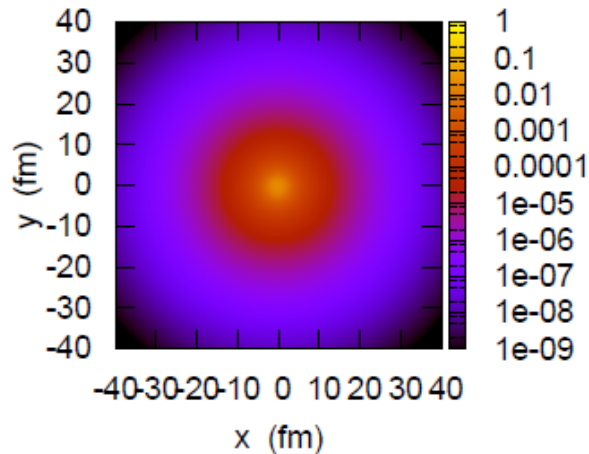


# Pairing correlations and odd-even staggering in reaction cross sections of weakly-bound nuclei



**K. Hagino (Tohoku U.)**

**H. Sagawa (U. of Aizu)**

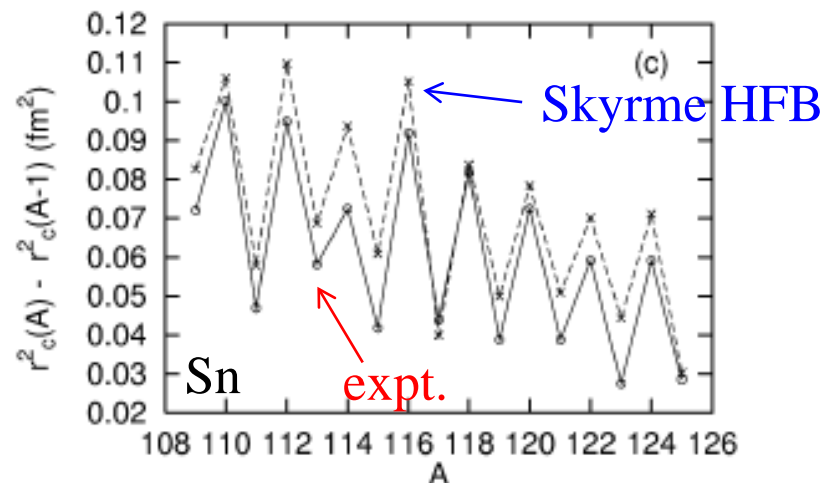
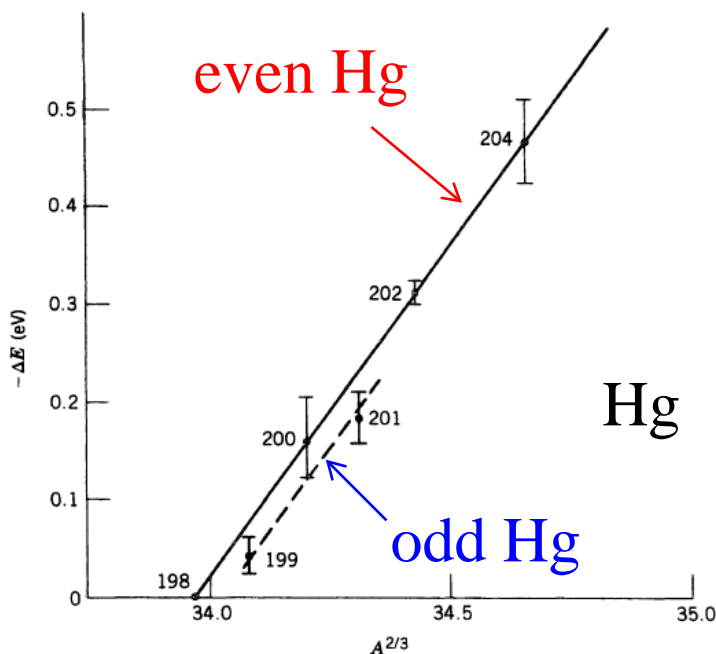
- PRC84('11)011303(R)
- PRC85('12)014303
- arXiv:1202.2725 [nucl-th]



- 1. Introduction: odd-even staggerings in atomic nuclei*
- 2. Odd-even staggering of reaction cross sections ( $\sigma_R$ )*
- 3. Pairing correlation in weakly-bound nuclei and  $\sigma_R$*
- 4. Staggering parameter*
- 5. Summary*

# Introduction: odd-even staggering in atomic nuclei

➤ isotope shifts: smaller charge radius for odd-A nuclei



S. Sakakihara and Y. Tanaka,  
NPA691('01)649

**Figure 3.6** K X-ray isotope shifts in Hg. The energy of the K X ray in Hg is about 100 keV, so the relative isotope shift is of the order of  $10^{-6}$ . The data show the

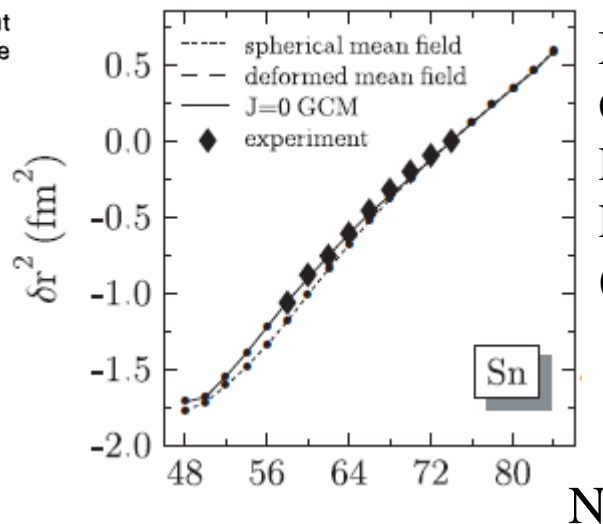
K.S. Krane, "Introductory Nuclear Physics"

$$\Delta E \sim -\frac{2}{5} \frac{Z^4 e^2}{a_0^3} (\langle r^2 \rangle_A - \langle r^2 \rangle_{A'})$$

cf. Bohr-Mottelson, eq. (2.85)

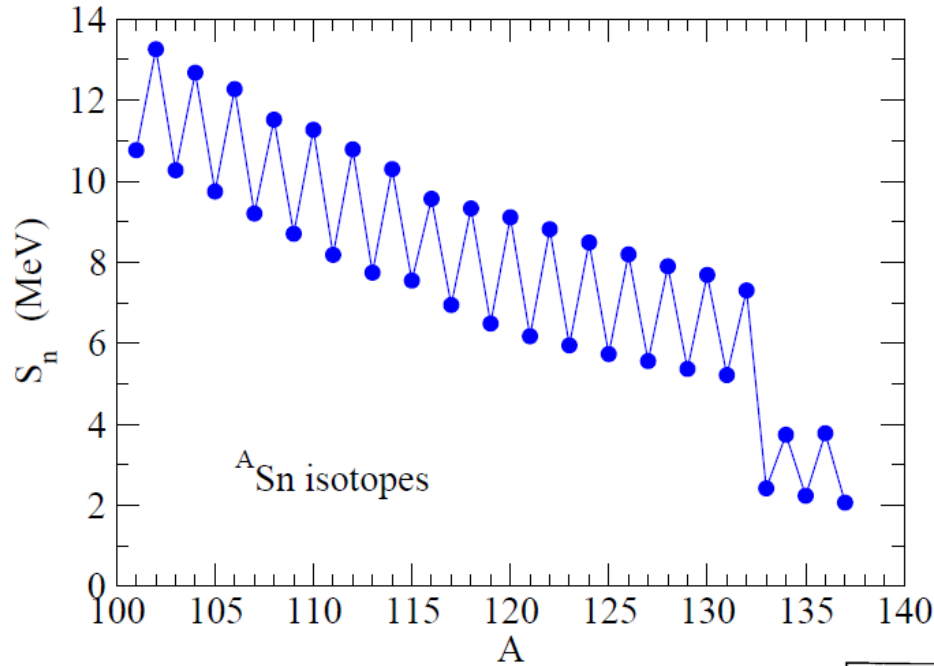
$$\gamma \equiv \frac{\langle r^2 \rangle_{A+1} - \langle r^2 \rangle_A}{\langle r^2 \rangle_{A+2} - \langle r^2 \rangle_A}$$

- deformation effect? - pairing effect?



M. Bender,  
G.F. Bertsch,  
P.-H. Heenen,  
PRC73('06)034322  
(even-even only)

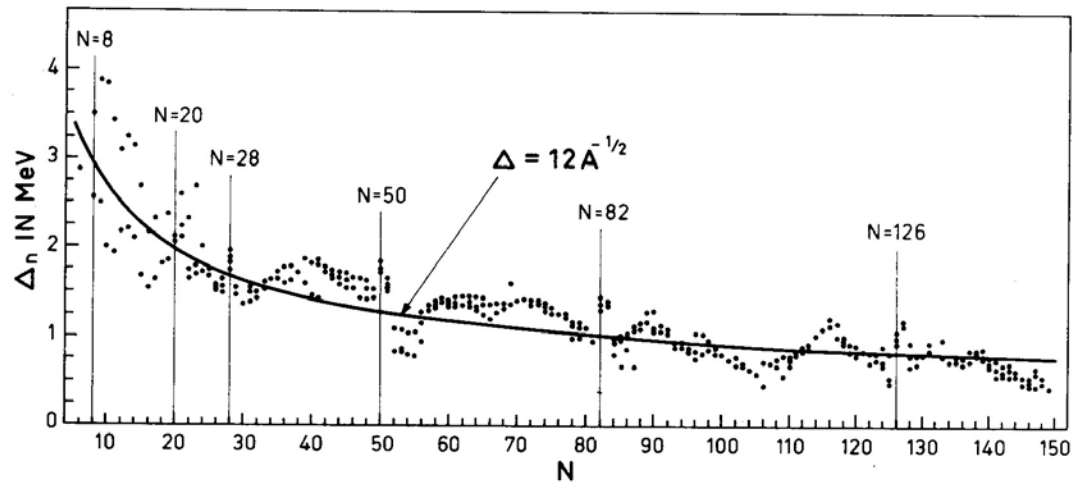
## ► binding energy



$$S_n(N) = B(N) - B(N-1)$$

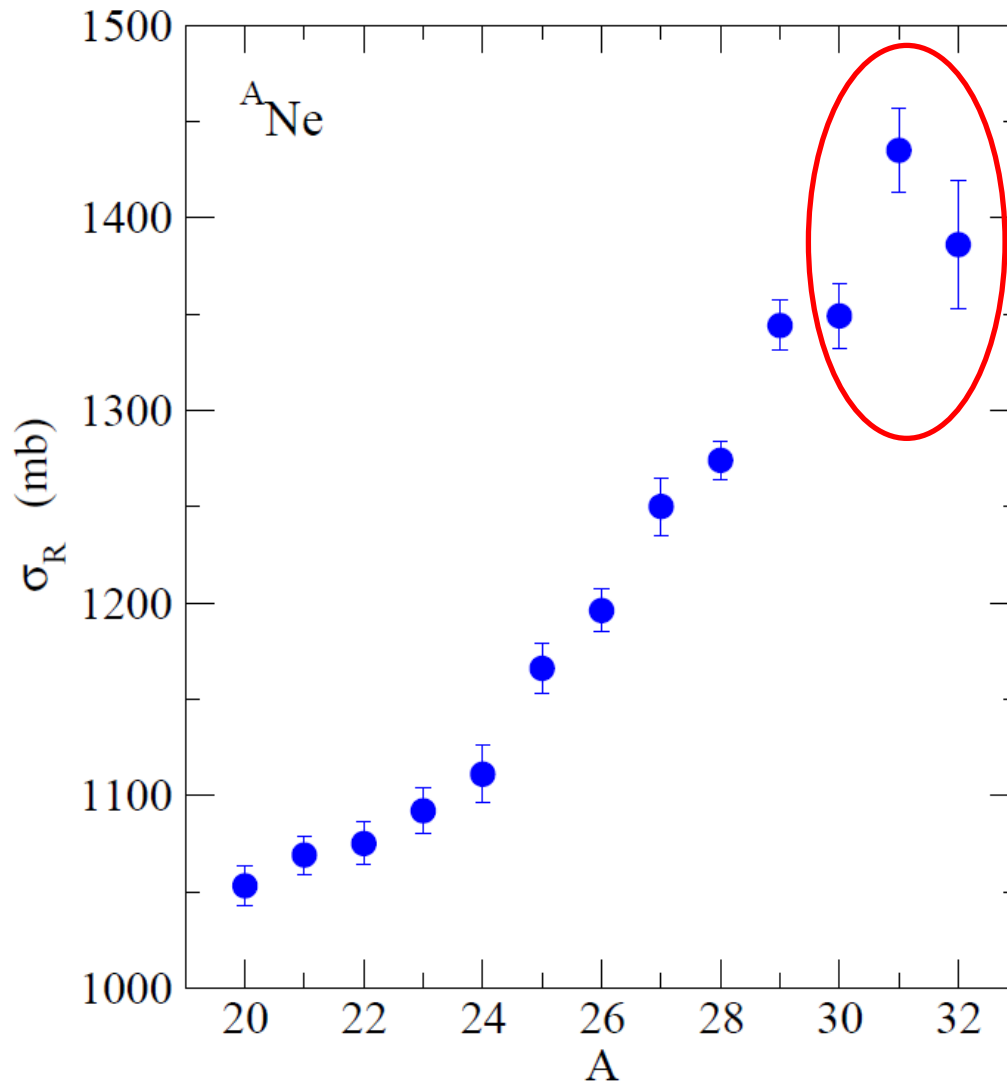
## pairing gap parameter

$$\begin{aligned} \Delta(N) &= \frac{(-)^N}{2} (B(N-1) - 2B(N) \\ &\quad + B(N+1)) \\ &= \frac{(-)^N}{2} (S_n(N-1) - S_n(N)) \end{aligned}$$



# Odd-even staggering of interaction cross sections

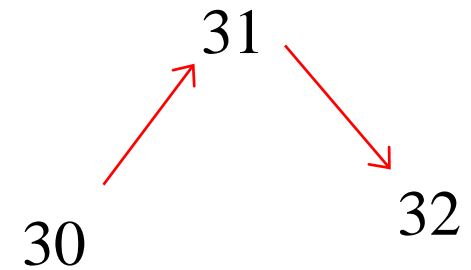
$\sigma_I$  of unstable nuclei: often show a large odd-even staggering



Typical example:

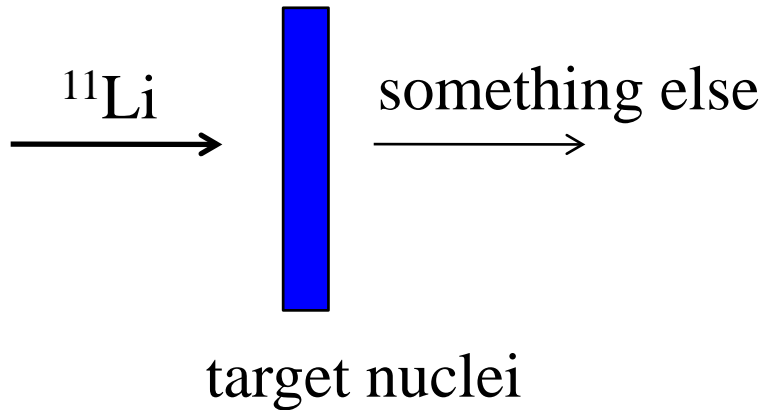
Recent experimental data  
on Ne isotopes

M. Takechi et al.,  
Phys. Lett. B707 ('12) 357

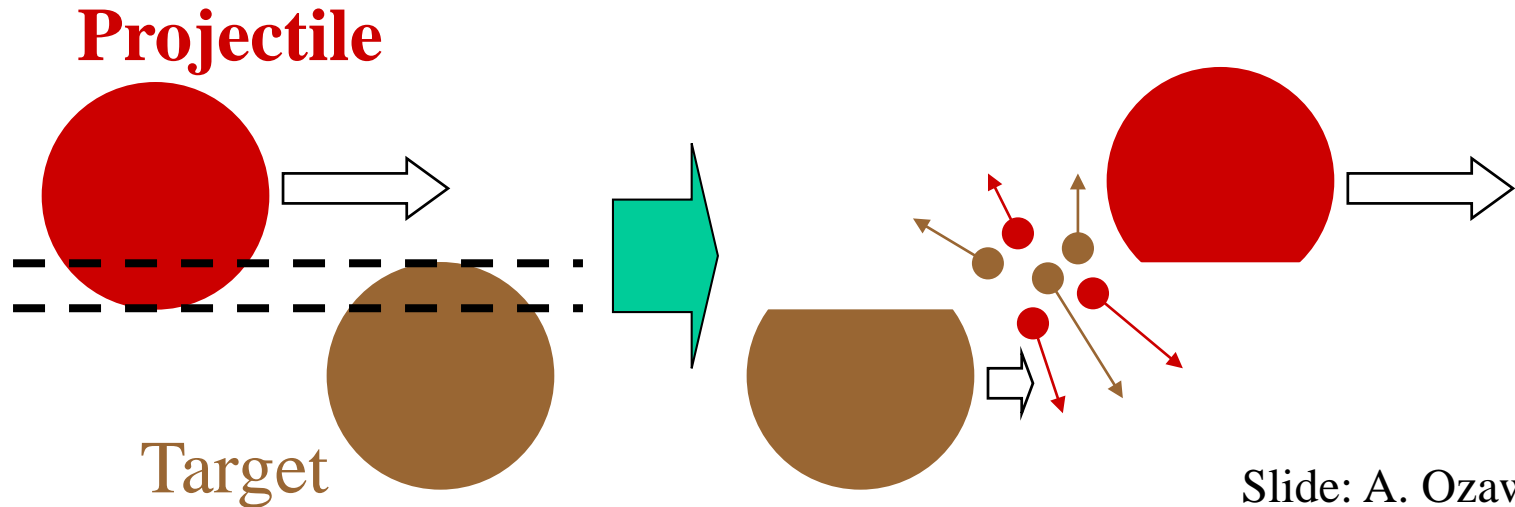


clear odd-even effect

# Introduction: interaction cross section



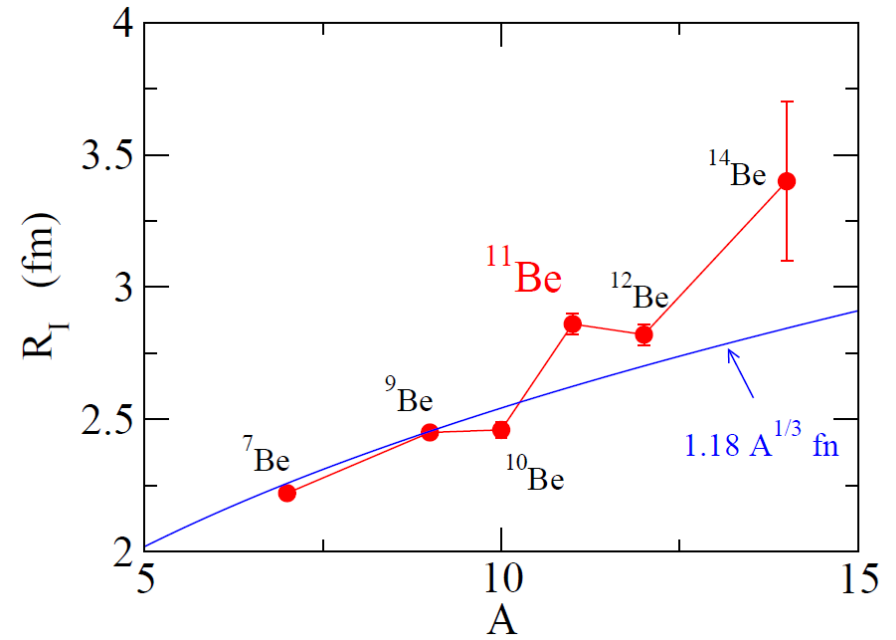
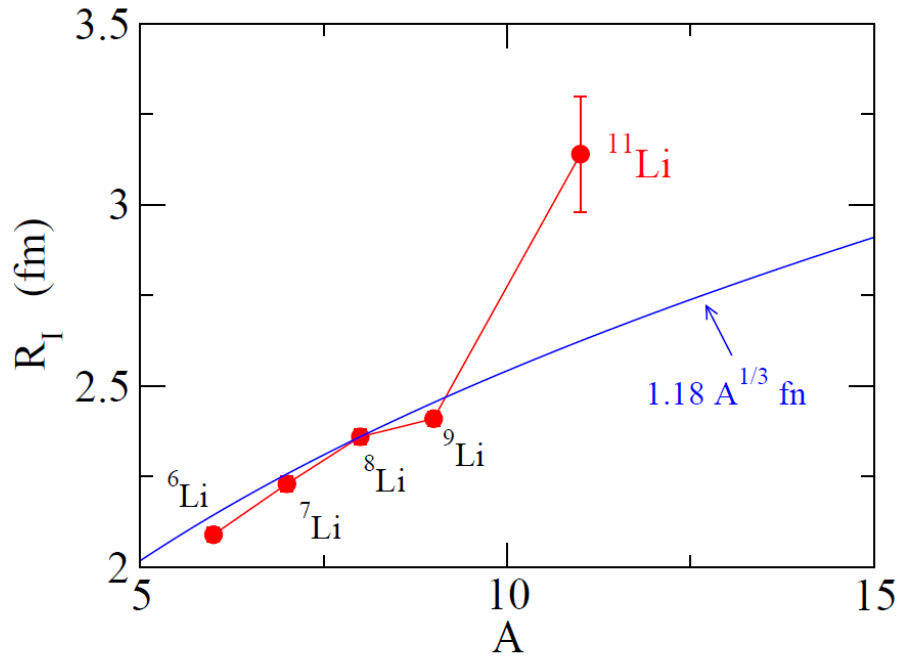
interaction cross section  $\sigma_I$   
= cross section for the change  
of Z a/o N in the incident nucleus



$$\sigma_I \sim \pi [R_I(P) + R_I(T)]^2$$

$$\longrightarrow R_I(P)$$

# Discovery of halo nuclei

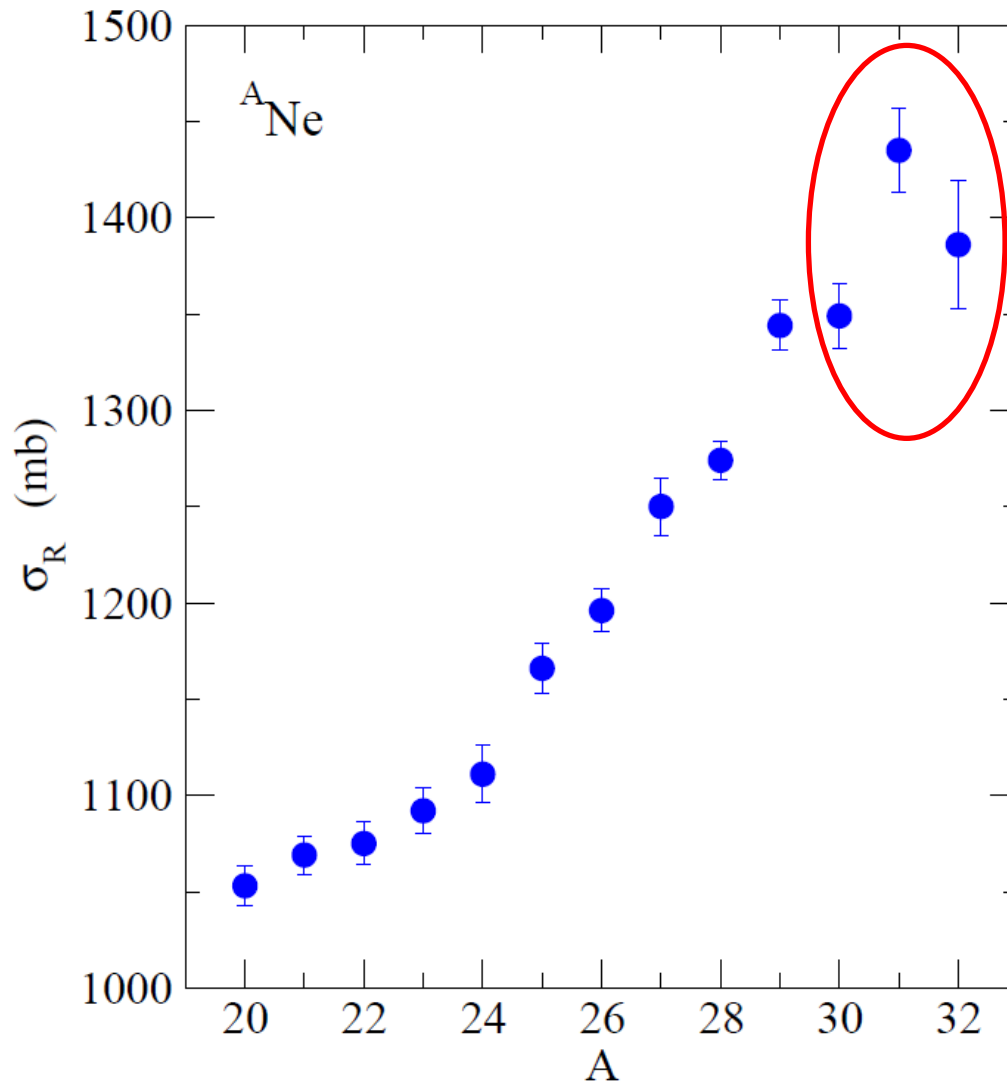


I. Tanihata, T. Kobayashi et al.,  
PRL55('85)2676; PLB206('88)592



# Odd-even staggering of interaction cross sections

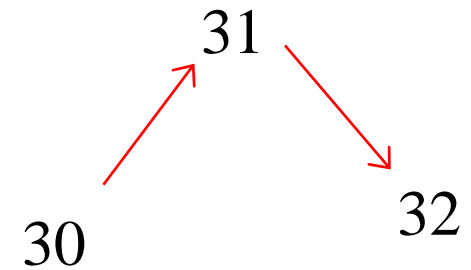
$\sigma_I$  of unstable nuclei: often show a large odd-even staggering



Typical example:

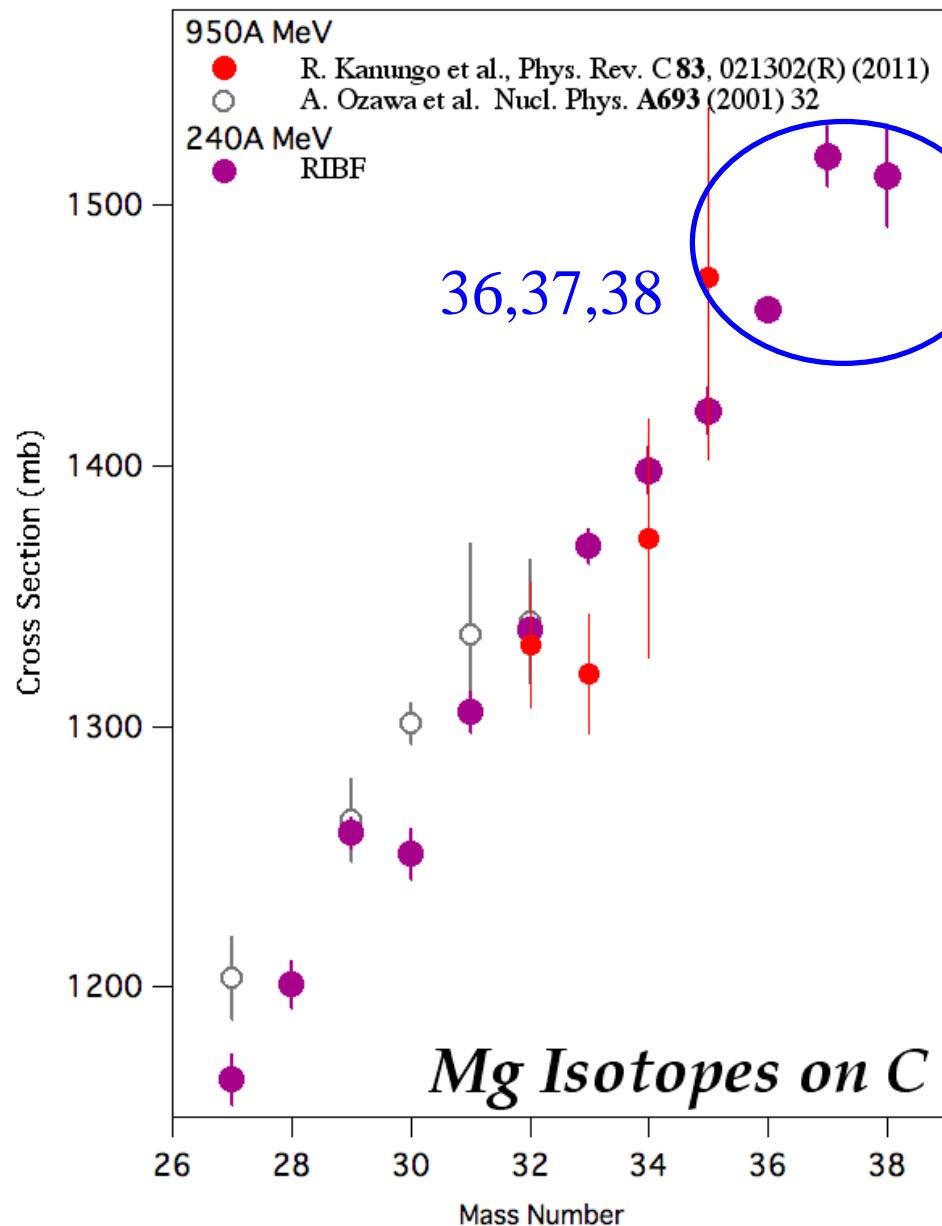
Recent experimental data  
on Ne isotopes

M. Takechi et al.,  
Phys. Lett. B707 ('12) 357

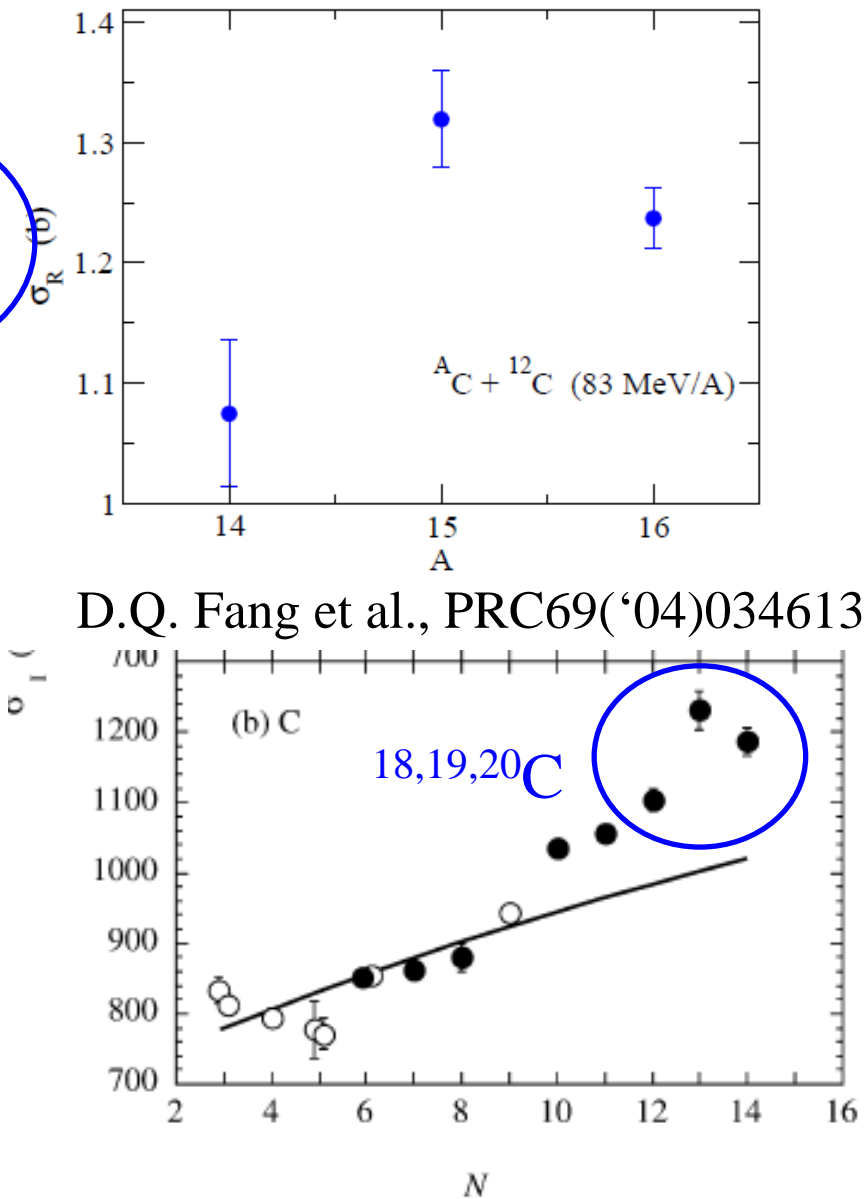


clear odd-even effect

# Other systems



M. Takechi, private communications



D.Q. Fang et al., PRC69('04)034613

A. Ozawa et al., NPA691('01)599



## Our motivation:

Relation between the odd-mass staggering (OES) of  $\sigma_R$  and pairing (anti-halo) effect?

➤ pairing anti-halo effect

K. Bennaceur, J. Dobaczewski,  
and M. Ploszajczak,  
PLB496('00)154

pairing

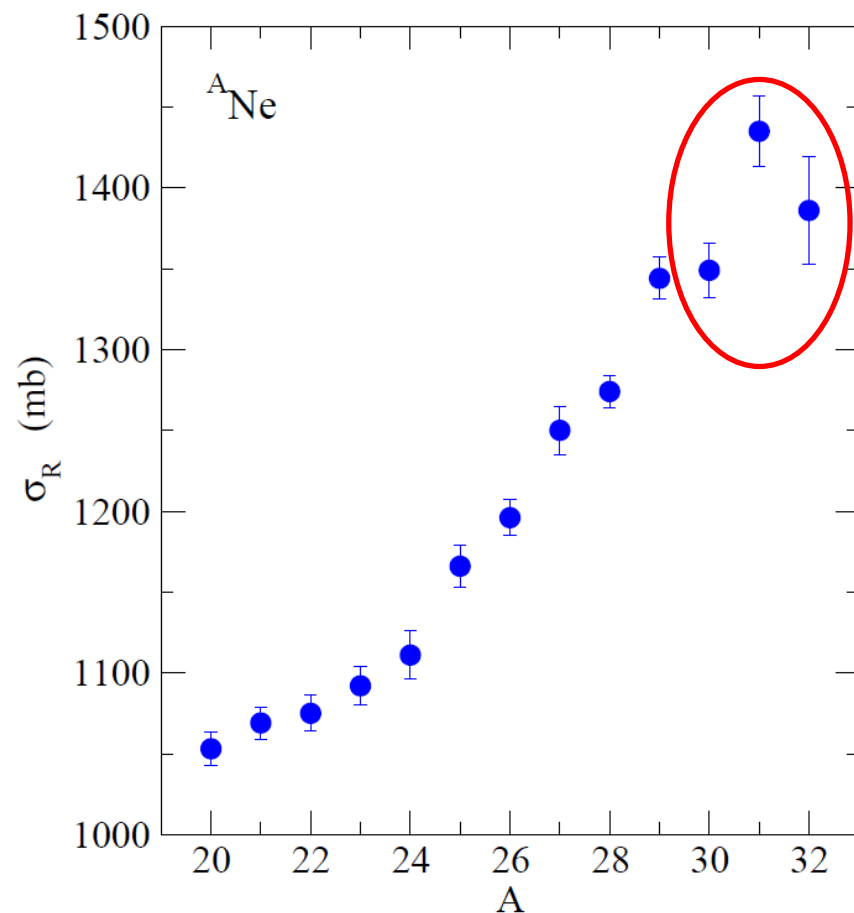


asymptotic behavior of s.p.  
wave functions



suppression of density distribution

➤ odd-even staggering of  $\sigma_R$



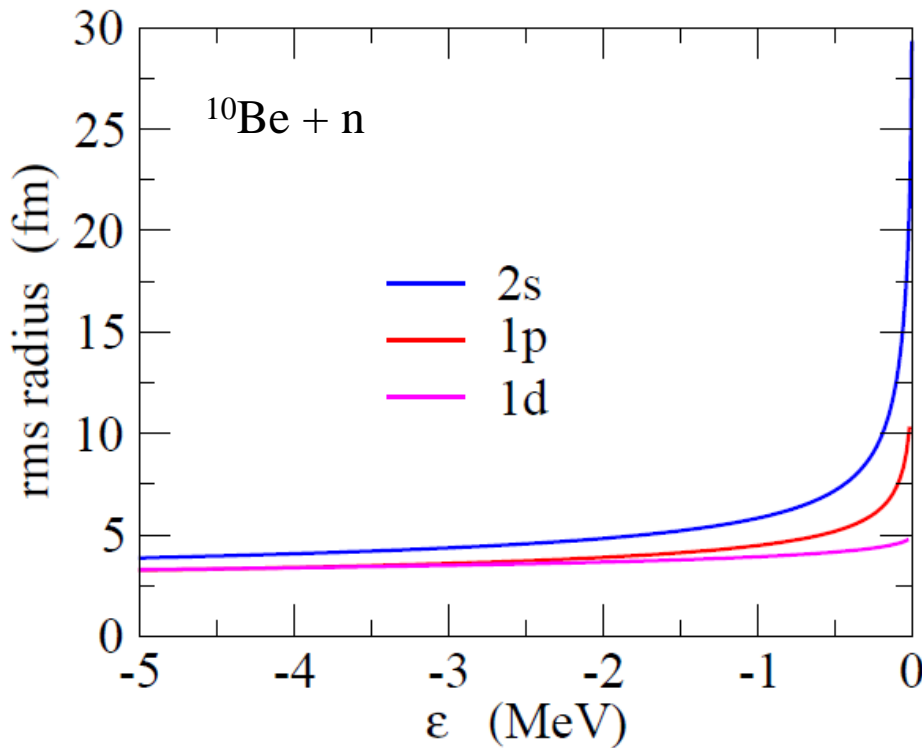
First experimental evidence for the anti-halo effect?

# Effect of pairing on radius of a weakly-bound orbit

asymptotic behavior of a s.p. wave function for s-wave:

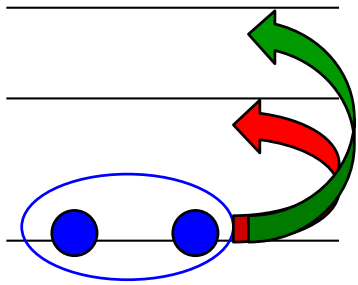
$$\psi(r) \sim \exp(-\kappa r) \quad \kappa = \sqrt{\frac{2m|\epsilon|}{\hbar^2}}$$

$$\langle r^2 \rangle_{\text{HF}} = \frac{\int r^2 |\psi(r)|^2 dr}{\int |\psi(r)|^2 dr} \propto \frac{1}{\kappa^2} = \frac{\hbar^2}{2m|\epsilon|} \rightarrow \infty$$



$$\langle r^2 \rangle \propto \begin{cases} \frac{1}{|\epsilon|} & (l = 0) \\ \frac{1}{\sqrt{|\epsilon|}} & (l = 1) \\ \text{const.} & (l = 2) \end{cases}$$

For even-mass system:



Cooper pair

Hartree-Fock-Bogoliubov (HFB) equations:

$$\begin{pmatrix} \hat{h} - \lambda & \Delta(r) \\ \Delta(r) & -\hat{h} + \lambda \end{pmatrix} \begin{pmatrix} U_k(r) \\ V_k(r) \end{pmatrix} = E_k \begin{pmatrix} U_k(r) \\ V_k(r) \end{pmatrix}$$

$\Delta(r)$ : pair potential  
 $\lambda$ : chemical potential

density:  $\rho(\mathbf{r}) = \sum_k |V_k(\mathbf{r})|^2$

Asymptotic form of  $V_k(r)$  :

$$V_k(r) \sim \exp(-\beta_k r)$$

$$\beta_k = \sqrt{\frac{2m(E_k - \lambda)}{\hbar^2}} \underset{\uparrow}{\sim} \sqrt{\frac{2m\Delta}{\hbar^2}}$$

$$E_k \sim \sqrt{(\epsilon - \lambda)^2 + \Delta^2} \sim \Delta$$

$(\epsilon, \lambda \rightarrow 0)$

$$\langle r^2 \rangle_{\text{HFB}} \propto \frac{\hbar^2}{2m\Delta}$$

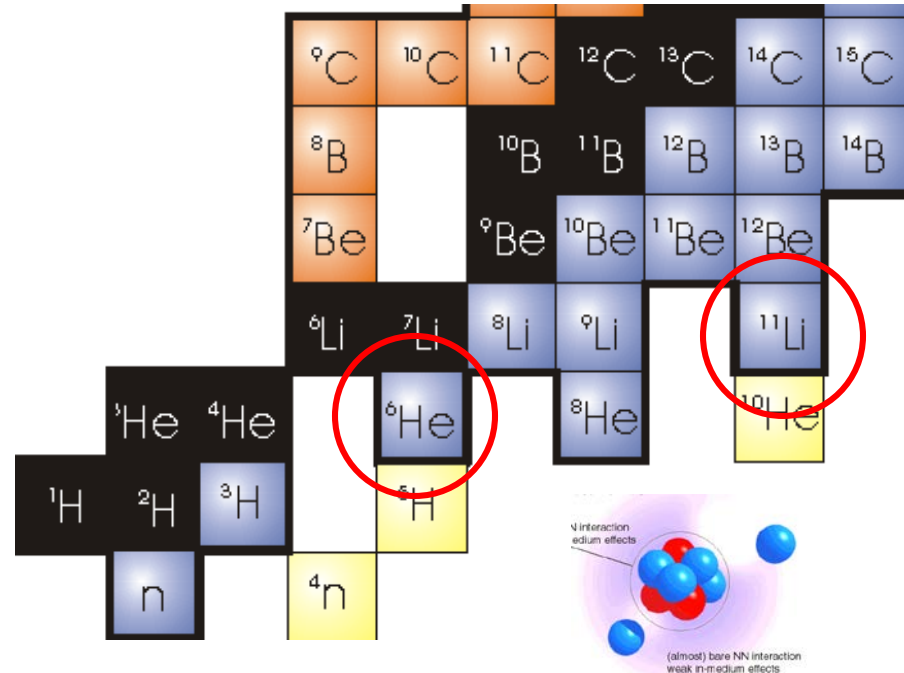
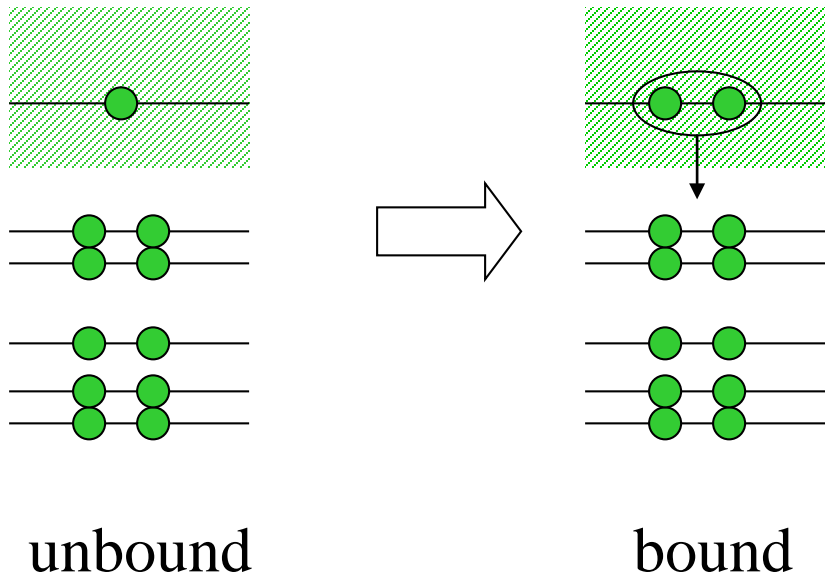
“pairing anti-halo effect”

# Pairing correlation in weakly-bound nuclei

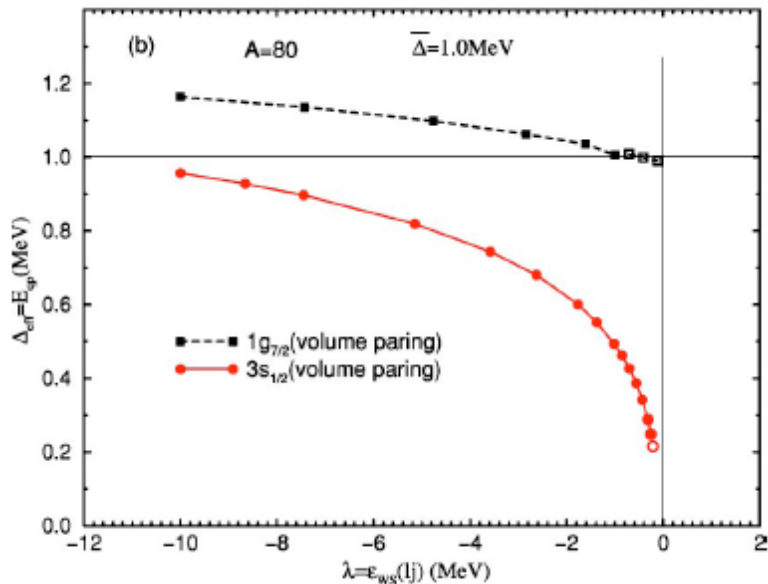
$$\langle r^2 \rangle_{\text{HFB}} \propto \frac{\hbar^2}{2m\Delta} \quad \text{“pairing anti-halo effect”}$$

$$\Delta \neq 0 \quad \text{as } \epsilon, \lambda \rightarrow 0?$$

cf. for light neutron-rich nuclei (Borromean nuclei)



# For heavier nuclei: controversial arguments based on HFB



I. Hamamoto and H. Sagawa,  
PRC70('04)034317

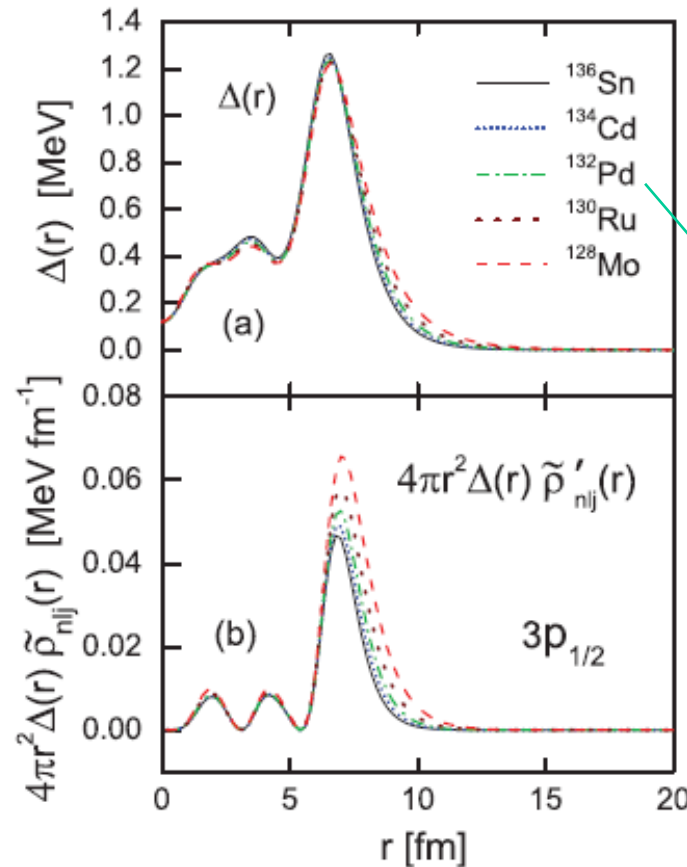
## ■ simplified HFB model

✓  $\Delta(r)$ : prefixed

✓ set  $\lambda = \epsilon_{\text{HF}}$

✓ define  $\Delta_{\text{eff}} = \text{lowest } E_{\text{qp}}$

➔  $\Delta_{\text{eff}} \rightarrow 0 \quad (\epsilon \rightarrow 0)$



e.g.  
 $\epsilon = -0.01$   
(MeV)  
 $\Delta = 0.57$   
(MeV)  
( $3p_{3/2}$ )

Y. Zhang, M. Matsuo, J. Meng,  
PRC83('11)054301

## ■ self-consistent HFB

$$\Delta_{\text{eff}} \neq 0 \quad (\epsilon \rightarrow 0)$$

see also M. Yamagami, PRC72('05)064308

# Model


## HFB with a Woods-Saxon mean-field potential

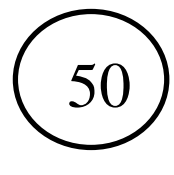
$$\begin{pmatrix} \hat{h} - \lambda & \Delta(r) \\ \Delta(r) & -\hat{h} + \lambda \end{pmatrix} \begin{pmatrix} U_k(r) \\ V_k(r) \end{pmatrix} = E_k \begin{pmatrix} U_k(r) \\ V_k(r) \end{pmatrix}$$

$$\hat{h} = -\frac{\hbar^2}{2m} \nabla^2 + V_{\text{WS}}(r)$$

$\uparrow$   
 ${}^{76}_{24}\text{Cr}_{52}$

-0.05 MeV —————  $2d_{5/2}$   
-0.26 MeV —————  $3s_{1/2}$





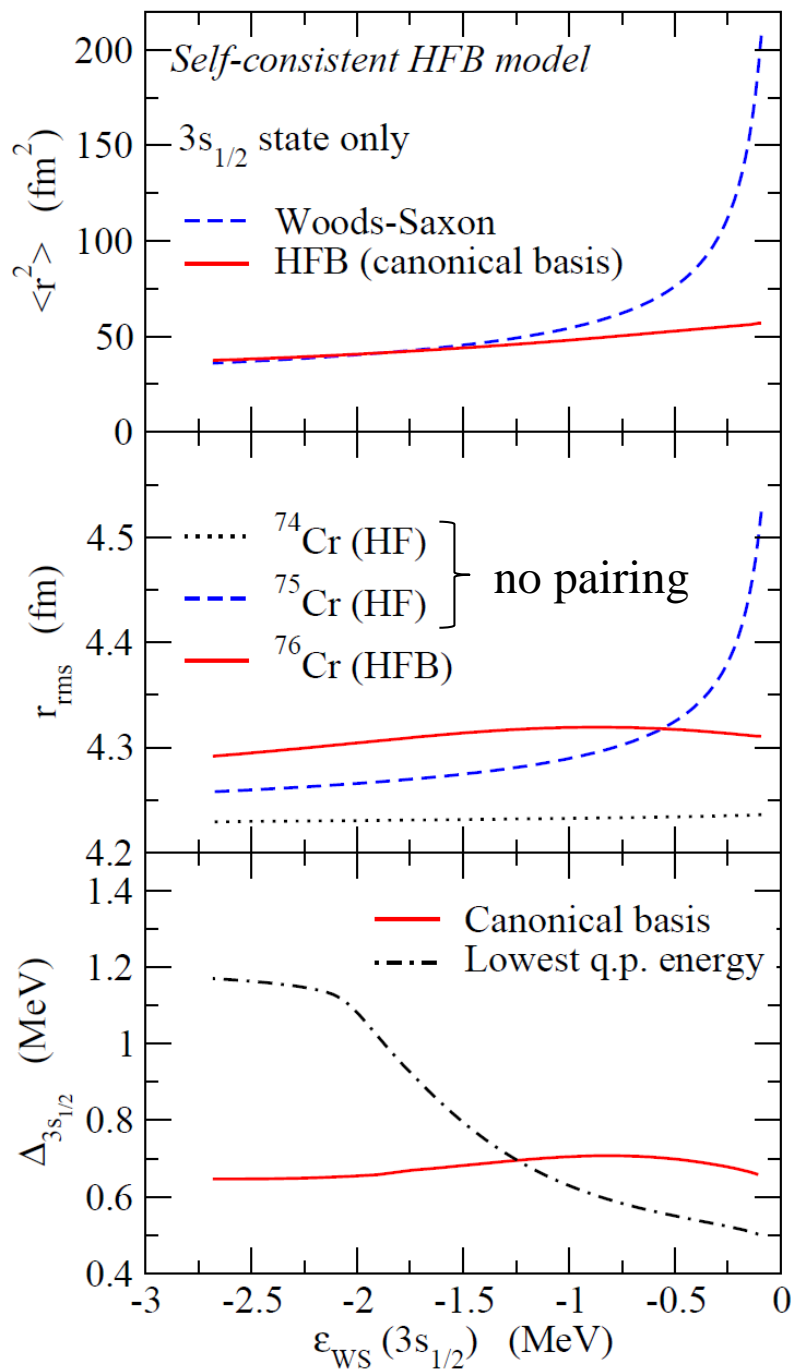
$$\Delta(r) = \frac{V_{\text{pair}}}{2} \left( 1 - \frac{\rho(r)}{\rho_0} \right) \tilde{\rho}_n(r) \quad V_{\text{pair}} \leftarrow \bar{\Delta} = 1.0 \text{ MeV}$$

$$\tilde{\rho}_n(r) = - \sum_{k=n} U_k^*(\mathbf{r}) V_k(\mathbf{r})$$

✓  $\lambda$ : self-consistently determined so that  $N=52$

✓  $E_{\text{cut}} = 50 \text{ MeV}$  above  $\lambda$

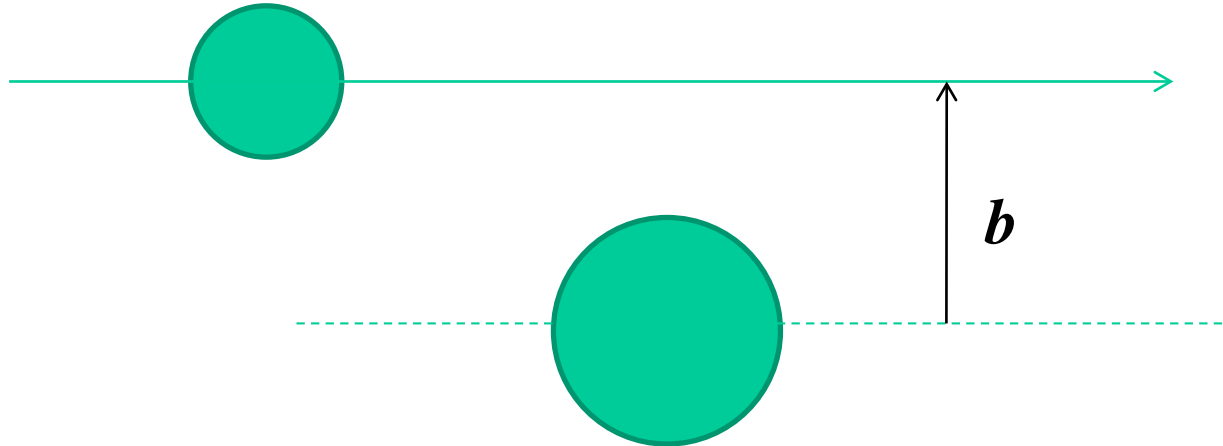
✓  $R_{\text{box}} = 60 \text{ fm}$



← suppression of the radius

← the effective pairing gap persists for both the definitions (agreement with Zhang-Matsuo-Meng)

# Reaction cross sections



Glauber theory (optical limit approximation:OLA)

$$\sigma_R = \int d^2b \left( 1 - |e^{i\chi(b)}|^2 \right)$$

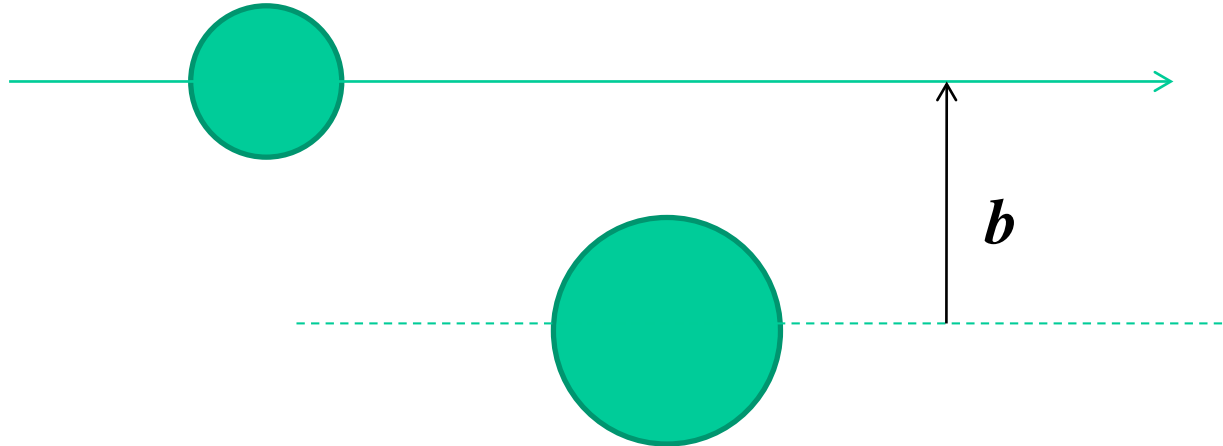
$$e^{i\chi(b)} = \exp \left[ - \int d\mathbf{r}_P d\mathbf{r}_T \rho_P(\mathbf{r}_P) \rho_T(\mathbf{r}_T) \Gamma(\mathbf{b} + \mathbf{s}_P - \mathbf{s}_T) \right]$$

$$\Gamma(\mathbf{b}) = \frac{1 - i\alpha}{4\pi\beta} \sigma_{NN}^{\text{tot}} \exp \left( -\frac{b^2}{2\beta} \right)$$

- straight-line trajectory
- adiabatic approximation
- simplified treatment for multiple scattering:  $(1 - x)^N \rightarrow e^{-Nx}$



# Reaction cross sections



Glauber theory (optical limit approximation:OLA)

$$\sigma_R = \int d^2b \left( 1 - |e^{i\chi(b)}|^2 \right)$$

$$e^{i\chi(b)} = \exp \left[ - \int d\mathbf{r}_P d\mathbf{r}_T \rho_P(\mathbf{r}_P) \rho_T(\mathbf{r}_T) \Gamma(\mathbf{b} + \mathbf{s}_P - \mathbf{s}_T) \right]$$

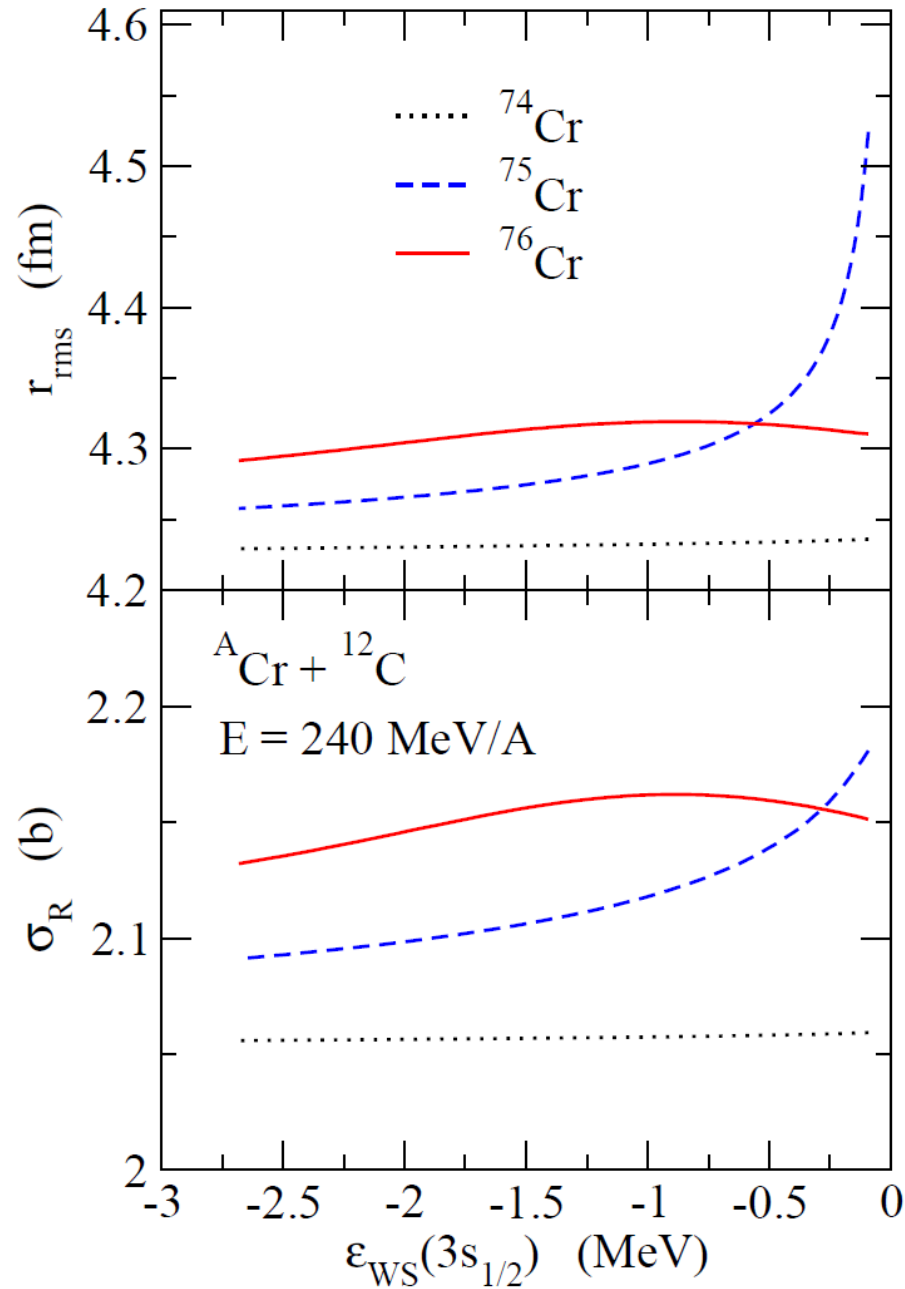
• Correction to the OLA

B. Abu-Ibrahim and Y. Suzuki, PRC61('00)051601(R)

$$i\chi(b) \rightarrow - \int d\mathbf{r}_P \rho_P(\mathbf{r}_P) \left[ 1 - e^{- \int d\mathbf{r}_T \rho_T(\mathbf{r}_T) \Gamma(\mathbf{b} + \mathbf{s}_P - \mathbf{s}_T)} \right]$$

$^{74,75,76}\text{Cr} + ^{12}\text{C}$  reactions  
at  $E=240 \text{ MeV/A}$

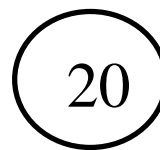
density of  $^{74,75,76}\text{Cr}$  : HFB  
density of  $^{12}\text{C}$  : Gaussian



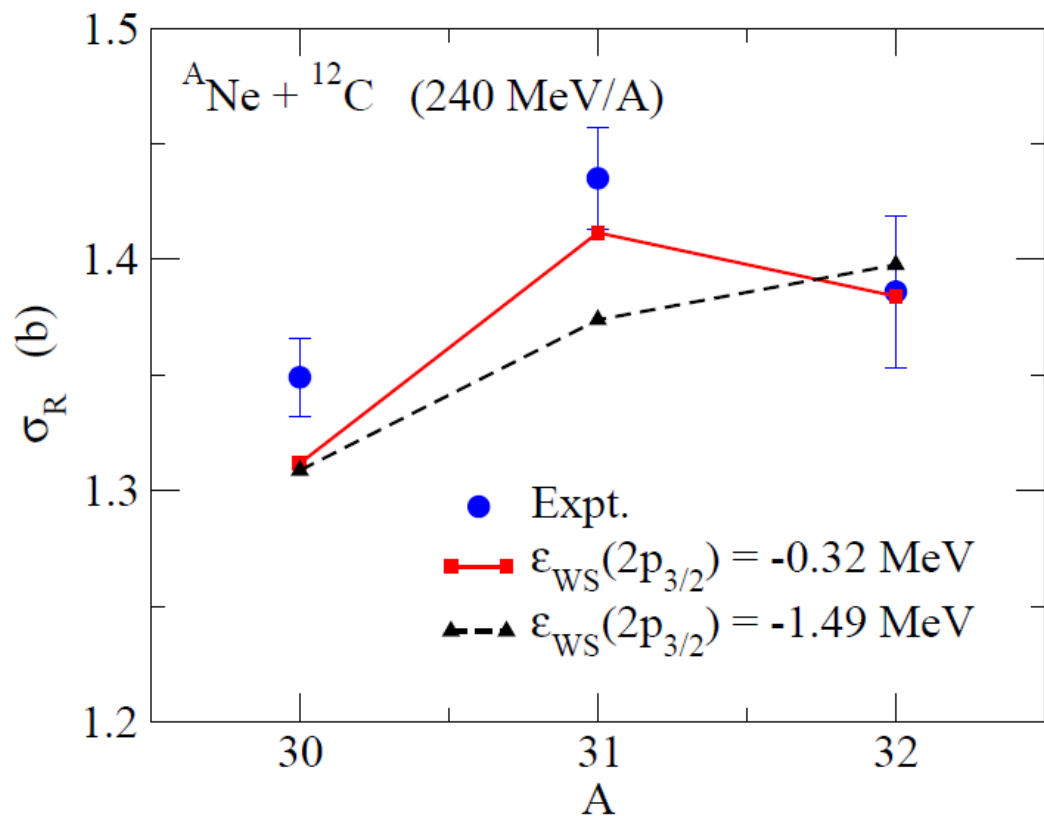
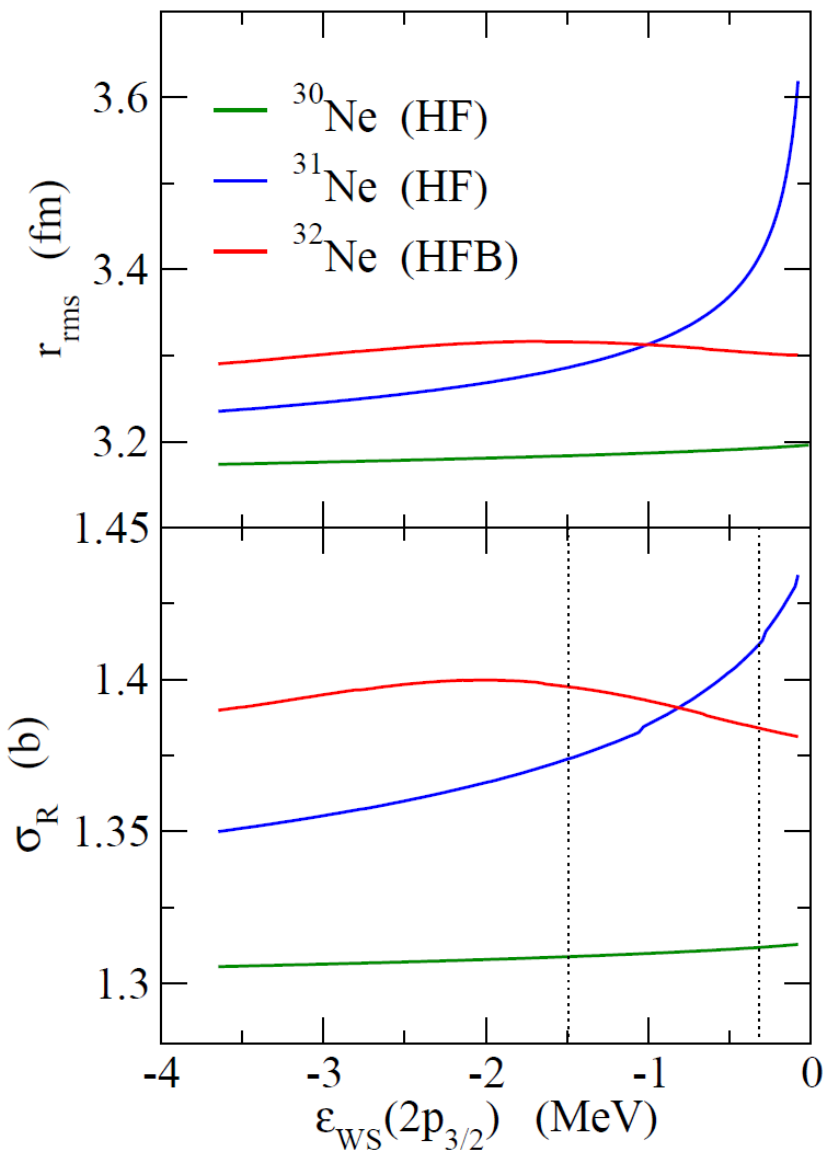
# Other system: $^{30,31,32}\text{Ne}$

HFB with a spherical Woods-Saxon

-0.066 MeV ———  $1f_{7/2}$   
-0.321 MeV ———  $2p_{3/2}$

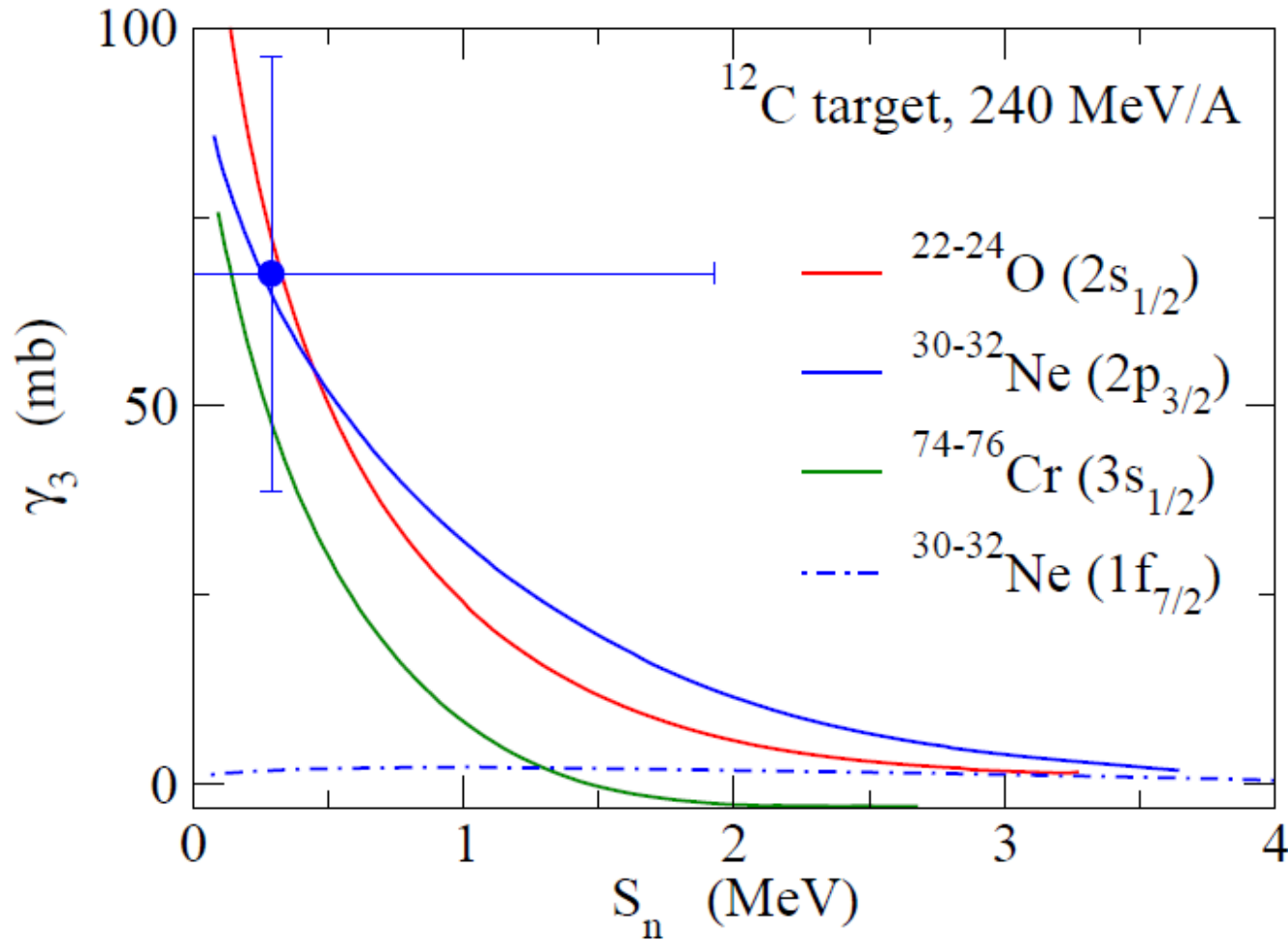


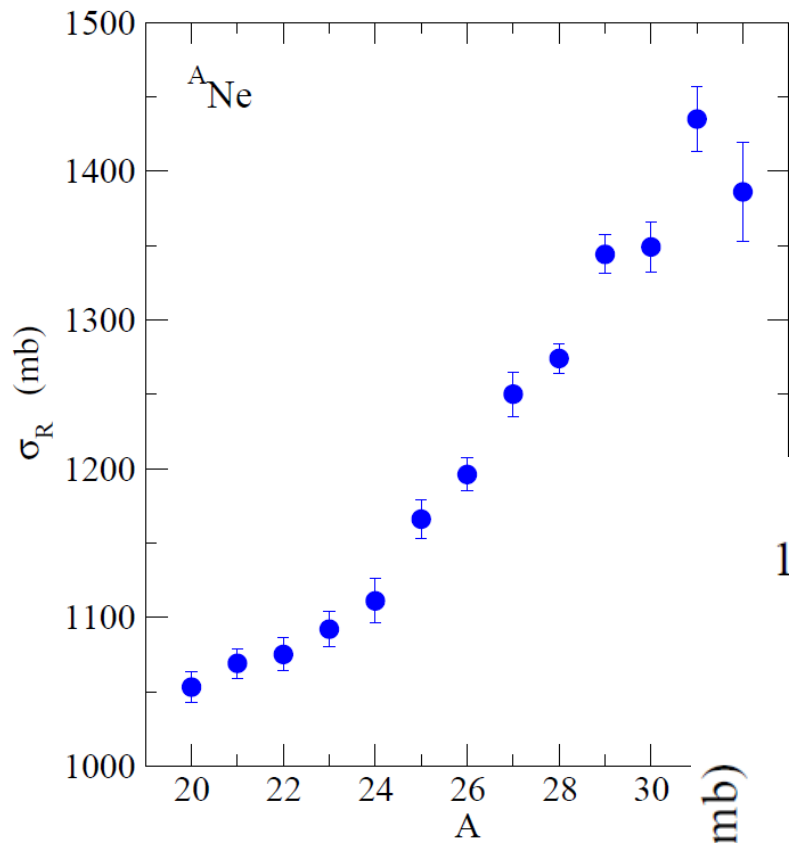
$^{31}\text{Ne}$  ( $a = 0.75$  fm)



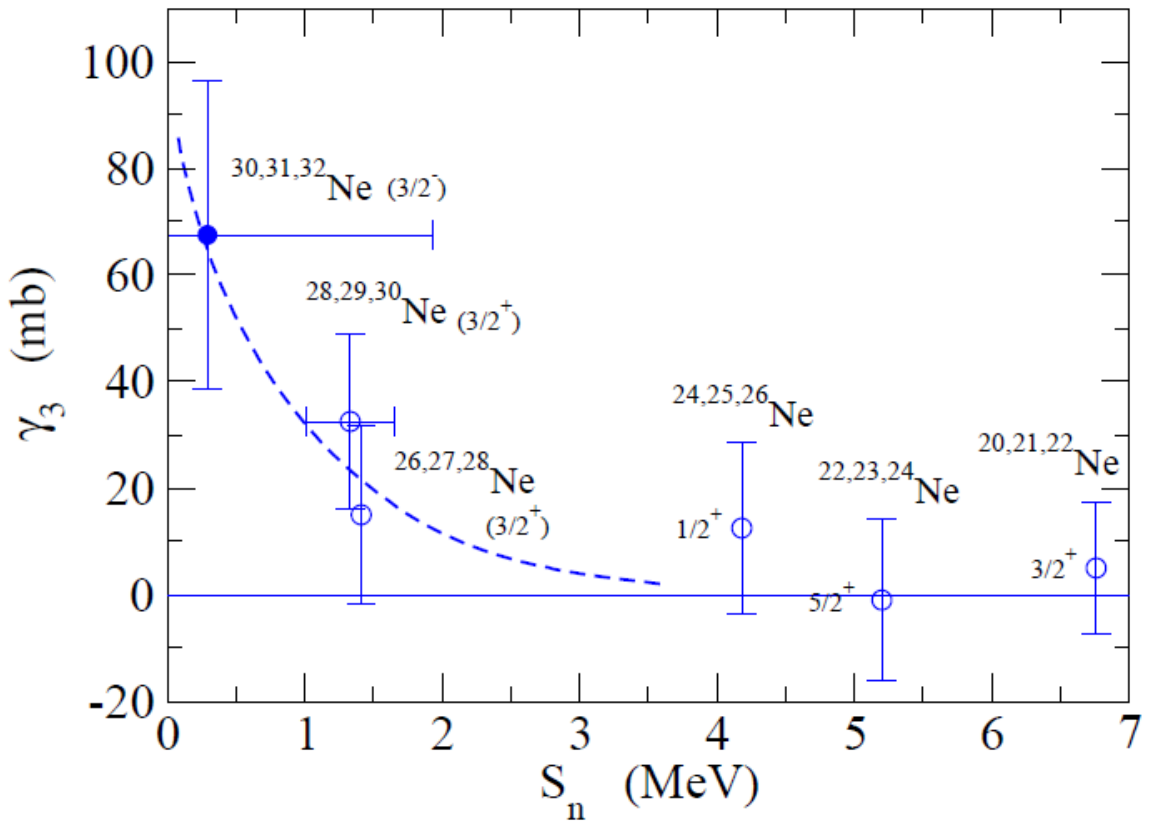
# Systematic study: OES parameter

$$\gamma_3 \equiv -\frac{1}{2}[\sigma_R(A+2) - 2\sigma_R(A+1) + \sigma_R(A)]$$

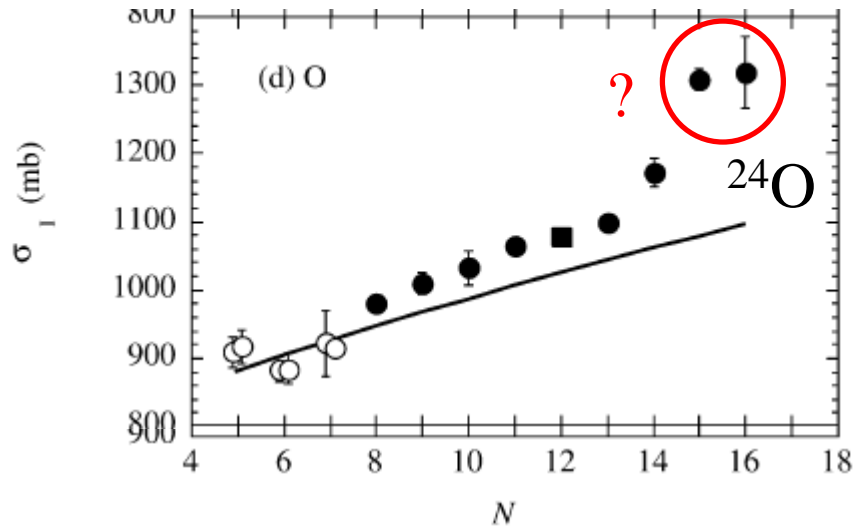




systematics

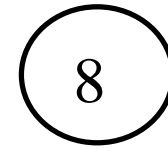


$^{22,23,24}\text{O} + ^{12}\text{C} @ 950 \text{ MeV/A}$

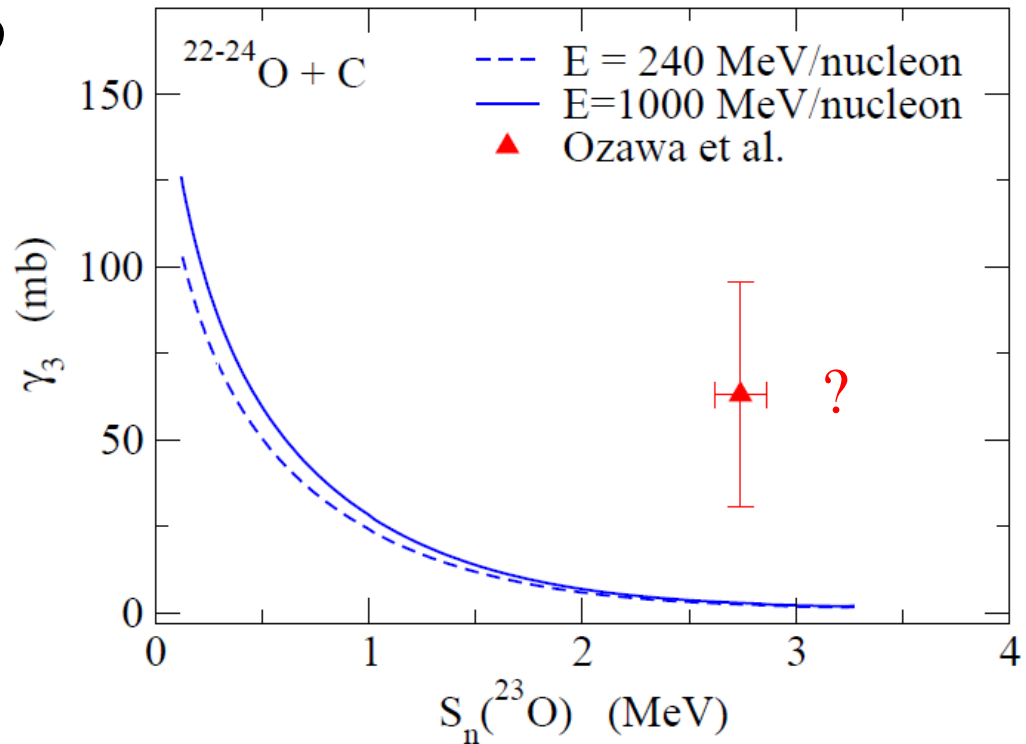


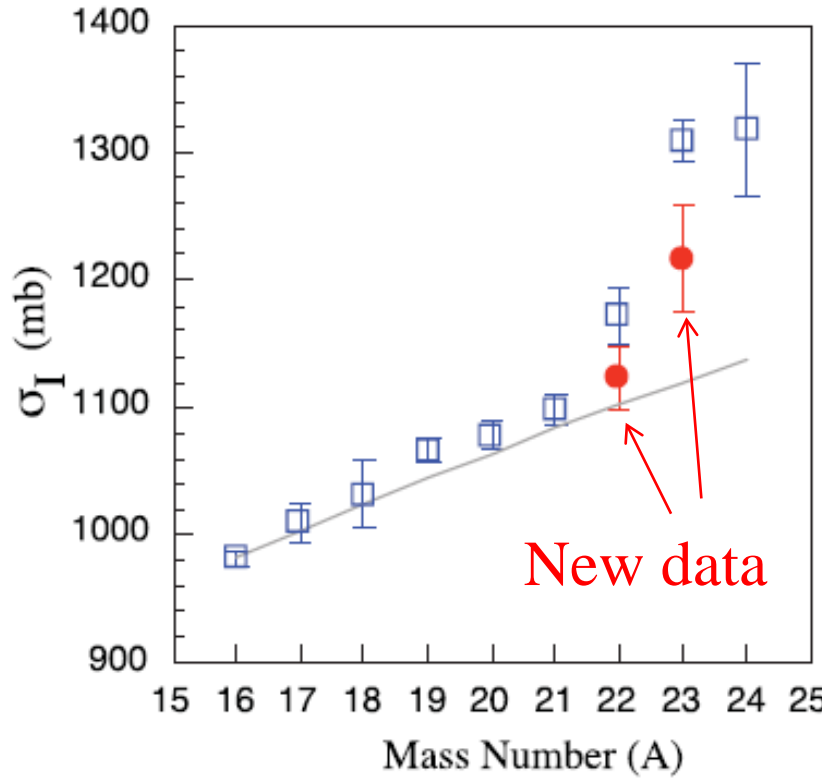
A. Ozawa et al., NPA691('01)599

-2.62 MeV ————  $2s_{1/2}$   
 -3.57 MeV ————  $1d_{5/2}$

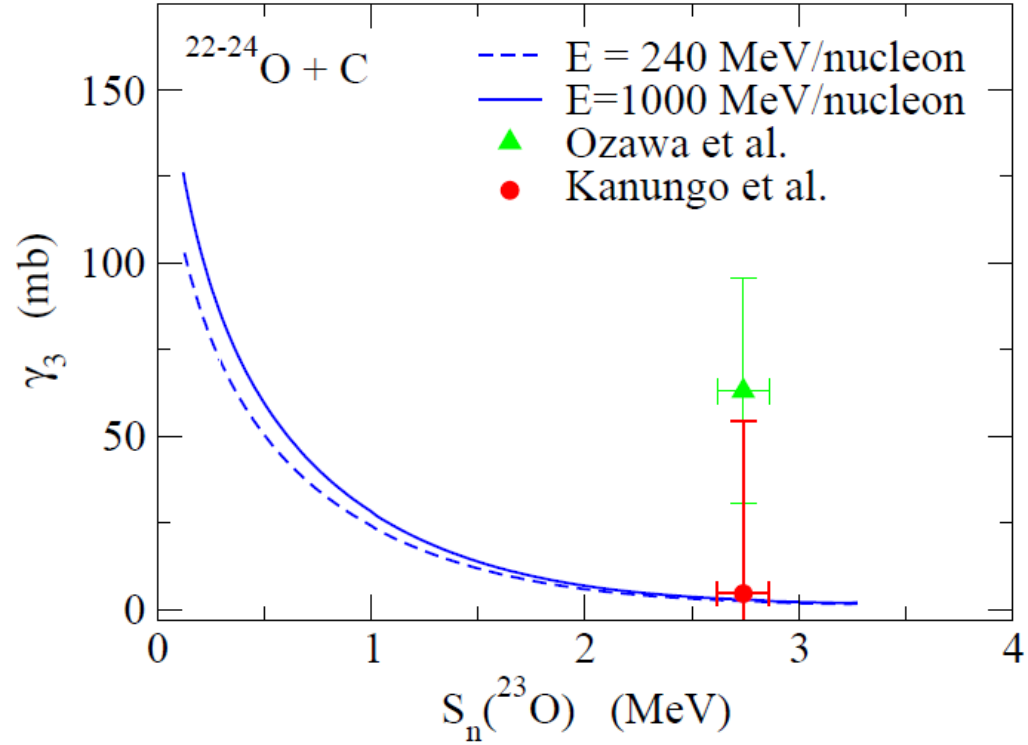


$^{24}\text{O}$





R. Kanungo et al.,  
 PRC84('11)061304(R)



K.H. and H. Sagawa,  
 arXiv:1202.2725 [nucl-th]

# Summary

## ➤ Analyses of $\sigma_R$ with HFB + Glauber

weakly-bound even-even nuclei:

- ✓ the pairing correlation persists even at the drip
- ✓ suppression of the radius due to *the pairing correlation ( $l = 0, 1$ )*

➡ reduction of  $\sigma_R$

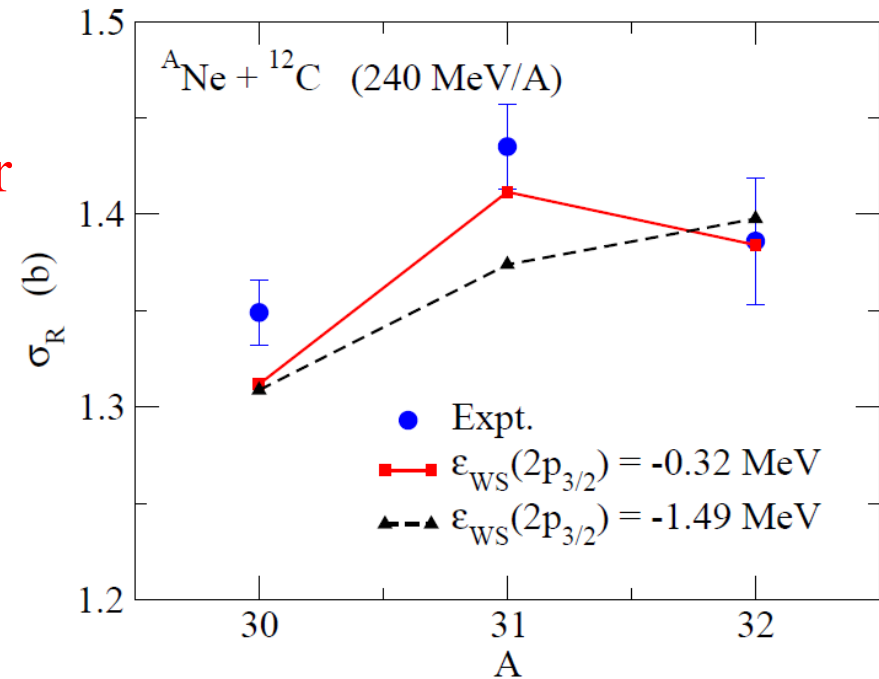


## Odd-even staggering of $\sigma_R$

### ➤ Odd-even staggering parameter

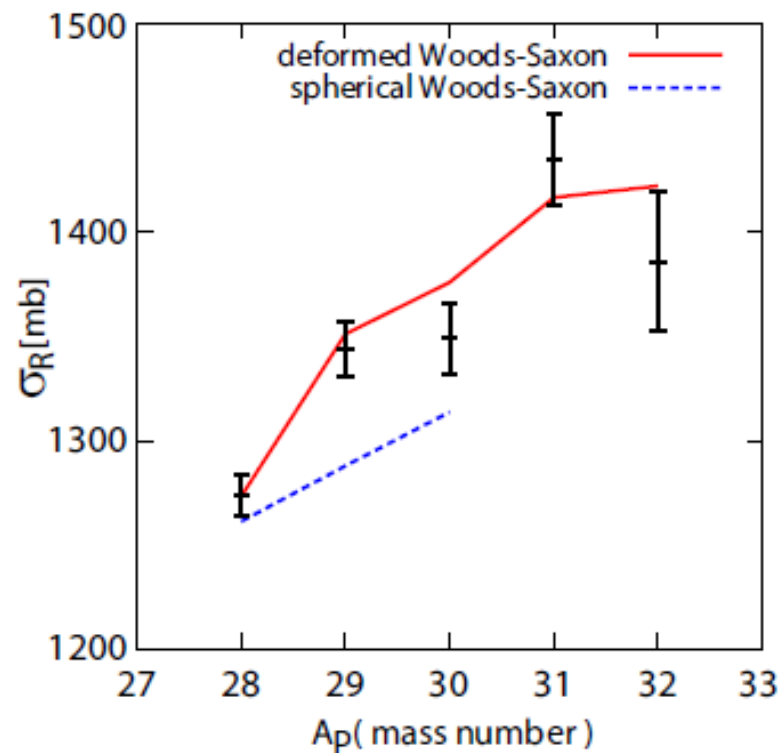
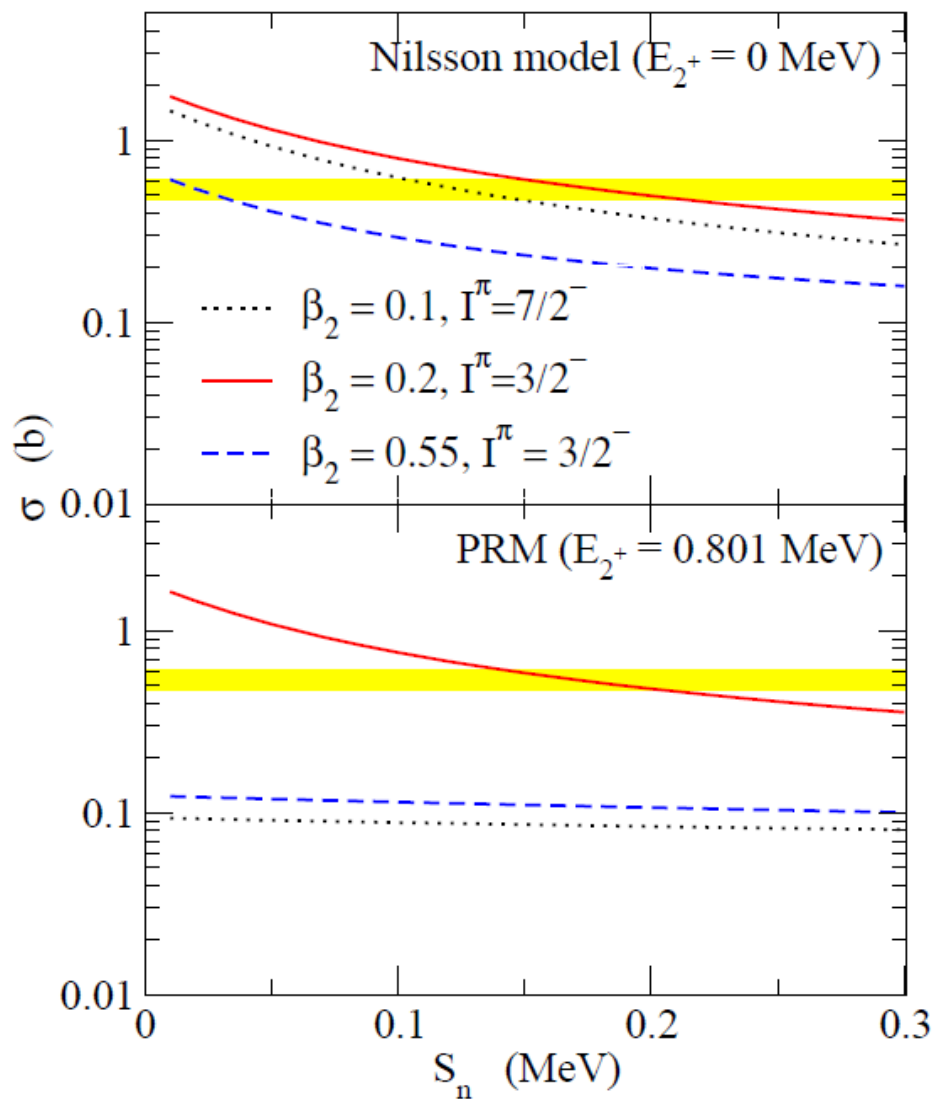
a good tool to investigate the pairing correlation in weakly bound nuclei

### ➤ Work in progress: deformation effects





# Deformation of $^{31}\text{Ne}$



K. Minomo et al.,  
PRC84('11)034602  
PRL108('12)052503

## Coulomb breakup cross sections

Y. Urata, K.H., H. Sagawa,  
PRC83('11)041303(R)