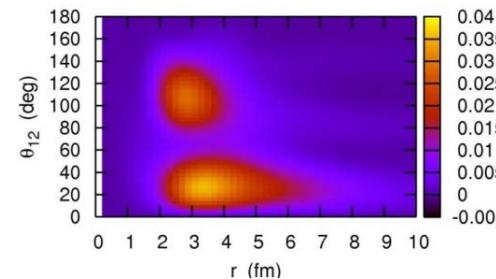
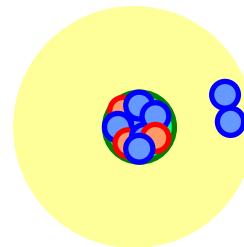


Di-neutron correlation and two-neutron decay of nuclei beyond the neutron drip line

Kouichi Hagino

Tohoku University, Sendai, Japan

Hiroyuki Sagawa (*U. of Aizu*)



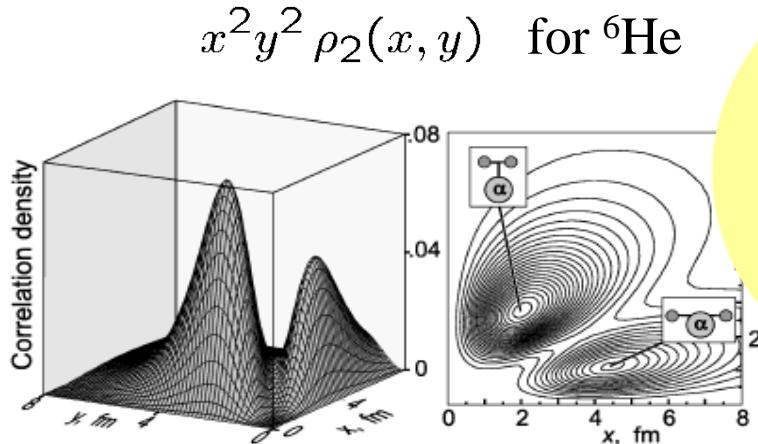
1. *Introduction: Di-neutron correlation*
2. *Coulomb breakup of Borromean nuclei*
3. *Two-neutron decay of unbound nucleus ^{26}O*
4. *Summary*

Borromean nuclei and Di-neutron correlation

Borromean nuclei: unique three-body systems

Three-body model calculations:

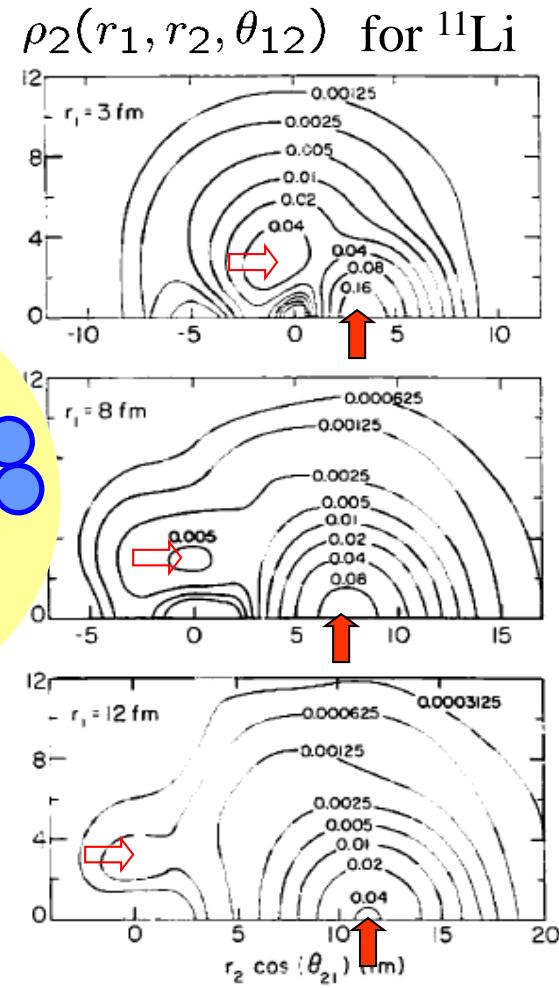
strong di-neutron correlation
in ^{11}Li and ^6He



Yu.Ts. Oganessian et al., *PRL*82('99)4996
M.V. Zhukov et al., *Phys. Rep.* 231('93)151

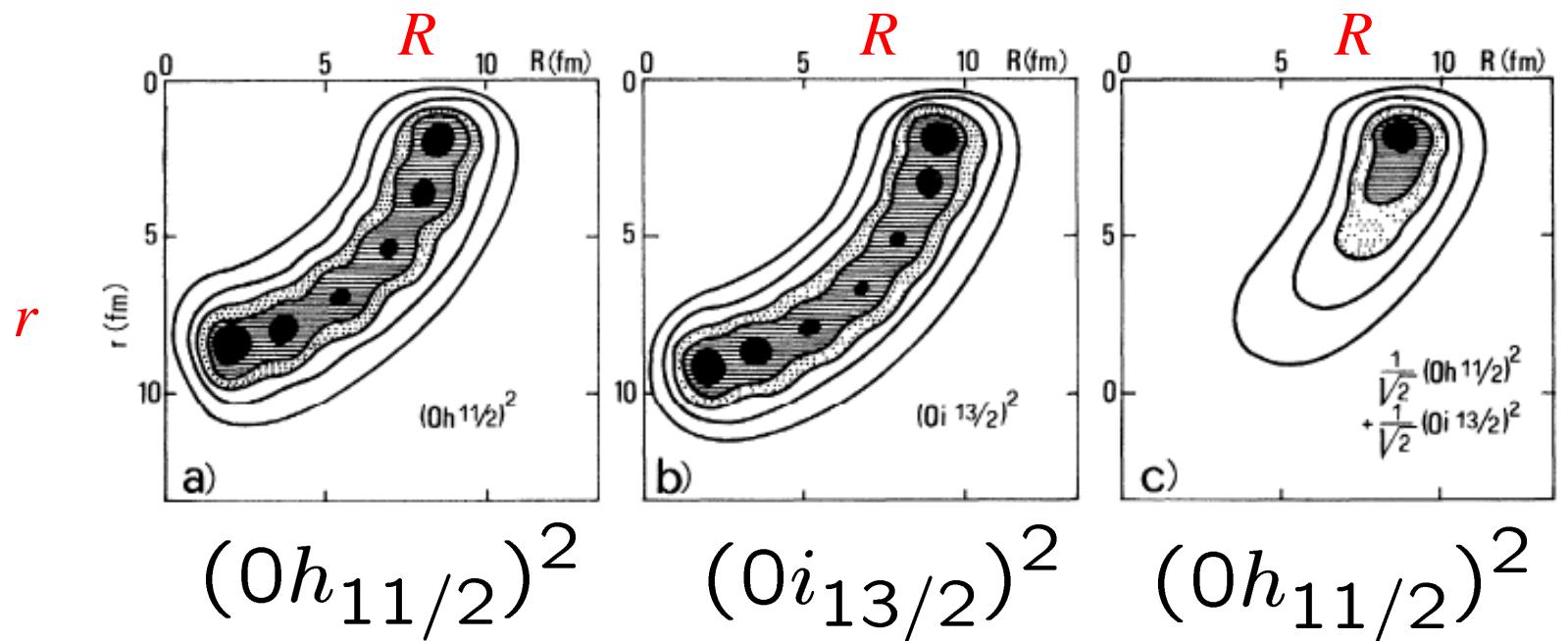
cf. earlier works

- ✓ A.B. Migdal ('73)
- ✓ P.G. Hansen and B. Jonson ('87)

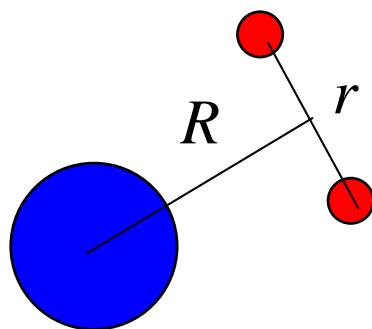


G.F. Bertsch, H. Esbensen,
Ann. of Phys., 209('91)327

dineutron correlation: caused by the admixture of different parity states



F. Catara, A. Insolia, E. Maglione,
and A. Vitturi, PRC29('84)1091

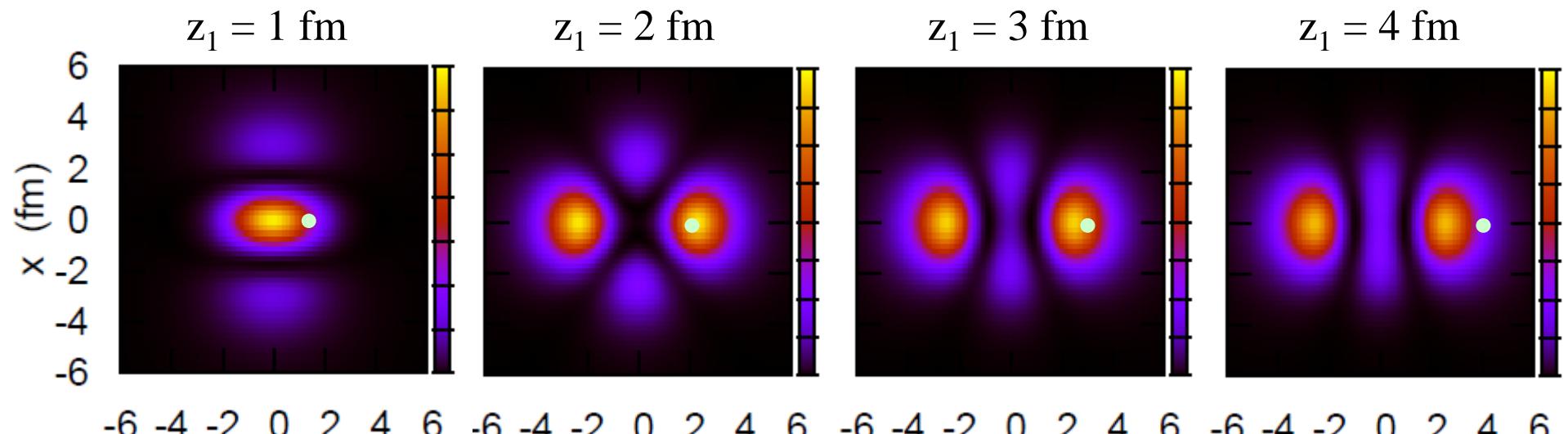


interference of even and odd partial waves

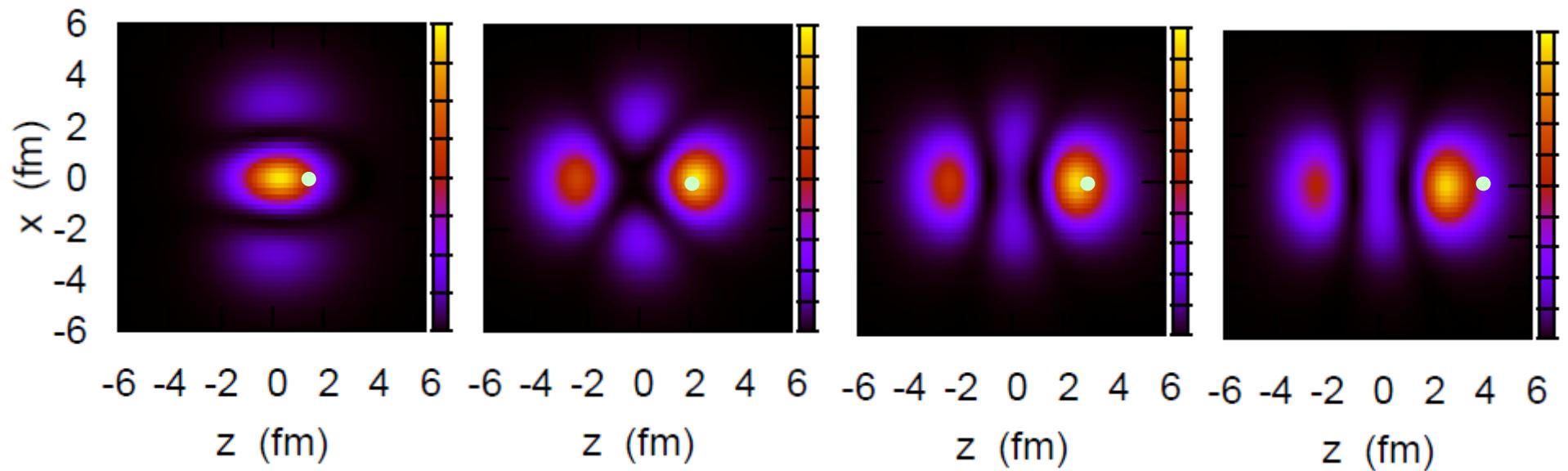
$$\rho_2(x_1, x_2) = |\Psi_{ee}(x_1, x_2)|^2 + |\Psi_{oo}(x_1, x_2)|^2 + 2\Psi_{ee}(x_1, x_2)\Psi_{oo}(x_1, x_2)$$

Example: $^{18}\text{O} = ^{16}\text{O} + \text{n} + \text{n}$ cf. ^{17}O : 3 bound states ($1\text{d}_{5/2}$, $2\text{s}_{1/2}$, $1\text{d}_{3/2}$)

i) even parity only



ii) both even and odd parities

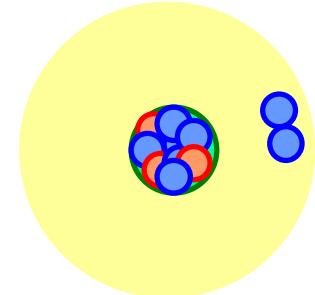


Dineutron correlation in the momentum space

$$\Psi(r, r') = \alpha \Psi_{s^2}(r, r') + \beta \Psi_{p^2}(r, r') \rightarrow \theta_r = 0: \text{enhanced}$$

→ Fourier transform

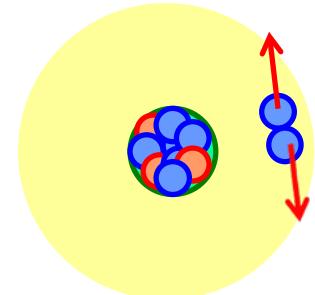
$$\tilde{\Psi}(k, k') = \int e^{ik \cdot r} e^{ik' \cdot r'} \Psi(r, r') dr dr'$$



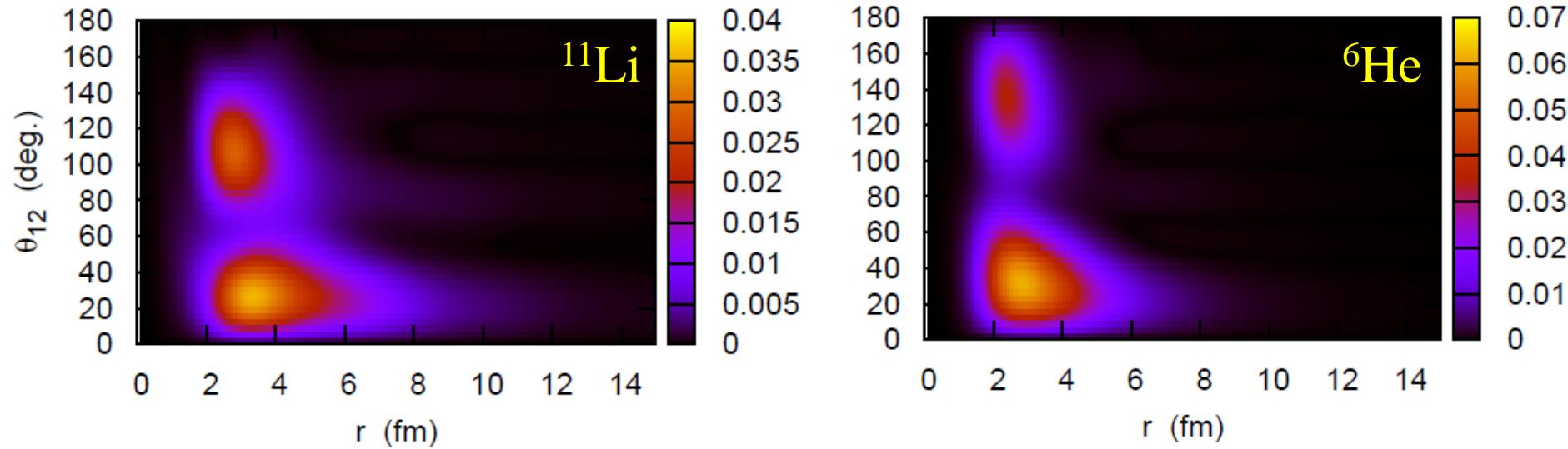
$$e^{ik \cdot r} = \sum_l (2l+1) i^l \dots \rightarrow i^l \cdot i^l = i^{2l} = (-)^l$$

$\uparrow \quad \uparrow$
 $r \quad r'$

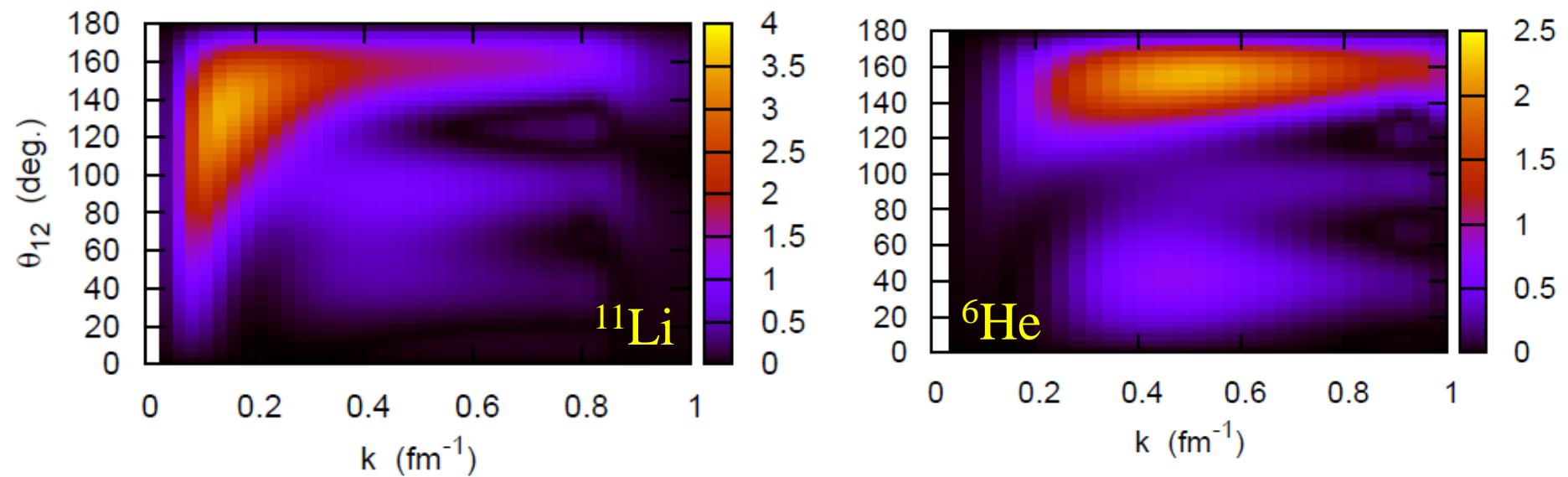
$$\tilde{\Psi}(k, k') = \alpha \tilde{\Psi}_{s^2}(k, k') - \beta \tilde{\Psi}_{p^2}(k, k') \rightarrow \theta_k = \pi: \text{enhanced}$$



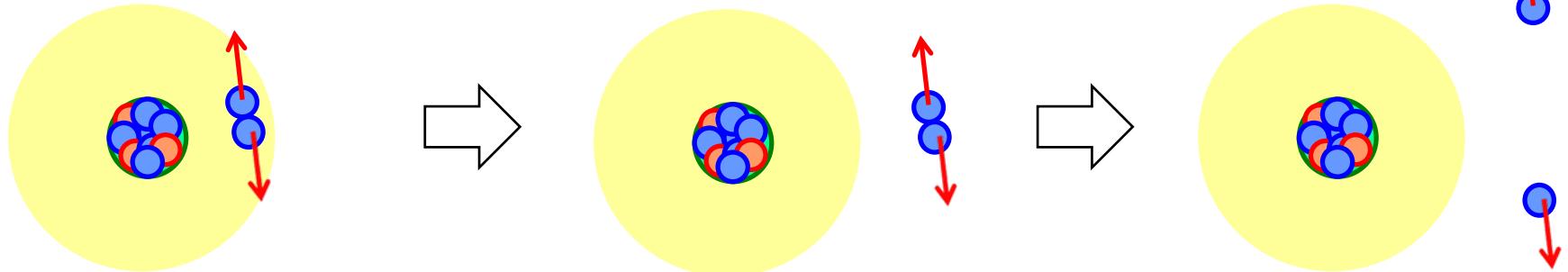
Two-particle density in the r space: $8\pi^2 r^4 \sin \theta \cdot \rho(r, r, \theta)$



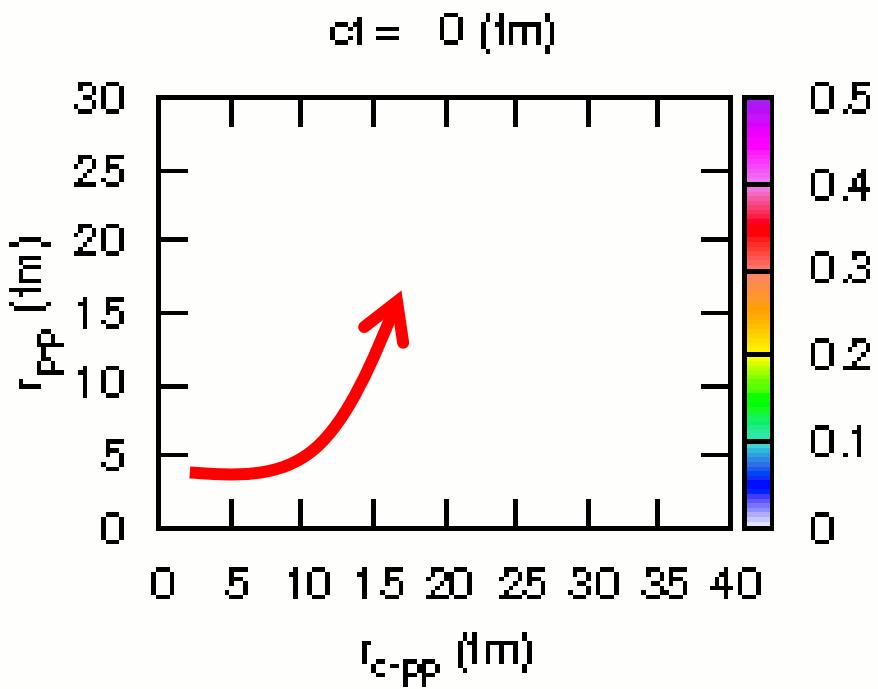
Two-particle density in the p space: $8\pi^2 k^4 \sin \theta \cdot \rho(k, k, \theta)$



Consequence to a two-nucleon emission decay

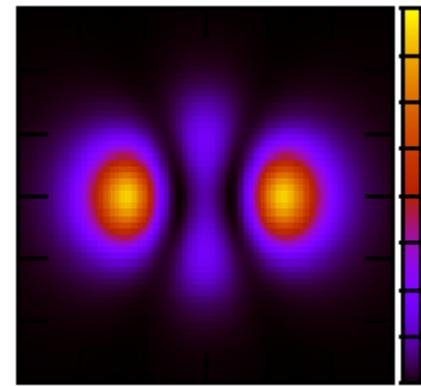


2p decay of ${}^6\text{Be}$
: time-dependent calculations



T. Oishi (Tohoku → Jyvaskyla),
K.H., H. Sagawa,
PRC90 ('14) 034303

Di-neutron correlation in weakly-bound exotic nuclei (WBEN)

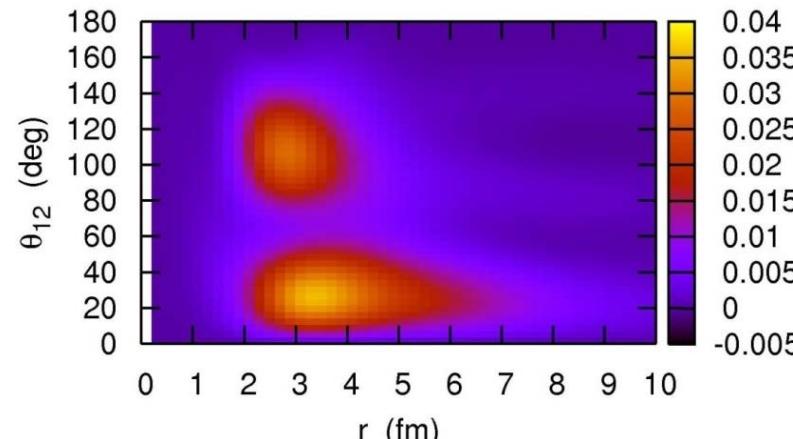
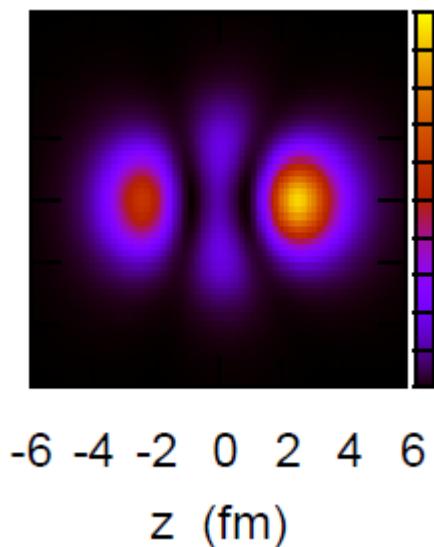


spatial localization of two neutrons
(dineutron correlation)

weakly bound systems

→ easy to mix different parity states due to
the continuum couplings
+ enhancement of pairing on the surface
→ dineutron correlation: enhanced

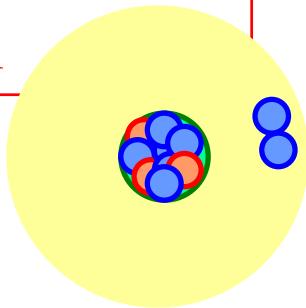
cf. - Bertsch, Esbensen, Ann. of Phys. 209('91)327
- M. Matsuo, K. Mizuyama, Y. Serizawa,
PRC71('05)064326



K.H. and H. Sagawa,
PRC72('05)044321

Di-neutron correlations in neutron-rich nuclei

Strong di-neutron correlations
in neutron-rich nuclei



✓ Borromean nuclei (3body calc.)

Bertsch-Esbensen ('91)

Zhukov et al. ('93)

Hagino-Sagawa ('05)

Kikuchi-Kato-Myo ('10)

✓ Heavier nuclei (HFB calc.)

Matsuo et al. ('05)

Pillet-Sandulescu-Schuck ('07)

How to probe it?

➤ Coulomb breakup

T. Nakamura et al.
cluster sum rule

(mean value of θ_{nn})

➤ pair transfer reactions

➤ two-proton decays

Coulomb 3-body problem

➤ two-neutron decays

3-body resonance due to
a centrifugal barrier

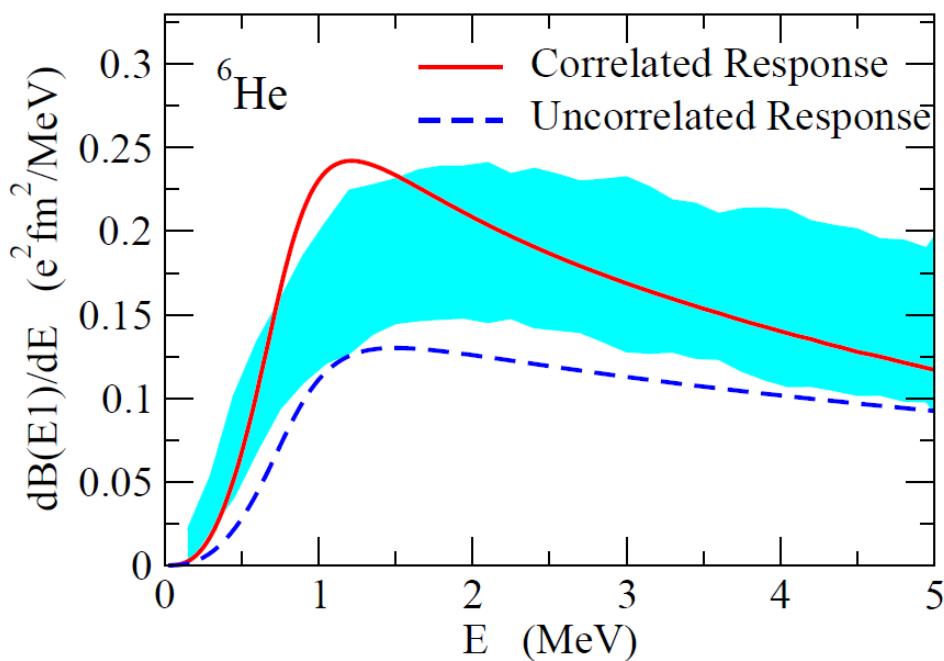
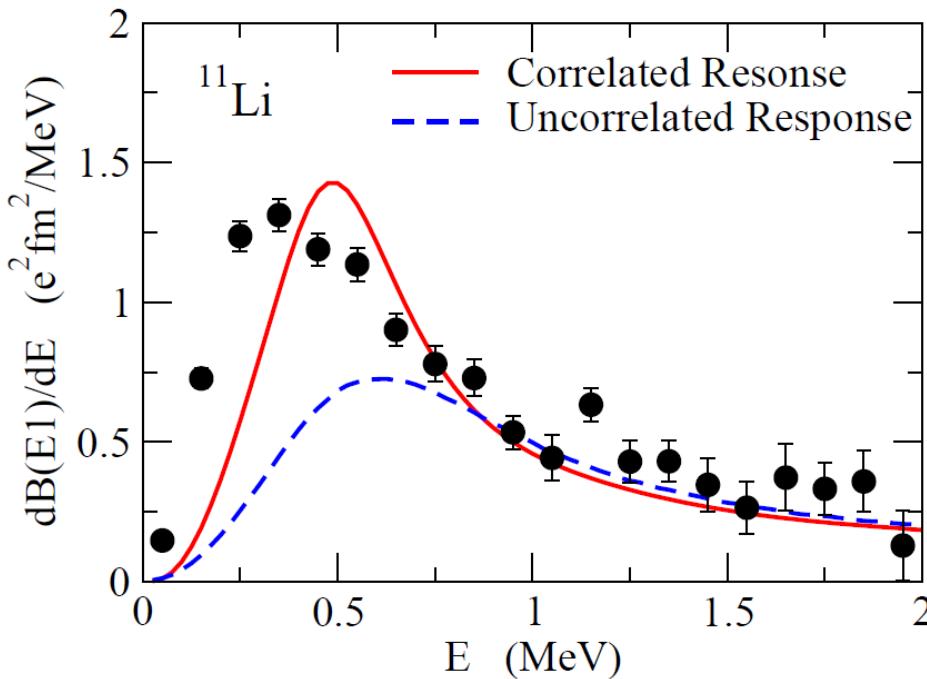
MoNA (^{16}Be , ^{13}Li , ^{26}O)

SAMURAI (^{26}O)

GSI (^{26}O)

Coulomb breakup of 2-neutron halo nuclei

How to probe the dineutron correlation? → Coulomb breakup



Experiments:

T. Nakamura et al., PRL96('06)252502

T. Aumann et al., PRC59('99)1252

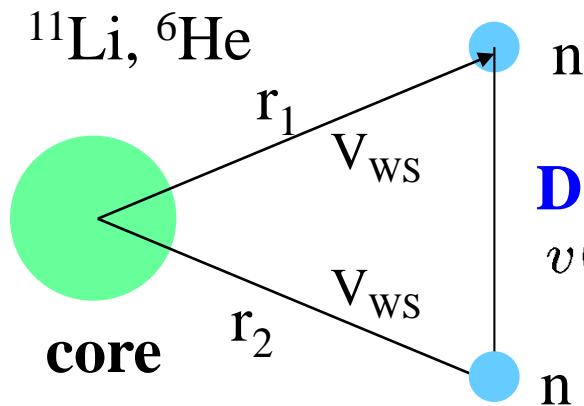
3-body model calculations:

K.H., H. Sagawa, T. Nakamura, S. Shimoura, PRC80('09)031301(R)

cf. Y. Kikuchi et al., PRC87('13)034606 ← structure of the core nucleus (^9Li)

also for ^{22}C , ^{14}Be , ^{19}B etc. (T. Nakamura et al.)

3-body model calculation for Borromean nuclei



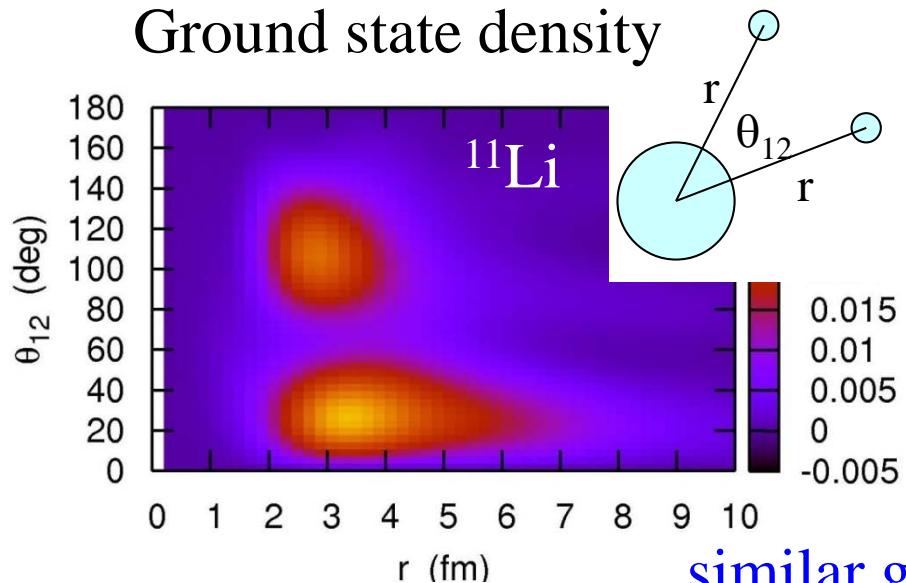
G.F. Bertsch and H. Esbensen,
Ann. of Phys. 209('91)327; *PRC*56('99)3054

Density-dependent delta-force

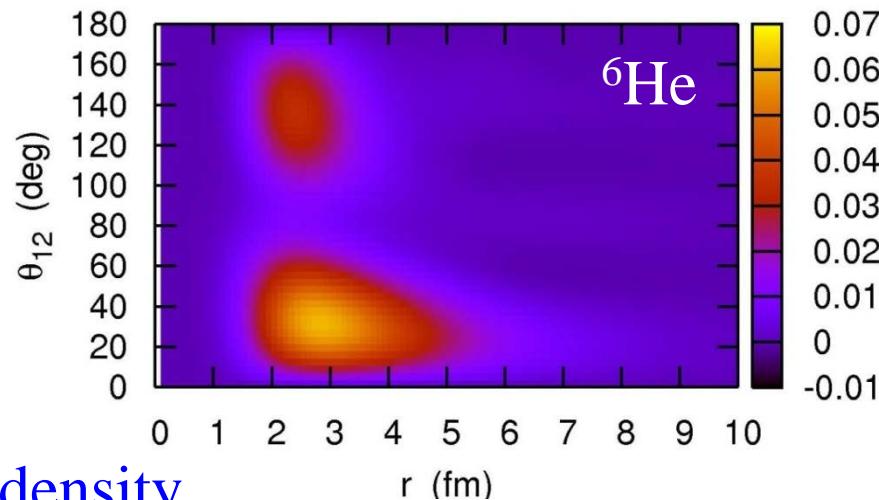
$$v(r_1, r_2) = v_0(1 + \alpha\rho(r)) \times \delta(r_1 - r_2)$$

$$H = \frac{p_1^2}{2m} + \frac{p_2^2}{2m} + V_{nC}(r_1) + V_{nC}(r_2) + V_{nn} + \frac{(p_1 + p_2)^2}{2A_c m}$$

Ground state density

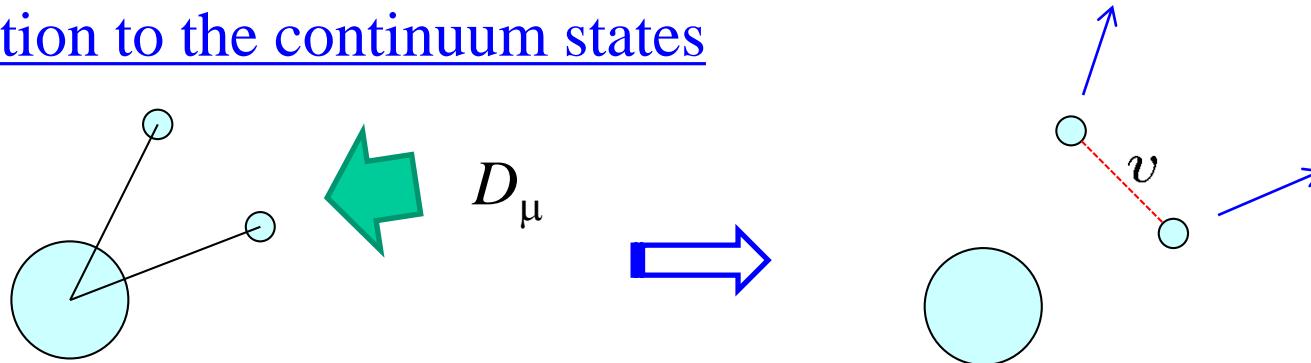


K.H. and H. Sagawa, *PRC*72('05)044321



similar g.s. density

E1 excitation to the continuum states

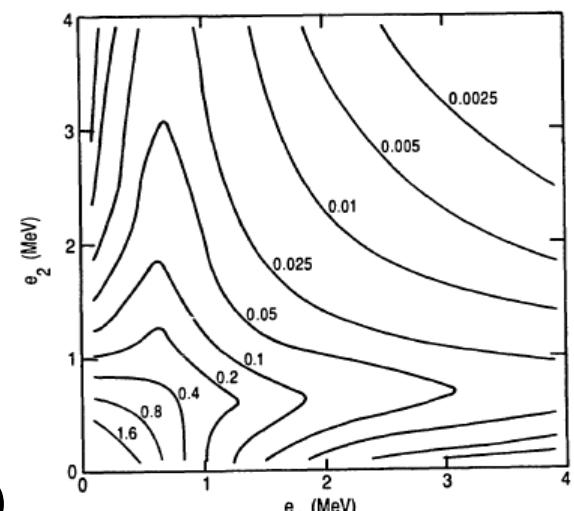


$$\begin{aligned}
 M(E1) &= \langle (j_1 j_2)_\mu^1 | (1 - vG_0 + vG_0 vG_0 - \dots) D_\mu | \Psi_{gs} \rangle \\
 &= \langle (j_1 j_2)_\mu^1 | \underbrace{(1 + vG_0)^{-1}}_{FSI} D_\mu | \Psi_{gs} \rangle
 \end{aligned}$$

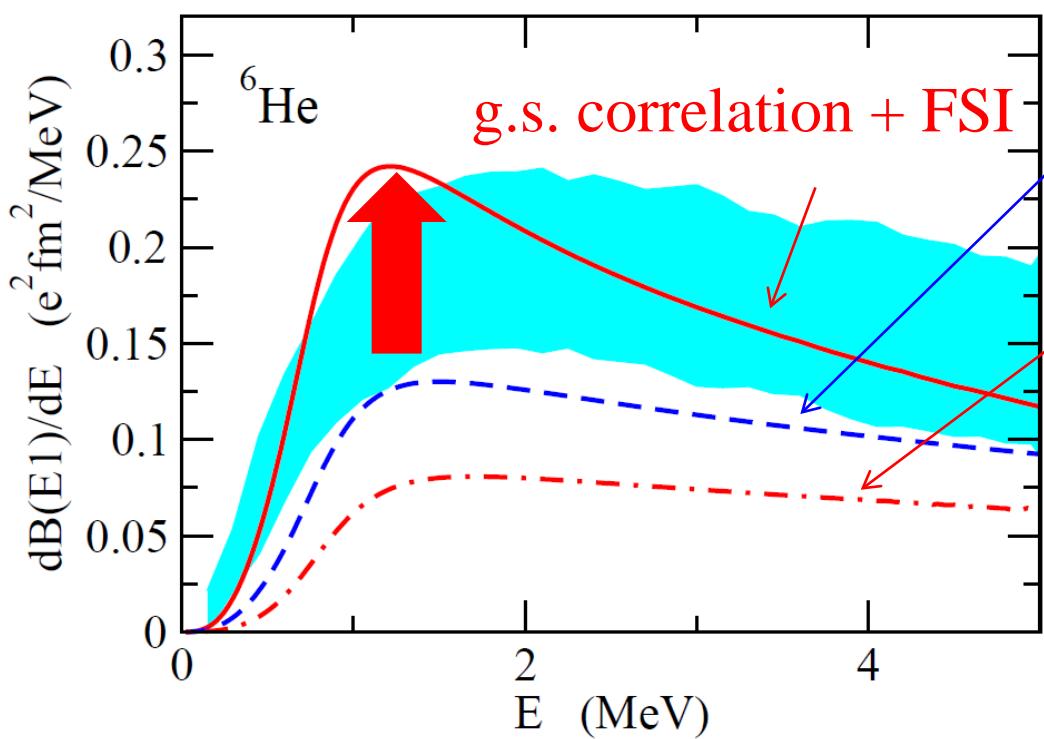
↑ ↑
 unperturbed continuum wf dipole operator

$$G_0(E) = \sum_{\mu, f.s.t.} \frac{\langle (j_1 j_2)_\mu^1 \rangle \langle (j_1 j_2)_\mu^1 |}{e_1 + e_2 - E - i\eta}$$

$$\frac{d^2 B(E1)}{de_1 de_2} = 3 \sum_{l_1 j_2 l_2 j_2} |M(E1)|^2 \frac{dk_1}{de_1} \frac{dk_2}{de_2}$$



g.s. correlation? or correlation in excited states?

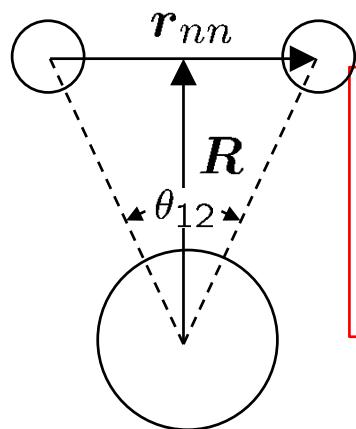


g.s. correlation only
(no nn interaction in
the final state)

g.s.: odd-l only
(no dineutron correlation)
+ FSI

- ✓ Both FSI and dineutron correlations: important role in E1 strength

Geometry of Borromean nuclei



Cluster sum rule

$$B_{\text{tot}}(E1) = \sum_f |\langle \Psi_f | \hat{T}_{E1} | \Psi_0 \rangle|^2$$

$$\sim \frac{3}{\pi} \left(\frac{Z_c e}{A_c + 2} \right)^2 \langle R^2 \rangle$$

reflects the g.s. correlation

“experimental data” for opening angle

$$\sqrt{\langle R^2 \rangle} \longleftrightarrow B_{\text{tot}}(E1)$$

$$\sqrt{\langle r_{nn}^2 \rangle} \longleftrightarrow \text{matter radius or HBT}$$

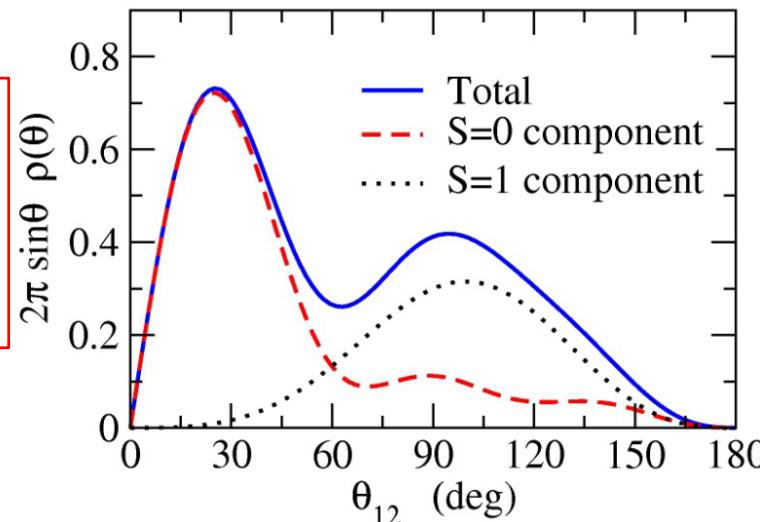
$$\begin{aligned} \langle \theta_{12} \rangle &= 65.2 \pm 12.2 \text{ } (^{11}\text{Li}) \\ &= 74.5 \pm 12.1 \text{ } (^6\text{He}) \end{aligned}$$

K.H. and H. Sagawa, PRC76('07)047302

cf. T. Nakamura et al., PRL96('06)252502

C.A. Bertulani and M.S. Hussein, PRC76('07)051602

3-body model calculations



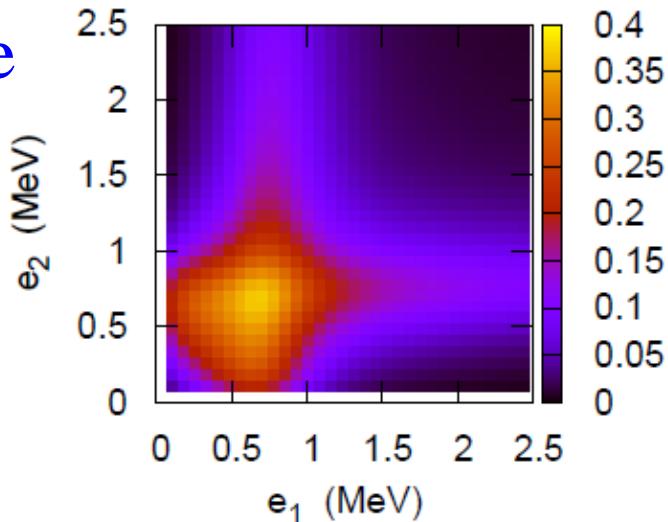
$$\langle \theta_{12} \rangle = 65.29 \text{ deg.}$$

$\langle \theta_{12} \rangle$: significantly smaller than 90 deg.

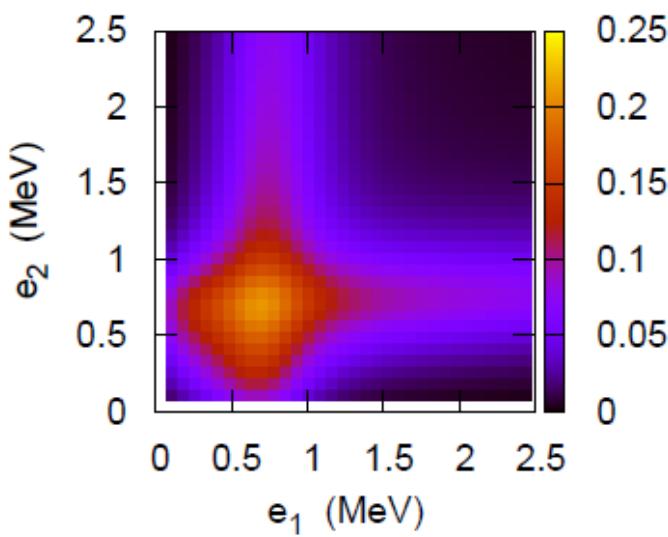
suggests dineutron corr.
(but, an average of small and large angles)

Energy distribution of emitted neutrons

^6He

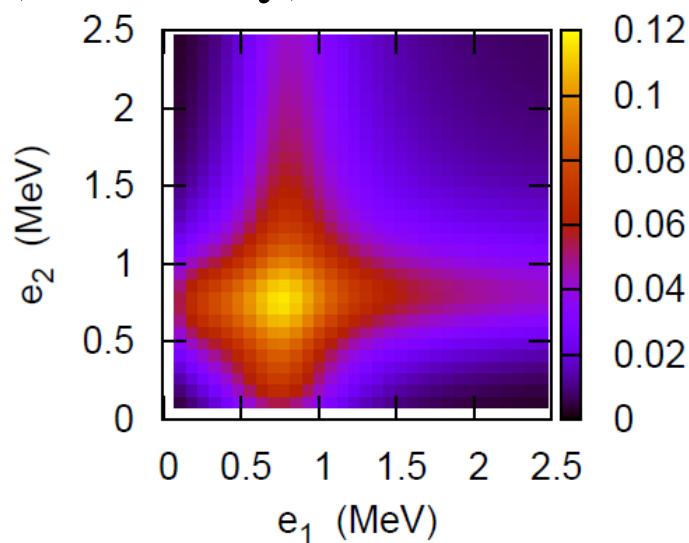


$$v_{nnn} = 0$$

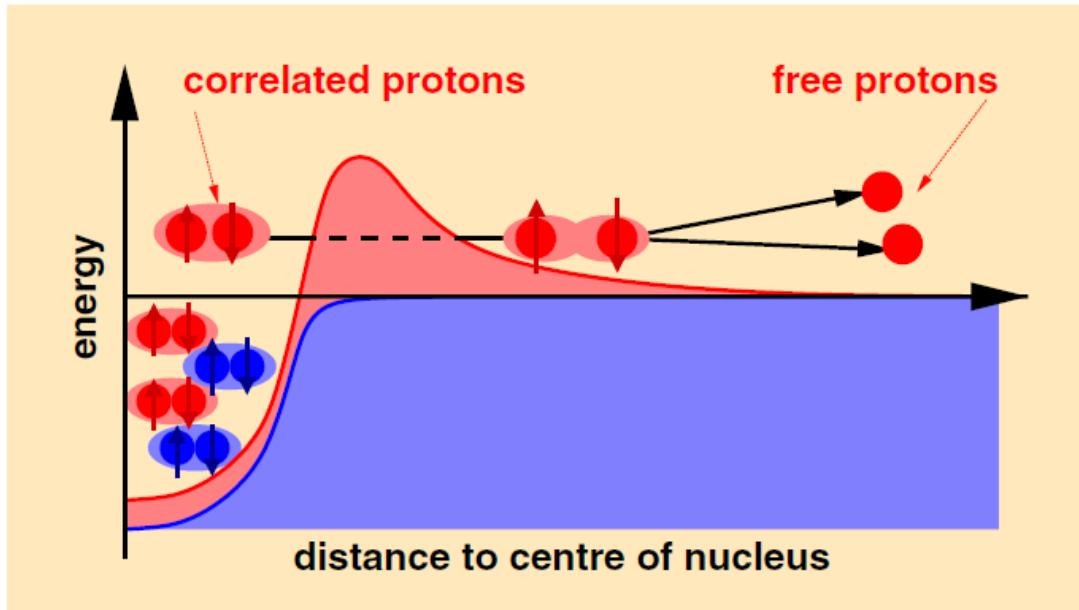


- ✓ shape of distribution: insensitive to the nn-interaction (except for the absolute value)
- ✓ strong sensitivity to V_{nC}
- ✓ similar situation in between ^{11}Li and ^6He

no di-neutron corr. in the g.s.
(odd- l only)



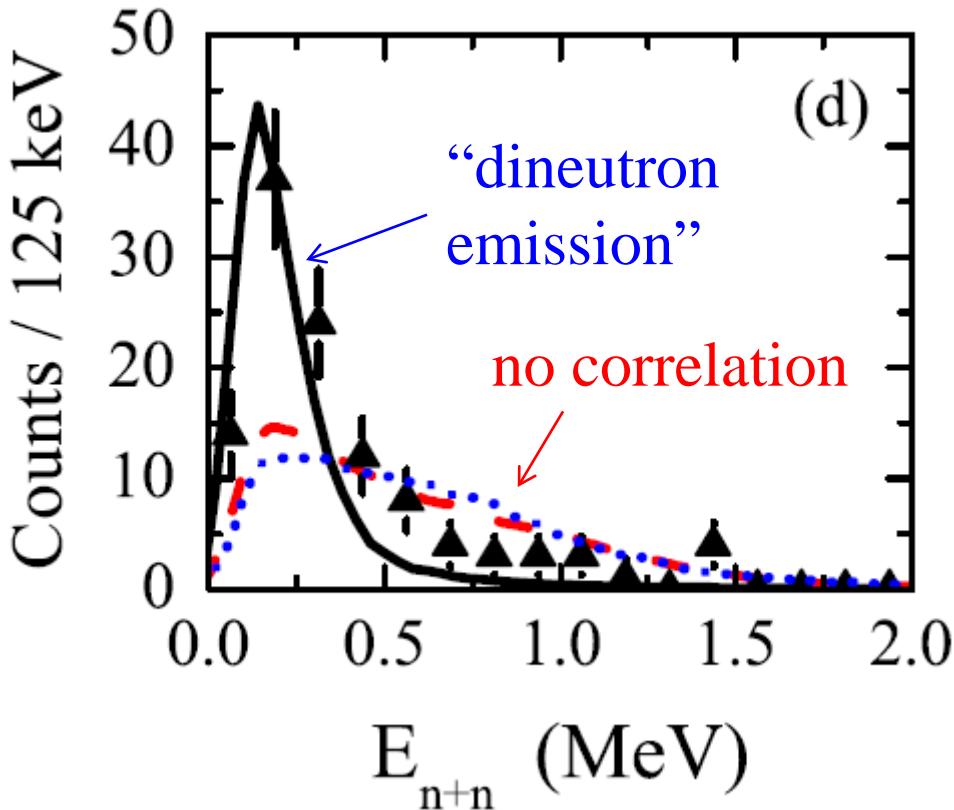
2-proton radio activity



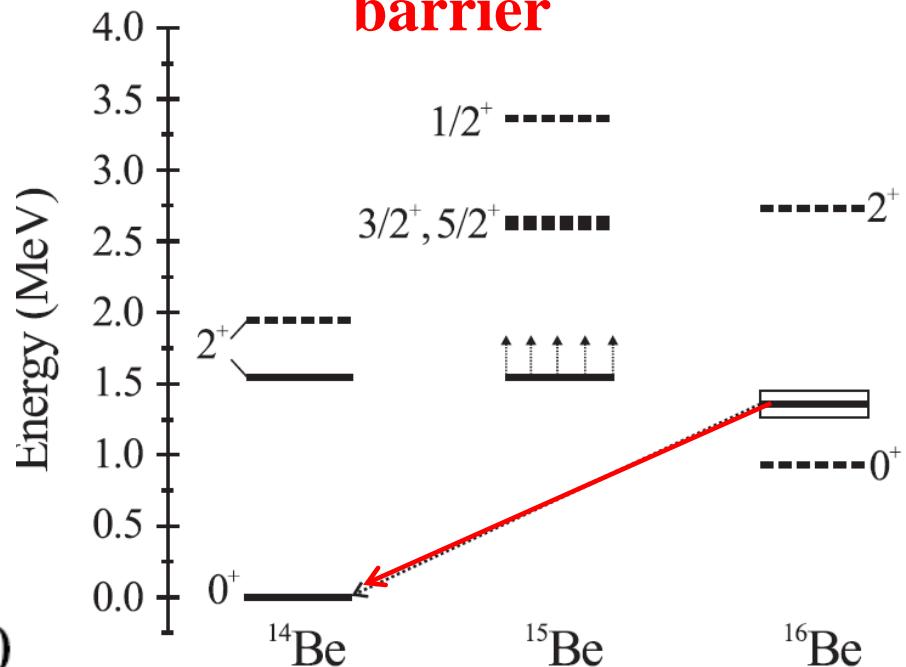
B. Blank and M. Ploszajczak, Rep. Prog. Phys. 71('08)046301

- ✓ probing correlations from energy and angle distributions of two emitted protons?
- ✓ Coulomb 3-body system
 - Theoretical treatment: difficult
 - how does FSI disturb the g.s. correlation?
 - diproton correlation: unclear in many systems
(theoretical calculations: not many)

2-neutron decay (MoNA@MSU)



3-body resonance due to the **centrifugal barrier**



A. Spyrou et al., PRL108('12) 102501

Other data:

^{13}Li (Z. Kohley et al., PRC87('13)011304(R))

^{26}O (E. Lunderbert et al., PRL108('12)142503)

$^{14}\text{Be} \rightarrow ^{13}\text{Li} \rightarrow ^{11}\text{Li} + 2\text{n}$

$^{27}\text{F} \rightarrow ^{26}\text{O} \rightarrow ^{24}\text{O} + 2\text{n}$

3-body model calculation with nn correlation: required

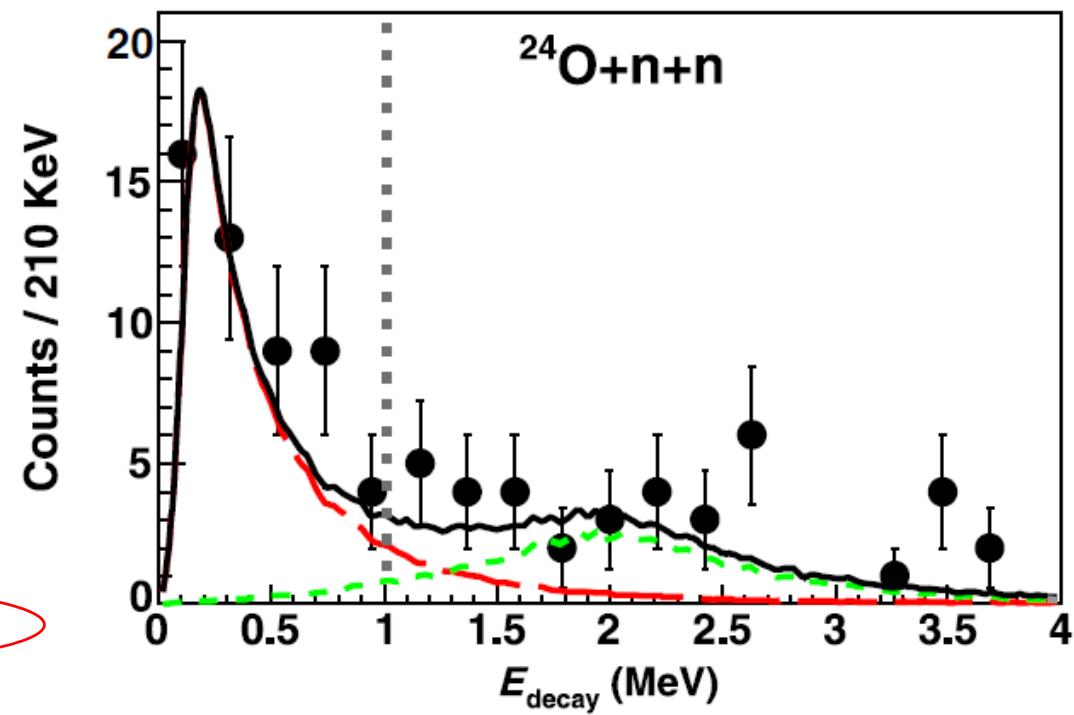
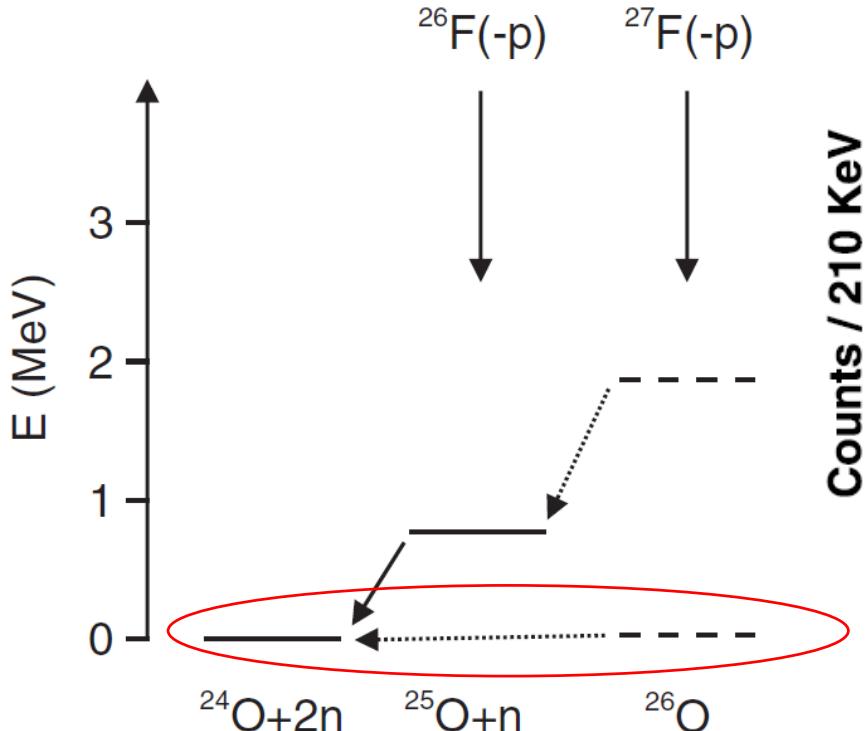
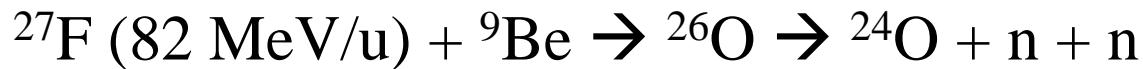
Two-neutron decay of ^{26}O

➤ the simplest among ^{16}Be , ^{13}Li , ^{26}O (MSU)

^{16}Be : deformation, ^{13}Li : treatment of ^{11}Li core

Experiment:

E. Lunderberg et al., PRL108 ('12) 142503
Z. Kohley et al., PRL 110 ('13) 152501



cf. C. Caesar et al., PRC88 ('13) 034313 (GSI exp.)

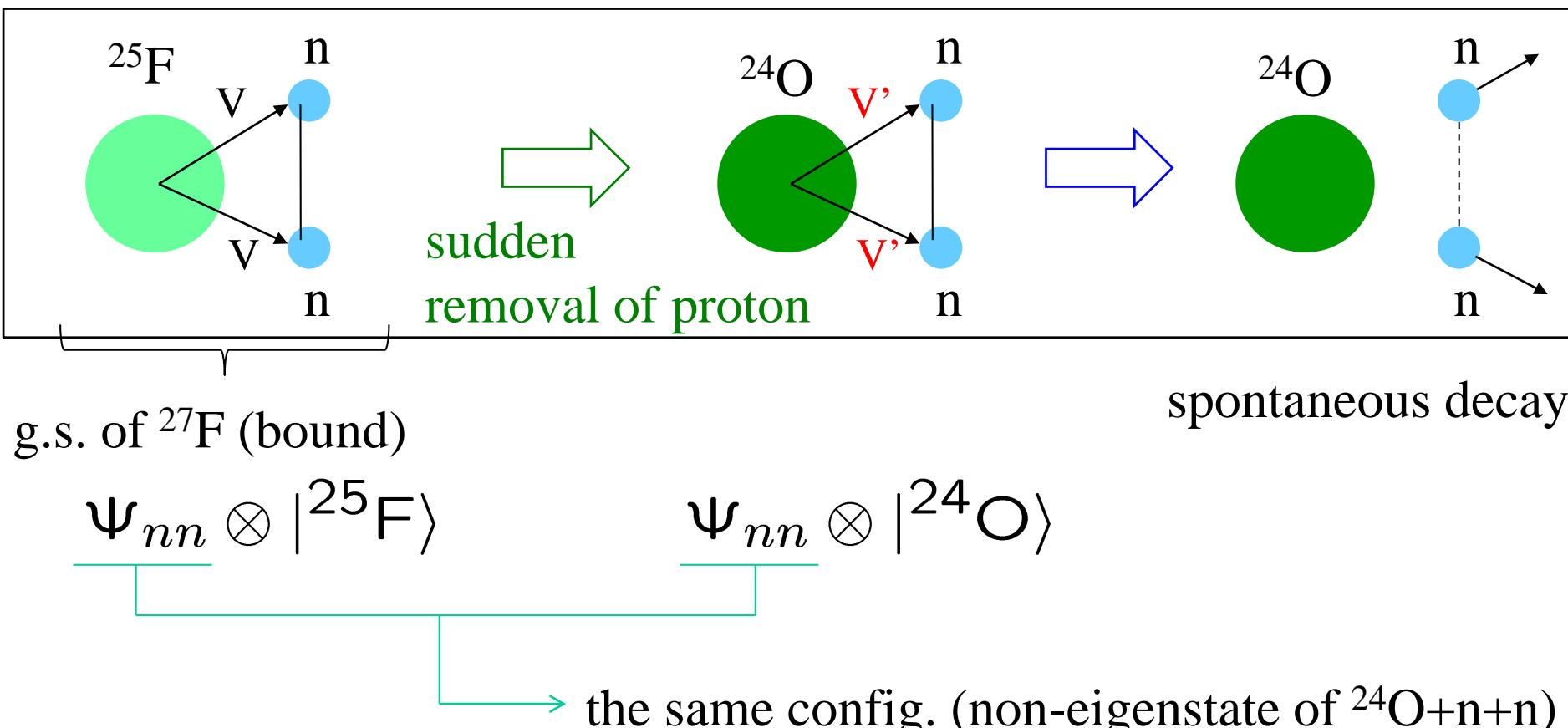
Y. Kondo et al., (SAMURAI)

$$E_{\text{decay}} = 150^{+50}_{-150} \text{ keV}$$

3-body model analysis for ^{26}O decay

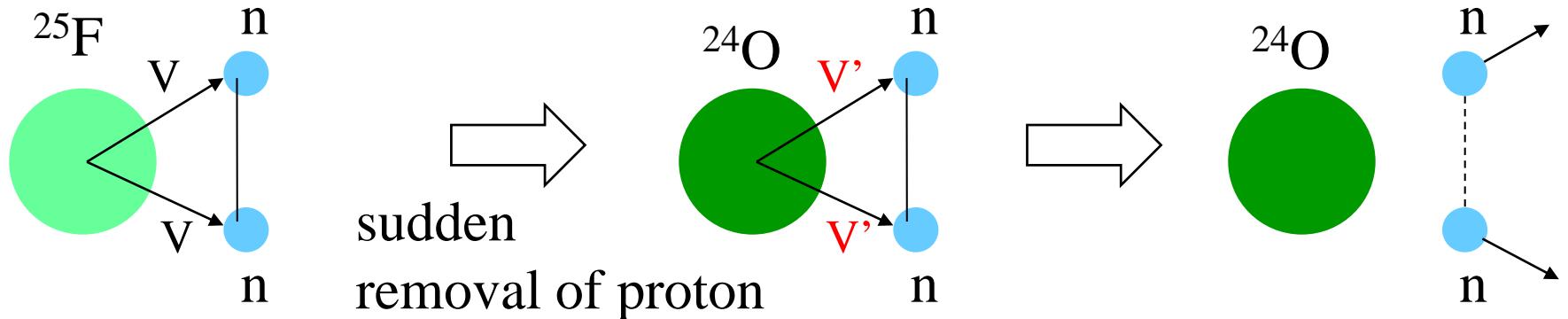
K.H. and H. Sagawa,
PRC89 ('14) 014331

cf. Expt. : ^{27}F (82 MeV/u) + $^9\text{Be} \rightarrow ^{26}\text{O} \rightarrow ^{24}\text{O} + n + n$



FSI → Green's function method ← continuum effects

$$\begin{aligned} M_{fi} &= \langle (j_1 j_2)^{J=0} | (1 - vG_0 + vG_0 vG_0 - \dots) | \Psi_i \rangle \\ &= \langle (j_1 j_2)^{J=0} | (1 + vG_0)^{-1} | \Psi_i \rangle \end{aligned}$$



➤ $^{24}\text{O} + \text{n}$ potential

Woods-Saxon potential

$$e_{2\text{s}1/2} = -4.09(13) \text{ MeV},$$

$$e_{1\text{d}3/2} = +770^{+20}_{-10} \text{ keV}, \quad \Gamma_{1\text{d}3/2} = 172(30) \text{ keV}$$

C.R. Hoffman et al.,
PRL100('08)152502

➤ $^{25}\text{F} + \text{n}$ potential

$(^{24}\text{O} + \text{n})$ potential + δV_{ls}

← pn tensor interaction

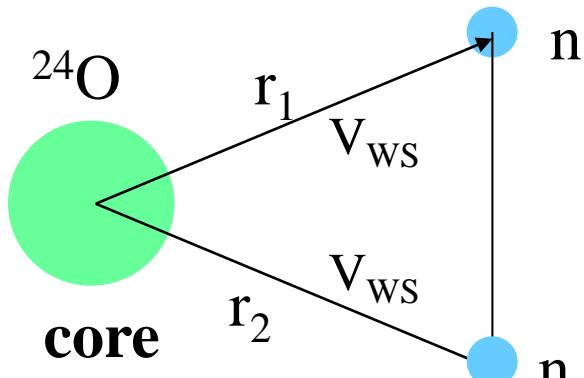
T. Otsuka et al., PRL95('05)232502

$$e_{1\text{d}3/2} (^{26}\text{F}) = -0.811 \text{ MeV}$$

➤ nn interaction (density-dependent zero-range interaction)

$$\leftarrow E_{\text{exp}} (^{27}\text{F}) = -2.80(18) \text{ MeV}$$

i) Decay energy spectrum



➤ $^{24}\text{O} + \text{n}$ potential

Woods-Saxon potential to reproduce

$$e_{2s1/2} = -4.09(13) \text{ MeV},$$

$$e_{1d3/2} = +770^{+20}_{-10} \text{ keV},$$

$$\Gamma_{1d3/2} = 172(30) \text{ keV}$$

➤ nn interaction

density-dep. contact interaction

$$E(^{27}\text{F}) = -2.69 \text{ MeV}$$

$$\frac{dP_I}{dE} = \sum_k |\langle \Psi_k^{(I)} | \Phi_{\text{ref}}^{(I)} \rangle|^2 \delta(E - E_k)$$

overlap with a ref.
state \leftarrow 2n config. with
 $^{25}\text{F} + \text{n} + \text{n}$

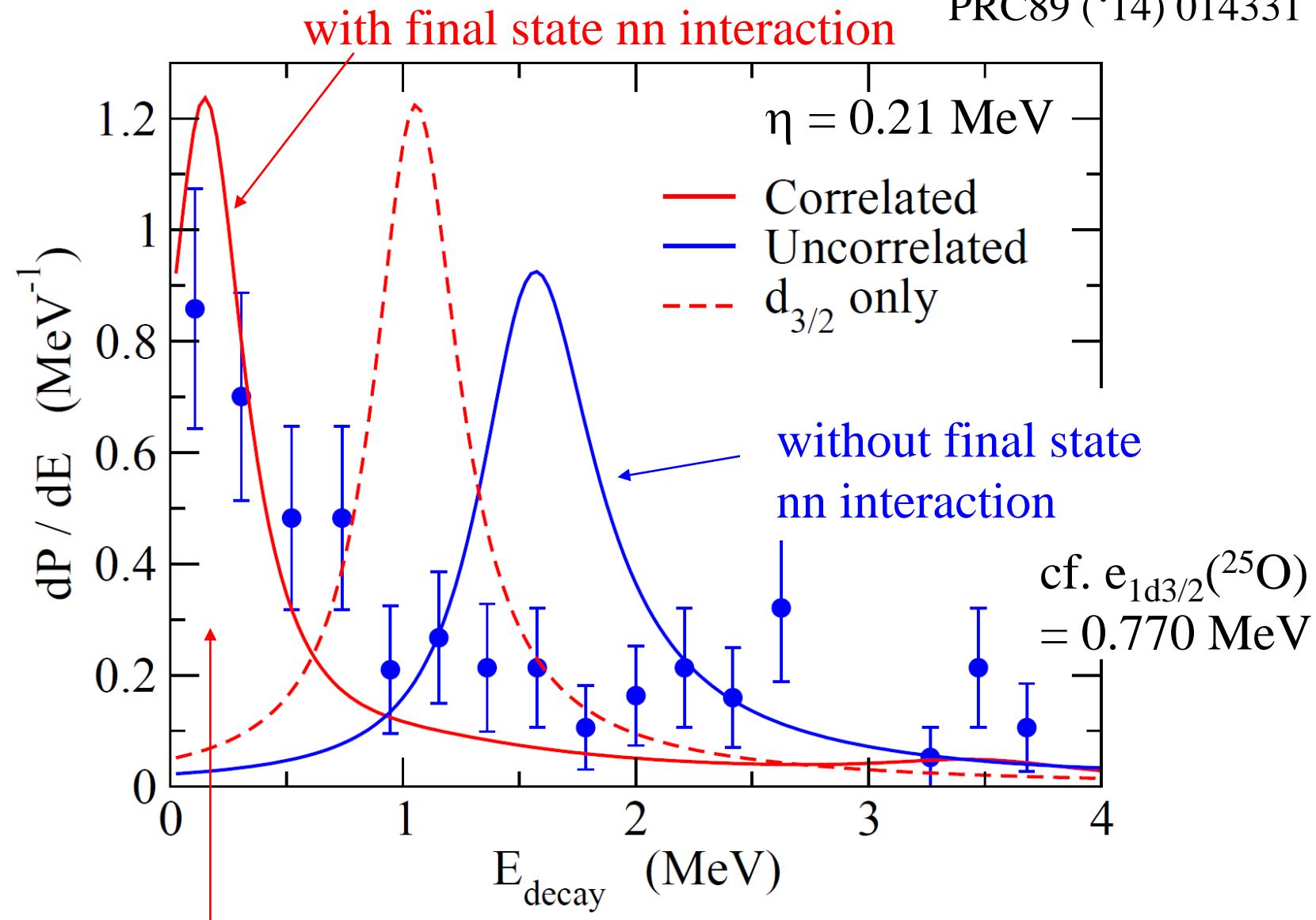
$$= -\frac{1}{\pi} \Im \langle \Phi_{\text{ref}}^{(I)} | G^{(I)}(E) | \Phi_{\text{ref}}^{(I)} \rangle,$$

$$G^{(I)}(E) = G_0^{(I)}(E) - G_0^{(I)}(E)v(1 + G_0^{(I)}(E)v)^{-1}G_0^{(I)}(E)$$

\leftarrow continuum effects

i) Decay energy spectrum

K.H. and H. Sagawa,
PRC89 ('14) 014331

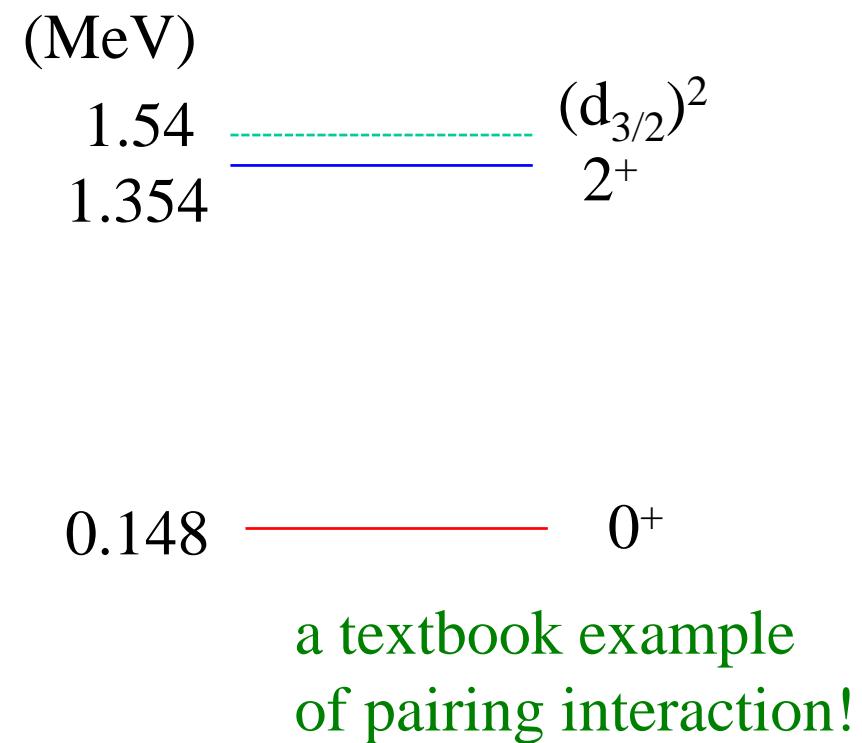
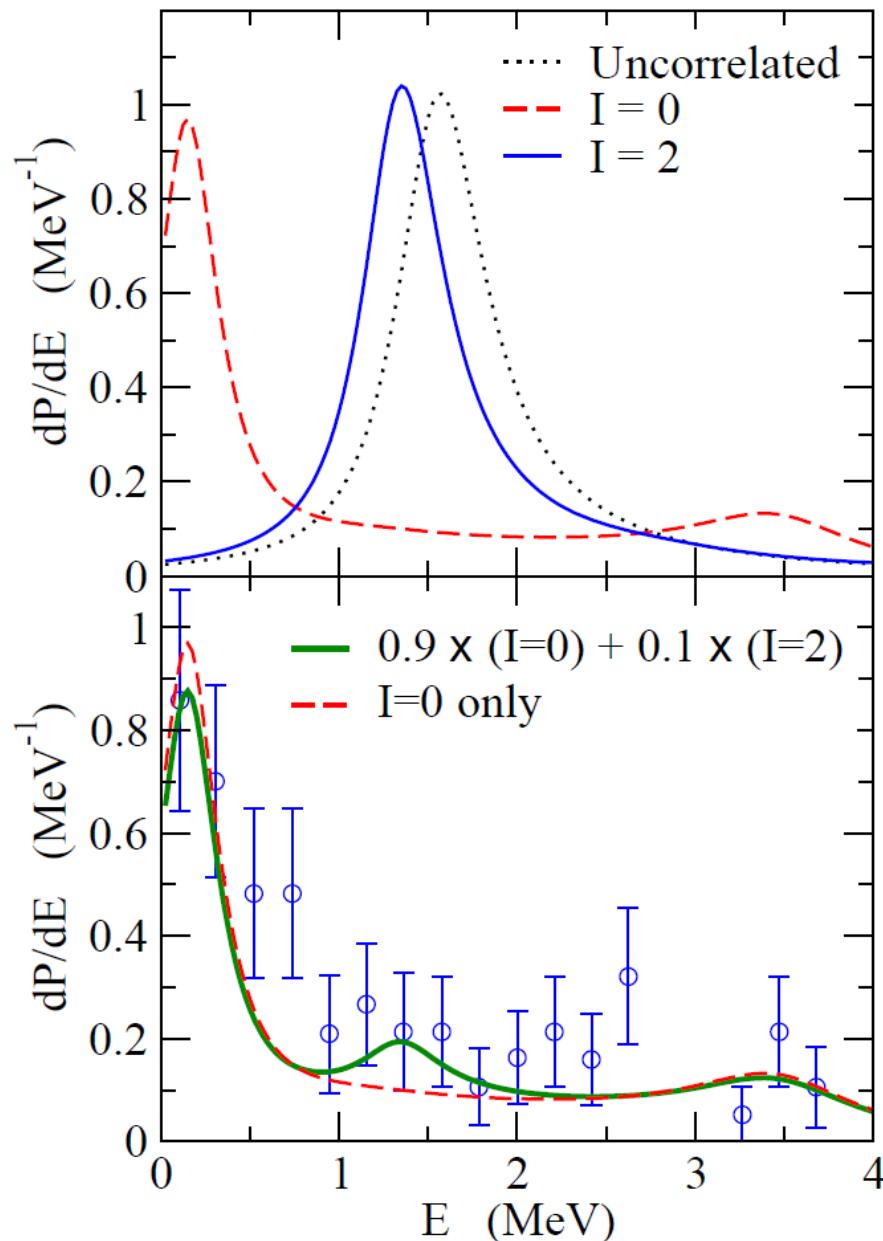


very narrow three-body resonance state ($\Gamma_{\text{exp}} \sim 10^{-10} \text{ MeV}$)

$E_{\text{peak}} = 0.14 \text{ MeV}$ with this setup for the Hamiltonian

2^+ state of ^{26}O

Kondo et al. : a prominent second peak at $E \sim 1.3$ MeV



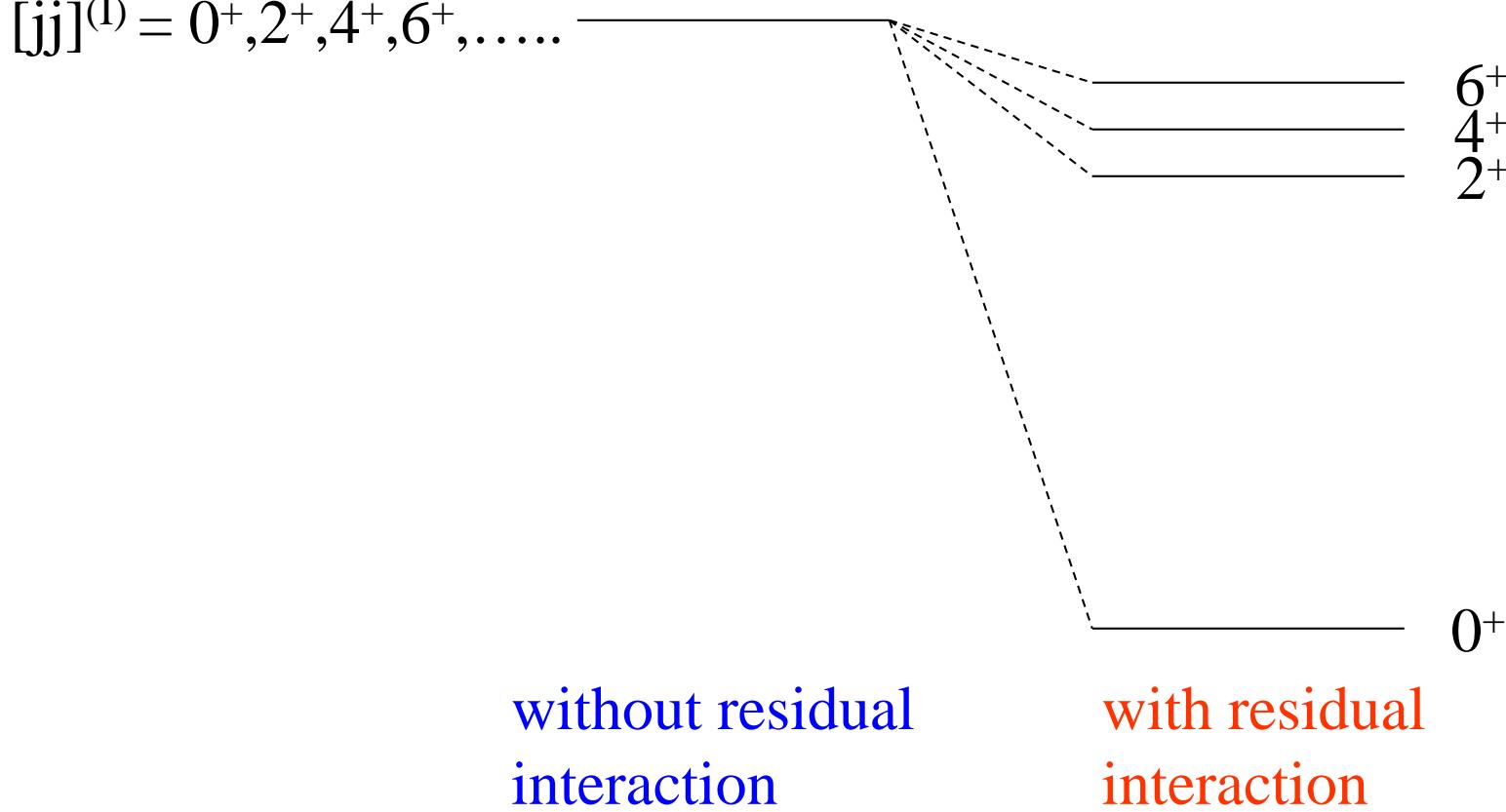
cf. another set of parameters:

$$E(0^+) = 5 \text{ keV}$$

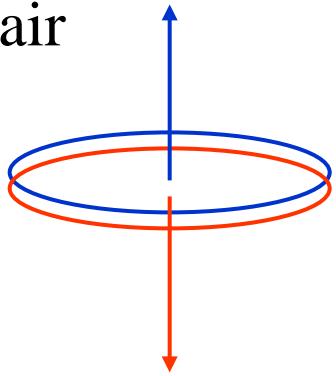
$$E(2^+) = 1.338 \text{ MeV}$$

K.H. and H. Sagawa,
PRC90('14)027303

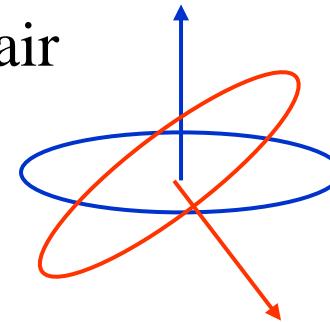
$[jj]^{(I)} = 0^+, 2^+, 4^+, 6^+, \dots$



$I=0$ pair

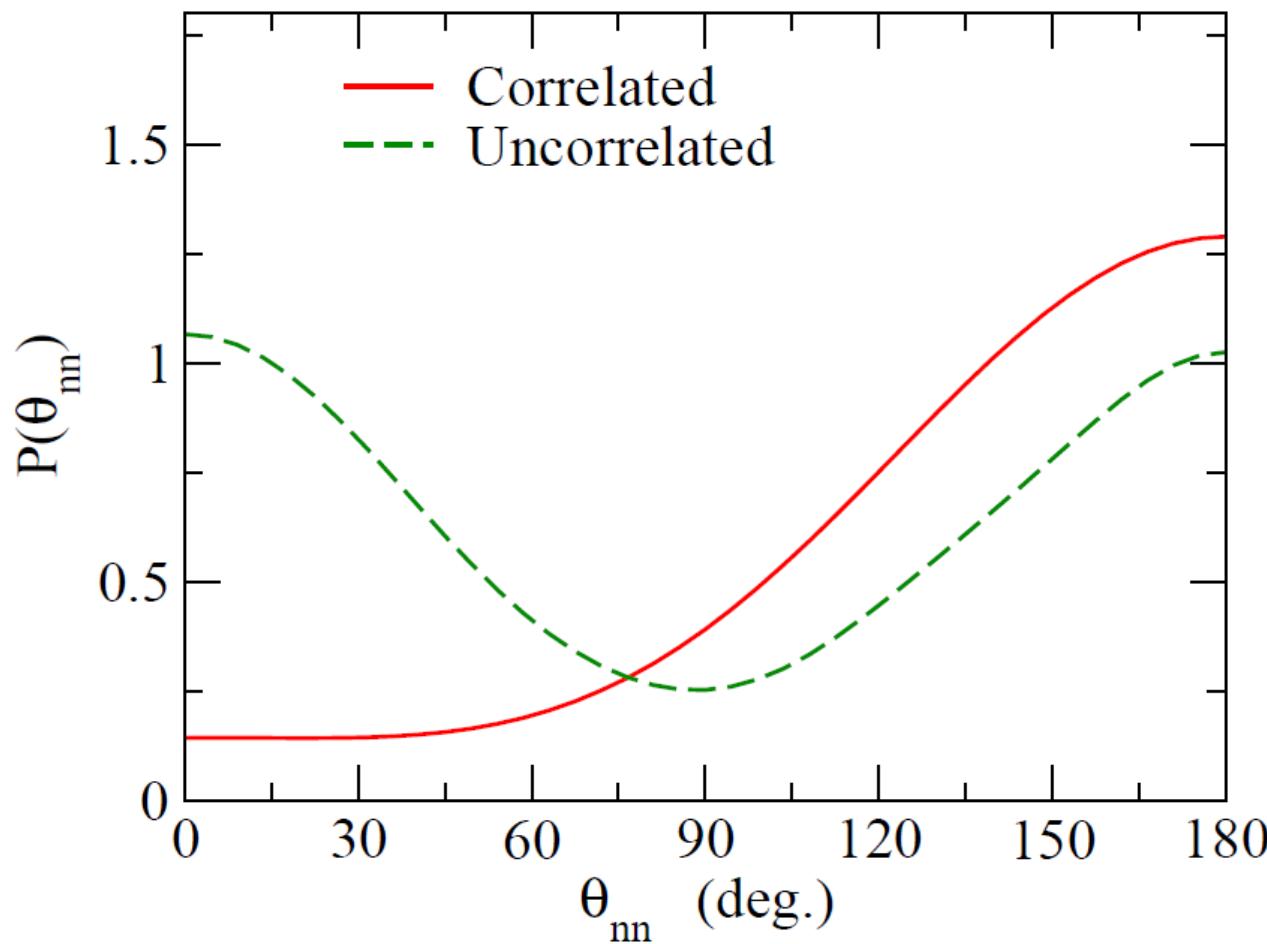


$I \neq 0$ pair



ii) angular correlations of the emitted neutrons

K.H. and H. Sagawa,
PRC89 ('14) 014331

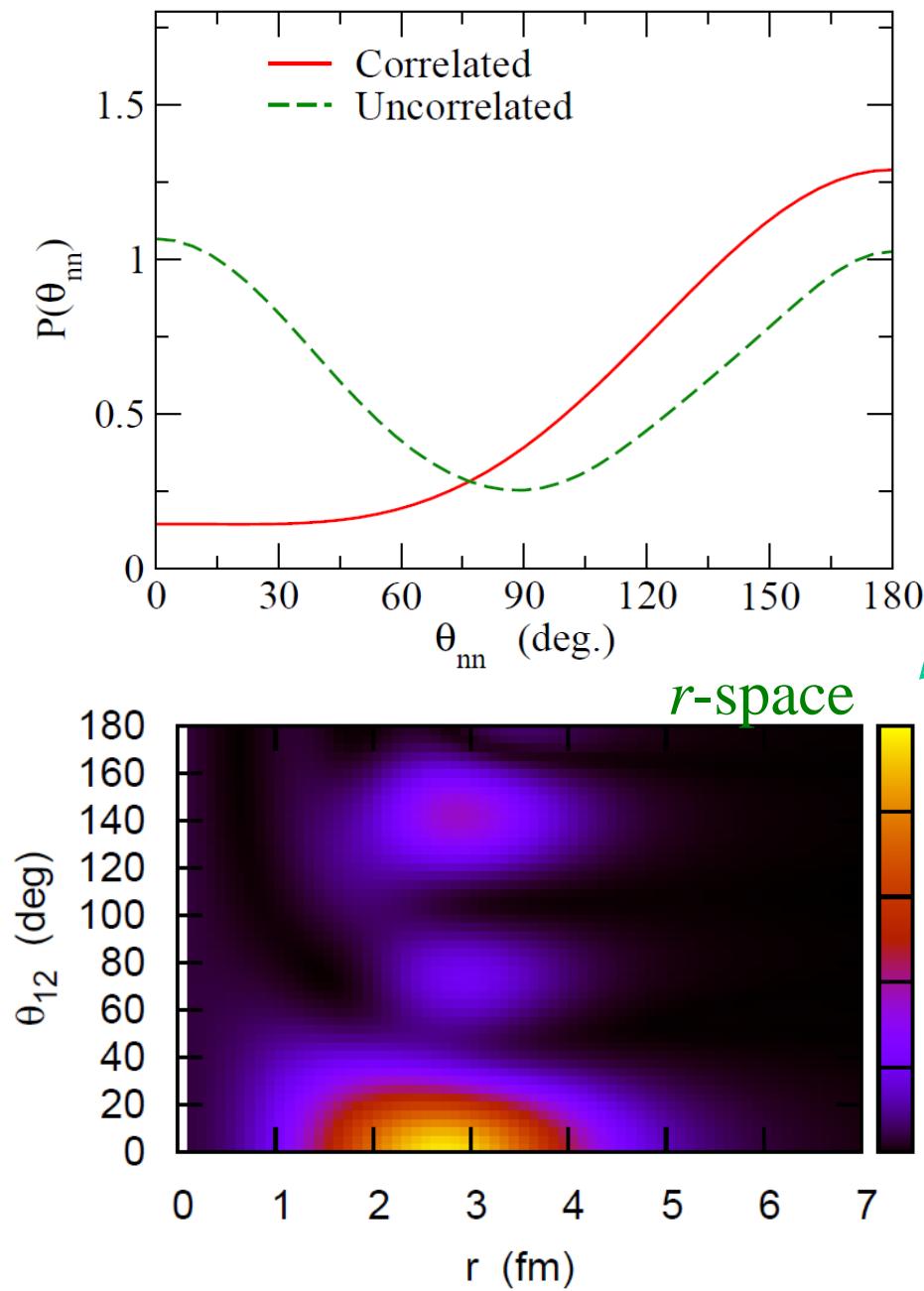


$$\langle \theta_{nn} \rangle = 115.3^\circ$$

correlation → enhancement of back-to-back emissions

cf. Similar conclusion: L.V. Grigorenko, I.G. Mukha, and M.V. Zhukov,
PRL 111 (2013) 042501

ii) distribution of opening angle for two-emitted neutrons

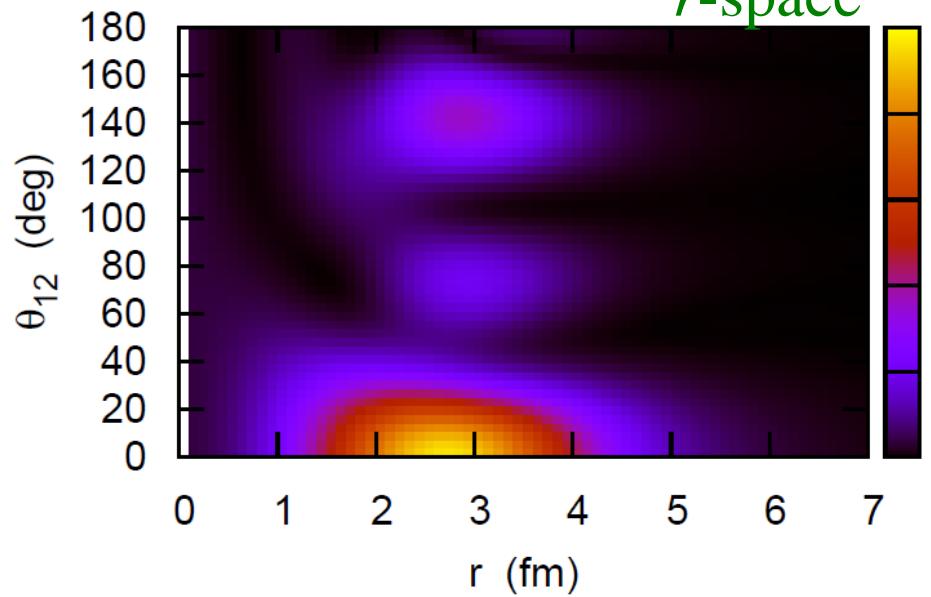


density of the resonance state (with the box b.c.)

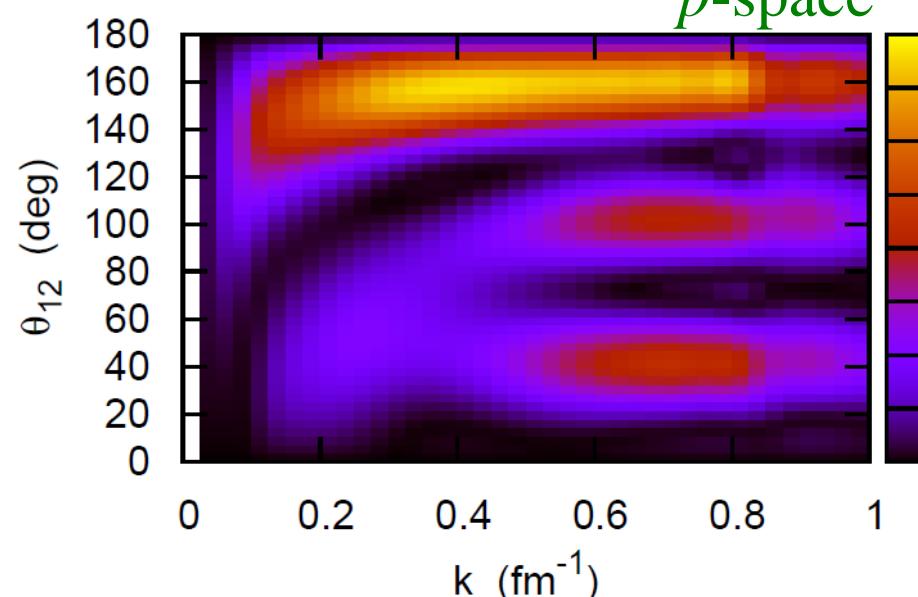
$$\rho(r, r, \theta)$$

$$8\pi^2 k^4 \sin \theta \cdot \rho(k, k, \theta)$$

r-space

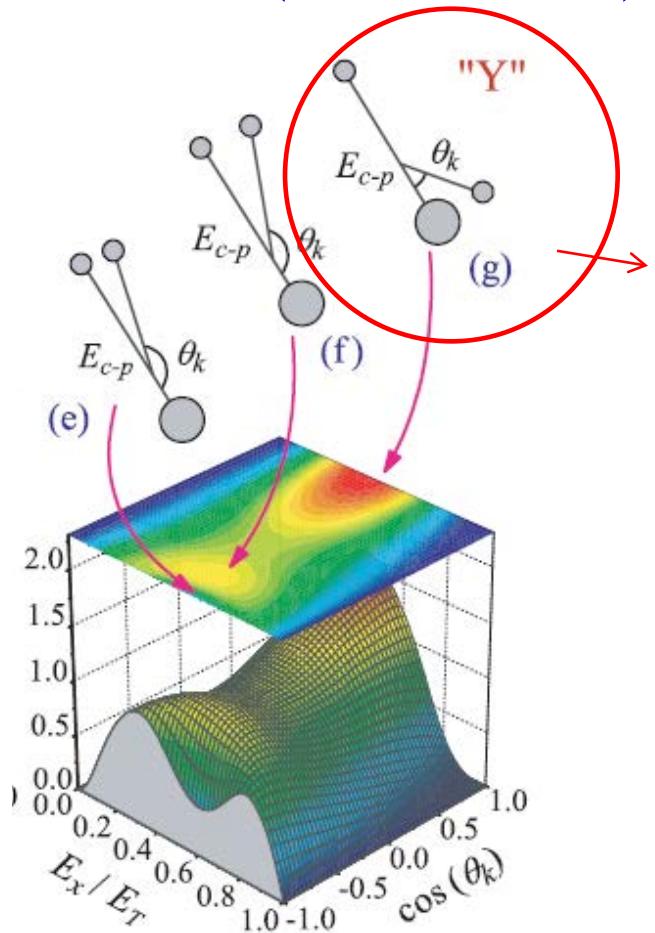


p-space



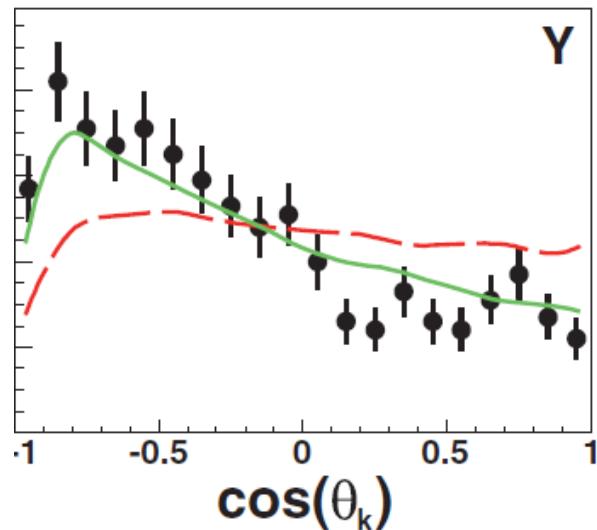
➤ Discussions: back-to-back? or forward angles?

two-proton decay
from ${}^6\text{Be}$ (back-to-back)



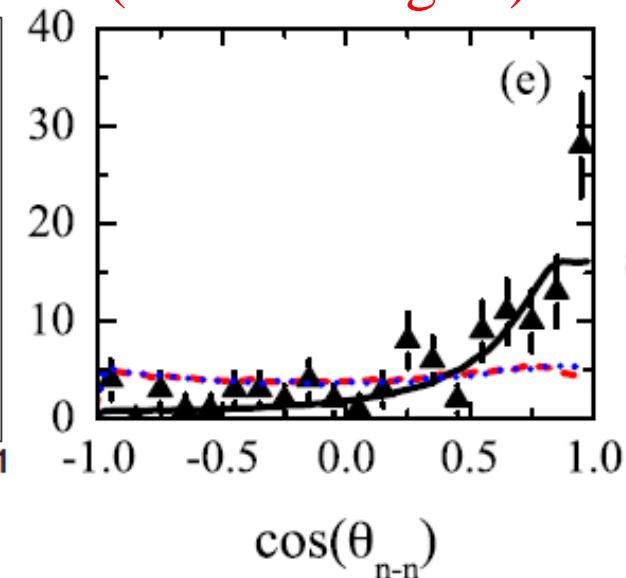
L.V. Grigorenko et al.,
PRC80 ('09) 034602

2n decay of ${}^{13}\text{Li}$
(forward angles)



Z. Kohley et al.,
PRC87('13)011304(R)

2n decay of ${}^{16}\text{Be}$
(forward angles)



A. Spyrou et al.,
PRL108('12) 102501

- ✓ Q-value effect? (cf. nuclear phase shifts)
- ✓ core excitations?



open problem

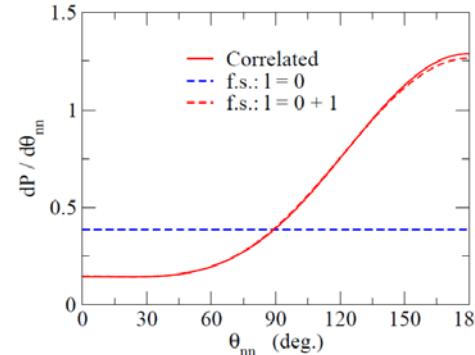
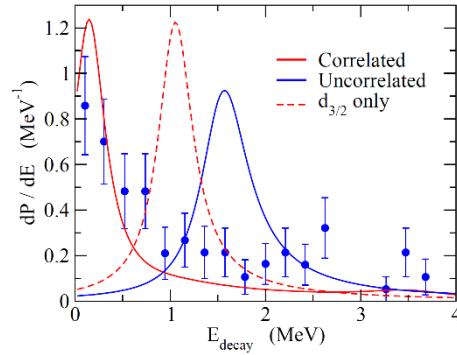
Summary

three-body model with
density-dependent zero-range interaction:
continuum calculations: relatively easy

- Coulomb breakup of Borromean nuclei
- 2n emission decay of ^{26}O

- ✓ Decay energy spectrum: strong low-energy peak
- ✓ 2^+ energy
- ✓ Angular distributions: enhanced back-to-back emission

↔
dineutron emission



□ open problems

- ✓ Analyses for ^{16}Be , ^{13}Li (especially angular distributions)
- ✓ Decay width?