Nonsupersymmetric Brane/Antibrane Configurations in Type IIA/M-theory

Masaki Shigemori (Caltech)

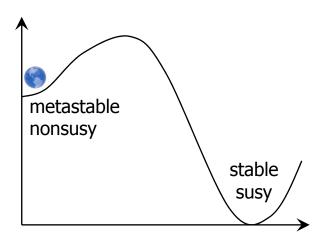
May 17, 2007 at KITP

This talk is based on...

- arXiv:0705.0983
   Joe Marsano, Kyriakos Papadodimas
   & MS (84 pages!)
- Joe gave a preview at KITP in March.
   This is continuation/supplement

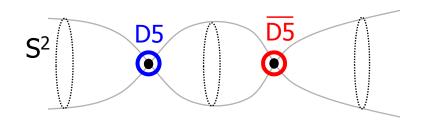
### Introduction

- Metastable nonsusy configs.
   in gauge theory / string theory
  - ISS
  - Landscape
  - Pheno./cosmo.



### Introduction

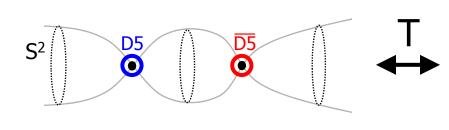
- Geometrically induced metastability
  - Aganagic-Beem-Seo-Vafa

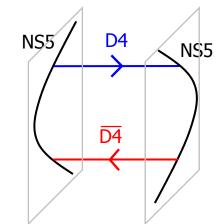


- Large N duality
- Fluxes softly break  $N=2 \rightarrow N=0$
- Calculational control

### Introduction

### T-dual: IIA/M





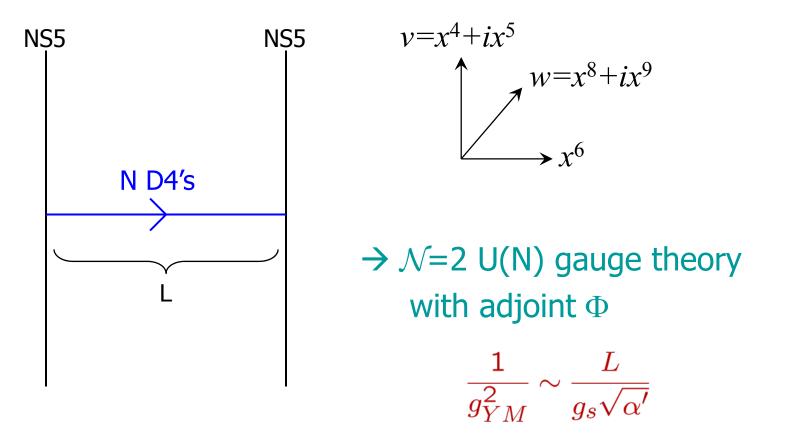
- IIA side: M-theory lift
- Unexplored param. regime in string/M-theory
- "Softness" of breaking
- Boundary conditions [Bena-Gorbatov-Hellerman-Seiberg-Shih]

### Plan of this talk

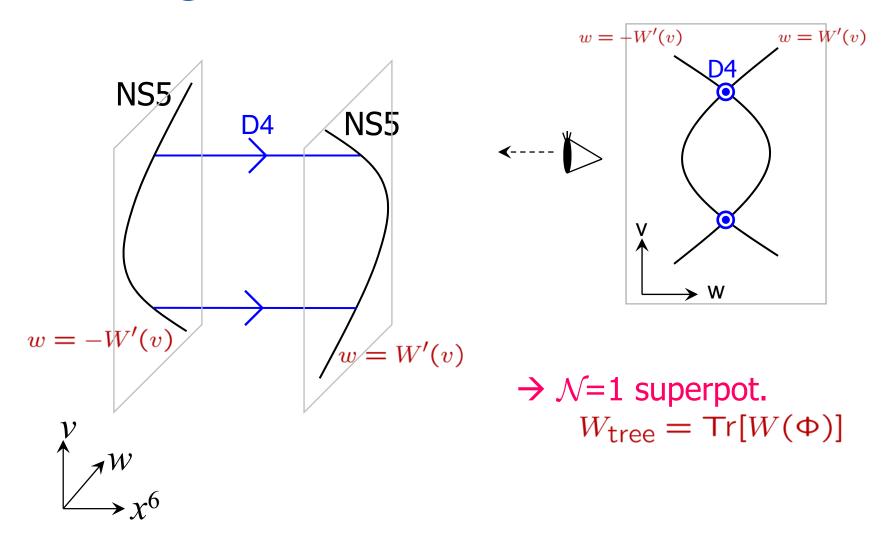
- Introduction
- Susy case
- Nonsusy curves
- "Soft" limit
- Boundary conditions

Review: Susy Case

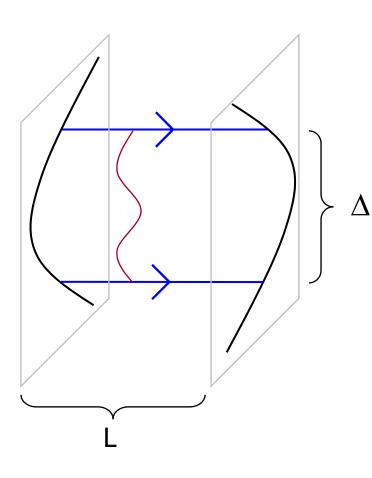
# Type IIA brane construction (a.k.a. Hanany-Witten)



Curving NS5



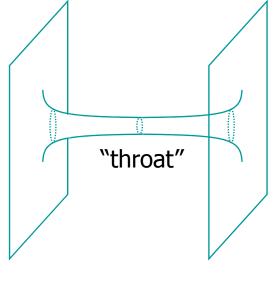
### Gauge theory limit



 $\alpha' \rightarrow 0$  $rac{1}{g_{YM}^2}\sim rac{L}{g_s\sqrt{lpha'}}$  : fixed  $M_{\rm F1} \sim \frac{\Delta}{\alpha'}$ : fixed  $L \sim g_s \sqrt{lpha'}$  $\Delta \sim lpha'$ 

#### System size is substringy

# Gauge dynamics: M-theory lift [Witten]



NS5/D4  $\rightarrow$  M5

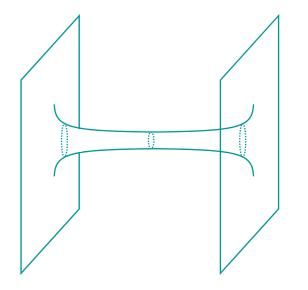
### NS5 & D4 both lift to M5

• Large  $g_s$  $\rightarrow$  NG action reliable

#### M5 curve:

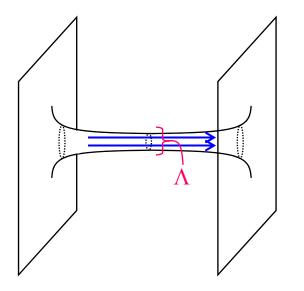
- (Relatively) easily obtained by holomorphicity
- SW curve itself!

# Strong coupling: M-theory lift



- Why does it work?
   Power of susy
  - Scales irrelevant for holomorphic quantities
  - NG & gauge theory give same curve

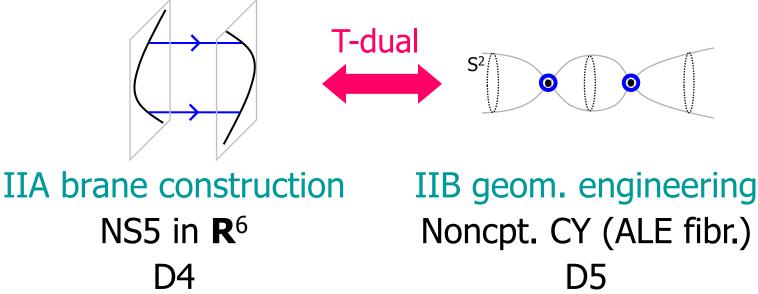
# When is classical curve *really* reliable?



### • As M5 curve ( $g_s \gg 1$ ):

- Curvature  $\ll l_{11}$
- Don't come within l<sub>11</sub>
   of self-intersecting
- As NS5 + flux ( $g_s \ll 1$ ):
  - Curvature  $\ll l_s$
  - Flux density  $\ll 1/g_s$
- Never met in gauge theory limit

### **T**-duality



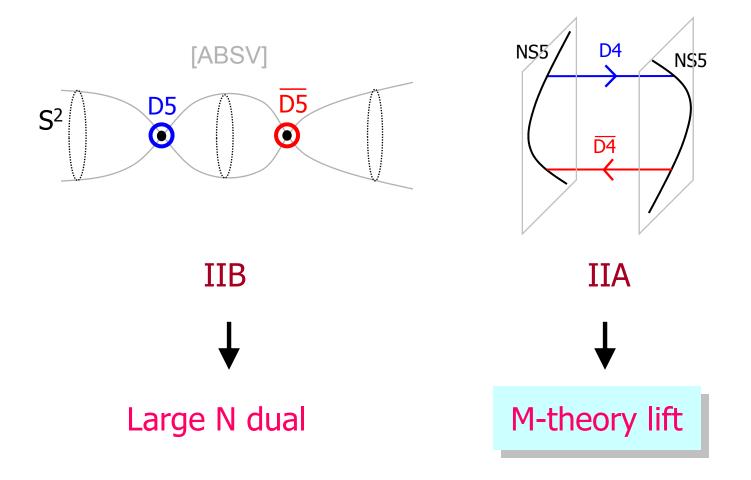
Tool: M-theory lift

Tool: large N duality

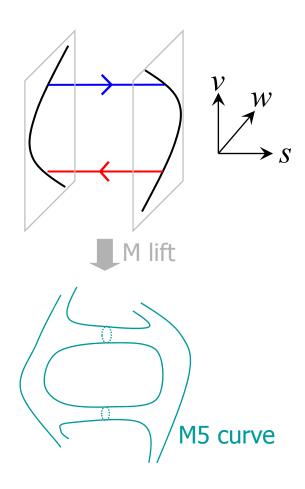
- Why work? power of susy
  - Scales irrelevant
  - Without susy, not expected to work, a priori

Nonsusy Curves

## Nonsusy configurations



# EOM's for nonsusy M5 curves



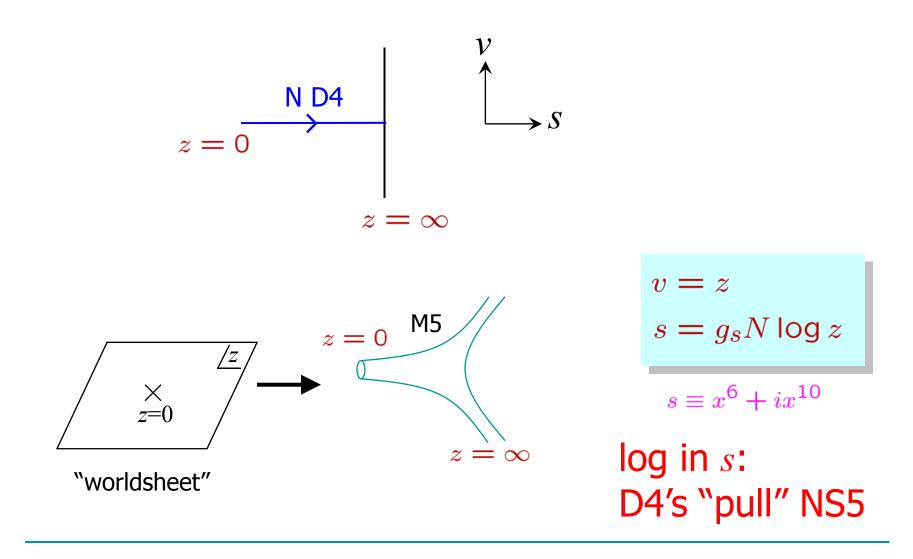
 Need to directly solve NG/Polyakov:

$$S_{NG}\sim \int \sqrt{g}$$

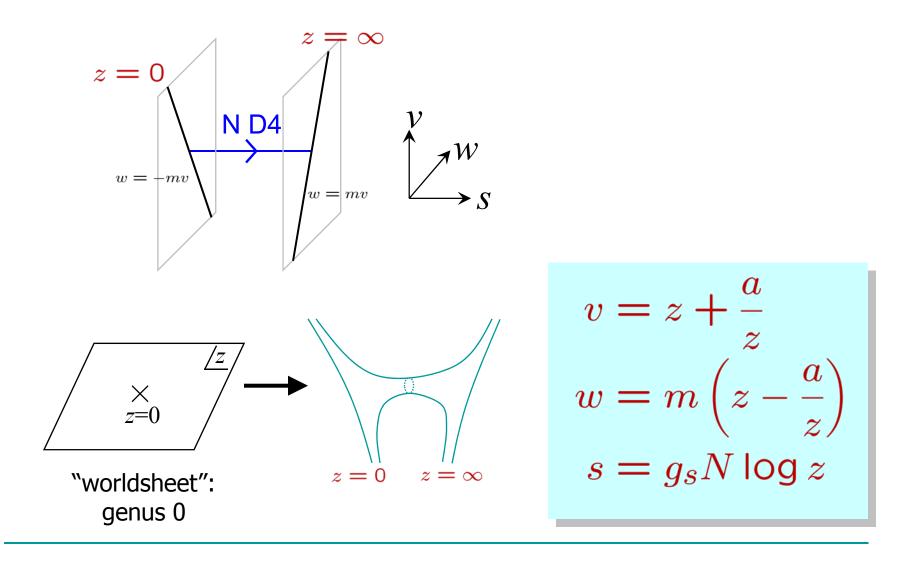
Harmonic:

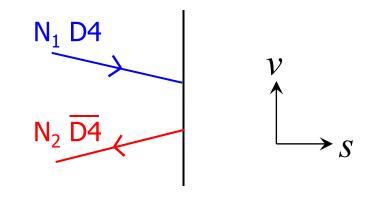
$$s = s_H(z) + \overline{s}_A(\overline{z})$$
$$v = v_H(z) + \overline{v}_A(\overline{z})$$
$$w = w_H(z) + \overline{w}_A(\overline{z})$$

Virasoro constraint:  $\partial s_H \partial s_A + \partial v_H \partial v_A + \partial w_H \partial w_A = 0$ 



### Practice 2: linearly curved NS5's

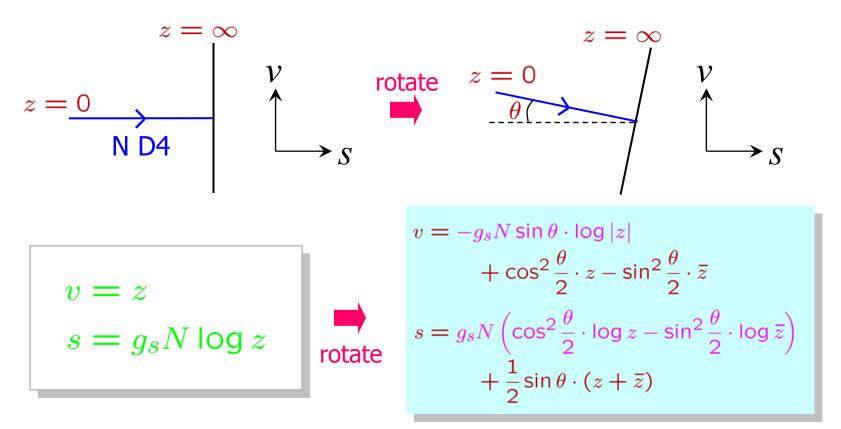




#### New features:

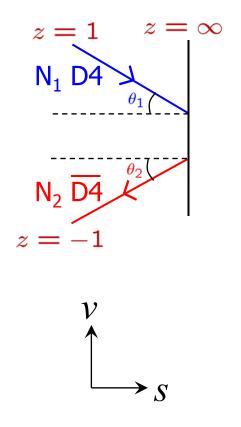
- D4's and D4's attract
- Need to "hold" back D4's
- D4's "tilt" in v direction

#### How to get tilted D4



log appears in tilted direction

## Now combine D4 and D4 parts

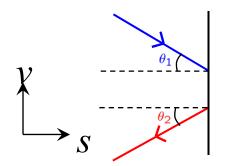


$$v = -g_s N_1 \sin \theta_1 \cdot \log |z - 1|$$
  
+  $g_s N_2 \sin \theta_2 \cdot \log |z + 1|$   
+  $\Delta v_H + \overline{\Delta} v_A$   
$$s = +g_s N_1 \left[ \cos^2 \frac{\theta_1}{2} \cdot \log(z - 1) - \sin^2 \frac{\theta_1}{2} \cdot \log(\overline{z} - 1) \right]$$
  
+  $g_s N_2 \left[ \cos^2 \frac{\theta_2}{2} \cdot \log(\overline{z} + 1) - \sin^2 \frac{\theta_2}{2} \cdot \log(z - 1) \right]$   
+  $\Delta s_H + \overline{\Delta} s_A$ 

 $\Delta v_{H,A}, \ \Delta s_{H,A}$  : determined by Virasoro

Virasoro

 $\partial s_{\mathbf{H}} \partial s_A + \partial v_{\mathbf{H}} \partial v_A + \partial w_{\mathbf{H}} \partial w_A = 0$ 

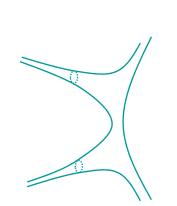


• Force balance  $N_1 \sin \theta_1 = N_2 \sin \theta_2$ 

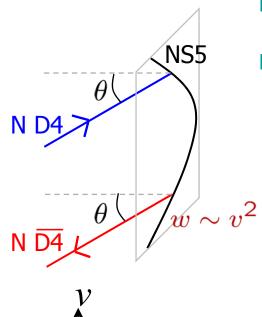
#### Full nonholo. curve obtained!

$$v = -g_s N_1 \sin \theta_1 \cdot \log |z - 1| + g_s N_2 \sin \theta_2 \cdot \log |z + 1| + az$$

$$s = +g_s N_1 \left[ \cos^2 \frac{\theta_1}{2} \cdot \log(z-1) - \sin^2 \frac{\theta_1}{2} \cdot \log(\bar{z}-1) \right] \\ +g_s N_2 \left[ \cos^2 \frac{\theta_2}{2} \cdot \log(\bar{z}+1) - \sin^2 \frac{\theta_2}{2} \cdot \log(z-1) \right] \\ a = \frac{2N_1 N_2 \cos^2 \left( \frac{\theta_1 + \theta_2}{2} \right)}{N_1 \sin \theta_1 + N_2 \sin \theta_2}$$



### Practice 4: quadratic curving ("half" of real thing)



Same method worksFull nonholo. curve:

$$v = Xz + \frac{(g_s N)^2}{\overline{X}}\overline{z}$$

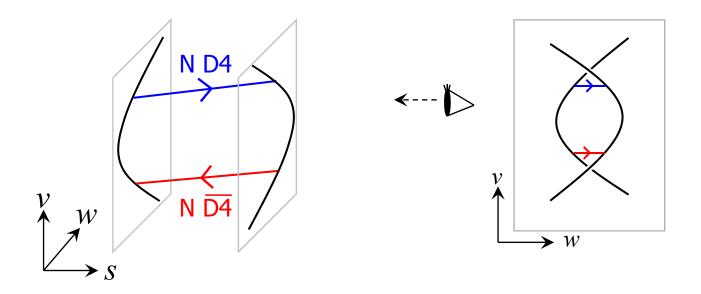
$$w = g_s N \sin \theta \cdot \log |z - 1| + g_s N \sin \theta \cdot \log |z + 1| - \frac{g_s N}{2 \sin \theta} z^2$$

$$s = +g_s N \Big[ \cos^2 \frac{\theta}{2} \cdot \log(z - 1) - \sin^2 \frac{\theta}{2} \cdot \log(\overline{z} - 1) \Big]$$

$$+g_s N \Big[ \cos^2 \frac{\theta}{2} \cdot \log(\overline{z} + 1) - \sin^2 \frac{\theta}{2} \cdot \log(z - 1) \Big]$$

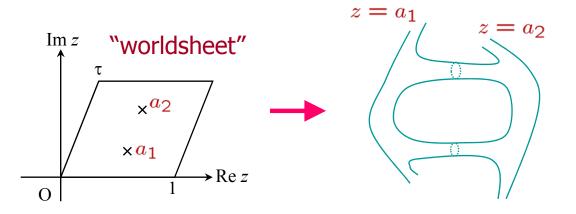
- ✤ Log bending in s, w
- B.C. at ∞ not holomorphic
- Minimal distance b/t D4 & D4

# The real problem



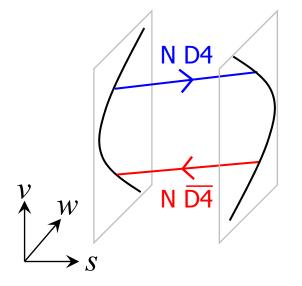
Worldsheet": genus 1
D4's tilt in w direction
→ w have "logs"

### The exact curve



$$\begin{cases} v = X \left[ F_1^{(1)} - F_2^{(1)} \right] + \frac{2\xi(g_s N)^2}{\overline{X}} \left[ F_1^{(1)} - F_2^{(1)} \right] \\ w = g_s N r_0 \sin \theta \left[ (F_2 - F_1 - i\pi)z + \text{cc} \right] + \frac{g_s N \xi}{r_0 \sin \theta} (F_1^{(2)} - F_2^{(2)}) \\ s = g_s N r_0 \cos \theta \left[ (F_1 - F_2 + i\pi z) + \text{cc} \right] + i\pi g_s N(z + \overline{z}) \end{cases}$$
$$F_i(z) = \log \left[ \theta \left( z - a_i - \frac{\tau + 1}{2} \right) \right], \qquad F_i^{(n)}(z) = \left( \frac{\partial}{\partial z} \right)^n F_i(z) \\ r_0^2 = \frac{3\pi^2 \wp(\tau/2)}{3\wp(\tau/2)^2 - g_2}, \qquad \xi \equiv \frac{\pi^2}{6\wp(\tau/2)^2 - 2g_2} \\ \text{Re}(\tau) = 0, \qquad \tau = 2(a_2 - a_1) \end{cases}$$

### The exact curve



#### Same features as the "half" curve:

- Log bending in s, w
- B.C. at ∞ not holomorphic
- Minimal distance
   between D4 & D4
- T-dual won't even be a complex manifold

Reliable as M5/NS5 curve

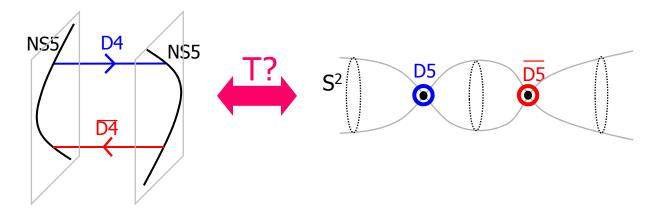
# The "Soft Limit"

What if  $g_s N$  is small?

Take 
$$\frac{g_s N}{\Delta} \to 0$$
 in the exact curve:  
 $v = X \left[ F_1^{(1)} - F_2^{(1)} \right]$   
 $w = C \left[ (F_2 - F_1 - i\pi)z + cc \right]$   
 $s = g_s N r_0 \cos \theta \left[ (F_1 - F_2 + i\pi z) + cc \right] + i\pi g_s N(z + \overline{z})$ 

- Tilting vanishes
- v, w : holomorphic, s : harmonic
- Corresponds to exact min. of IIB potential
   : unexpected!

### "T-duality"



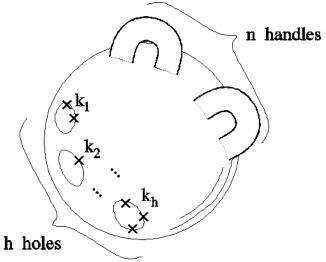
- Holds more generally
- Why work without susy?
  - There must be some protection
  - IIB [ABSV]: assuming soft breaking
  - → g<sub>s</sub>N is controlling "softness"

### The "soft limit"

Small g<sub>s</sub>N : consistent with reliability of curve

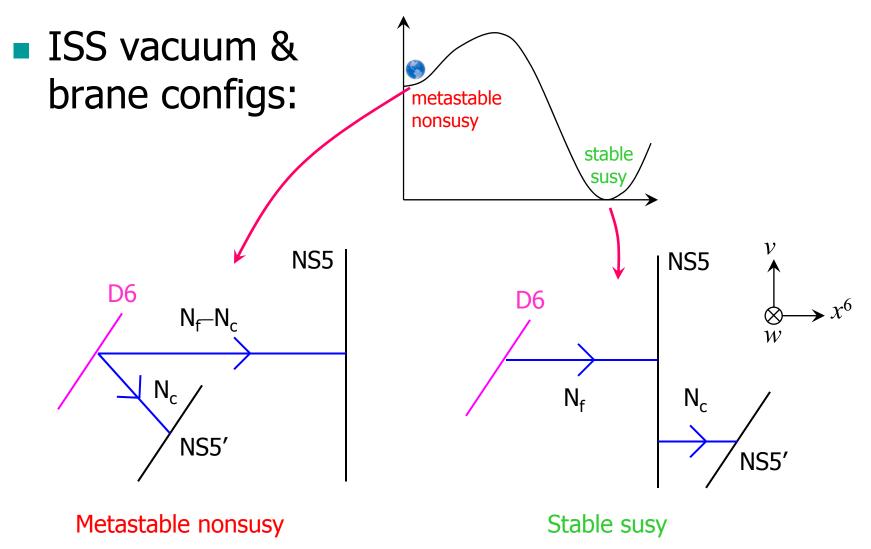
- Why is g<sub>s</sub>N controlling "softness"?
  - Tree level  $\rightarrow$  soft
  - Higher genus can destroy this structure:

$$\sim g_s^{2n}(g_sN)^{h-1}$$



# Boundary Conditions

### The issue

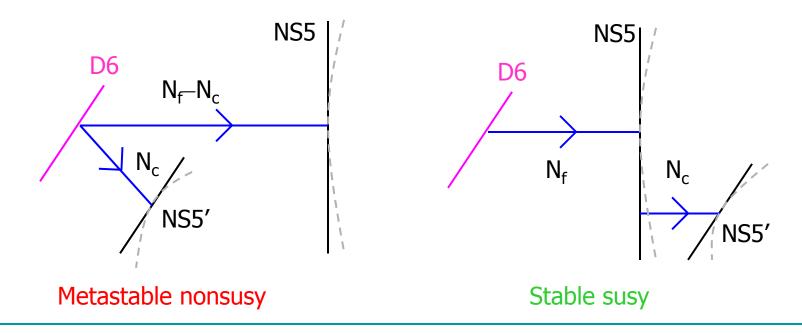


### The issue

#### BGHSS:

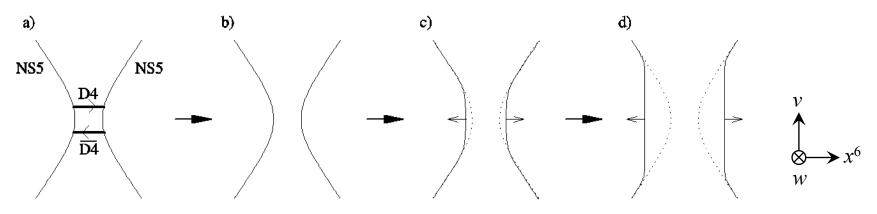
#### Boundary conds. different.

- Two brane configs are in different theories
- One can't decay into other

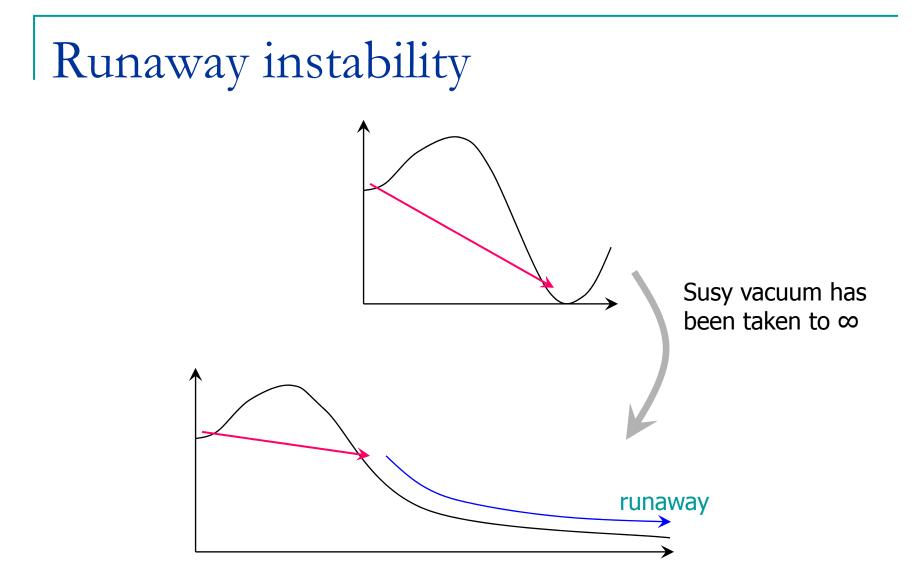


# Nonsusy configs decay

 Does not mean our nonsusy config is stable: It decays

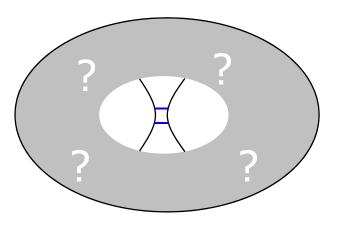


- D4/D4 pair annihilate by quantum tunneling (Cf. β-decay in EM field)
- NS5's straighten
- Takes ∞ time (runaway)



Relevance of nonsusy brane configs

- Gauge theory is approximation of string ③
- Embed in compact CY: BC arises from dynamics of "the rest"



- In the whole sys, runaway ends
- Local "building block" for model building

## Conclusions

- M5/NS5 curve:
  - Can explore nonsusy landscape of string/M-theory
  - Can be easily generalized (ADE, etc.)
- $g_s N$  controls "softness" of breaking
  - Can study string/M landscape in a controlled way
- Boundary conditions
- Building blocks for model building