

Heavy-ion fusion reactions and superheavy elements

Kouichi Hagino
Tohoku University, Sendai, Japan



1. H.I. fusion reactions: why are they interesting?
2. Coupled-channels approach
3. Future perspectives: superheavy elements

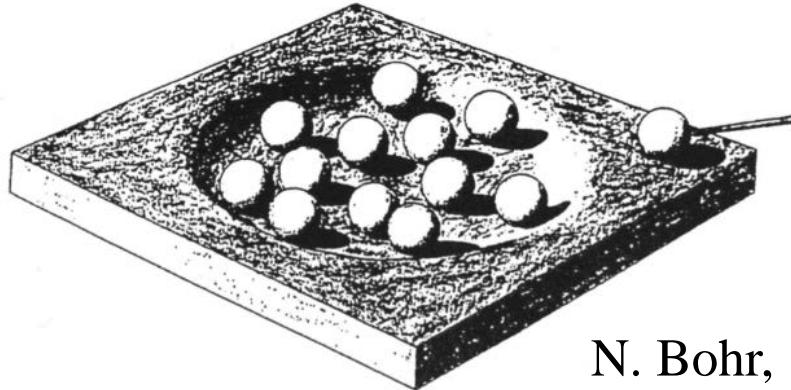
Recent review article:

K. Hagino and N. Takigawa, Prog. Theo. Phys. 128 ('12)1061.

Fusion reactions: compound nucleus formation

Niels Bohr (1936)

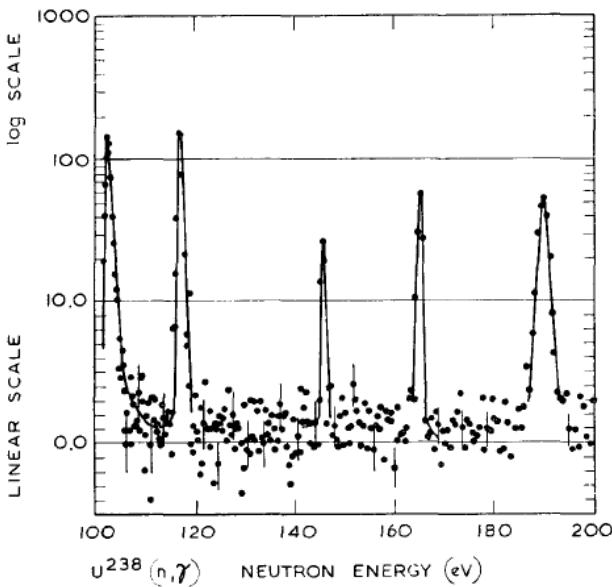
Neutron capture of nuclei → compound nucleus



N. Bohr,
Nature 137 ('36) 351



Wikipedia



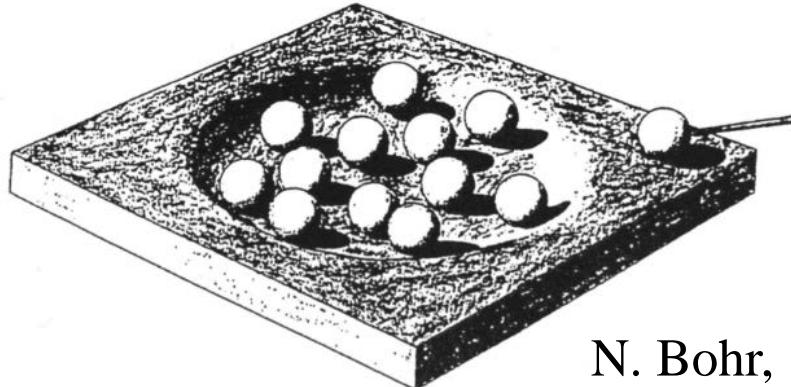
cf. Experiment of Enrico Fermi (1935)
many very narrow (=long life-time)
resonances (width ~ eV)

M. Asghar et al., Nucl. Phys. 85 ('66) 305

Fusion reactions: compound nucleus formation

Niels Bohr (1936)

Neutron capture of nuclei → compound nucleus



N. Bohr,
Nature 137 ('36) 351

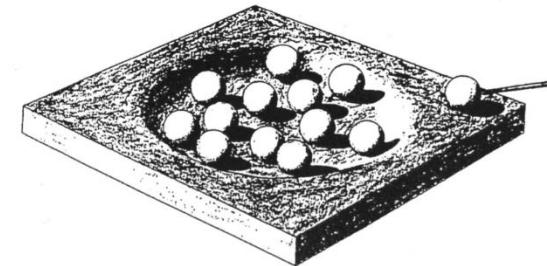
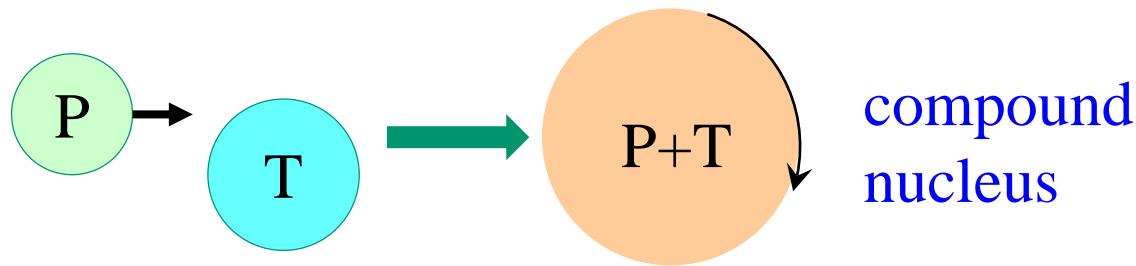


Wikipedia

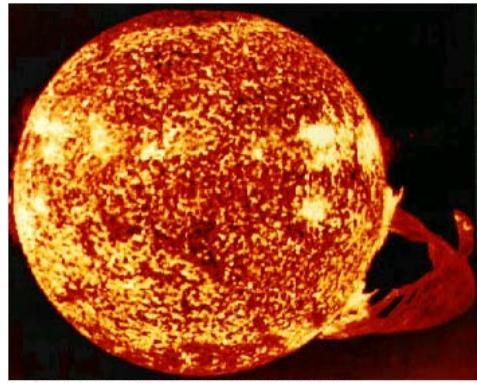
forming a compound nucleus with heavy-ion reactions = H.I. fusion



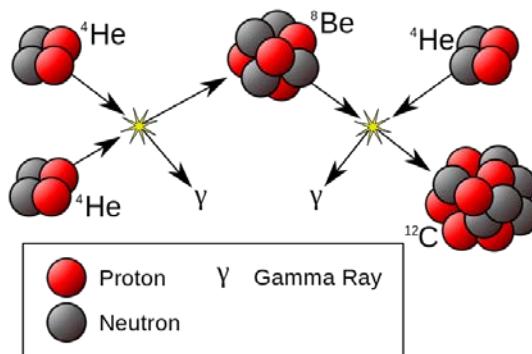
Fusion reactions: compound nucleus formation



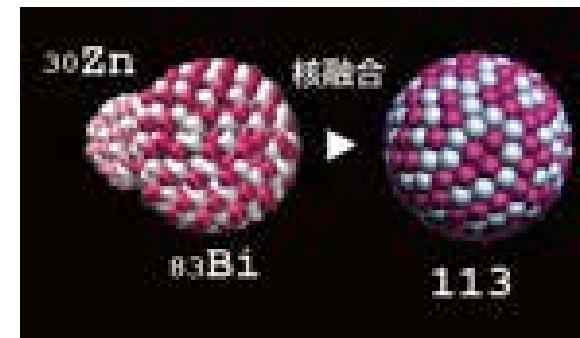
cf. Bohr '36



energy production
in stars (Bethe '39)



nucleosynthesis

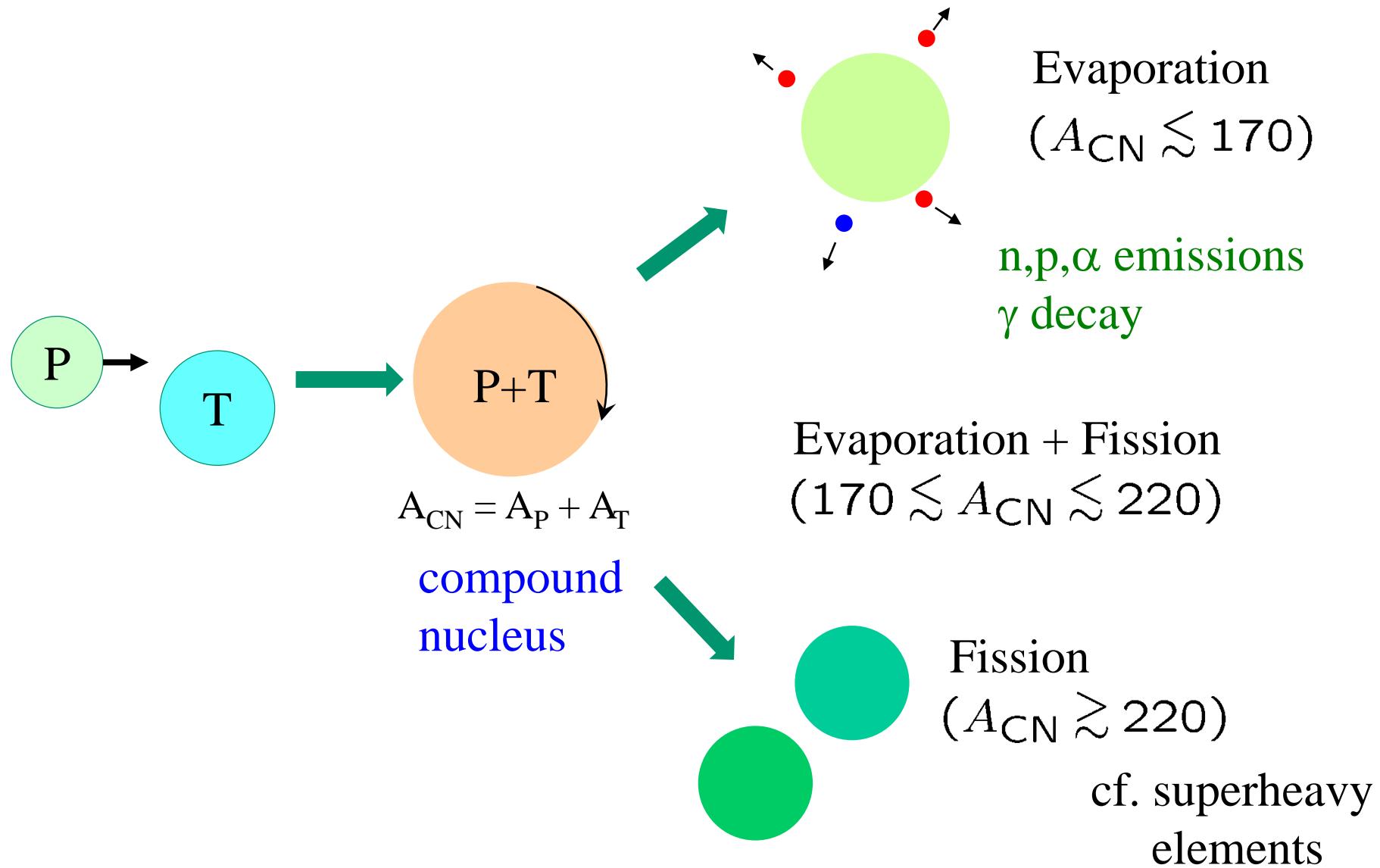


superheavy elements

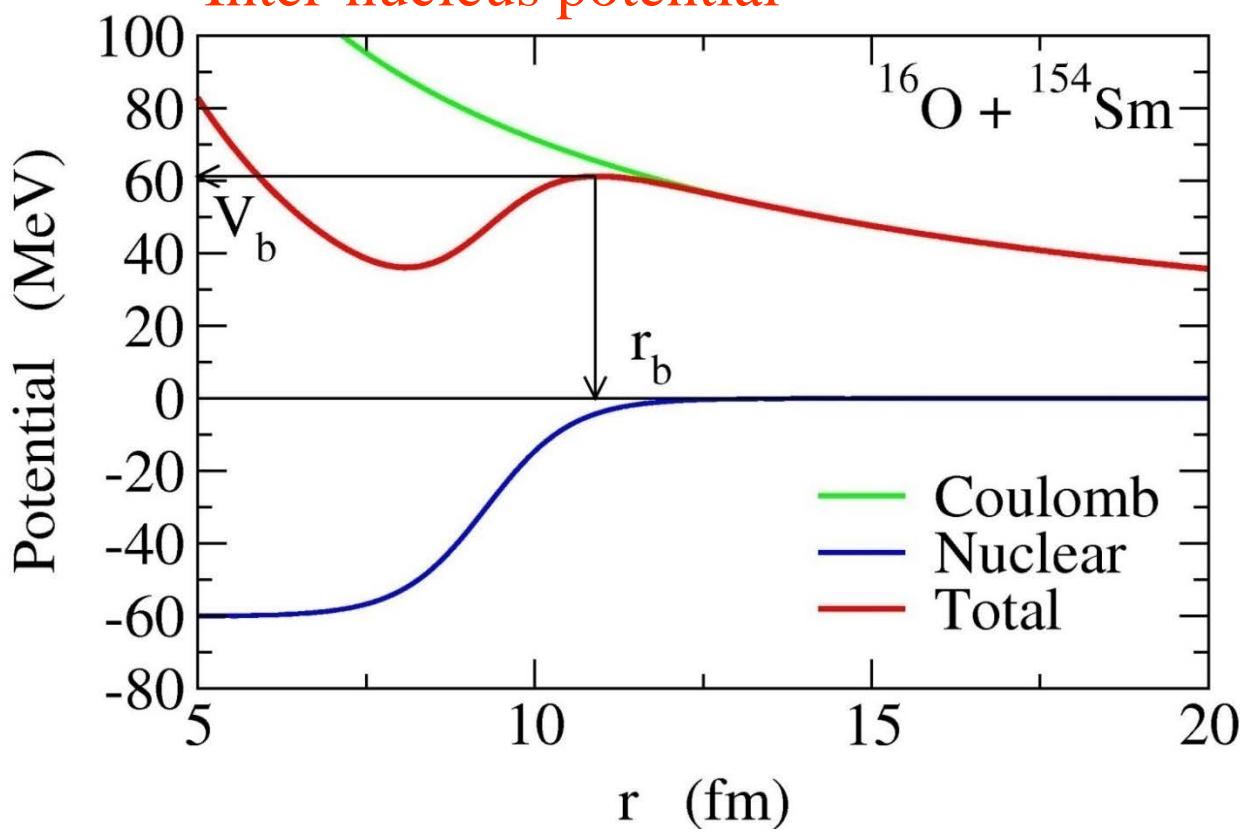
Fusion and fission: large amplitude motions of quantum many-body systems with strong interaction

← microscopic understanding: an ultimate goal of nuclear physics

Fusion reactions: compound nucleus formation



Inter-nucleus potential



Two interactions:

1. Coulomb force
long range repulsion
2. Nuclear force
short range attraction

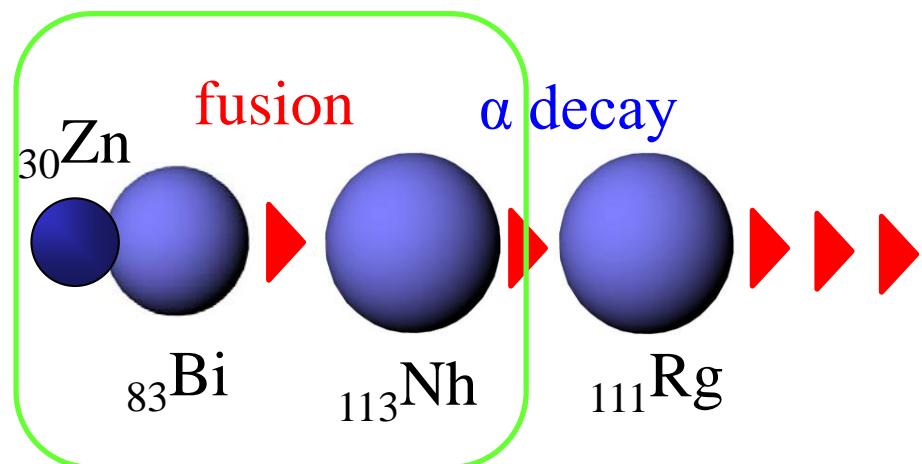
potential barrier
due to a cancellation
between the two
(Coulomb barrier)

- Above-barrier energies
- • Sub-barrier energies
(energies around the Coulomb barrier)
- Deep sub-barrier energies

Why sub-barrier fusion?

two obvious reasons:

113 Nh nihonium	115 Mc moscovium
117 Ts tennessine	118 Og oganesson



superheavy elements

cf. ^{209}Bi ($^{70}\text{Zn}, \text{n}$) ^{278}Nh

$$V_B \sim 260 \text{ MeV}$$

$$E_{\text{cm}}^{\text{(exp)}} \sim 262 \text{ MeV}$$

Why sub-barrier fusion?

two obvious reasons:

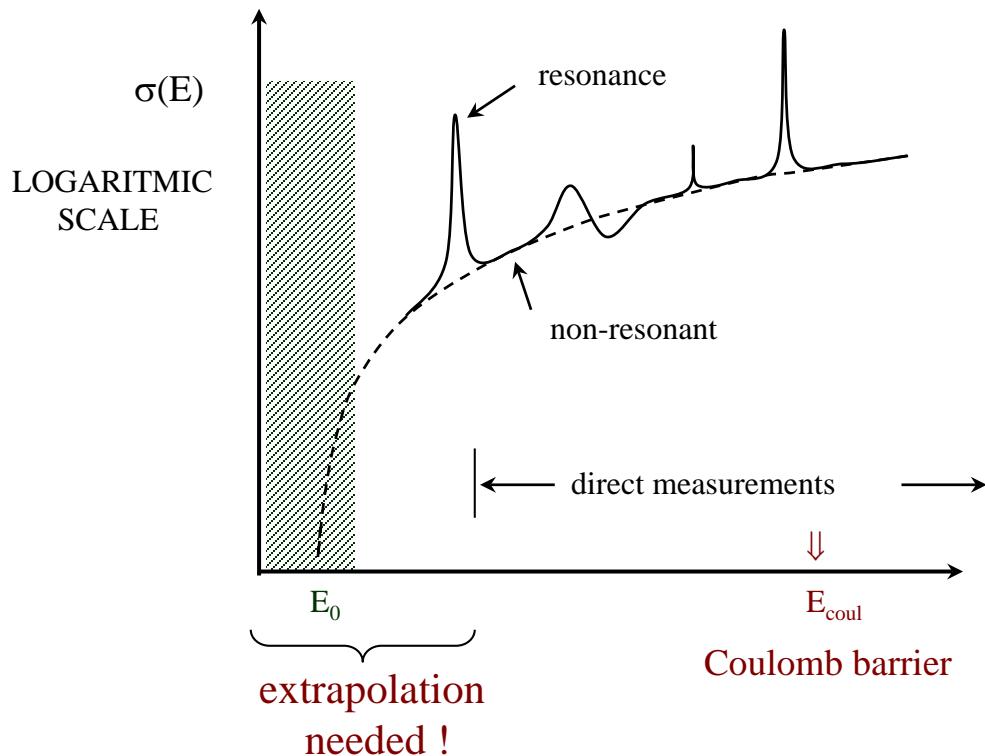
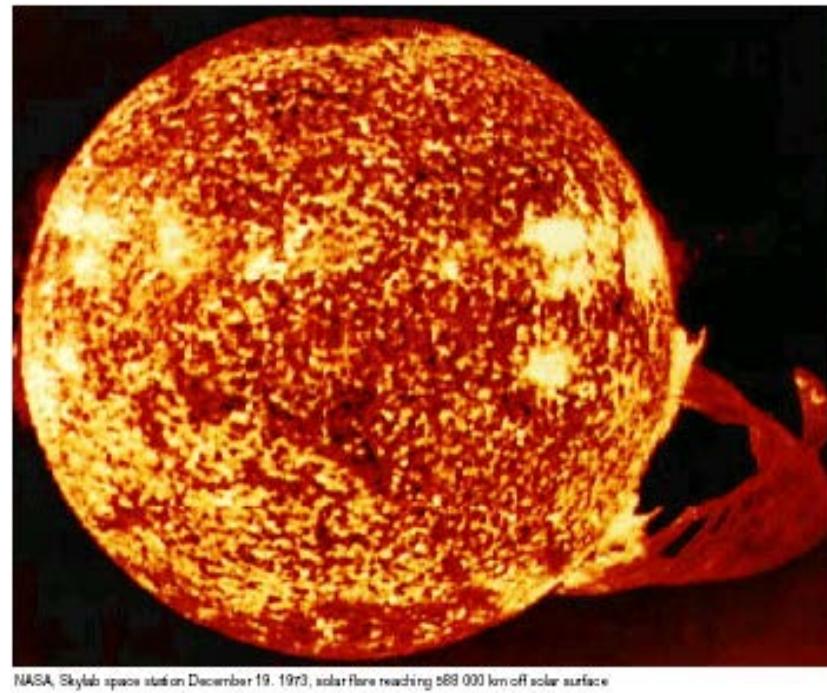


figure: M. Aliotta



nuclear astrophysics
(nuclear fusion in stars)

cf. extrapolation of data

Why sub-barrier fusion?

two obvious reasons:

- ✓ superheavy elements
- ✓ nuclear astrophysics

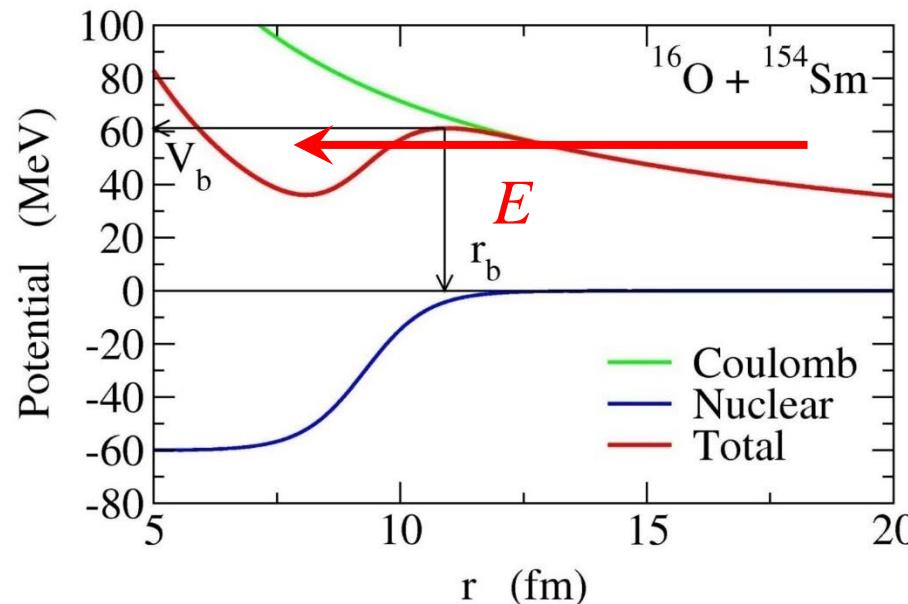
other reasons:

- ✓ reaction dynamics

strong interplay between reaction and structure

cf. high E reactions: much simpler reaction mechanisms

- ✓ many-particle tunneling



Why sub-barrier fusion?

two obvious reasons:

- ✓ superheavy elements
- ✓ nuclear astrophysics

other reasons:

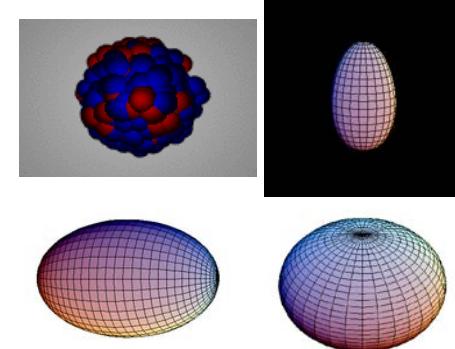
- ✓ reaction dynamics

strong interplay between reaction and structure

cf. high E reactions: much simpler reaction mechanisms

- ✓ many-particle tunneling

- many types of intrinsic degrees of freedom
(several types of collective vibrations,
deformation with several multipolarities)
- energy dependence of tunneling probability
cf. alpha decay: fixed energy



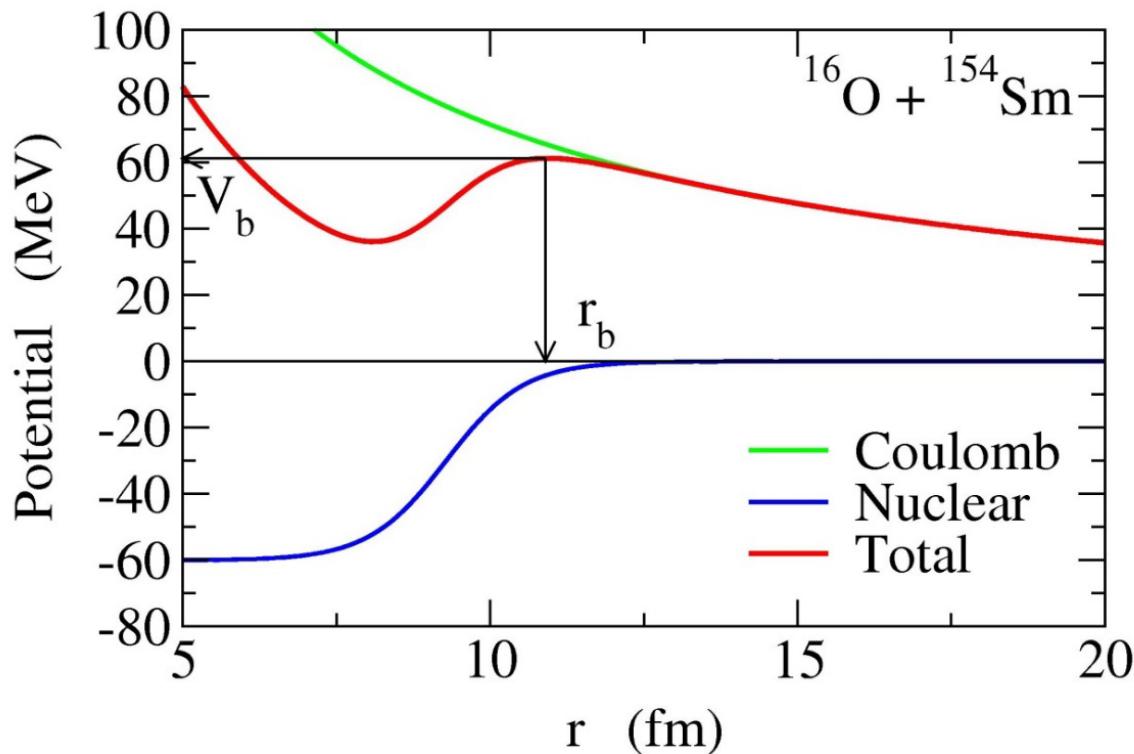
H.I. fusion reaction = an ideal playground to study quantum
tunneling with many degrees of freedom

The simplest approach to fusion: potential model

Potential model: $V(r)$ + absorption

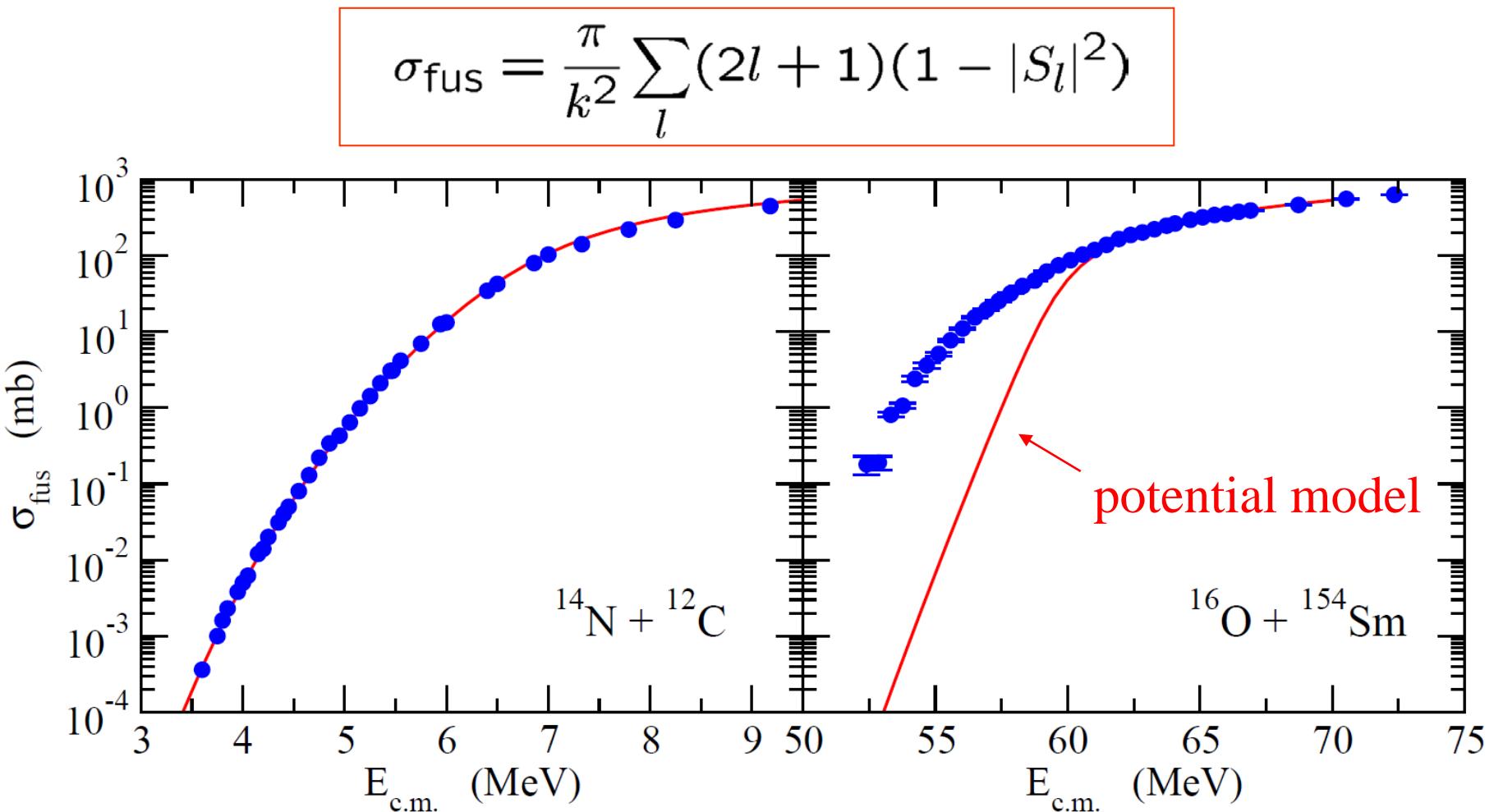
$$\sigma_{\text{fus}}(E) = \frac{\pi}{k^2} \sum_l (2l + 1) P_l(E)$$

$P_l(E)$: barrier penetrability



Comparison with experimental data: large enhancement of σ_{fus}

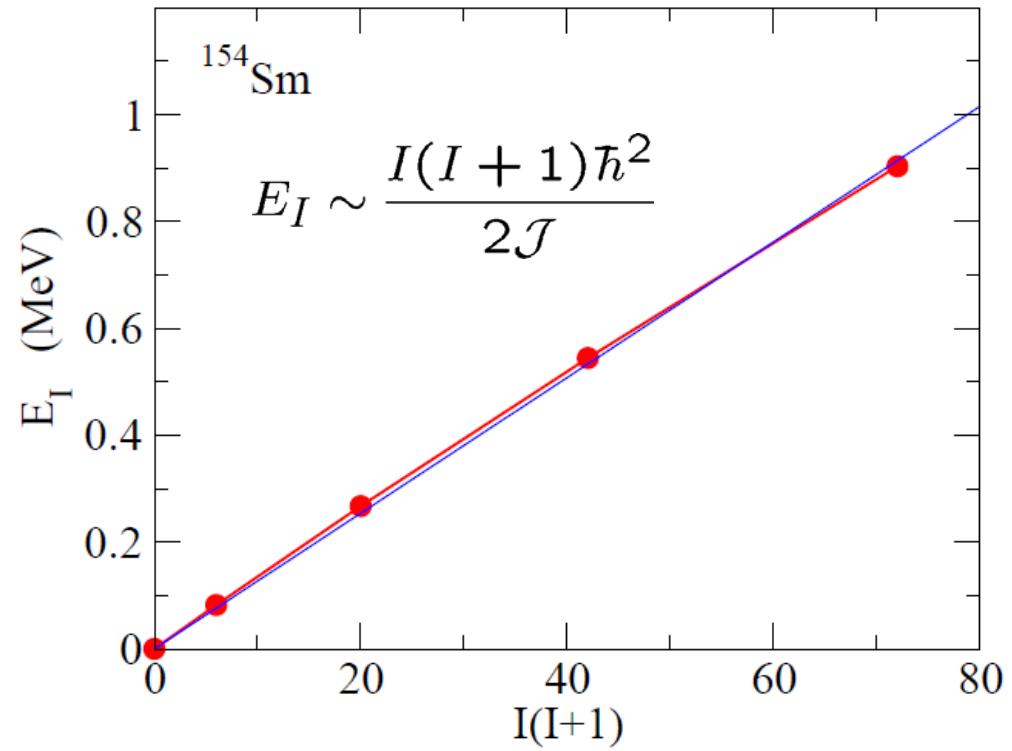
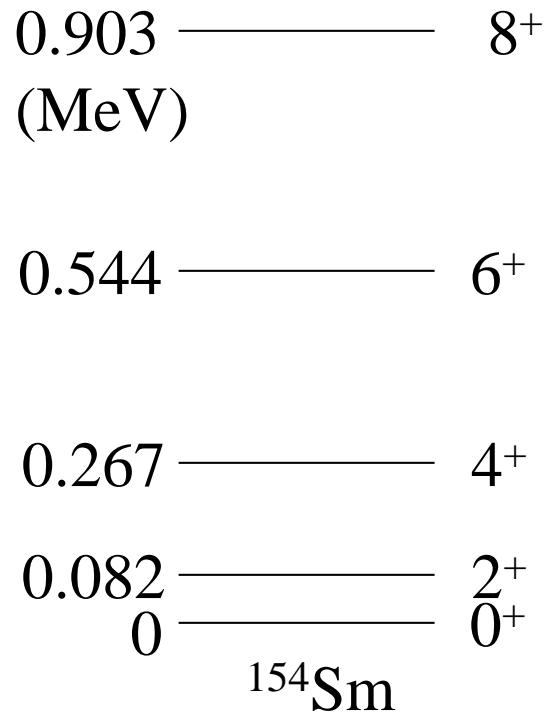
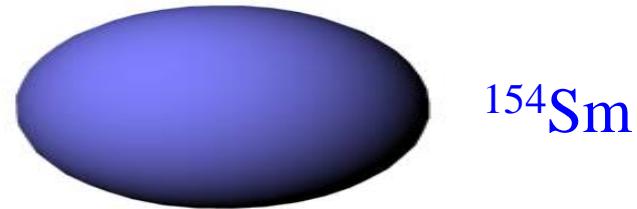
Potential model: $V(r)$ + absorption



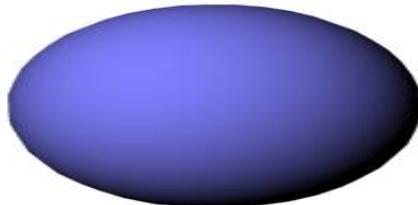
cf. seminal work:

R.G. Stokstad et al., PRL41('78) 465

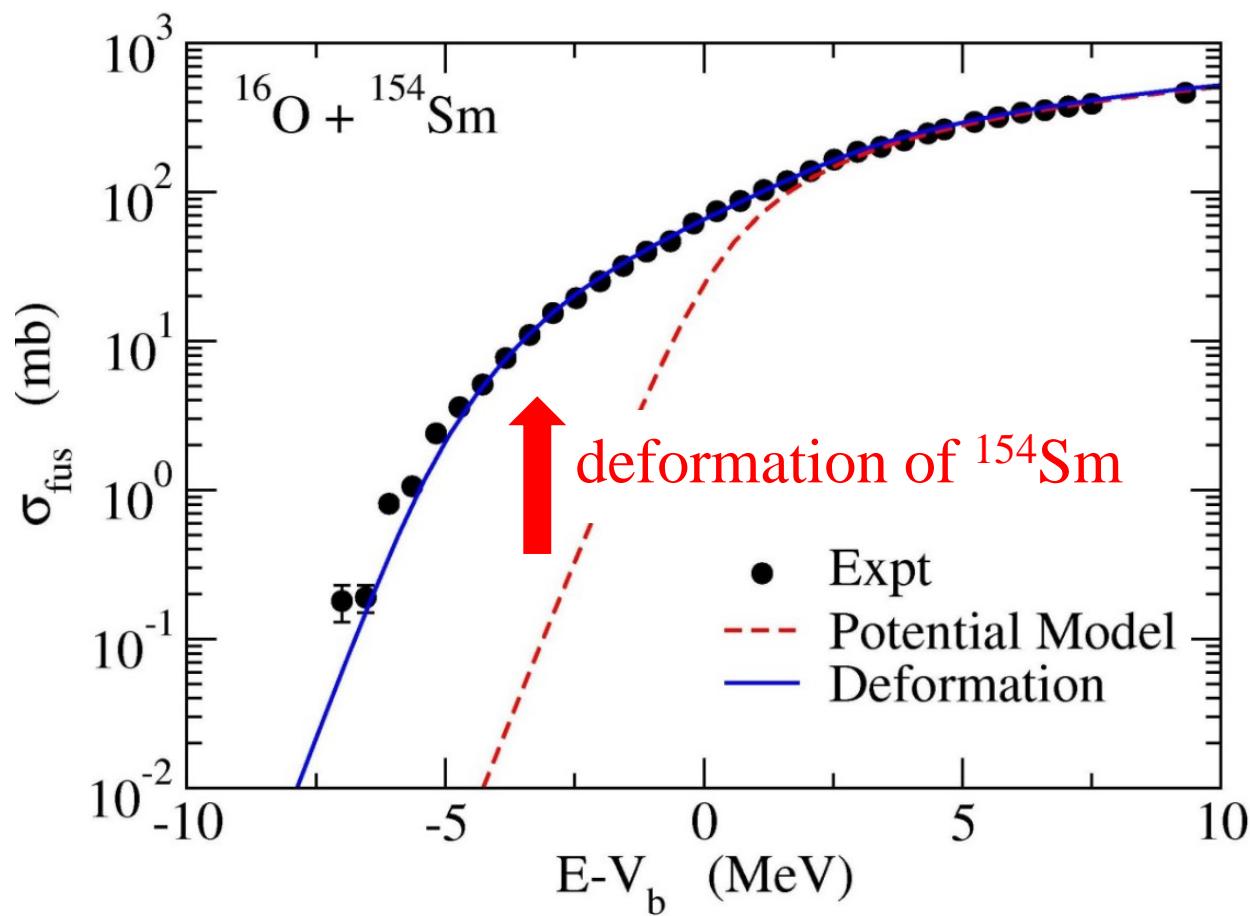
^{154}Sm : a typical deformed nucleus



^{154}Sm : a typical deformed nucleus

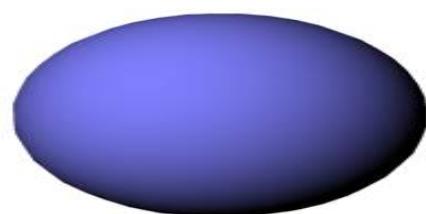


^{154}Sm

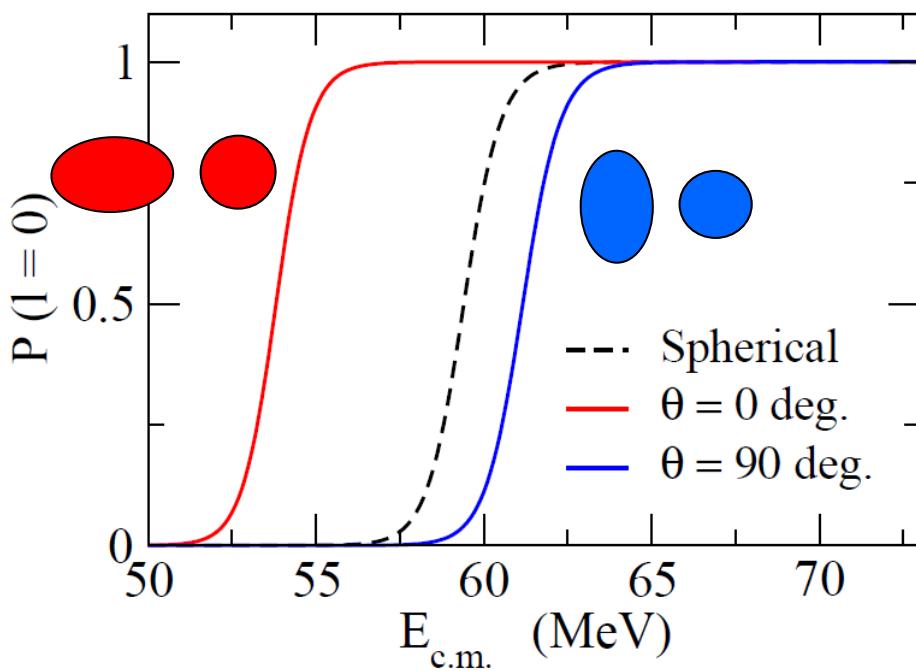
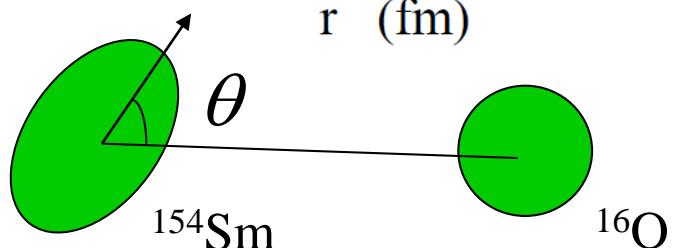
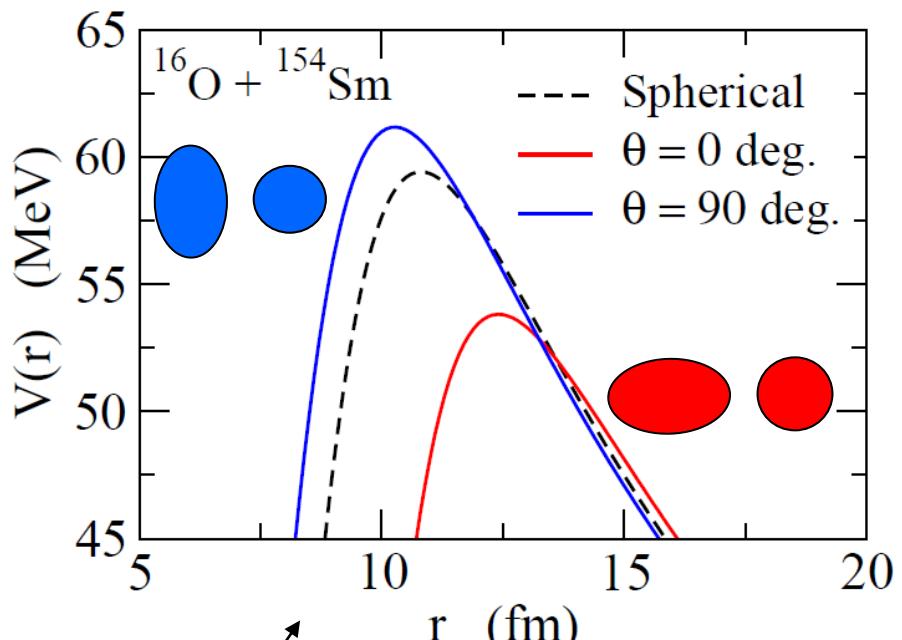


Effects of nuclear deformation

^{154}Sm : a typical deformed nucleus

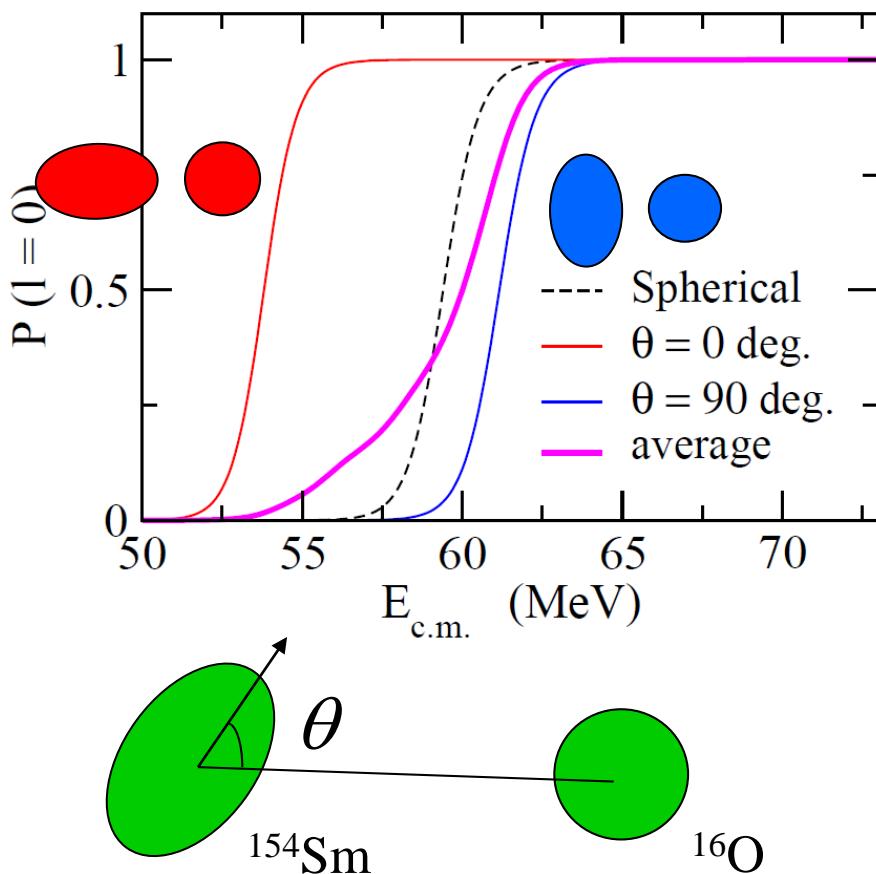


^{154}Sm

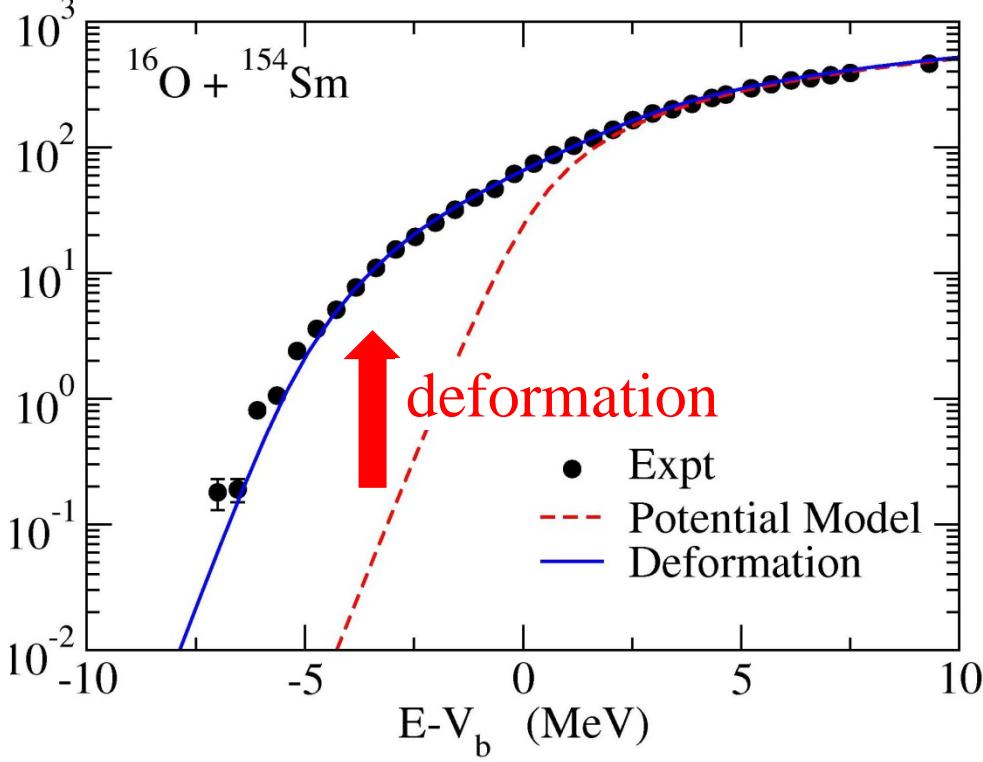


Effects of nuclear deformation

^{154}Sm : a typical deformed nucleus

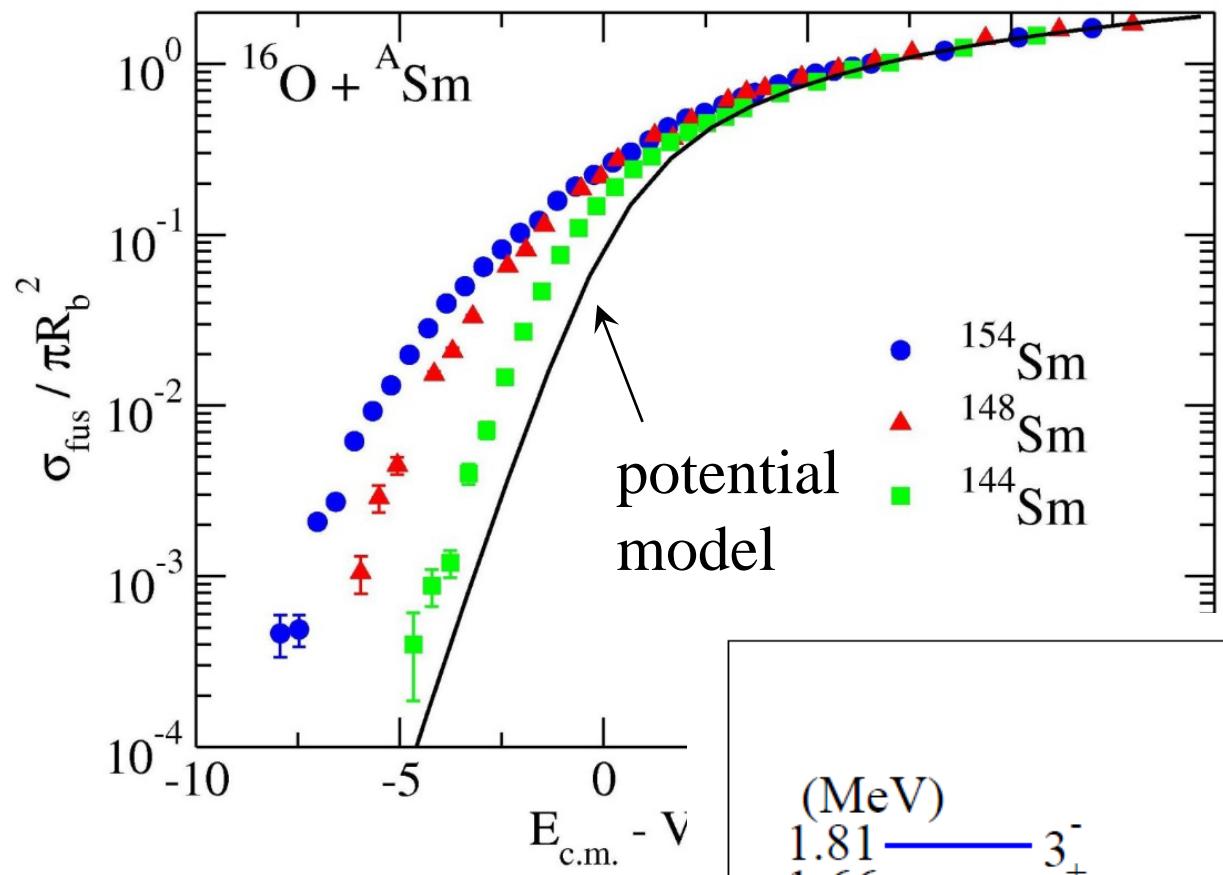


$$\sigma_{\text{fus}}(E) = \int_0^1 d(\cos \theta) \sigma_{\text{fus}}(E; \theta)$$



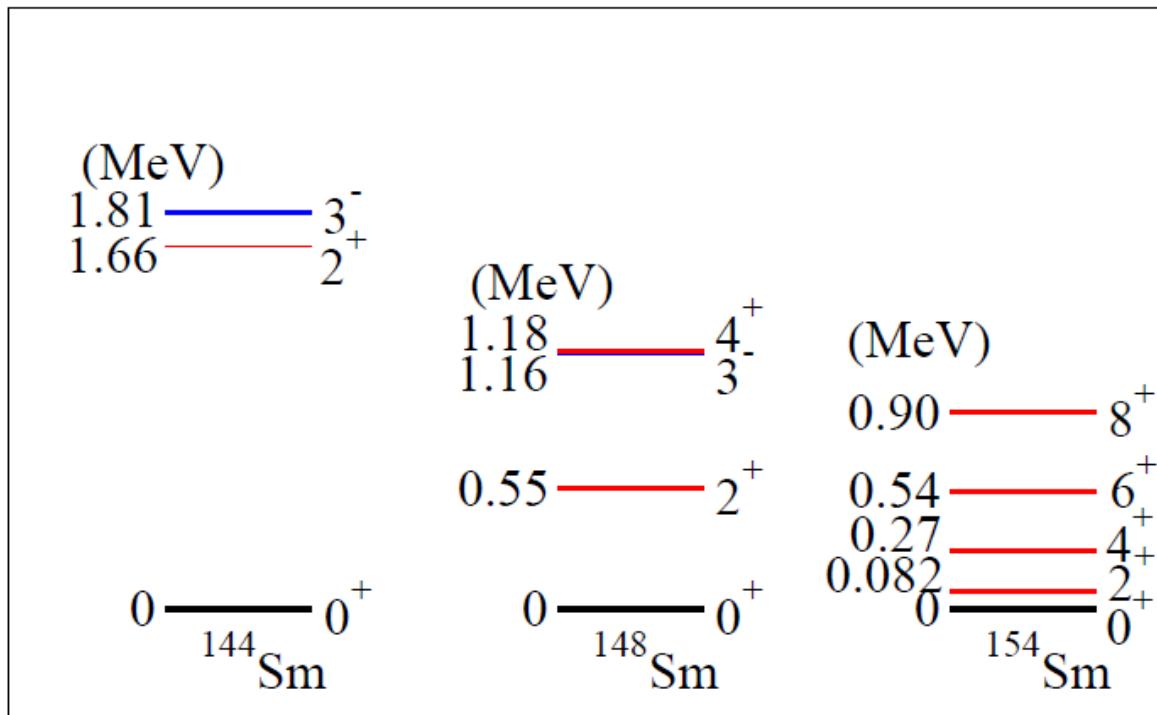
Fusion: strong interplay between nuclear structure and reaction

* Sub-barrier enhancement also for non-deformed targets:
couplings to low-lying collective excitations → coupling assisted tunneling



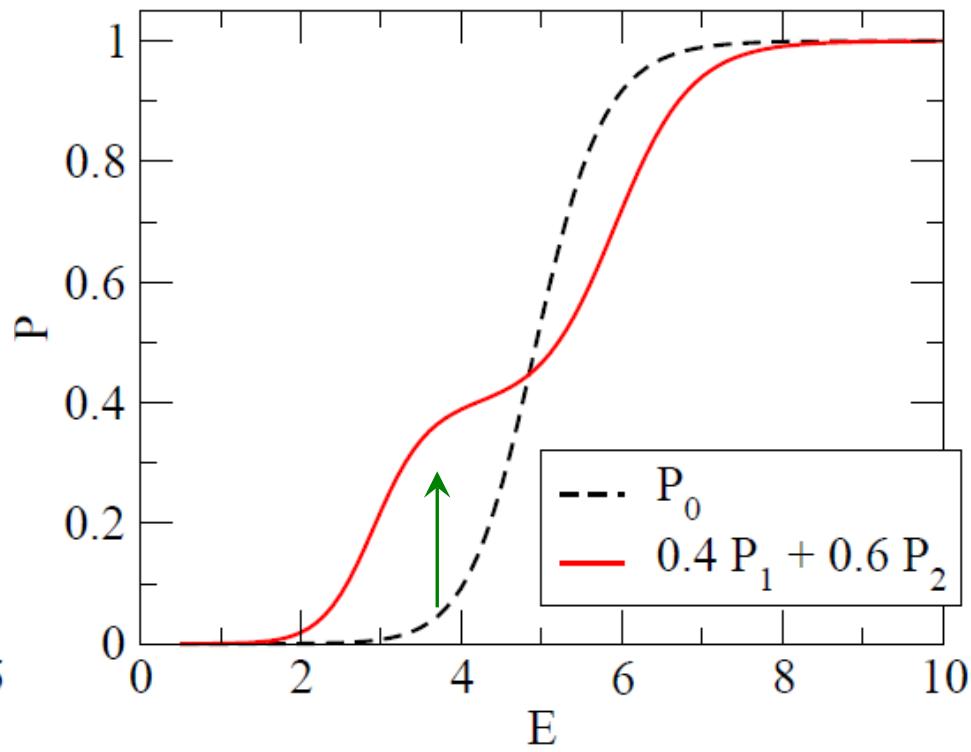
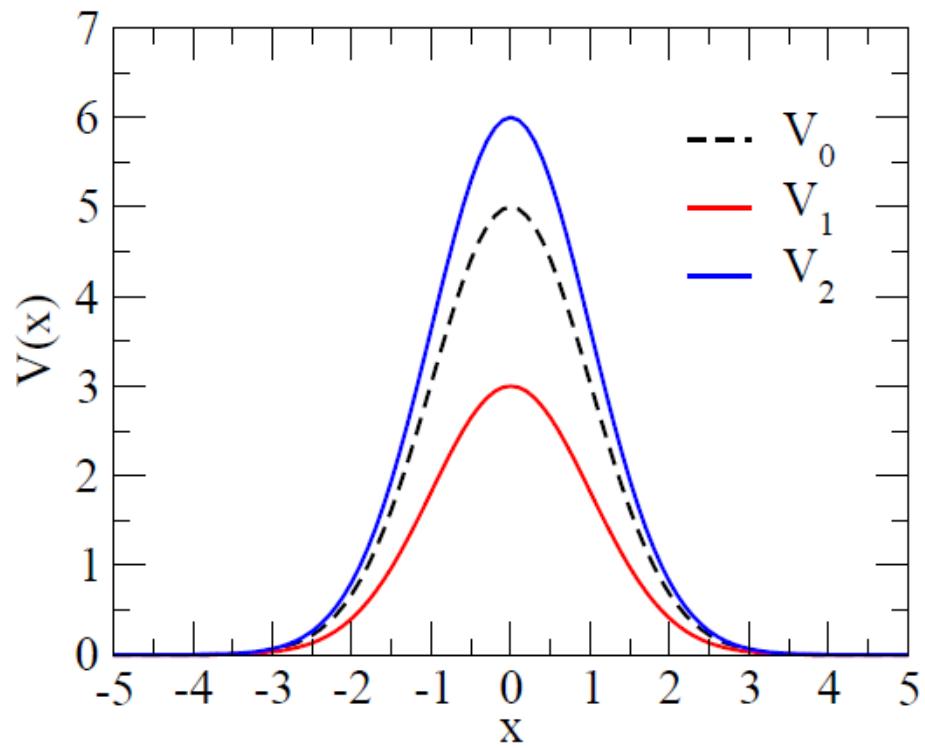
enhancement of fusion cross sections
: a general phenomenon

strong correlation with nuclear spectrum
→ coupling assisted tunneling



Enhancement of tunneling probability : a problem of two potential barriers

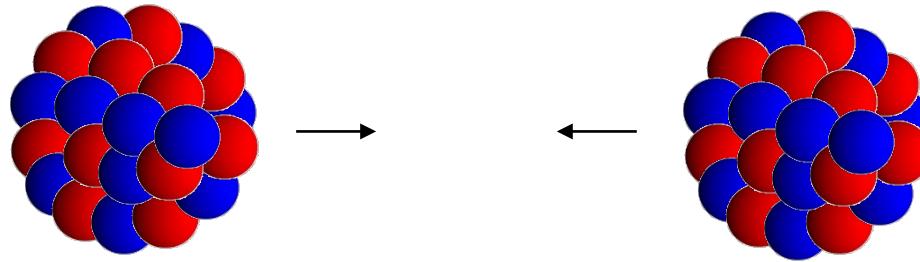
$$P(E) = P(E; V_0) \rightarrow w_1 P(E; V_1) + w_2 P(E; V_2)$$



“barrier distribution” due to couplings to excited states
in projectile/target nuclei

Coupled-channels method: a quantal scattering theory with excitations

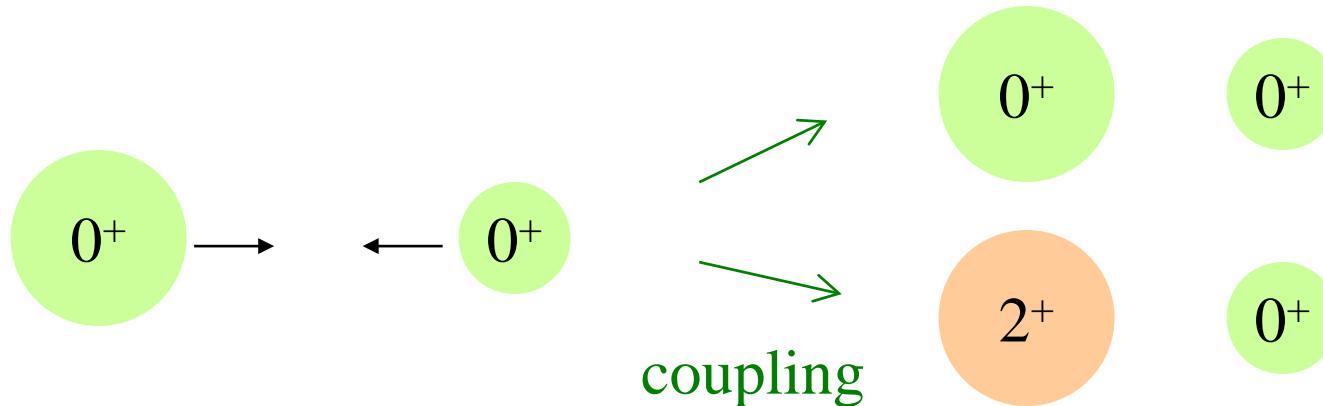
many-body problem



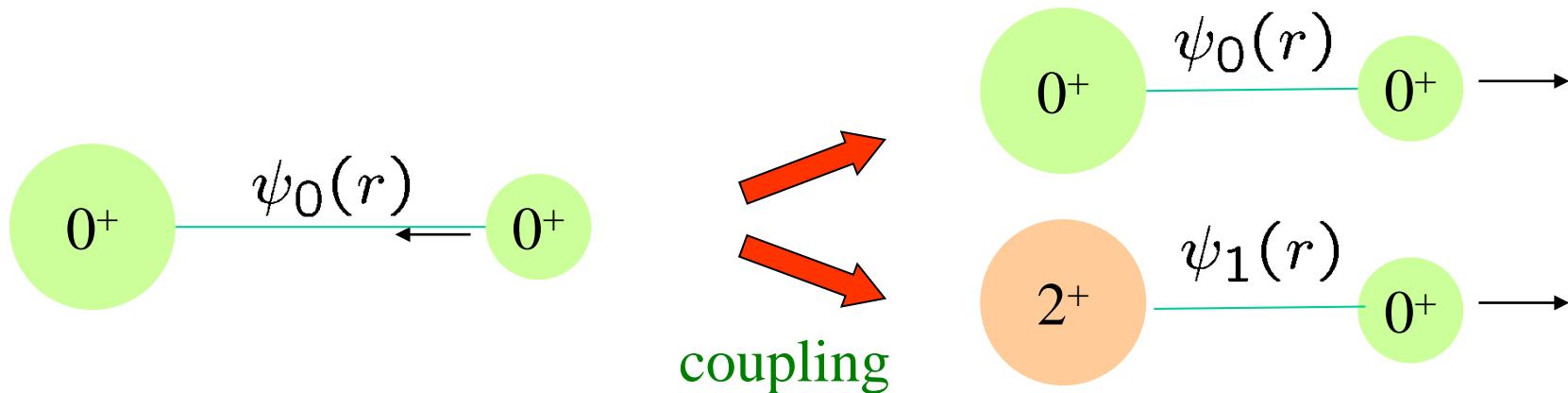
still very challenging



two-body problem, but with excitations
(coupled-channels approach)



Coupled-channels method: a quantal scattering theory with excitations



$$\left[-\frac{\hbar^2}{2\mu} \nabla^2 + \stackrel{\leftrightarrow}{V}(r) - \stackrel{\leftrightarrow}{E} \right] \vec{\psi}(r) = 0$$

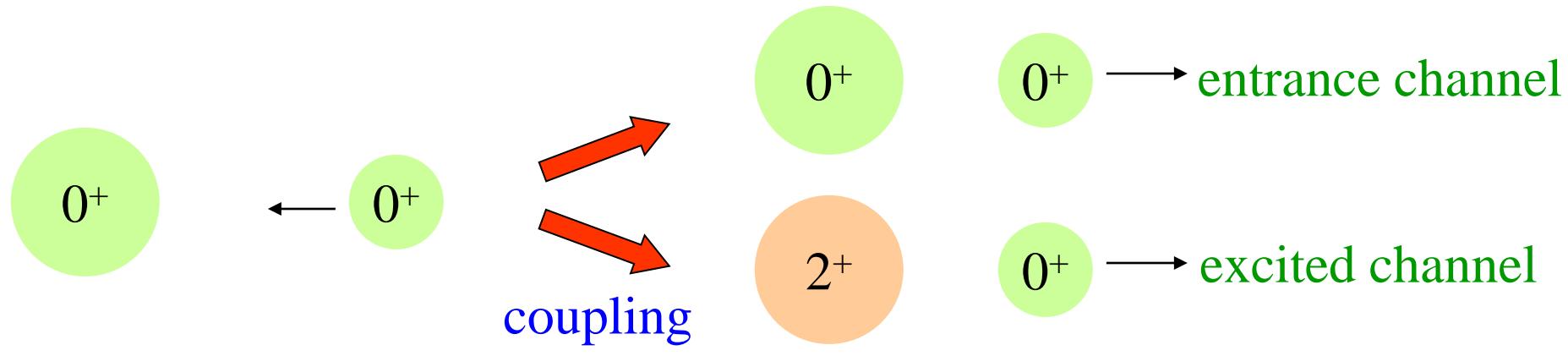
if written down more explicitly:

$$\left[-\frac{\hbar^2}{2\mu} \nabla^2 + V_0(r) + \epsilon_k - E \right] \psi_k(r) + \sum_{k'} \langle \phi_k | V_{\text{coup}} | \phi_{k'} \rangle \psi_{k'}(r) = 0$$

excitation energy

excitation operator

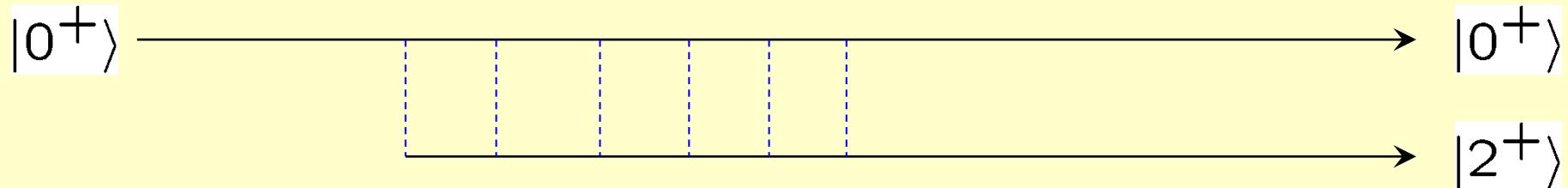
Coupled-channels method: a quantal scattering theory with excitations



$$\left[-\frac{\hbar^2}{2\mu} \nabla^2 + V_0(r) + \epsilon_k - E \right] \psi_k(\mathbf{r}) + \sum_{k'} \langle \phi_k | V_{\text{coup}} | \phi_{k'} \rangle \psi_{k'}(\mathbf{r}) = 0$$

excitation energy

excitation operator



full order treatment of excitation/de-excitation dynamics during reaction

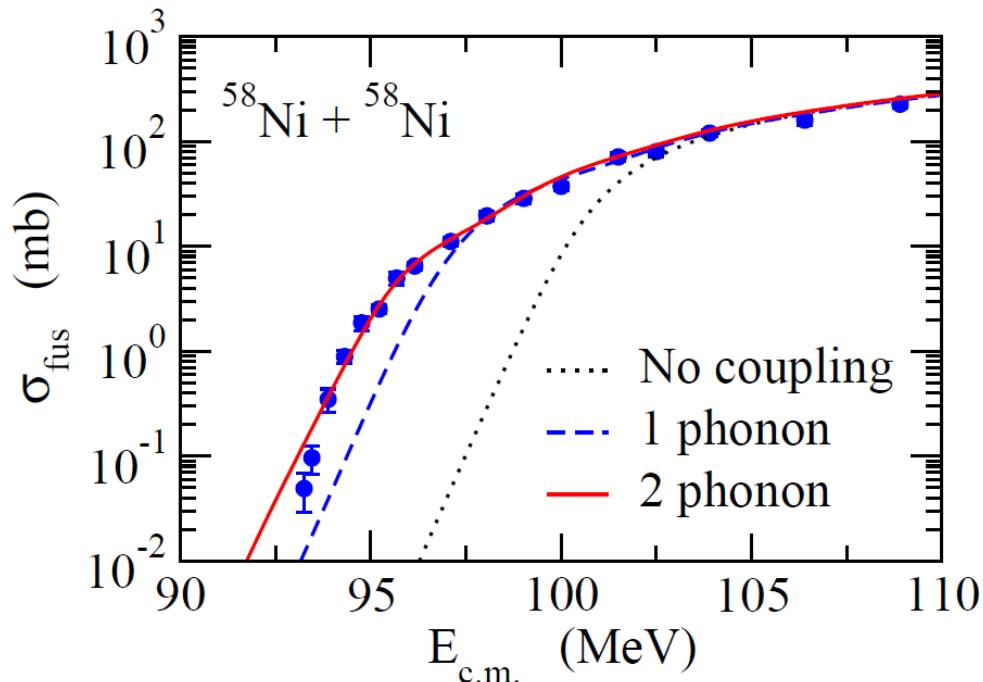
Inputs for C.C. calculations

i) Inter-nuclear potential

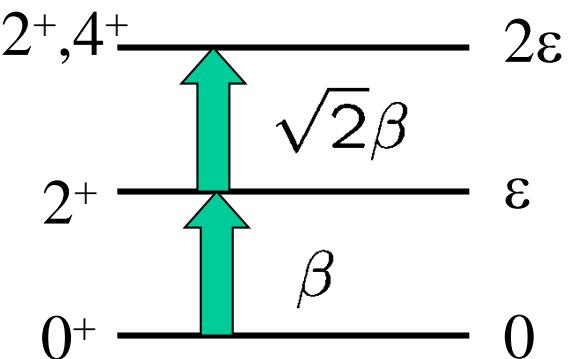
a fit to experimental data at above barrier energies

ii) Intrinsic degrees of freedom

in most of cases, (macroscopic) collective model
(rigid rotor / harmonic oscillator)



simple harmonic oscillator



Further development: semi-microscopic modelling

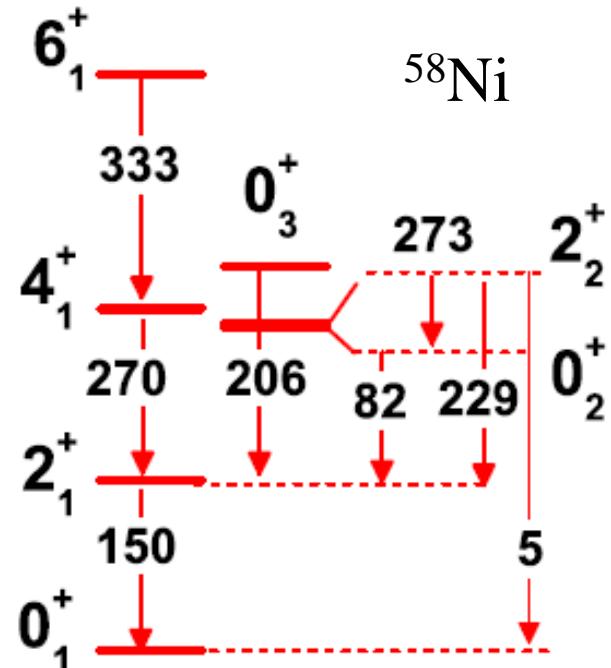
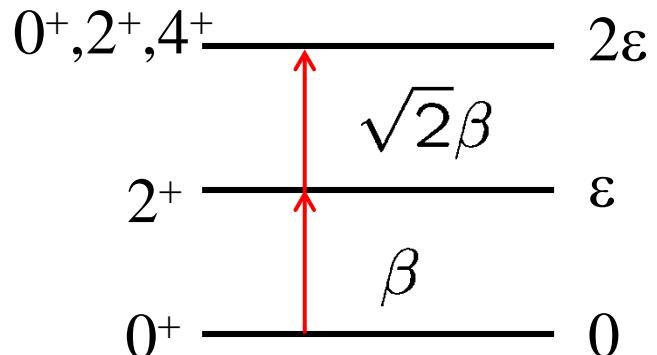
K.H. and J.M. Yao, PRC91('15) 064606

CCFULL

+ microscopic nuclear structure
calculations
(GCM, Shell Model, IBM.....)



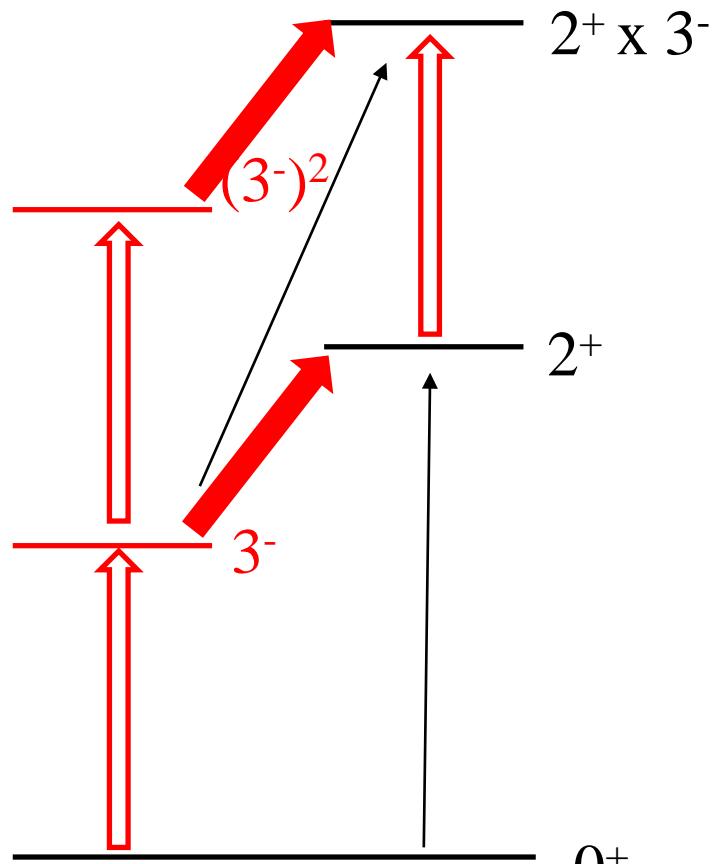
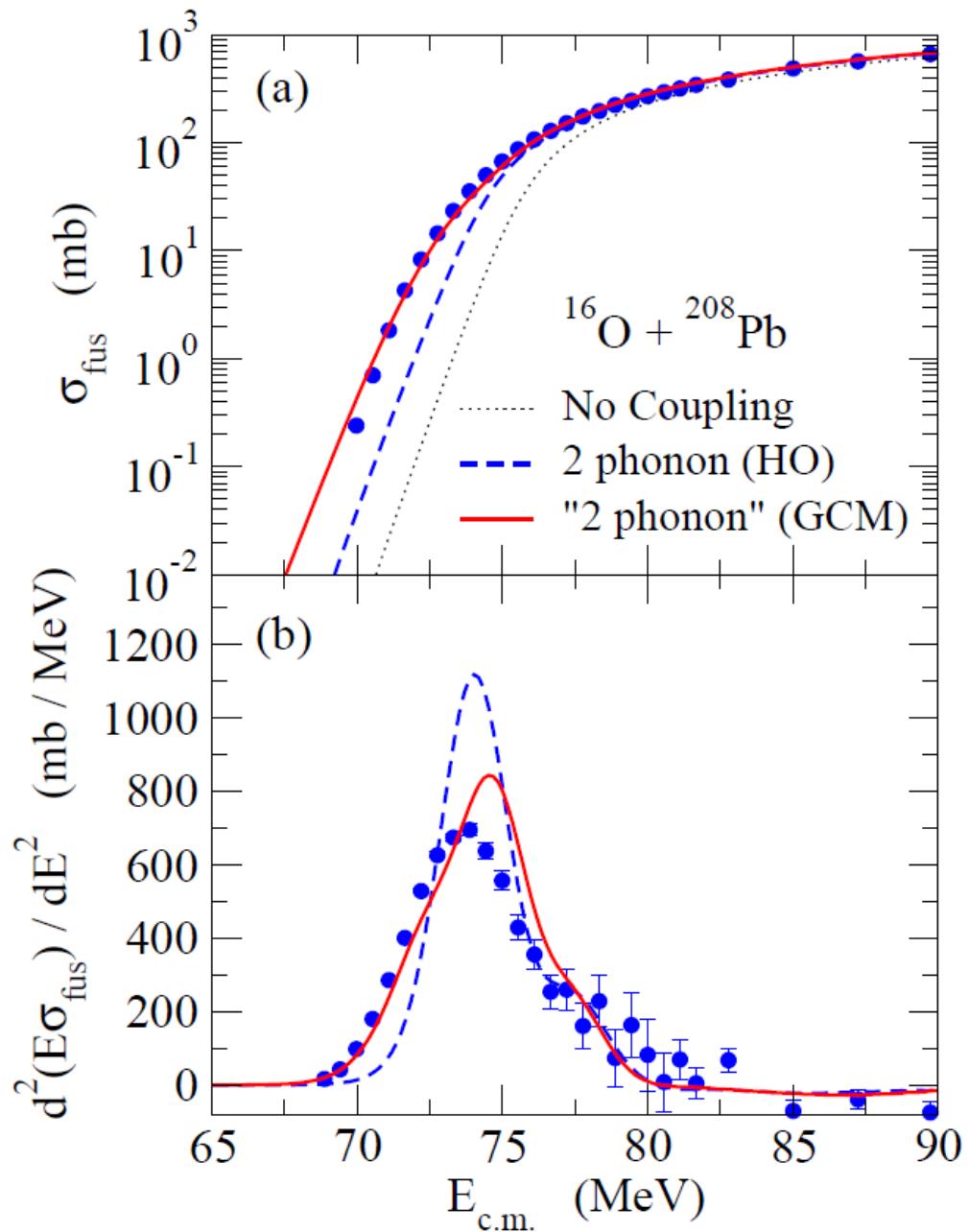
simple harmonic
oscillator



relativistic MF + GCM

anharmonicity of phonon spectra

CCFULL with RMF+GCM



J.M. Yao and K.H.,
PRC94 ('16) 11303(R)

From phenomenological approach to microscopic approach

Macroscopic (phenomenological)

C.C. with collective model

C.C. with inputs from
microscopic nuclear
structure calculations

- * Hagino-Yao
- * Ichikawa-Matsuyanagi

C.C. with inputs based
on TDHF

- * Umar et al. (DC-TDHF)
- * Washiyama-Lacroix
- * Simenel et al.

TDHF simulations

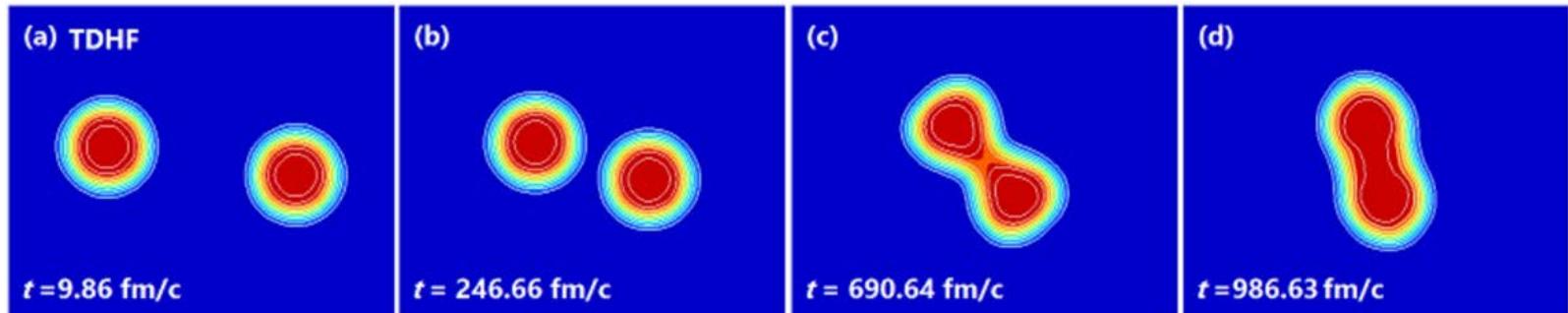
- * Simenel
- * Sekizawa
- * Washiyama
- * Iwata-Otsuka
- * Maruhn, Stevenson

Microscopic

“ab initio”, but no tunneling

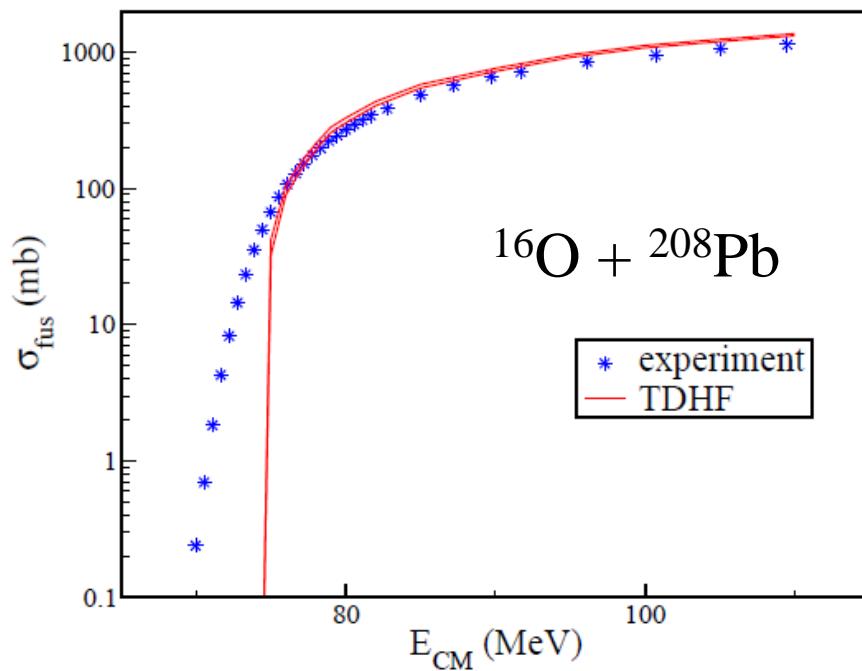
TDHF simulation

TDHF = Time Dependent Hartree-Fock



S. Ebata, T. Nakatsukasa, JPC Conf. Proc. 6 ('15) 020056

“ab-initio”, but no tunneling



C. Simenel,
EPJA48 ('12) 152

TDHF simulation

“ab-initio”, but no tunneling

- ✓ DC-TDHF (Umar, Oberacker, Maruhn)

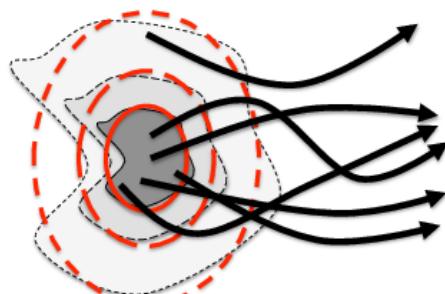
➤ Beyond mean-field approximation

- ✓ Collective Hamiltonian and requantization

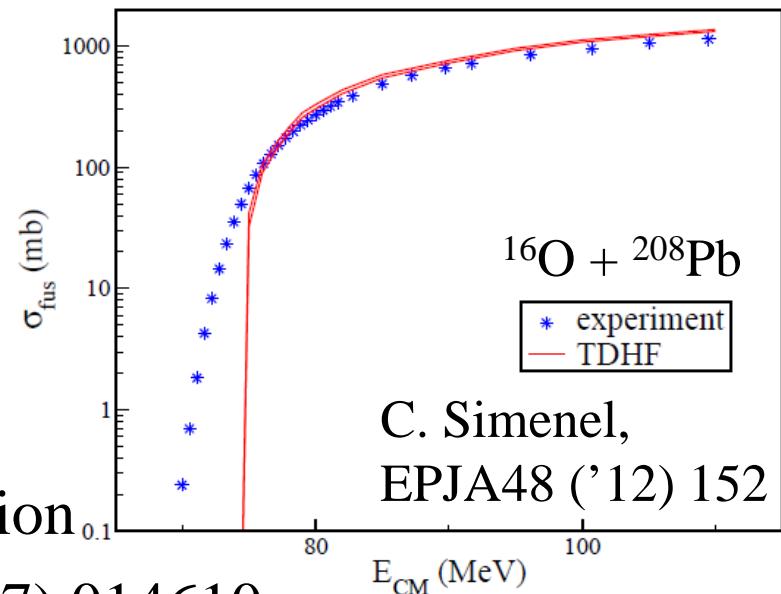
K. Wen and T. Nakatsukasa, PRC96 ('17) 014610

- ✓ Time-dependent Generator Coordinate Method (TD-GCM)

$$|\Psi(t)\rangle = \int dq f(q, t) |\Phi_q(t)\rangle$$



an important future direction



a linear superposition of many
TDHF trajectories (Slater determinants)

cf. Stochastic mean-field method
B. Yilmaz et al.,
PRC90 ('14) 054617

Fusion reactions for SHE

the element 113: Nh

113 Nh nihonium	115 Mc moscovium
117 Ts tennessine	118 Og oganesson

November, 2016



Group → 1 Period ↓		2	3	4	5	6	7	8	9	10	11	12	13	14	15	16	17	18	
1	1 H																2 He		
2	3 Li	4 Be															10 Ne		
3	11 Na	12 Mg															18 Ar		
4	19 K	20 Ca	21 Sc	22 Ti	23 V	24 Cr	25 Mn	26 Fe	27 Co	28 Ni	29 Cu	30 Zn	31 Ga	32 Ge	33 As	34 Se	36 Kr		
5	37 Rb	38 Sr	39 Y	40 Zr	41 Nb	42 Mo	43 Tc	44 Ru	45 Rh	46 Pd	47 Ag	48 Cd	49 In	50 Sn	51 Sb	52 Te	54 Xe		
6	55 Cs	56 Ba	57 La	*	72 Hf	73 Ta	74 W	75 Re	76 Os	77 Ir	78 Pt	79 Au	80 Hg	81 Tl	82 Pb	83 Bi	85 At	86 Rn	
7	87 Fr	88 Ra	89 Ac	*	104 Rf	105 Db	106 Sg	107 Bh	108 Hs	109 Mt	110 Ds	111 Rg	112 Cn	113 Nh	114 Fl	115 Mc	116 Lv	117 Ts	118 Og
* 58 Ce 59 Pr 60 Nd 61 Pm 62 Sm 63 Eu 64 Gd 65 Tb 66 Dy 67 Ho 68 Er 69 Tm 70 Yb 71 Lu																			
* 90 Th 91 Pa 92 U 93 Np 94 Pu 95 Am 96 Cm 97 Bk 98 Cf 99 Es 100 Fm 101 Md 102 No 103 Lr																			

Wikipedia

Fusion reactions for SHE

the element 113: Nh

113 鉻 nihonium	115 镆 moscovium
117 石田 tennessine	118 氣 oganesson

核データニュース, No.118 (2017)

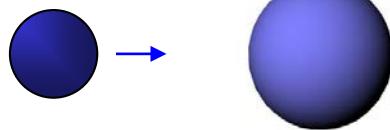
話題・解説(I)

Chinese Names of New Elements with $Z = 113, 115, 117 \& 118$

Shan-Gui Zhou (周善貴/周善貴)¹

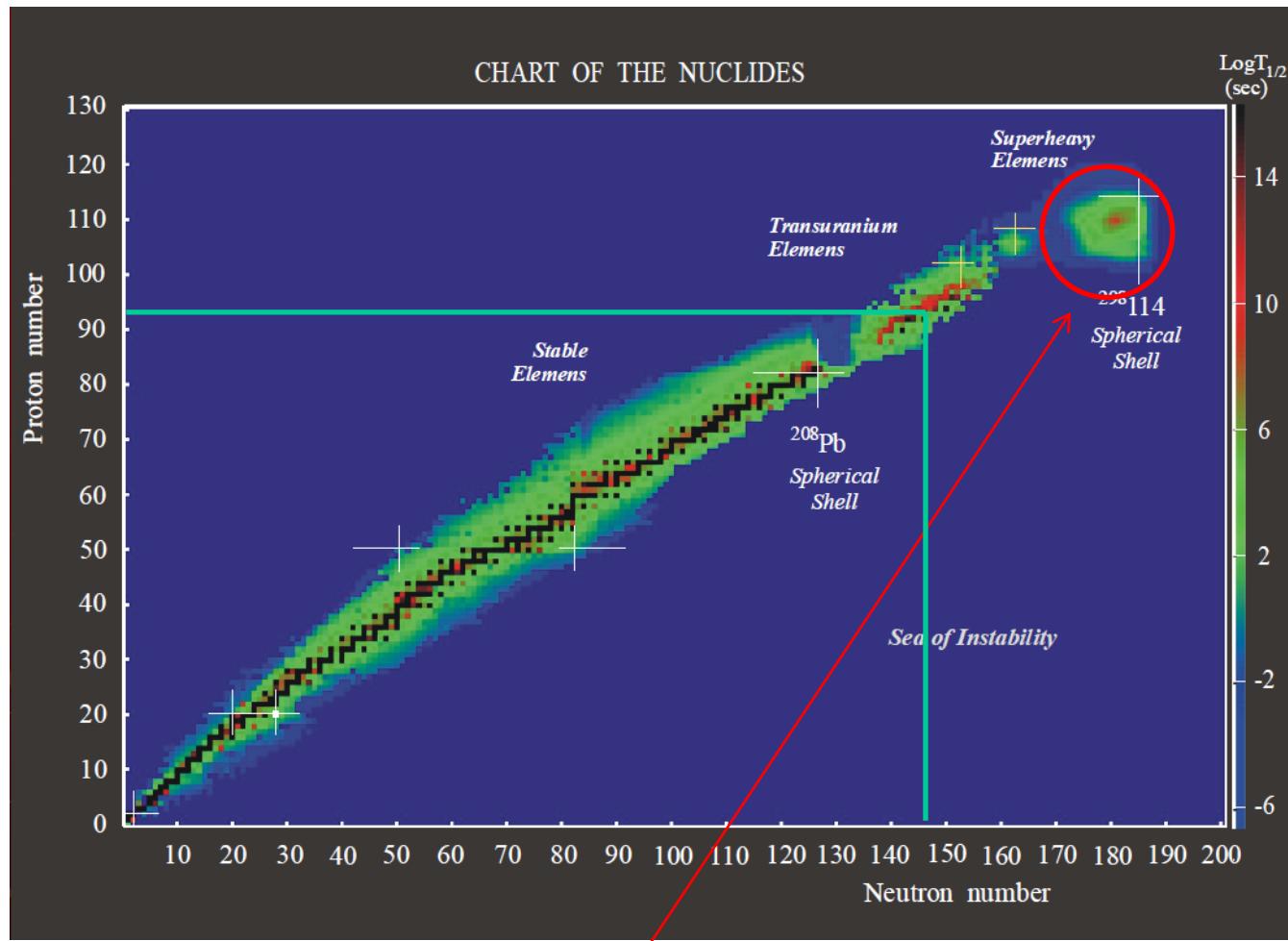
Institute of Theoretical Physics, Chinese Academy of Sciences, Beijing 100190, China
(中国科学院理论物理研究所/中國科學院理論物理研究所)

November, 2016



Heavy-ion fusion reaction

Future perspectives: Superheavy elements



island of stability around $Z=114, N=184$

W.D. Myers and W.J. Swiatecki (1966), A. Sobiczewski et al. (1966)

Yuri Oganessian

who is she?

	Cs	Ba		Hf	Ta	W	Ru	Tc		Pt	Uu	Hg		Pt	Ru	Pt	W		Ru
7	87 Fr	88 Ra		104 Rf	105 Db	106 Sg	107 Bh	108 Hs	109 Mt	110 Ds	111 Rg	112 Cn	113 Uut	114 Fl	115 Uup	116 Lv	117 Uus	118 Uuo	

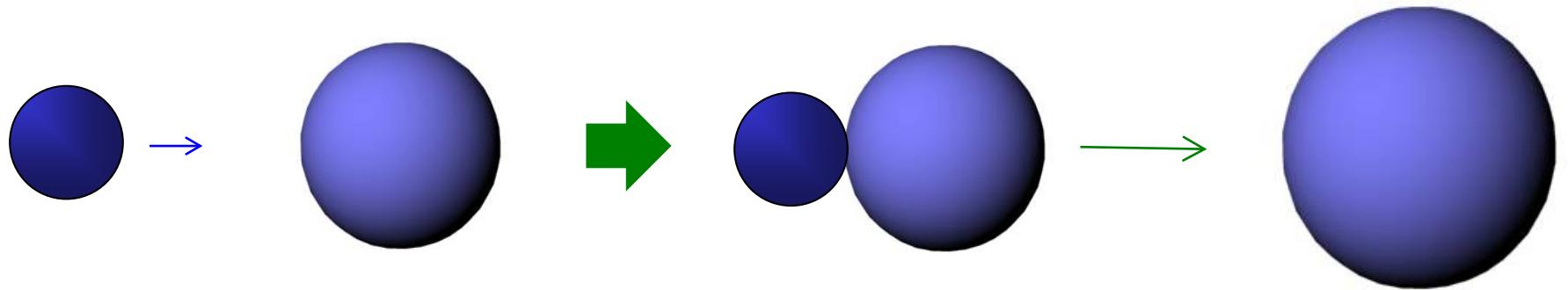
- Z=110 Darmstadtium (Ds) 1994 Germany
Z=111 Roentgenium (Rg) 1994 Germany
Z=112 Copernicium (Cn) 1996 Germany
Z=113 **Nihonium (Nh)** 2003 Russia / 2004 Japan
Z=114 Flerovium (Fl) 1999 Russia
Z=115 **Moscovium (Mc)** 2003 Russia
Z=116 Livermorium (Lv) 2000 Russia
Z=117 **Tennessine (Ts)** 2010 Russia
Z=118 **Oganesson (Og)** 2002 Russia

113 Nh nihonium	115 Mc moscovium
117 Ts tennessine	118 Og oganesson

How to synthesize SHE?

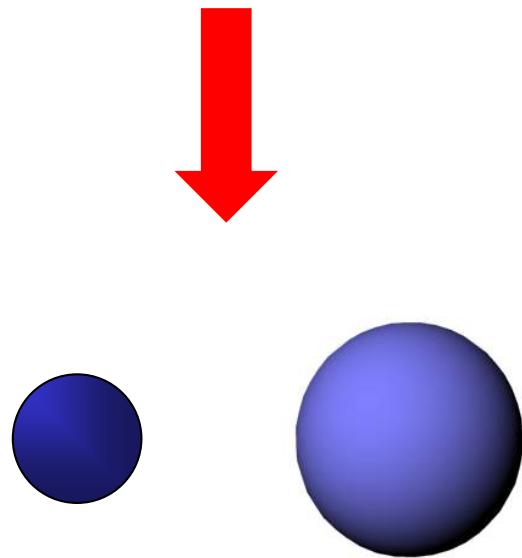
Nuclear fusion reactions

e.g.,

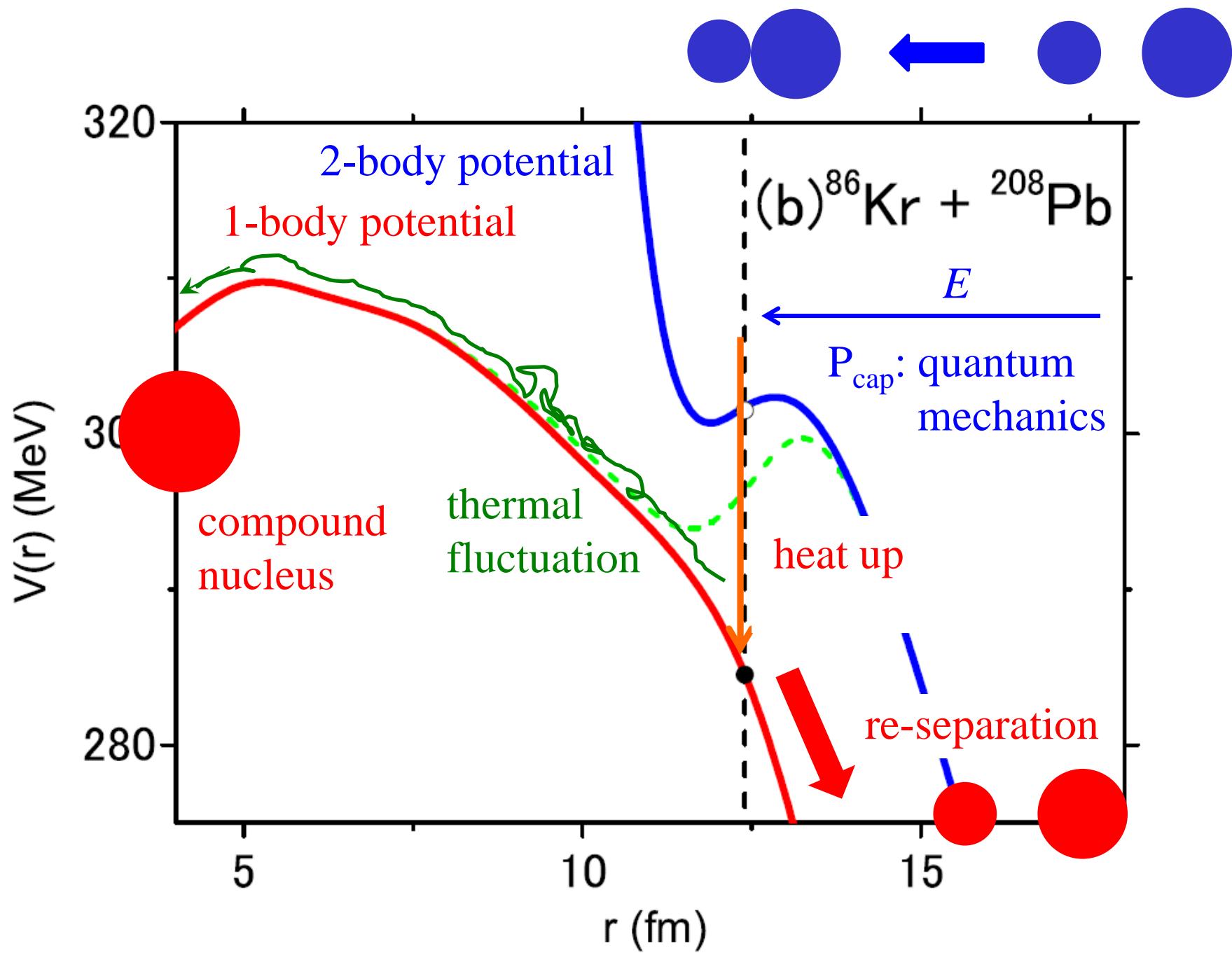


two positive charges
repel each other

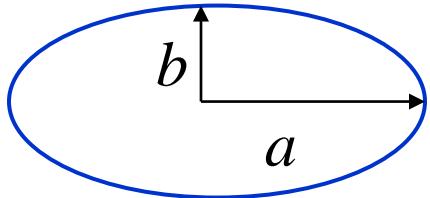
compound
nucleus



re-separation

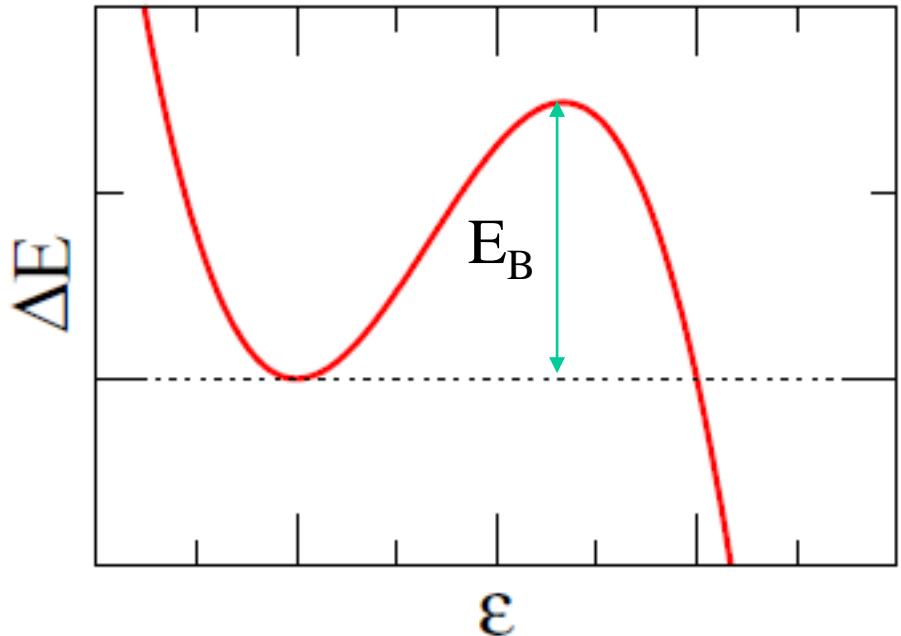


(note) fission barrier in the liquid drop model



$$\begin{aligned} a &= R \cdot (1 + \epsilon) \\ b &= R \cdot (1 + \epsilon)^{-1/2} \\ ab^2 &= R^3 = \text{constant} \end{aligned}$$

$$\begin{aligned} \Delta E &= \Delta E_{\text{surf}} + \Delta E_{\text{coul}} \\ &= E_S^{(0)} \left\{ \frac{2}{5}(1-x)\epsilon^2 - \frac{4}{105}(1+2x)\epsilon^3 + \dots \right\} \end{aligned}$$

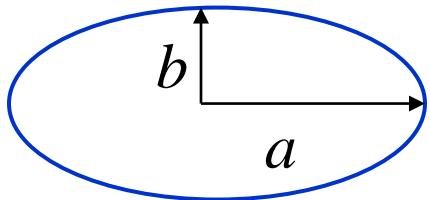


$$E_S^{(0)} = +a_S A^{2/3}$$

$$x \equiv \frac{E_C^{(0)}}{2E_S^{(0)}} = \frac{a_C}{2a_S} \cdot \frac{Z^2}{A} \sim \frac{1}{53.3} \cdot \frac{Z^2}{A}$$

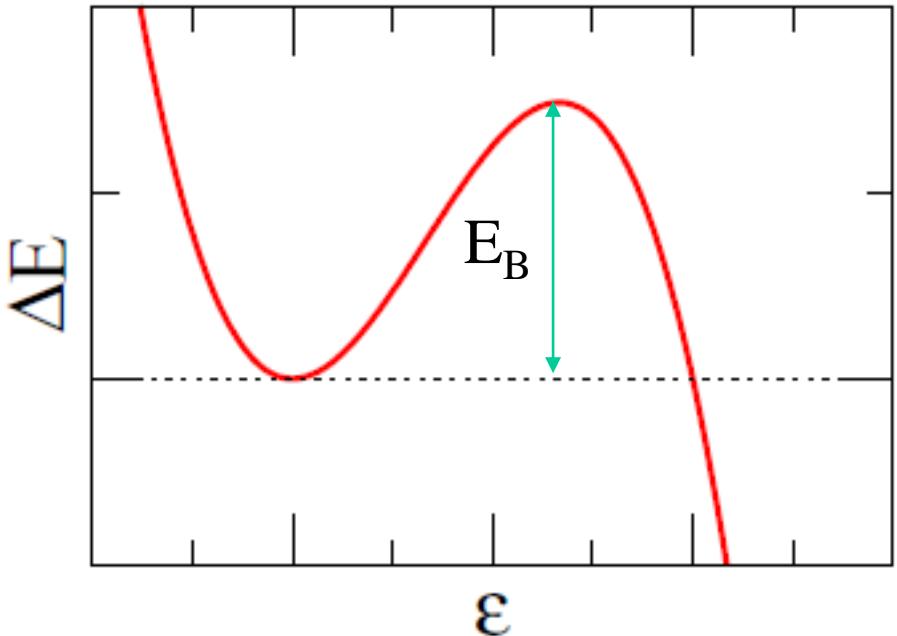
$$E_C^{(0)} = a_C Z(Z-1)/A^{1/3}$$

(note) fission barrier in the liquid drop model



$$\begin{aligned} a &= R \cdot (1 + \epsilon) \\ b &= R \cdot (1 + \epsilon)^{-1/2} \\ ab^2 &= R^3 = \text{constant} \end{aligned}$$

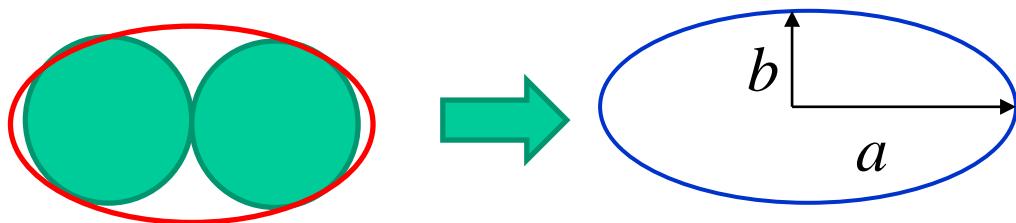
$$\begin{aligned} \Delta E &= \Delta E_{\text{surf}} + \Delta E_{\text{coul}} \\ &= E_S^{(0)} \left\{ \frac{2}{5}(1-x)\epsilon^2 - \frac{4}{105}(1+2x)\epsilon^3 + \dots \right\} \end{aligned}$$



fission barrier:

$$\begin{aligned} \epsilon_B &= \frac{21(1-x)}{3(1+2x)} \\ E_B &= \frac{98}{15} \cdot \frac{(1-x)^3}{(1+2x)^2} \cdot E_S^{(0)} \end{aligned}$$

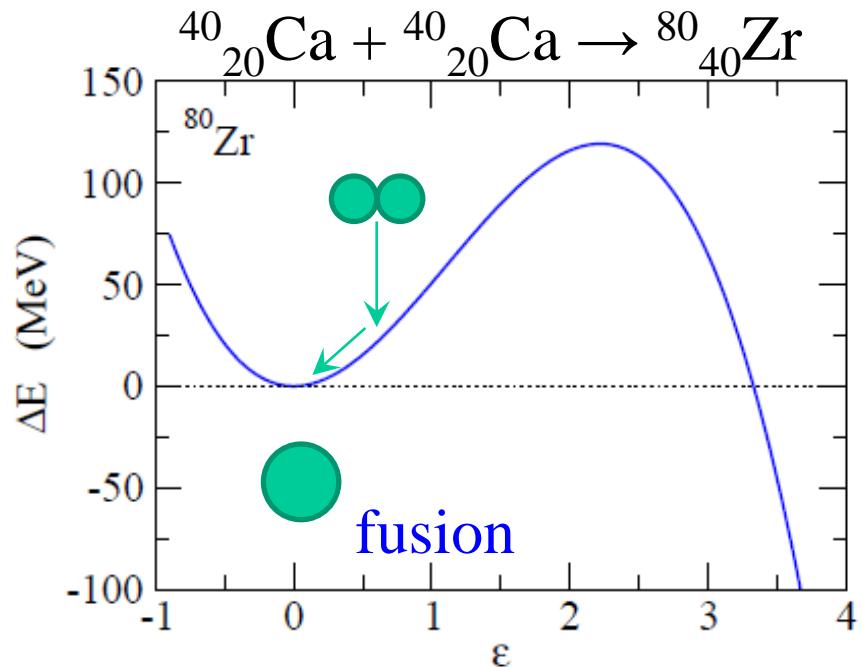
if two identical nuclei contact:



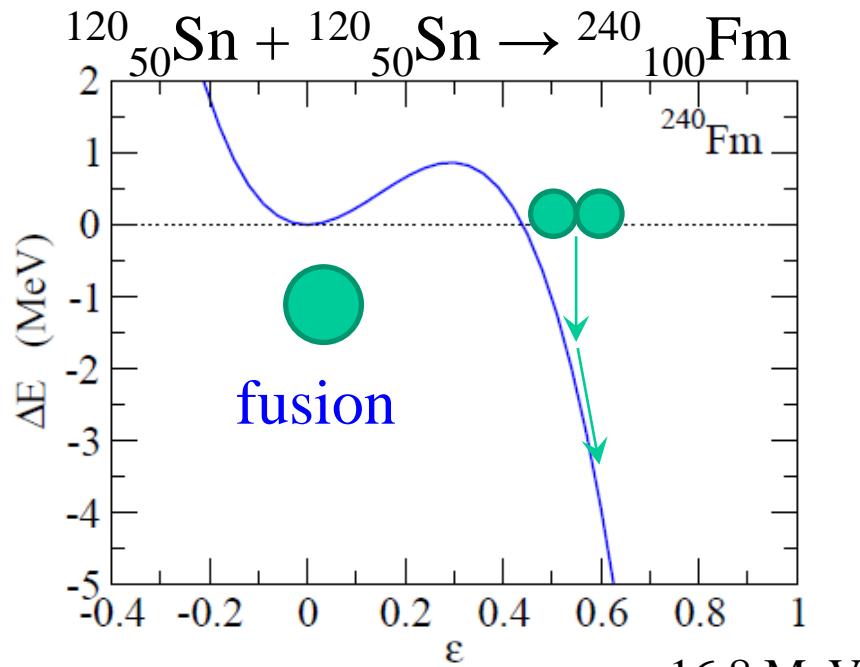
$$a = R_0 \cdot (1 + \epsilon)$$

$$b = R_0 \cdot (1 + \epsilon)^{-1/2}$$

$$\frac{a}{b} \sim \frac{2R}{R} = 2 \rightarrow \epsilon \sim 0.587$$

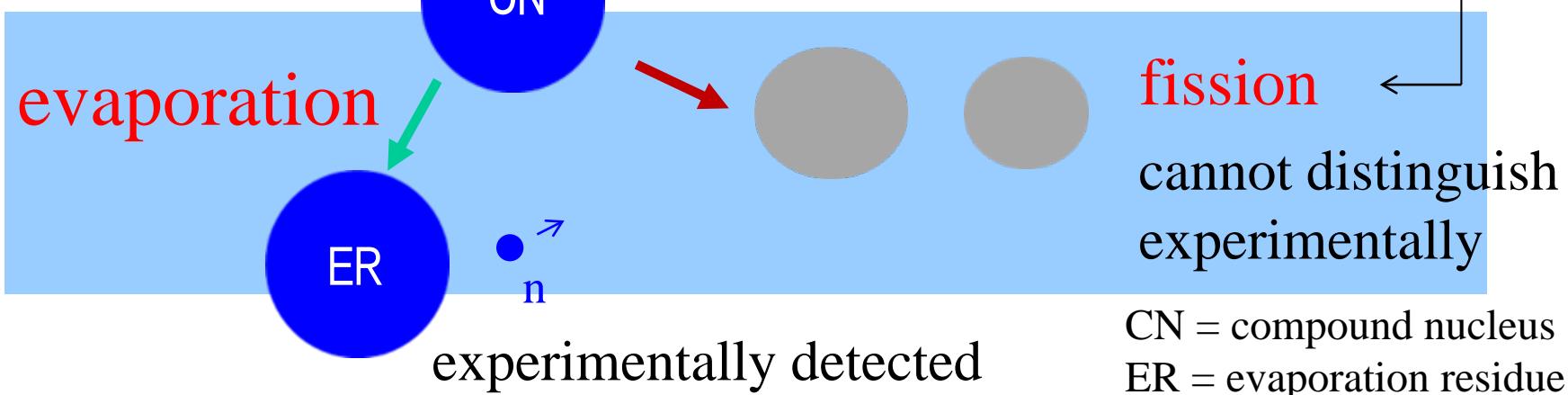
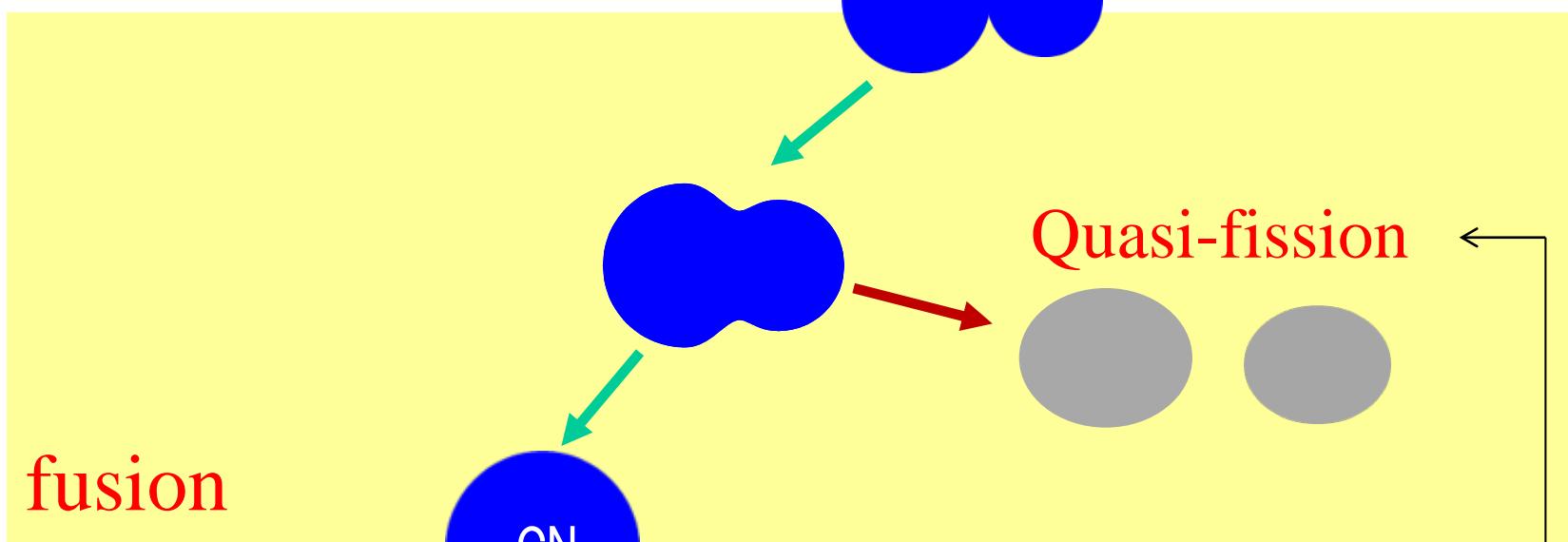
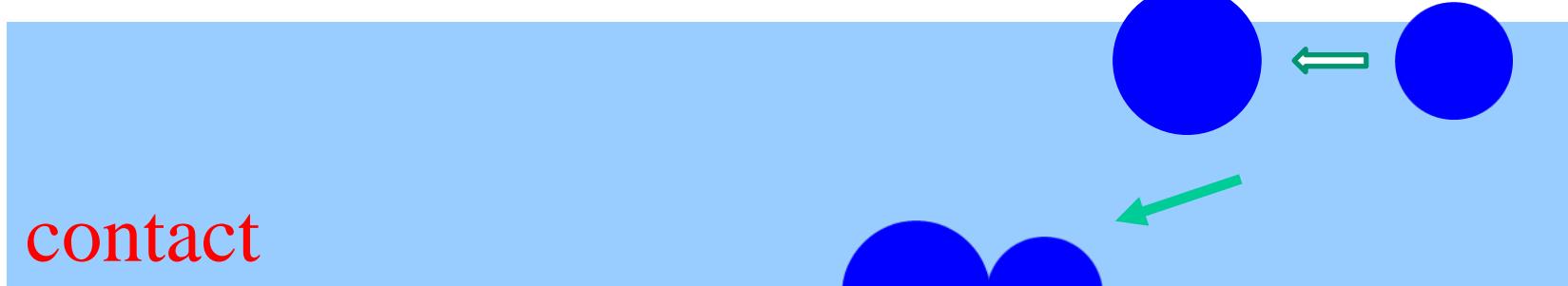


threshold: $Z_1^* Z_2 = 1600 \sim 1800$

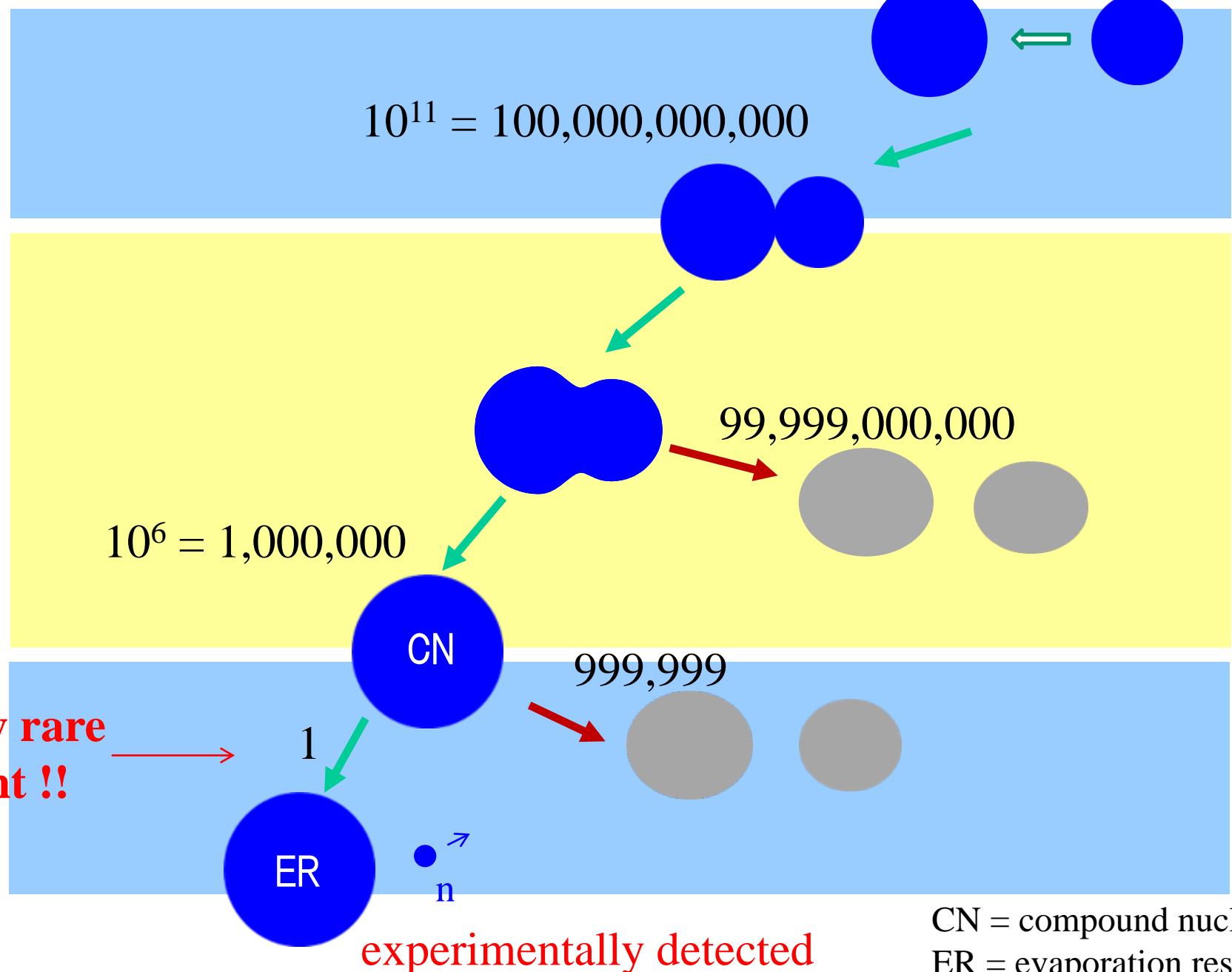


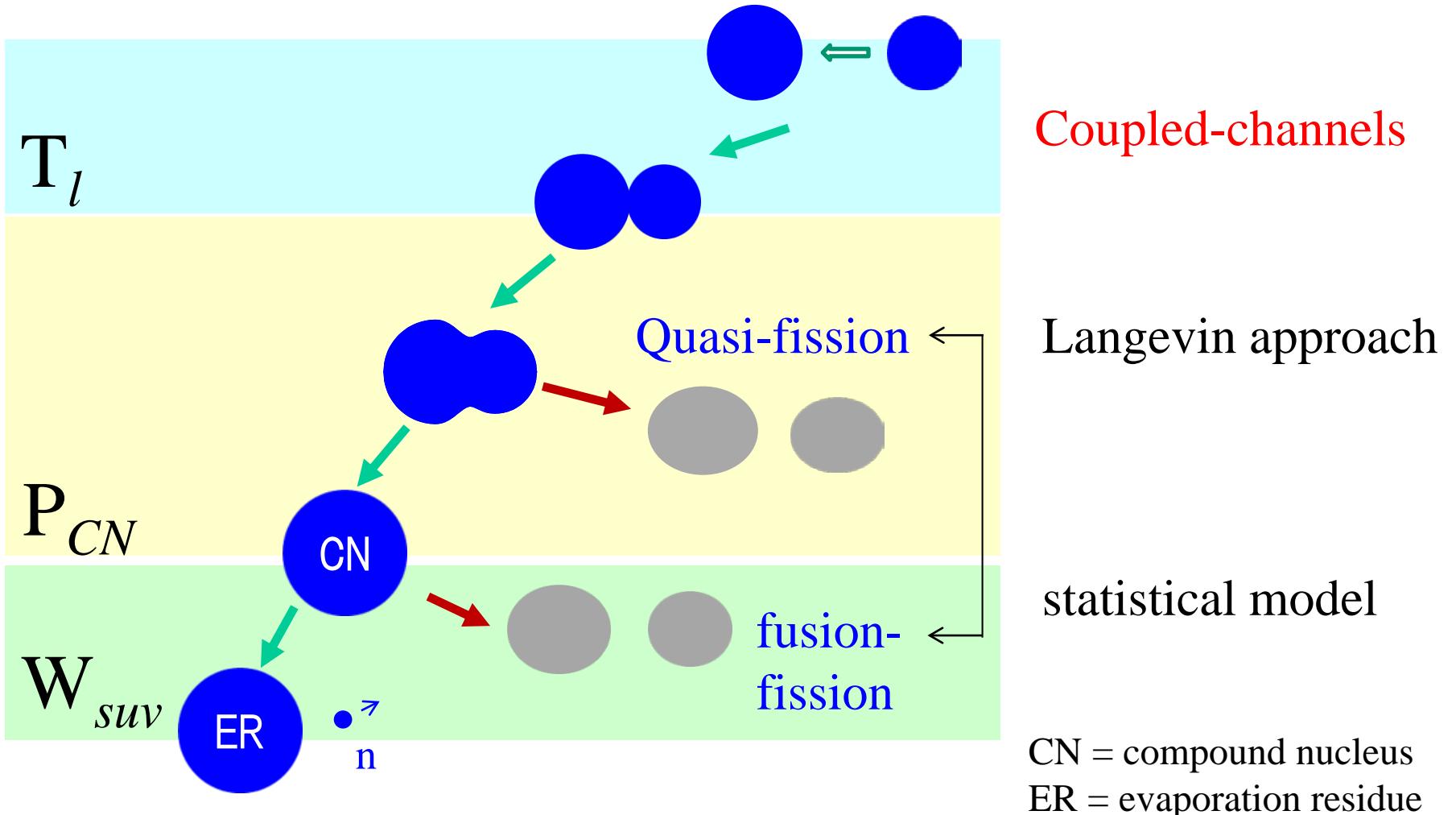
$$a_S = 16.8 \text{ MeV}$$

$$a_C = 0.72 \text{ MeV}$$

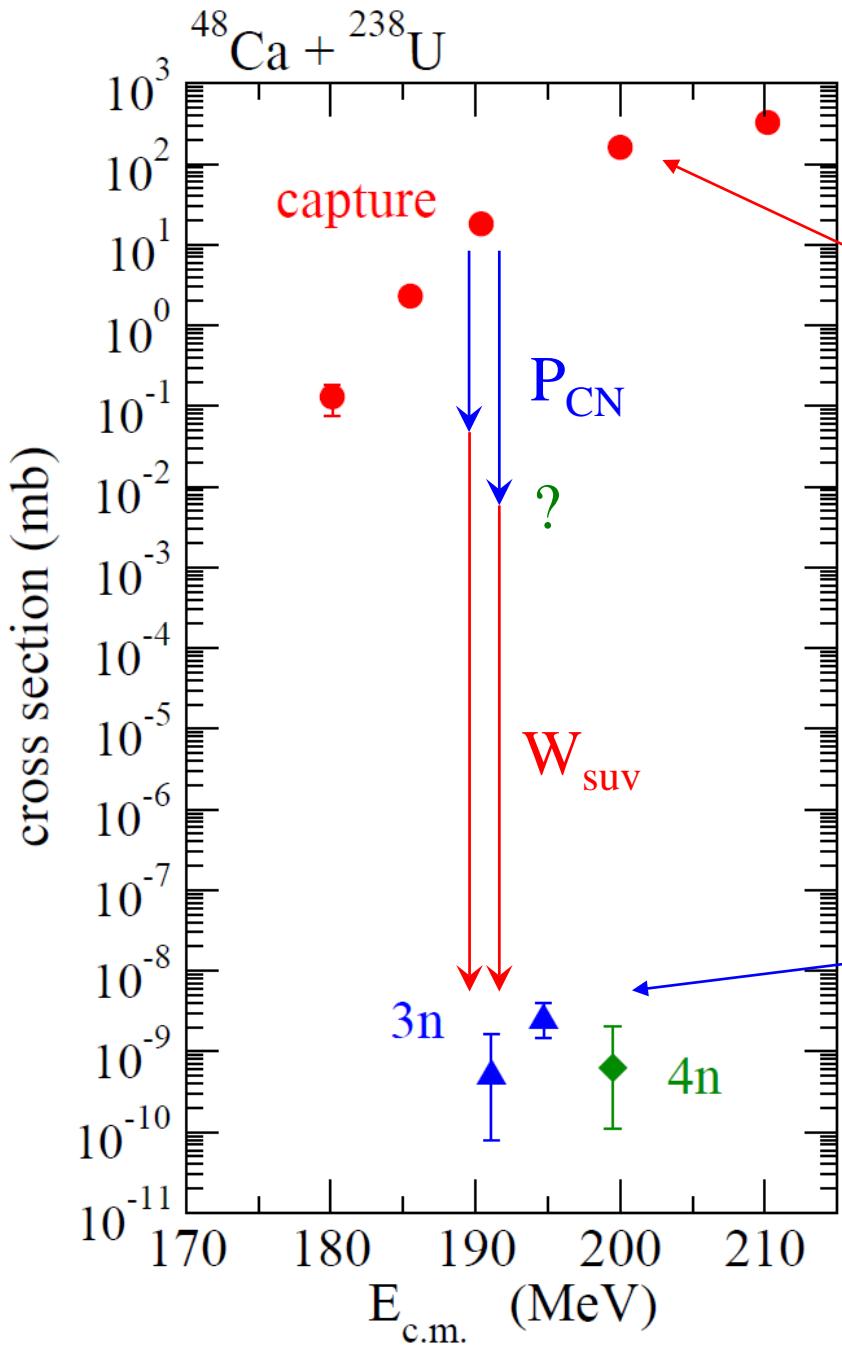


typical values for Ni + Pb reaction





$$\sigma_{ER}(E) = \frac{\pi}{k^2} \sum_l (2l + 1) T_l(E) P_{CN}(E, l) W_{suv}(E^*, l)$$



no experimental data for P_{CN}

$$\sigma_{\text{cap}}(E) = \frac{\pi}{k^2} \sum_l (2l + 1) T_l(E)$$

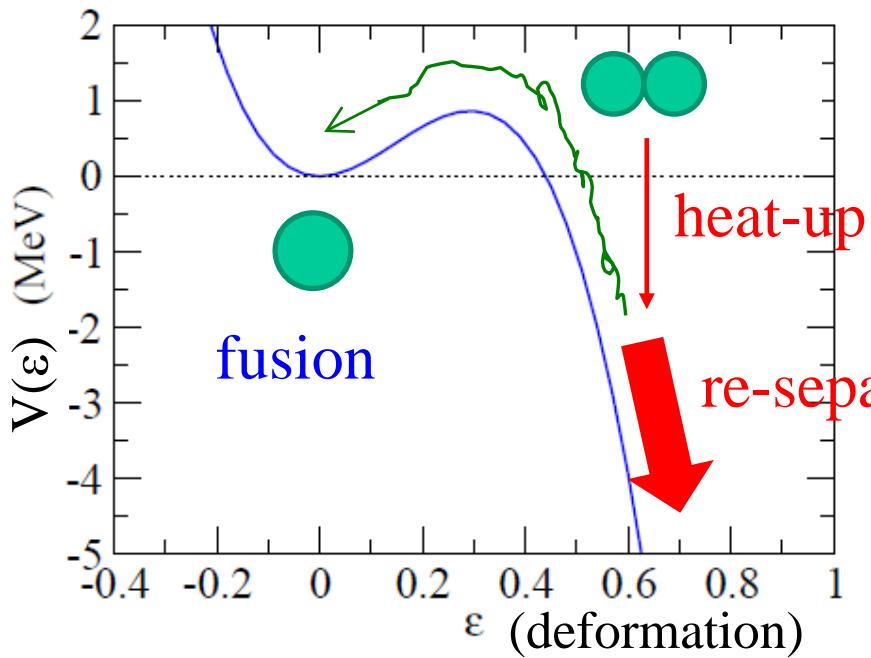
$$\sigma_{\text{CN}}(E) = \frac{\pi}{k^2} \sum_l (2l + 1) T_l(E) \times P_{\text{CN}}$$

not available

$$\sigma_{\text{ER}}(E) = \frac{\pi}{k^2} \sum_l (2l + 1) T_l(E) \times P_{\text{CN}} \cdot W_{\text{suv}}$$

large uncertainties

Langevin approach



thermal fluctuation

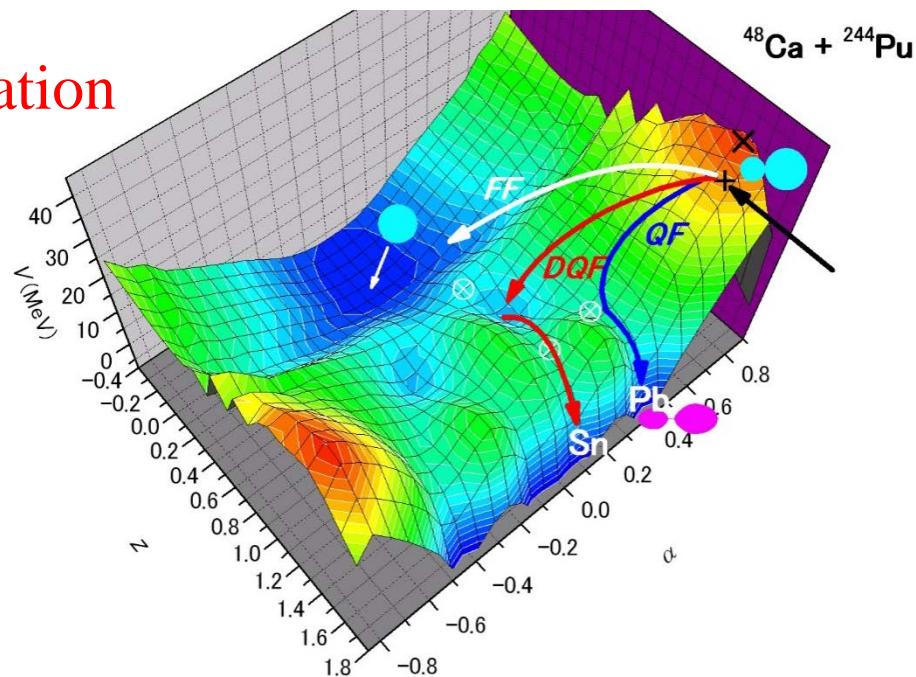
→ Langevin method
(Brownian method)

$$m \frac{d^2q}{dt^2} = -\frac{dV(q)}{dq} - \gamma \frac{dq}{dt} + R(t)$$

γ : friction coefficient
 $R(t)$: random force

multi-dimensional extention

- q :
 - internuclear separation,
 - deformation,
 - asymmetry of the two fragments



- ✓ Abe, Wada, Bao
- ✓ Aritomo, Ohta
- ✓ Zagrebaev

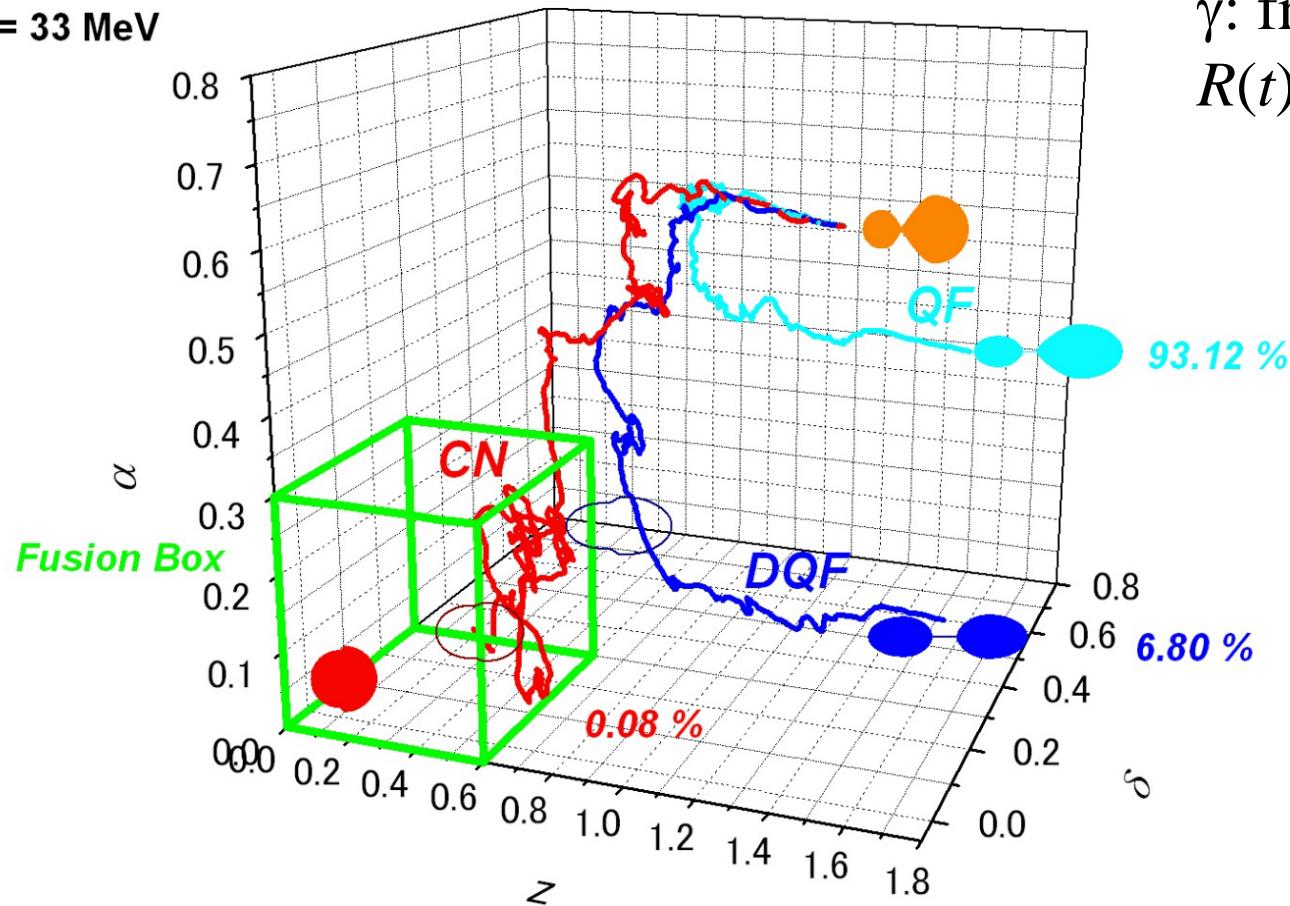
Theory: Lagenvin approach

multi-dimensional extension of:

$$m \frac{d^2q}{dt^2} = -\frac{dV(q)}{dq} - \gamma \frac{dq}{dt} + R(t)$$



$E^* = 33 \text{ MeV}$

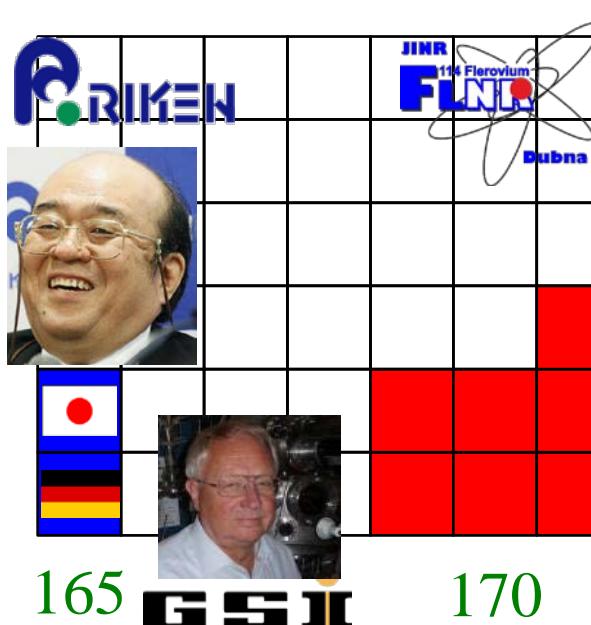


γ : friction coefficient

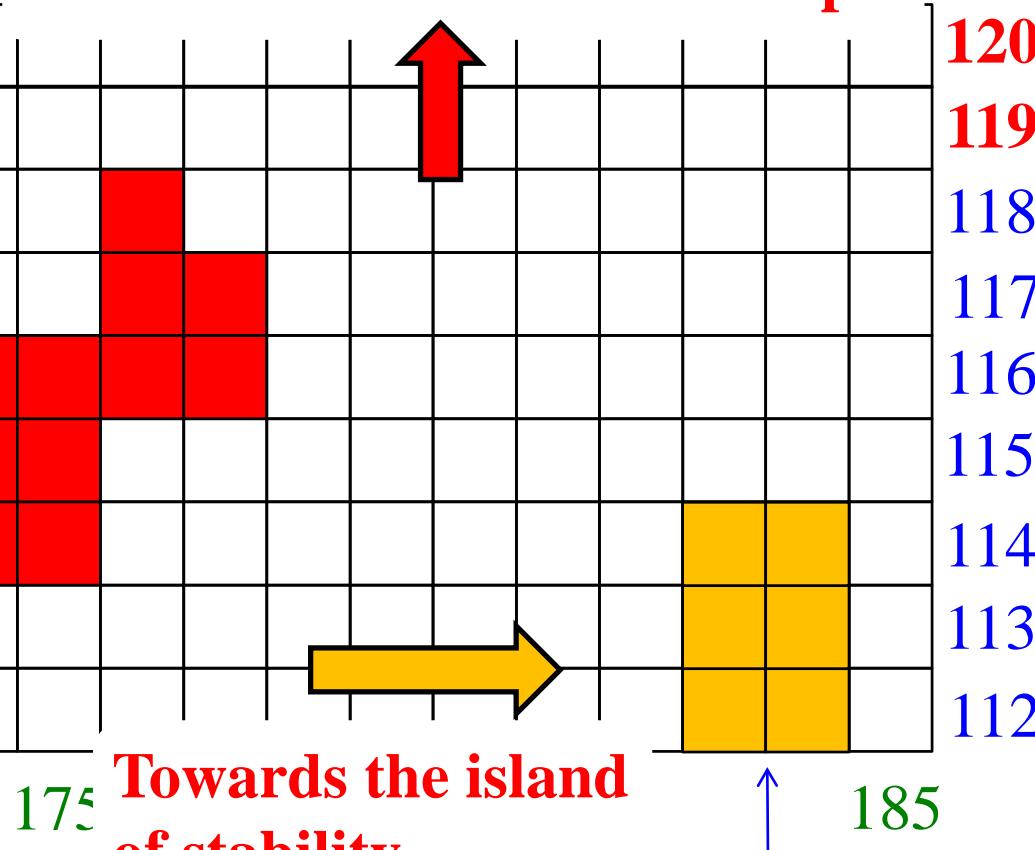
$R(t)$: random force

Future directions

Superheavy elements
synthesized so far



Towards Z=119 and 120 isotopes

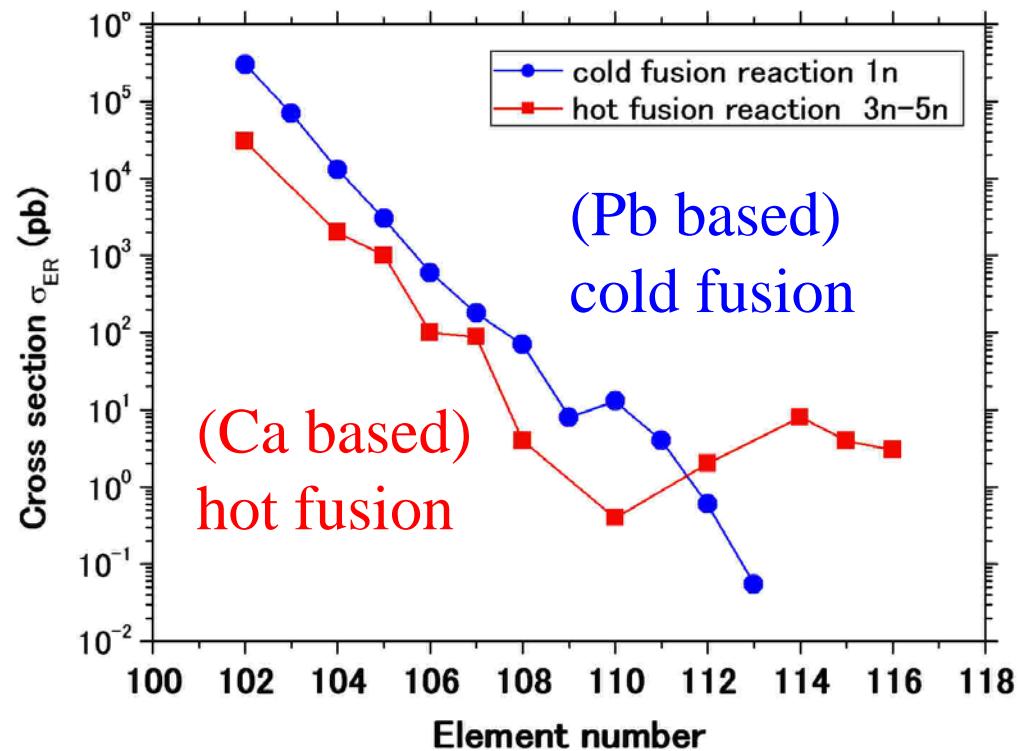
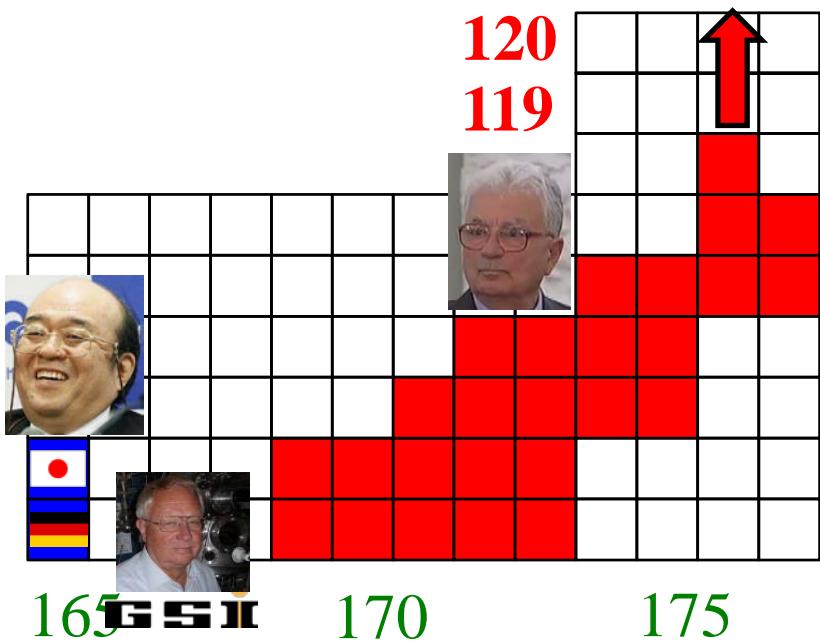


Theoretical issues:

- to understand the reaction dynamics
- to make a reliable theoretical prediction for fusion cross sections

Hot fusion for Z = 119 and 120

Towards Z=119 and 120 isotopes



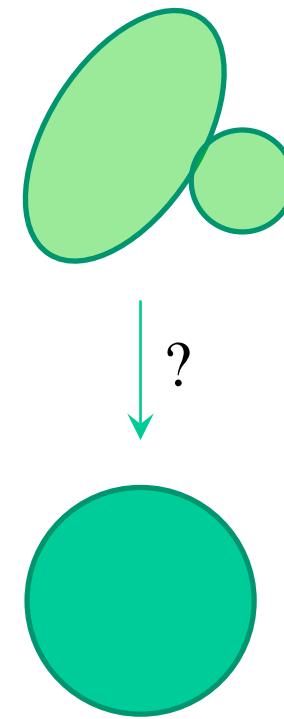
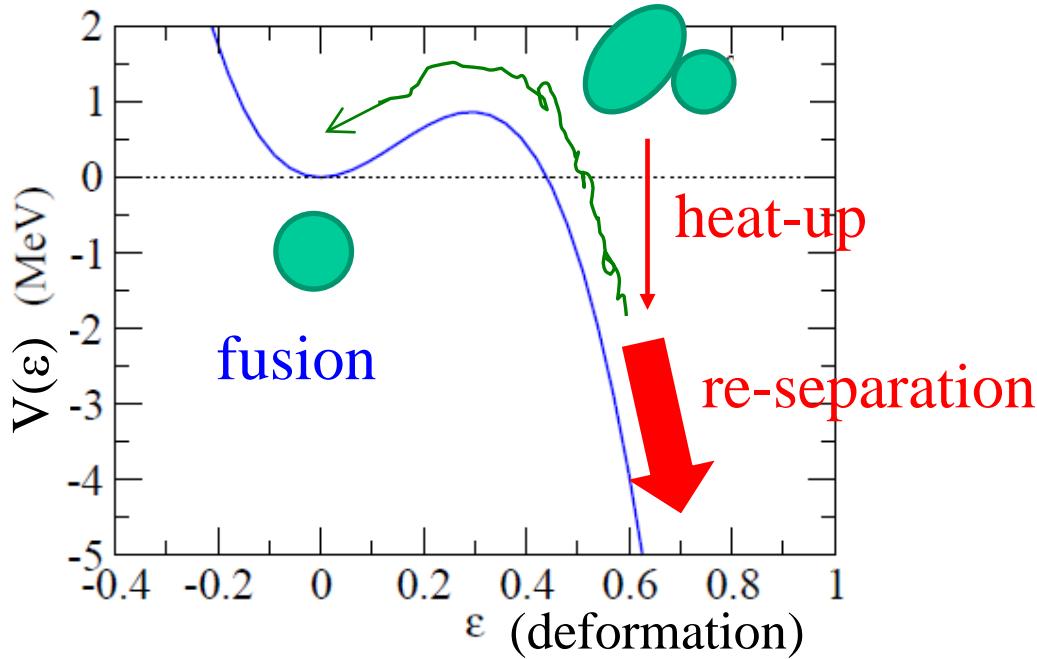
hot fusion: $^{48}\text{Ca} + \text{actinide targets}$

Dubna: $^{48}\text{Ca} + ^{249}\text{Cf}$ ($\beta_2 = 0.235$) $\rightarrow ^{297-x}\text{Og}$ (Z=118) + xn

role of deformation?

Hot fusion: $^{48}\text{Ca} +$ deformed actinide target

Effect of deformation



Open problems

- how is the shape evolved to a compound nucleus?
- Deformation: a quantum effect
 - how does the deformation disappear during heat-up?

Quantum friction

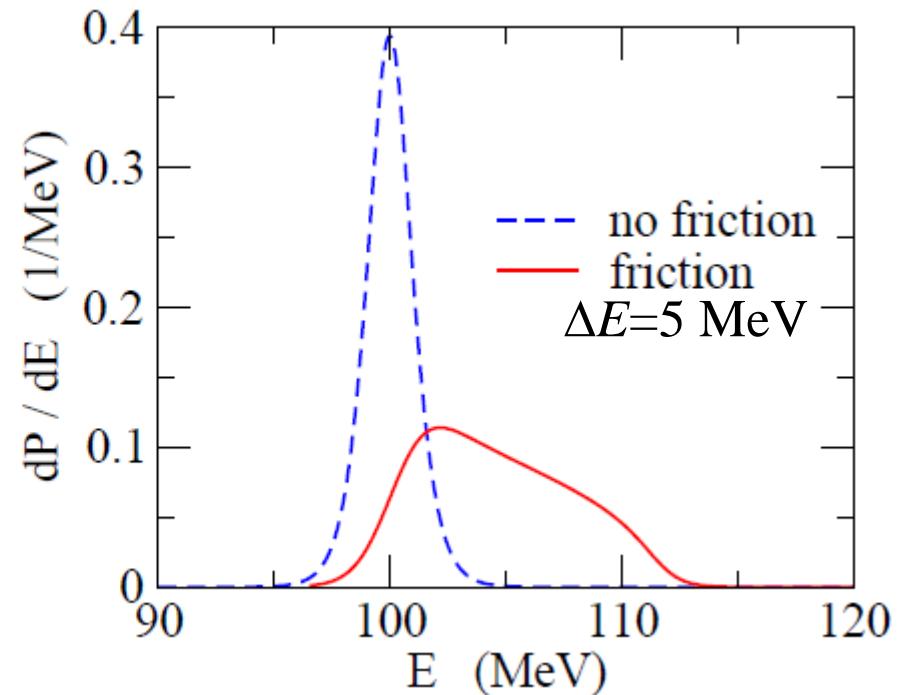
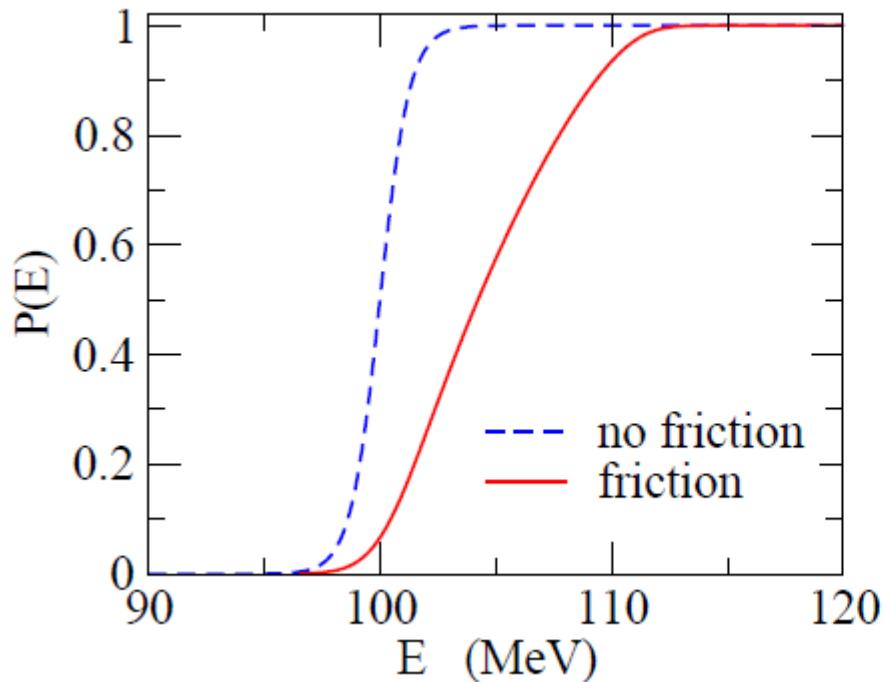
Quantum friction classical eq. of motion $\dot{p} = -V'(x) - \gamma p$

a quantization: Kanai model E. Kanai, PTP 3 (1948) 440)

$$H = \frac{p^2}{2m} + V(x) \rightarrow \frac{\pi^2}{2m} e^{-\gamma t} + e^{\gamma t} V(x) \quad (\pi = e^{\gamma t} p)$$

(a quantal Hamiltonian which reproduces the classical eq. of motion)

time-dep. wave packet approach



➤ Towards Z=119 and 120 nuclei

Another issue



the targets: not available with sufficient amounts

Dubna: ${}^{48}\text{Ca} + {}^{249}_{98}\text{Cf} \rightarrow {}^{297-x}\text{Og}$ (Z=118) + xn

${}^{249}_{98}\text{Cf}$ (351 year)

${}^{252}_{99}\text{Es}$ (471.7 day)

${}^{257}_{100}\text{Fm}$ (100.5 day)



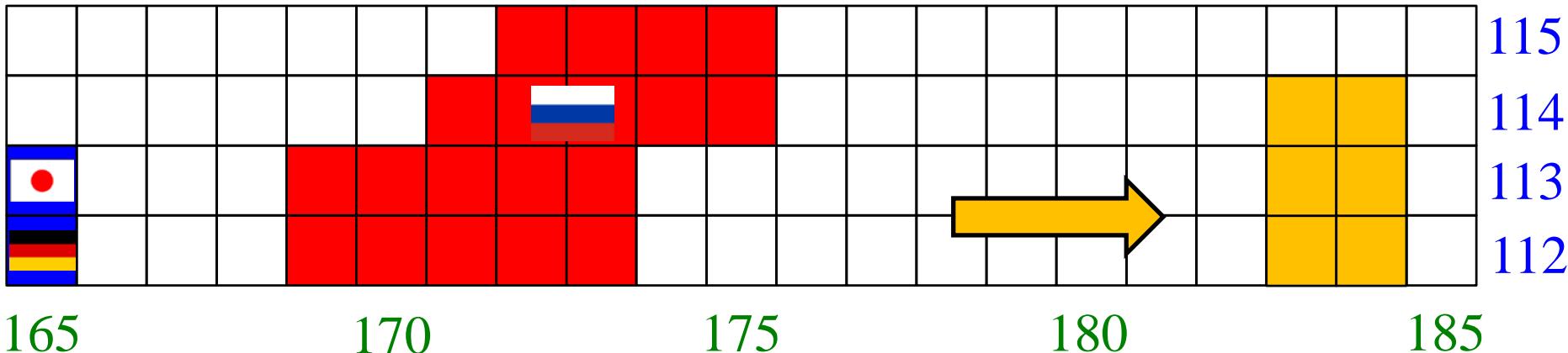
${}^{48}\text{Ca} \rightarrow {}^{50}_{22}\text{Ti}, {}^{51}_{23}\text{V}, {}^{54}_{24}\text{Cr}$ projectiles

cf. ${}^{46}_{21}\text{Sc}_{25}$: relatively small neutron number

how much will fusion cross sections be reduced?

nobody still knows

Towards the island of stability



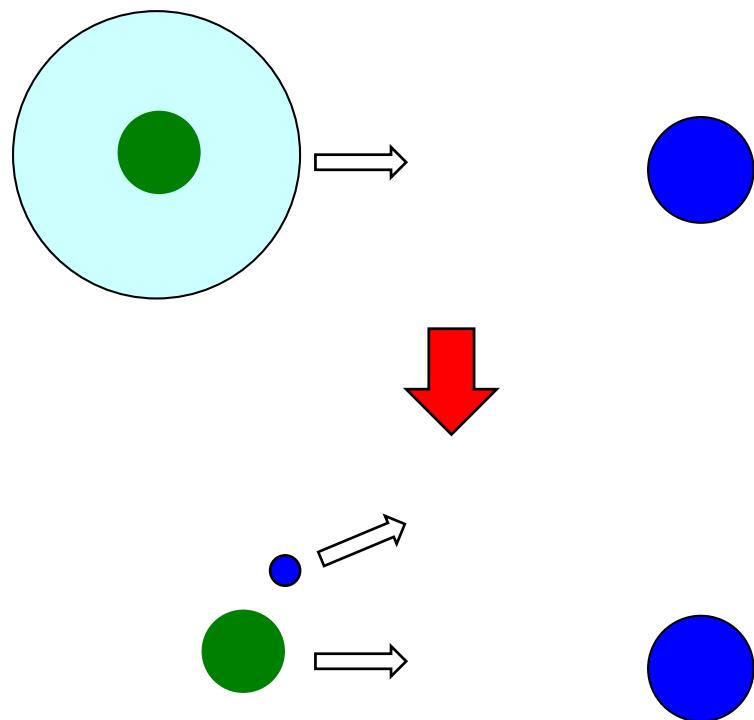
neutron-rich beams: indispensable

- how to deal with low beam intensity?
- reaction dynamics of neutron-rich beams?
 - ✓ capture: role of breakup and (multi-neutron) transfer?
 - ✓ diffusion: neutron emission during a shape evolution?
 - ✓ survival: validity of the statistical model?

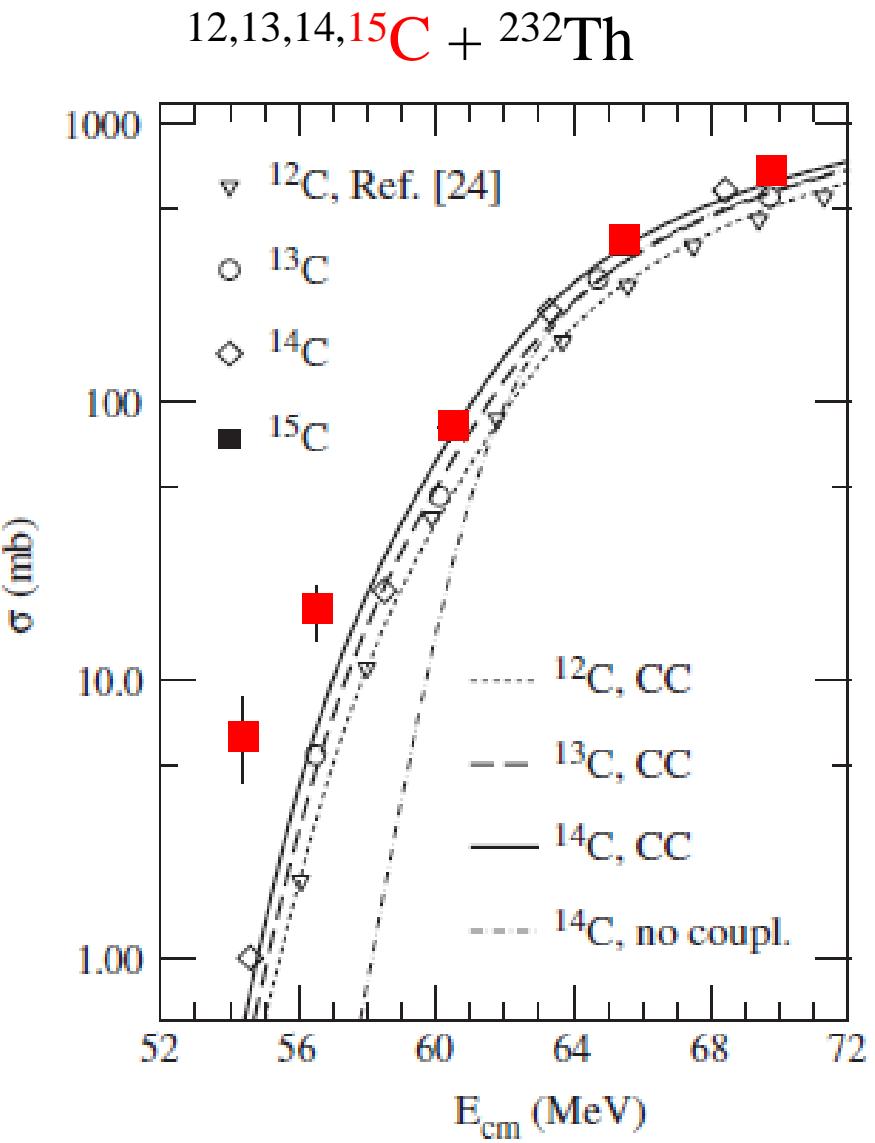
structure of exotic nuclei

more studies are required

Fusion of halo nuclei



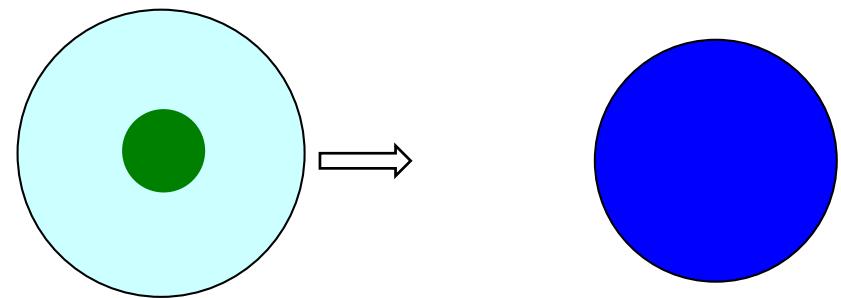
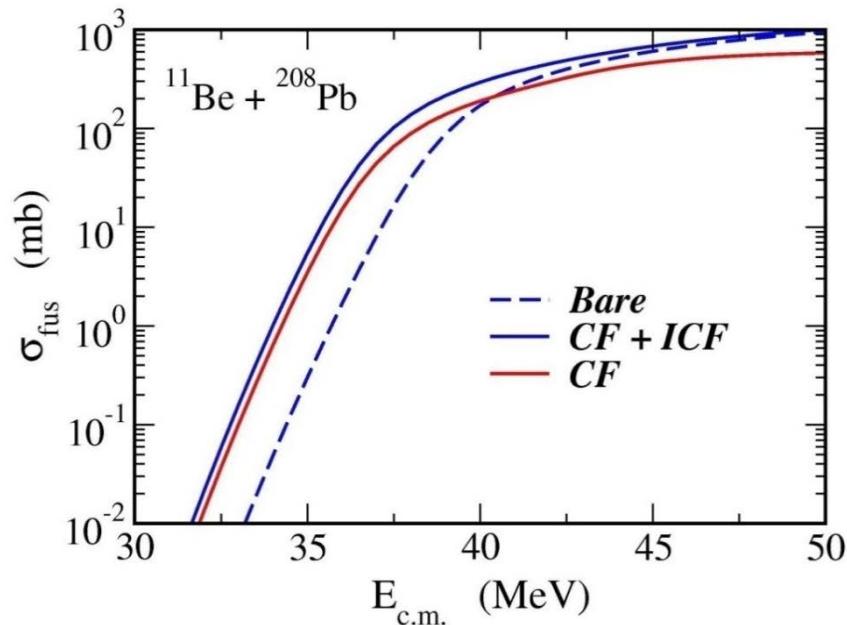
1. Lowering of potential barrier
due to a halo structure
→ enhancement
2. effect of breakup
3. effect of transfer: equally important



M. Alcorta et al.,
PRL106('11)172701

Future directions - 2

➤ Towards the island of stability
neutron-rich beams: indispensable



simultaneous treatment
of **breakup** and **transfer**

→ an important future problem

K. Hagino, A. Vitturi, C.H. Dasso,
and S.M. Lenzi, Phys. Rev. C61 ('00) 037602

Summary

Heavy-ion fusion reactions around the Coulomb barrier

- ✓ Strong interplay between nuclear structure and reaction
- ✓ Quantum tunneling with various intrinsic degrees of freedom
- ✓ coupled-channels approach

Remaining challenges

- ✓ microscopic understanding of heavy-ion fusion reactions

Future perspectives: superheavy elements

- ✓ how to reduce theoretical uncertainties?
- ✓ Towards heavier SHE ($Z = 119, 120$)
- ✓ Towards the island of stability

investigations of physics of SHE with neutron-rich nuclei as a keyword