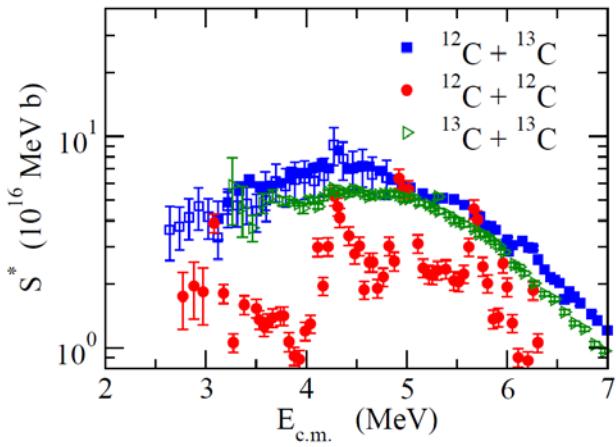


Subbarrier fusion of carbon isotopes ~ from resonance structure to fusion oscillations ~



Kouichi Hagino, *Tohoku University*
Neil Rowley, *IPN Orsay*



1. *Introduction: $^{12}\text{C} + ^{12}\text{C}$ fusion*
2. *Molecular resonances at subbarrier energies*
3. *Fusion oscillations at above barrier energies*
4. *Summary*

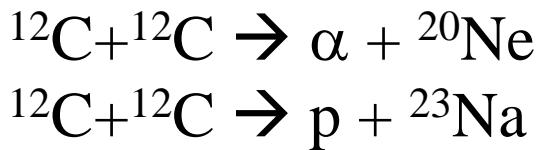
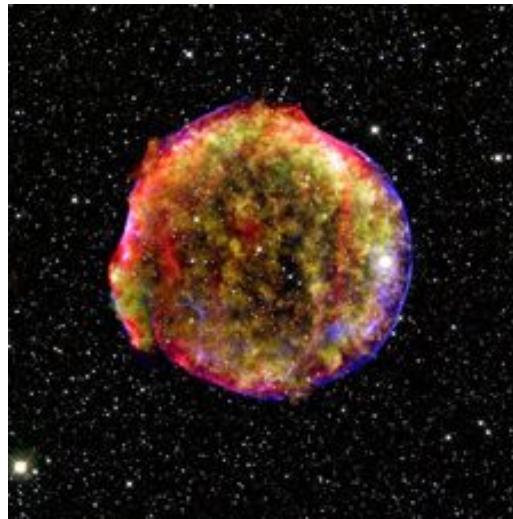
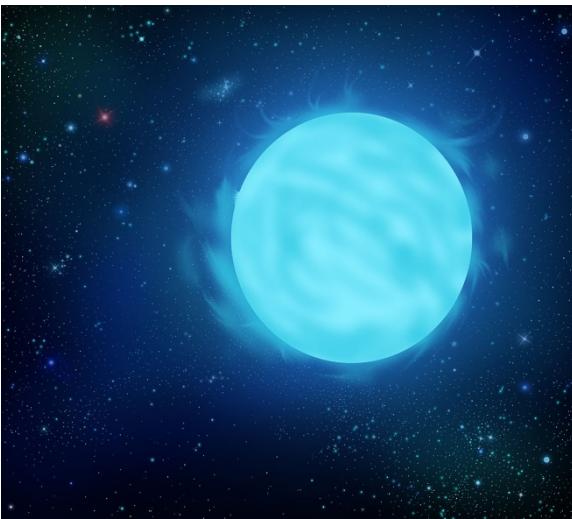
Introduction: $^{12}\text{C} + ^{12}\text{C}$ fusion

$^{12}\text{C} + ^{12}\text{C}$ fusion : a key reaction in nuclear astrophysics

Carbon burning
in massive stars

Type Ia supernovae

X-ray superburst

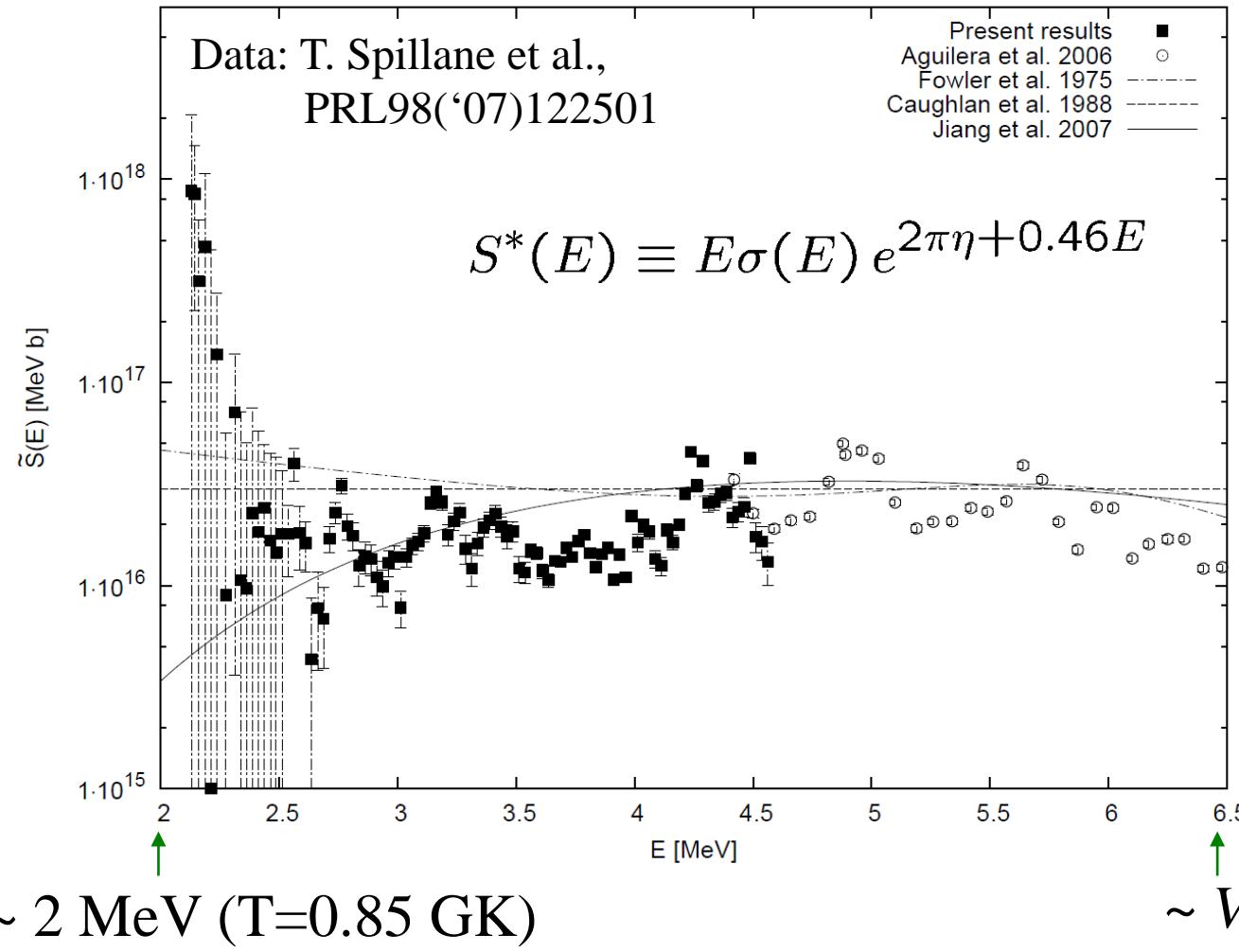


stellar evolution

deep layer of the outer
crust in accreting neutron
stars

important to understand $^{12}\text{C} + ^{12}\text{C}$ fusion at deep subbarrier energies

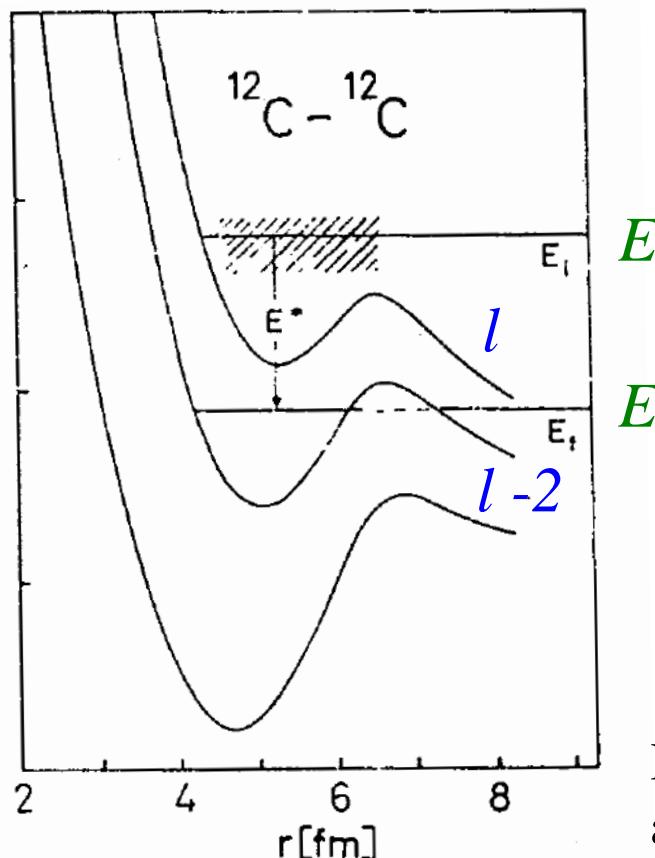
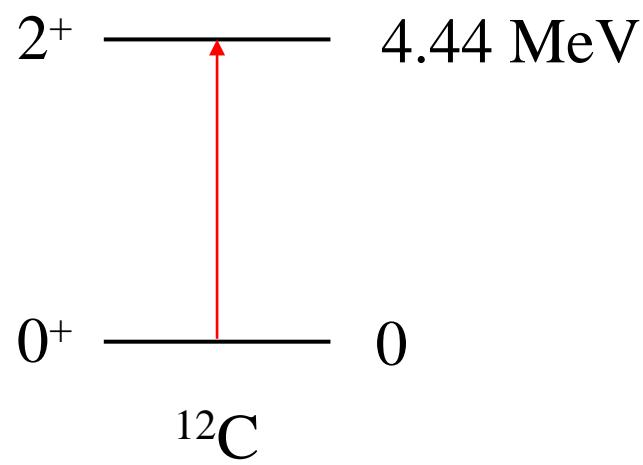
Experimental data at low energies



- ✓ pronounced resonance structures
“molecular resonances”
- ✓ difficult to extrapolate down to E_G

Theoretical calculations

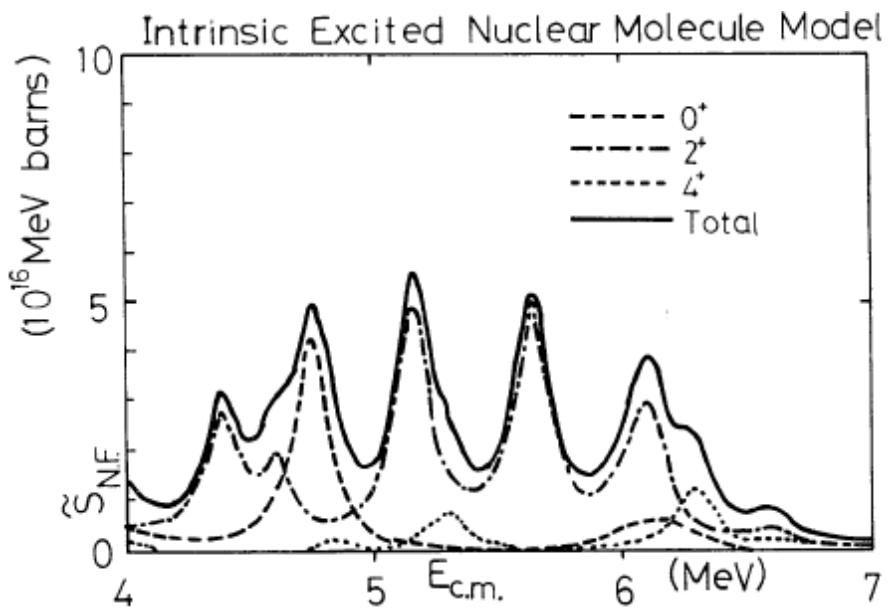
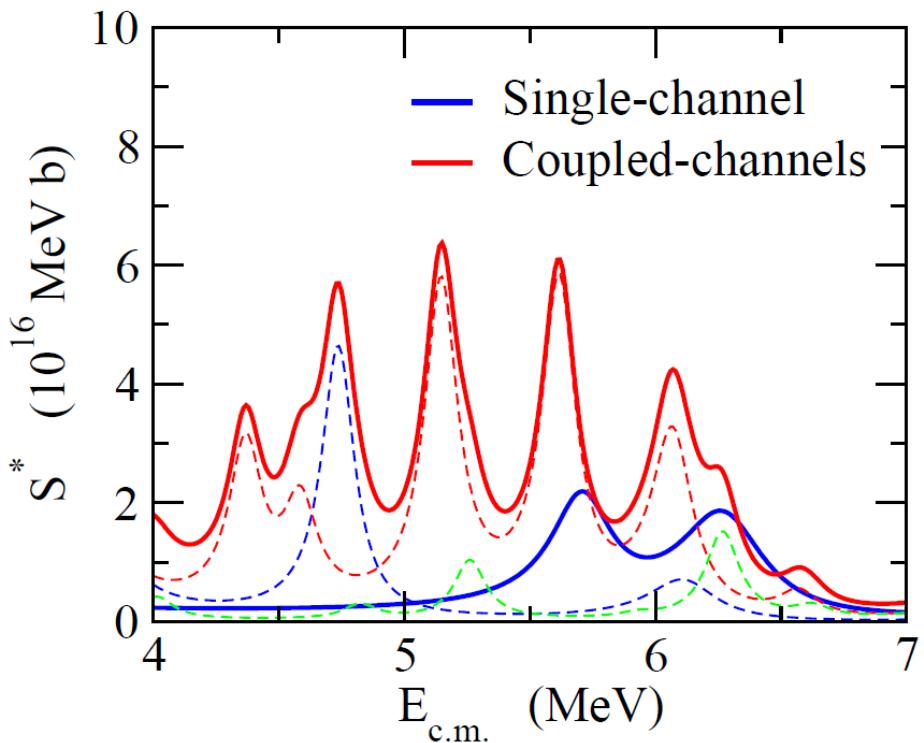
- Nogami-Imanishi model (B. Imanishi, PL 27B ('68) 267, NPA125 ('69) 33)
- Band-crossing model (Y. Kondo, T. Matsuse, Y. Abe, PTP59 ('78) 465)
- Double resonance model (W. Scheid, W. Greiner, R. Lemmer, PRL25 ('70) 176)
* the basic concept is all same



H.-J. Fink, W. Scheid,
and W. Greiner,
NPA188 ('72) 259

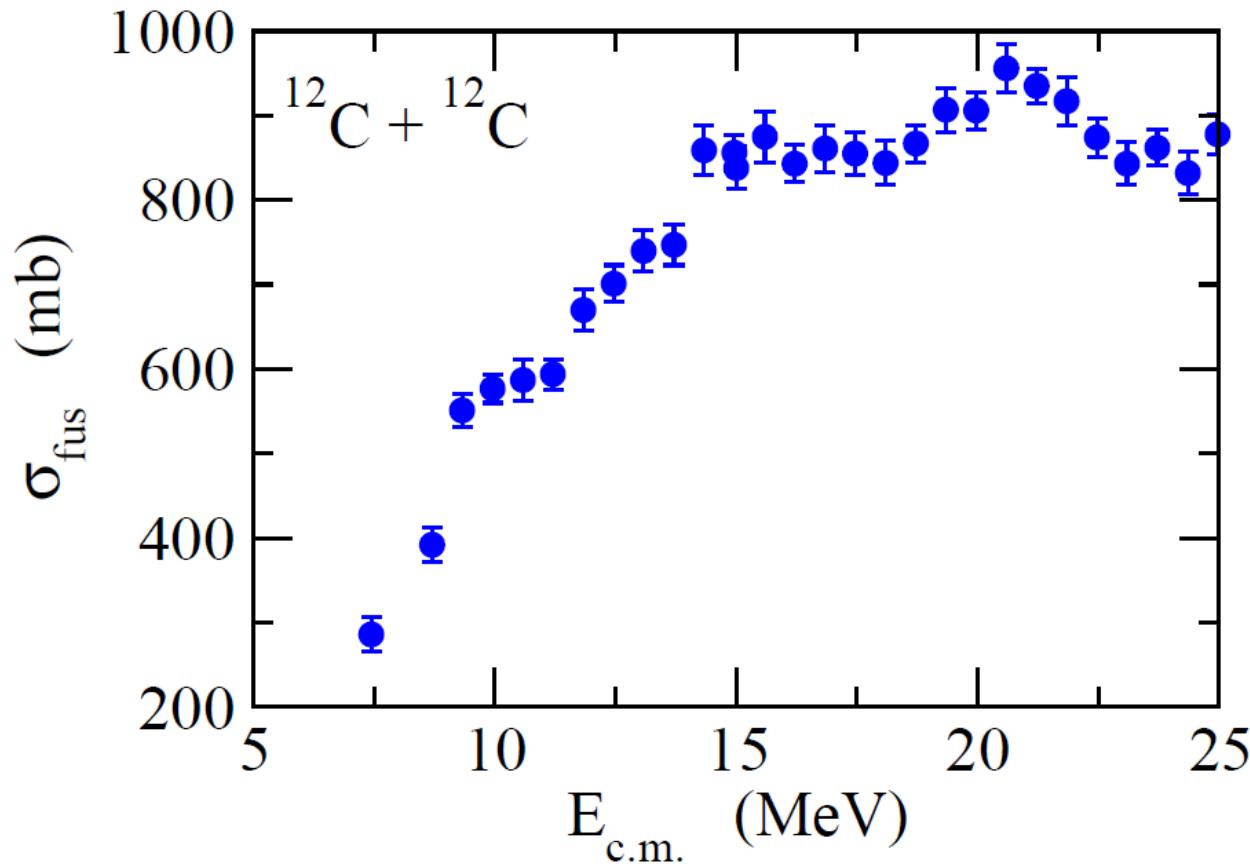
Theoretical calculations

- Nogami-Imanishi model (B. Imanishi, PL 27B ('68) 267, NPA125 ('69) 33)
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Y. Kondo, T. Matsuse, and Y. Abe,
PTP 59 ('78) 465

Experimental data at above barrier energies



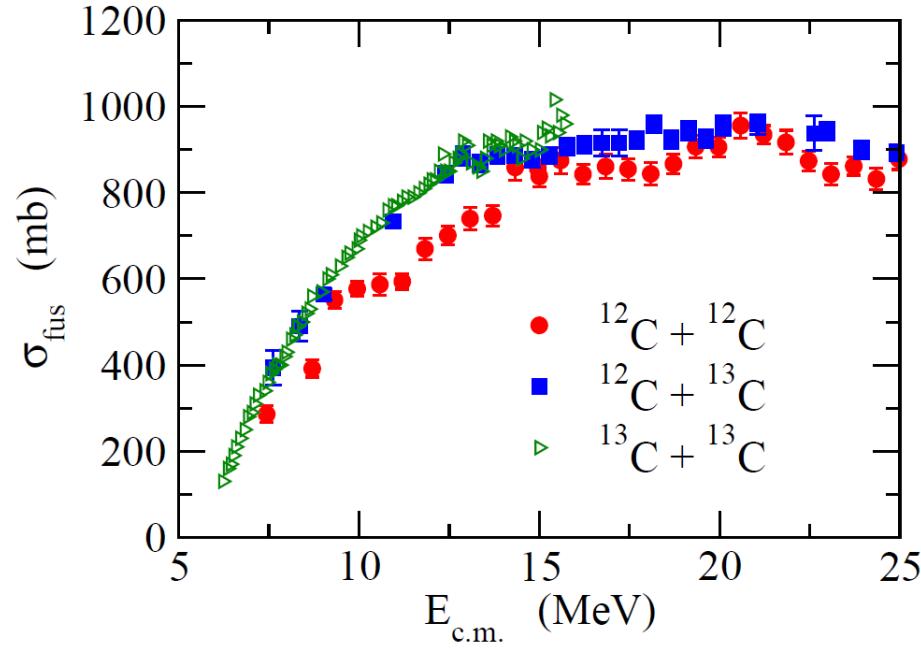
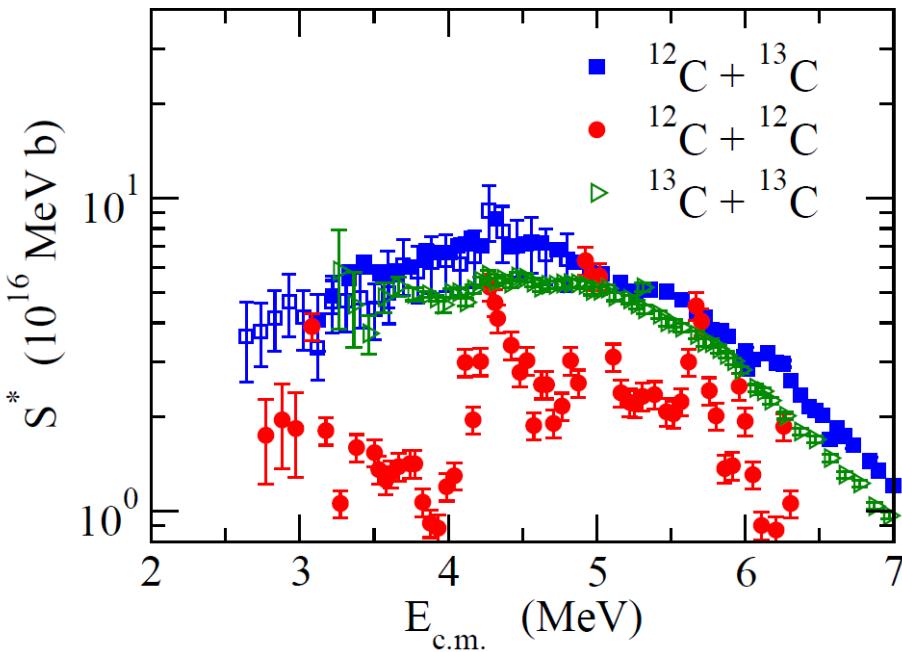
Data: D.G. Kovar et al., PRC20 ('79) 1305

✓ fusion oscillations

← successive contributions of individual partial waves

(N. Poffe, N. Rowley, and R. Lindsay, NPA410 ('83) 498)

Comparison with other C+C systems



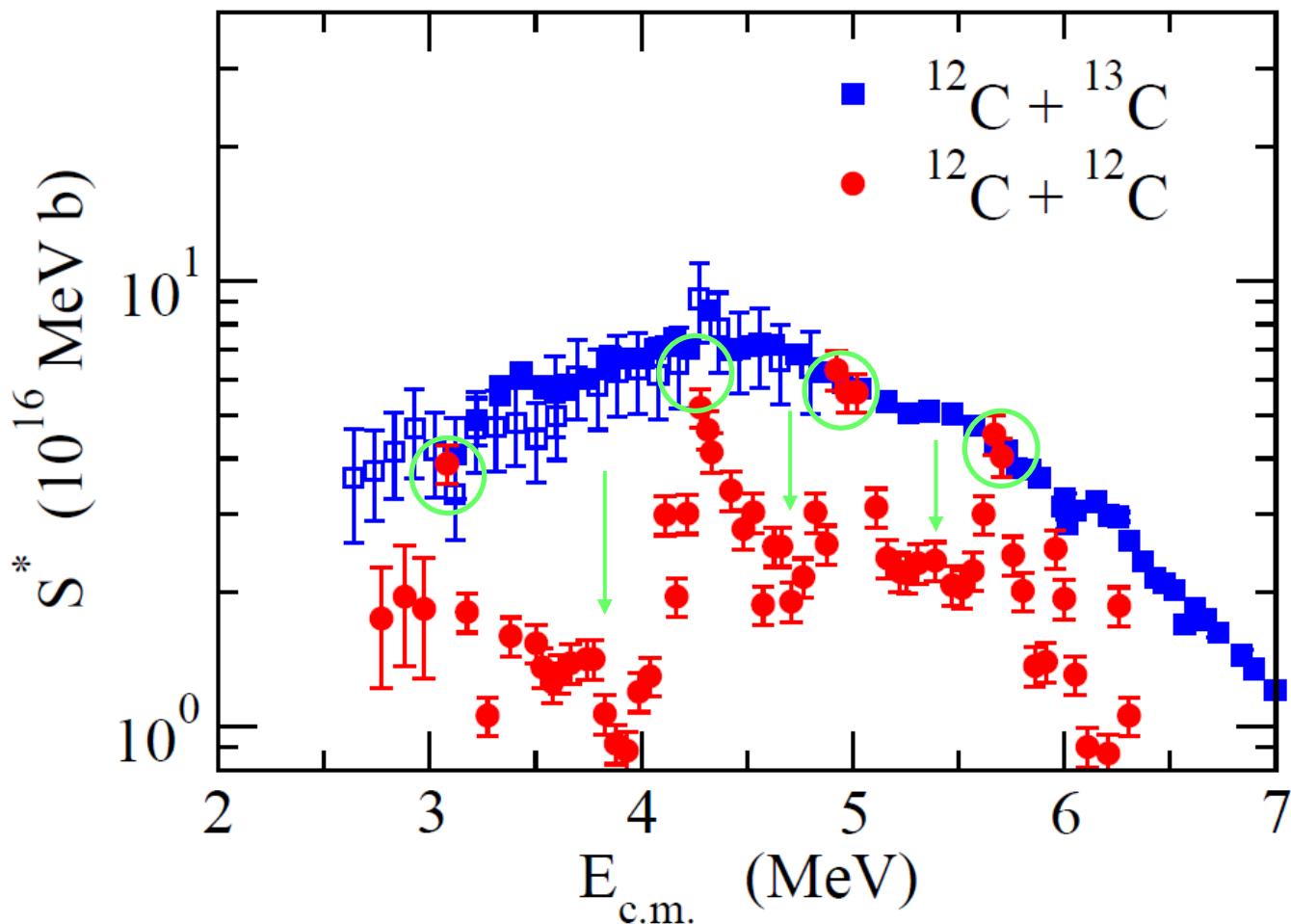
fusion cross sections for $^{12}\text{C} + ^{13}\text{C}$, $^{13}\text{C} + ^{13}\text{C}$: much less structured

How can one understand the systematics?

- from $^{12}\text{C} + ^{12}\text{C}$ to $^{12}\text{C} + ^{13}\text{C}$, $^{13}\text{C} + ^{13}\text{C}$
- origins for the resonances/oscillations?
- from low to high energies

cf. most of the previous studies: $^{12}\text{C} + ^{12}\text{C}$ only

Molecular resonances at subbarrier energies



M. Notani, X.D. Tang
et al.,
PRC85('12)014607

off-resonance: fusion inhibition
on-resonance: match with $^{12}\text{C} + ^{13}\text{C}$

properties of compound nucleus (^{24}Mg)?

$^{12}\text{C} + ^{12}\text{C}$ reaction:

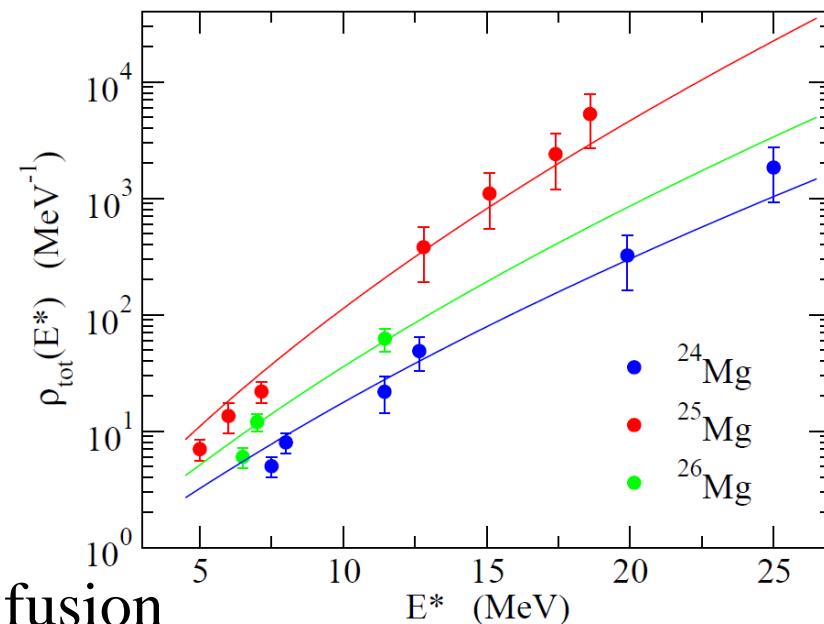
- ✓ level density of ^{24}Mg : small (e-e)
- ✓ small fusion Q-value

$$Q = +13.9 \text{ MeV } (^{12}\text{C} + ^{12}\text{C})$$

$$+16.3 \text{ MeV } (^{12}\text{C} + ^{13}\text{C})$$

$$+22.5 \text{ MeV } (^{13}\text{C} + ^{13}\text{C})$$

→ small E^* for ^{24}Mg in $^{12}\text{C} + ^{12}\text{C}$ fusion



→ $\sigma \sim \sum_J \sigma_{\text{cap}}^J \left[1 - e^{-2\pi \Gamma_J / D_J} \right]$

large hindrance factor

$$D_J = 1/\rho_J$$

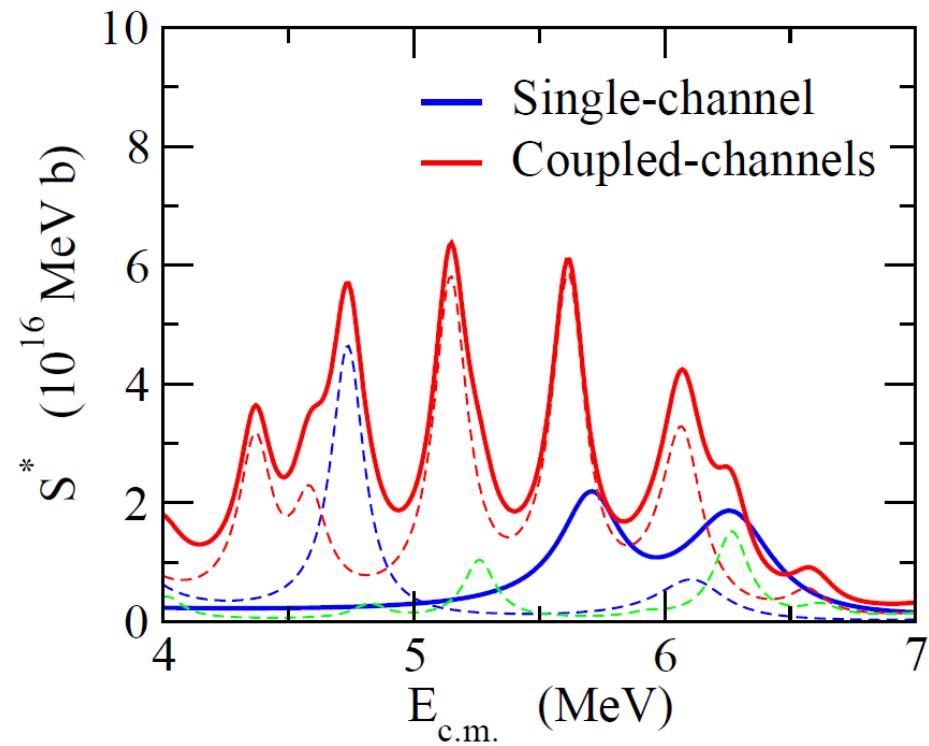
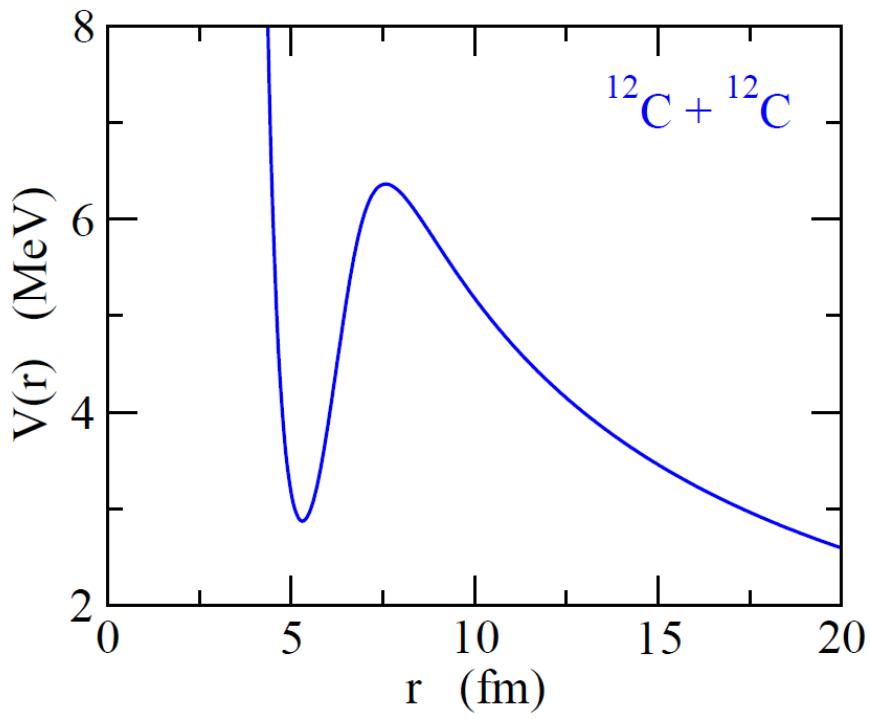
$$\Gamma_J : \text{CN width}$$

incorporate this idea in the coupled-channels calculations?

C.C. calculations with level-density-dependent imaginary potential

^{12}C - ^{12}C potential (Kondo, Matsuse, Abe, PTP('78))

- ✓ two-range Woods-Saxon + Coulomb for the real part
- ✓ a Woods-Saxon for the imaginary part



C.C. calculations with level-density-dependent imaginary potential

^{12}C - ^{12}C potential (Kondo, Matsuse, Abe, PTP('78))

- ✓ two-range Woods-Saxon + Coulomb for the real part
- ✓ a Woods-Saxon for the imaginary part


$$W(r) = -W_0 \cdot f_{\text{WS}}(r) \rightarrow -w_0 \rho_J(E^*) \cdot f_{\text{WS}}(r)$$

G. Helling, W. Scheid, W. Greiner, PL 36B ('71) 64

H.-J. Fink, W. Scheid, W. Greiner, NPA188 ('72) 259

J.M. Quesada, M. Lozano, G. Madurga, PLB125 ('83) 14

M.V. Andres, Quesada, Lozano, Madurga, NPA443 ('85) 380

- ✓ E and J dependent imaginary potential
- ✓ system dependence through $\rho(E)$

cf. Fermi's golden rule

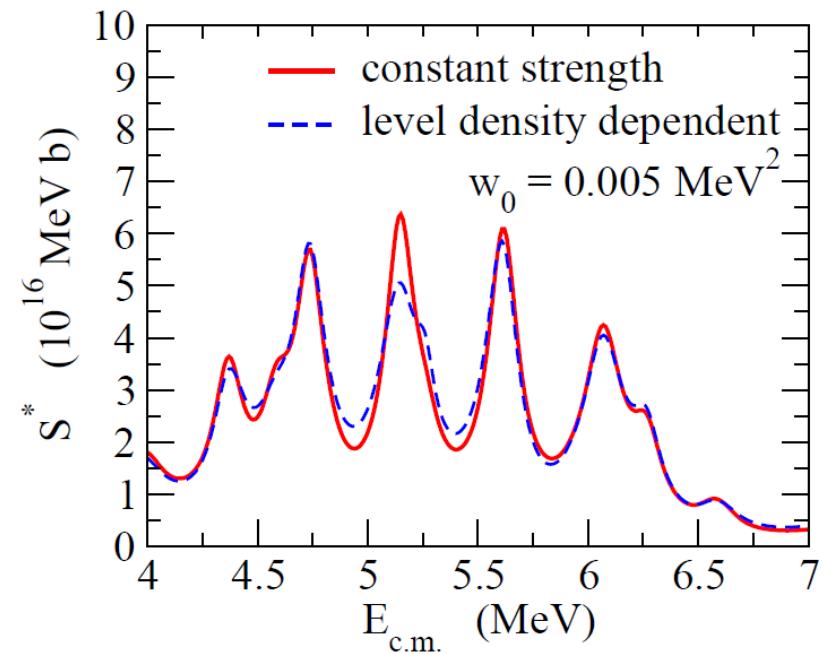
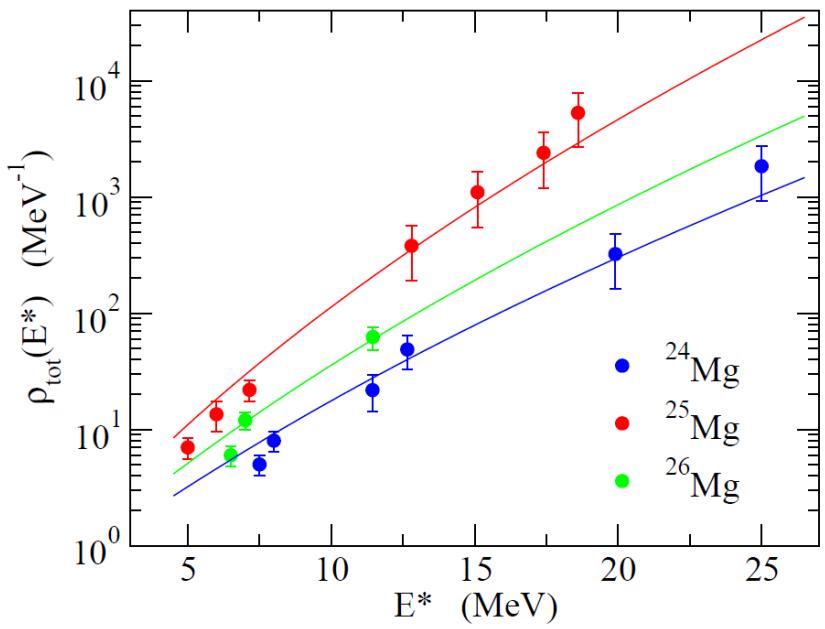
$$\frac{dw}{dt} = \frac{2\pi}{\hbar} |\langle \psi_{\text{CN}} | V_{\text{int}} | \psi_{\text{elastic}} \rangle|^2 \rho_J(E^*)$$

C.C. calculations with level-density-dependent imaginary potential

$^{12}\text{C}-^{12}\text{C}$ potential (Kondo, Matsuse, Abe, PTP('78))

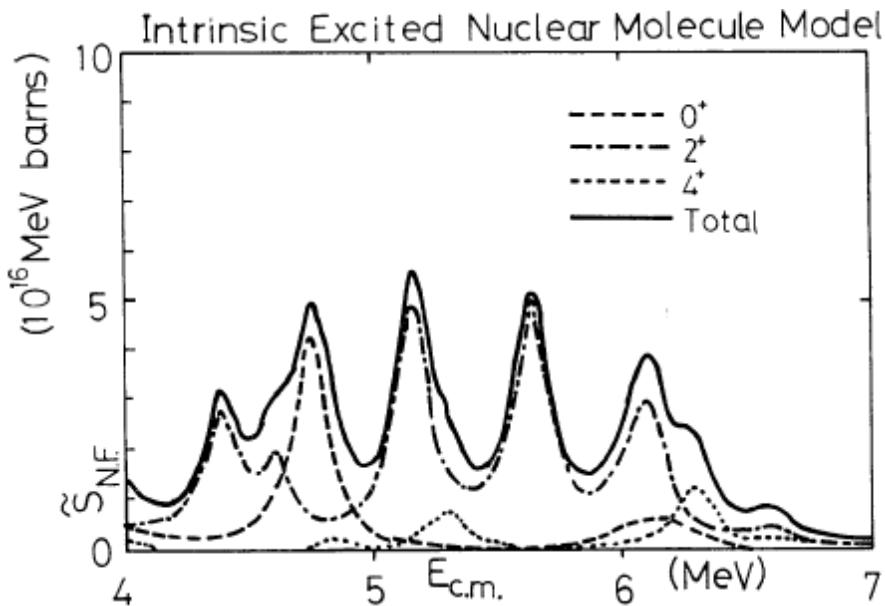
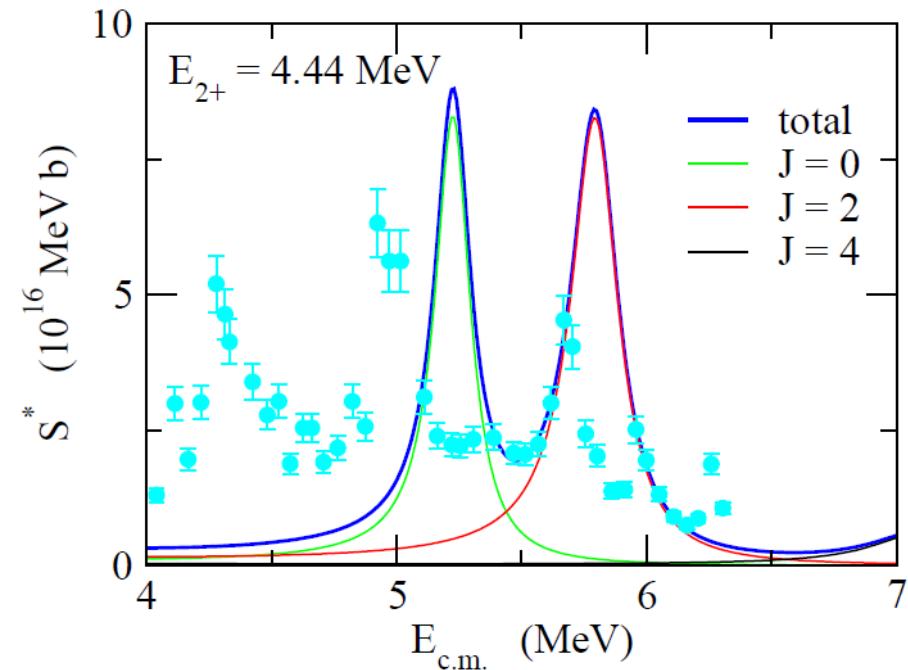
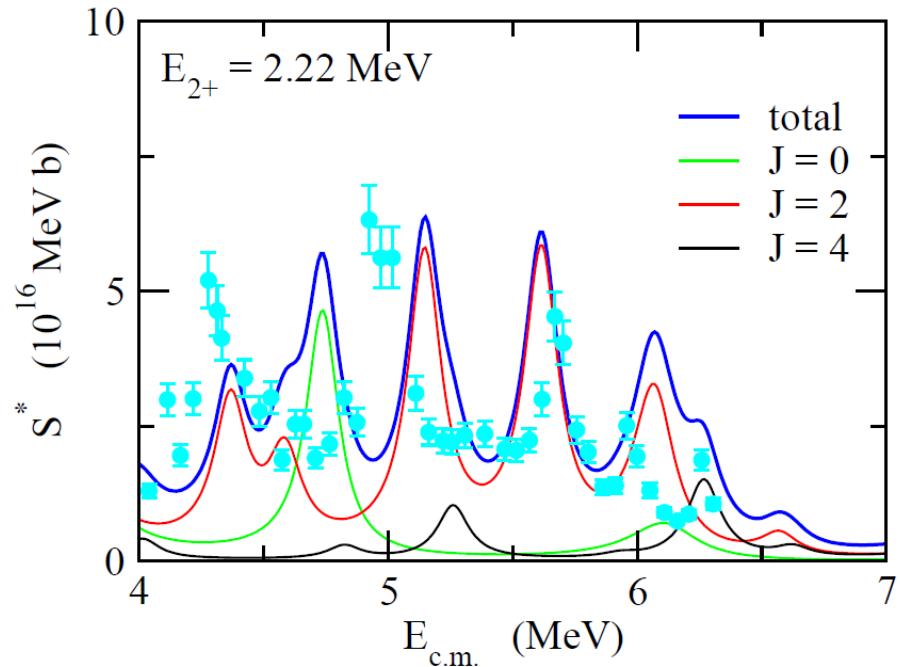
- ✓ two-range Woods-Saxon + Coulomb for the real part
- ✓ a Woods-Saxon for the imaginary part

→
$$W(r) = -W_0 \cdot f_{\text{WS}}(r) \rightarrow -w_0 \rho_J(E^*) \cdot f_{\text{WS}}(r)$$



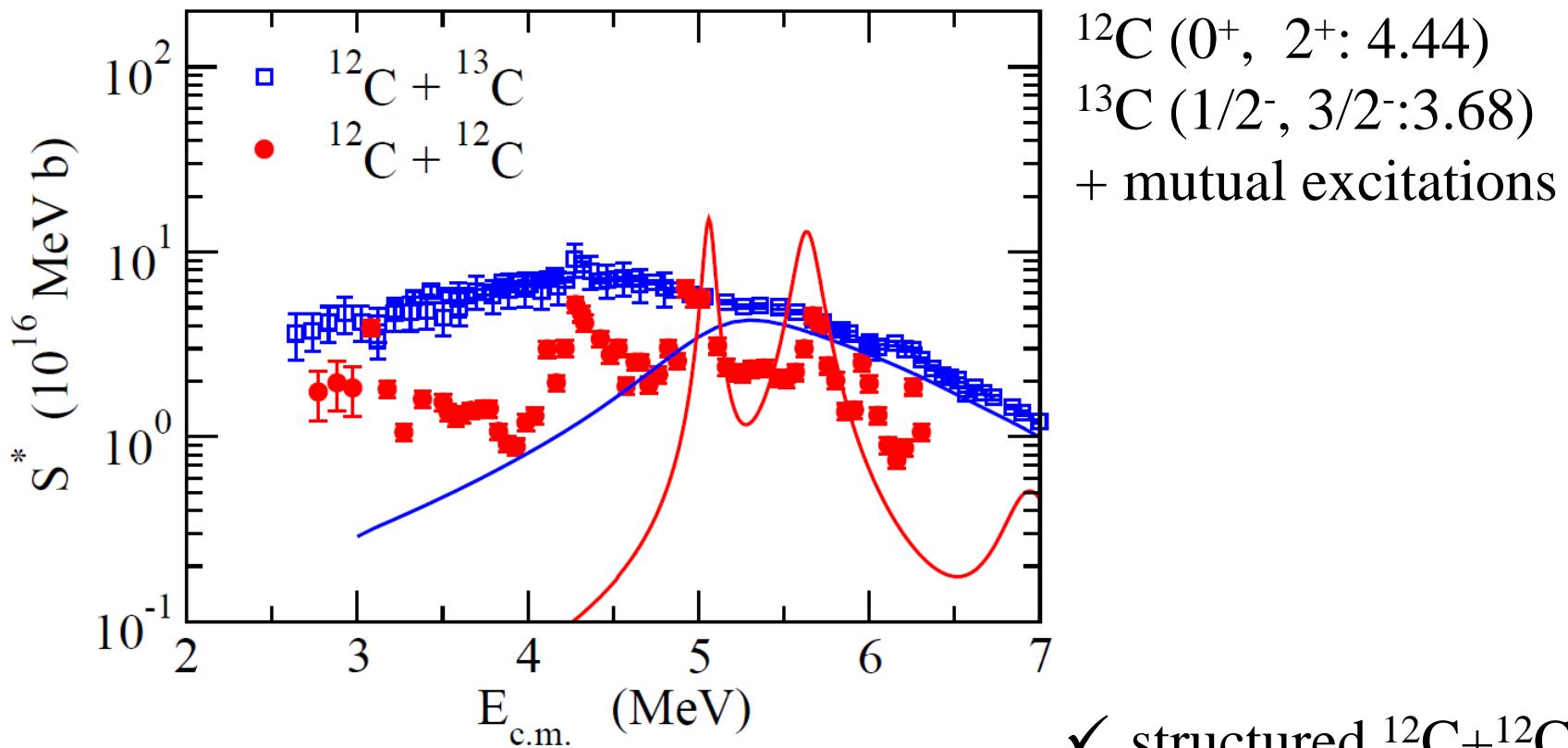
$$\rho_J(E^*) = \frac{(2J+1)e^{-(J+1/2)^2/2\sigma^2}}{4\sigma^3\sqrt{2\pi}} \frac{\sqrt{\pi}}{12} \frac{e^{2\sqrt{aE^*}}}{a^{1/4}(E^*)^{5/4}} \quad \left(\sigma^2 = 0.088 a A^{2/3} \sqrt{\frac{E^*}{a}} \right)$$

some mystery?



use slightly different input parameters from Kondo-Matsuse-Abe, but keep the physical value of E_{2+} ($= 4.44$ MeV)

Results of coupled-channels calculations



system dependence: qualitatively reproduced

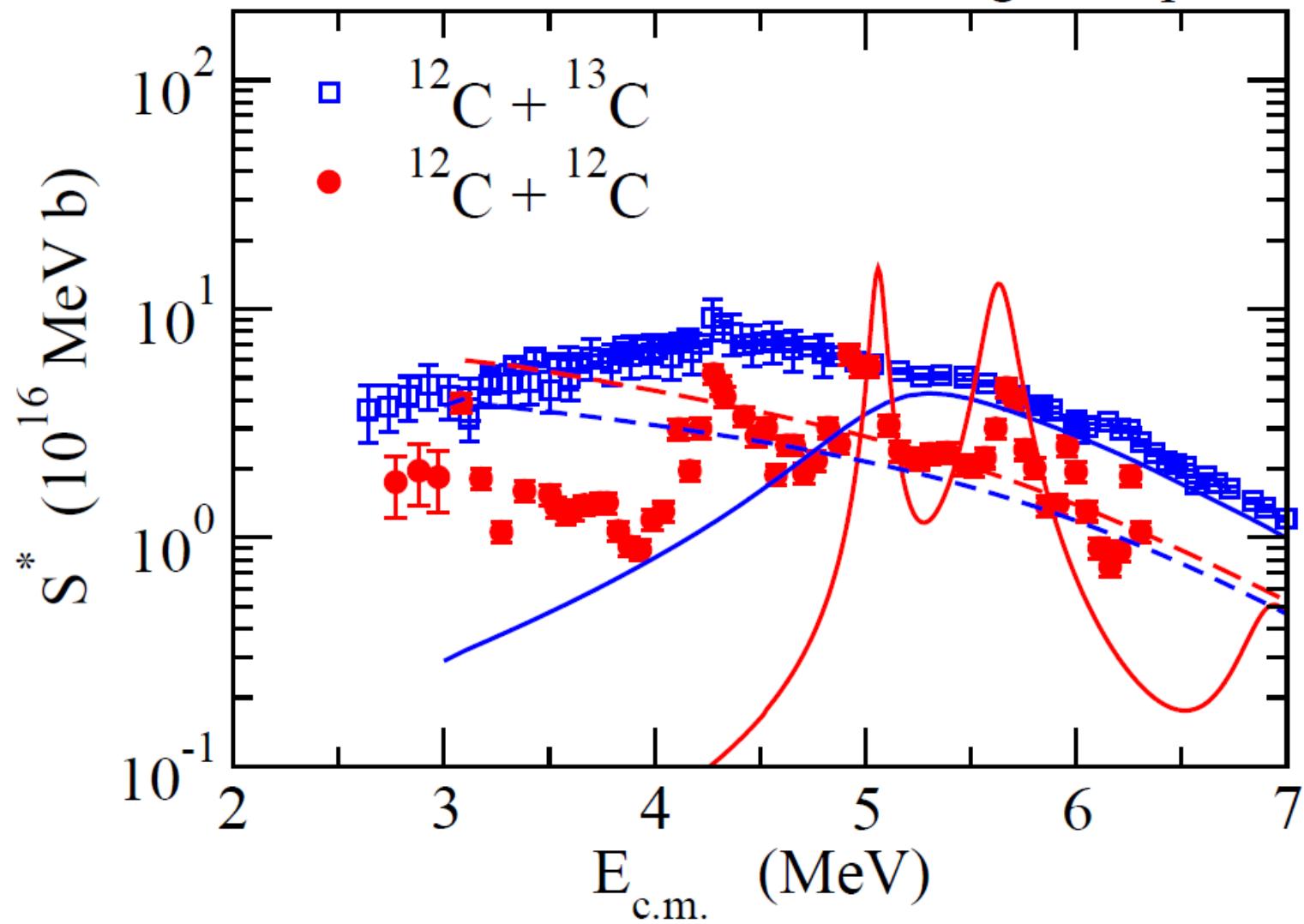
- ✓ structured $^{12}\text{C} + ^{12}\text{C}$
- ✓ smooth $^{12}\text{C} + ^{13}\text{C}$

underestimate of fusion cross sections at deep subbarrier energies:
→ couplings to 3^- and 0_2^+ (Hoyle state)?

cf. role of Hoyle state in $^{12}\text{C} + ^{12}\text{C}$:

M. Assuncao and P. Descouvemont, PLB723 ('13) 355

Dashed curves: strong absorption



Fusion oscillations at above barrier energies

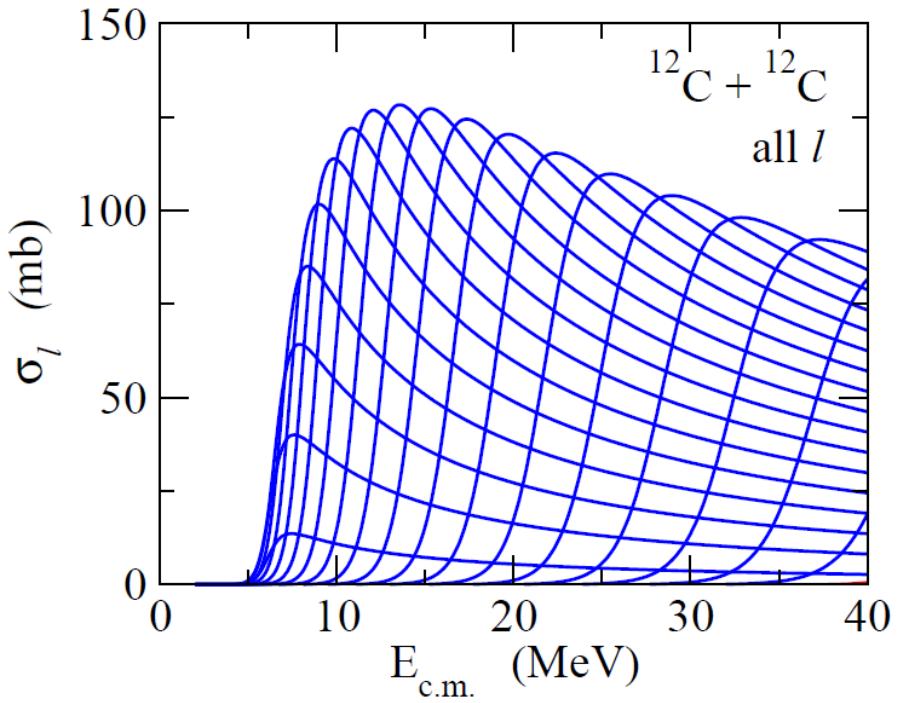
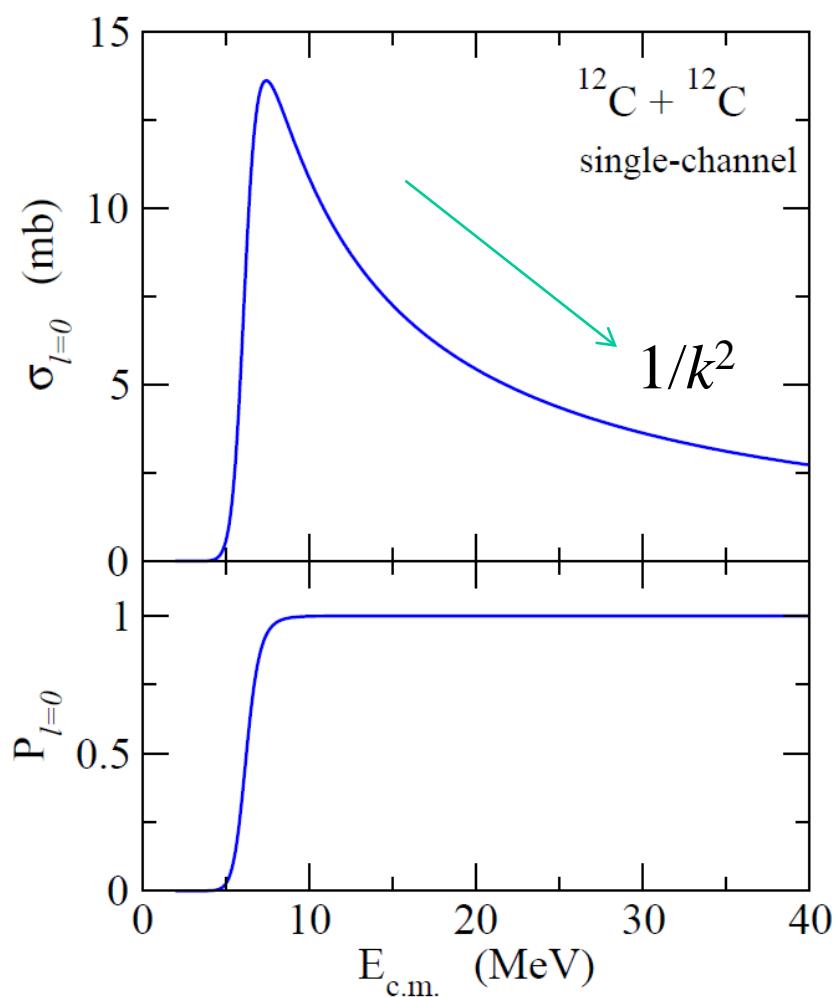
high- E : high level density of CN \longrightarrow overlapping resonances
 \longrightarrow strong absorption

$$\sigma_{\text{fus}}(E) = \frac{\pi}{k^2} \sum_l (2l + 1) P_l(E)$$

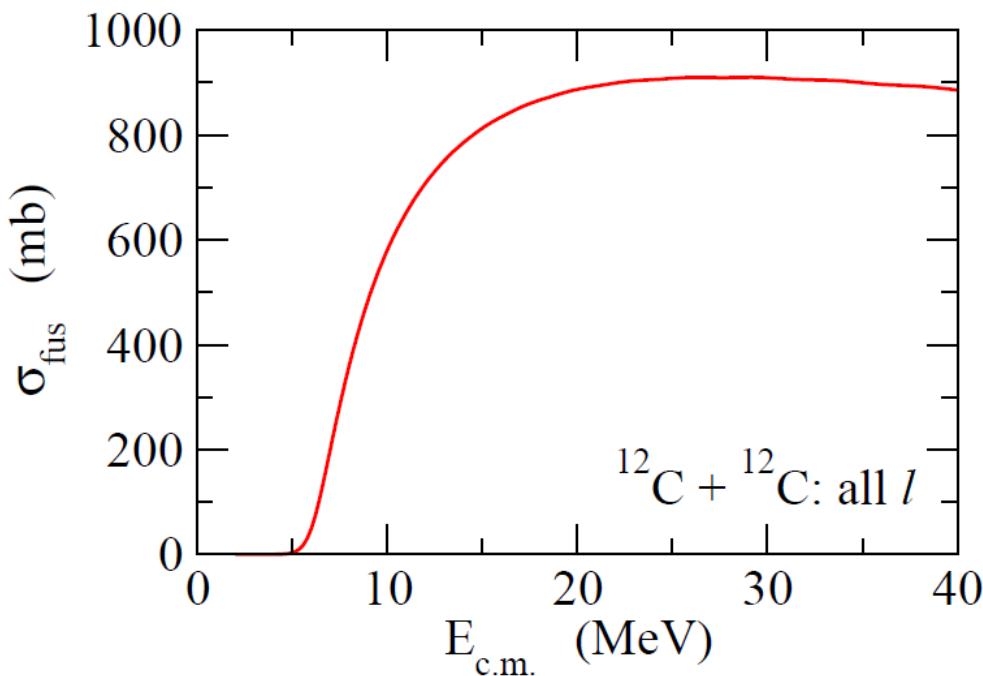
$P_l(E)$: barrier penetrability

Fusion oscillations at above barrier energies

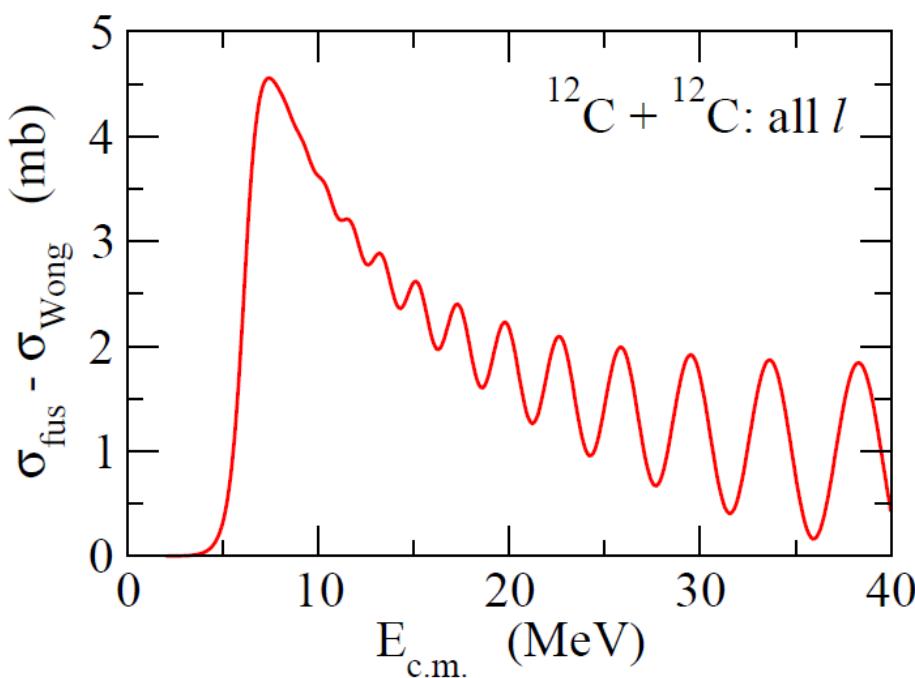
$$\sigma_{\text{fus}}(E) = \frac{\pi}{k^2} \sum_l (2l + 1) P_l(E)$$



discrete l -sum
→ (oscillatory) structure in
 σ_{fus}



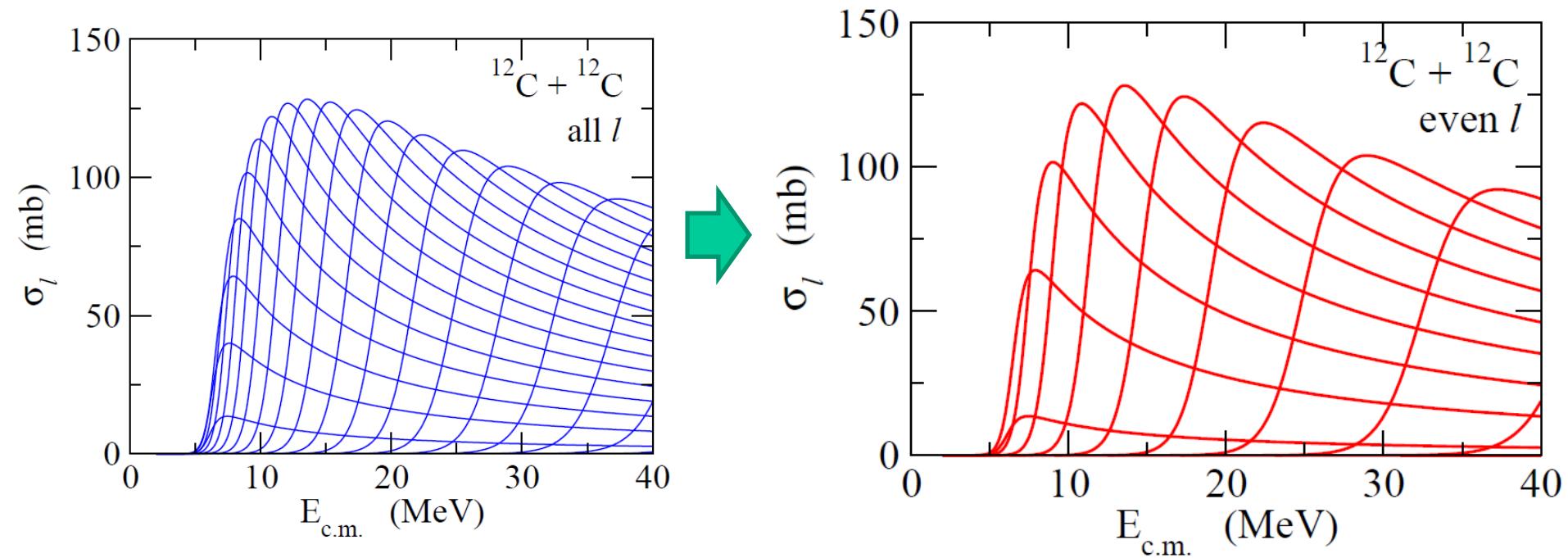
* practically, the oscillations
are invisible ($\Delta\sigma \sim 1$ mb)
if all- l are summed



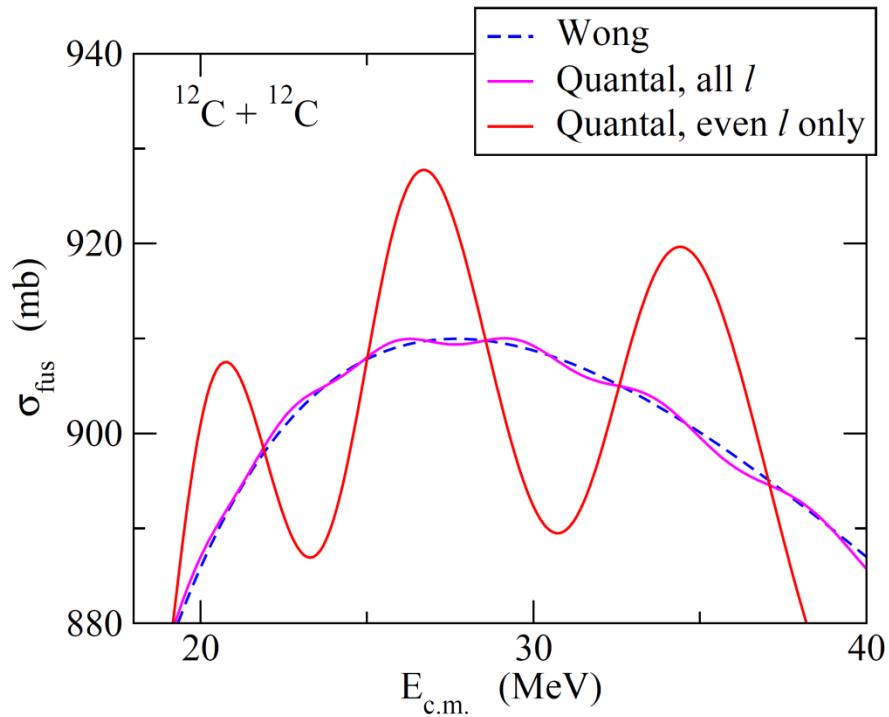
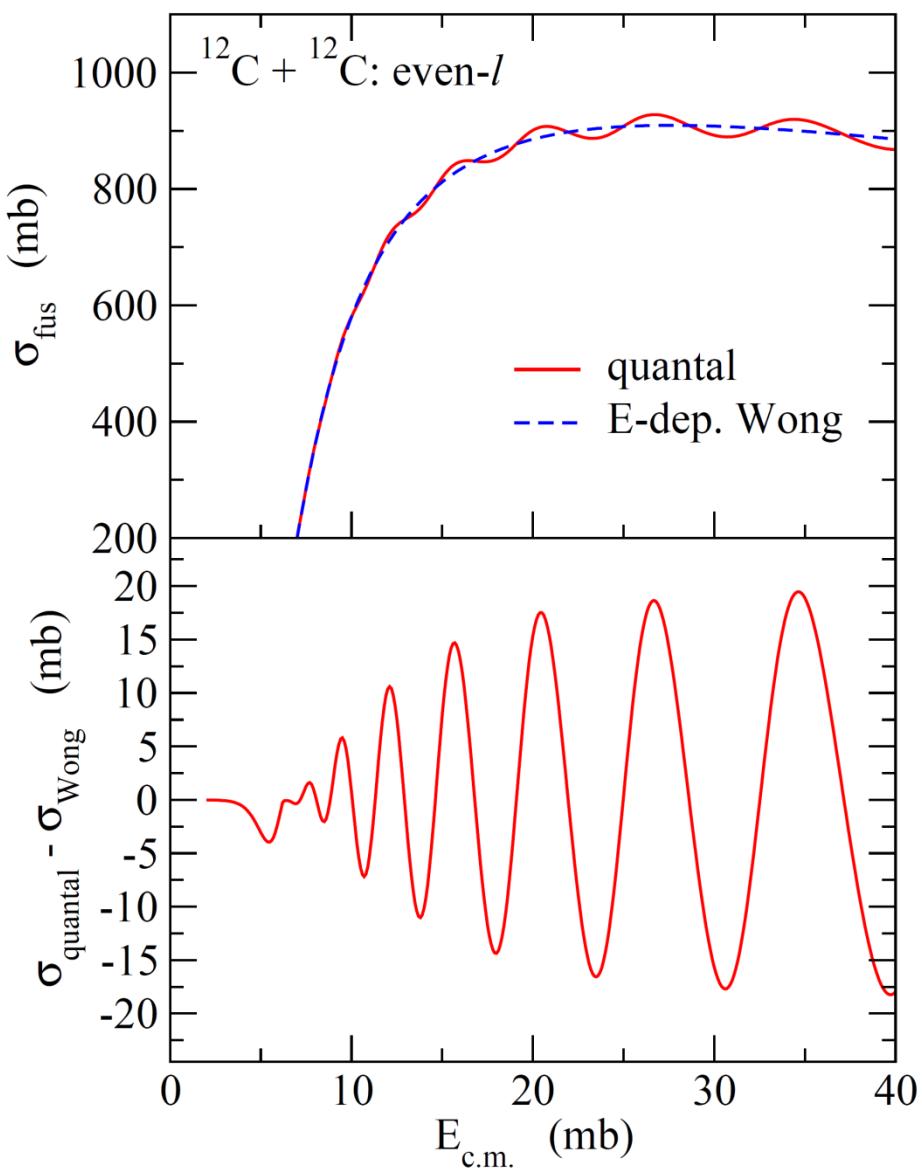
effect of symmetrization: fusion oscillations in light symmetric systems

fusion of identical spin-zero bosons: wf has to be symmetric

$$\sigma_{\text{fus}}(E) = \frac{\pi}{k^2} \sum_l (2l+1) P_l(E) \rightarrow \frac{\pi}{k^2} \sum_l (1 + (-)^l)(2l+1) P_l(E)$$



- ✓ the angular mom. is quantized in units of $2\text{-}\hbar$
- ✓ a larger amplitude of fusion oscillations



Analytic formula for fusion oscillations

N. Poffe, N. Rowley, and R. Lindsay, Nucl. Phys. A410 ('83) 498

N. Rowley and K. Hagino, in preparation

$$\sigma_{\text{fus}}(E) = \frac{\pi}{k^2} \sum_l (1 \pm (-)^l)(2l+1) P_l(E)$$

$$\sim \sigma_{\text{E-Wong}} \pm 2\pi R_E^2 \frac{\hbar\Omega_E}{E} e^{-\xi} \sin(\pi l_g)$$

← Poisson
sum
formula

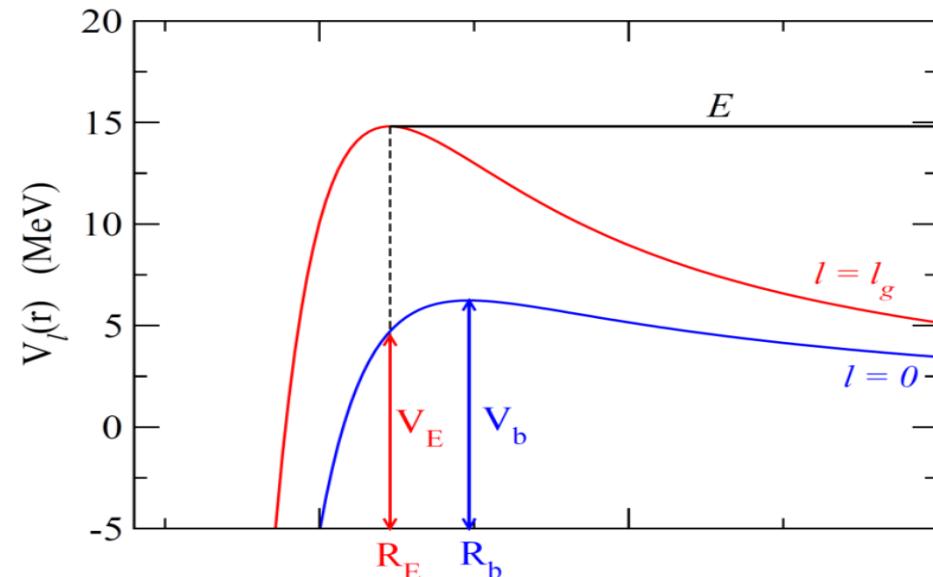
$$\xi = \pi \cdot \frac{\hbar\Omega_E}{2l_g + 1} \cdot \frac{\mu R_E^2}{\hbar^2}$$

(note)

↻ $2l_g + 1 \gg \pi \hbar \Omega_E \cdot \frac{\mu R_E^2}{\hbar^2}$

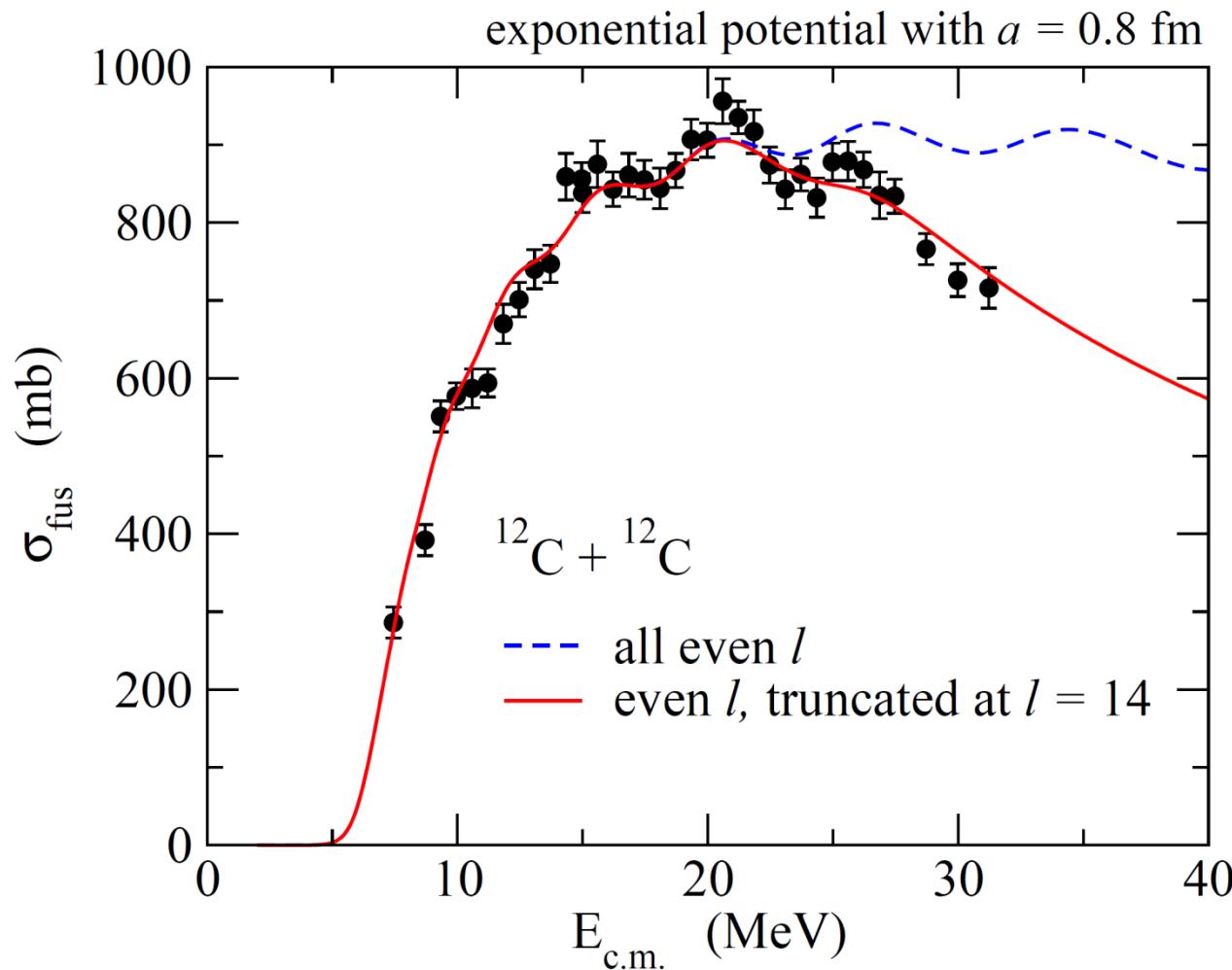
in order for the osc. to be visible

→ light symmetric systems



i) Comparison with the experimental data: $^{12}\text{C} + ^{12}\text{C}$

$^{12}\text{C}_{\text{g.s.}} : 0^+ \rightarrow$ the relative w.f. has to be spatially symmetric



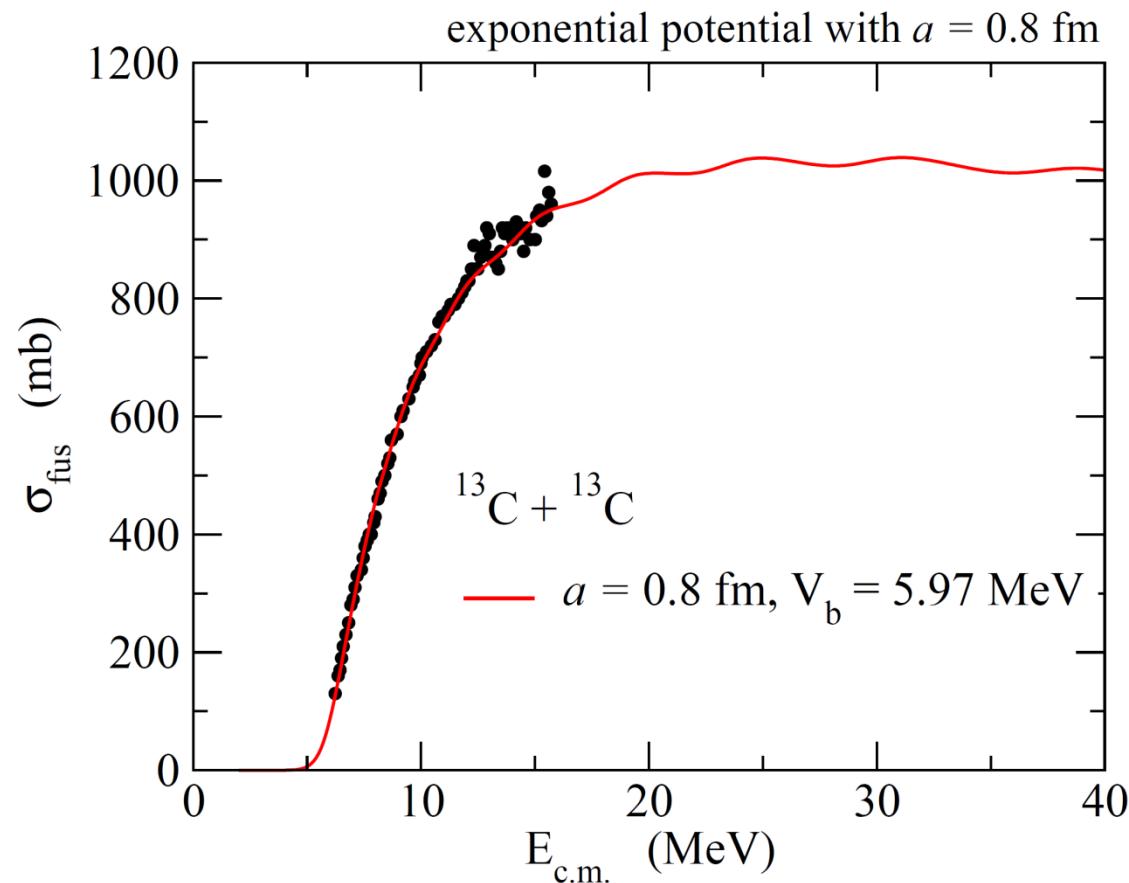
ii) $^{13}\text{C} + ^{13}\text{C}$

$^{13}\text{C}_{\text{g.s.}} : 1/2^- \rightarrow$ the relative w.f. has to be spatially symmetric for $S = 0$
 spatially anti-symmetric for $S = 1$

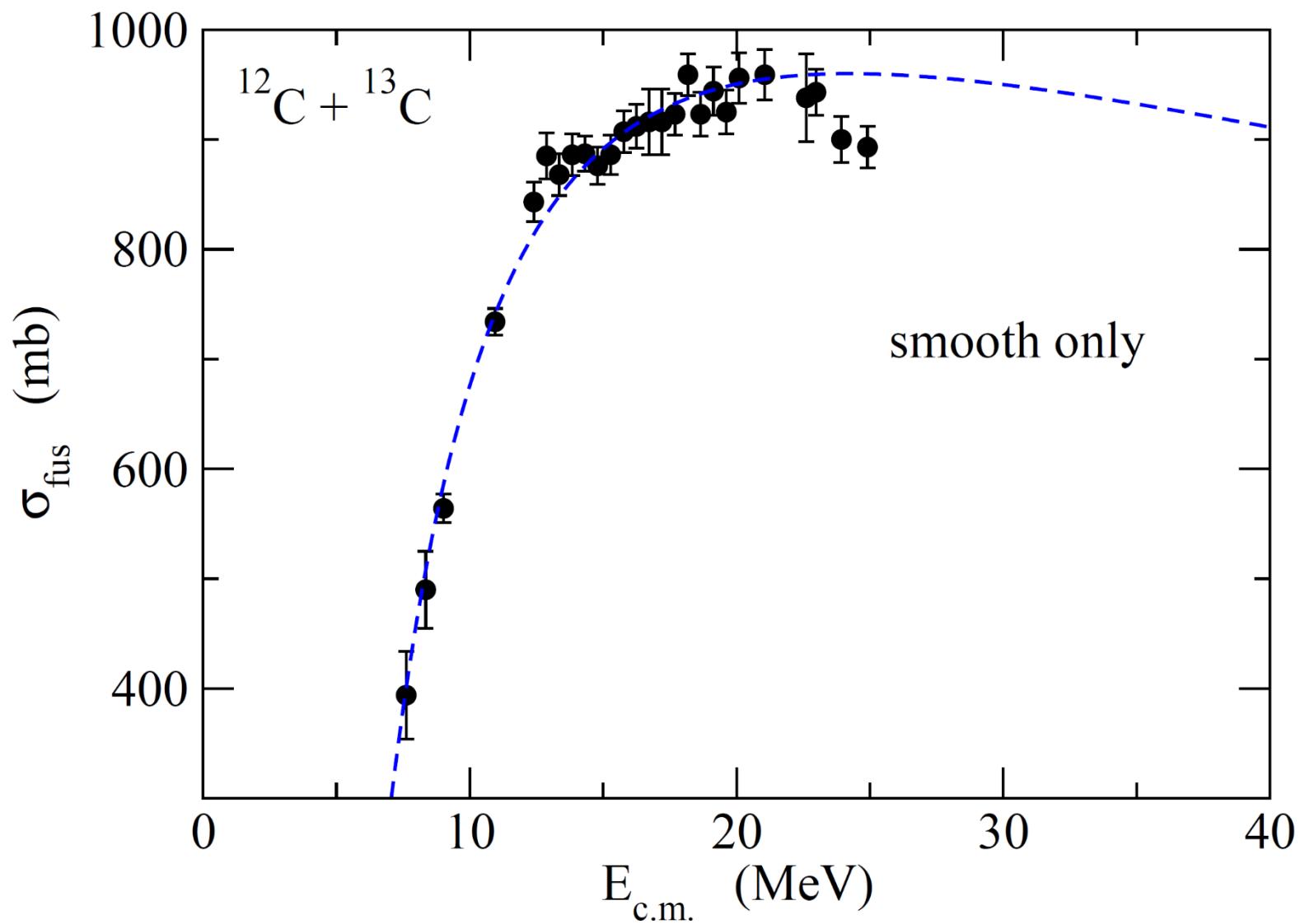
$$\sum_l \rightarrow \frac{1}{4} \sum_l (1 + (-1)^l) + \frac{3}{4} \sum_l (1 - (-1)^l)$$



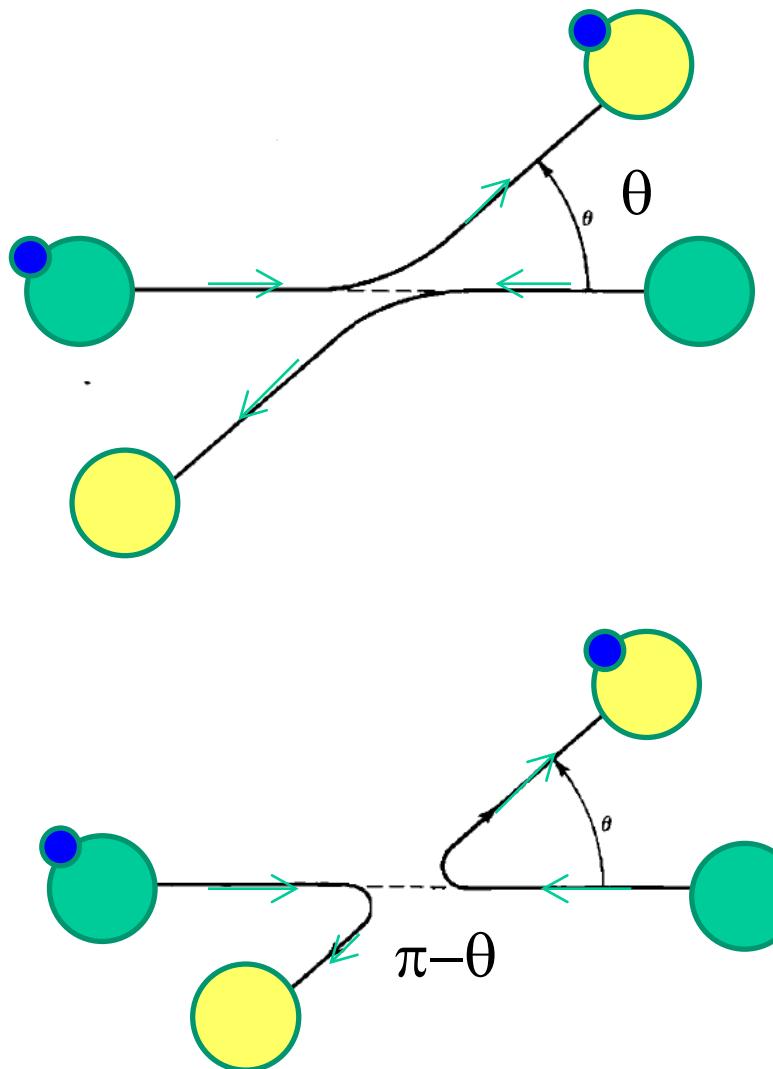
$$\sigma_{\text{osc}} = \frac{1}{2} \sigma_{\text{osc}} (\text{odd} - 1)$$



iii) $^{12}\text{C} + ^{13}\text{C}$



role of elastic transfer



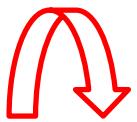
elastic scattering

$$f_{\text{el}}(\theta)$$

indistinguishable

transfer

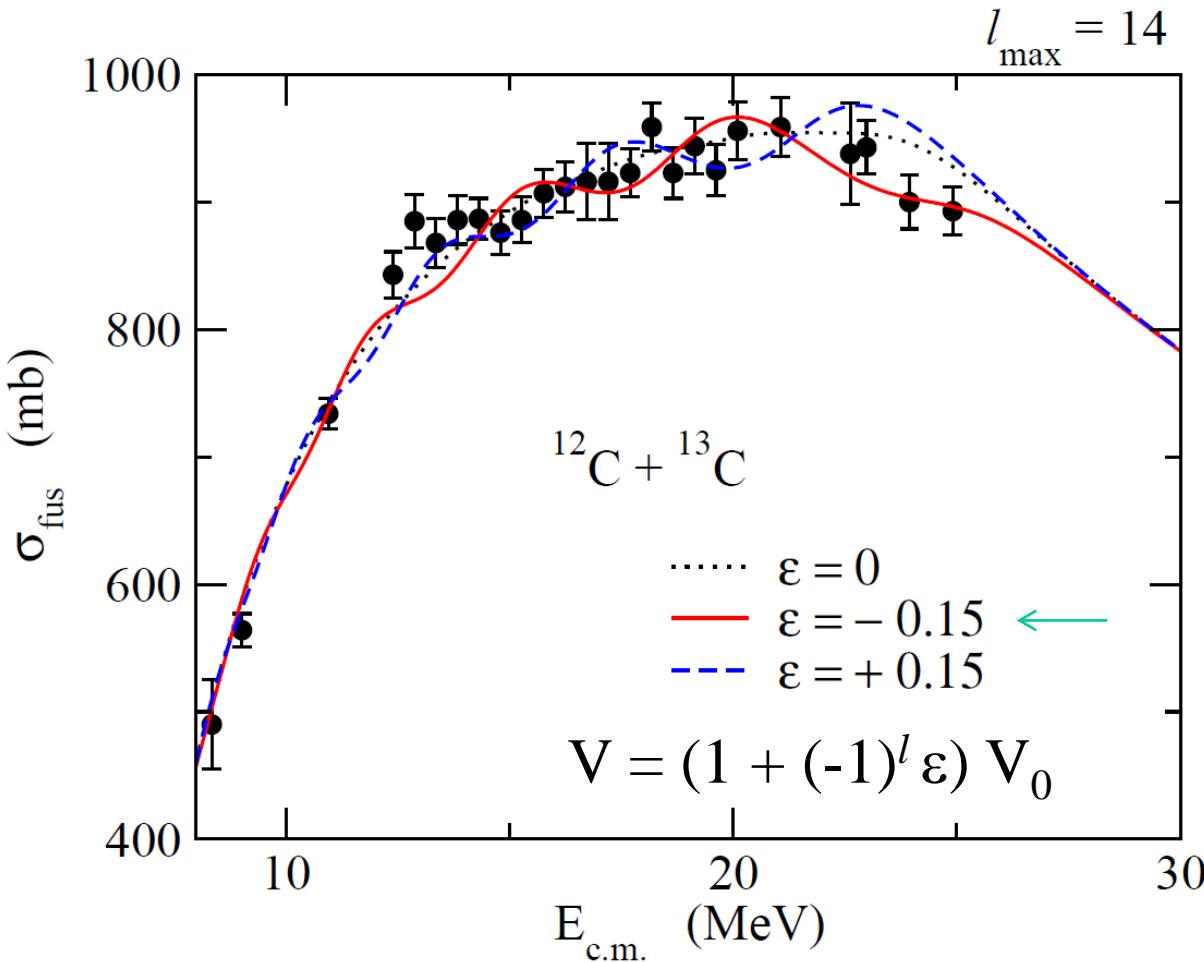
$$f_{\text{trans}}(\pi - \theta)$$



$$f(\theta) \rightarrow f_{\text{el}}(\theta) + f_{\text{trans}}(\pi - \theta)$$

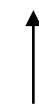
parity-dependent potential

- ✓ W. von Oertzen and H.G. Bohlen, Phys. Rep. 19C('75) 1
- ✓ A. Vitturi and C.H. Dasso, Nucl. Phys. A458 ('86) 157
- ✓ A. Kabir, M.W. Kermode and N. Rowley, Nucl. Phys. A481('88) 94



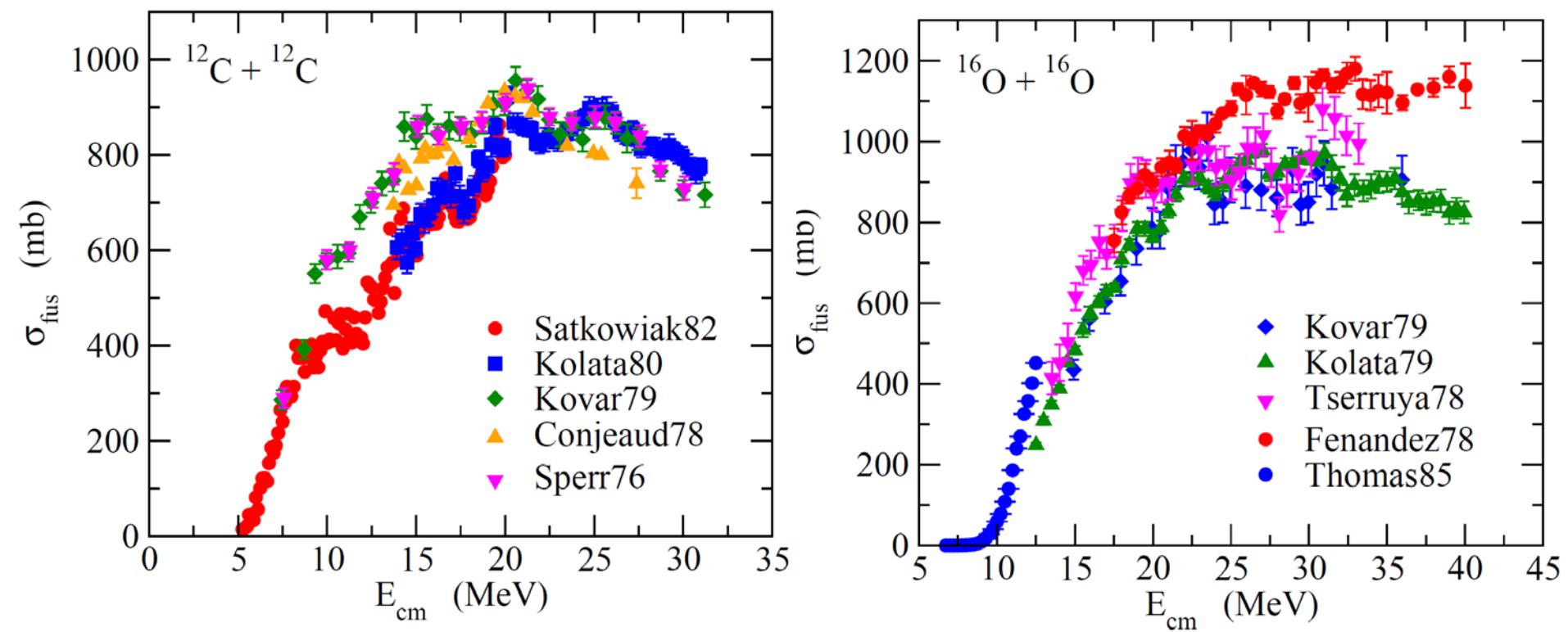
the sign of ϵ :
consistent with Baye's
simple rule

D. Baye, J. Deenen, and
Y. Salmon,
Nucl. Phys. A289('77) 511
D. Baye,
Nucl. Phys. A460 ('86) 581



RGM with two-center
HO shell model

Fusion oscillations with SOD method



The expt. data: rather scattered

- ✓ systematic errors
- ✓ missing evaporation channels

→ σ_{fus} from Sum-of-Differences (SOD) method?

Sum-of-differences (SOD) method

J.T. Holdeman and R.M. Thaler, PRL14('65)81, PR139('65)B1186
C. Marty, Z. Phys. A309('83)261, A322('85)499

$$\sigma_R \sim 2\pi \int_{\theta_{\min}}^{\pi} \sin \theta d\theta (\sigma_{\text{Ruth}}(\theta) - \sigma_{\text{el}}(\theta))$$

expt.: H. Wojciechowski et al., PRC16('77)1767
H. Oeschler et al., NPA325('79)463
T. Yamaya et al., PLB417('98)7 etc.

generalization (K.H. and N. Rowley, in preparation)

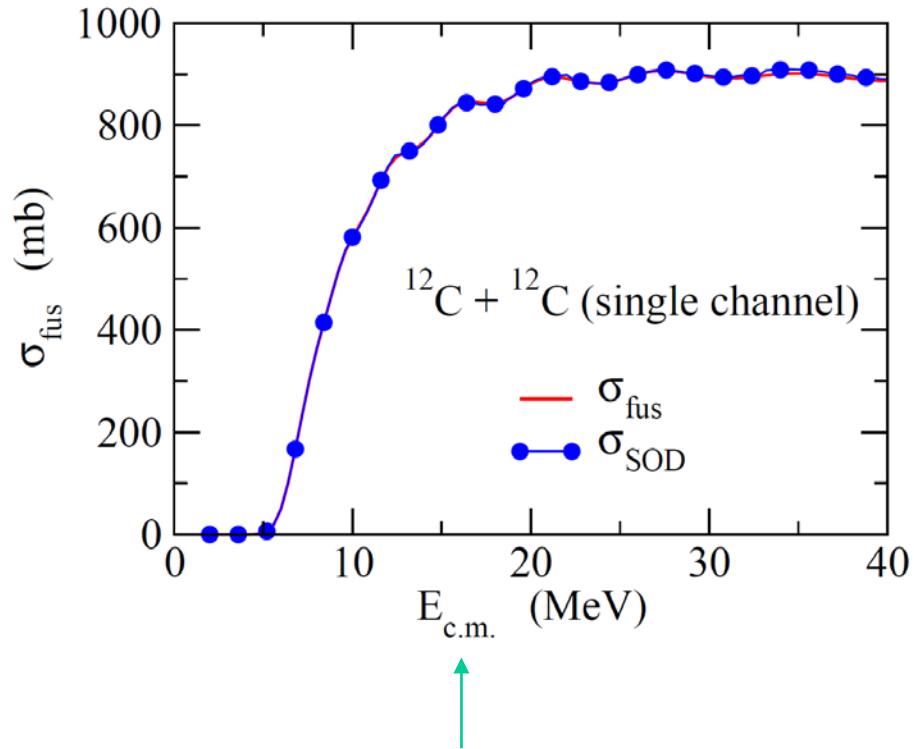
$$\sigma_R = \sigma_{\text{fus}} + \sigma_{\text{inel}} + \sigma_{\text{tr}}$$



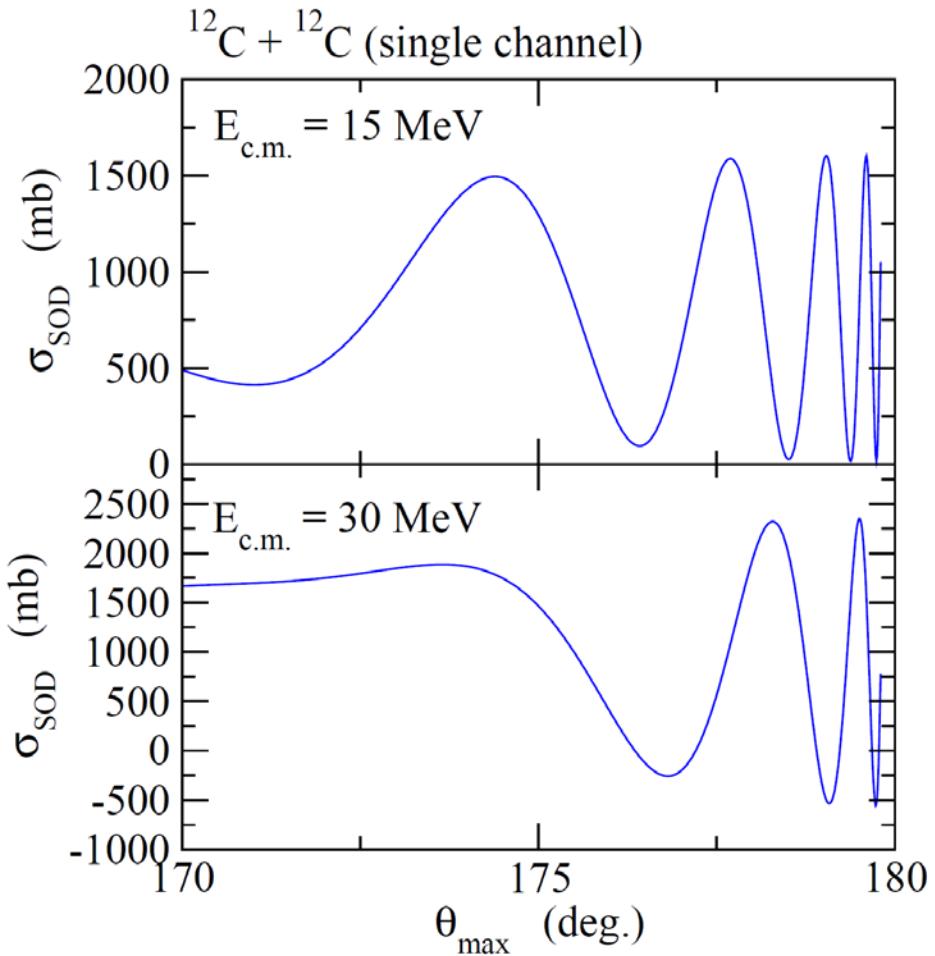
$$\sigma_{\text{fus}} \sim 2\pi \int_{\theta_{\min}}^{\pi} \sin \theta d\theta (\sigma_{\text{Ruth}}(\theta) - \sigma_{\text{qel}}(\theta))$$

for symm. systems: $\sigma_{\text{fus}} \sim 2\pi \int_{\pi/2}^{\theta_{\max}} \sin \theta d\theta (\sigma_{\text{Mott}}(\theta) - \sigma_{\text{qel}}(\theta))$

$$\sigma_{\text{SOD}} = 2\pi \int_{\pi/2}^{\theta_{\max}} \sin \theta d\theta (\sigma_{\text{Mott}}(\theta) - \sigma_{\text{qel}}(\theta))$$



average of a maximum and
a minimum in σ_{SOD}



Summary

sub-barrier fusion of C+C systems

➤ Molecular resonances at subbarrier energies

$^{12}\text{C} + ^{12}\text{C}$: well pronounced resonance structure

$^{13}\text{C} + ^{13}\text{C}$, $^{12}\text{C} + ^{13}\text{C}$: rather smooth

← CN ^{24}Mg : low level density (low Q-value, e-e nucleus)

cf. Jiang's conjecture

➤ Fusion oscillations: successive contribution of discrete centrifugal barriers

$^{12}\text{C}(0^+) + ^{12}\text{C}(0^+)$
 $^{13}\text{C}(1/2^-) + ^{13}\text{C}(1/2^-)$
 $^{12}\text{C} + ^{13}\text{C}$

} symmetrization of relative wave function
--- elastic transfer

➤ Sum-of-differences (SOD) method

- an alternative way to obtain sfus
- application to fusion oscillations?