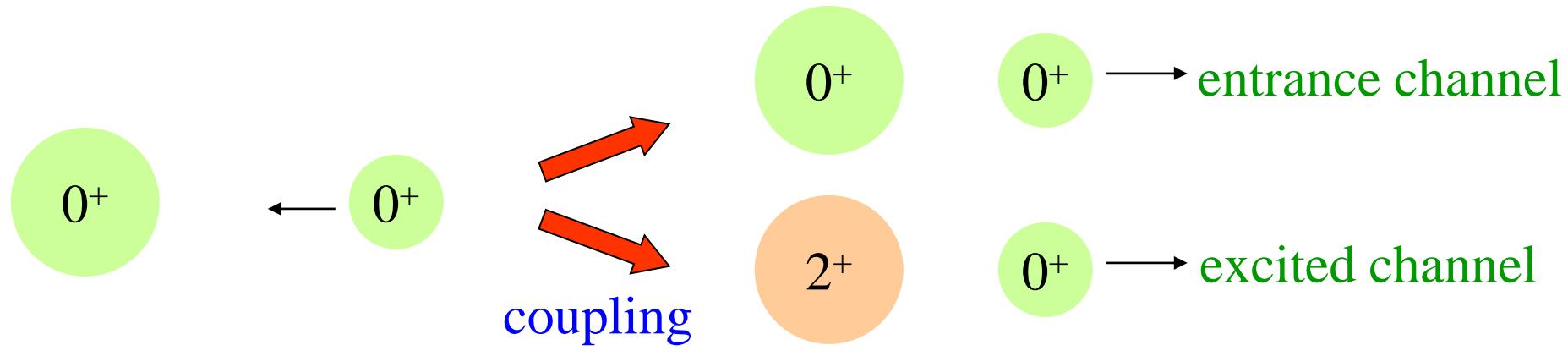
if written down more explicitly:

$$\left[-\frac{\hbar^2}{2\mu} \nabla^2 + V_0(r) + \epsilon_k - E \right] \psi_k(r) + \sum_{k'} \langle \phi_k | V_{\text{coup}} | \phi_{k'} \rangle \psi_{k'}(r) = 0$$

excitation energy

excitation operator

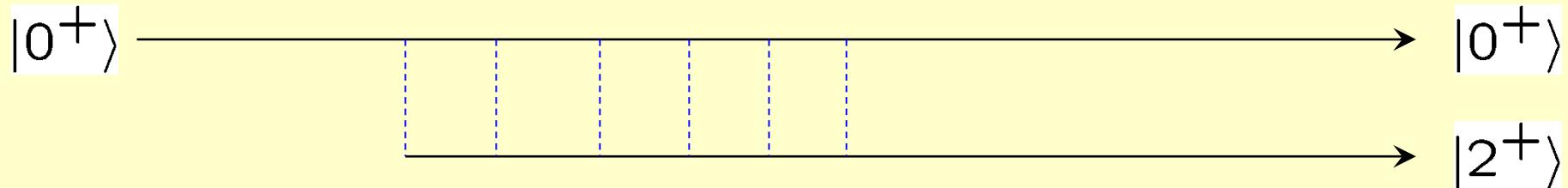
Coupled-channels method: a quantal scattering theory with excitations



$$\left[-\frac{\hbar^2}{2\mu} \nabla^2 + V_0(r) + \epsilon_k - E \right] \psi_k(\mathbf{r}) + \sum_{k'} \langle \phi_k | V_{\text{coup}} | \phi_{k'} \rangle \psi_{k'}(\mathbf{r}) = 0$$

excitation energy

excitation operator



full order treatment of excitation/de-excitation dynamics during reaction

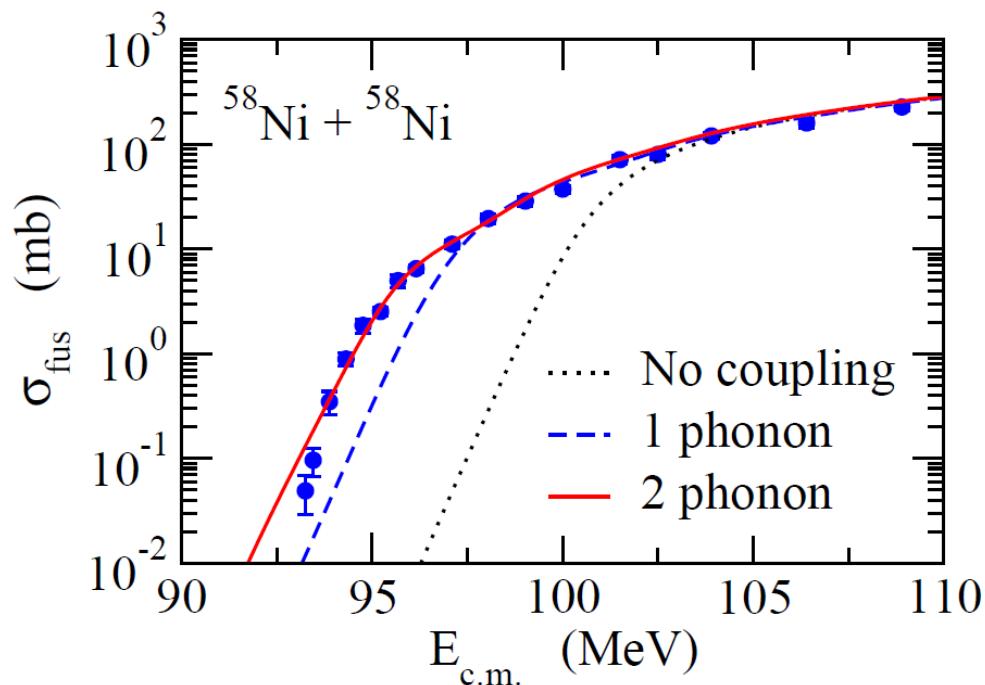
Inputs for C.C. calculations

i) Inter-nuclear potential

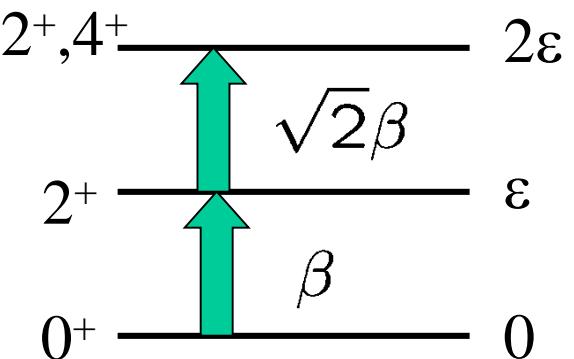
a fit to experimental data at above barrier energies

ii) Intrinsic degrees of freedom

in most of cases, (macroscopic) collective model
(rigid rotor / harmonic oscillator)



simple harmonic oscillator



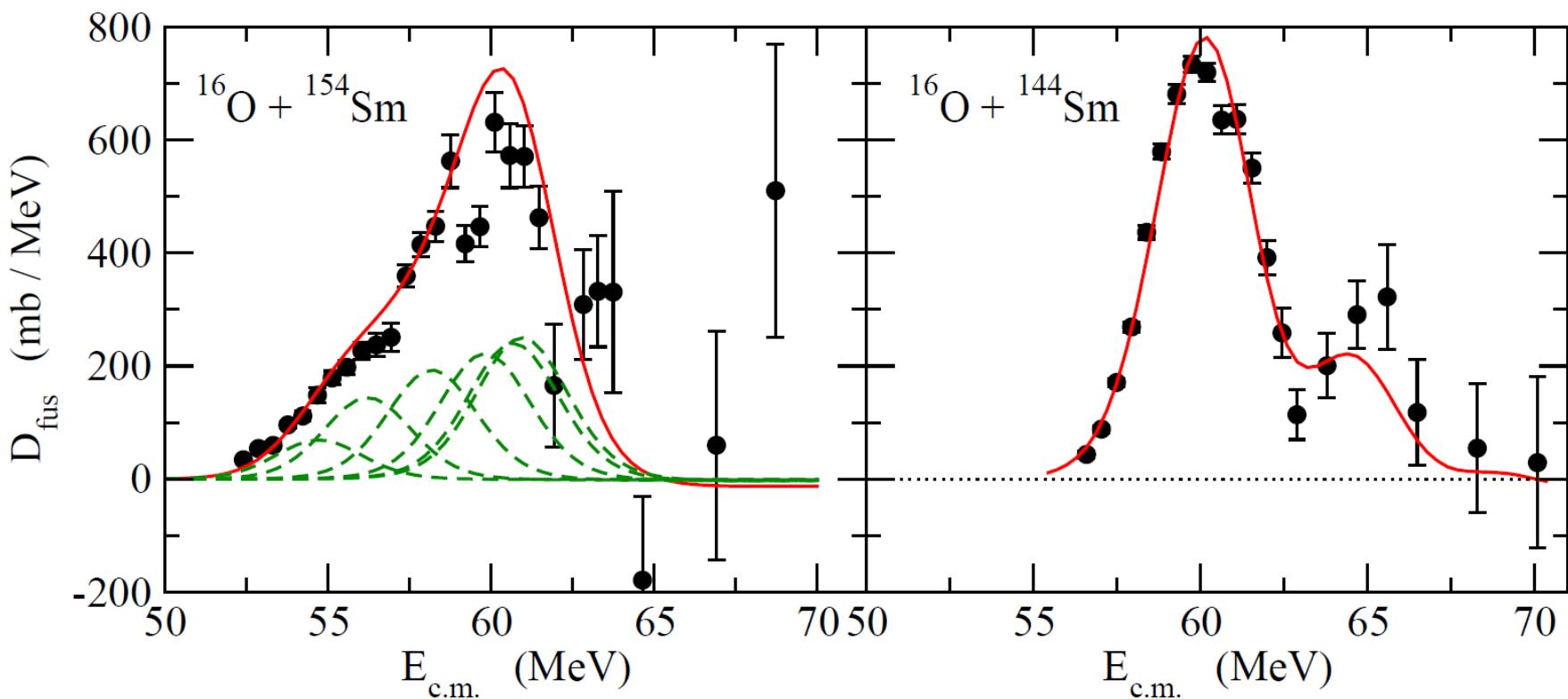
C.C. approach: a standard tool for sub-barrier fusion reactions

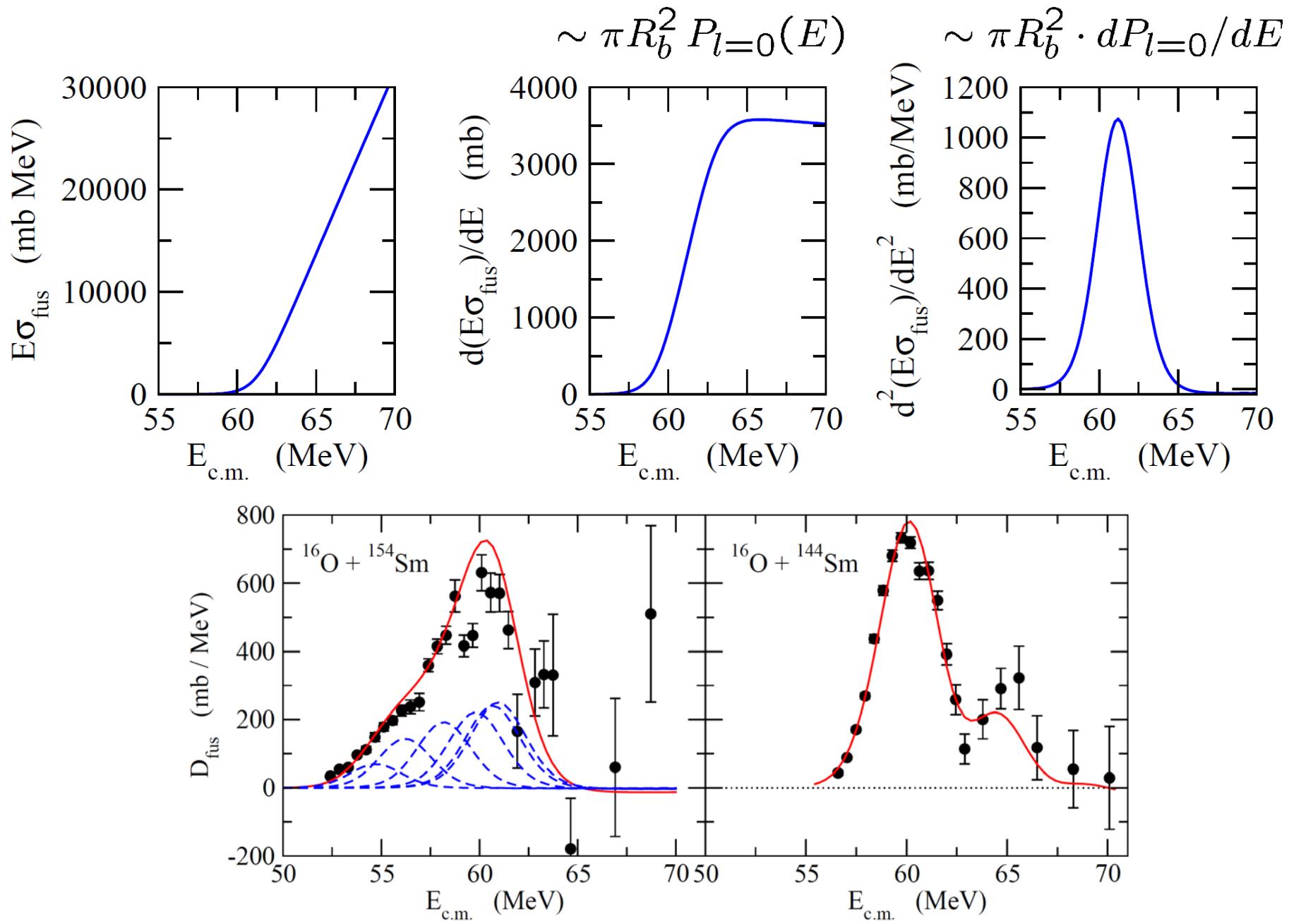
cf. CCFULL (K.H., N. Rowley, A.T. Kruppa, CPC123 ('99) 143)

✓ Fusion barrier distribution (Rowley, Satchler, Stelson, PLB254('91))

$$D_{\text{fus}}(E) = \frac{d^2(E\sigma_{\text{fus}})}{dE^2}$$

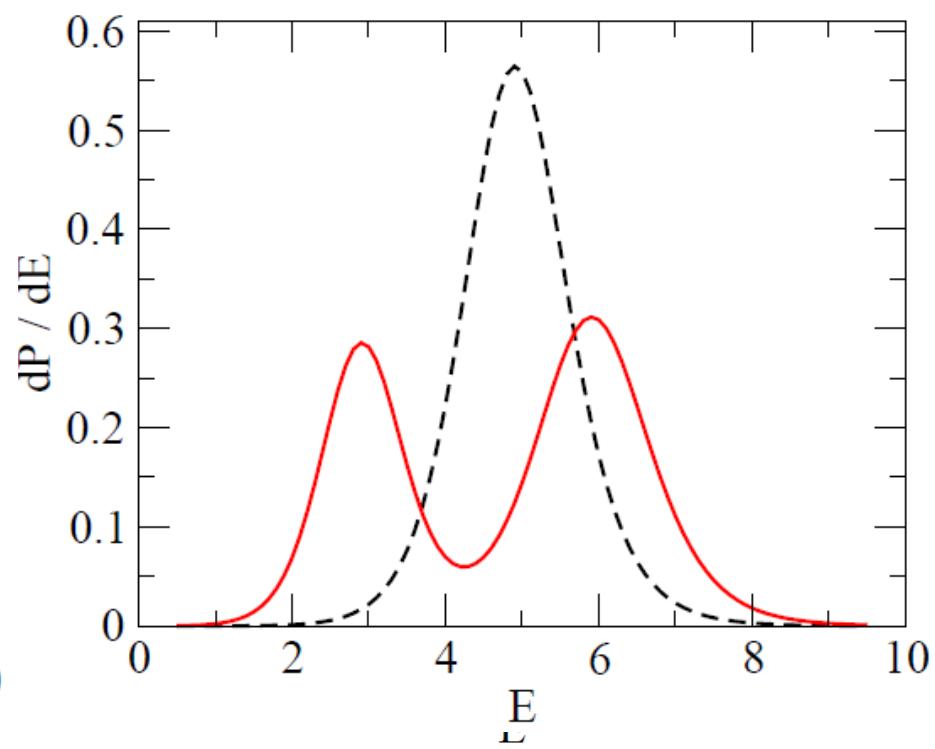
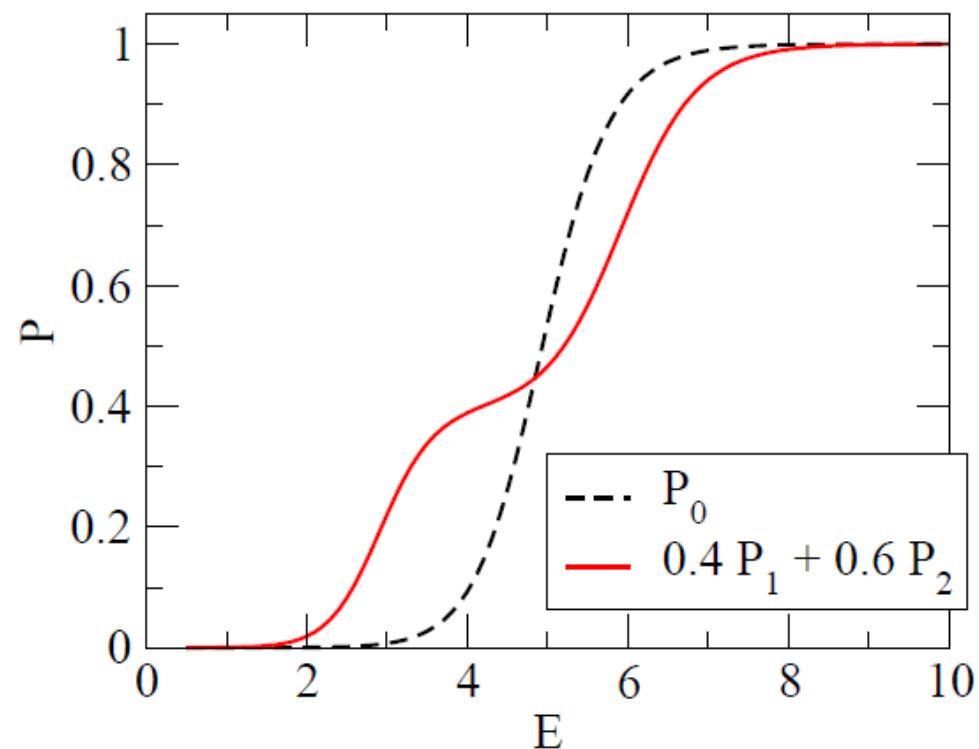
— c.c. calculations





barrier distribution: a problem of two potential barriers

$$P(E) = P(E; V_0) \rightarrow w_1 P(E; V_1) + w_2 P(E; V_2)$$



Fusion barrier distribution

$$D_{\text{fus}}(E) = \frac{d^2(E\sigma_{\text{fus}})}{dE^2}$$

- ◆ N. Rowley, G.R. Satchler, and P.H. Stelson, PLB254 ('91) 25
- ◆ J.X. Wei, J.R. Leigh et al., PRL67 ('91) 3368
- ◆ M. Dasgupta et al., Annu. Rev. Nucl. Part. Sci. 48 ('98) 401
- ◆ A.M. Stefanini et al., Phys. Rev. Lett. 74 ('95) 864

