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# **Spectral Problem in N=4 Super Yang-Mills Theory and Integrable Spin-Chains**

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# N=4 SYM

- ◆ Interesting field theory
- ◆ The main motivation is that it plays an important role in **gauge/string duality** (AdS/CFT correspondence)
- ◆ Maximally supersymmetric field theory in 4D
- ◆ **Conformally invariant** even at quantum level
- ◆ Conformal dimensions are fundamental quantities

$$\langle \mathcal{O}_n(x) \mathcal{O}_m(y) \rangle = \frac{c_n \delta_{nm}}{|x - y|^{2\Delta_n}} \leftarrow \text{conformal dim.}$$

- ◆ **Integrable structure** appears in computing conformal dimensions (today's talk)

- ◆ **Action** ( $\mu, \nu = 1, \dots, 4; i, j = 1, \dots, 6$ )

$$S = \frac{1}{2g_{\text{YM}}^2} \int d^4x \text{Tr} \left[ -\frac{1}{2} (F_{\mu\nu})^2 + (D_\mu \Phi_i)^2 + \frac{1}{2} [\Phi_i, \Phi_j]^2 \right. \\ \left. + i\bar{\psi} \Gamma^\mu D_\mu \psi + \bar{\psi} \Gamma^i [\Phi_i, \psi] \right]$$

- ◆ All appearing fields belong to the **adjoint representation** of the gauge group  $SU(N)$
- ◆ The large  $N$  limit is taken:  $N \rightarrow \infty$  with  $\lambda \equiv g_{\text{YM}}^2 N$  fixed  
 ➔ Only **planar diagrams** contribute ↖ 't Hooft coupling
- ◆ Conformal dims. are **eigenvalues of the dilatation operator**
- ◆ The dilatation operator is expanded in  $\lambda$  perturbatively
 
$$D = D_0 + g^2 D_1 + g^4 D_2 + g^6 D_3 + \dots, \quad (g \equiv \sqrt{\lambda}/4\pi)$$
- ◆ Generic gauge inv. op. is not an eigenstate of  $D$   
**(operator mixing problem)**
- ◆ We have to **diagonalize** it to know conformal dimensions

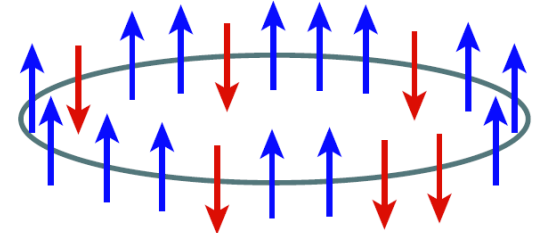
# Integrable Spin-Chain

- ◆ Minahan and Zarembo found that the one-loop dilatation operator is identified with the **Hamiltonian of certain integrable spin-chain** Minahan, Zarembo '02

$$D_1 = 2 \sum_{k=1}^L (I_{k,k+1} - P_{k,k+1}) \leftarrow \text{Hamiltonian of XXX}_{1/2} \text{ spin-chain}$$

- ◆ **Operator-state mapping**  $Z = \Phi_1 + i\Phi_2$   
 $W = \Phi_3 + i\Phi_4$

$$\text{Tr}(ZZZZWZZWWZZZWZZZWZZW)$$



- ◆ This Hamiltonian is diagonalized by the **Bethe ansatz method**

$$e^{ip_j L} = \prod_{k(\neq j)}^M \frac{u(p_j) - u(p_k) - i}{u(p_j) - u(p_k) + i}$$

**two-body S-matrix**

$$\Delta = L + \sum_{j=1}^M 8g^2 \sin^2 \left( \frac{p_j}{2} \right)$$

**dispersion relation**

$$u(p) \equiv \frac{1}{2} \cot \left( \frac{p}{2} \right)$$

- ◆ Dilatation operators at **higher loops** are also mapped to integrable spin-chains with **long-range interactions**

$$D_2 = -8\{\} + 12\{1\} - 2(\{1, 2\} + \{2, 1\}),$$

$$D_3 = 60\{\} - 104\{1\} + 4\{1, 3\} + 24(\{1, 2\} + \{2, 1\}) - 4i\epsilon_2(\{1, 3, 2\} - \{2, 1, 3\}) - 4(\{1, 2, 3\} + \{3, 2, 1\}), \quad \{a, b, c, \dots\} \equiv \sum_{r=0}^{L-1} P_{a+r, a+r+1} P_{b+r, b+r+1} P_{c+r, c+r+1} \dots$$

- ◆ Higher-loop operators have very complicated form, so it is very hard (but possible in principle) to compute them
- ◆ If we **assume** that these operators are also diagonalized by Bethe ansatz method, the problem is to determine the **exact S-matrix** and **dispersion relation**

$$e^{ip_j L} = \prod_{k(\neq j)}^M S(p_j, p_k; \lambda) \quad \Delta = L - M + \sum_{j=1}^M \epsilon(p_j; \lambda)$$

including higher-loop corrections

- ◆ This problem has been solved

Beisert '05  
Beisert, Eden, Staudacher '06

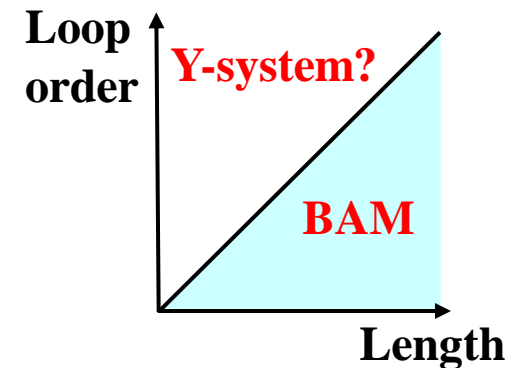
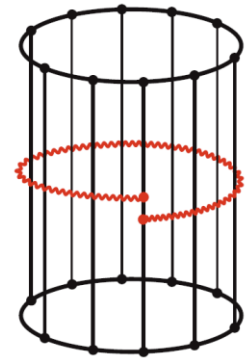
# Finite-Size Problem (Wrapping Problem)

- ◆ Recall that higher-loop dilatation ops. contain **long-range** int.
- ◆ If (loop order)  $>$  (length of spin-chain), interactions ‘wrap’ the spin-chain
- ◆ In this situation, we cannot define asymptotic states and S-matrix

➔ cannot use Bethe ansatz method  
(**wrapping problem**)

- ◆ We need a new approach to resolve this problem
- ◆ Recently a resolution was proposed by using AdS/CFT **Gromov, Kazakov, Vieira ‘09**

➔ **AdS/CFT Y-system**



# Good Example: Konishi Operator

- ◆ Shortest operator with the non-trivial anomalous dim.

$$\mathcal{O}_K \equiv \text{Tr}(Z^2 W^2 - ZWZW)$$

- ◆ Konishi operator has length four, so BA gives correct conformal dimension **up to three loop order**

$$\Delta_K = \underbrace{4 + 12g^2 - 48g^4 + 336g^6}_{\text{BA}} + \underbrace{\Delta^{(4)}g^8 + \Delta^{(5)}g^{10} + \dots}_{\cancel{\text{BA}}}$$

$$\Delta^{(4)} = -2496 + 576\zeta(3) - 1440\zeta(5) \quad \text{via AdS/CFT } \text{Bajnok, Janik '08}$$

Diagrammatic computations  
Fiamberti, Santambrogio, Sieg, Zanon '07  
Velizhanin '08

$$\Delta^{(5)} = 15168 + 6912\zeta(3) - 5184\zeta(3)^2 - 8640\zeta(5) + 30240\zeta(7)$$

via AdS/CFT **Bajnok, Janik, Hegedus, Lukowski '09**  
No field theoretical computations

# Konishi Dimension at Any Coupling

## ◆ One result using AdS/CFT

Gromov, Kazakov, Vieira, arXiv:0906.4240 [hep-th]

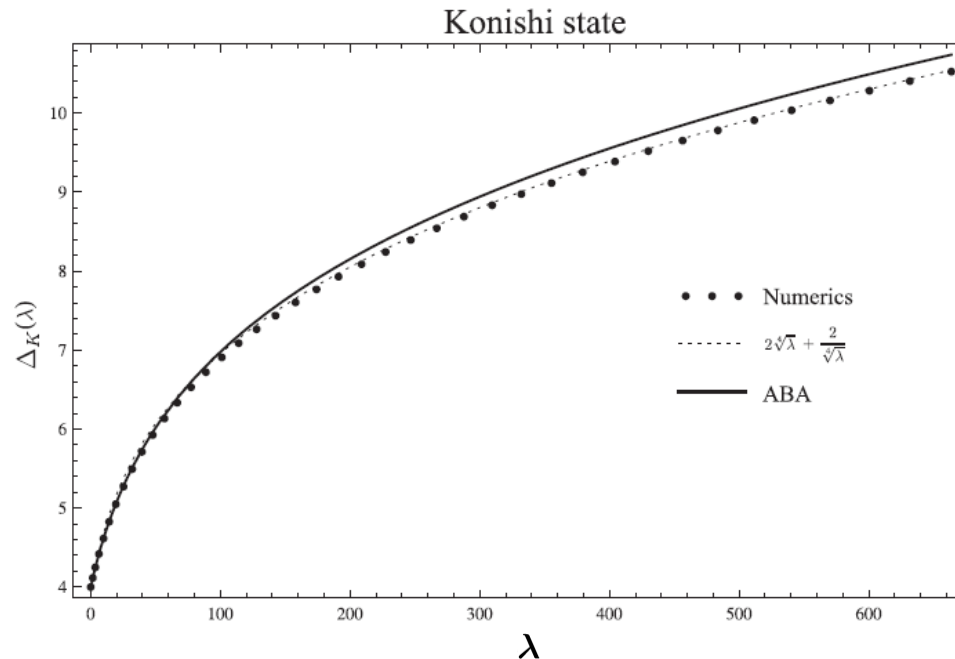


Figure 1: Numerical solution of exact finite size integral Y-system equations for the Konishi dimension  $\Delta_K(\lambda)$  in a wide range of 't Hooft couplings  $\lambda$ , compared to the asymptotic Bethe ansatz curve and to the predicted large  $\lambda$  asymptotics  $\Delta_K(\lambda) \simeq 2\lambda^{1/4} + 2/\lambda^{1/4}$  obtained by fit.



# Summary

- ◆ **The dilatation operator is mapped to that of certain spin-chain Hamiltonian**
  - ◆ **For long operators (or spin-chains), the conformal dims. are given by the Bethe ansatz method**
  - ◆ **The Bethe ansatz method breaks down for finite-size ops.**
  - ◆ **A resolution was proposed by using AdS/CFT**
  - ◆ **Open problems**
    - **Derivation of Y-system purely in field theoretical framework (not using AdS/CFT)**
    - **Exact dimension for the Konishi operator?**
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