



Thermodynamics of Gauge Theory from Gradient Flow

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For FlowQCD Collaboration

*T* *μν*

*T*<sub>μν</sub>

$T_{\mu\nu}$

Poincare  
symmetry

$T_{\mu\nu}$

Poincare  
symmetry

$T_{\mu\nu}$

	momentum		
energy	$T_{01}$	$T_{02}$	$T_{03}$
$T_{10}$	$T_{11}$	$T_{12}$	$T_{13}$
$T_{20}$	$T_{21}$	$T_{22}$	$T_{23}$
$T_{30}$	$T_{31}$	$T_{32}$	$T_{33}$
	stress	pressure	

Poincare  
symmetry

$T_{\mu\nu}$

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energy	$T_{01}$	$T_{02}$	$T_{03}$
$T_{10}$	$T_{11}$	$T_{12}$	$T_{13}$
$T_{20}$	$T_{21}$	$T_{22}$	$T_{23}$
$T_{30}$	$T_{31}$	$T_{32}$	$T_{33}$
	stress	pressure	

Einstein Equation

$$G_{\mu\nu} + \Lambda g_{\mu\nu} = \kappa T_{\mu\nu}$$

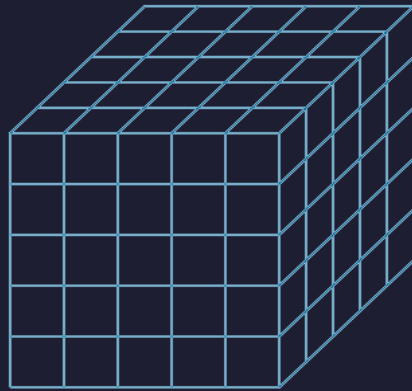
Hydrodynamic Eq.

$$\partial_{\mu} T_{\mu\nu} = 0$$



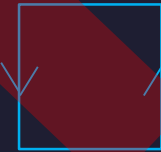
$T_{\mu\nu}$  : nontrivial observable  
on the lattice

- ① Definition of the operator is nontrivial  
because of the explicit breaking of Lorentz symmetry



ex:  $T_{\mu\nu} = F_{\mu\rho}F_{\nu\rho} - \frac{1}{4}\delta_{\mu\nu}FF$

$F_{\mu\nu} =$



- ② Its measurement is extremely noisy  
due to high dimensionality and etc.

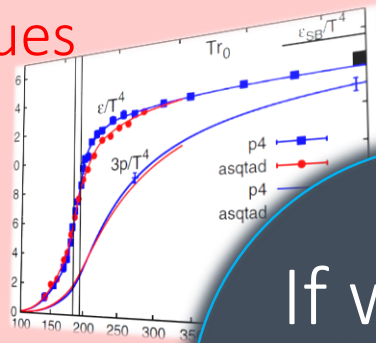
If we have

$$T_{\mu\nu}$$

# Thermodynamics

direct measurement of  
expectation values

$$\langle T_{00} \rangle, \langle T_{ii} \rangle$$



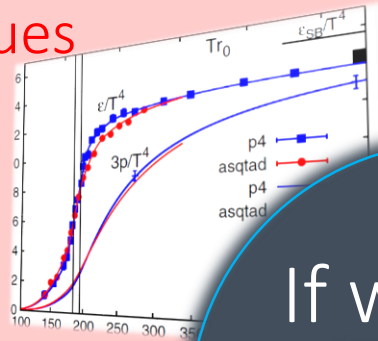
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$$\langle T_{00} \rangle, \langle T_{ii} \rangle$$



# Fluctuations and Correlations

viscosity, specific heat, ...

$$c_V \sim \langle \delta T_{00}^2 \rangle$$

$$\eta = \langle T_{12}; T_{12} \rangle$$

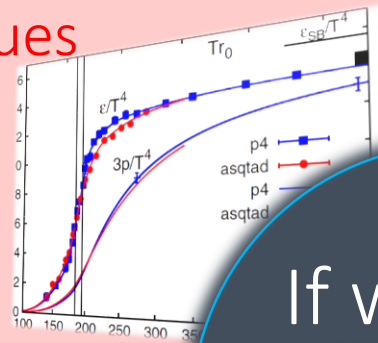
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$$T_{\mu\nu}$$

# Thermodynamics

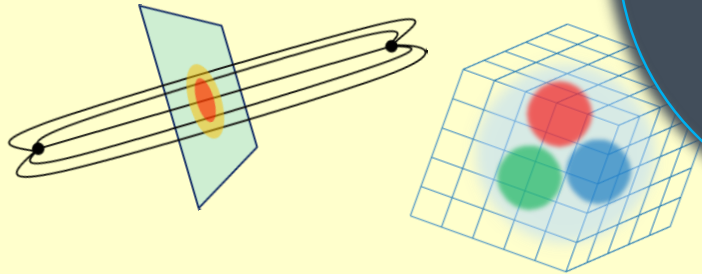
direct measurement of  
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$$\langle T_{00} \rangle, \langle T_{ii} \rangle$$



If we have

$$T_{\mu\nu}$$



- confinement string
- EM distribution in hadrons

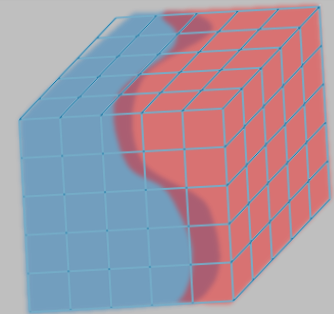
## Hadron Structure

# Fluctuations and Correlations

viscosity, specific heat, ...

$$c_V \sim \langle \delta T_{00}^2 \rangle$$

$$\eta = \langle T_{12}; T_{12} \rangle$$



- vacuum configuration
- mixed state on 1<sup>st</sup> transition

## Vacuum Structure

# Measurement of Thermodynamics using gradient flow

Masakiyo Kitazawa (Osaka U.)

for FlowQCD Collaboration

Asakawa, Hatsuda, Iritani, Itou, MK, Suzuki

FlowQCD, arXiv:1312.7492[hep-lat]; to appear in PRD

# Flow QCD Collaboration



# Flow QCD Collaboration

from

# Gradient

# flow

$$\partial_t B_\mu = D_\nu G_{\nu\mu}$$





# Flow QCD Collaboration

from

**Gradient  
flow**

$$\partial_t B_\mu = D_\nu G_{\nu\mu}$$

to

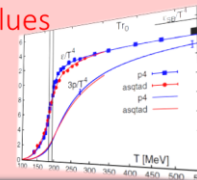
**Hydrodynamic  
flow**

$$\partial_\mu T_{\mu\nu} = 0$$

## Thermodynamics

direct measurement of  
expectation values

$$\langle T_{00} \rangle, \langle T_{ii} \rangle$$



## Fluctuations and Correlations

viscosity, specific heat, ...

$$c_V \sim \langle \delta T_{00}^2 \rangle$$

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# Gradient Flow

# YM Gradient Flow

Luescher, 2010

A. Ramos, plenary, Friday

$$\partial_t A_\mu(t, x) = - \frac{\partial S_{\text{YM}}}{\partial A_\mu}$$

t: "flow time"  
dim:[length<sup>2</sup>]

$$A_\mu(0, x) = A_\mu(x)$$

# YM Gradient Flow

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t: "flow time"  
dim:[length<sup>2</sup>]

- transform gauge field like diffusion equation

$$\partial_t A_\mu = D_\nu G_{\mu\nu} = \partial_\nu \partial_\nu A_\mu + \dots$$

- diffusion length  $d \sim \sqrt{8t}$

- This is **NOT** the standard cooling/smearing

- All composite operators at  $t > 0$  are UV finite Luescher,Weisz,2011

# Applications of Gradient Flow

A. Ramos, plenary, Friday

- ① scale setting
- ② running coupling
- ③ topology
- ④ operator relation
- ⑤ autocorrelation, etc.

# Small Flow Time Expansion of Operators and EMT

# Operator Relation

Luescher, Weisz, 2011

$$\tilde{\mathcal{O}}(t, x) \xrightarrow[t \rightarrow 0]{} \sum_i c_i(t) \mathcal{O}_i^R(x)$$

an operator at  $t > 0$

remormalized operators  
of original theory

# Operator Relations

$t$

Luescher, Weisz, 2011

  
 $\tilde{\mathcal{O}}(t, x)$

$$\tilde{\mathcal{O}}(t, x) \xrightarrow{t \rightarrow 0} \sum_i c_i(t) \mathcal{O}_i^R(x)$$

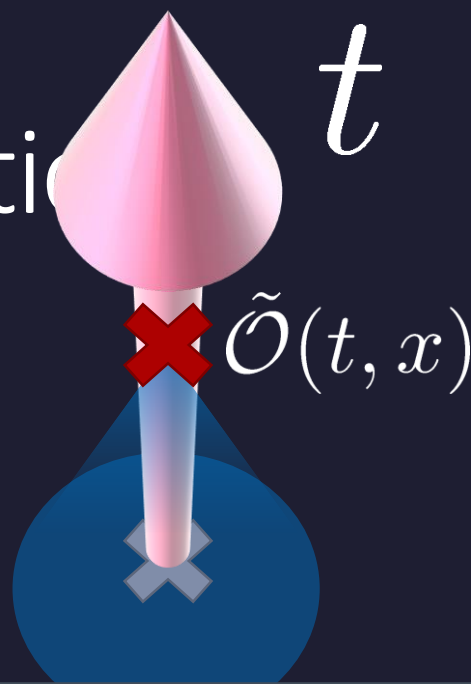
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# Operator Relations

Luescher, Weisz, 2011



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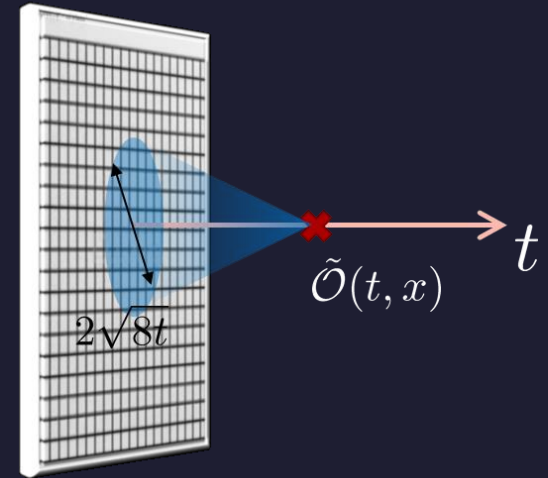
renormalized operators  
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# Constructing EMT

Suzuki, 2013

DelDebbio, Patella, Rago, 2013

$$\tilde{\mathcal{O}}(t, x) \xrightarrow{t \rightarrow 0} \sum_i c_i(t) \mathcal{O}_i^R(x)$$



## □ gauge-invariant dimension 4 operators

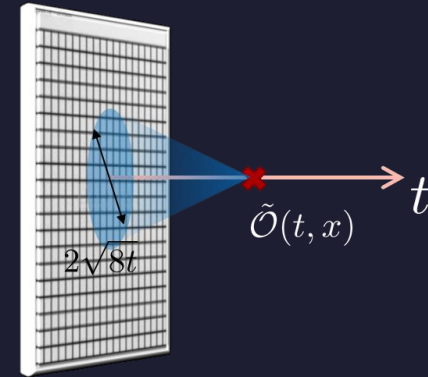
$$\left\{ \begin{array}{l} U_{\mu\nu}(t, x) = G_{\mu\rho}(t, x)G_{\nu\rho}(t, x) - \frac{1}{4}\delta_{\mu\nu}G_{\mu\nu}(t, x)G_{\mu\nu}(t, x) \\ E(t, x) = \frac{1}{4}\delta_{\mu\nu}G_{\mu\nu}(t, x)G_{\mu\nu}(t, x) \end{array} \right.$$

# Constructing EMT 2

Suzuki, 2013

$$U_{\mu\nu}(t, x) = \alpha_U(t) \left[ T_{\mu\nu}^R(x) - \frac{1}{4} \delta_{\mu\nu} T_{\rho\rho}^R(x) \right] + \mathcal{O}(t)$$

$$E(t, x) = \langle E(t, x) \rangle + \alpha_E(t) T_{\rho\rho}^R(x) + \mathcal{O}(t)$$



Suzuki coeffs.  $\left\{ \begin{array}{l} \alpha_U(t) = g^2 [1 + 2b_0 s_1 g^2 + O(g^4)] \\ \alpha_E(t) = \frac{1}{2b_0} [1 + 2b_0 s_2 g^2 + O(g^4)] \end{array} \right.$

$$\begin{aligned} g &= g(1/\sqrt{8t}) \\ s_1 &= 0.03296\dots \\ s_2 &= 0.19783\dots \end{aligned}$$

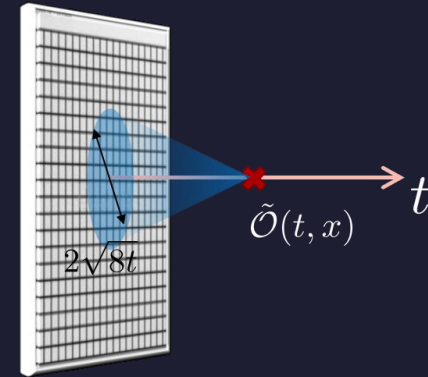
See also, Patella, Parallel7E, Thu.

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$$s_1 = 0.03296 \dots$$

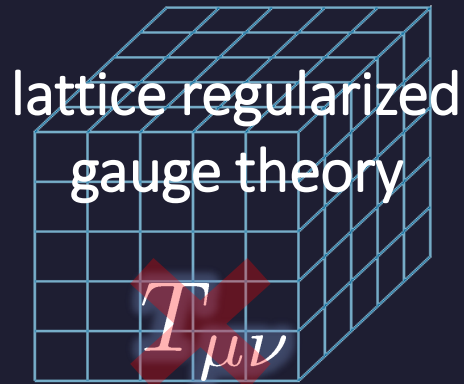
$$s_2 = 0.19783 \dots$$

## Remormalized EMT

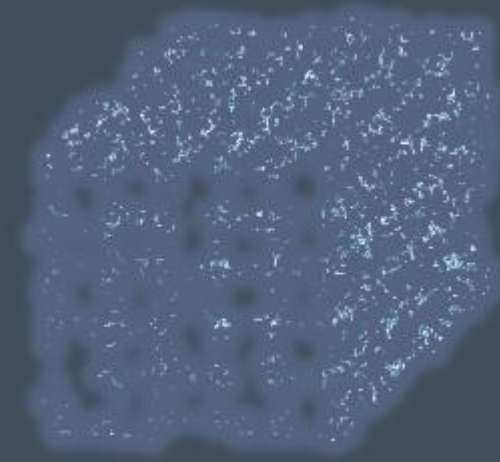
$$T_{\mu\nu}^R(x) = \lim_{t \rightarrow 0} \left[ \frac{1}{\alpha_U(t)} U_{\mu\nu}(t, x) + \frac{\delta_{\mu\nu}}{4\alpha_E(t)} E(t, x)_{\text{subt.}} \right]$$

# Numerical Analysis on the Lattice

# Gradient Flow Method



gradient flow



$$T_{\mu\nu}^R$$

continuum theory  
(with dim. reg.)

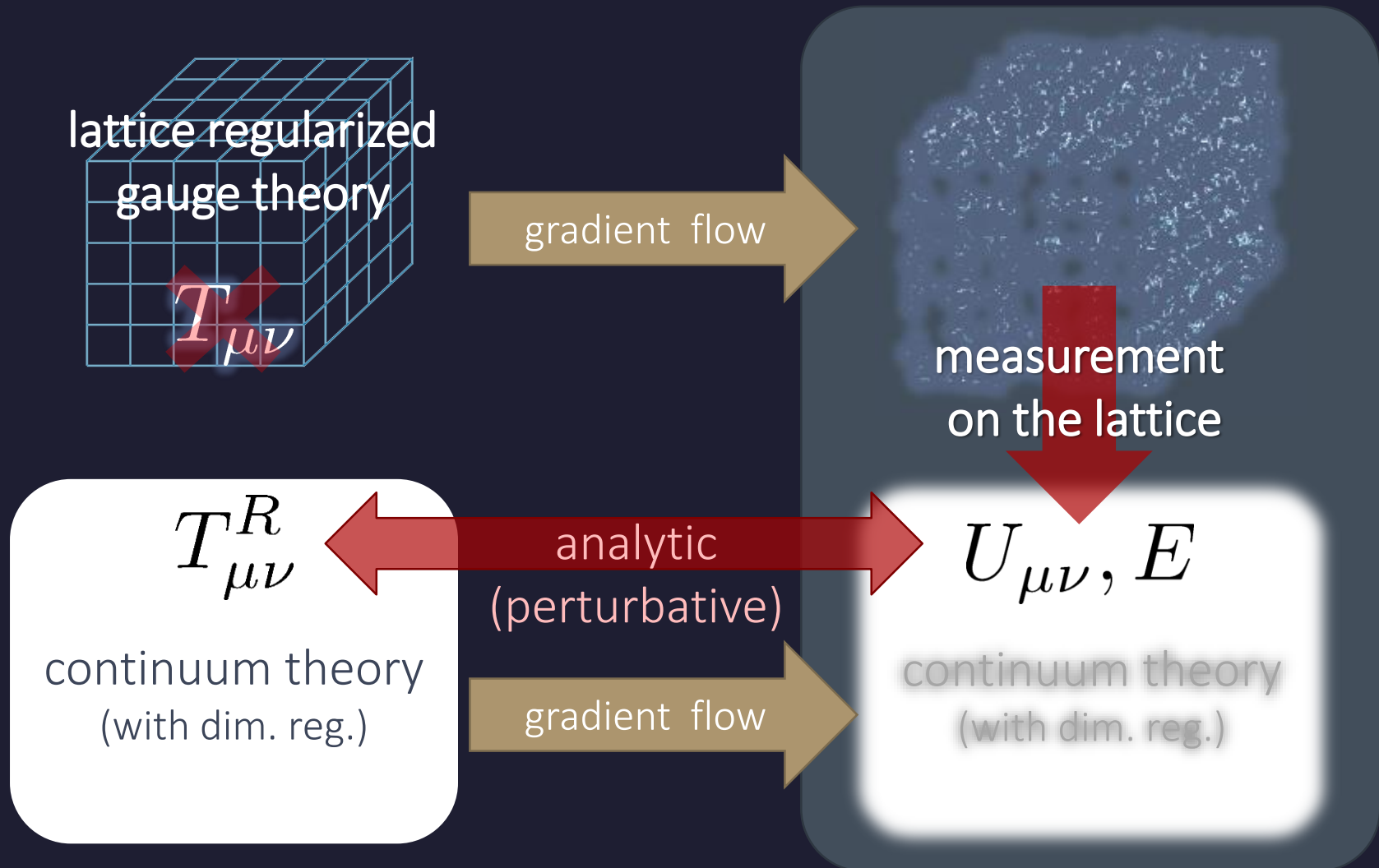
analytic  
(perturbative)

$$U_{\mu\nu}, E$$

continuum theory  
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gradient flow

# Gradient Flow Method

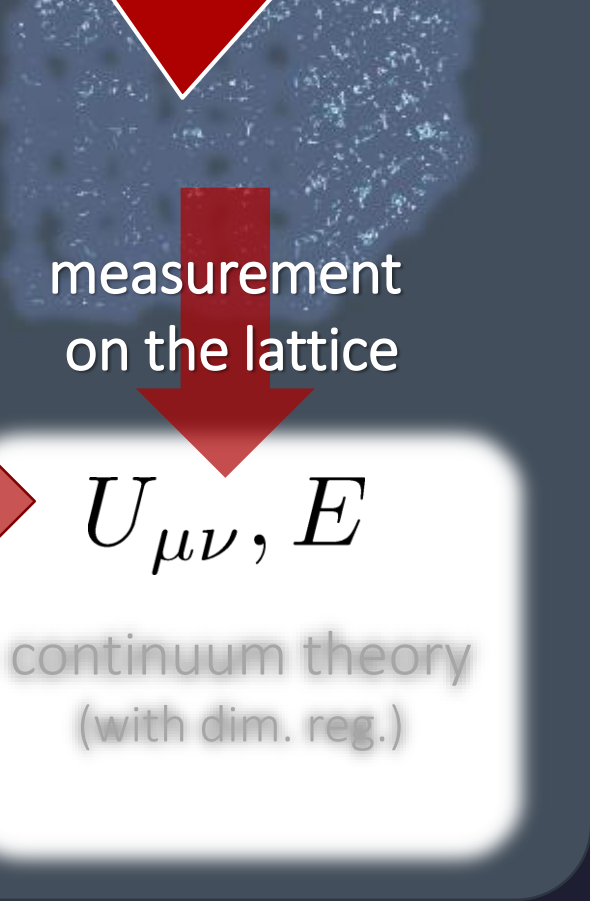


# Caveats



Perturbative relation has to be applicable!  
 $\sqrt{8t} \ll \Lambda^{-1}, T^{-1}$

Gauge field has to be sufficiently smeared!  
 $a \ll \sqrt{8t}$



measurement on the lattice

$T^R_{\mu\nu}$   
continuum theory (with dim. reg.)

analytic (perturbative)

$U_{\mu\nu}, E$

continuum theory (with dim. reg.)

gradient flow



# Caveats



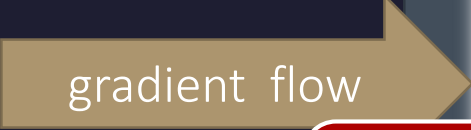
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$T R_{\mu\nu}$   
 continuum theory  
 (with dim. reg.)

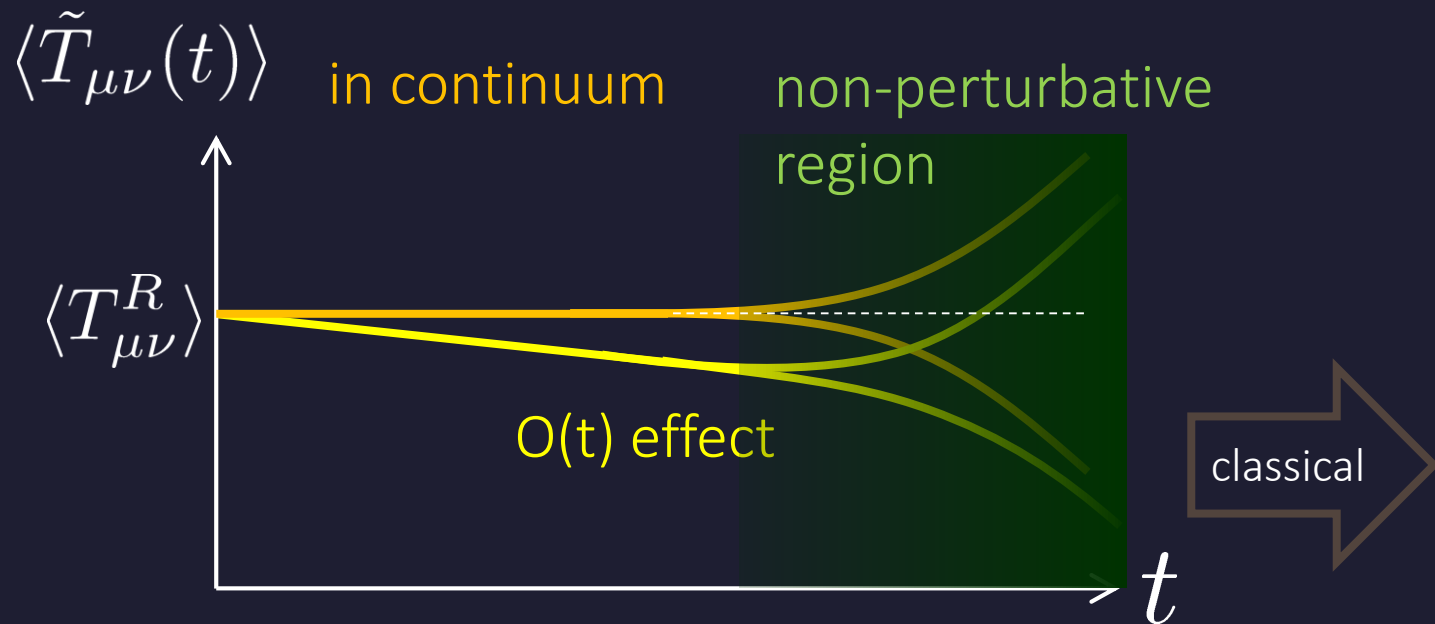
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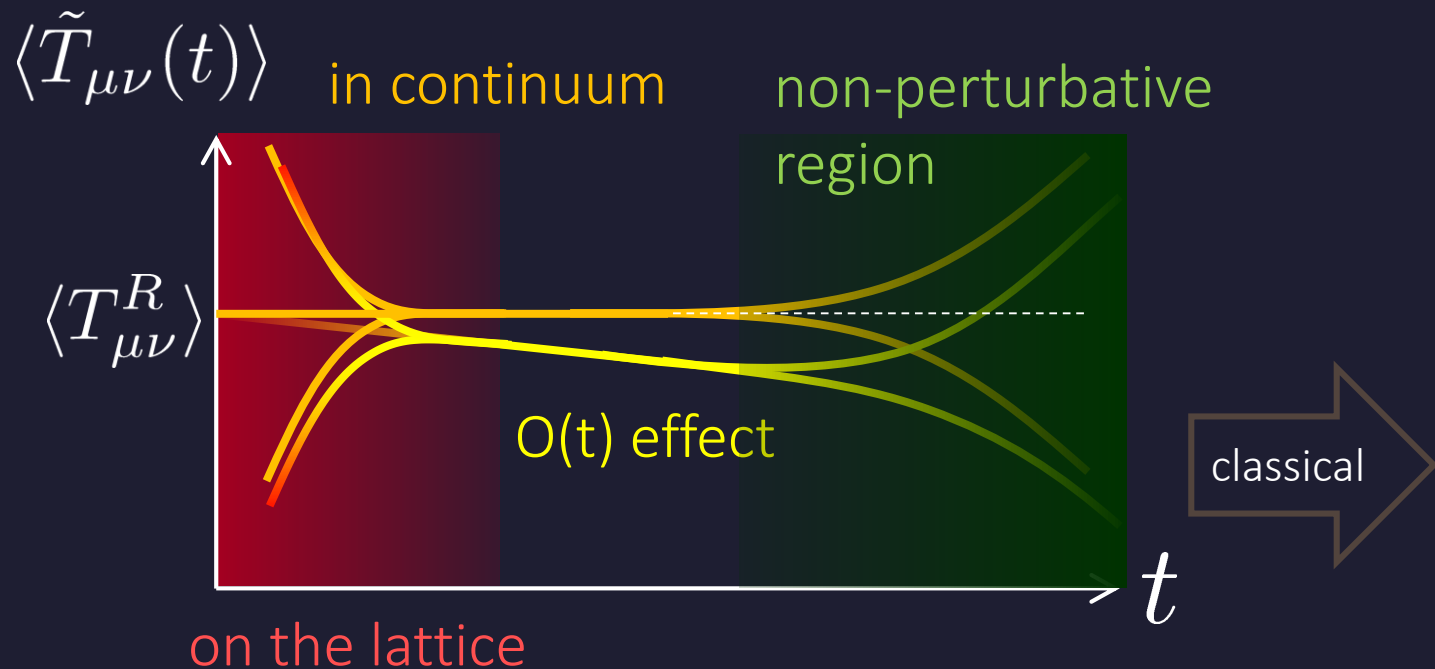
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$a \ll \sqrt{8t} \ll \Lambda^{-1}$

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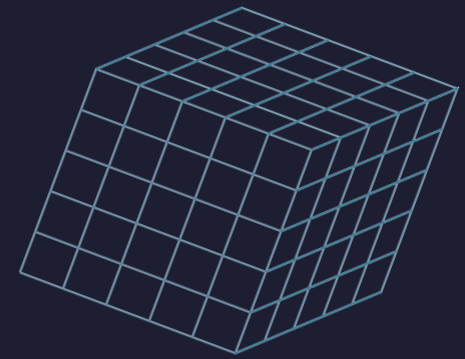
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□  $t \rightarrow 0$  limit with keeping  $t \gg a^2$

# Numerical Simulation

- SU(3) YM theory
- Wilson gauge action



*twice finer lattice!*

## Simulation 1

(arXiv:1312.7492)

- lattice size:  $32^3 \times N_t$
- $N_t = 6, 8, 10$
- $\beta = 5.89 - 6.56$
- $\sim 300$  configurations



## Simulation 2

(*new*, preliminary)

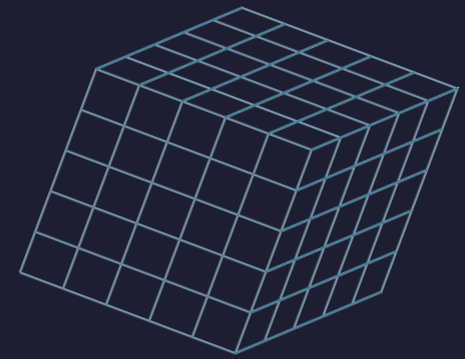
- lattice size:  $64^3 \times N_t$
- $N_t = 10, 12, 14, 16$
- $\beta = 6.4 - 7.4$
- $\sim 2000$  configurations

using SX8 @ RCNP  
SR16000 @ KEK

using BlueGeneQ @ KEK  
efficiency  $\sim 40\%$

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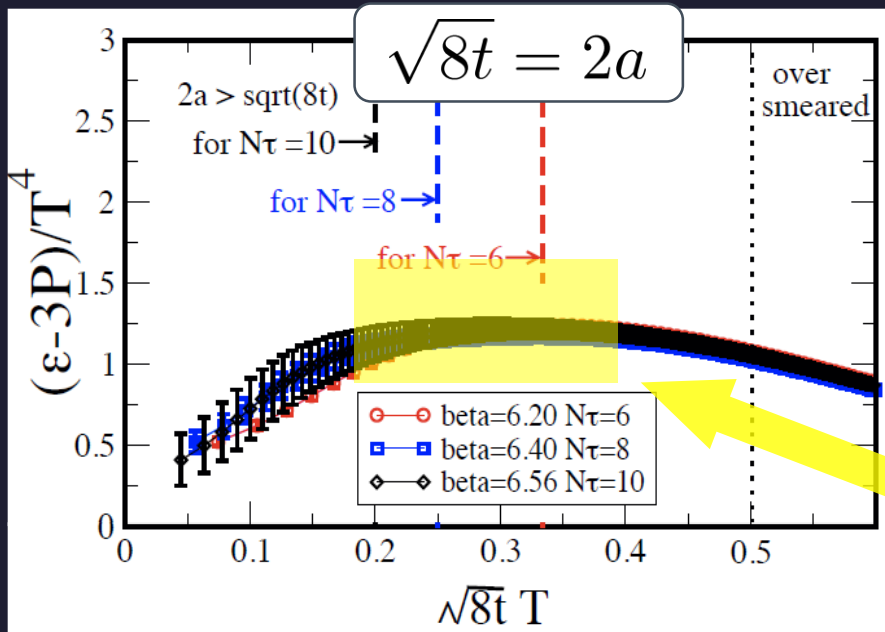
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# $\varepsilon$ -3p at $T=1.65T_c$

$$\tilde{T}_{\mu\nu}(t) = \frac{1}{\alpha_U(t)} U_{\mu\nu}(t) + \frac{\delta_{\mu\nu}}{4\alpha_E(t)} E(t)_{\text{subt.}}$$

$$T_{\mu\nu}^R = \lim_{t \rightarrow 0} \tilde{T}_{\mu\nu}(t)$$



Emergent plateau!

$$2a \lesssim \sqrt{8t} \lesssim 0.4T^{-1}$$

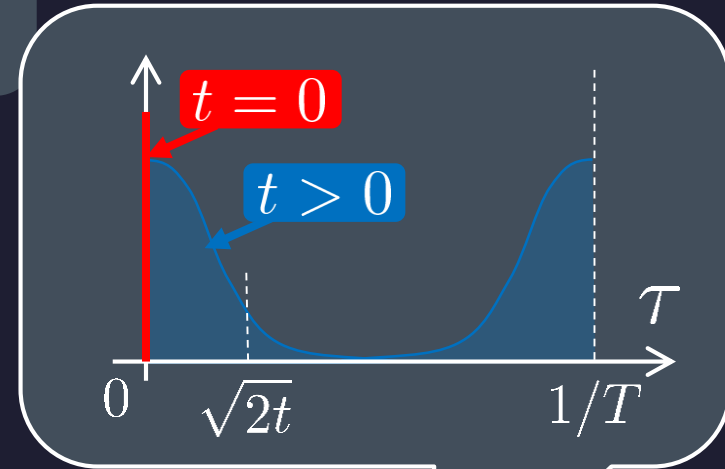
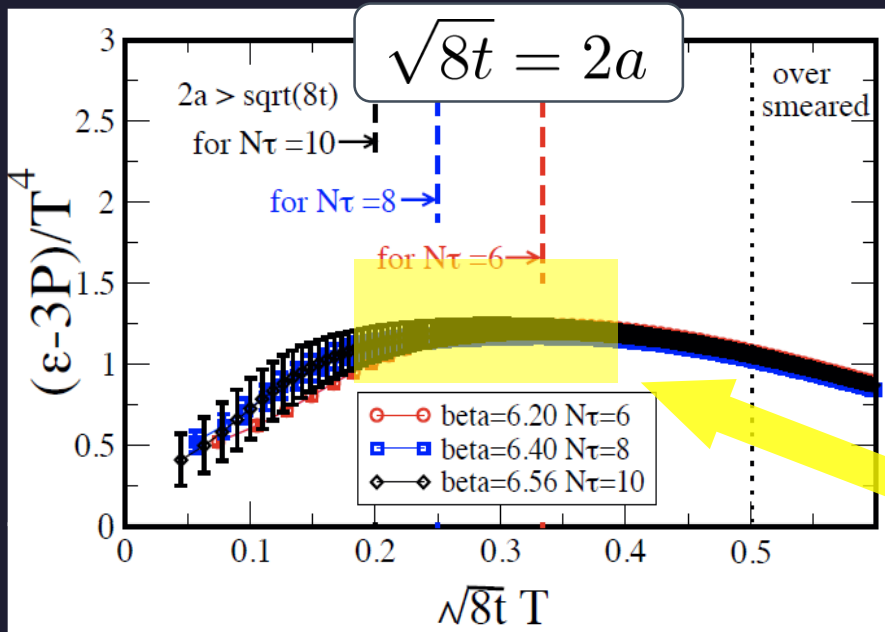
$Nt=6,8,10$   
 $\sim 300$  confs.

the range of  $t$  where the EMT formula is successfully used!

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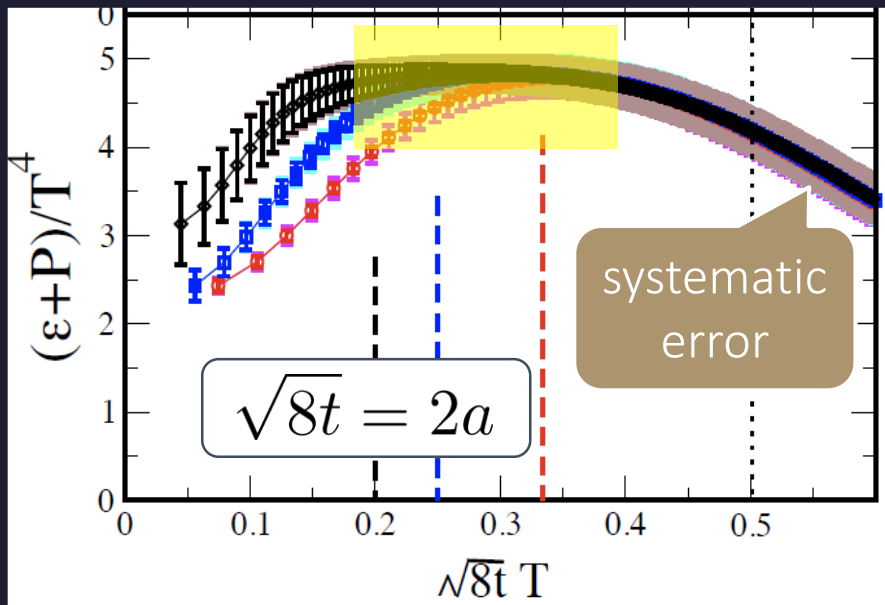
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 $\sim 300$  confs.

the range of  $t$  where the EMT formula is successfully used!

# Entropy Density at $T=1.65T_c$



Emergent plateau!

$$2a \lesssim \sqrt{8t} \lesssim 0.4T^{-1}$$

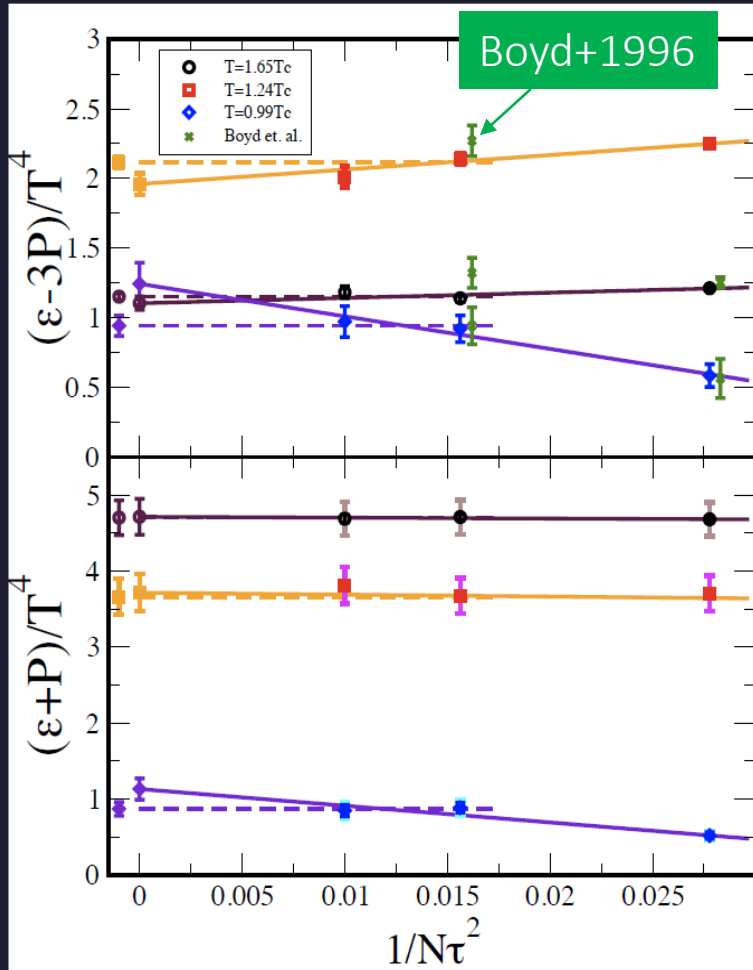
$N_t=6,8,10$   
 $\sim 300$  confs.

Direct measurement of  $e+p$  on a given  $T$ !

**NO** integral / **NO** vacuum subtraction



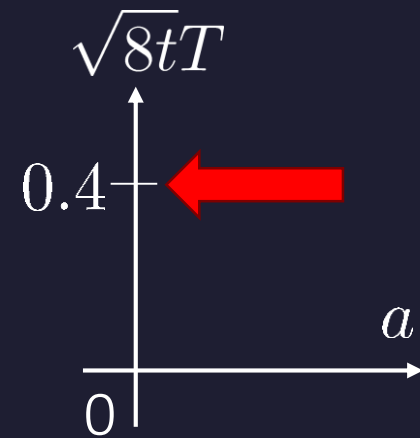
# Continuum Limit



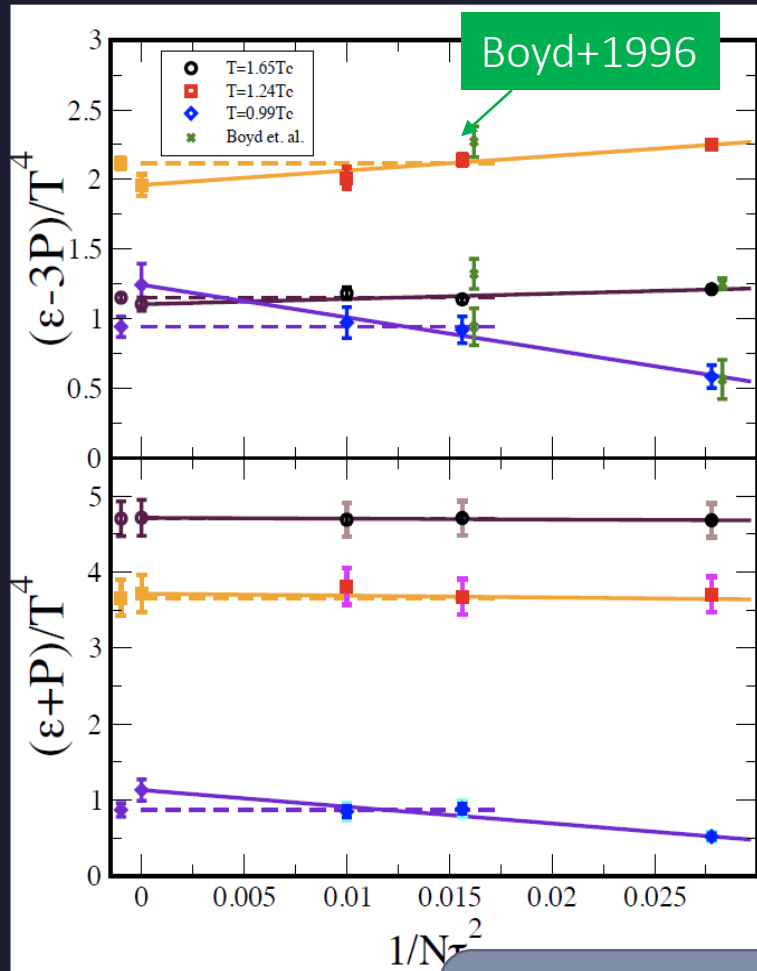
$32^3 \times N_t$

$N_t = 6, 8, 10$

$T/T_c = 0.99, 1.24, 1.65$



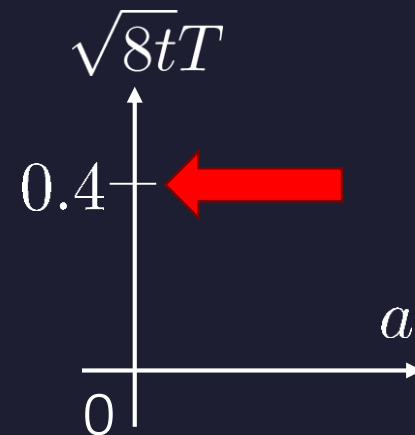
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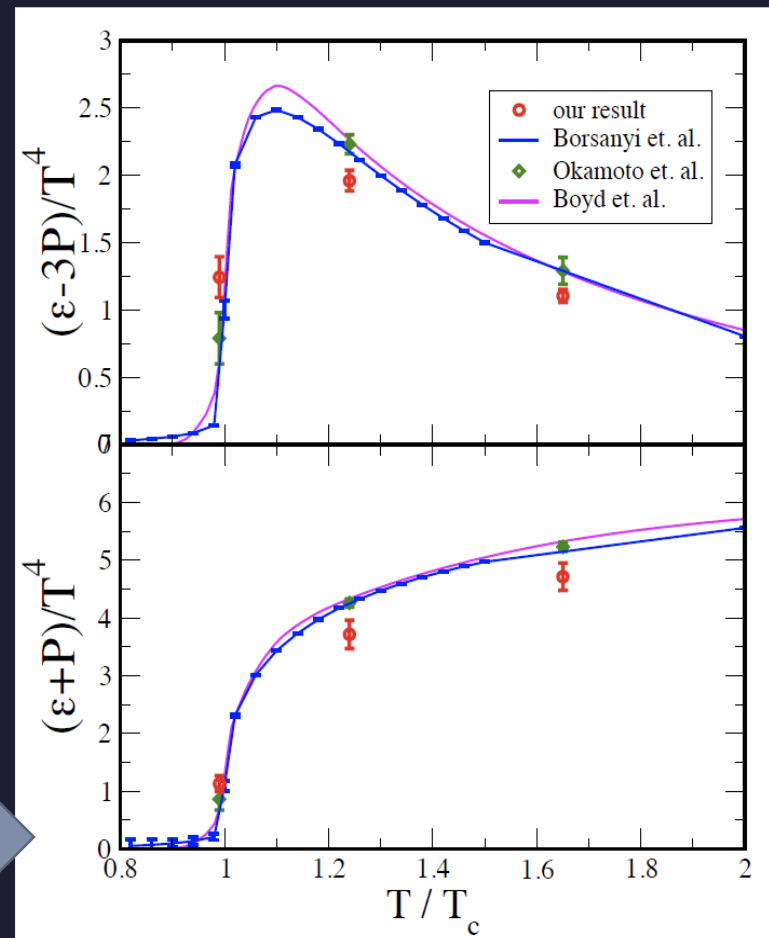
$32^3xNt$

$Nt = 6, 8, 10$

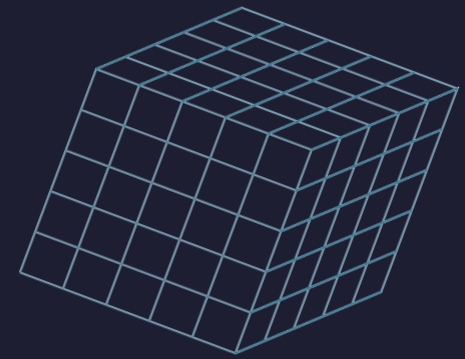
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Comparison with previous studies



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## Simulation 2

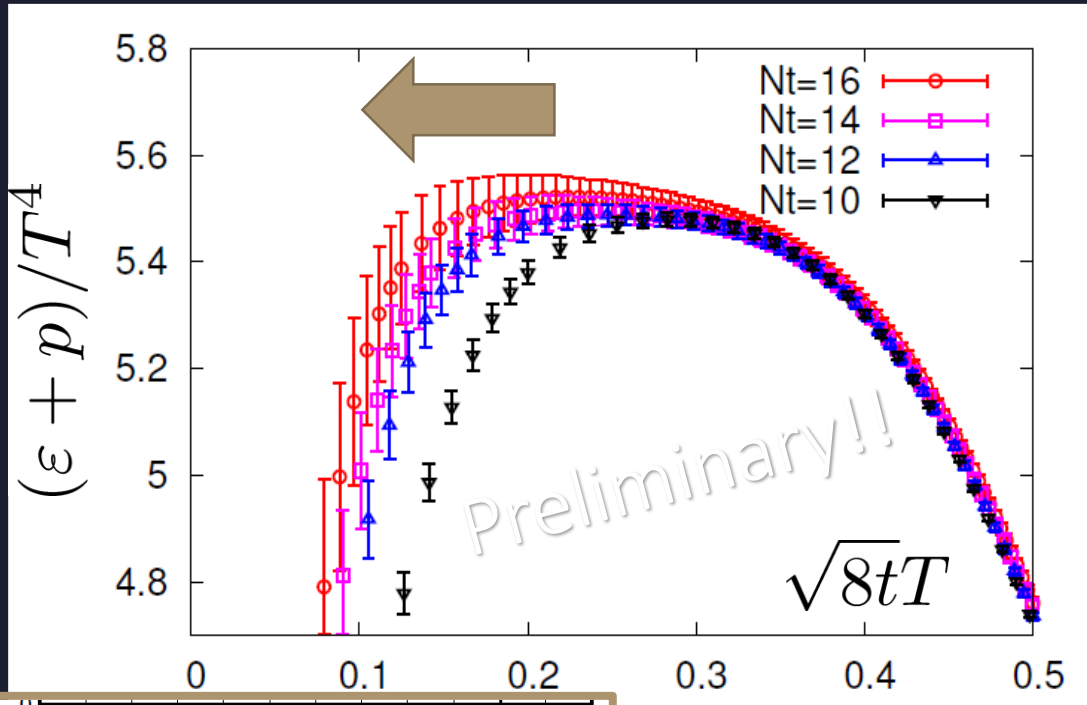
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- lattice size:  $64^3 \times N_t$
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efficiency ~40%

*twice finer lattice!*

# Entropy Density on Finer Lattices

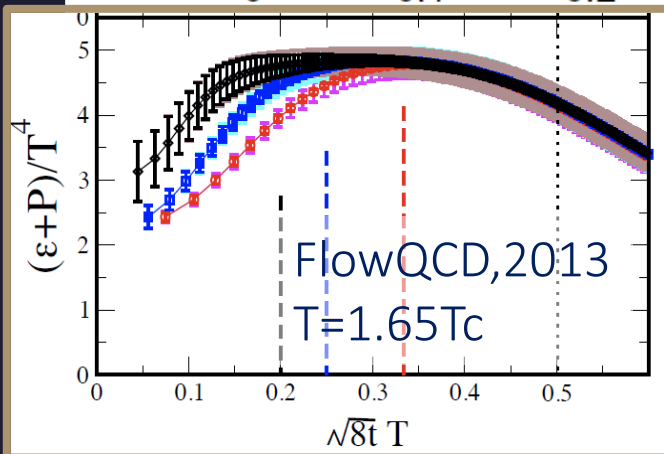


$T = 2.31T_c$

$64^3 \times N_t$

$N_t = 10, 12, 14, 16$

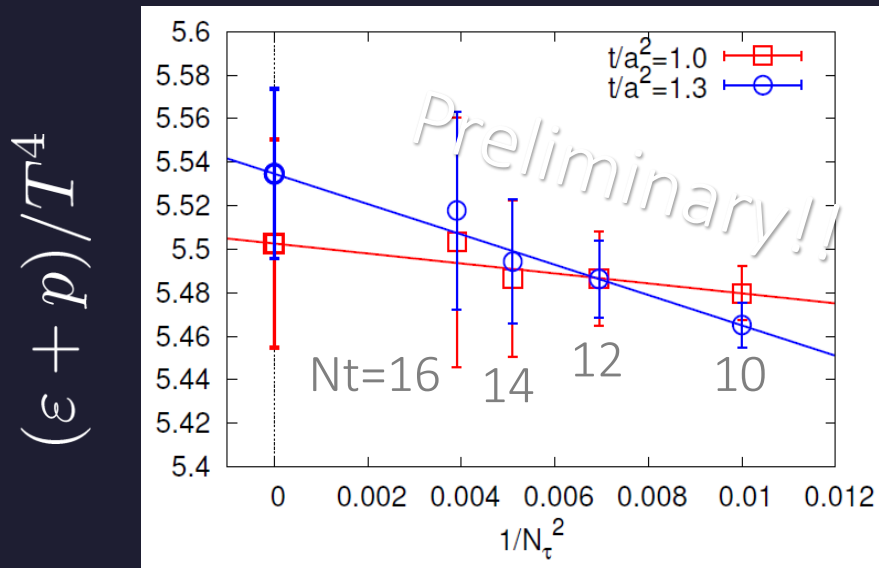
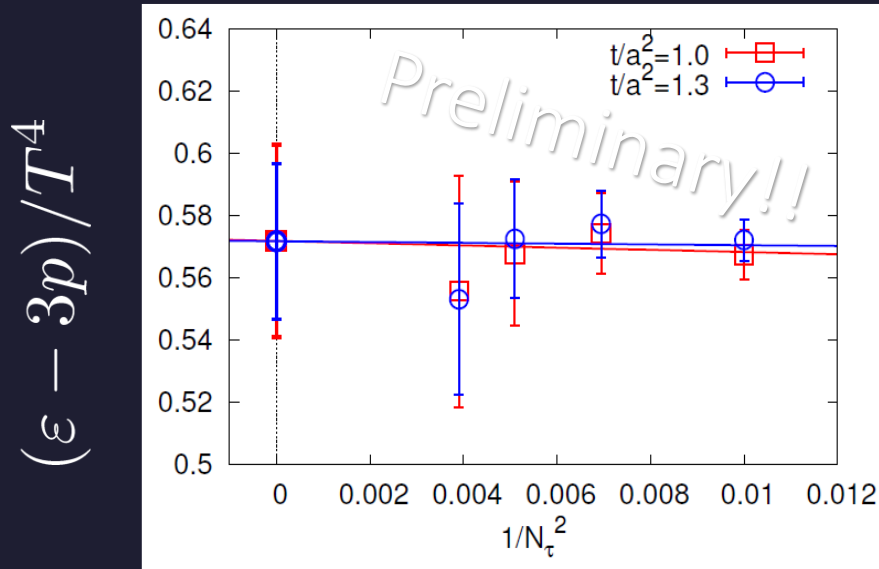
2000 confs.



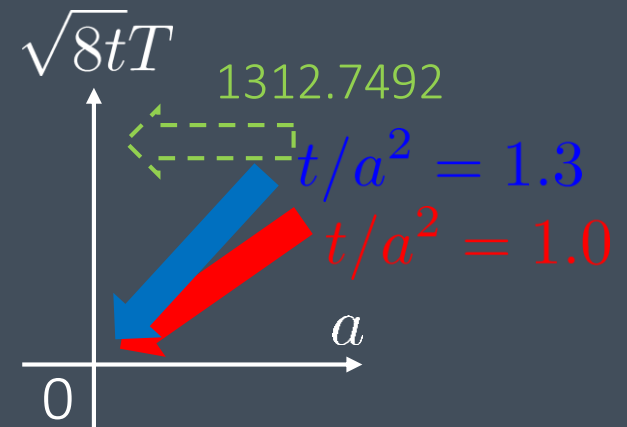
- The wider plateau on the finer lattices
- Plateau may have a nonzero slope

# Continuum Extrapolation

- $T=2.31T_c$
- 2000 confs
- $Nt = 10 \sim 16$



$a \rightarrow 0$  limit with fixed  $t/a^2$



Continuum extrapolation  
is stable

# Summary

$$T_{\mu\nu}^R(x)$$

# Summary

## EMT formula from gradient flow

$$T_{\mu\nu}^R(x) = \lim_{t \rightarrow 0} \left[ \frac{1}{\alpha_U(t)} U_{\mu\nu}(t, x) + \frac{\delta_{\mu\nu}}{4\alpha_E(t)} E(t, x)_{\text{subt.}} \right]$$

This formula can successfully define and calculate the EMT on the lattice

This operator provides us with novel approaches to measure various observables on the lattice!

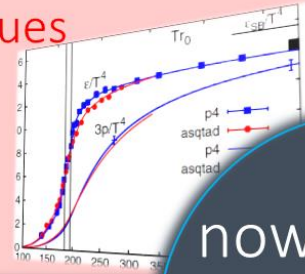
This method is direct, intuitive and less noisy

# Many Future Studies!!

## Thermodynamics

direct measurement of expectation values

$$\langle T_{00} \rangle, \langle T_{ii} \rangle$$



## Fluctuations and Correlations

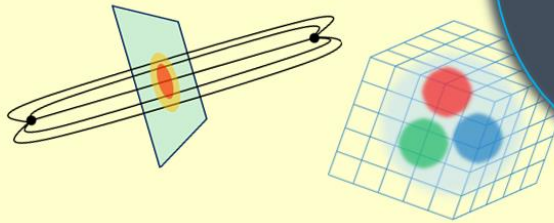
viscosity, specific heat, ...

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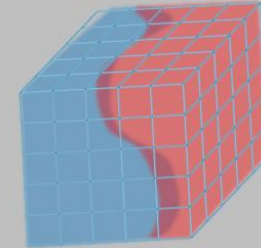
now we have

$$T_{\mu\nu}$$



- confinement string
- EM distribution in hadrons

## Hadron Structure



- vacuum configuration
- mixed state on 1<sup>st</sup> transition

## Vacuum Structure

Other observables

full QCD Makino, Suzuki, 2014

non-pert. improvement Patella 7E(Thu)

O(a) improvement

Nogradi, 7E(Thu); Sint, 7E(Thu)

Monahan, 7E(Thu)

and etc.



One More Thing...

# One More Thing...

## Fluctuations and Correlations

viscosity, specific heat, ...

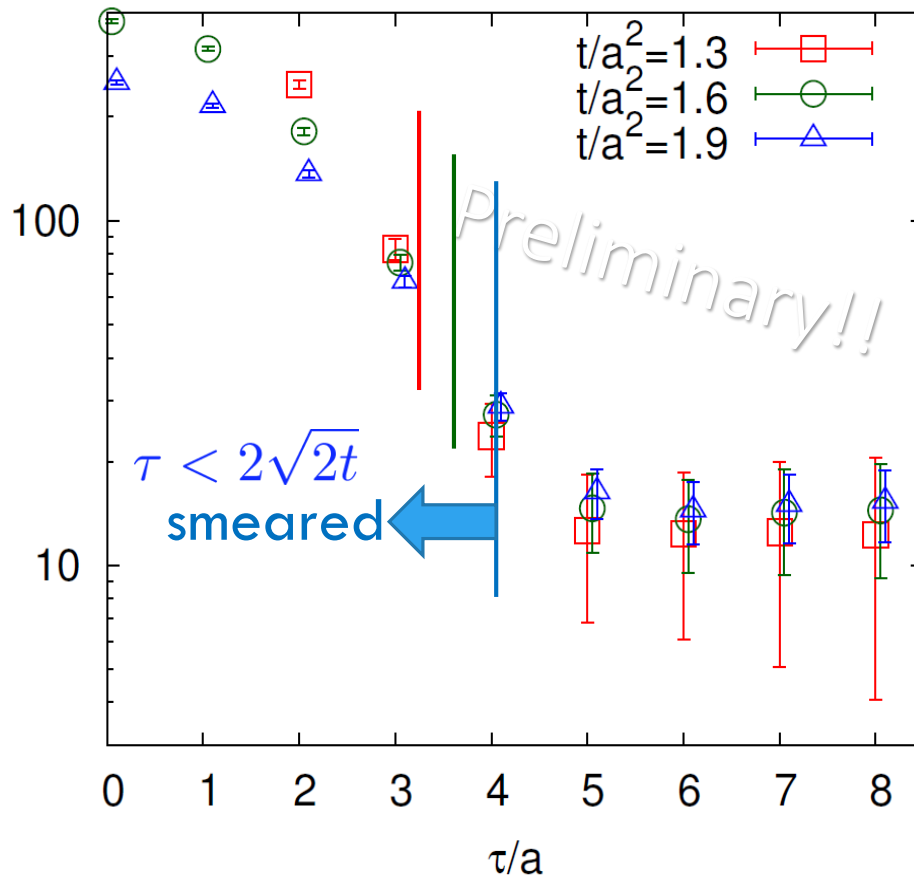
$$c_V \sim \langle \delta T_{00}^2 \rangle$$

$$\eta = \langle T_{12}; T_{12} \rangle$$

# Energy Correlation Function

$$\langle T_{00}(\tau)T_{00}(0) \rangle$$

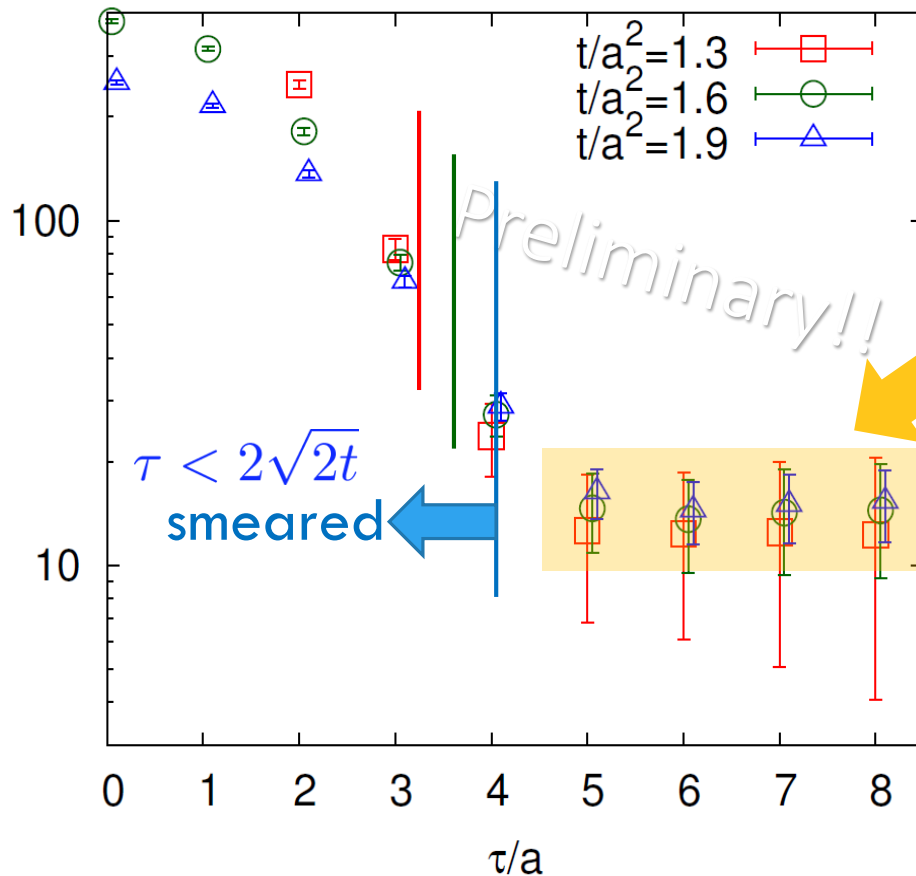
$T=2.31T_c$   
 $b=7.2, Nt=16$   
2000 confs  
 $p=0$  correlator



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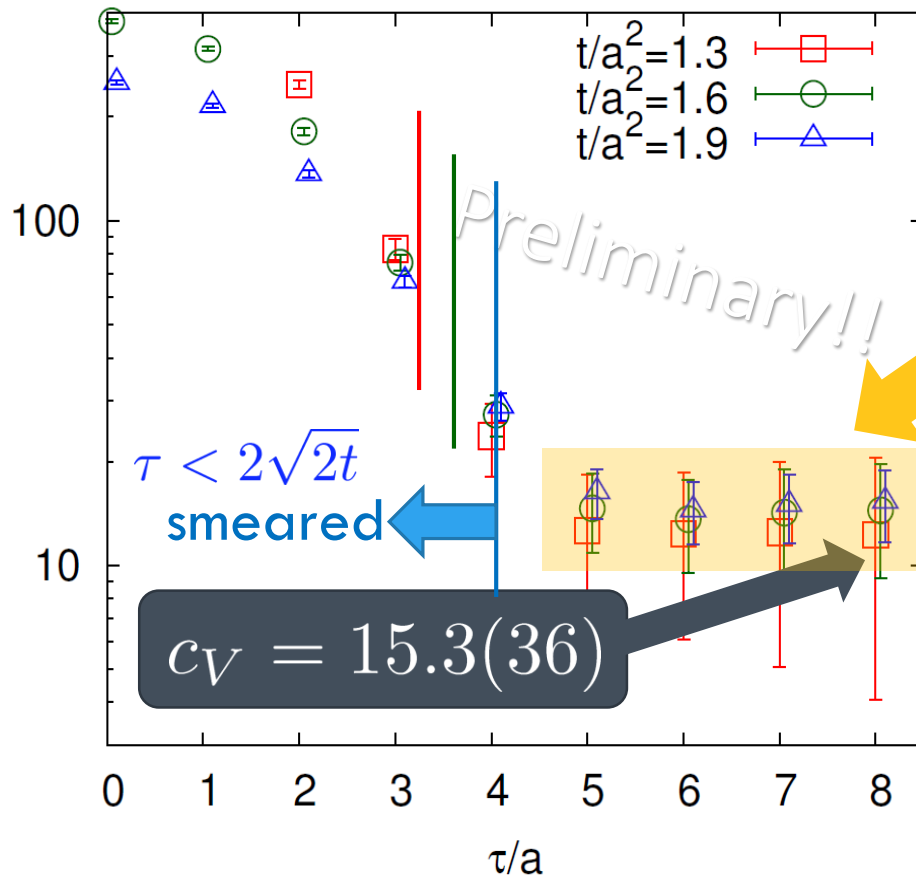


□  $\tau$  independent const.  
→ energy conservation

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$\square$   $\tau$  independent const.  
 $\rightarrow$  energy conservation

$\square$  specific heat

$$c_V = \frac{\langle \delta E^2 \rangle}{VT^2}$$

$\rightarrow$  Novel approach to  
 measure specific heat!

Keep your attention to this new **flow**

just like...

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Thank you for your attention!