集団励起の微視的理論

原子核の励起状態

- ✓ 一粒子励起(一つの核子が励起に関与)
- ✓ 集団励起(多くの核子が集団として励起に関与)

集団励起を微視的に理解してみる
(集団励起をミクロに見てみるとどうなっているのか?)

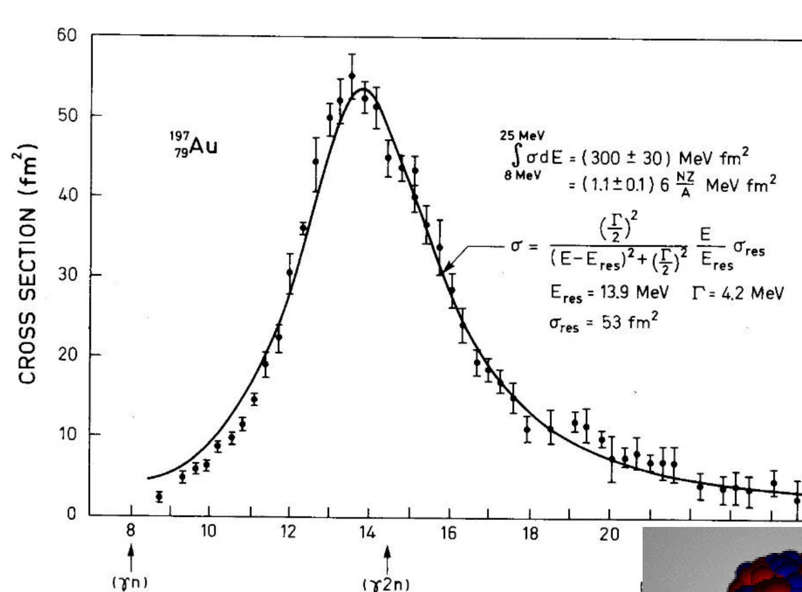
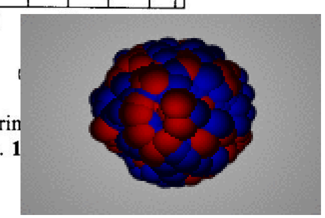
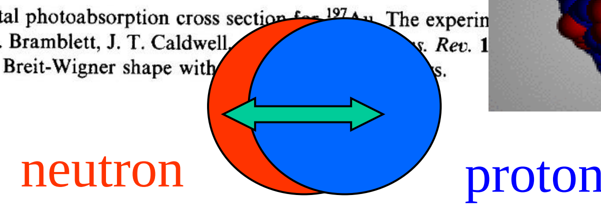


Figure 6-18 Total photoabsorption cross section ... The experim...
S. C. Fultz, R. L. Bramblett, J. T. Caldwell ... *Rev.* 1...
solid curve is of Breit-Wigner shape with ...





集団励起の例:巨大双極子共鳴

一般に:　$$|\Psi_k\rangle = \sum_i C_i|\Phi(SD)_i\rangle$$

タム・ダンコフ近似

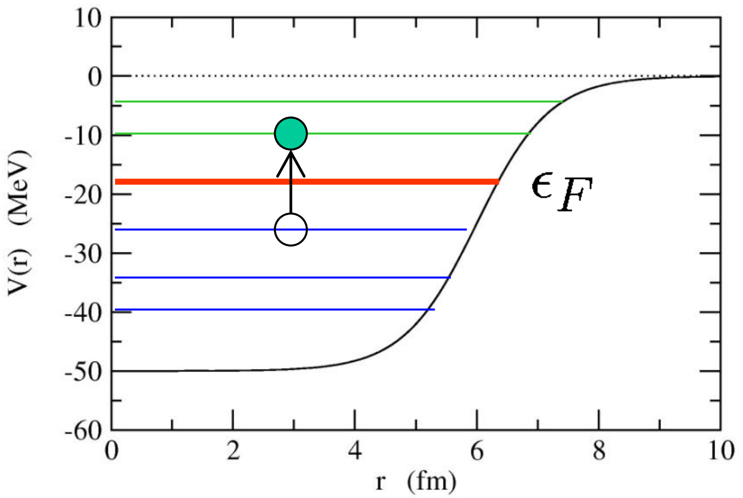
基底状態:　$|HF\rangle$

励起状態:　$$|\nu\rangle = Q_\nu^\dagger|HF\rangle \;=\; \sum_{ph} X_{ph}\, a_p^\dagger a_h|HF\rangle$$

$$\equiv \sum_{ph} X_{ph}|ph^{-1}\rangle$$

1 particle-1 hole (1p1h) state

（1p1h 状態の重ね合わせ）

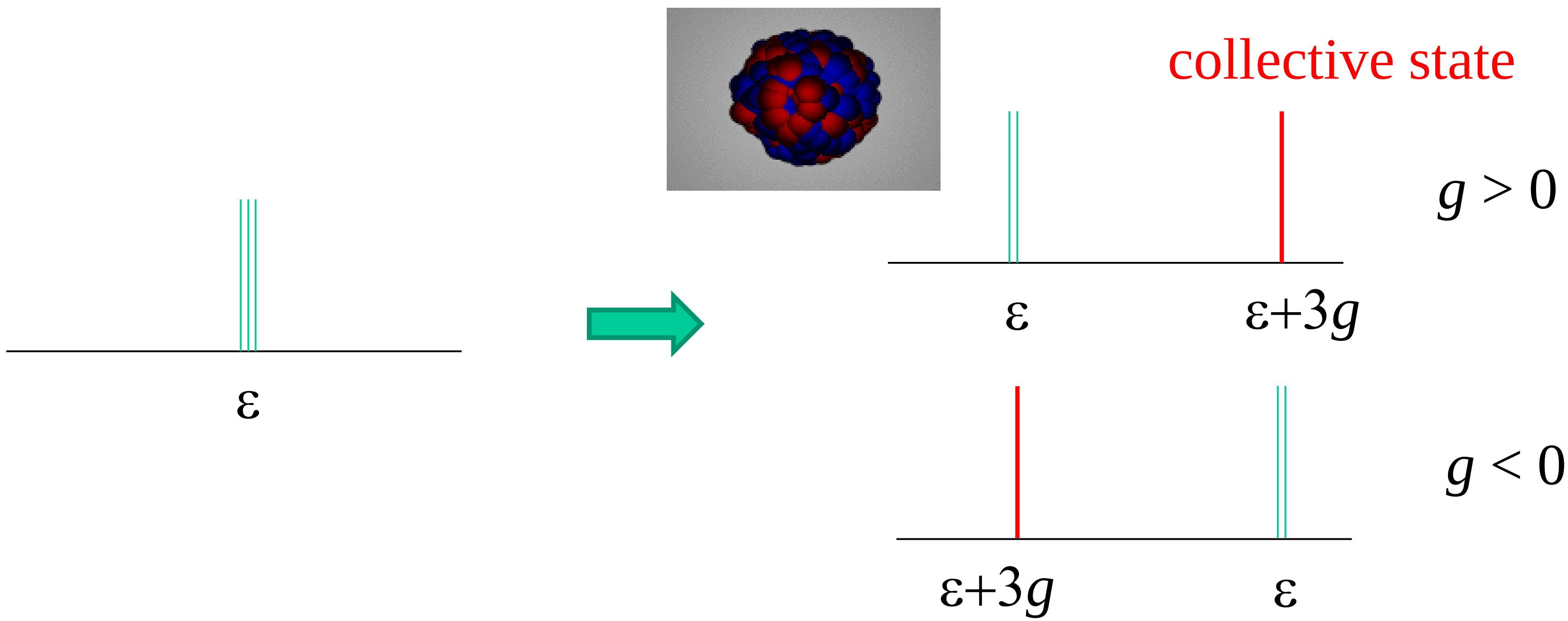


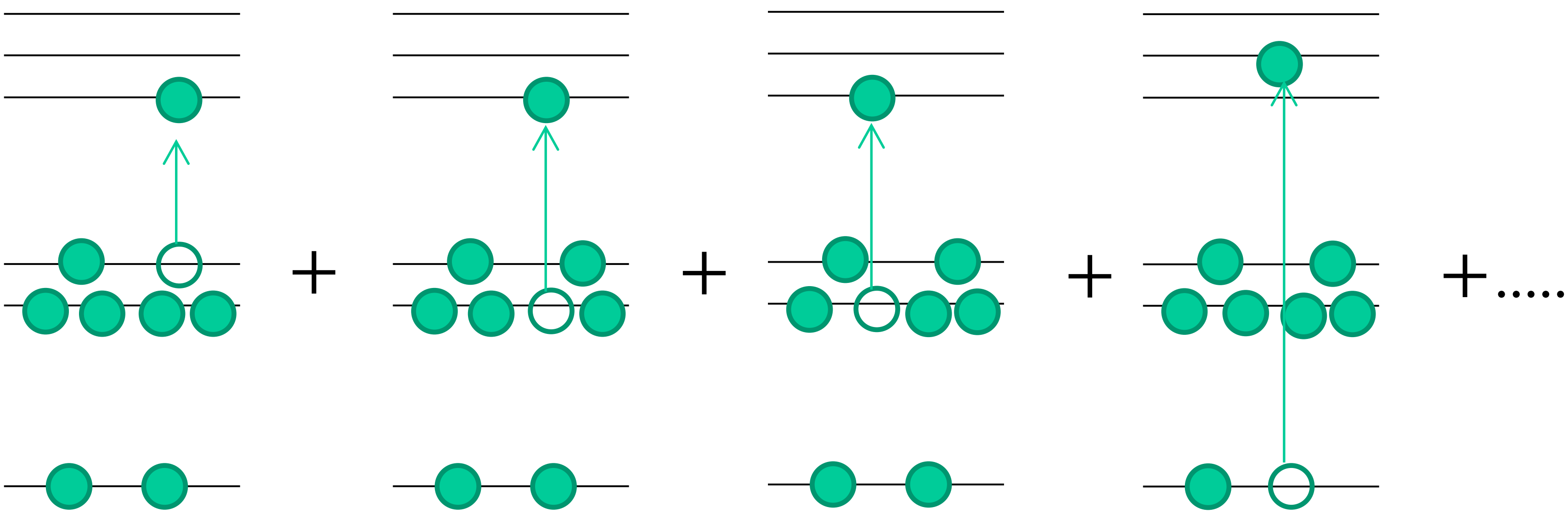
$$a_p^\dagger a_h|HF\rangle$$

TDA on a schematic model

$$H = \epsilon \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix} + g \begin{pmatrix} 1 & 1 & 1 \\ 1 & 1 & 1 \\ 1 & 1 & 1 \end{pmatrix}$$

→ Diagonalization: $\lambda = \epsilon, \epsilon, \underline{\epsilon + 3g}$





複数の粒子・空孔状態をコヒーレントに重ね合わせることによって
多数の核子が励起に関与していることを表現する

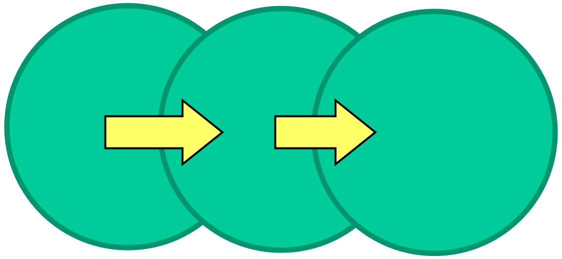
Spurious motion and RPA

Mean-Field Approximation $\longleftrightarrow$ Broken symmetiries

- Center of mass localization (single center)
- Rotational motion

Restoration of broken symmetries

$\Rightarrow$ Zero energy mode (Nambu-Goldstone mode)



does not require an extra energy → zero energy mode

A drawback of TDA:

Zero modes appear at finite excitation energies.

$$|\nu\rangle = Q_\nu^\dagger|0\rangle = \sum_{ph} X_{ph}\, a_p^\dagger a_h|0\rangle \qquad \text{(TDA)}$$

A better approximation:

the random phase approximation (RPA)

$$\boxed{|\nu\rangle = Q_\nu^\dagger|0\rangle = \sum_{ph}\left(X_{ph}\, a_p^\dagger a_h - Y_{ph}\, a_h^\dagger a_p\right)|0\rangle}$$

(superposition of 1p1h states)

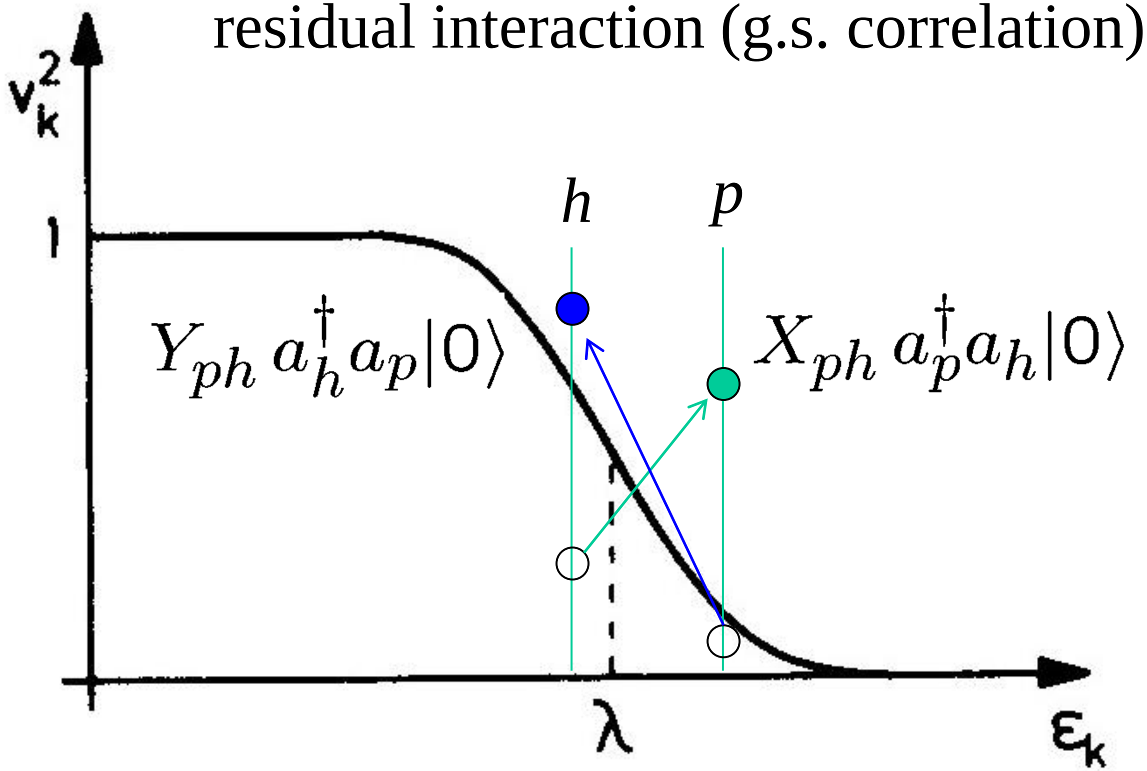
基底状態:　$Q_\nu|0\rangle = 0$　で定義。

A better approximation: **the random phase approximation (RPA)**

$$\boxed{|\nu\rangle = Q_\nu^\dagger|0\rangle = \sum_{ph}\left(X_{ph}\,a_p^\dagger a_h - Y_{ph}\,a_h^\dagger a_p\right)|0\rangle}$$

(superposition of 1p1h states)

smearing of Fermi surface due to the residual interaction (g.s. correlation)



A better approximation: **the random phase approximation (RPA)**

$$\boxed{|\nu\rangle = Q_\nu^\dagger|0\rangle = \sum_{ph}\left(X_{ph}\, a_p^\dagger a_h - Y_{ph}\, a_h^\dagger a_p\right)|0\rangle}$$

(superposition of 1p1h states)

$$[H, Q_\nu^\dagger] \sim E_\nu Q_\nu^\dagger$$

$$\Downarrow$$

$$\langle HF|[\delta Q, [H, Q_\nu^\dagger]]|HF\rangle = E_\nu \langle HF|[\delta Q, Q_\nu^\dagger]|HF\rangle$$

→ coupled equations for *X* and *Y*

$$\langle HF|[\delta Q,[H,Q_\nu^\dagger]]|HF\rangle = E_\nu \langle HF|[\delta Q,Q_\nu^\dagger]|HF\rangle$$

$$Q_\nu^\dagger = \sum_{ph} X_{ph}\, a_p^\dagger a_h - Y_{ph}\, a_h^\dagger a_p \qquad \delta Q = a_h^\dagger a_p, \qquad a_p^\dagger a_h$$

RPA equation:

$$\begin{aligned} \sum_{p'h'} A_{ph,p'h'} X_{p'h'} + B_{ph,p'h'} Y_{p'h'} &= E_\nu X_{ph} \\ \sum_{p'h'} B^*_{ph,p'h'} X_{p'h'} + A^*_{ph,p'h'} Y_{p'h'} &= -E_\nu Y_{ph} \end{aligned}$$

$$\begin{aligned} A_{ph,p'h'} &= (\epsilon_p - \epsilon_h)\delta_{ph,p'h'} + \langle ph'|\bar{v}|hp'\rangle \\ B_{ph,p'h'} &= \langle pp'|\bar{v}|hh'\rangle \end{aligned}$$

or

$$\begin{pmatrix} A & B \\ -B^* & -A^* \end{pmatrix} \begin{pmatrix} X \\ Y \end{pmatrix} = E_\nu \begin{pmatrix} X \\ Y \end{pmatrix}$$

Spurious motion in RPA

Mean-Field Approximation $\Longleftrightarrow$ Broken symmetiries

- Center of mass localization (single center)
- Rotational motion

Restoration of broken symmetries

$\Rightarrow$ Zero mode (Nambu-Goldstone mode)

$$[H, \hat{O}] = 0$$

RPA

$$[H, Q_\nu^\dagger] \sim E_\nu Q_\nu^\dagger$$

$\hat{O}$ is a solution of RPA with E=0

$$Q^\dagger = \hat{O} = \sum_{ph} (O_{ph} a_p^\dagger a_h + O_{hp} a_h^\dagger a_p)$$

(note) $Q_{\mathrm{TDA}}^\dagger = \sum_{ph} O_{ph} a_p^\dagger a_h$ $\Rightarrow$ $[H, Q_{\mathrm{TDA}}^\dagger] \neq 0$

Spurious motion in RPA

Mean-Field Approximation $\Longleftrightarrow$ Broken symmetiries

- Center of mass localization (single center)
- Rotational motion

Restoration of broken symmetries

$\Rightarrow$ Zero mode (Nambu-Goldstone mode)

RPA

$$[H, Q_\nu^\dagger] \sim E_\nu Q_\nu^\dagger$$

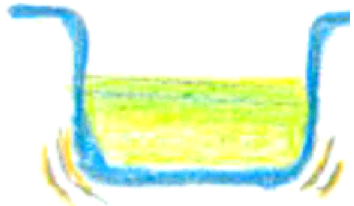
if $[H, \hat{O}] = 0$

Then $\hat{O}$ is a solution of RPA with E=0

The physical solutions are completely separated out from the spurious modes.

他のRPAの定式化

- 線形応答理論

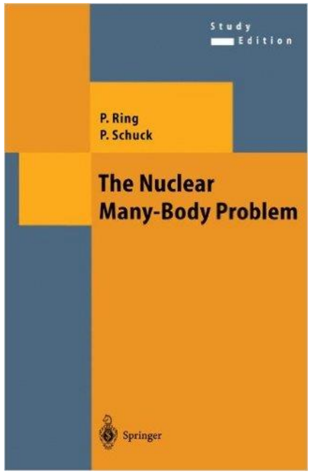


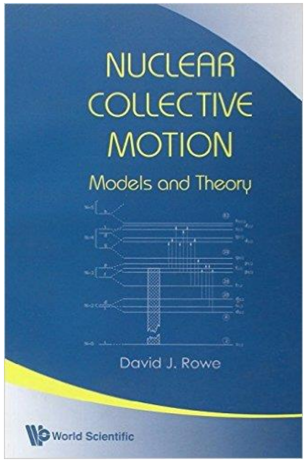
外場で原子核を揺すった時に、
原子核がどのように応答するか摂動論
を使って議論する
→固有モードを見つける

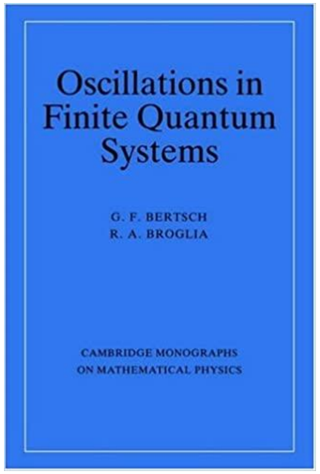
- 時間に依存するハートリー・フォック(TDHF)方程式を線形化

$$i\hbar\dot{\rho}(t) = [h[\rho], \rho] \qquad \longleftarrow \qquad \rho(t) = \rho_0 + \delta\rho(t)$$

詳しくは：







Comparison between Skyrme-(Q)RPA calculation and exp. data

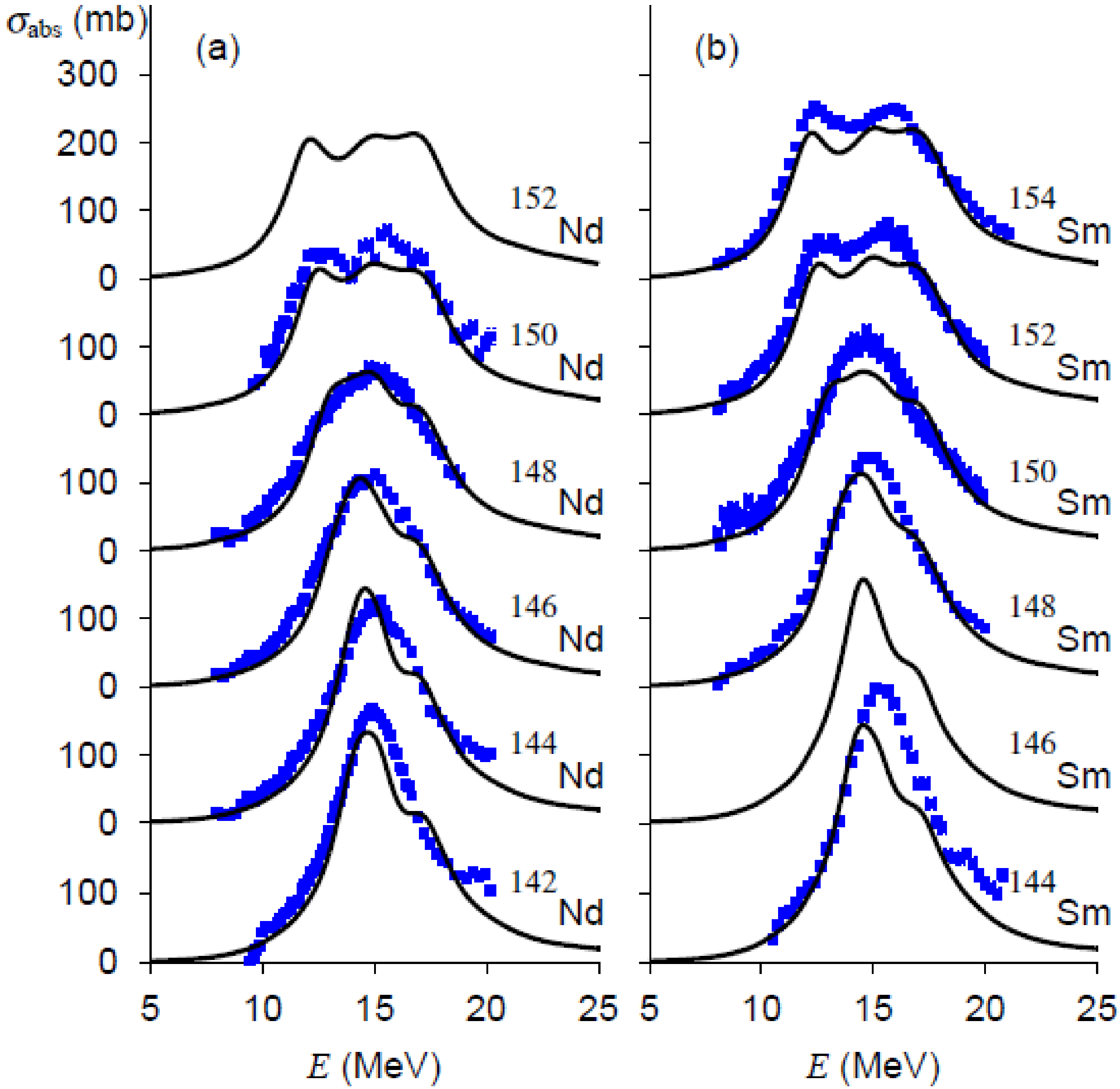
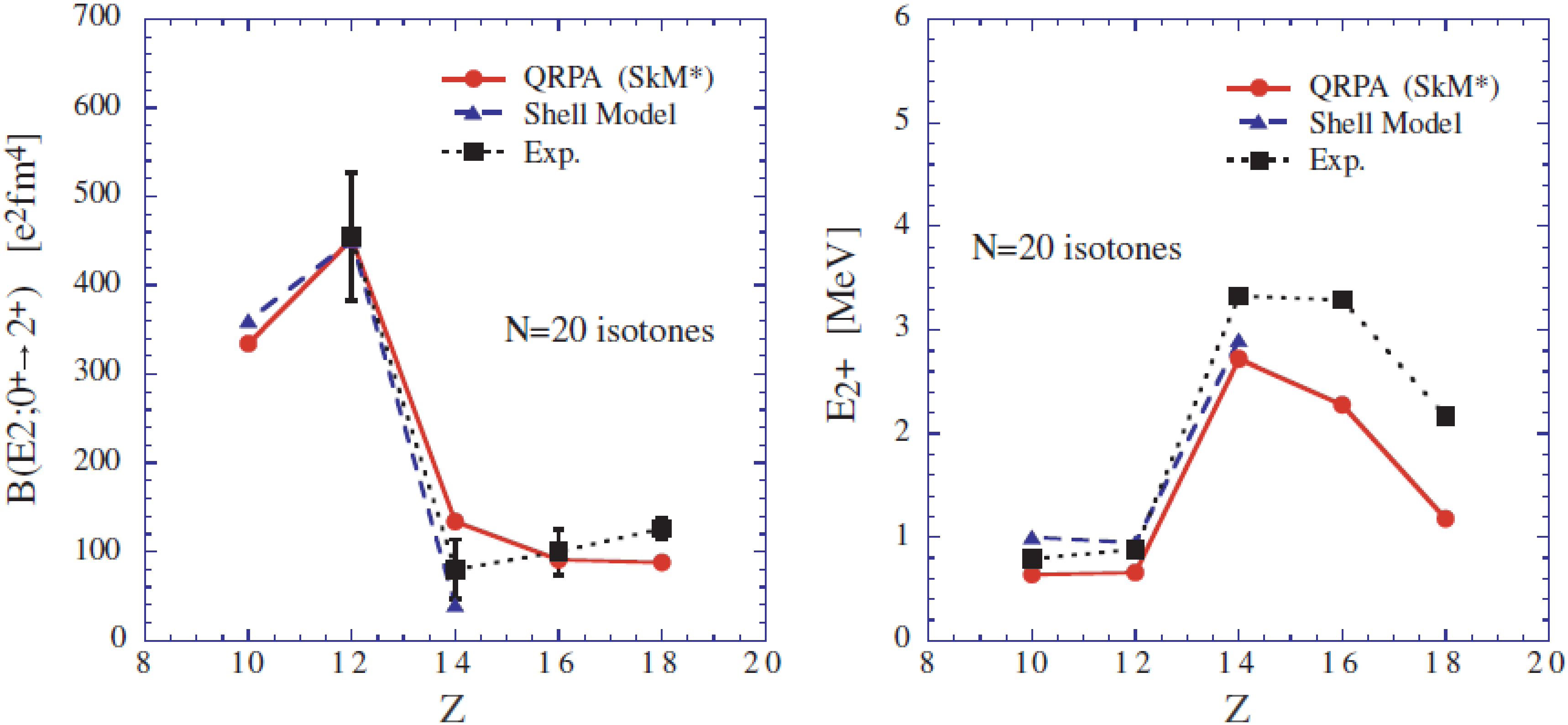


photo-absorption
cross section
(GDR)

K. Yoshida
and T. Nakatsukasa,
PRC83('11)021304



M. Yamagami and Nguyen Van Giai, PRC69 ('04) 034301

RPA on a schematic model

$$\boxed{\begin{pmatrix} A & B \\ -B^* & -A^* \end{pmatrix}\begin{pmatrix} X \\ Y \end{pmatrix} = E\begin{pmatrix} X \\ Y \end{pmatrix}}$$

$$\begin{aligned} A_{ph,p'h'} &= (\epsilon_p - \epsilon_h)\delta_{ph,p'h'} + \langle ph'|\bar{v}|hp'\rangle \\ B_{ph,p'h'} &= \langle pp'|\bar{v}|hh'\rangle \end{aligned}$$

Separable interaction:

$$\begin{aligned} \langle ph'|\bar{v}|hp'\rangle &= \lambda D_{ph} D^*_{p'h'} \\ \langle pp'|\bar{v}|hh'\rangle &= \lambda D_{ph} D_{p'h'} \end{aligned}$$

$$\boxed{\frac{1}{\lambda} = \sum_{ph} \frac{|D_{ph}|^2}{E - \epsilon_{ph}} - \frac{|D_{ph}|^2}{E + \epsilon_{ph}}}$$

(RPA dispersion relation)

Cf. TDA dispersion relation: $\frac{1}{\lambda} = \sum_{ph} \frac{|D_{ph}|^2}{E - \epsilon_{ph}}$

cf. タム・ダンコフ近似の場合:

Separable interaction: $\langle ph'|\bar{v}|hp'\rangle = \lambda D_{ph} D^*_{p'h'}$

Tamm-Dancoff equation: $\sum_{p'h'} A_{ph,p'h'} X_{p'h'} = E X_{ph}$

$$A_{ph,p'h'} = (\epsilon_p - \epsilon_h)\delta_{ph,p'h'} + \langle ph'|\bar{v}|hp'\rangle$$

$$(\epsilon_{ph} - E) X_{ph} + \lambda D_{ph} \cdot T = 0 \qquad T \equiv \sum_{ph} D^*_{ph} X_{ph}$$

$$X_{ph} = -\lambda \frac{D_{ph} T}{\epsilon_{ph} - E}$$

$$T = \sum_{ph} D^*_{ph} X_{ph} = -\lambda \sum_{ph} \frac{|D_{ph}|^2}{\epsilon_{ph} - E} T$$

or $$\boxed{\frac{1}{\lambda} = \sum_{ph} \frac{|D_{ph}|^2}{E - \epsilon_{ph}}}$$

(TDA dispersion relation)

RPA on a schematic model

$$\langle ph'|\bar{v}|hp'\rangle = \lambda D_{ph} D^*_{p'h'}$$
$$\langle pp'|\bar{v}|hh'\rangle = \lambda D_{ph} D_{p'h'}$$

Separable interaction:

$$\boxed{\frac{1}{\lambda} = \sum_{ph} \frac{|D_{ph}|^2}{E - \epsilon_{ph}} - \frac{|D_{ph}|^2}{E + \epsilon_{ph}}}$$

(RPA dispersion relation)

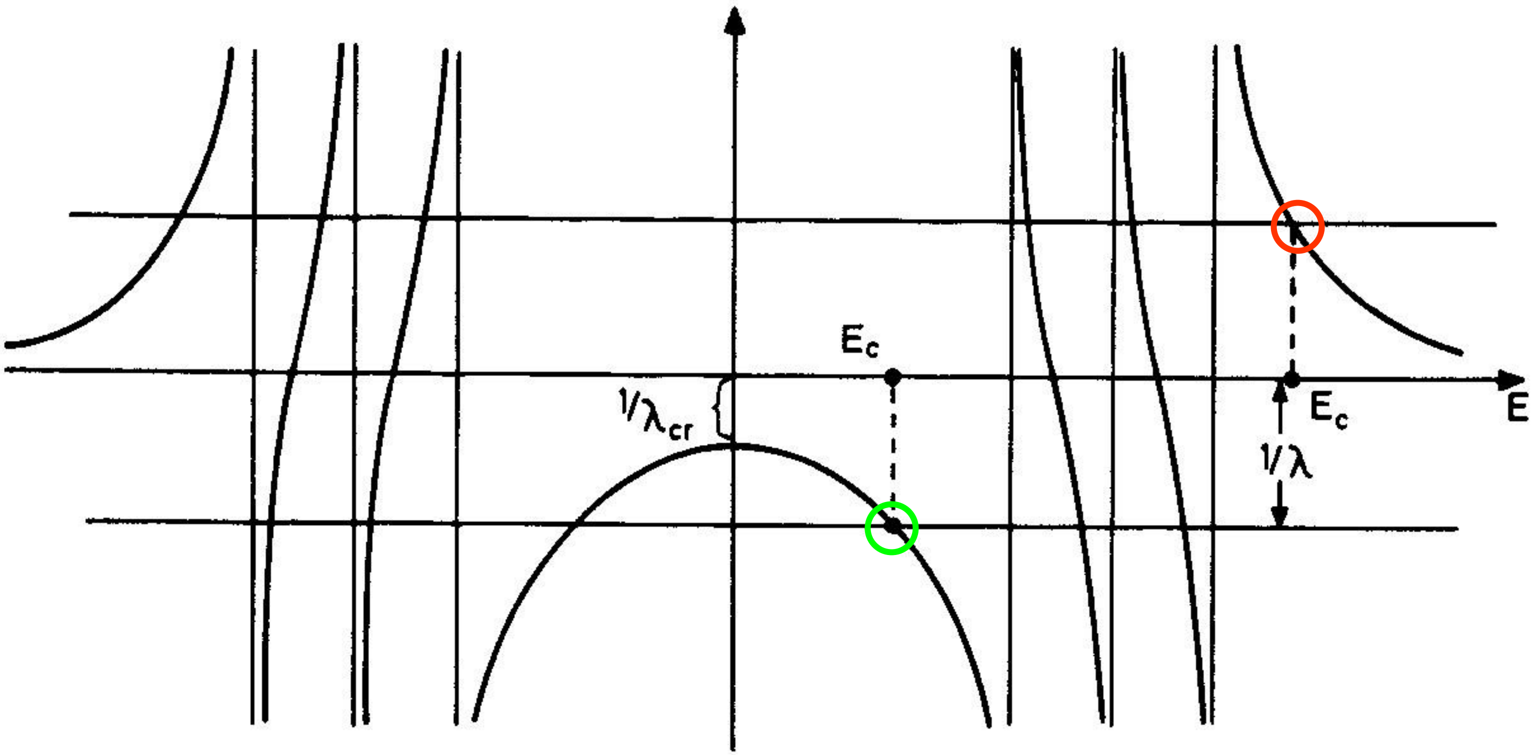


Figure 8.11. Graphical solution of the dispersion relation (8.135).

RPA on a schematic model

$$\langle ph'|\bar{v}|hp'\rangle = \lambda D_{ph} D^*_{p'h'}$$
$$\langle pp'|\bar{v}|hh'\rangle = \lambda D_{ph} D_{p'h'}$$

Separable interaction:

$$\boxed{\frac{1}{\lambda} = \sum_{ph} \frac{|D_{ph}|^2}{E - \epsilon_{ph}} - \frac{|D_{ph}|^2}{E + \epsilon_{ph}}}$$

(RPA dispersion relation)

$\epsilon_{ph} = \epsilon$ 　のとき、

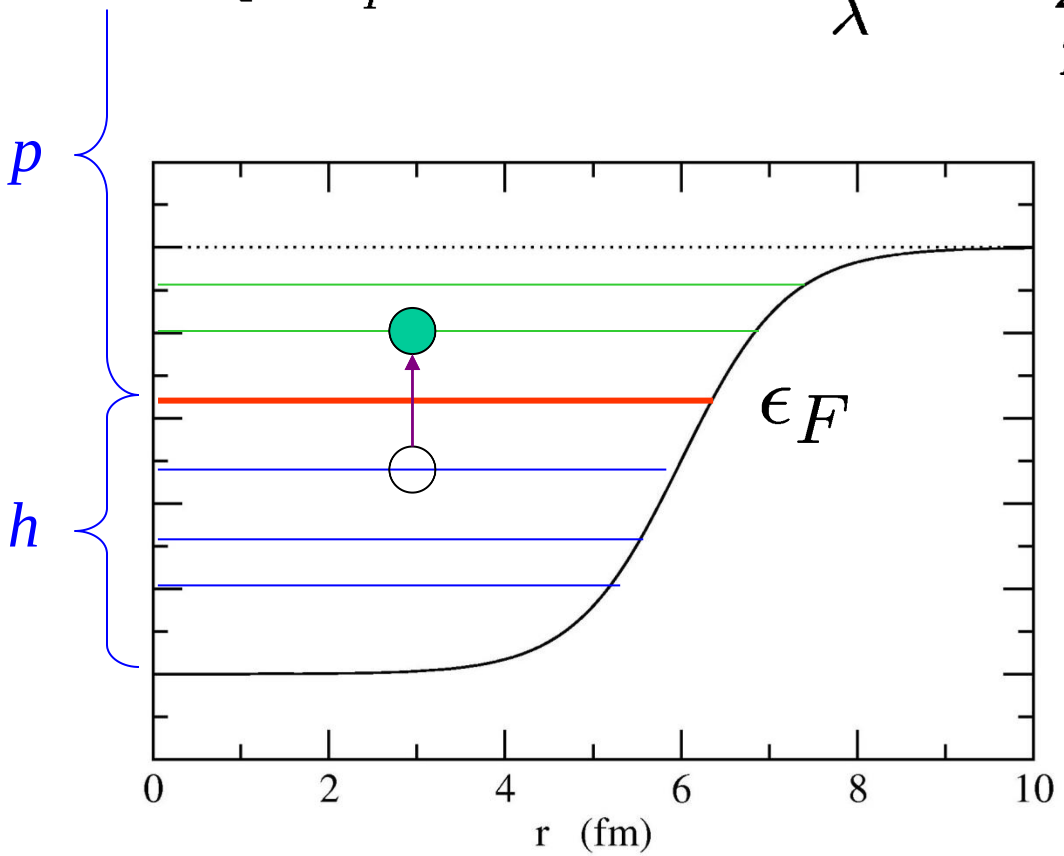
$$\frac{1}{\lambda} = \sum_{ph} |D_{ph}|^2 \left(\frac{1}{E - \epsilon} - \frac{1}{E + \epsilon} \right) = \sum_{ph} |D_{ph}|^2 \frac{2\epsilon}{E^2 - \epsilon^2}$$

$$\longrightarrow \qquad E^2 = \epsilon^2 + 2\epsilon\lambda \sum_{ph} |D_{ph}|^2$$

λ が負（引力）だと、どこかで $E^2 < 0$ となる

Continuum Excitations

$$\langle ph'|\bar{v}|hp'\rangle = \lambda D_{ph} D^*_{p'h'} \quad \Longrightarrow \quad \frac{1}{\lambda} = -\sum_{ph} \frac{|D_{ph}|^2}{\epsilon_p - \epsilon_h - E}$$



h: all the occupied (bound) states

p: the bound excited states + continuum states

$$\frac{1}{\lambda} = -\sum_{ph} \frac{|D_{ph}|^2}{\epsilon_p - \epsilon_h - E} = -\sum_{ph} \langle \phi_h | D^\dagger | \phi_p \rangle \frac{1}{\epsilon_p - \epsilon_h - E} \langle \phi_p | D | \phi \rangle$$

(note) $\hat{h}\phi_p = \epsilon_p \phi_p$

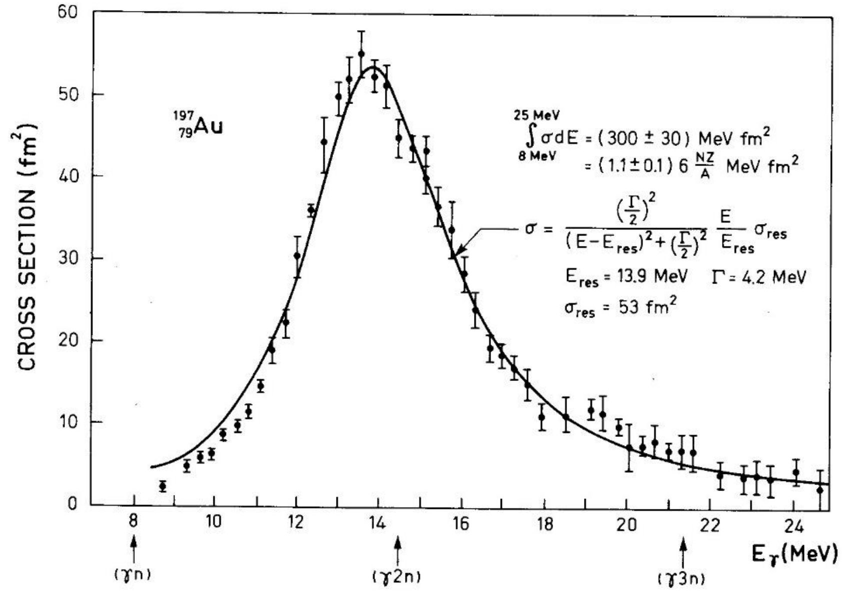
$$\frac{1}{\lambda} = -\sum_{ph} \langle \phi_h | D^\dagger \frac{1}{\hat{h} - \epsilon_h - E} | \phi_p \rangle \langle \phi_p | D | \phi \rangle$$

$$1 = \sum_i |\phi_i\rangle\langle\phi_i| = \sum_p |\phi_p\rangle\langle\phi_p| + \sum_h |\phi_h\rangle\langle\phi_h|$$

$$\frac{1}{\lambda} = -\sum_{h} \langle \phi_h | D^\dagger \frac{1}{\hat{h} - \epsilon_h - E} \left[1 - \sum_{h'} |\phi_{h'}\rangle\langle\phi_{h'}| \right] D | \phi \rangle$$

particle 状態の和がなくなった→連続状態もすべて自動的に入る

巨大共鳴の幅



i) 連続状態との結合（粒子放出）

escape width $\Gamma^{\uparrow}$

continuum RPA

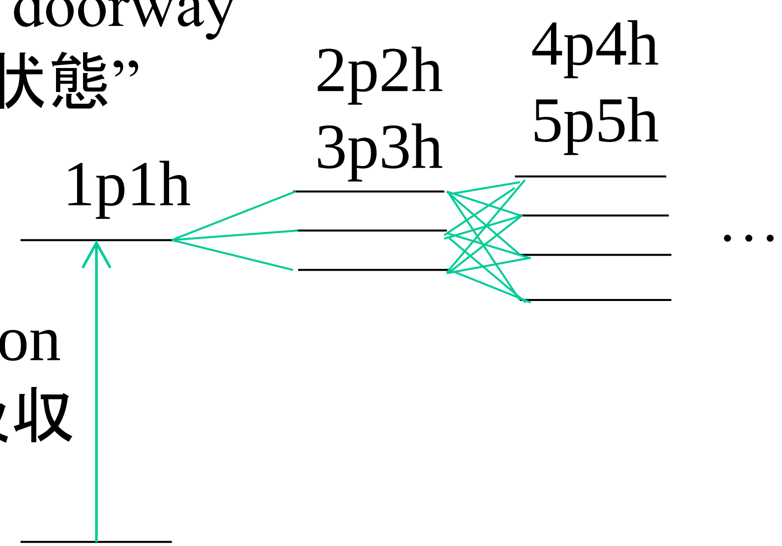
ii) より複雑な状態との結合

spreading width $\Gamma^{\downarrow}$

軽い核を除き
幅の主成分

1体演算子

“doorway
状態”



photon
の吸収

レポート問題４（〆切：12月4日（土））

１．分離型相互作用

$$
\begin{aligned}
\langle ph'|\bar{v}|hp'\rangle &= \lambda D_{ph} D^*_{p'h'} \\
\langle pp'|\bar{v}|hh'\rangle &= \lambda D_{ph} D_{p'h'}
\end{aligned}
$$

の場合にRPA方程式を解き、RPA dispersion relation を導け。

２．RPA のA 行列、B行列が

$$
A = \begin{pmatrix} \epsilon & g \\ g & \epsilon \end{pmatrix}, \qquad B = \begin{pmatrix} 0 & g \\ g & 0 \end{pmatrix}
$$

で与えられているときに RPA方程式を解き、（正の）固有値を求めよ。