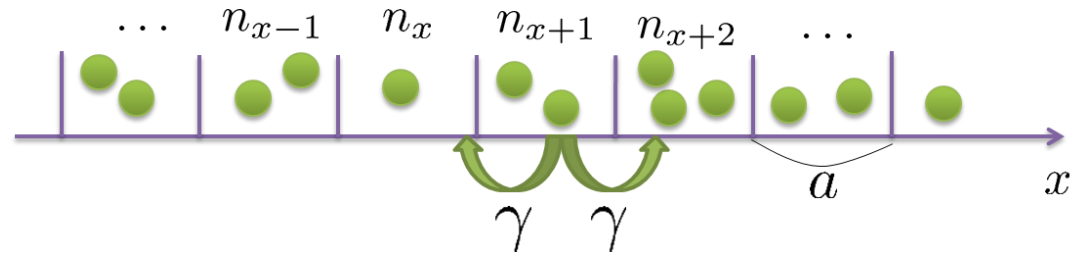
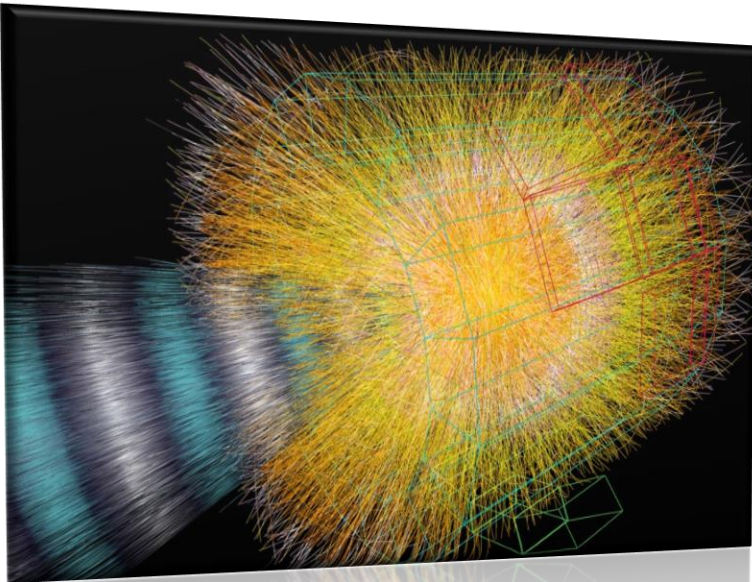


重イオン衝突実験による クォーク物質の生成と 非ガウスゆらぎの時間発展



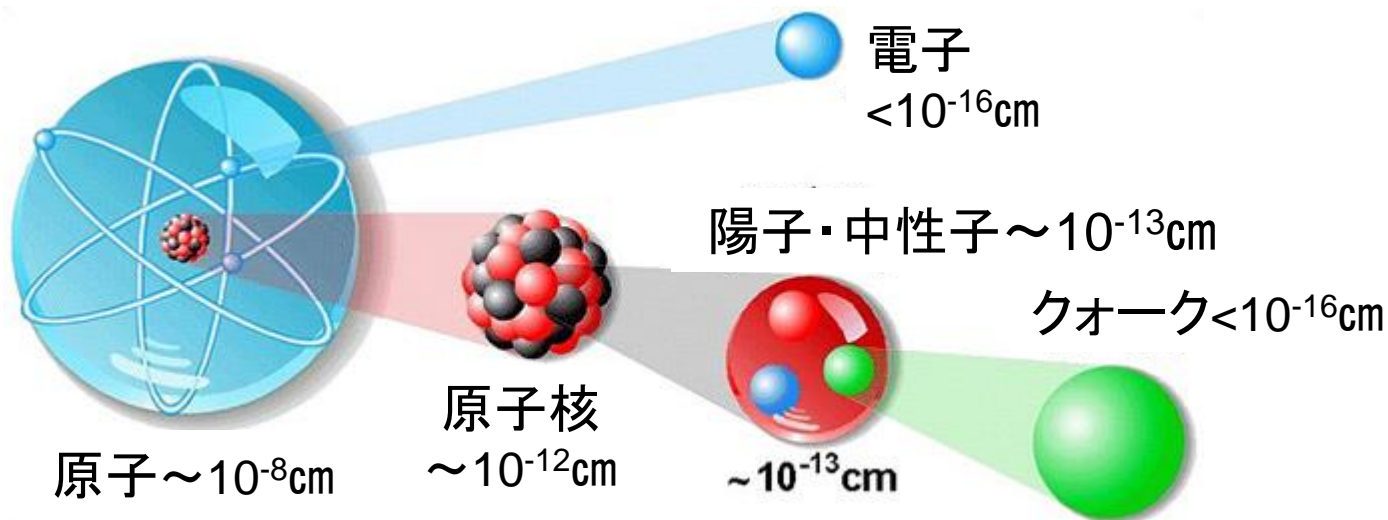
北沢 正清

(阪大物理・原子核理論)

MK, Asakawa, Ono, Physics Letters B, **728**, 386 (2013)

交通流シンポ、名古屋大学、2013/12/16

ハドロンとクォーク

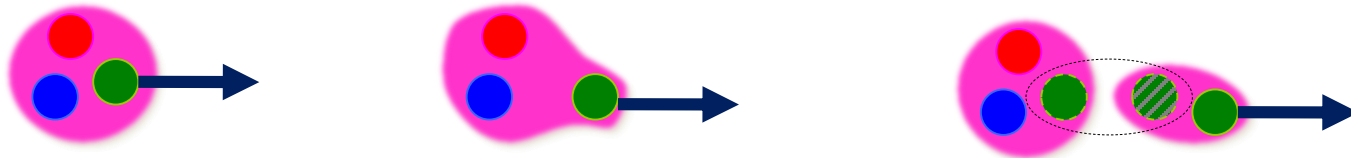


□ 量子色力学 (QCD)

$$\mathcal{L} = \bar{\psi}(i\not{D} - m)\psi - \frac{1}{4}F_{\mu\nu,a}F_a^{\mu\nu}$$

クォークの閉じ込め

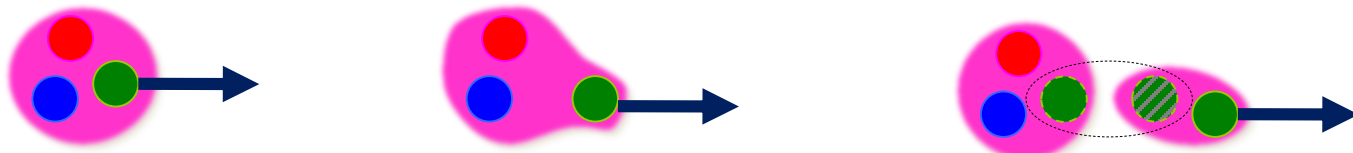
- クォークは、単独では観測できない
- 無理矢理引っ張り出そうとすると...



クォークの閉じ込め

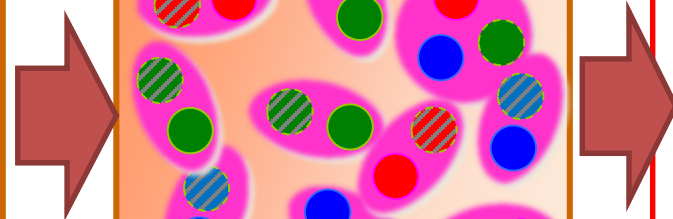
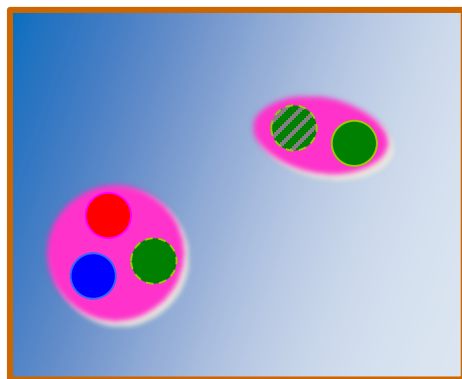
□ クォークは、単独では観測できない

□ 無理矢理引っ張り出そうとすると...



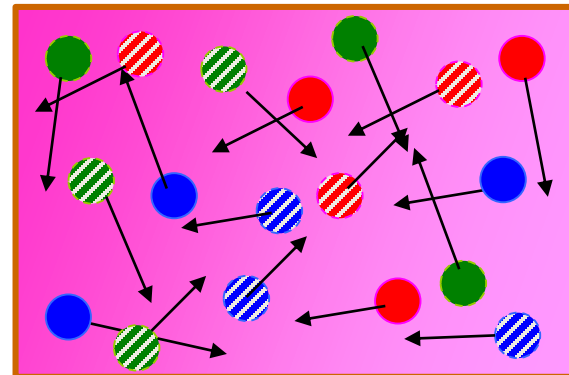
□ クォーク・グルオン・プラズマ (QGP)

□ 物質の温度を上昇させていくと...



空間がハドロンで
埋め尽くされてきて

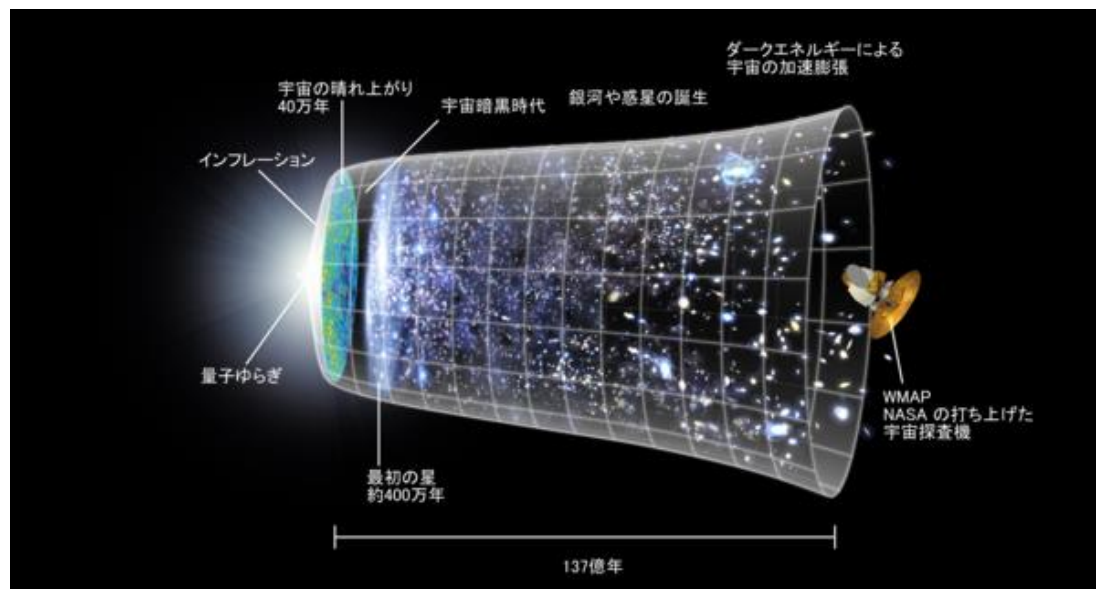
約2兆度に達すると



クォークが自由に飛び
回る世界が実現する！

クォークの閉じ込め

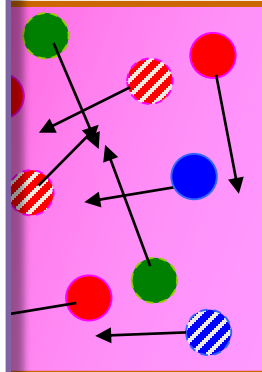
- クォークは、単独では観測できない
 - 無理矢理引っ張り出そうとすると...



- クォーク
- 物

QGP状態は、誕生直後の宇宙の姿
($t < 10^{-5} \text{s}$)

二達すると



空間がハドロンで埋め尽くされてきて

クォークが自由に飛び回る世界が実現する！

加速器による原子核衝突実験



RHIC

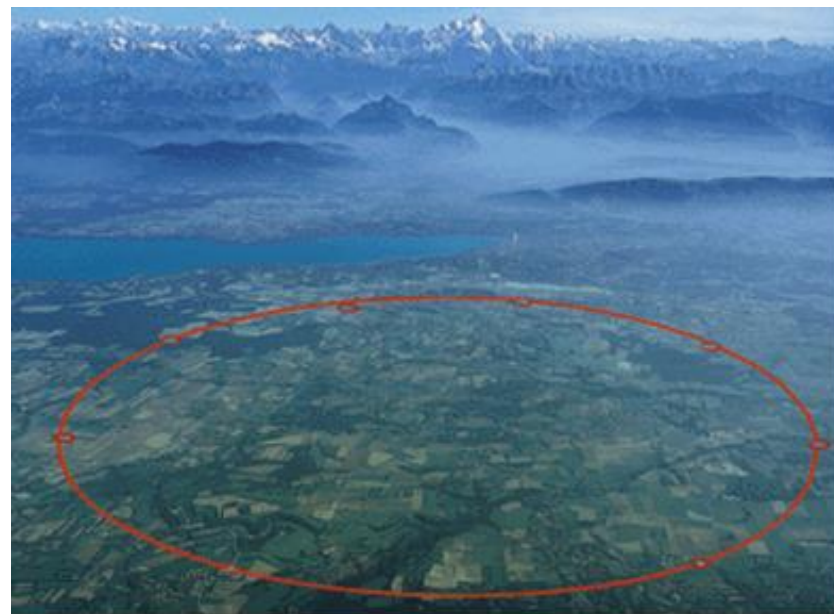
アメリカ

2000年～

全長6km

光速の99.996%

約4兆度



LHC

スイス・フランス

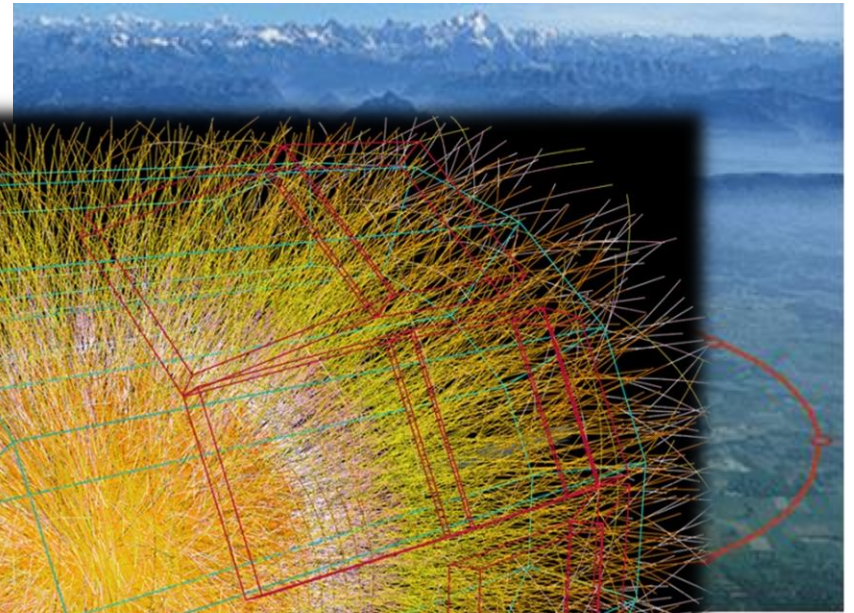
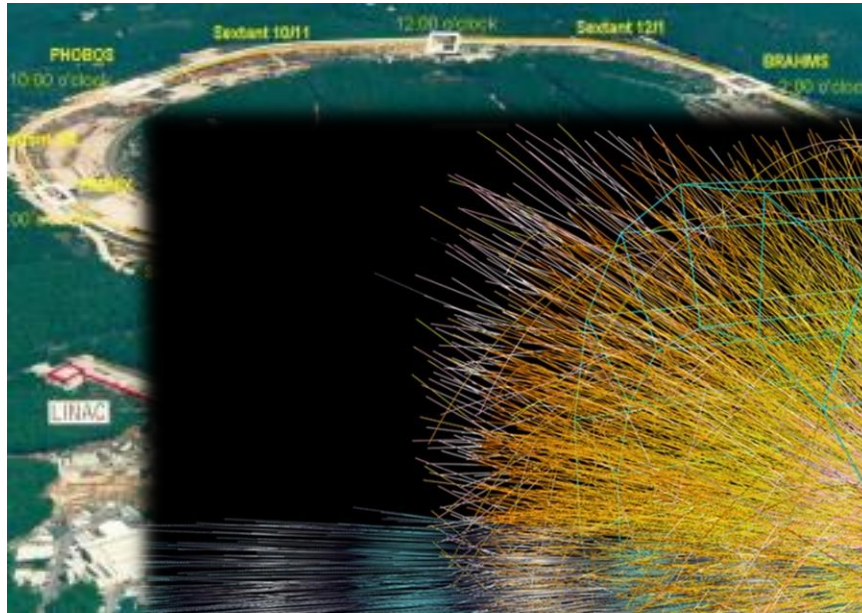
2010年～

全長30km

光速の99.9999%

約6兆度

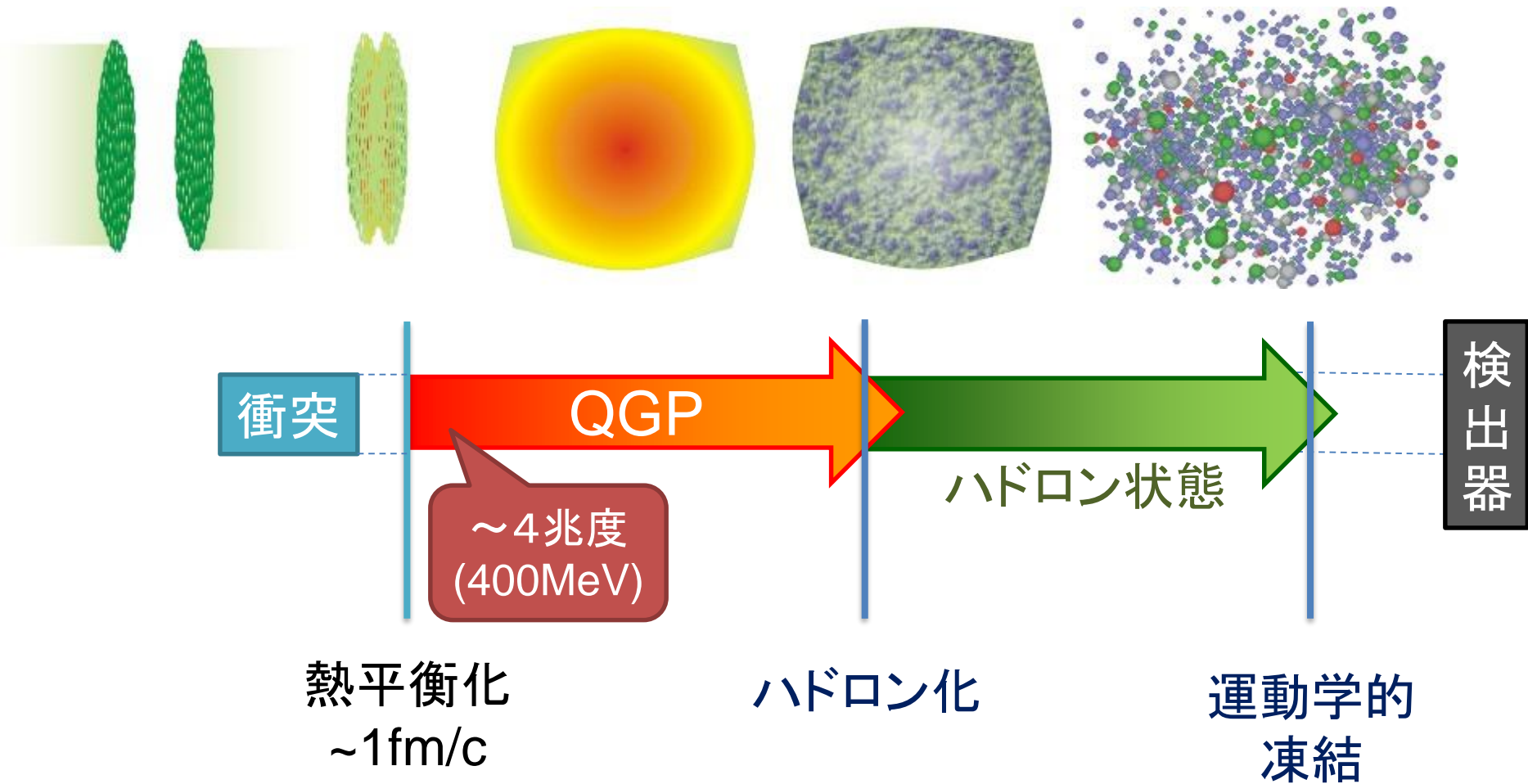
加速器による原子核衝突実験



RHIC
アメリカ
2000年
全長
光速
約4兆度

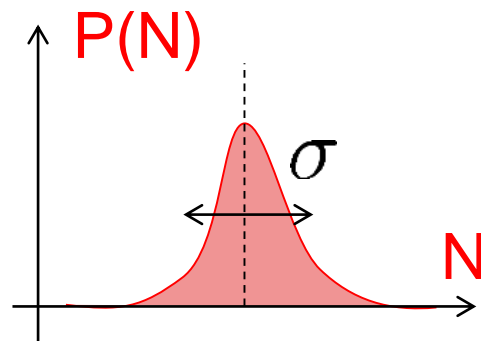
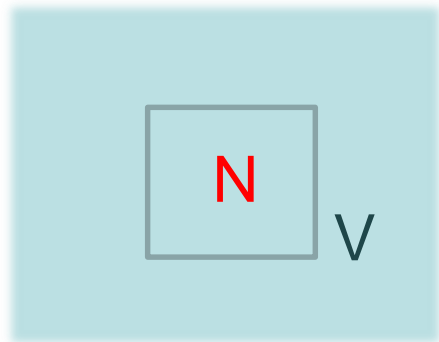
LHC
フランス
2008年～
全長
光速
約6兆度

高温物質の時間発展



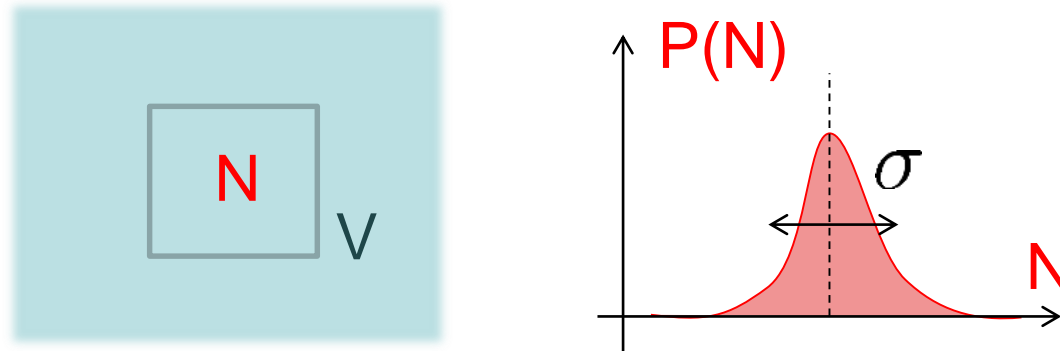
ゆらぎ

熱平衡状態にある物質の観測量は、ゆらいでいる。



ゆらぎ

熱平衡状態にある物質の観測量は、ゆらいでいる。



➤ Variance: $\langle \delta N^2 \rangle = V \chi_2 = \sigma^2$

$$\delta N = N - \langle N \rangle$$

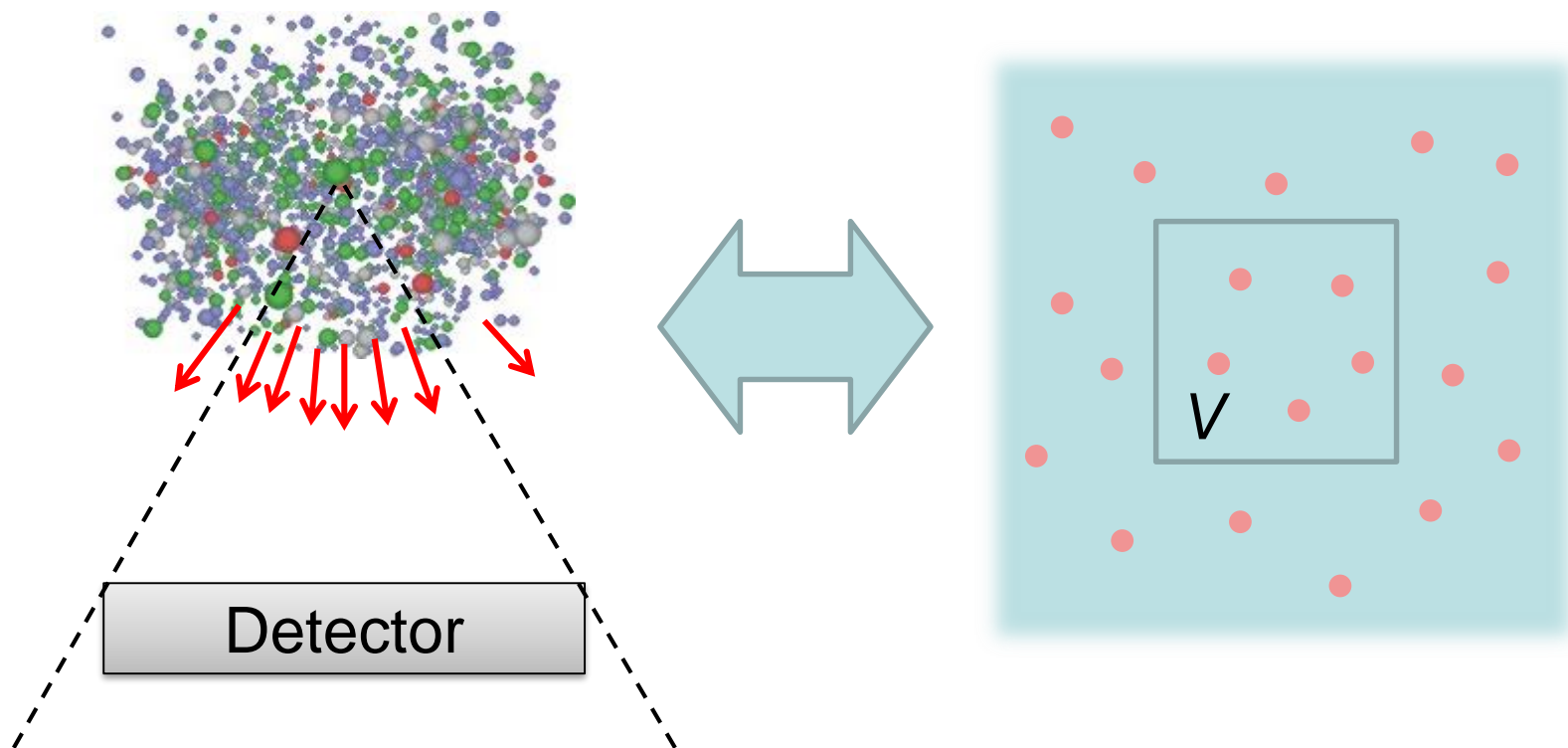
➤ Skewness: $S = \frac{\langle \delta N^3 \rangle}{\sigma^3}$

➤ Kurtosis: $\kappa = \frac{\langle \delta N^4 \rangle - 3\langle \delta N^2 \rangle^2}{\chi_2 \sigma^2}$

Non-Gaussianity

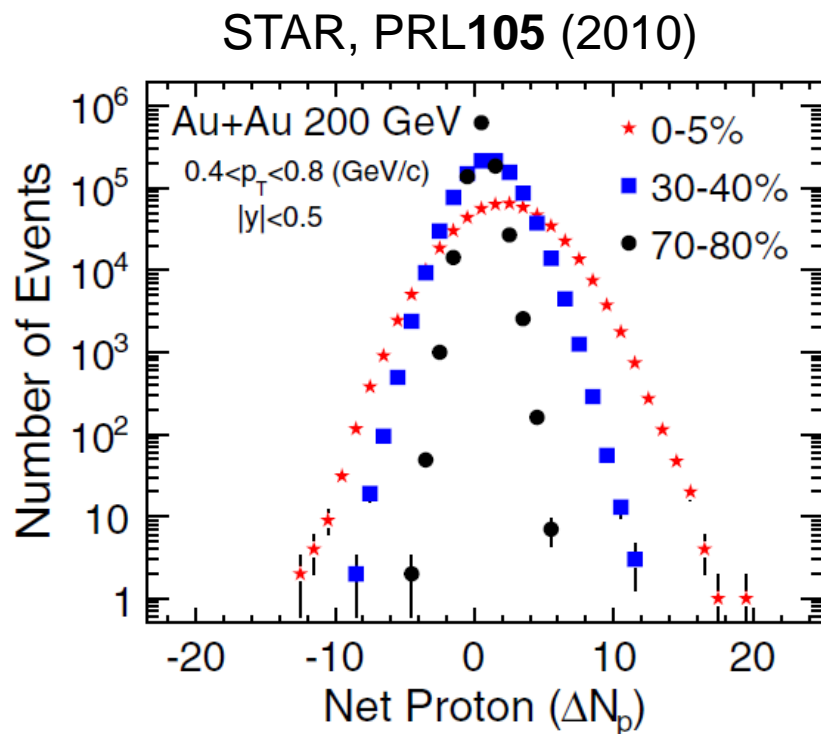
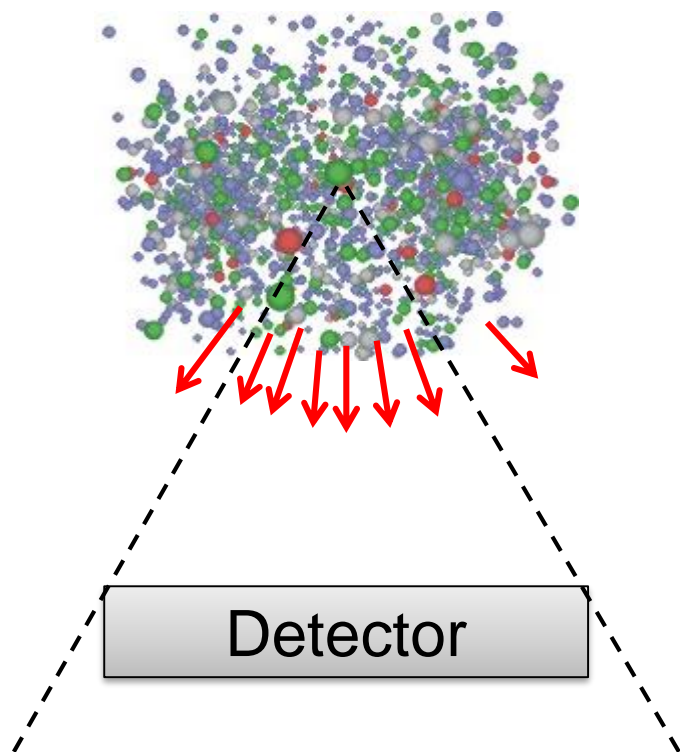
イベント-バイ-イベント解析 @ HIC

- 重イオン衝突の、各衝突イベントごとの終状態ゆらぎを測定する。



イベント-バイ-イベント解析 @ HIC

- 重イオン衝突の、各衝突イベントごとの終状態ゆらぎを測定する。

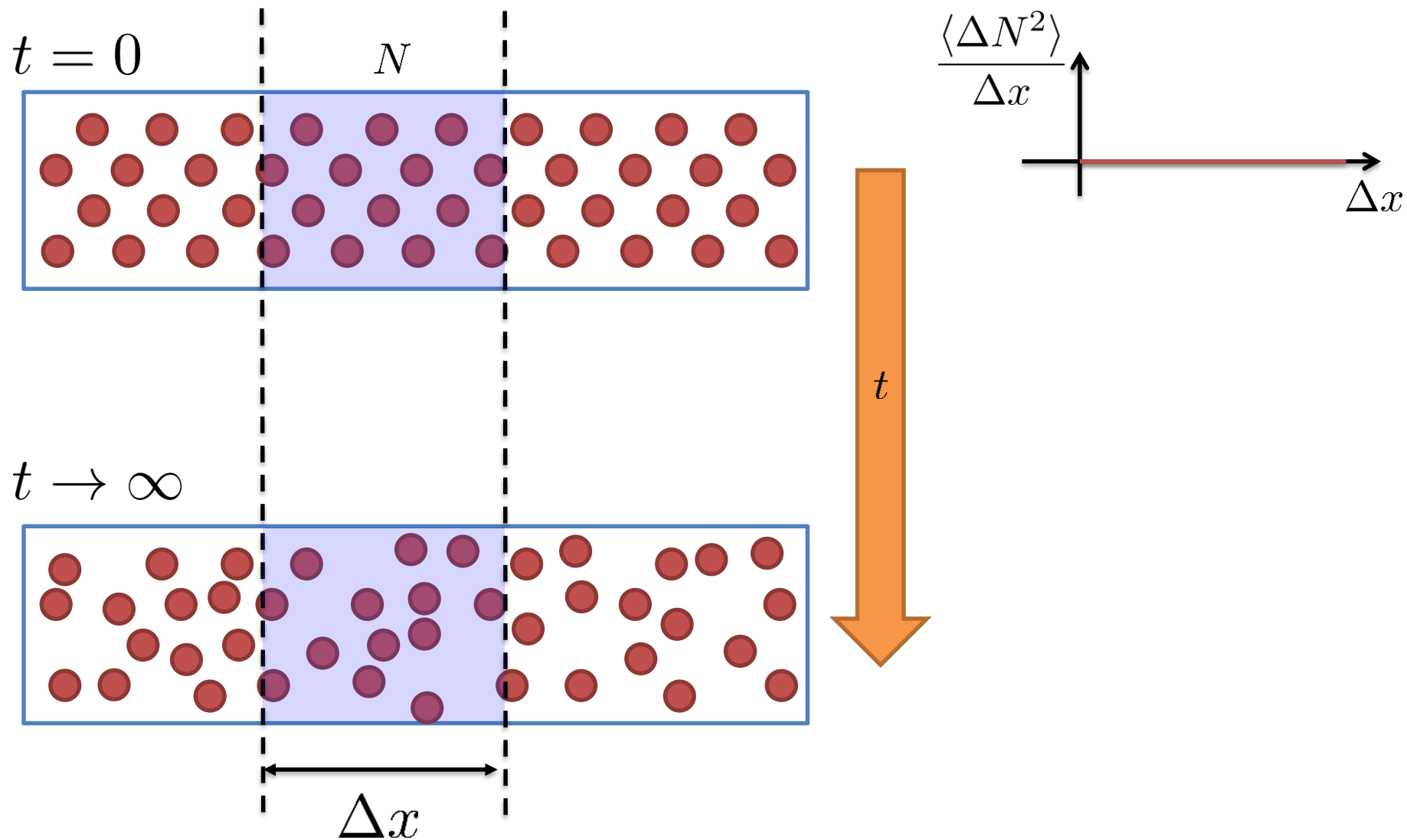


以下では、
電磁電荷、バリオン数を考察

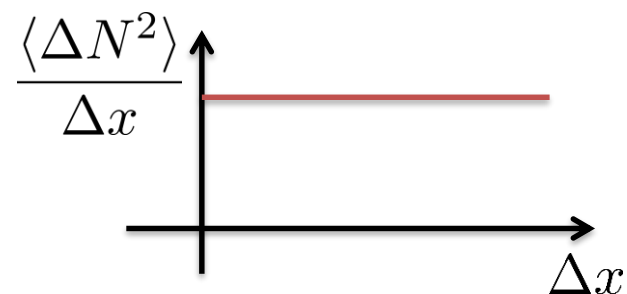
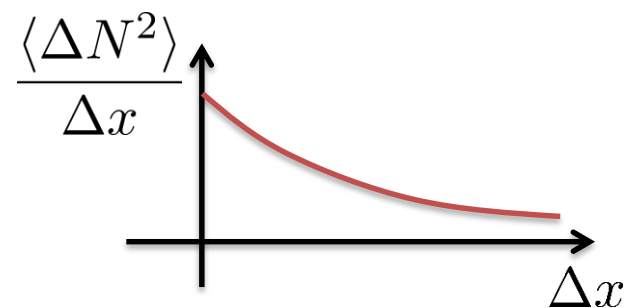
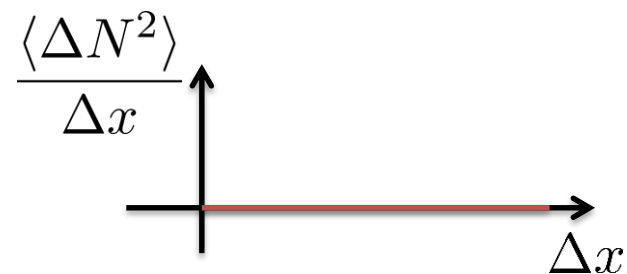
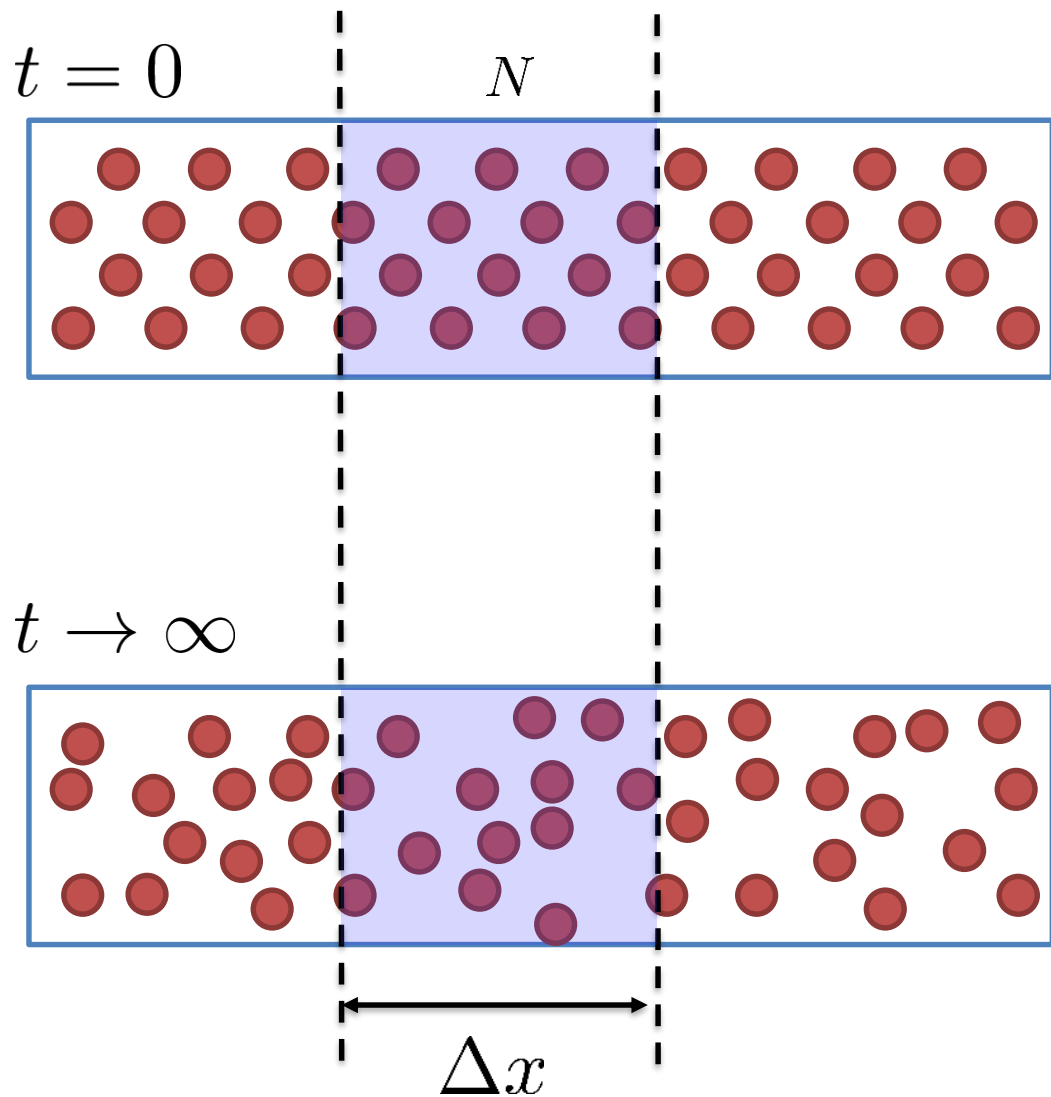
↓

$$\langle \delta N^2 \rangle, \langle \delta N^3 \rangle, \langle \delta N^4 \rangle_c, \dots$$

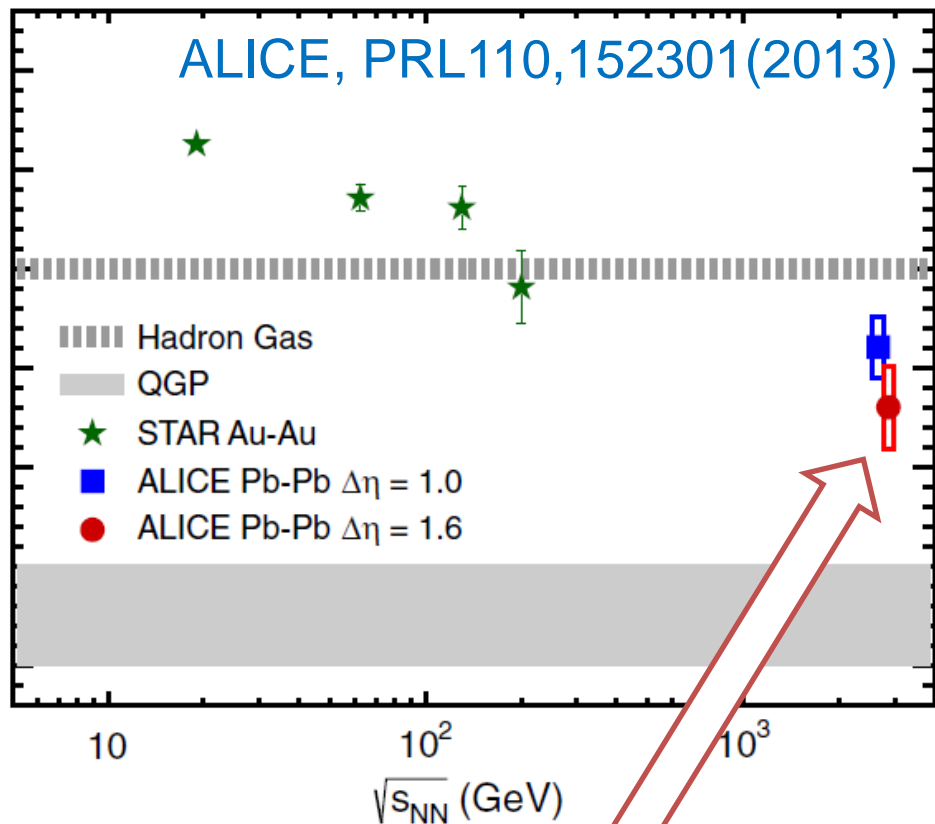
保存電荷ゆらぎの散逸過程



保存電荷ゆらぎの散逸過程



電磁電荷2次ゆらぎ @ LHC



ハドロン状態
熱平衡値

D-measure

$$D = 4 \frac{\langle \delta N_Q^2 \rangle}{\langle N_Q^+ + N_Q^- \rangle}$$

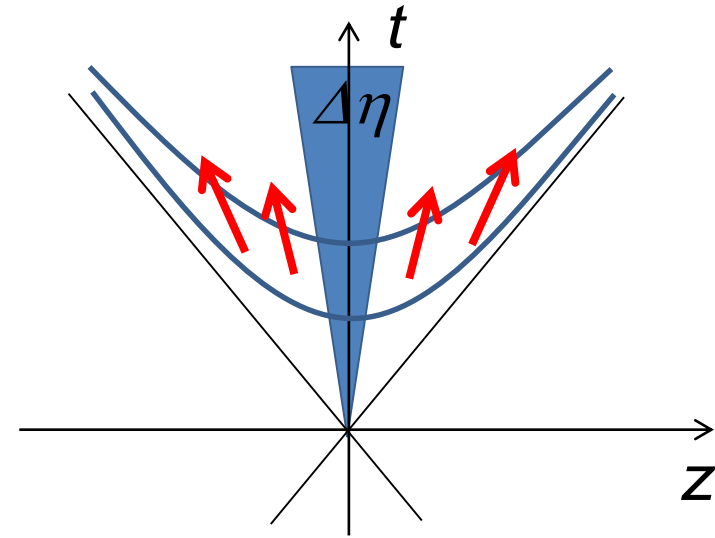
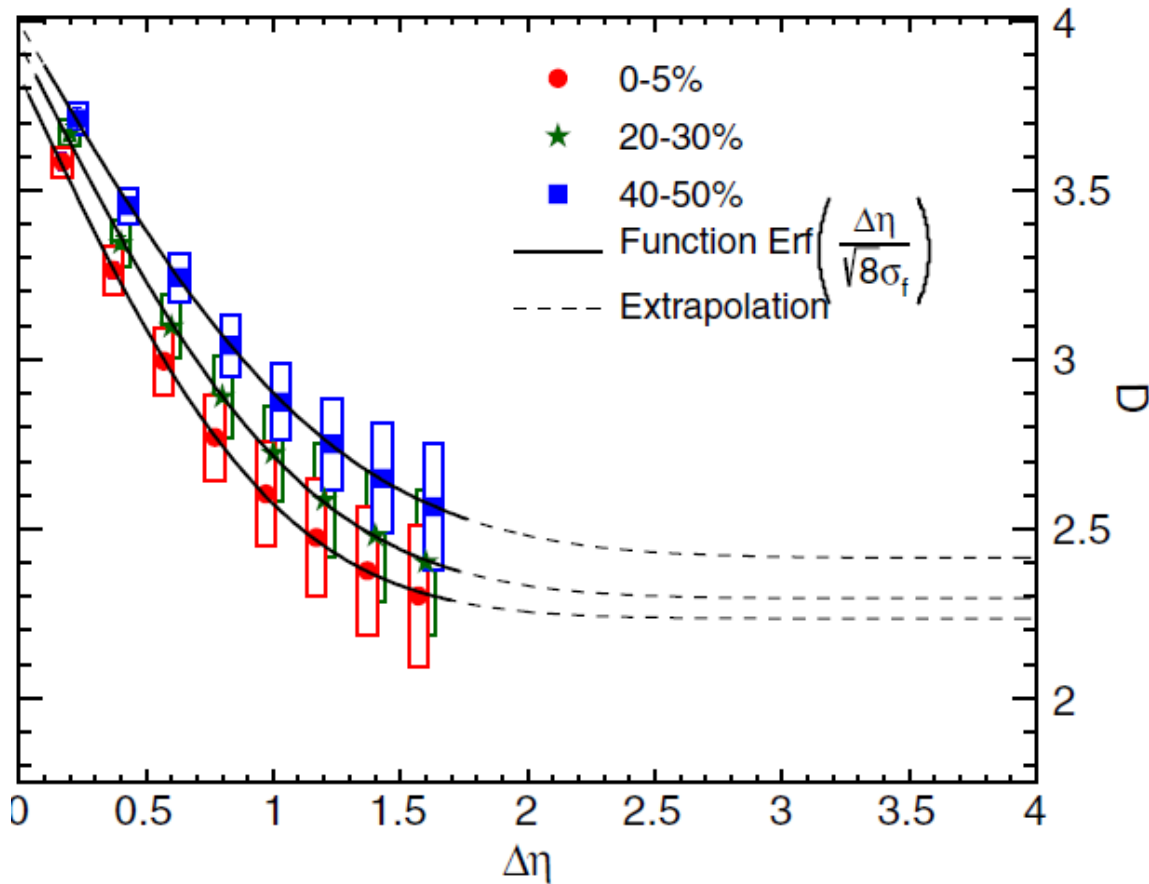
- $D \sim 3-4$ Hadronic
- $D \sim 1-1.5$ Quark

QGP状態
熱平衡値

LHCでの電磁電荷ゆらぎは
ハドロン状態熱平衡値より
抑制されている！！

$\Delta\eta$ Dependence @ ALICE

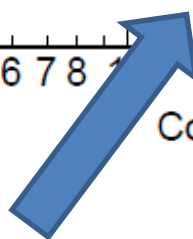
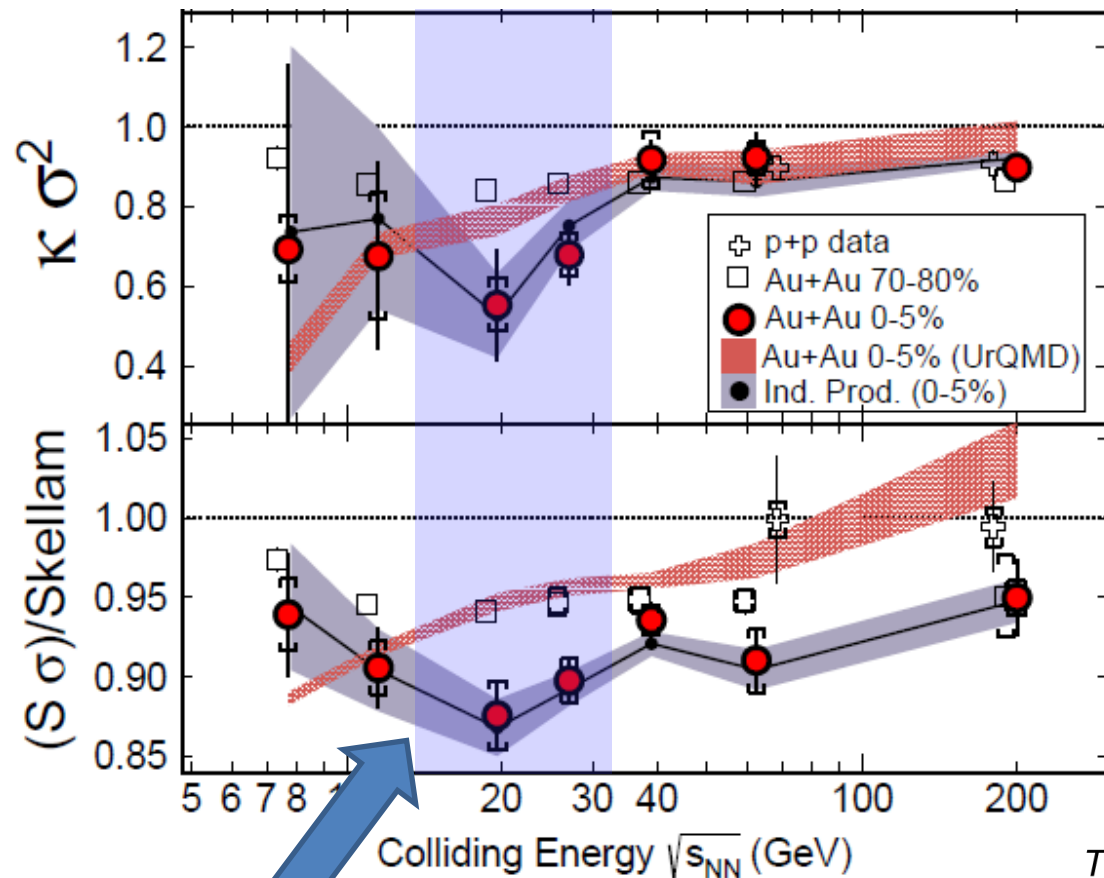
ALICE
PRL 2013



$\Delta\eta$: 粒子を数えるラピデティ幅

陽子数ゆらぎ @ RHIC

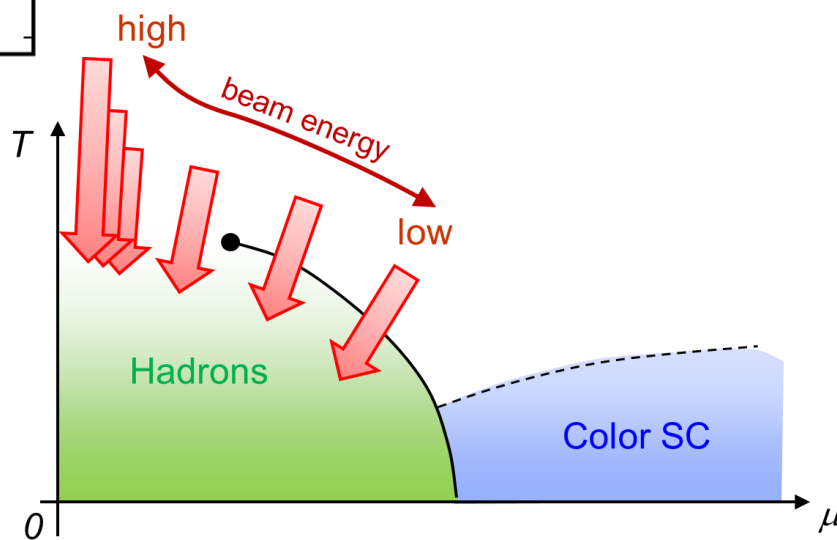
STAR, 1309.5681



Something interesting??

$$\frac{\langle \delta N_p^4 \rangle_c}{\langle \delta N_p^2 \rangle}$$

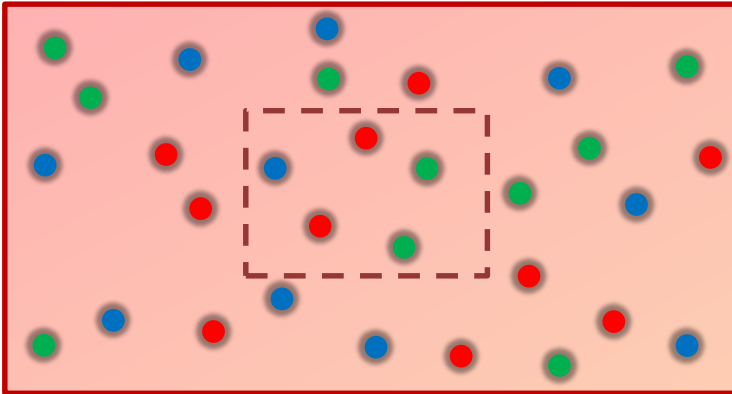
$$\frac{\langle \delta N_p^3 \rangle}{\langle N_p \rangle}$$



Fluctuations

Free Boltzmann \rightarrow Poisson

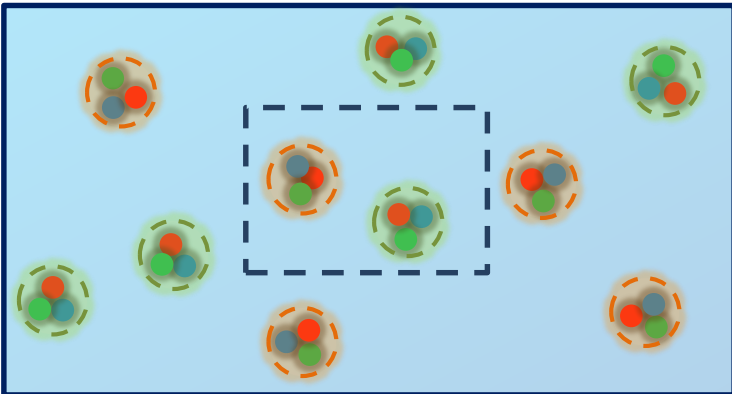
$$\langle \delta N^n \rangle_c = \langle N \rangle$$



$$\langle \delta N_q^n \rangle_c = \langle N_q \rangle$$

$$\rightarrow \langle \delta N_B^n \rangle_c = \frac{1}{3^{n-1}} \langle N_B \rangle$$

$$3N_B = N_q$$

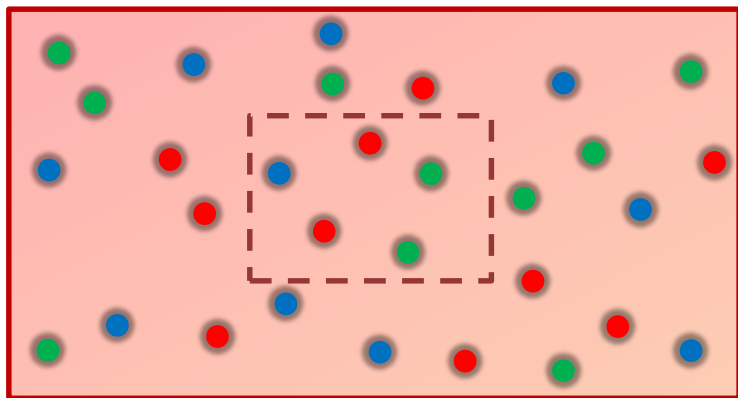


$$\langle \delta N_B^n \rangle_c = \langle N_B \rangle$$

Fluctuations

Free Boltzmann \rightarrow Poisson

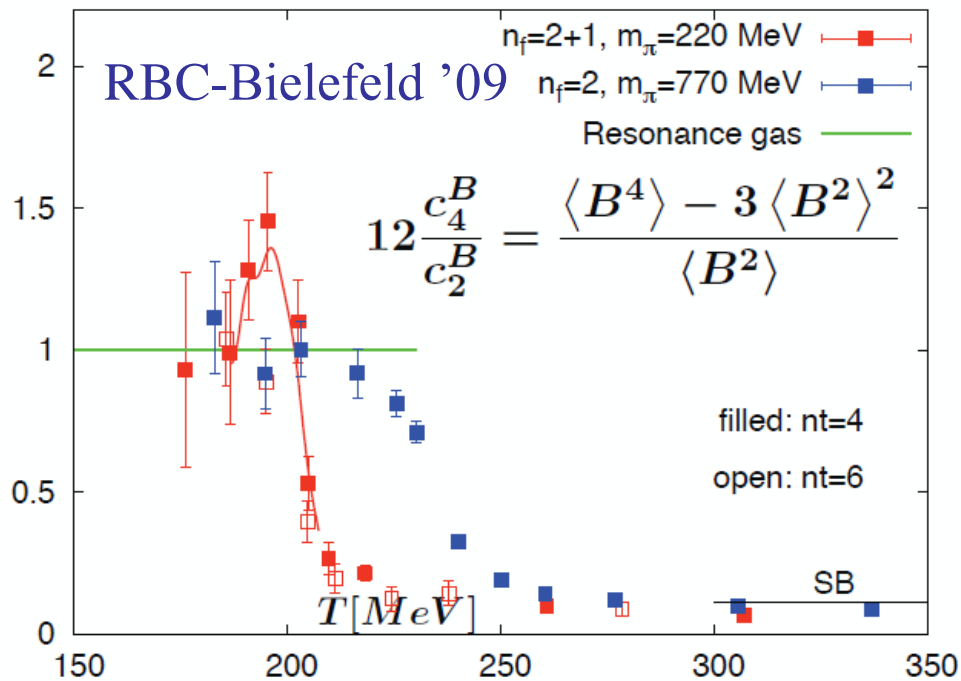
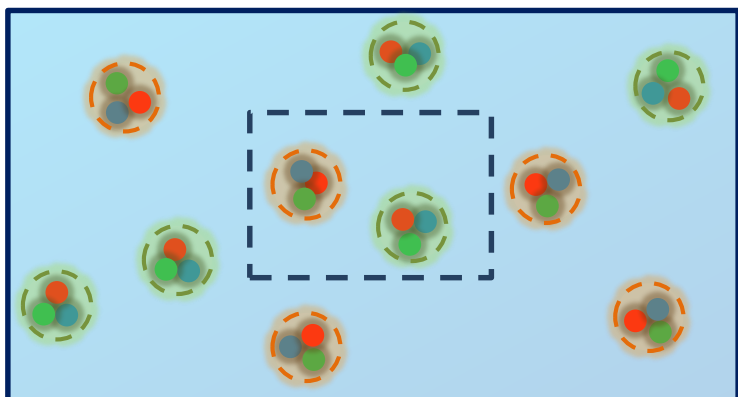
$$\langle \delta N^n \rangle_c = \langle N \rangle$$



$$\langle \delta N_q^n \rangle_c = \langle N_q \rangle$$

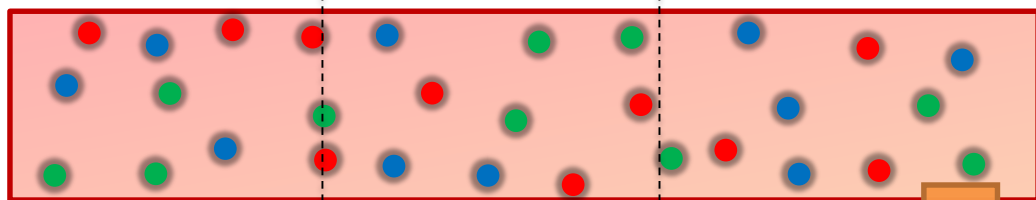
$$\Rightarrow \langle \delta N_B^n \rangle_c = \frac{1}{3^{n-1}} \langle N_B \rangle$$

$$3N_B = N_q$$

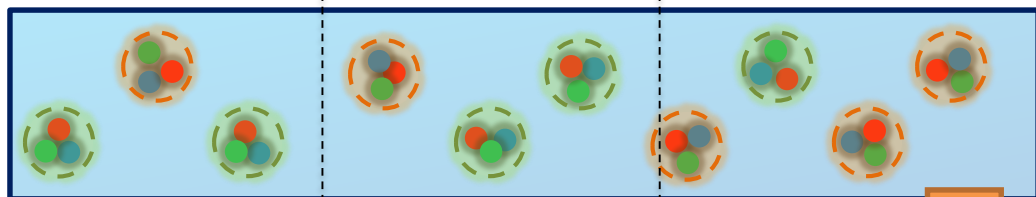


Time Evolution in HIC

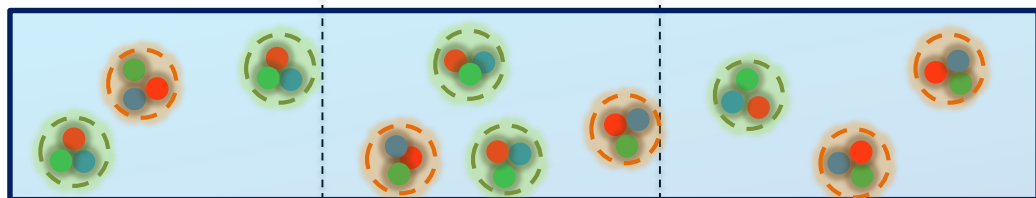
Quark-Gluon Plasma



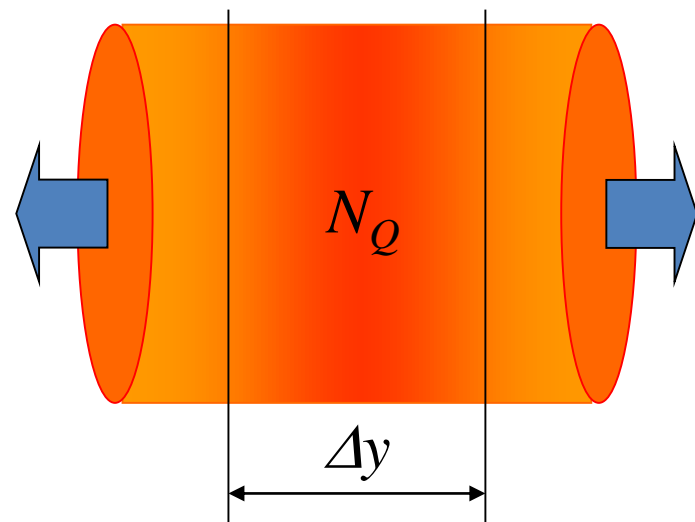
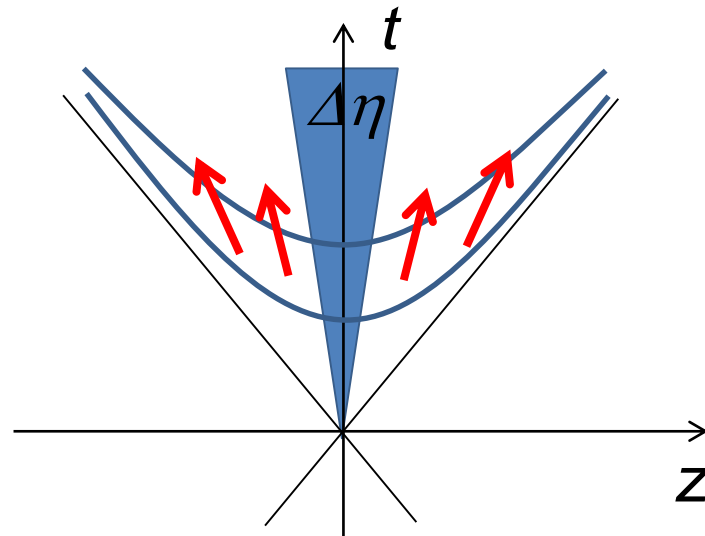
Hadronization



Freezeout

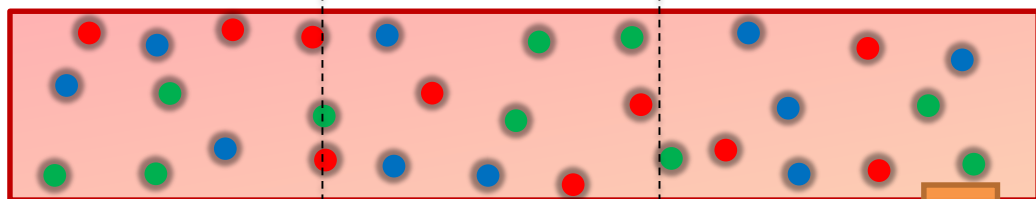


$\Delta\eta$

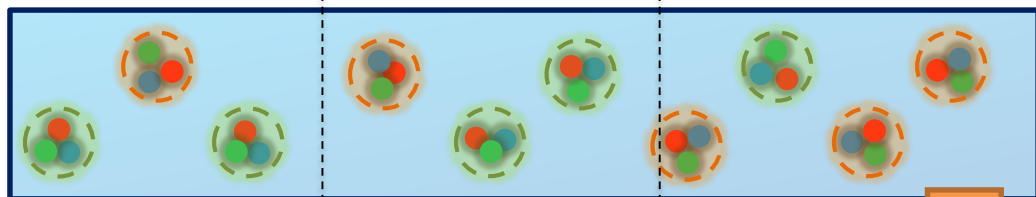


Time Evolution in HIC

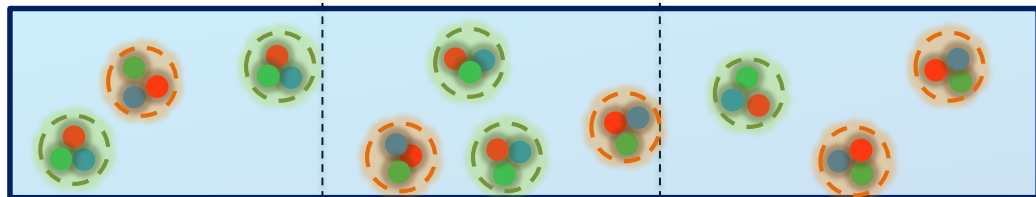
Quark-Gluon Plasma



Hadronization



Freezeout



$\Delta\eta$

$$\frac{\langle \Delta N^2 \rangle}{\Delta\eta}$$

χ_{HAD}

χ_{QGP}

$\Delta\eta$

χ_{HAD}

χ_{QGP}

$\Delta\eta$

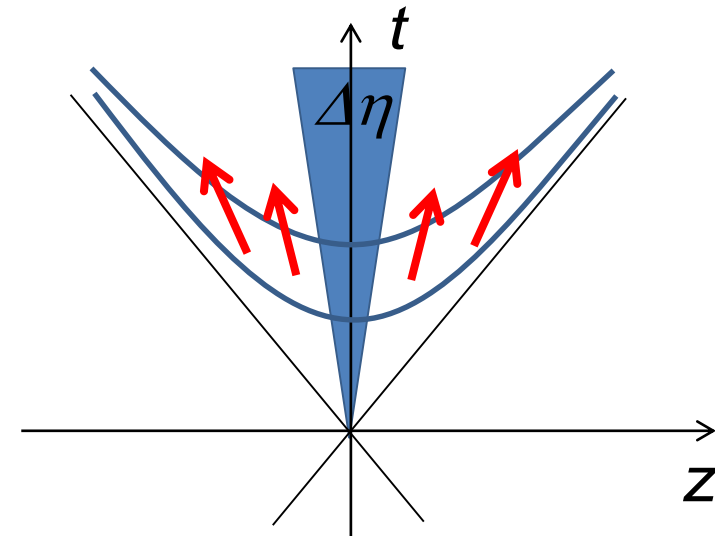
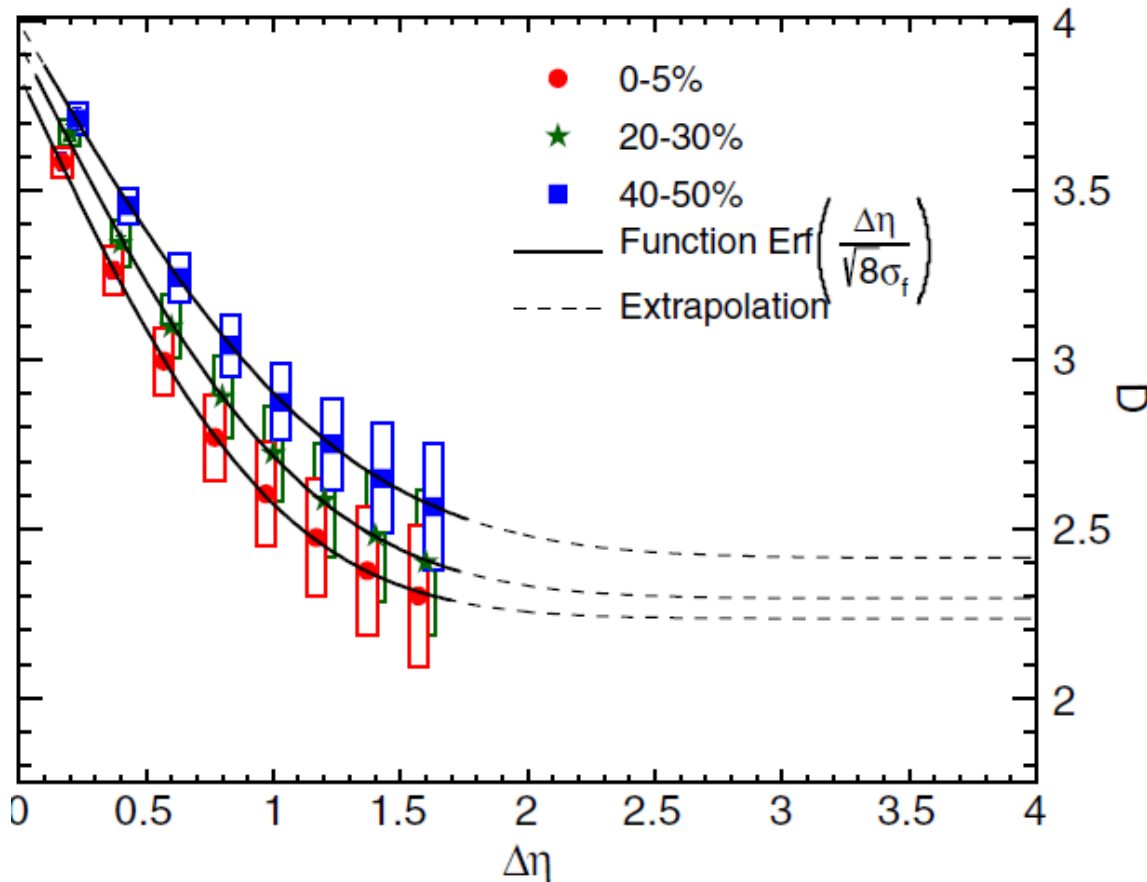
χ_{HAD}

χ_{QGP}

$\Delta\eta$

$\Delta\eta$ Dependence @ ALICE

ALICE
PRL 2013



$\Delta\eta$: ラピッドティ幅

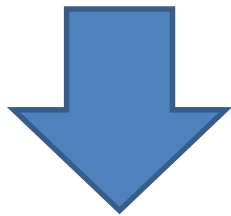
2次ゆらぎの $\Delta\eta$ 依存性はOK。
非ガウスゆらぎの非平衡性は、
 $\Delta\eta$ 依存性にどのように反映するか？

流体ゆらぎの理論

Landau, Lifshitz, Statistical Mechanics II
Stephanov, Shuryak, 2001

確率論的拡散方程式

$$\partial_{\tau} n = D \partial_{\eta}^2 n + \partial_{\eta} \xi(\eta, \tau)$$



Gaussian fluctuation
in equilibrium

Markov (temporary local)
+
continuity

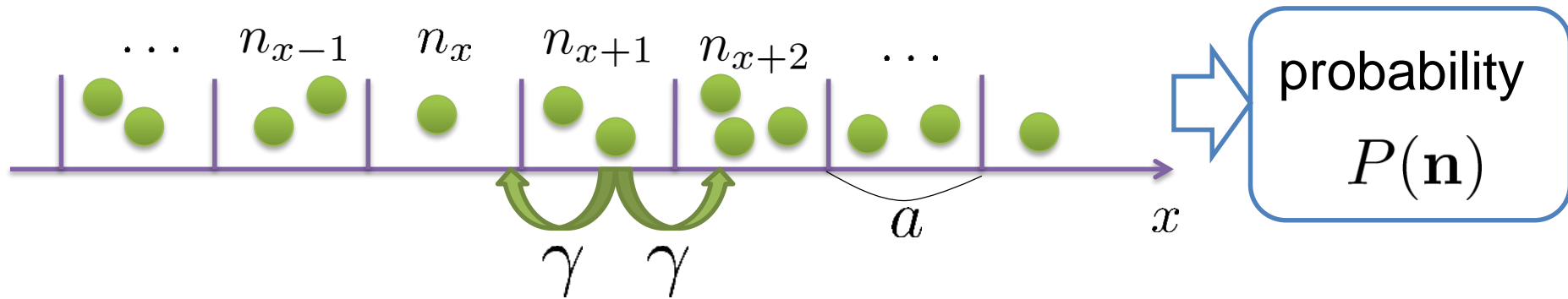


Gaussian

cf) Gardiner, "Stochastic Methods"

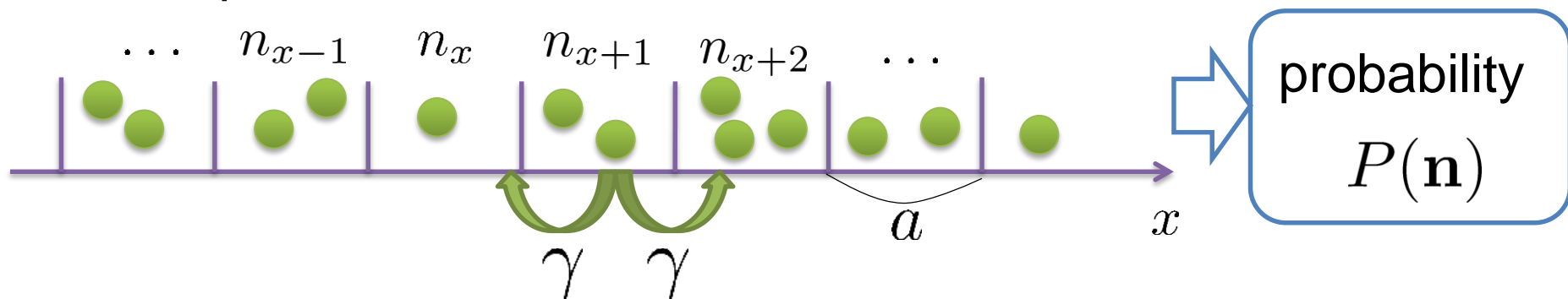
拡散マスター方程式

Divide spatial coordinate into discrete cells



拡散マスター方程式

Divide spatial coordinate into discrete cells



Master Equation for $P(n)$

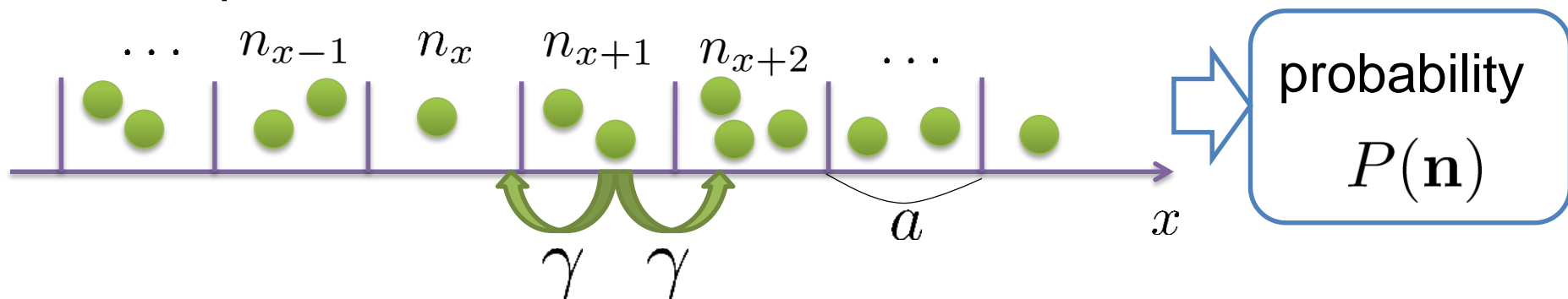
$$\frac{\partial}{\partial t} P(\mathbf{n}) = \gamma \sum_x [(n_x + 1) \{P(\mathbf{n} + \mathbf{e}_x - \mathbf{e}_{x+1}) + P(\mathbf{n} + \mathbf{e}_x - \mathbf{e}_{x-1})\} - 2n_x P(\mathbf{n})]$$

Solve the DME **exactly**, and take $a \rightarrow 0$ limit

No approx., ex. van Kampen's system size expansion

拡散マスター方程式

Divide spatial coordinate into discrete cells



□ 拡散マスター方程式の解

□ 期待値 $\langle n \rangle$

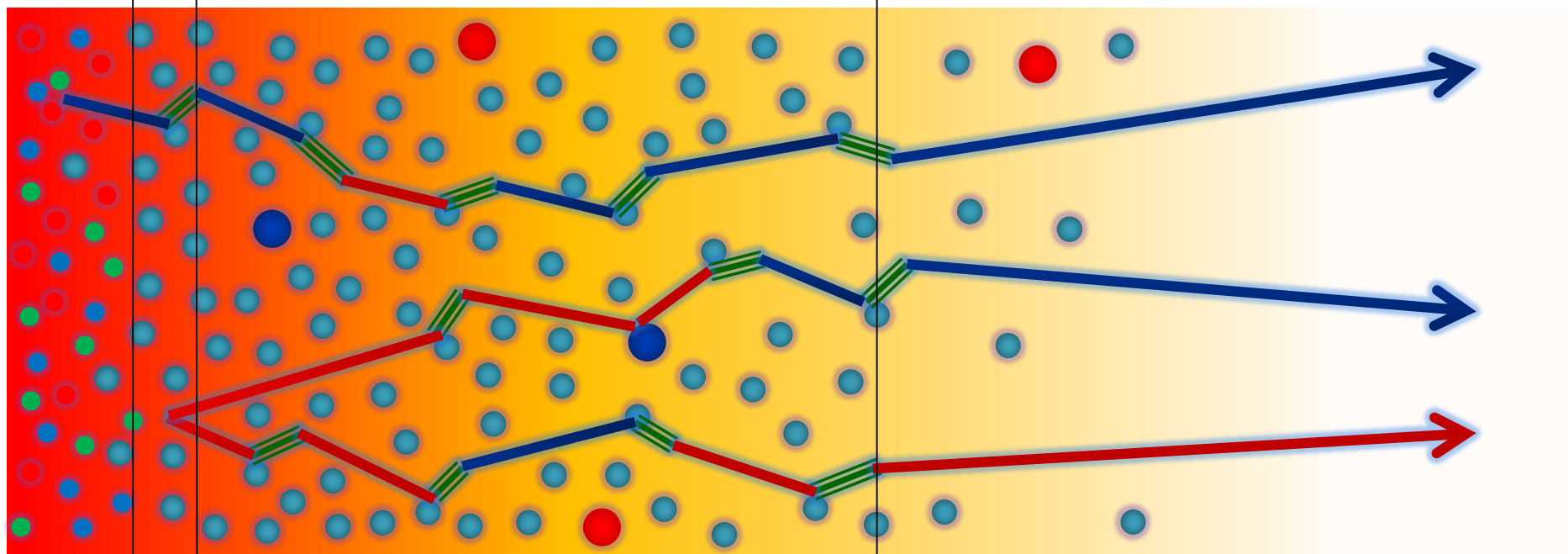
➡ $D = \gamma a^2$ とおけば、拡散方程式と一致

□ 2次(ガウス)ゆらぎ $\langle \delta n^2 \rangle$

➡ 確率論的拡散方程式と一致

Baryons in Hadronic Phase






time →



hadronize
chem. f.o.

← 10~20fm →

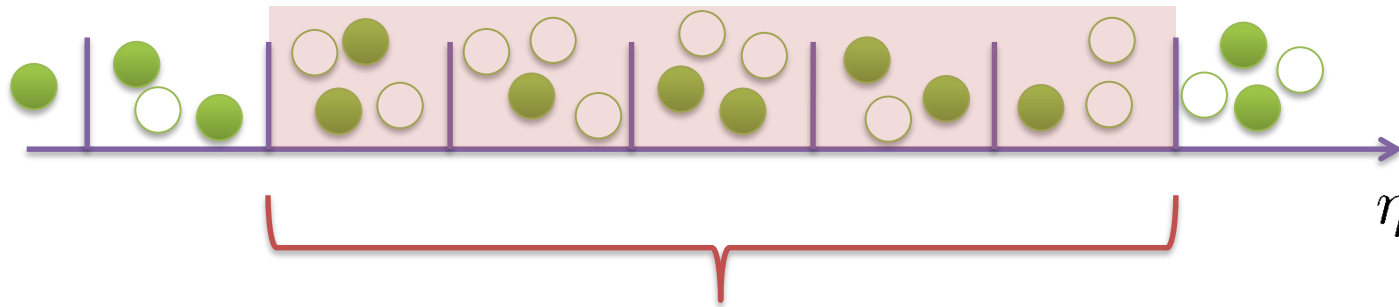
kinetic f.o.

- | | | | |
|--|----------------|---|---------|
|  | p, \bar{p} |  | mesons |
|  | n, \bar{n} |  | baryons |
|  | $\Delta(1232)$ | | |

Baryons behave like
Brownian pollens in water

Net Charge Number

Prepare 2 species of (non-interacting) particles



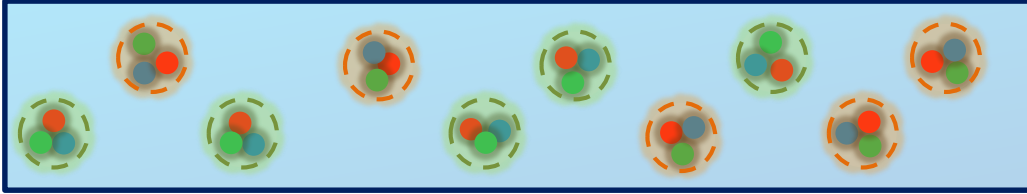
$$\bar{Q}(\tau) = \int_0^{\Delta\eta} d\eta (n_1(\eta, \tau) - n_2(\eta, \tau))$$

Let us investigate

$$\langle \bar{Q}^2 \rangle_c \quad \langle \bar{Q}^4 \rangle_c \quad \text{at freezeout time } t$$

Time Evolution in Hadronic Phase

Hadronization (initial condition)



- Boost invariance / infinitely long system
- Local equilibration / local correlation

$$\langle \bar{Q}^2 \rangle_c$$

$$\langle \bar{Q}^4 \rangle_c$$

$$\langle \bar{Q}^2 Q_{(\text{tot})} \rangle_c$$

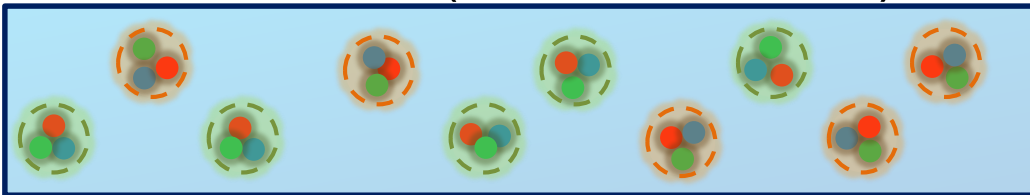
$$\langle Q_{(\text{tot})}^2 \rangle_c$$

suppression owing to
local charge conservation

strongly dependent on
hadronization mechanism

Time Evolution in Hadronic Phase

Hadronization (initial condition)



Time evolution via DME

- Boost invariance / infinitely long system
- Local equilibration / local correlation

$$\langle \bar{Q}^2 \rangle_c$$

$$\langle \bar{Q}^4 \rangle_c$$

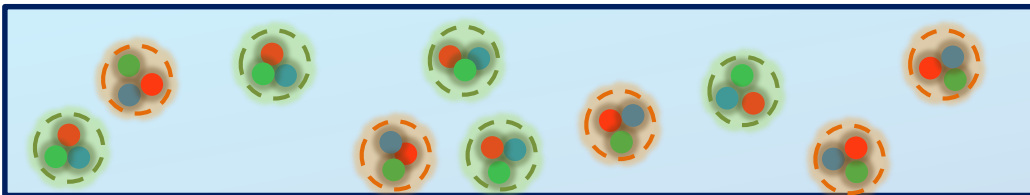
$$\langle \bar{Q}^2 Q_{(\text{tot})} \rangle_c$$

$$\langle Q_{(\text{tot})}^2 \rangle_c$$

suppression owing to local charge conservation

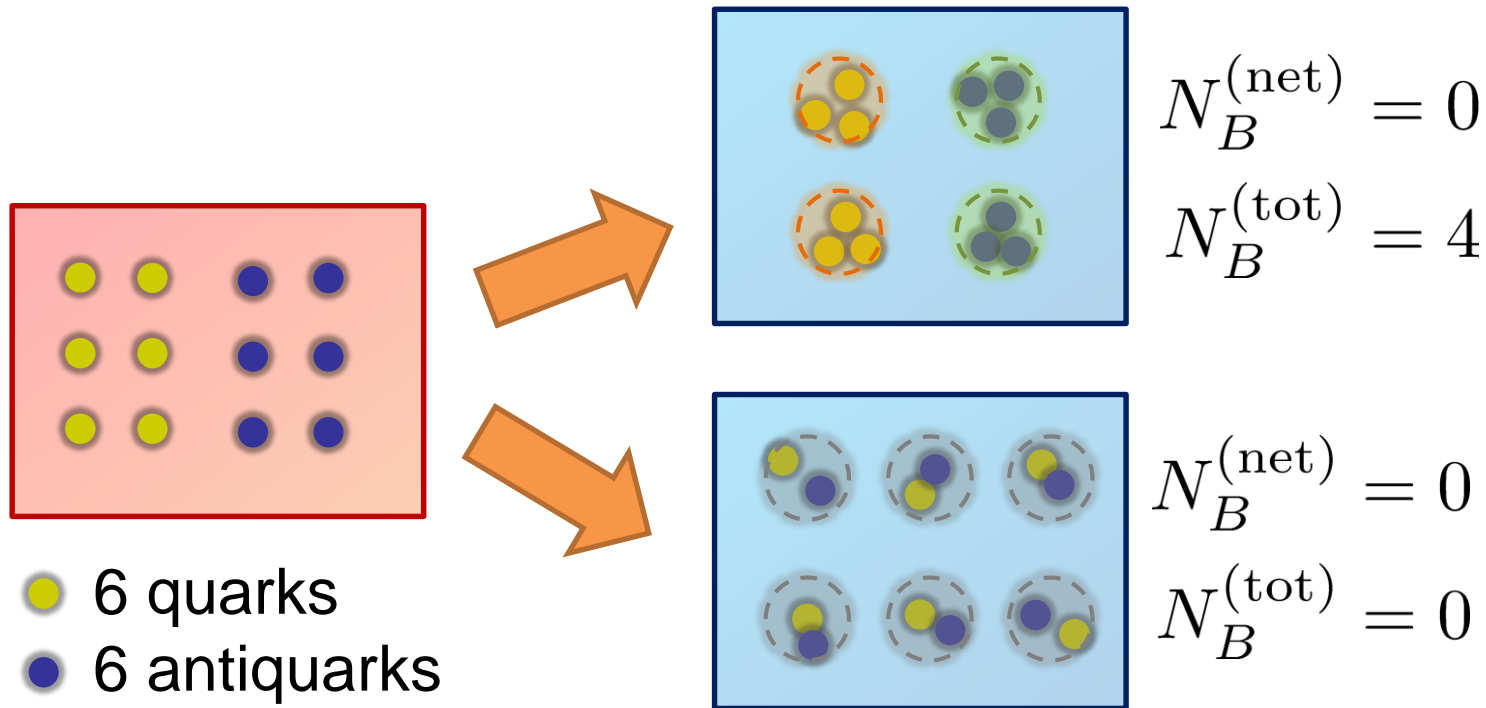
strongly dependent on hadronization mechanism

Freezeout



Total Charge Number Fluctuation

In recombination model,

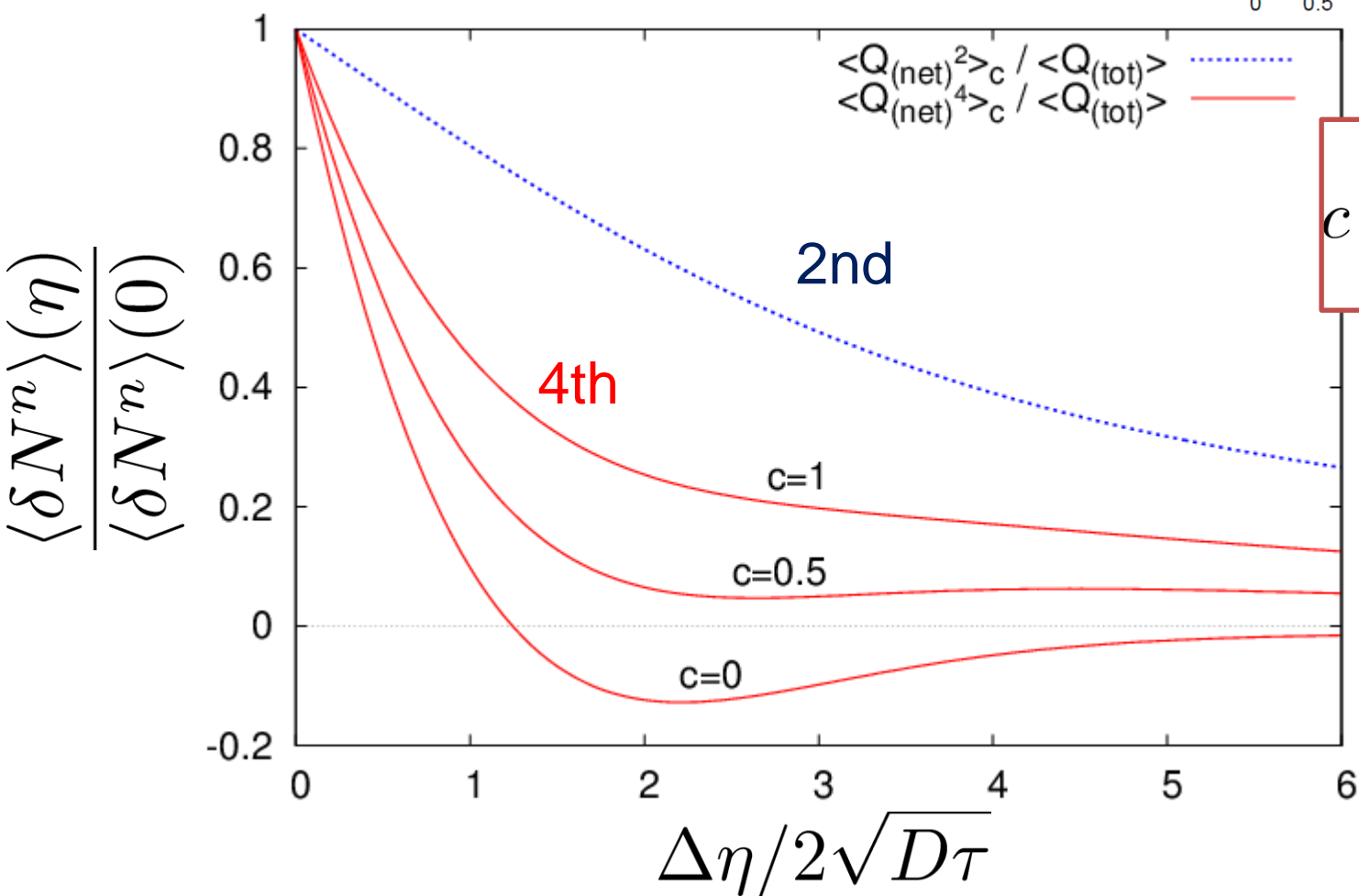
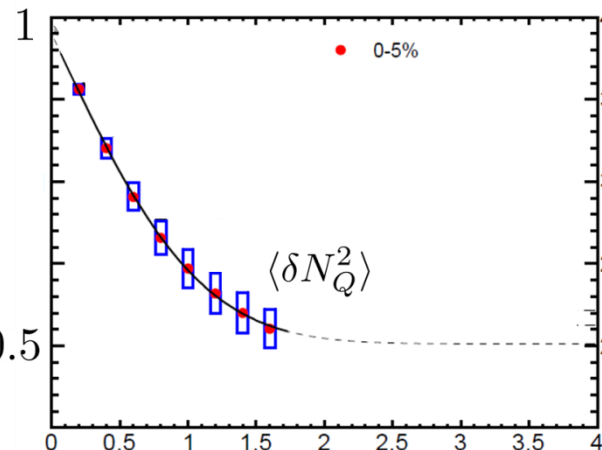


□ $N_B^{(\text{tot})}$ can fluctuate, while $N_B^{(\text{net})}$ does not.

$\Delta\eta$ Dependence at Freezeout

Initial fluctuations:

$$\langle \bar{Q}^2 \rangle_c = \langle \bar{Q}^4 \rangle_c = \langle \bar{Q}^2 Q_{(\text{tot})} \rangle_c = 0$$



$$c = \frac{\langle Q_{(\text{tot})}^2 \rangle_c}{\langle Q_{(\text{tot})} \rangle}$$

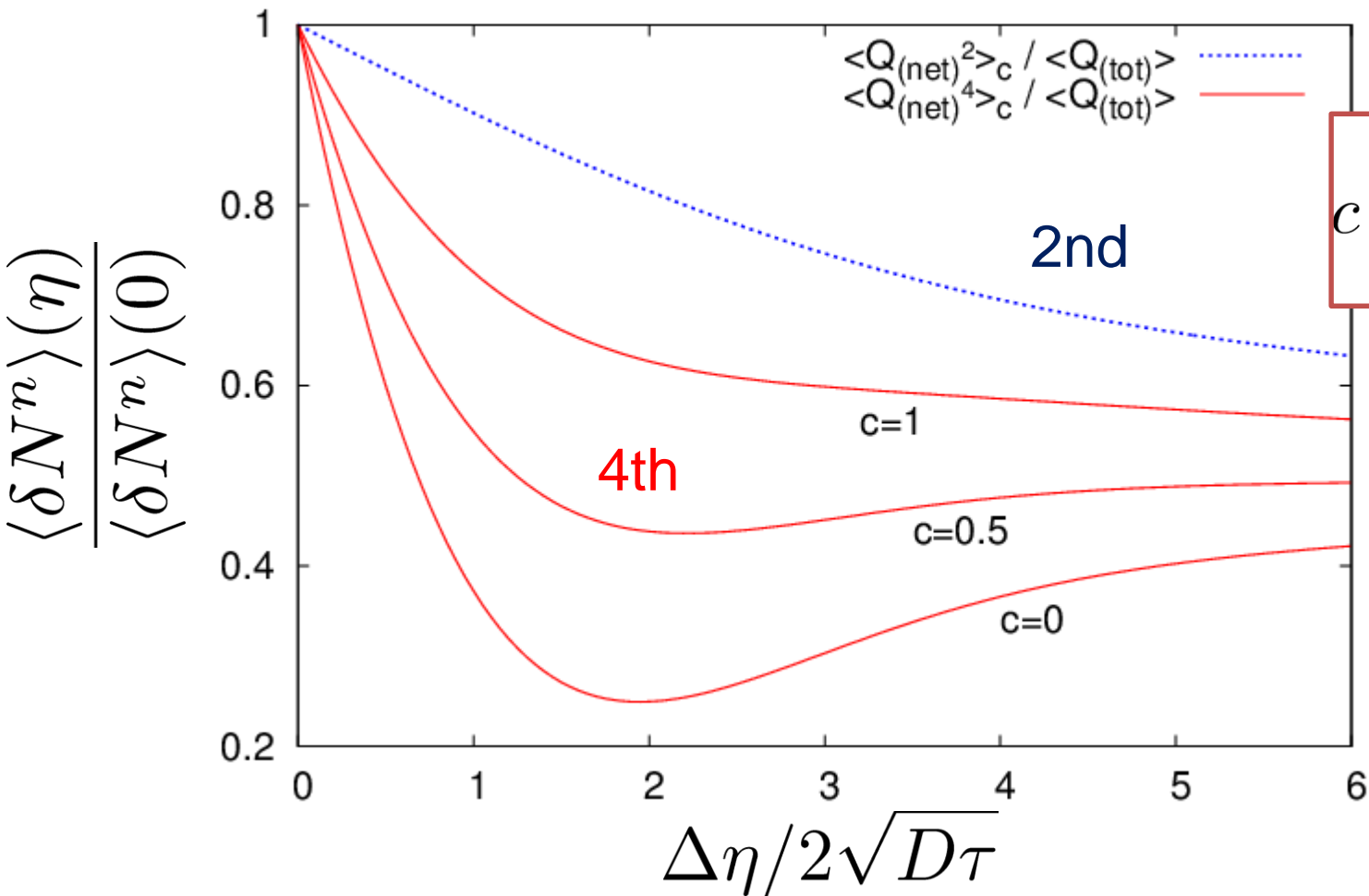


parameter sensitive to hadronization

$\Delta\eta$ Dependence at Freezeout

Initial fluctuations:

$$\langle \bar{Q}^2 \rangle_c = \langle \bar{Q}^4 \rangle_c = \langle \bar{Q}^2 Q_{(\text{tot})} \rangle_c = 0.5 \langle Q_{(\text{tot})} \rangle$$



$$c = \frac{\langle Q_{(\text{tot})}^2 \rangle_c}{\langle Q_{(\text{tot})} \rangle}$$

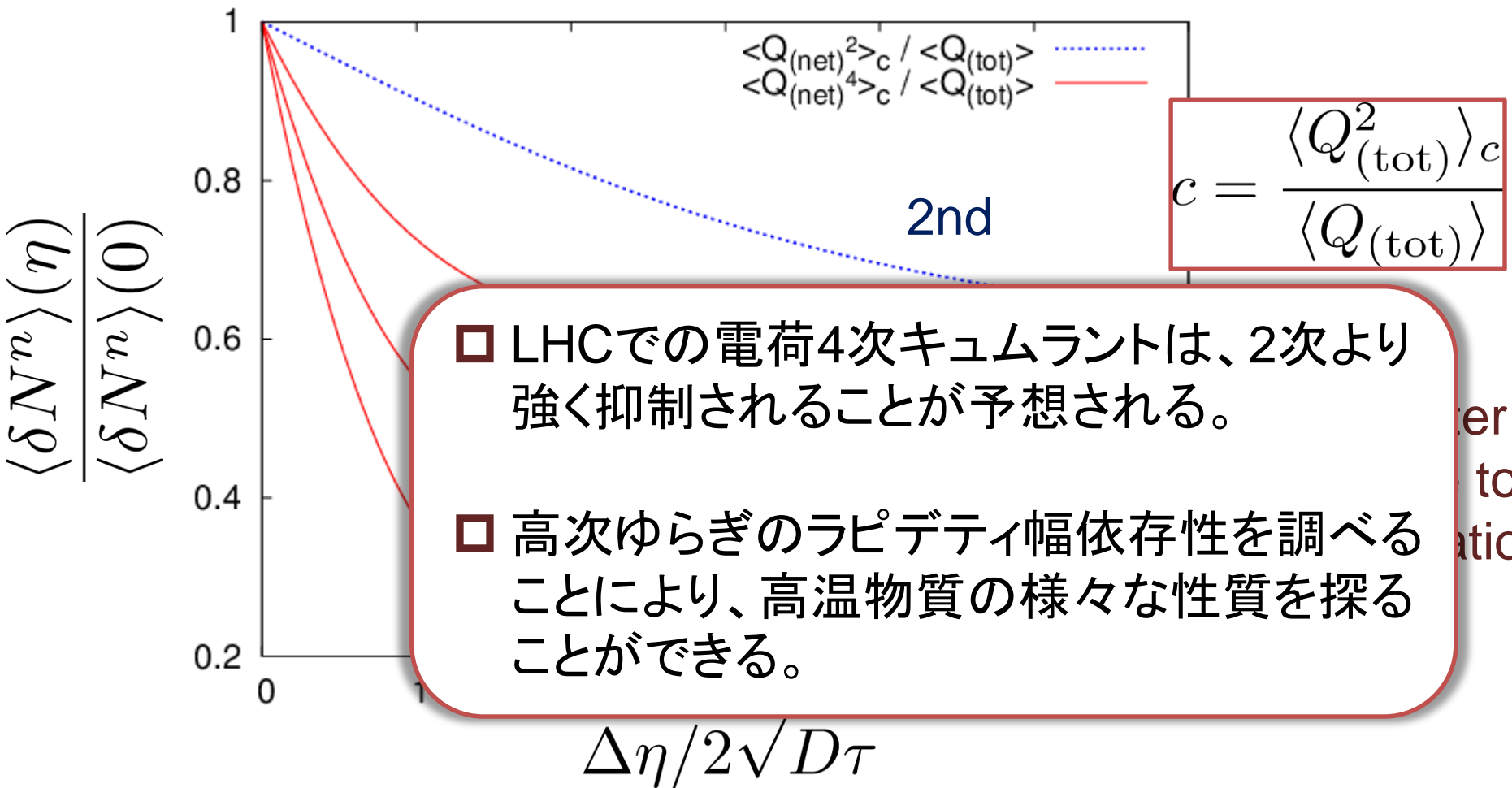


parameter
sensitive to
hadronization

$\Delta\eta$ Dependence at Freezeout

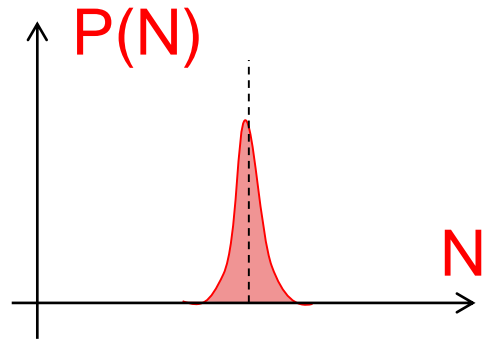
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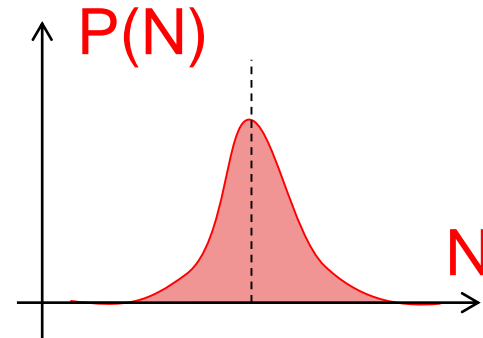


er
to
ation

高次キュムラントが抑制される理由



初期条件: ゆらぎ小



熱平衡: ゆらぎ大

- 高次キュムラントの方が、分布の裾の構造に強く依存
- 分布の裾が熱平衡化するには時間がかかる

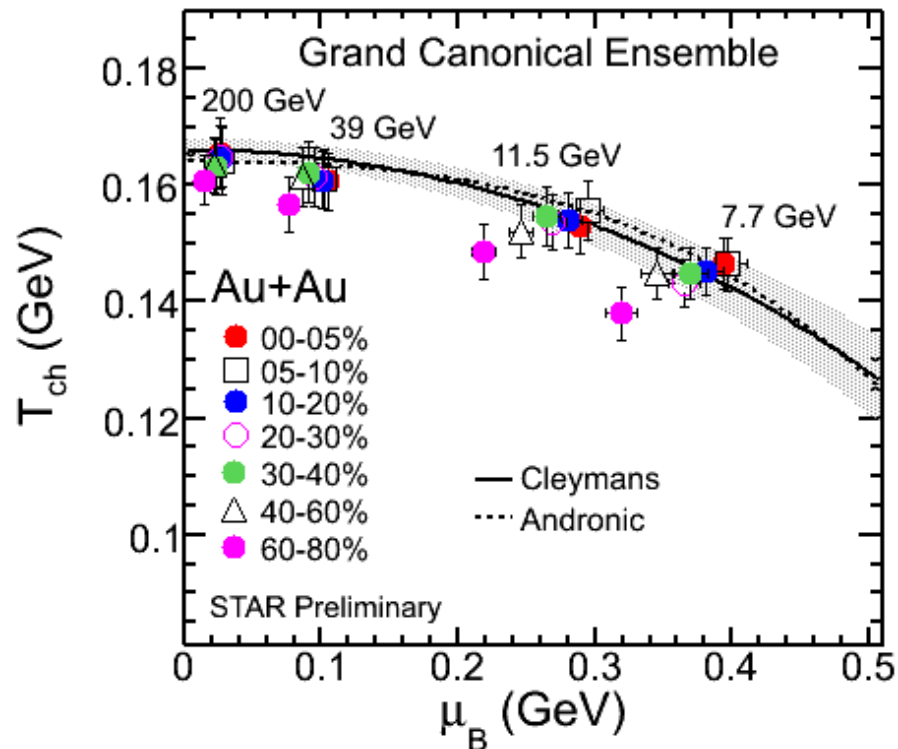
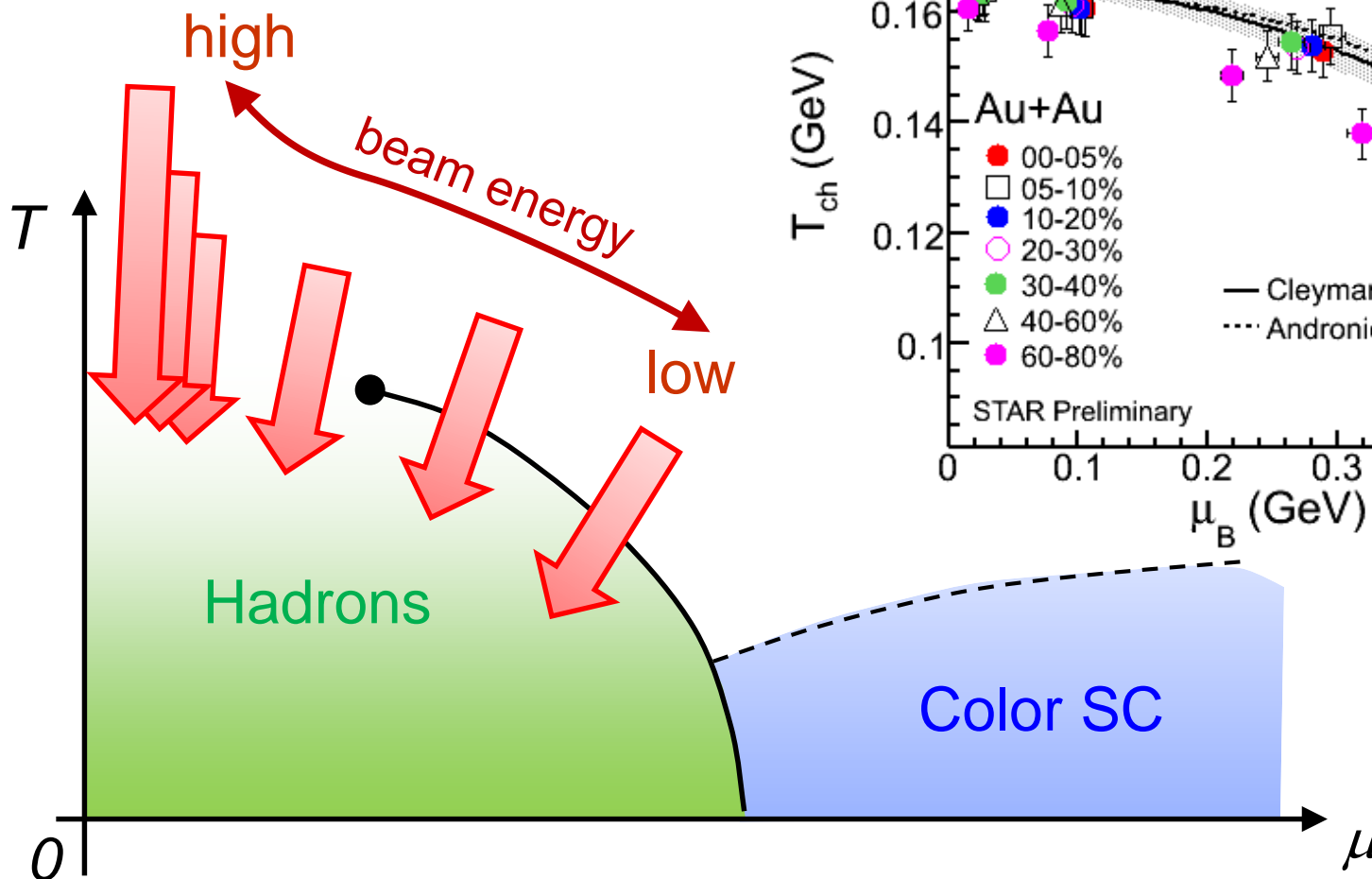
➡ 高次キュムラントは、緩和により長時間を要する。

まとめ

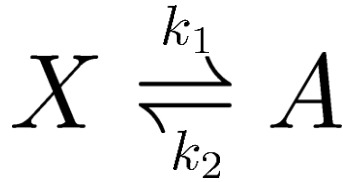
- 保存電荷ゆらぎは、重イオン衝突実験における重要な観測量の一つである。RHICおよびLHCでは4次までの高次キュムラントが解析されている。
- 実験で観測されるゆらぎはハドロン状態の熱平衡値ではなく、QGP状態で生成されたゆらぎが緩和していく過程のある瞬間だと考えられる。
- 本研究では、非ガウスゆらぎの時間発展を調べるため、拡散マスター方程式を解析的に解き、連続極限を取ることで、高次キュムラントのラピディティ幅依存性の振る舞いに対する予言を行った。

Beam-Energy Scan

STAR 2012



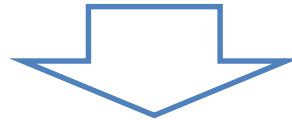
Chemical Reaction 1



x: # of X

a: # of A (**fixed**)

Master eq.:
$$\frac{\partial}{\partial t} P(x, t) = k_2 a P(x - 1, t) + k_1 (x + 1) P(x + 1, t) - (k_1 x + k_2 a) P(x, t)$$



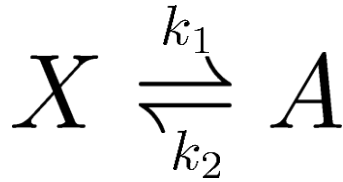
Cumulants with fixed initial condition $P(x, 0) = \delta_{x, N_0}$

$$\langle x(t) \rangle = N_0 e^{-k_1 t} + N_{eq} (1 - e^{-k_1 t})$$

$$\langle \delta x(t)^2 \rangle = N_0 (e^{-k_1 t} - e^{-2k_1 t}) + N_{eq} (1 - e^{-k_1 t})$$

$$\langle \delta x(t)^3 \rangle = \underbrace{N_0 (e^{-k_1 t} - 3e^{-2k_1 t} + 2e^{-3k_1 t})}_{\text{initial}} + \underbrace{N_{eq} (1 - e^{-k_1 t})}_{\text{equilibrium}}$$

Chemical Reaction 2

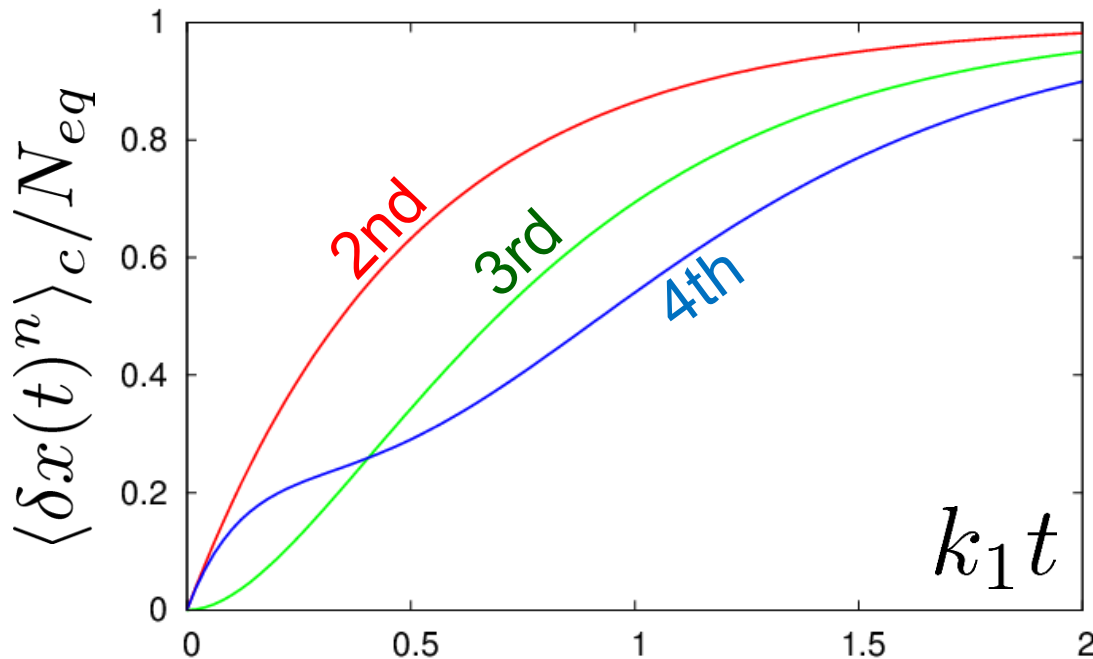


$$N_0 = N_{eq}$$

$$\langle x(t) \rangle = N_{eq}$$

$$\langle \delta x(t)^2 \rangle = N_{eq}(1 - e^{-2k_1 t})$$

$$\langle \delta x(t)^3 \rangle = N_{eq}(1 - 3e^{-2k_1 t} + 2e^{-3k_1 t})$$



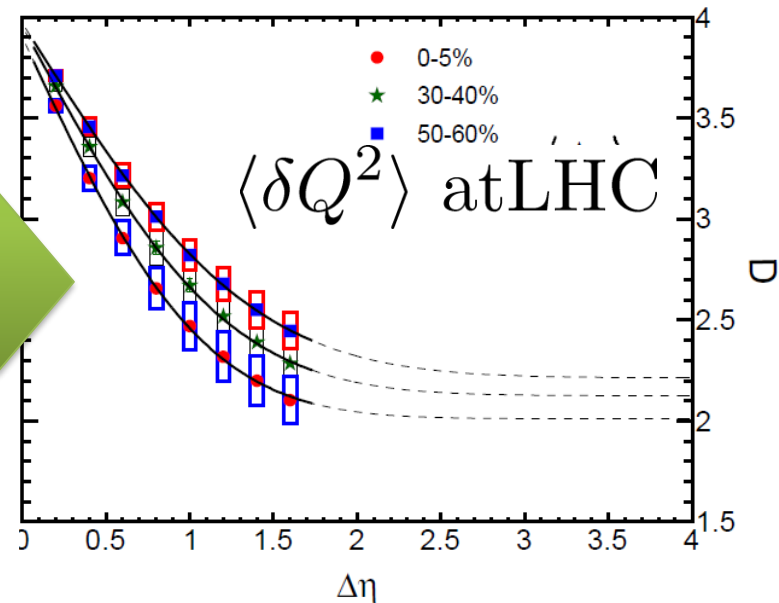
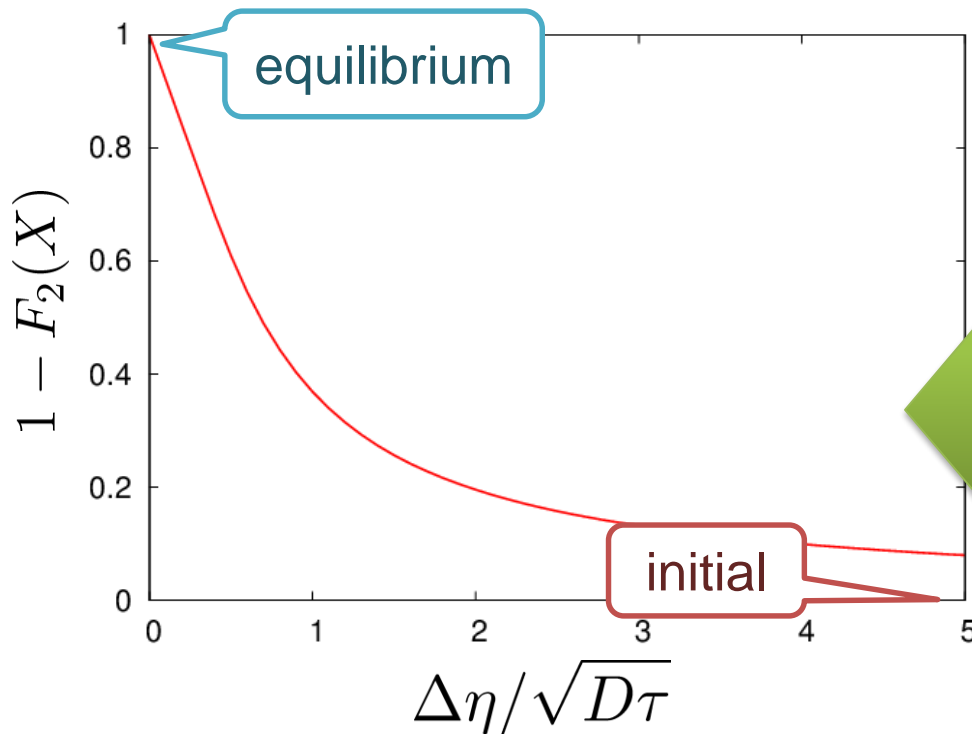
Higher-order
cumulants
grow slower.

$\Delta\eta$ Dependence

Shuryak, Stephanov, 2001

- Initial condition: $\langle \delta n(\eta_1, 0) \delta n(\eta_2, 0) \rangle = \sigma_2 \delta(\eta_1 - \eta_2)$
- Translational invariance

$$Q(\tau) = \int_0^{\Delta\eta} d\eta n(\eta, \tau) \quad \rightarrow \quad \langle \delta Q(\tau)^2 \rangle = \underbrace{\sigma_2 F_2(X)}_{\text{initial}} + \underbrace{\chi_2(1 - F_2(X))}_{\text{equilibrium}}$$

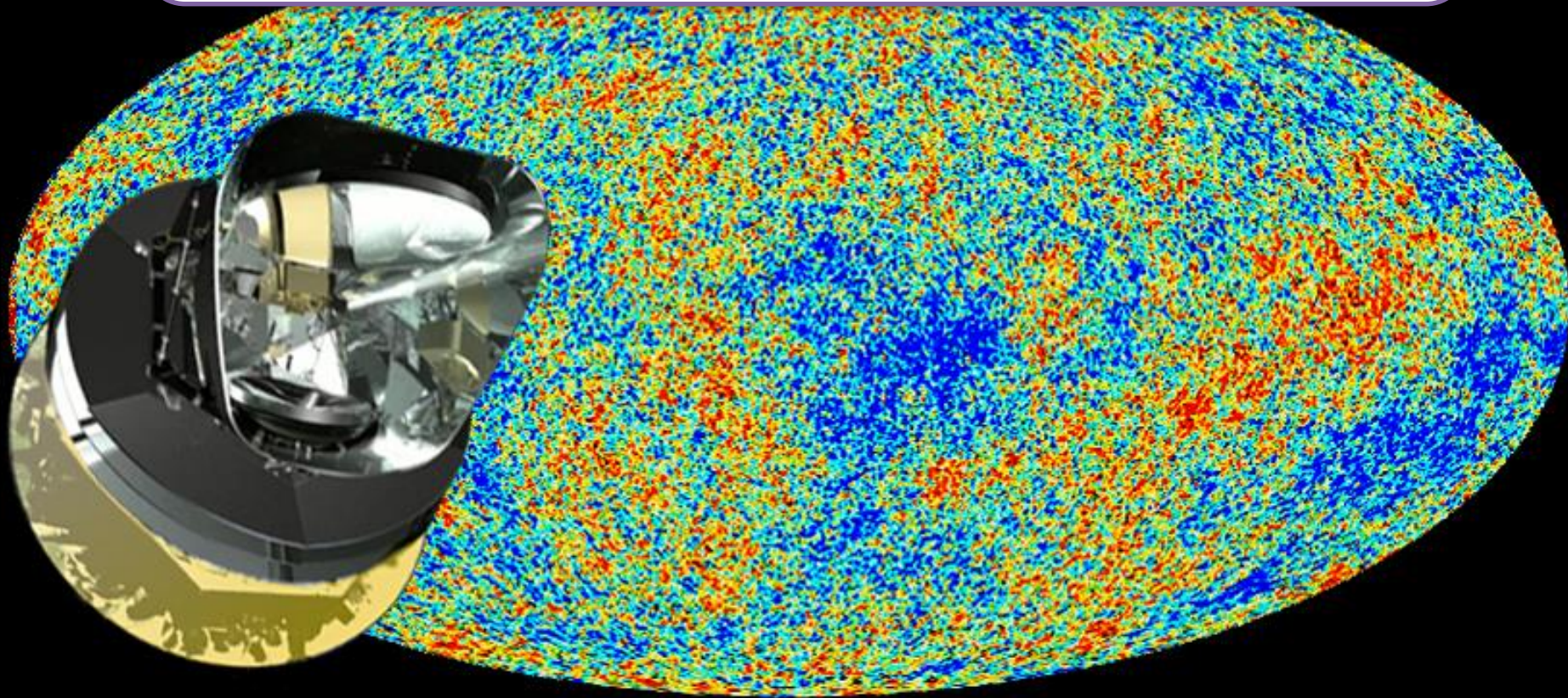


Non-Gaussianity

fluctuations (correlations)

$$\langle \delta n_1 \delta n_2 \rangle, \langle \delta n_1 \delta n_2 \delta n_3 \rangle, \langle \delta n_1 \delta n_2 \delta n_3 \delta n_4 \rangle_c, \dots$$

→ Non-Gaussianity



PLANCK : statistics insufficient to see non-Gaussianity...(2013)

Fluctuations

Fluctuations reflect properties of matter.

Enhancement near the critical point

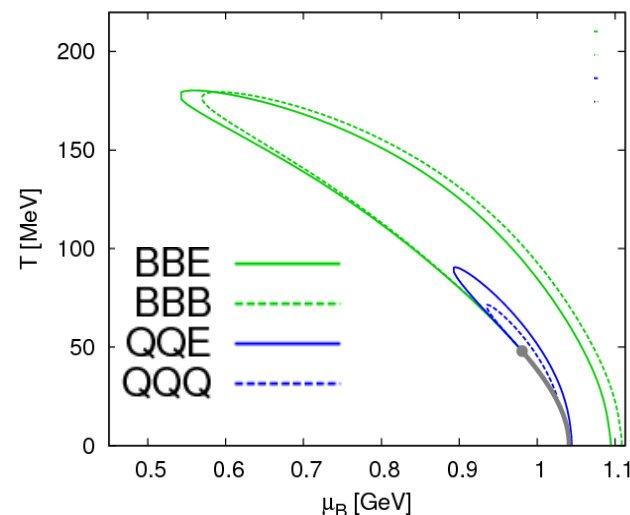
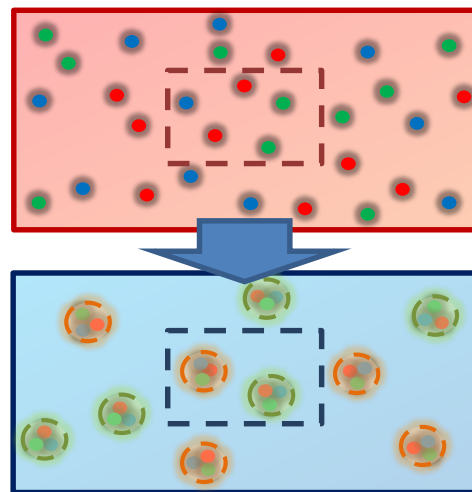
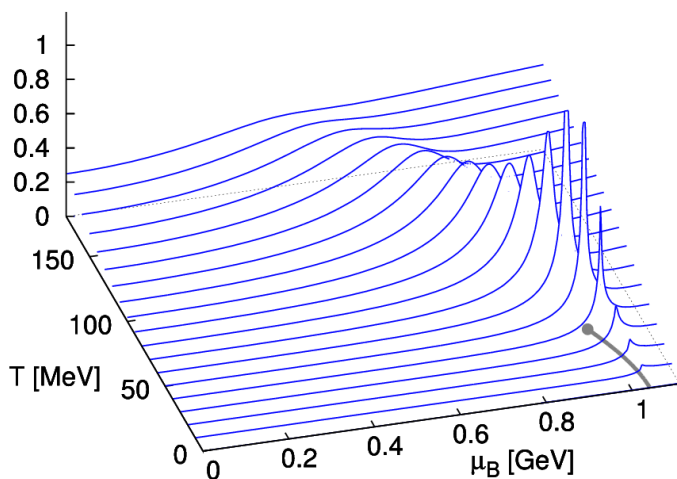
Stephanov,Rajagopal,Shuryak('98); Hatta,Stephanov('02); Stephanov('09);...

Ratios between cumulants of conserved charges

Asakawa,Heinz,Muller('00); Jeon, Koch('00); Ejiri,Karsch,Redlich('06)

Signs of higher order cumulants

Asakawa,Ejiri,MK('09); Friman,et al.('11); Stephanov('11)

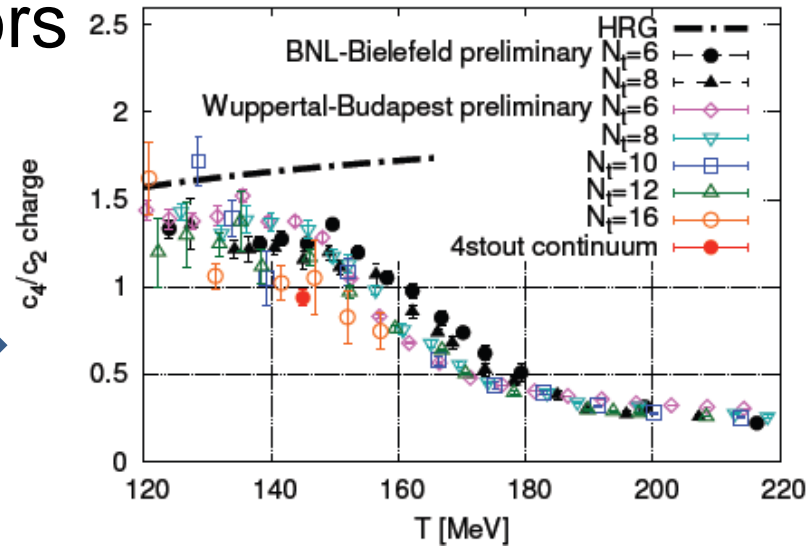


Conserved Charges : Theoretical Advantage

□ Definite definition for operators

- as a Noether current
- calculable on any theory

ex: on the lattice

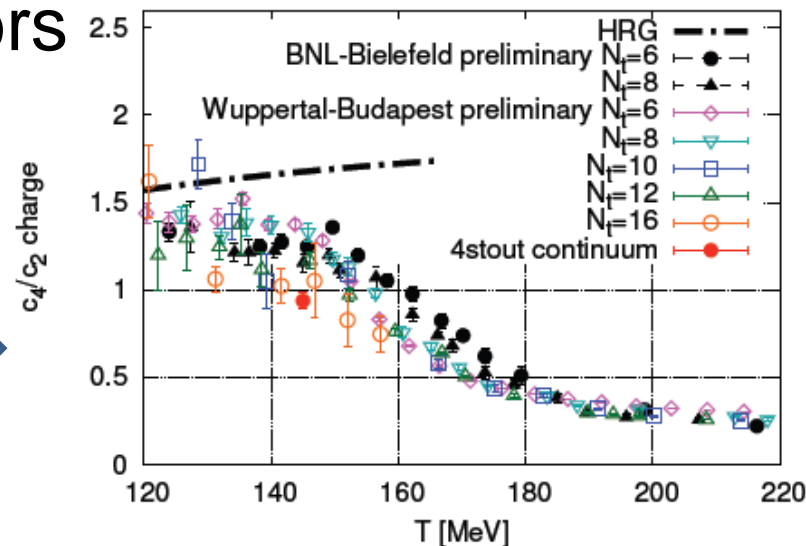


Conserved Charges : Theoretical Advantage

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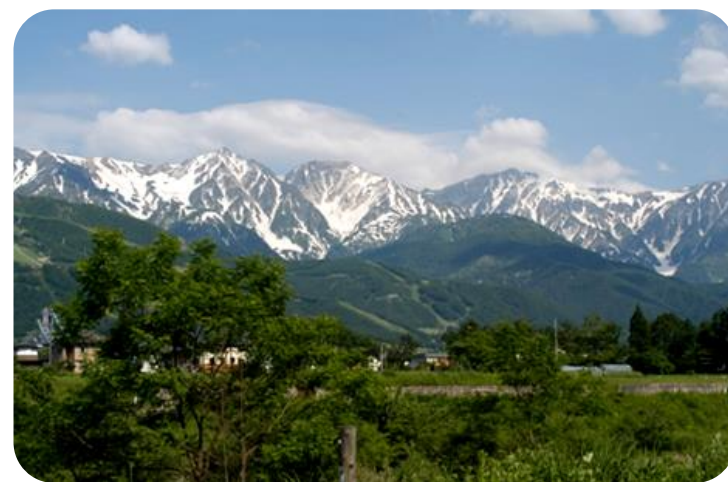


□ Simple thermodynamic relations

$$\langle \delta N_c^n \rangle = \frac{1}{VT^{n-1}} \frac{\partial^n \Omega}{\partial \mu_c^n}$$

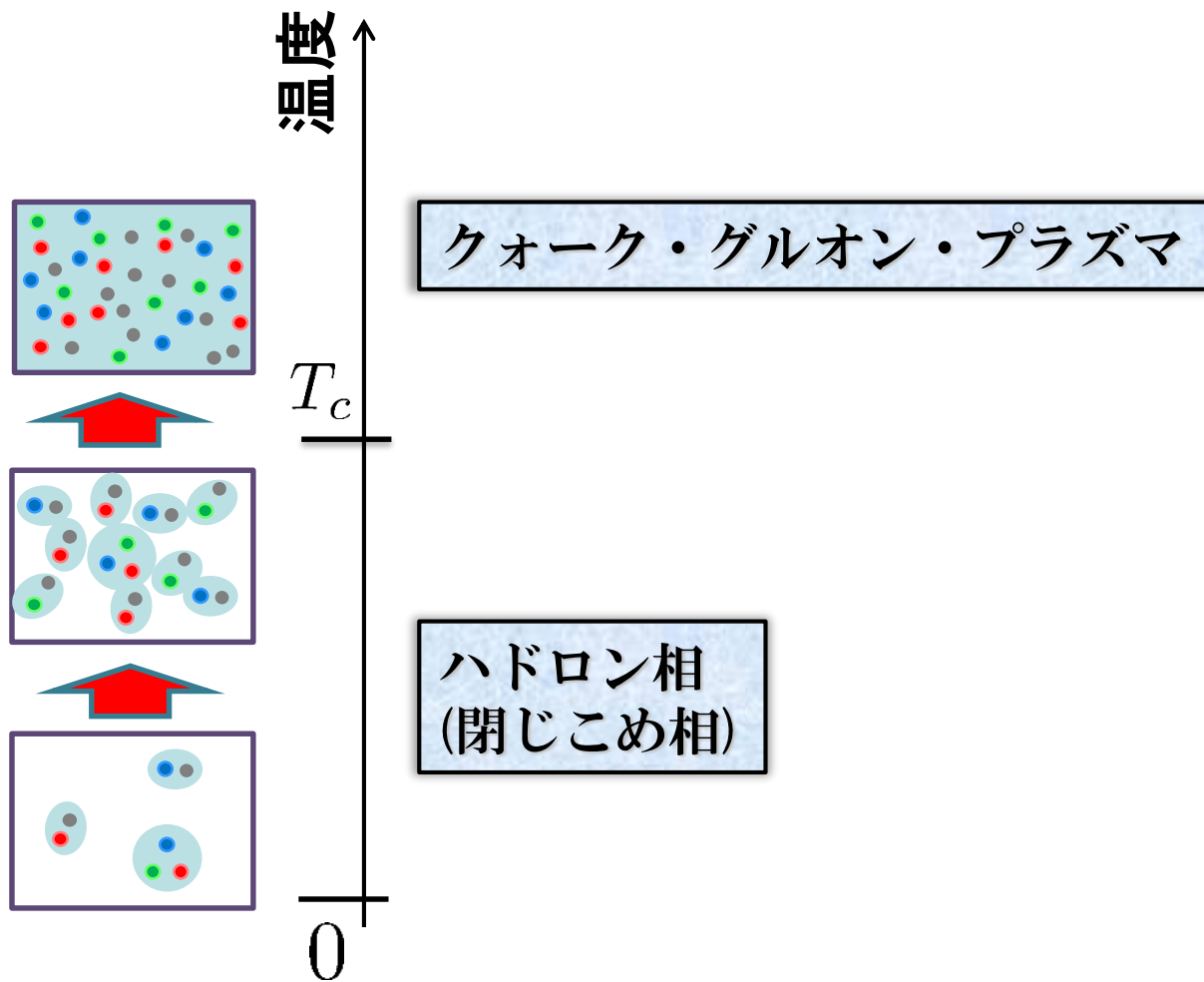
- Intuitive interpretation for the behaviors of cumulants

ex: $\langle \delta N_B^3 \rangle = \frac{1}{VT^2} \frac{\partial \langle \delta N_B^2 \rangle}{\partial \mu_B}$

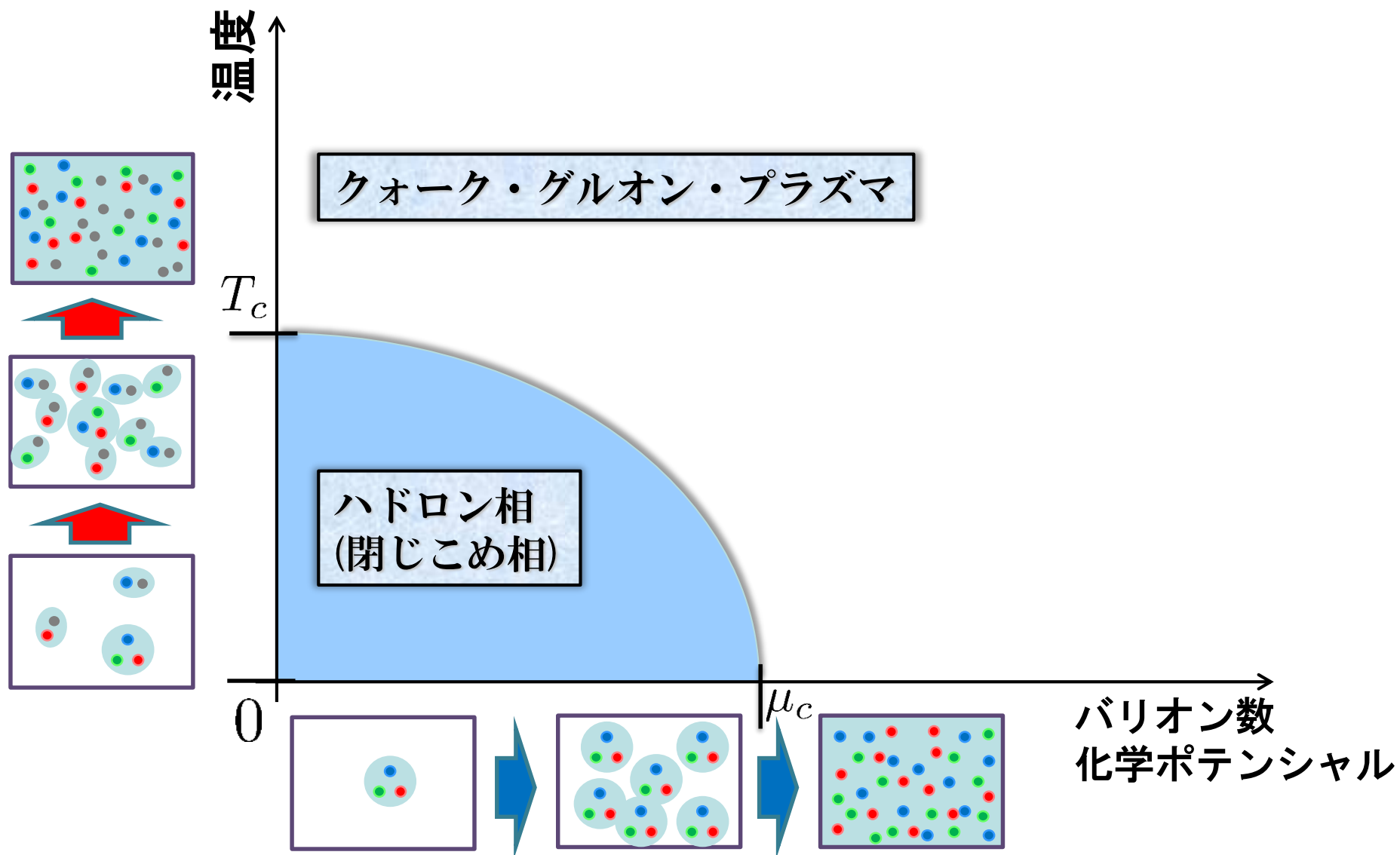


Asakawa, Ejiri, MK, 2009

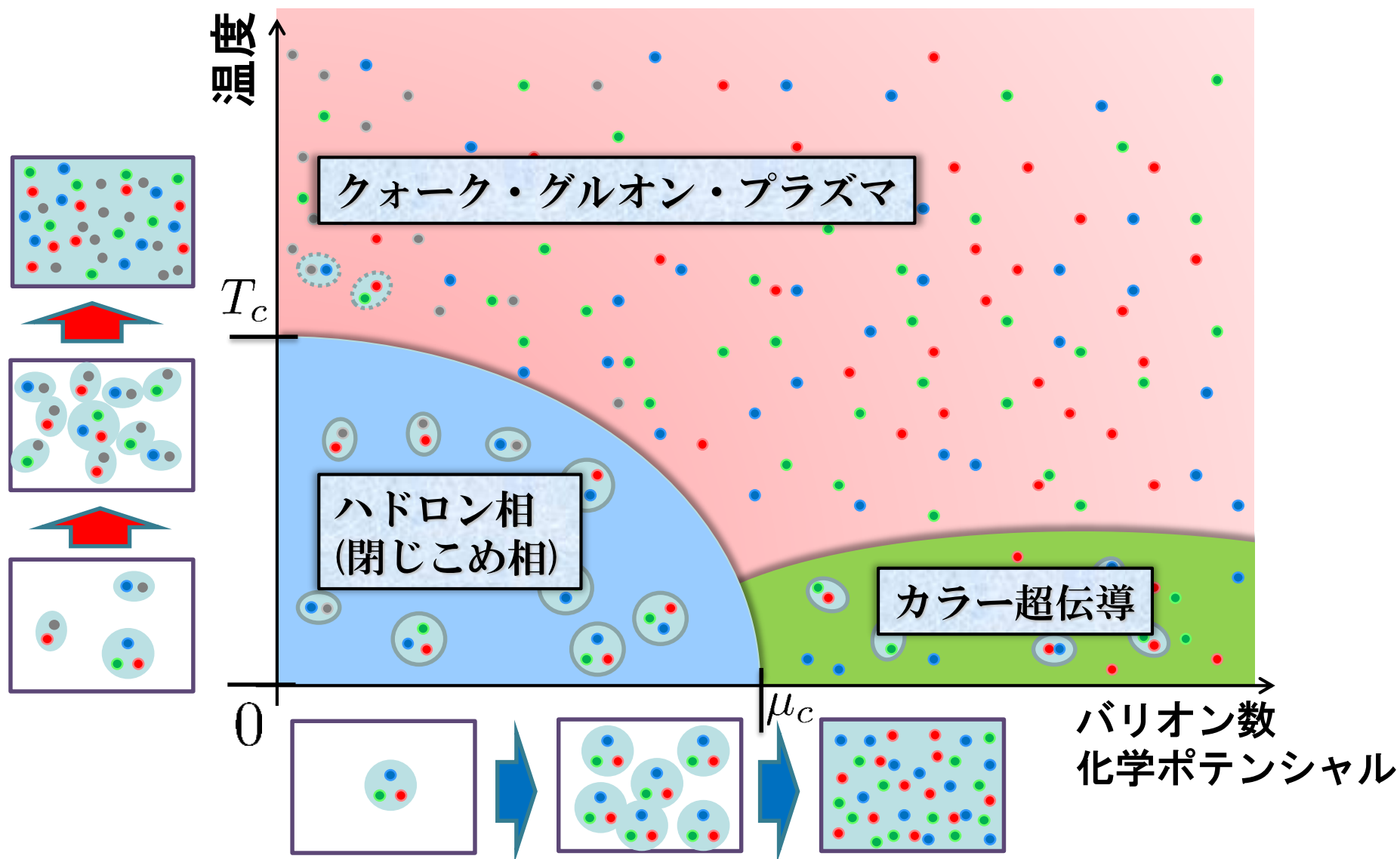
QCDの相転移



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