

Lunch Seminar 2015/06/10

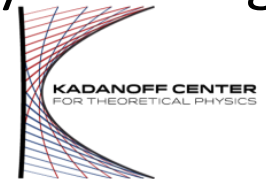
Quantum Entanglement of Local Operators

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1. arXiv:1401.0539 [hep-th]
2. arXiv:1405.5875 [hep-th]
3. arXiv:1405.5946 [hep-th]
4. arXiv.15xx.xxxxx[hep-th]

Introduction

Recently, (Renyi) entanglement entropy ((R)EE) has a center of wide interest in a broad array of theoretical physics.

- It is useful to study the distinctive features of various quantum state in condensed matter physics. (*Quantum Order Parameter*)
- (Renyi) entanglement entropy is expected to be an important quantity which may shed light on the mechanism behind the AdS/CFT correspondence. (*Gravity* \leftrightarrow *Entanglement*)

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- In the lattice gauge theory, it is expected that entanglement entropy is a new order parameter which helps us study QCD more.
- **But** entanglement entropy in the gauge theory is ill-defined.

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 - **But** entanglement entropy in the gauge theory is ill-defined.

It is important to study the properties of (Renyi) entanglement entropy.

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 - **But** entanglement entropy in the gauge theory is ill-defined.

In this work, we investigate the time dependent property of (Renyi) entanglement entropy.

The Definition of (Renyi) Entanglement Entropy

- Definition of Entanglement Entropy

We divide the total Hilbert space into A and B: $H_{tot} = H_A \otimes H_B$.

The reduced density matrix ρ_A is defined by $\rho_A \equiv Tr_B \rho_{tot}$

This means the D O F in B are traced out.

The entanglement entropy is defined by von Neumann entropy S_A .

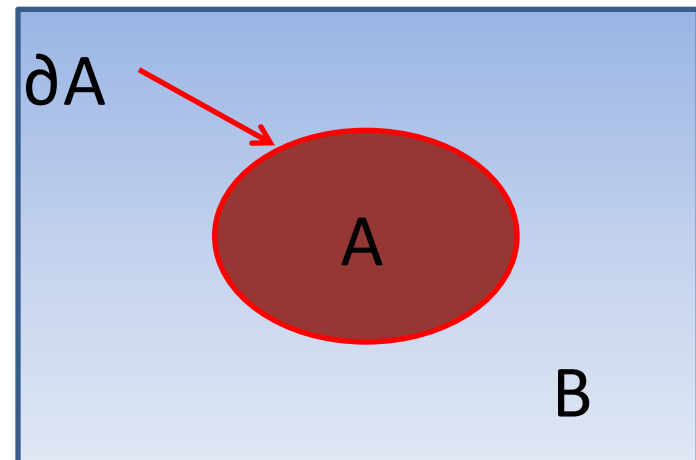
(Renyi) Entanglement Entropy (REE)

$$S_A^{(n)} = \frac{\log \text{tr}[\rho_A^n]}{1 - n}$$

↓ $n \rightarrow 1$

Entanglement Entropy (EE)

$$S_A = -\text{tr}_A \rho_A \log \rho_A$$



on a certain time slice

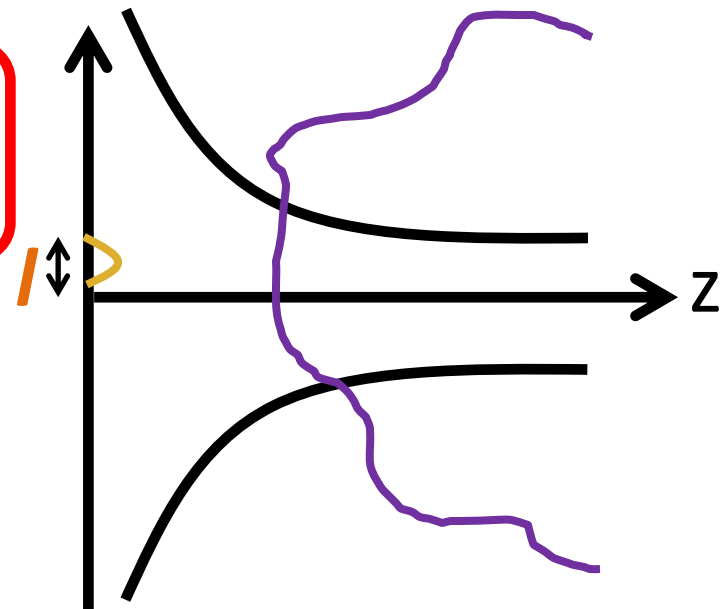
Motivation

Previously, we studied the property of EE for the subsystem whose size (l) is *very small* in d CFT.

$$l \ll (\text{The Excess of Energy Density})^{-d}$$

$$\Delta E_A = T_{ent} \Delta S_A$$

This temperature is universal.



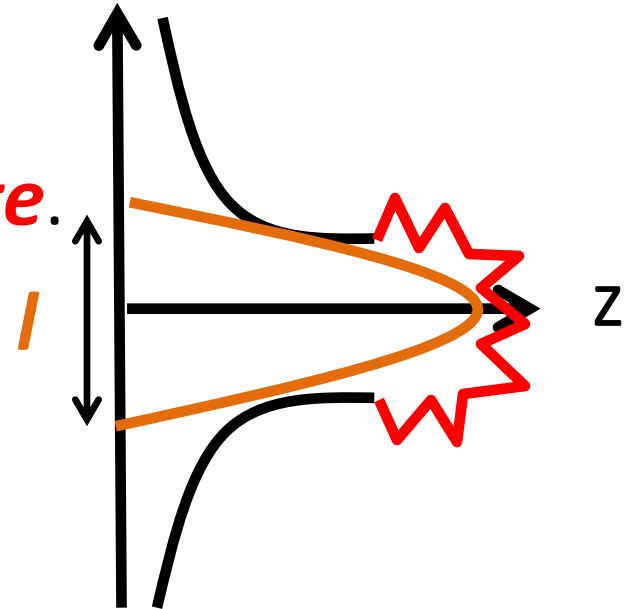
Setup

We study the property of (R)EE for

1. The size of subsystem is *infinite*.

A half of the total system:

$$x^1 \geq 0$$



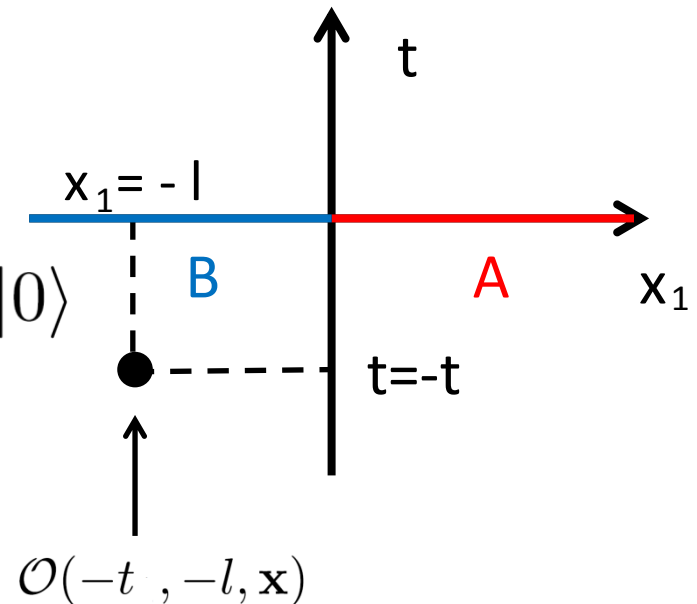
Setup

We study the property of (R)EE for

2. A state is defined by acting a local operator

on the ground state:

$$|\Psi\rangle = \mathcal{N} \mathcal{O}(t = -t, x^1 = -l, \mathbf{x}) |0\rangle$$



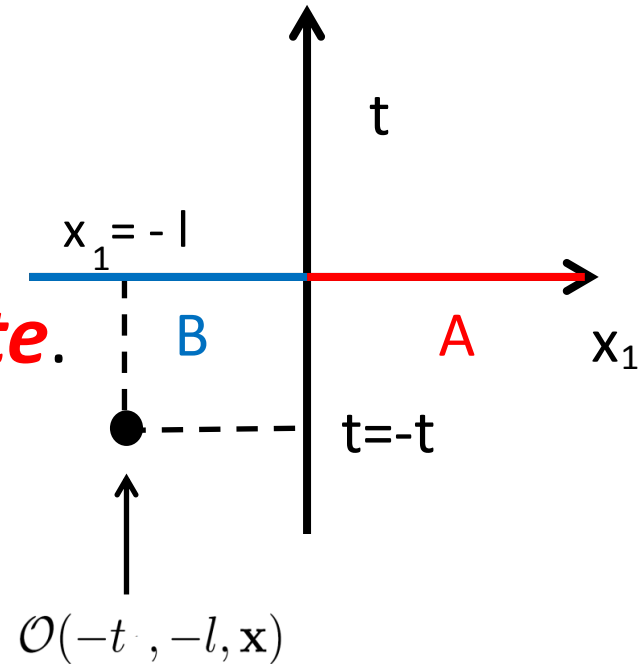
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$$|\Psi\rangle = \mathcal{N} \mathcal{O}(t, x^i) |0\rangle$$

Motivation

We would like to focus on the time evolution of the (R)EE.

We define $\Delta S_A^{(n)}$ the excess of the (R)EE:

$$\Delta S_A^{(n)} = S_A^{(n)Ex} - S_A^{(n)G}$$

$S_A^{(n)Ex}$: (R)EE for $\hat{\rho}_A$ (Reduced Density Matrix for $|\Psi\rangle = \mathcal{N}\mathcal{O}(t, x^i) |0\rangle$)

$S_A^{(n)G}$: (R)EE for the ground state

Field Theory

1. Free massless scalar field theory
2. $U(N)$ or $SU(N)$ free massless scalar field theory
in Large N limit
3. Free massless fermionic field theory
4. Holographic field theory

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Field Theory

1. Free massless scalar field theory


2. $\Delta S_A^{(n)}$  *Some Constants*
At late time

3. Free massless fermionic field theory

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Field Theory

1. Free massless scalar field theory

2. $\Delta S_A^{(n)}$  ***Logarithmically grows***
At late time

3. *[Faded text]*

4. Holographic field theory

Results

$\Delta S_A^{(n)}$ At Late Time ($t \gg \ell$)

F.T.	Operator	$n \geq 2$	$n = 1$
Free Massless Scalar	$:(\partial^m \phi)^k:$	$\frac{1}{1-n} \log \left(\frac{1}{2^{nk}} \sum_{j=0}^k \binom{k}{j} C_j^n \right)$	$k \cdot \log 2 - \frac{1}{2^k} \sum_{j=0}^k \binom{k}{j} C_j \log \binom{k}{j}$
$U(N), SU(N)$ Free Massless Scaler	$Tr(\phi_1 + i\phi_2)^J$	$\frac{2n-1}{n-1} \cdot \log 2$ $J=2$	$\log \left(2\sqrt{2}N \right)$ $J=2$
Free Massless Fermion	$\bar{\psi}\psi$	$\frac{1}{1-n} \log \left[\frac{2 \cdot 12^n + 4 \cdot 9^n + 4}{2^{6n}} \right]$	$\frac{3}{4} \log \left(\frac{128}{9} \right)$
AdS/CFT	Gauge invariant Operator \mathcal{O}_Δ	$\frac{4n\Delta}{d(n-1)} \log t$ $1 \ll \Delta \ll c$	$\frac{c}{6} \log t$ 2d CFT $\Delta \simeq c$

Results

$\Delta S_A^{(n)}$ At Late Time ($t \gg l$)

F.T.	Operator	$n \geq 2$	$n = 1$
Free Massless Scalar	$: (\partial^m \phi)^k :$	<div style="border: 2px solid blue; border-radius: 20px; padding: 20px; display: inline-block;"> <p>Constant</p> </div>	
$U(N), SU(N)$ Free Massless Scaler	$Tr(\phi_1 + i\phi_2)^J$		
Free Massless Fermion	$\bar{\psi}\psi$		
AdS/CFT	Gauge invariant Operator \mathcal{O}_Δ	$\frac{4n\Delta}{d(n-1)} \log t$ $1 \ll \Delta \ll c$	$\frac{c}{6} \log t$ <p>2d CFT $\Delta \simeq c$</p>

Example

We consider *free massless scalar* field theory in *$d+1$ dim.*

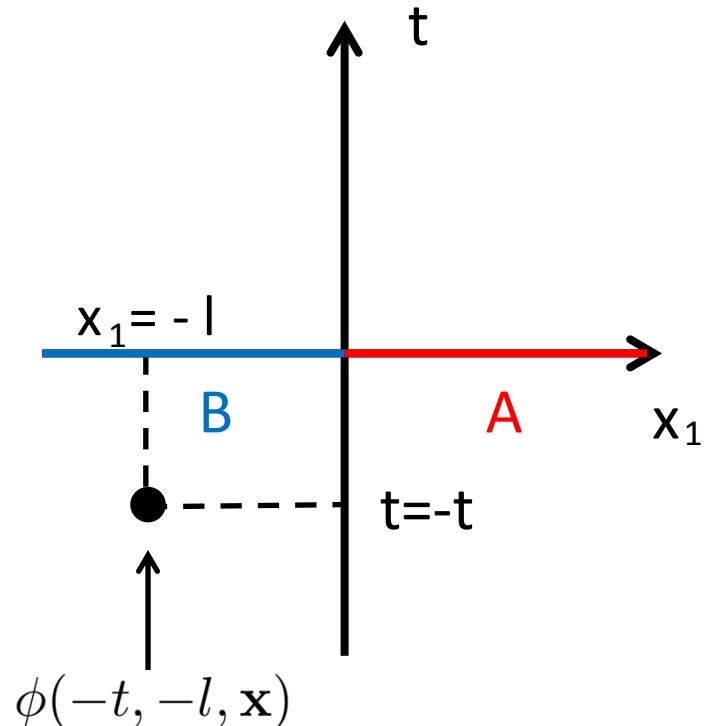
Especially, we focus on that in *4 dim.*

We act a local operator $\phi(-t, -l, \mathbf{x})$ on the ground state:

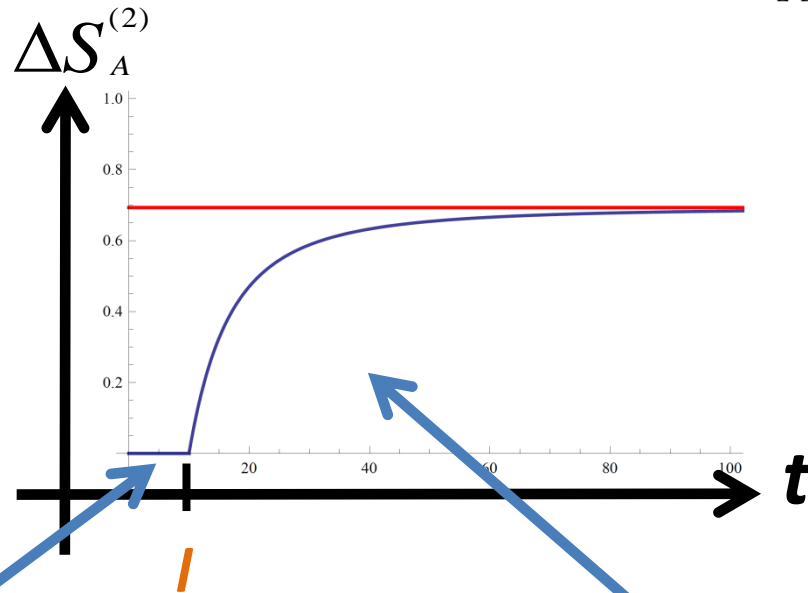
$$|\Psi\rangle = \mathcal{N} \phi(-t, -l, \mathbf{x}) |0\rangle .$$

We measure the (Renyi)
entanglement entropies at $t=0$.

➡ *Time evolution!!*



Time Evolution of $\Delta S_A^{(2)}$

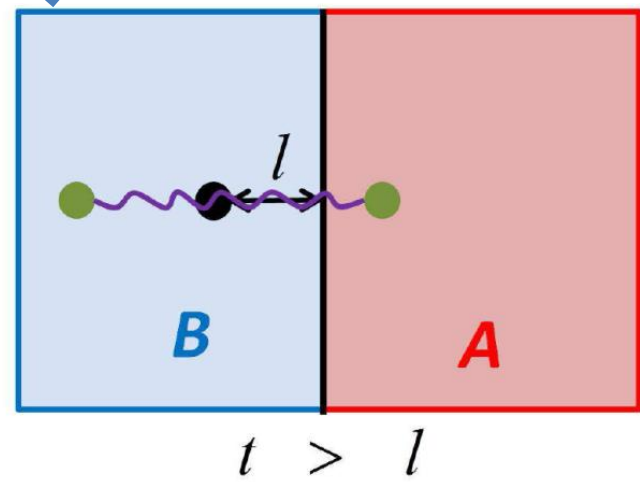
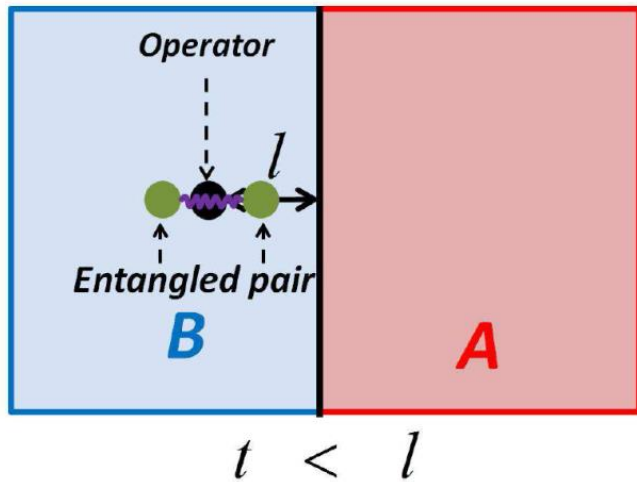


$$t < l$$

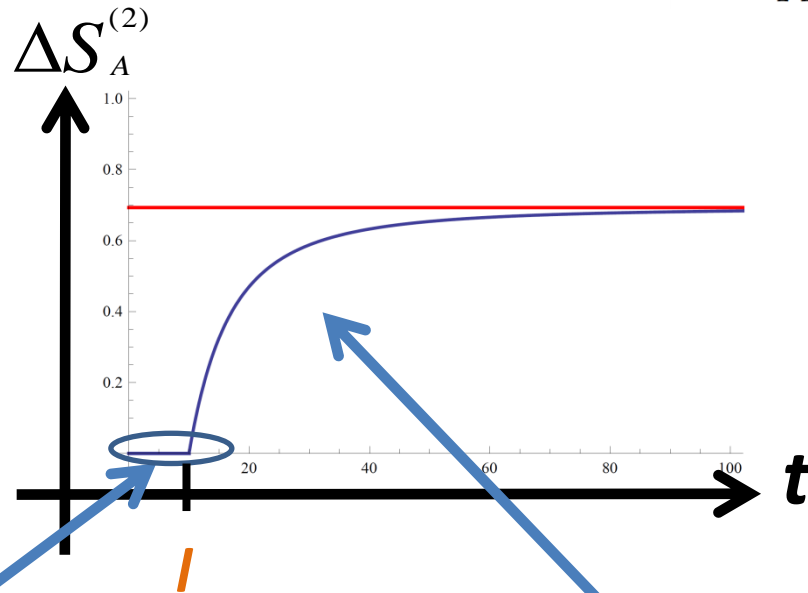
$$\Delta S_A^{(2)} = 0$$

$$t \geq l$$

$$\Delta S_A^{(2)} = \log \left[\frac{2t^2}{t^2 + l^2} \right]$$



Time Evolution of $\Delta S_A^{(2)}$

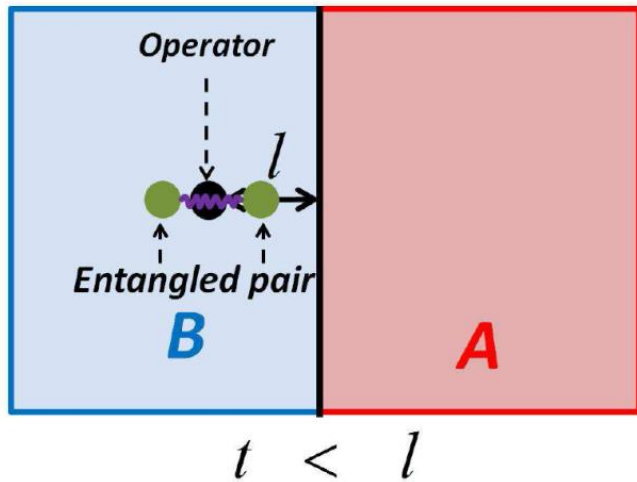


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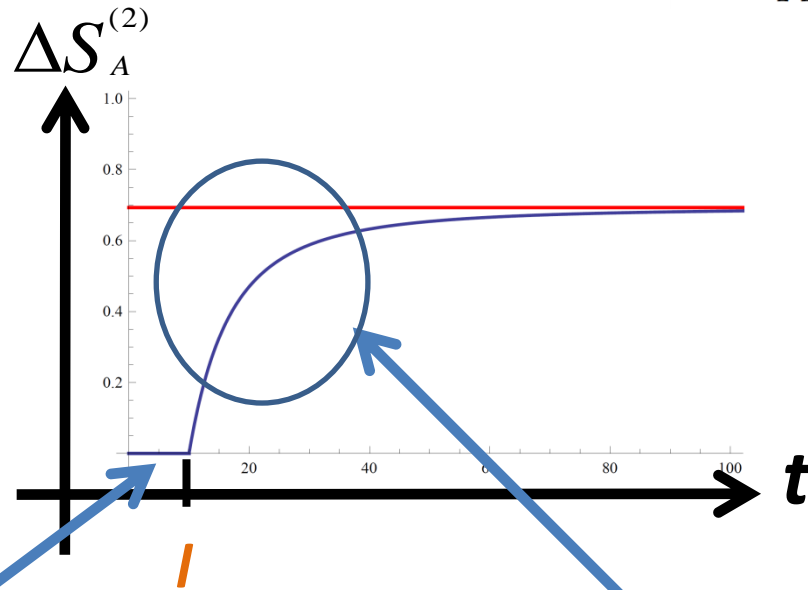
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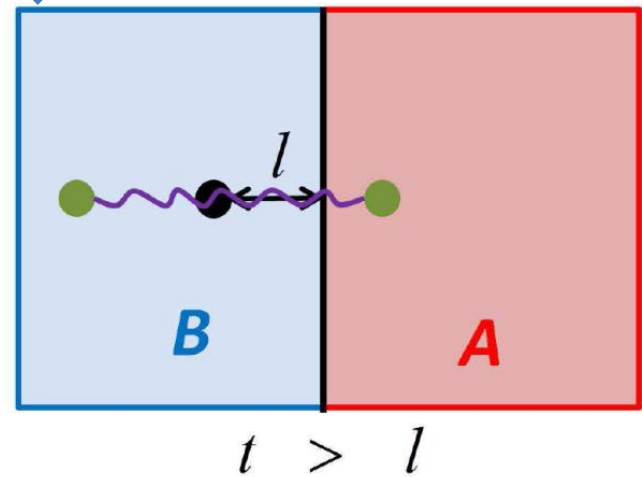
An entangled pair appears.
Each of pair is included in
the region B.

Time Evolution of $\Delta S_A^{(2)}$

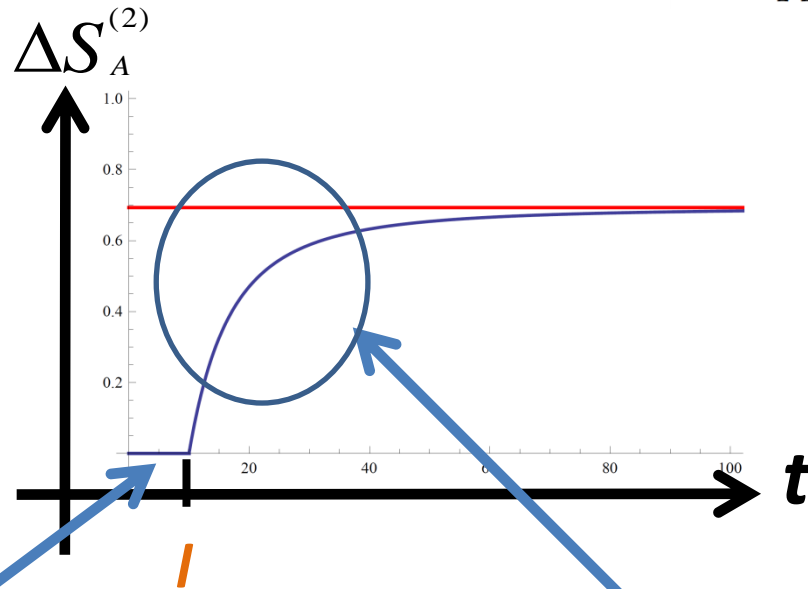


$$t \geq l$$

In this region, two quanta is included in A and B respectively .



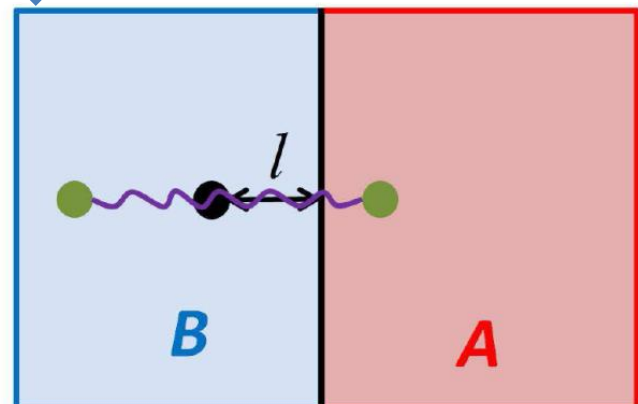
Time Evolution of $\Delta S_A^{(2)}$



$$\Delta S_A^{(2)} = \log \left[\frac{2t^2}{t^2 + l^2} \right]$$

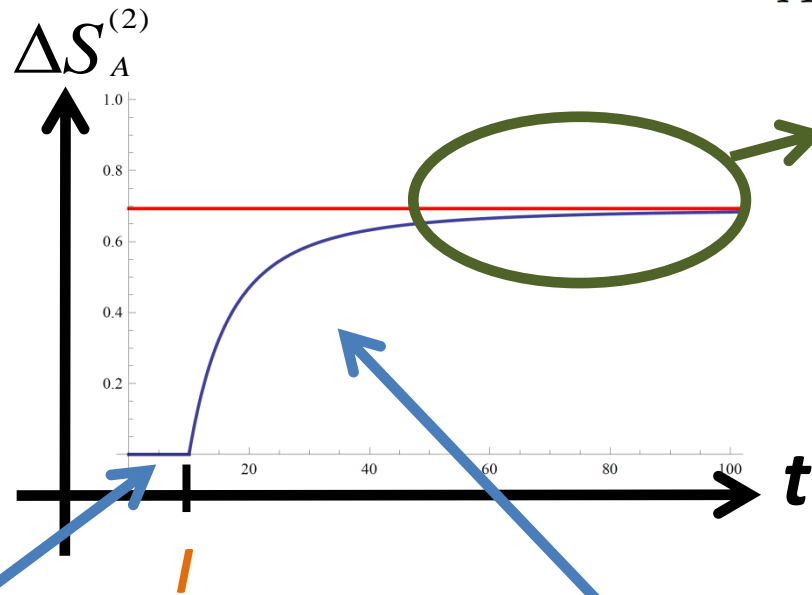
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In this region, two quanta is included in A and B respectively .



Entanglement between quanta can contribute to $\Delta S_A^{(2)}$.

Time Evolution of $\Delta S_A^{(2)}$



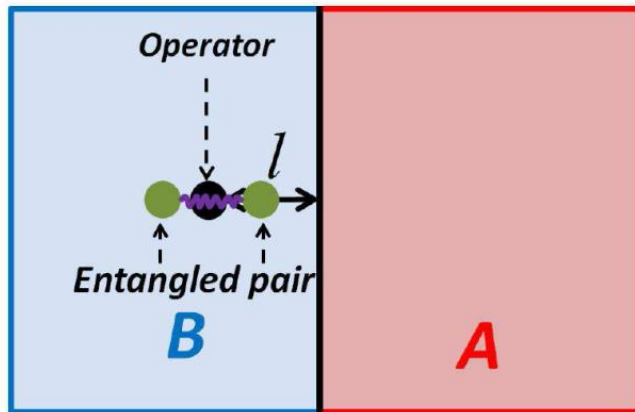
Subsystem

= a half of the
total space

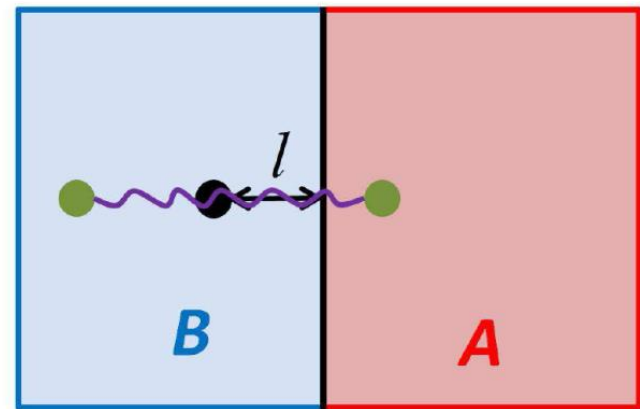


$\Delta S_A^{(2)}$ approaches

Constant!!

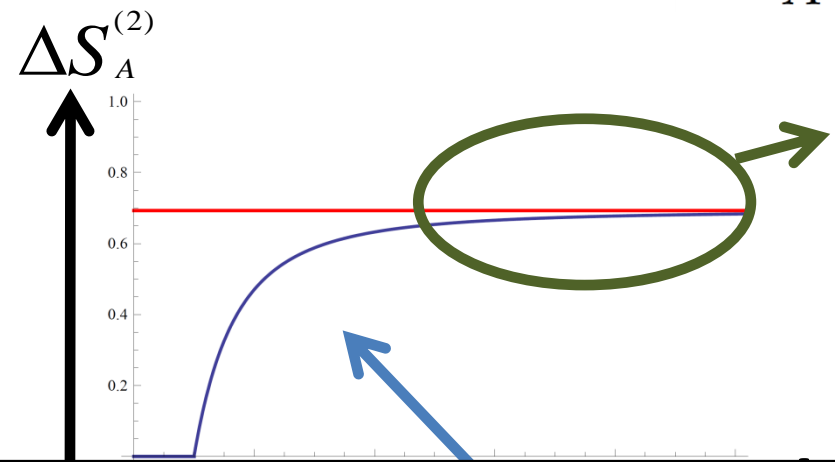


$t < l$



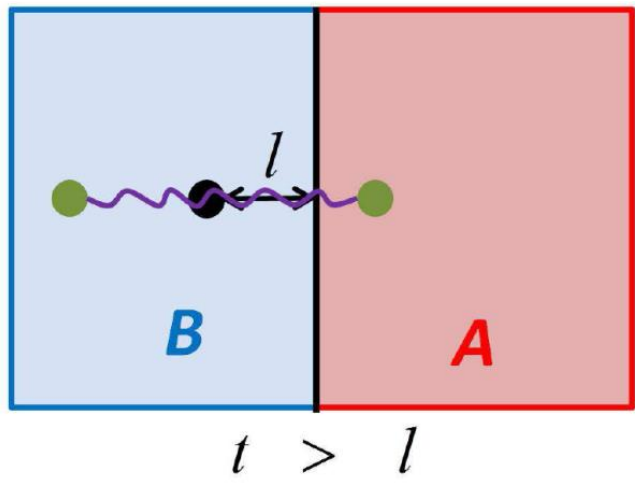
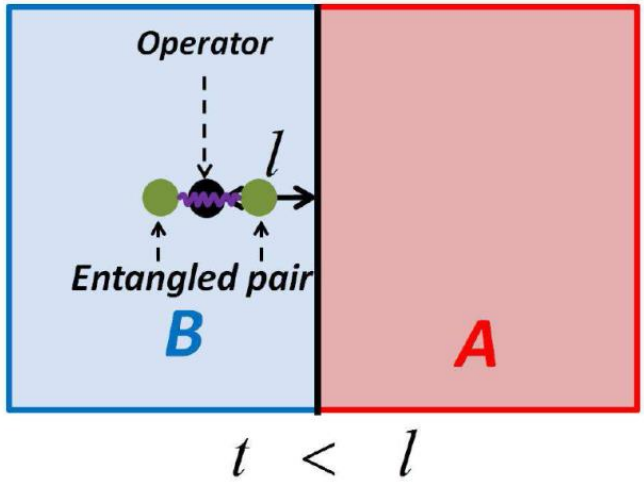
$t > l$

Time Evolution of $\Delta S_A^{(2)}$

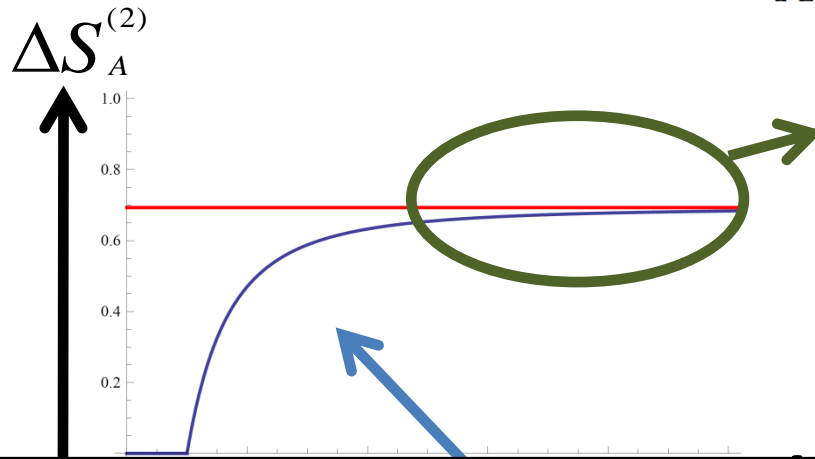


Subsystem
 = a half of the
 total space
 ↓
 $\Delta S_A^{(2)}$ approaches

$\Delta S_A^{(2)}$ = Entanglement between Quasi-particles !!!



Time Evolution of $\Delta S_A^{(2)}$



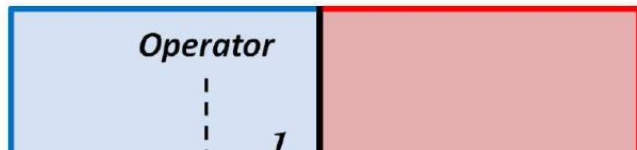
Subsystem

= a half of the
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$\Delta S_A^{(2)}$ approaches

$\Delta S_A^{(2)}$ = Entanglement between Quasi-particles !!!



(Renyi) entanglement entropies of (local) operators

Entangled pair

B

A

$t < l$

B

A

$t > l$

Results

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F.T.	Operator	$n \geq 2$	$n = 1$
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Summary

- If non-perturbative effect for \mathbf{c} or \mathbf{N} , $\Delta S_A^{(n)}$ can possibly approach constant.

• *Free field theory (Infinite conserved charges) \rightarrow Quasi-particle picture holds*

$$\Rightarrow \Delta S_A^{(n)} \rightarrow \text{Constant}$$

Holographic CFT \rightarrow Quasi-particle picture breaks down

$$\Rightarrow \Delta S_A^{(n)} \rightarrow \text{Not constant}$$

Future Problems

- In non-relativistic case, the time evolution of $\Delta S_A^{(n)}$
- In gauge field theory, $\Delta S_A^{(n)}$
- The (Renyi) entanglement entropies of operators in the interacting field theory . (also massive and charged Renyi.)
- Beyond large N, we investigate $\Delta S_A^{(n)}$.
 - approach constant?
 - diverge?