

Gradient Flow and Energy-Momentum Tensor in Lattice Gauge Theory

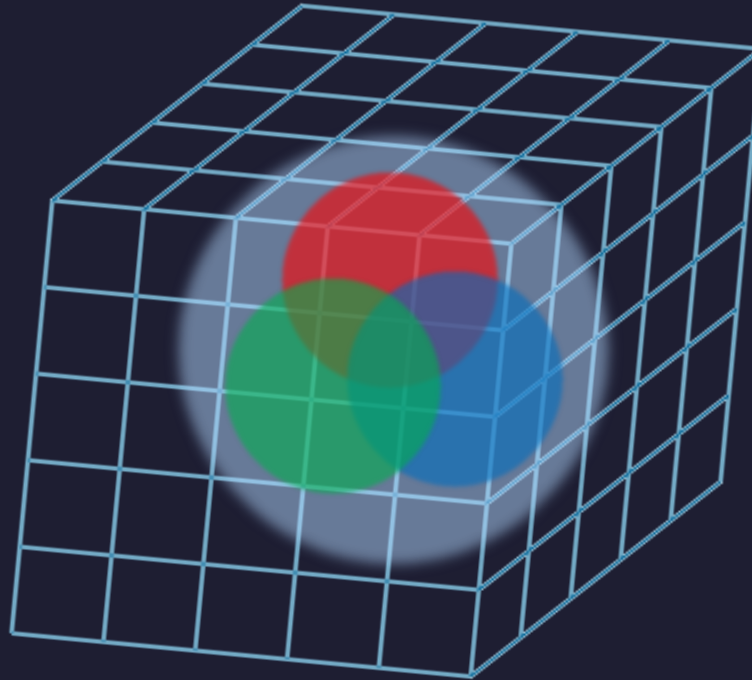
Masakiyo Kitazawa (Osaka U.)

for FlowQCD Collaboration

Asakawa, Hatsuda, Iritani, Itou, MK, Suzuki

FlowQCD, arXiv:1312.7492[hep-lat]; to appear in PRD

Lattice QCD



First principle calculation of QCD
Monte Carlo for path integral

hadron spectra, chiral symmetry, phase transition, etc.

Gradient Flow

$$\partial_t A_\mu(t, x) = -\frac{\partial S_{\text{YM}}}{\partial A_\mu}$$

Luscher, 2010

A powerful tool for various analyses on the lattice

Gradient Flow

$$\partial_t A_\mu(t, x) = - \frac{\partial S_{\text{YM}}}{\partial A_\mu}$$

Luscher, 2010

A powerful tool for various analyses on the lattice

Why care?

D. Negradi, LATTICE2014,7B

- Tuesday 14:55 – Nathan Brown – Gradient Flow Analysis on MILC HISQ Ensembles
- Tuesday 14:35 – Andrea Shindler – Beyond the Standard Model Matrix Elements with the gradient flow
- Tuesday 14:35 – Liam Keegan – TEK twisted gradient flow running coupling
- Wednesday 09:00 – Anna Hasenfratz – Improved gradient flow for step scaling function and scale setting
- Wednesday 09:20 – Jarno Rantaharju – The gradient flow running coupling in SU2 with 8 flavors
- Wednesday 11:10 – Marco Ce – Testing the WittenVeneziano mechanism with the YangMills gradient flow on the lattice
- Thursday 14:55 – Agostino Patella – Energy-momentum tensor on the lattice and Wilson flow
- Thursday 15:15 – Masanori Okawa – String tension from smearing and Wilson flow methods
- Thursday 15:55 – Stefan Sint – How to reduce $O(a^2)$ effects in gradient flow observables
- Friday 10:15 – Alberto Ramos – Wilson flow and renormalization
- Saturday 09:30 – Kitazawa Masakiyo – Measurement of thermodynamics using Gradient Flow

Gradient Flow

$$\partial_t A_\mu(t, x) = - \frac{\partial S_{\text{YM}}}{\partial A_\mu}$$

Luscher, 2010

A powerful tool for various analyses on the lattice

Applications:

- ① scale setting
- ② running coupling
- ③ topology
- ④ operator relation
- ⑤ autocorrelation
- ⑥ etc.

thermodynamics

EMT correlator

Lattice Scale Setting

□ $a(\beta)$

□ previous references

- string tension
- Sommer scale

SU(3) pure YM

Edwards, Heller, Klassen, 1998

$\beta < 6.56$

Wilson gauge

Alpha-Collab., 1998

$\beta < 6.57$

Necco, Sommer, 2002

$\beta < 6.92$

(Durr, Fodor, Hoelbling, 2007

$\beta < 6.92$)

We perform the precision scale setting of
SU(3) YM theory up to $\beta=7.5$ using gradient flow

$T_{\mu\nu}$

Poincare
symmetry

$T_{\mu\nu}$

	momentum		
energy	T_{01}	T_{02}	T_{03}
T_{10}	T_{11}	T_{12}	T_{13}
T_{20}	T_{21}	T_{22}	T_{23}
T_{30}	T_{31}	T_{32}	T_{33}
	stress	pressure	

Einstein Equation

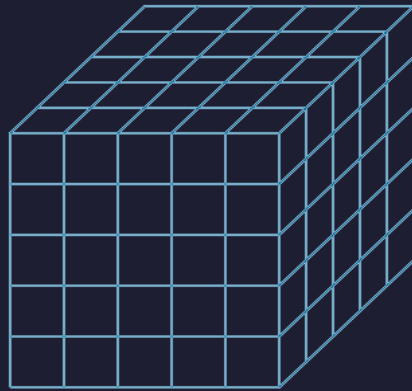
$$G_{\mu\nu} + \Lambda g_{\mu\nu} = \kappa T_{\mu\nu}$$

Hydrodynamic Eq.

$$\partial_{\mu} T_{\mu\nu} = 0$$

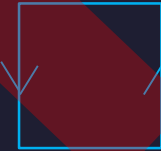
$T_{\mu\nu}$: nontrivial observable
on the lattice

- ① Definition of the operator is nontrivial
because of the explicit breaking of Lorentz symmetry



ex: $T_{\mu\nu} = F_{\mu\rho}F_{\nu\rho} - \frac{1}{4}\delta_{\mu\nu}FF$

$F_{\mu\nu} =$



- ② Its measurement is extremely noisy
due to high dimensionality and etc.

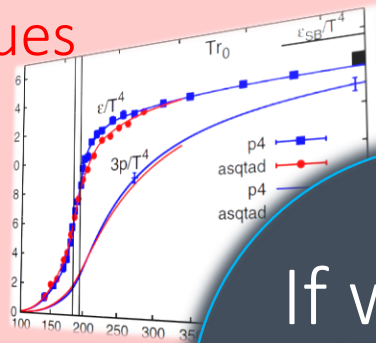
If we have

$$T_{\mu\nu}$$

Thermodynamics

direct measurement of
expectation values

$$\langle T_{00} \rangle, \langle T_{ii} \rangle$$



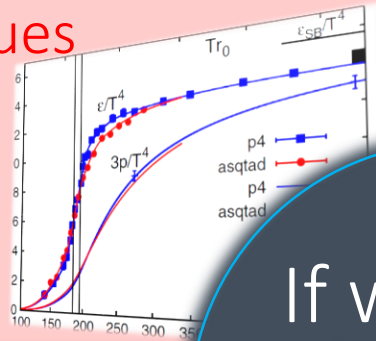
If we have

$$T_{\mu\nu}$$

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$$\langle T_{00} \rangle, \langle T_{ii} \rangle$$



Fluctuations and Correlations

viscosity, specific heat, ...

$$\eta = \int_0^\infty dt \langle T_{12}; T_{12} \rangle$$

$$c_V \sim \langle \delta T_{00}^2 \rangle$$

If we have

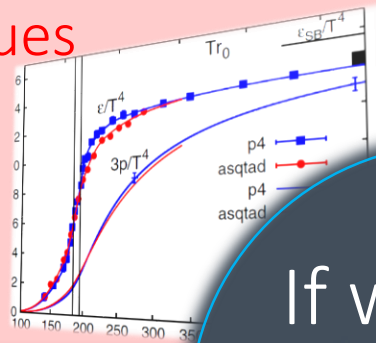
$$T_{\mu\nu}$$

We construct the EMT using gradient flow
and measure these quantities

Thermodynamics

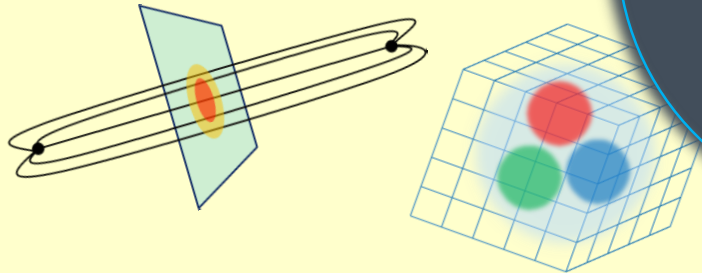
direct measurement of
expectation values

$$\langle T_{00} \rangle, \langle T_{ii} \rangle$$



If we have

$$T_{\mu\nu}$$



- confinement string
- EM distribution in hadrons

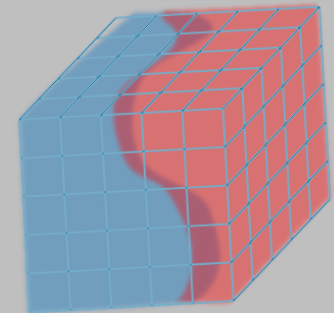
Hadron Structure

Fluctuations and Correlations

viscosity, specific heat, ...

$$\eta = \int_0^\infty dt \langle T_{12}; T_{12} \rangle$$

$$c_V \sim \langle \delta T_{00}^2 \rangle$$



- vacuum configuration
- mixed state on 1st transition

Vacuum Structure

Gradient Flow

YM Gradient Flow

Luscher, 2010

$$\partial_t A_\mu(t, x) = - \frac{\partial S_{\text{YM}}}{\partial A_\mu}$$

t: "flow time"
dim:[length²]

$$A_\mu(0, x) = A_\mu(x)$$

YM Gradient Flow

Luscher, 2010

$$\partial_t A_\mu(t, x) = - \frac{\partial S_{\text{YM}}}{\partial A_\mu}$$

$$A_\mu(0, x) = A_\mu(x)$$

t: "flow time"
dim:[length²]

- transform gauge field like diffusion equation

$$\partial_t A_\mu = D_\nu G_{\mu\nu} = \partial_\nu \partial_\nu A_\mu + \dots$$

- diffusion length $d \sim \sqrt{8t}$

- This is **NOT** the standard cooling/smearing

- All composite operators at $t > 0$ are UV finite Luescher,Weisz,2011

Lattice Scale Setting

Flow Time Dep. of an Observable

Luscher, 2010

$\langle \mathcal{O}(t) \rangle$: universal function of t

\mathcal{O} : an observable

use this function used to determine $a(\beta)$

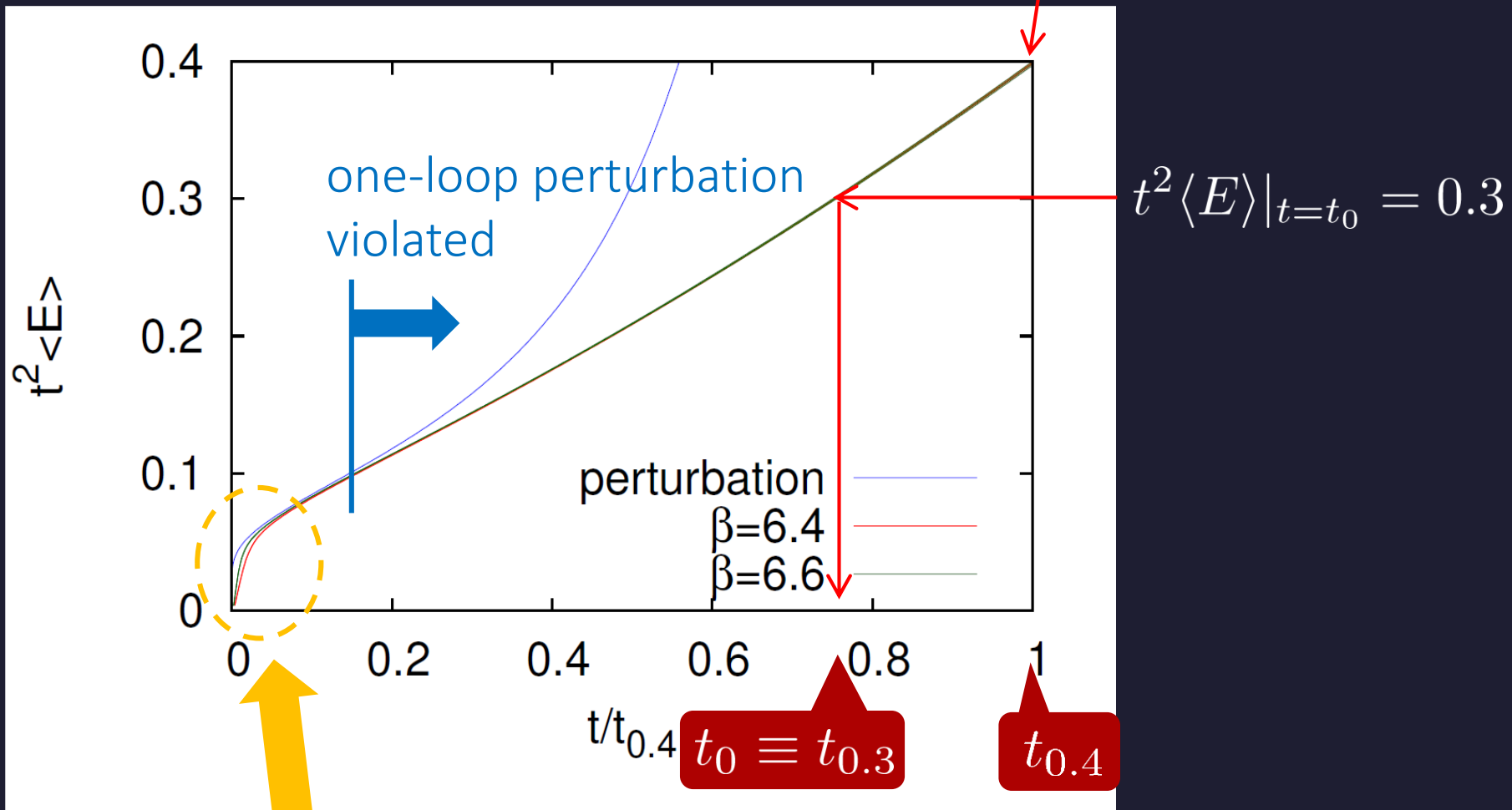
$$\langle \mathcal{O}(t_0) \rangle = \text{const} \Rightarrow t_0 = \hat{t} a^2$$

□ standard choice of \mathcal{O} : $\mathcal{O}(t) = \frac{1}{4} F_{\mu\nu} F^{\mu\nu} \equiv E$

□ perturbative formula: $t^2 \langle E \rangle = \frac{3}{(4\pi)^2} g^2 (1 + k_1 g^2 + \dots)$ $g = g(1/\sqrt{8t})$

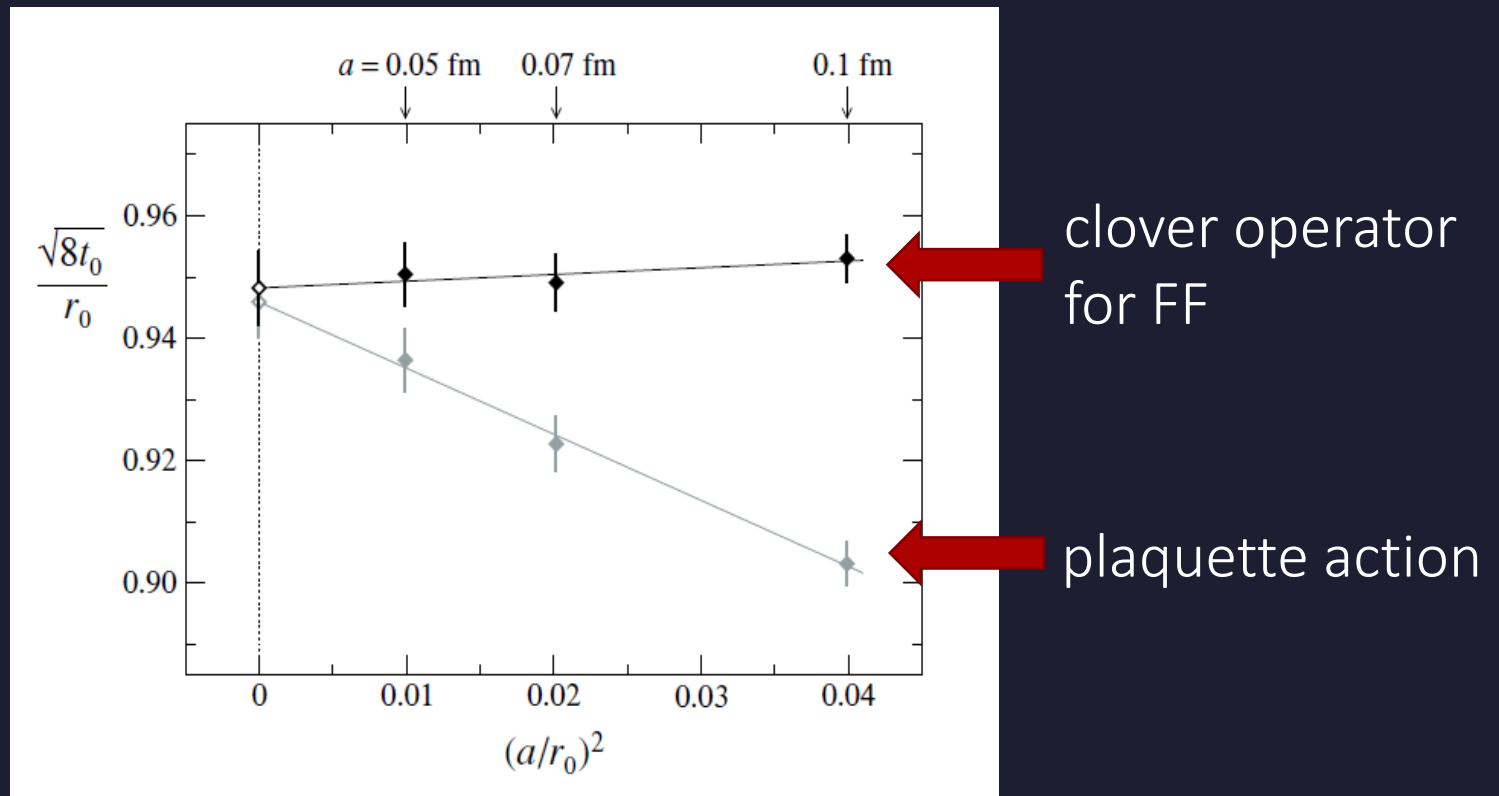
t dep of $t^2 \langle E \rangle$

$$t^2 \langle E \rangle|_{t=t_{0.4}} = 0.4$$



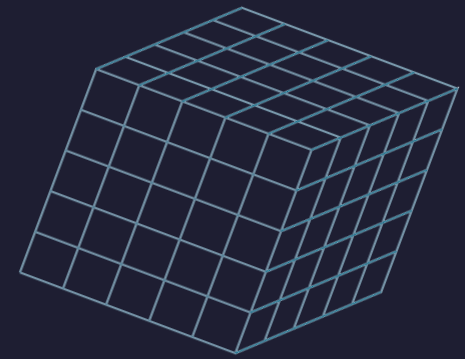
lattice discretization effect

$t_{\#}$ Scale Setting



good agreement with previous scales

Numerical Analysis

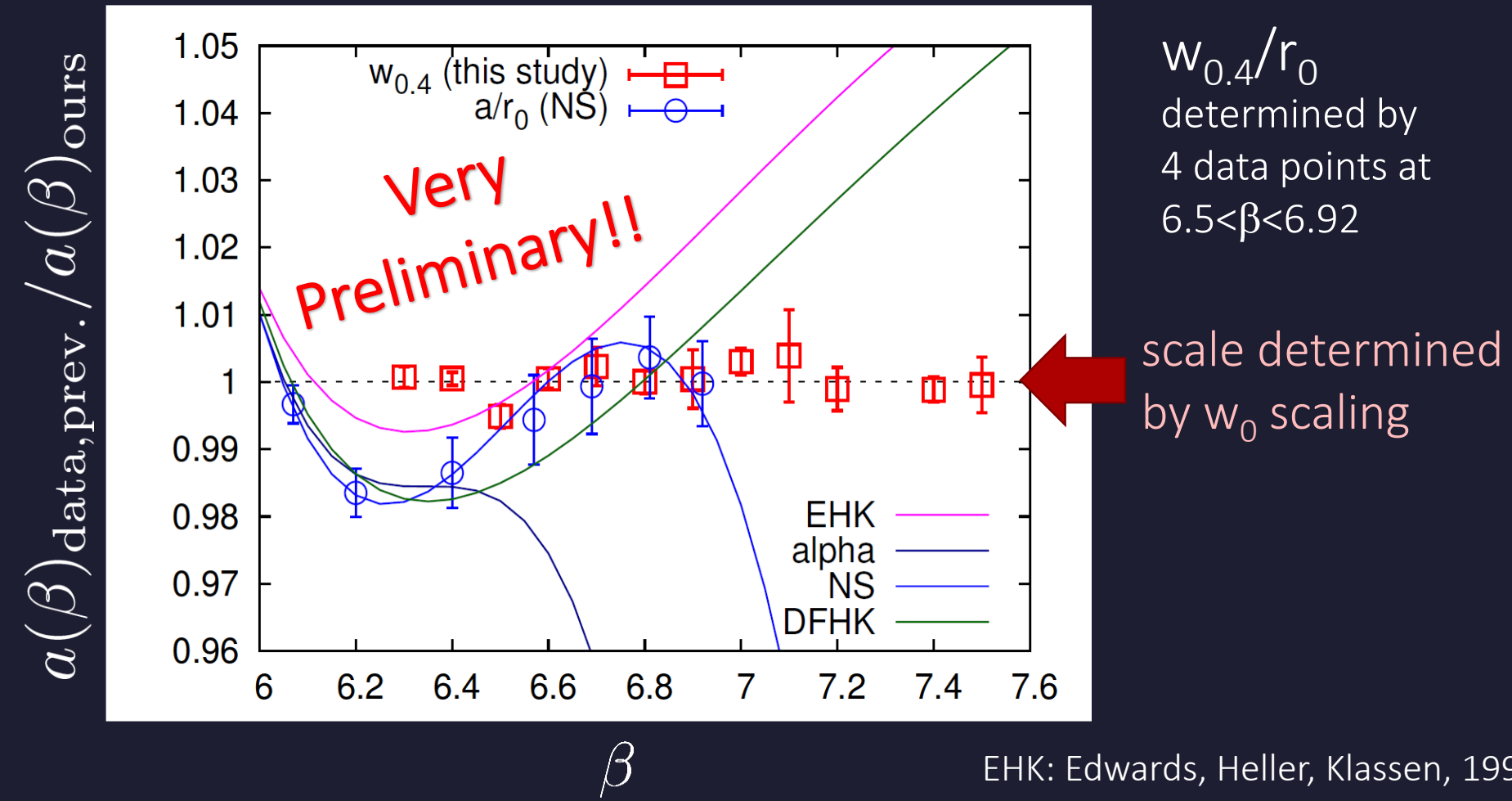


- SU(3) YM theory
- Wilson gauge action
- w_0 scaling

β	size	N_{conf}	β	size	N_{conf}
6.3	64^4	30	6.9	64^4	30
6.4	64^4	100	7.0	96^4	100
6.5	64^4	30	7.2	96^4	50
6.6	64^4	100	7.4	128^4	40
6.7	64^4	30	7.5	128^4	40
6.8	64^4	100			

each configuration is separated by 1000 gauge updates (HB+OR⁵)

Comparison with Previous Studies



$w_{0.4}/r_0$
determined by
4 data points at
 $6.5 < \beta < 6.92$

EHK: Edwards, Heller, Klassen, 1998
Alpha-Collaboration, 1998
NS: Necco, Sommer, 2002
DFHK: Durr, Fodor, Hoelbling, 2007

Small Flow Time Expansion of Operators and EMT



Operator Relation

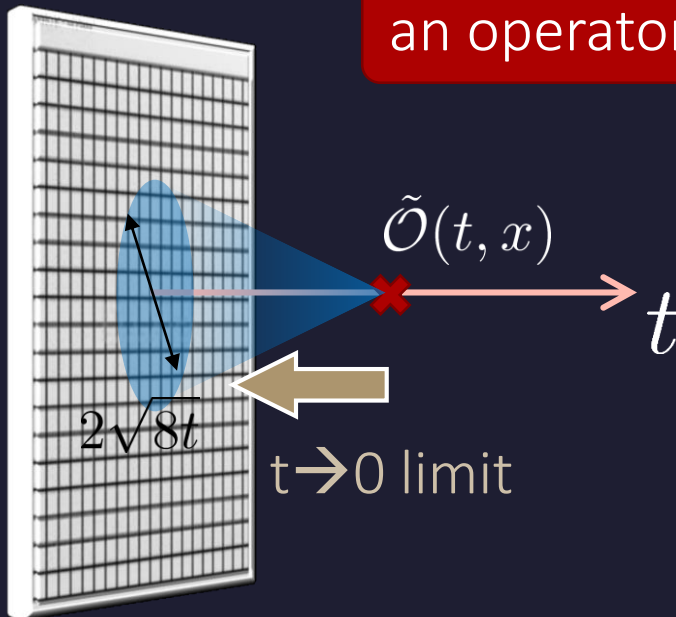
Luescher, Weisz, 2011

$$\tilde{\mathcal{O}}(t, x) \xrightarrow[t \rightarrow 0]{} \sum_i c_i(t) \mathcal{O}_i^R(x)$$

an operator at $t > 0$

renormalized operators
of original theory

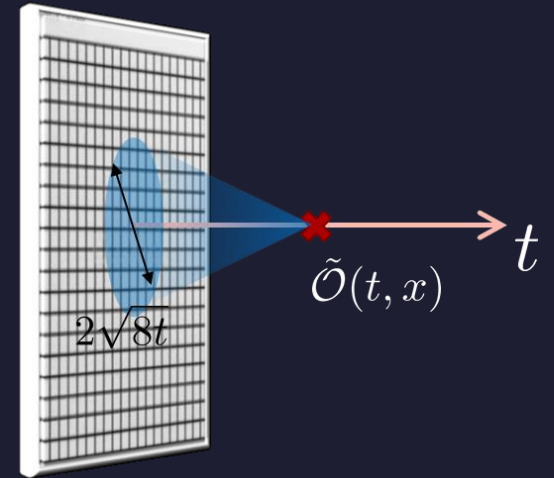
original 4-dim theory



Constructing EMT

Suzuki, 2013
DelDebbio, Patella, Rago, 2013

$$\tilde{\mathcal{O}}(t, x) \xrightarrow{t \rightarrow 0} \sum_i c_i(t) \mathcal{O}_i^R(x)$$



□ gauge-invariant dimension 4 operators

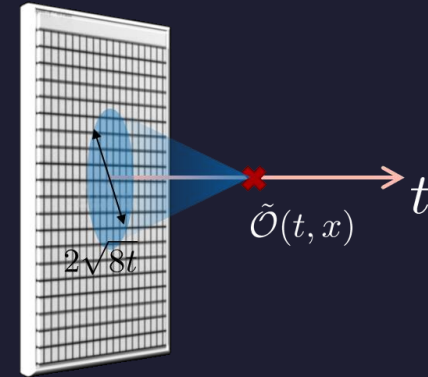
$$\left\{ \begin{array}{l} U_{\mu\nu}(t, x) = G_{\mu\rho}(t, x)G_{\nu\rho}(t, x) - \frac{1}{4}\delta_{\mu\nu}G_{\mu\nu}(t, x)G_{\mu\nu}(t, x) \\ E(t, x) = \frac{1}{4}\delta_{\mu\nu}G_{\mu\nu}(t, x)G_{\mu\nu}(t, x) \end{array} \right.$$

Constructing EMT 2

Suzuki, 2013

$$U_{\mu\nu}(t, x) = \alpha_U(t) \left[T_{\mu\nu}^R(x) - \frac{1}{4} \delta_{\mu\nu} T_{\rho\rho}^R(x) \right] + \mathcal{O}(t)$$

$$E(t, x) = \langle E(t, x) \rangle + \alpha_E(t) T_{\rho\rho}^R(x) + \mathcal{O}(t)$$



Suzuki coeffs. $\left\{ \begin{array}{l} \alpha_U(t) = g^2 [1 + 2b_0 s_1 g^2 + O(g^4)] \\ \alpha_E(t) = \frac{1}{2b_0} [1 + 2b_0 s_2 g^2 + O(g^4)] \end{array} \right.$

$$\begin{aligned} g &= g(1/\sqrt{8t}) \\ s_1 &= 0.03296\dots \\ s_2 &= 0.19783\dots \end{aligned}$$

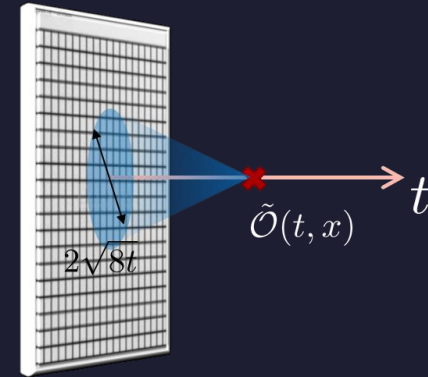
See also, Patella, Parallel7E, Thu.

Constructing EMT 2

Suzuki, 2013

$$U_{\mu\nu}(t, x) = \alpha_U(t) \left[T_{\mu\nu}^R(x) - \frac{1}{4} \delta_{\mu\nu} T_{\rho\rho}^R(x) \right] + \mathcal{O}(t)$$

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$$g = g(1/\sqrt{8t})$$

$$s_1 = 0.03296 \dots$$

$$s_2 = 0.19783 \dots$$

Remormalized EMT

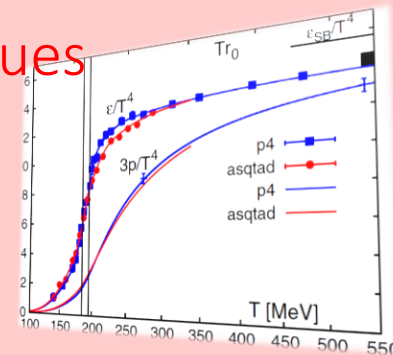
$$T_{\mu\nu}^R(x) = \lim_{t \rightarrow 0} \left[\frac{1}{\alpha_U(t)} U_{\mu\nu}(t, x) + \frac{\delta_{\mu\nu}}{4\alpha_E(t)} E(t, x)_{\text{subt.}} \right]$$

Numerical Analysis: thermodynamics

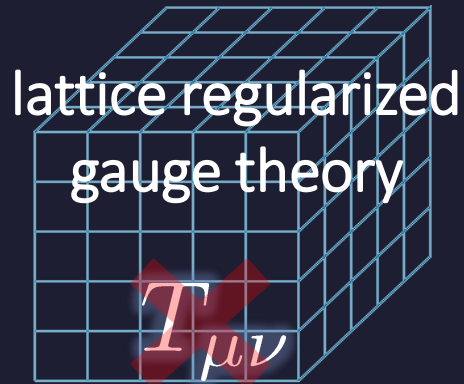
Thermodynamics

direct measurement of
expectation values

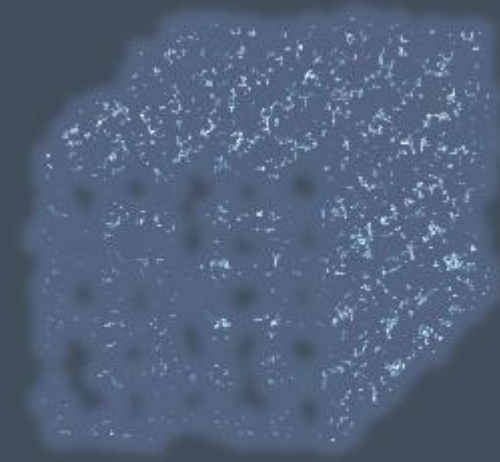
$$\langle T_{00} \rangle, \langle T_{ii} \rangle$$



Gradient Flow Method



gradient flow



$$T_{\mu\nu}^R$$

continuum theory
(with dim. reg.)

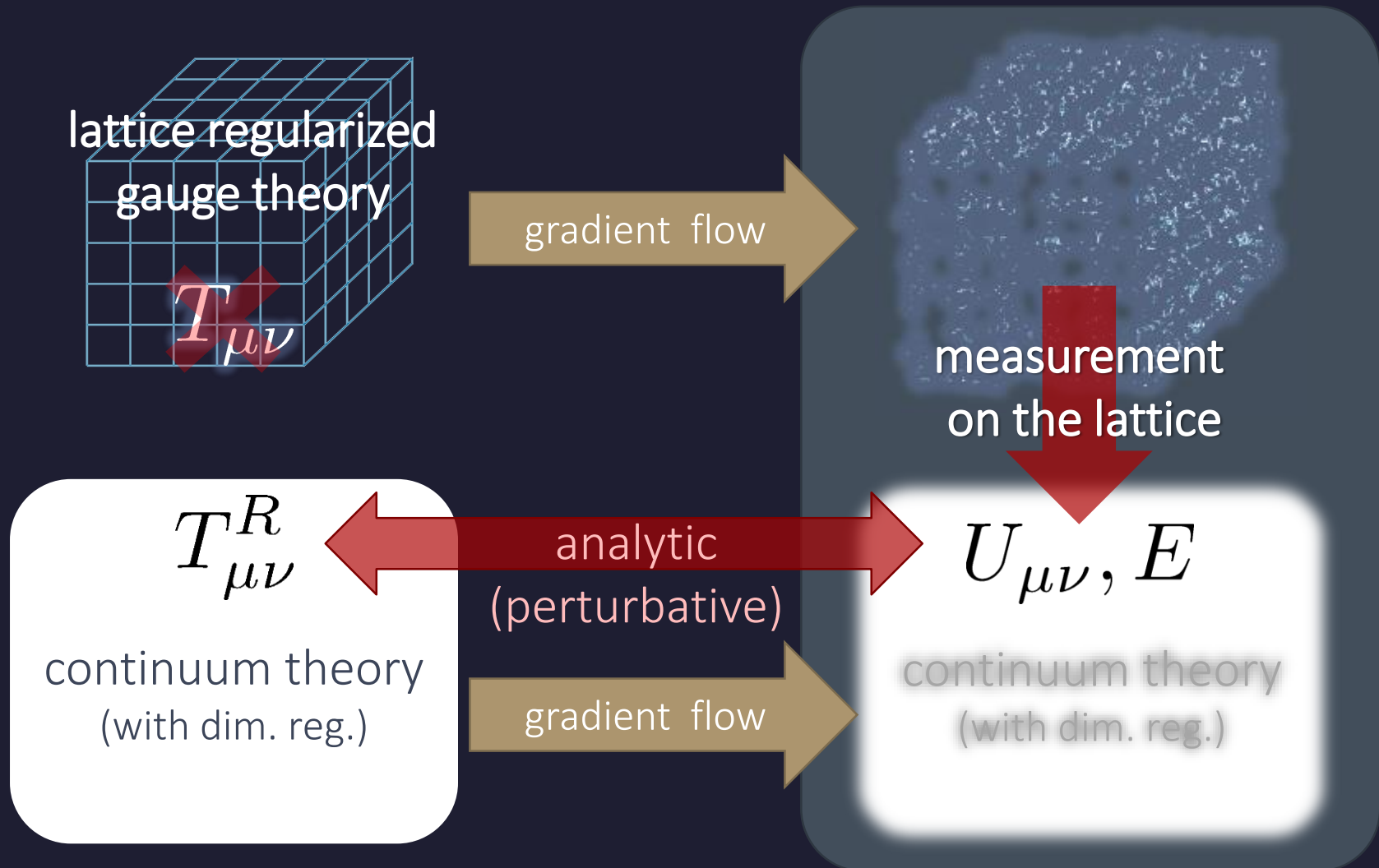
analytic
(perturbative)

$$U_{\mu\nu}, E$$

continuum theory
(with dim. reg.)

gradient flow

Gradient Flow Method



Caveats



Perturbative relation
has to be applicable!
 $\sqrt{8t} \ll \Lambda^{-1}, T^{-1}$

Gauge field has to be
sufficiently smeared!
 $a \ll \sqrt{8t}$



measurement
on the lattice



$T R_{\mu\nu}$
continuum theory
(with dim. reg.)

analytic
(perturbative)

$U_{\mu\nu}, E$

continuum theory
(with dim. reg.)

gradient flow

Caveats



Perturbative relation has to be applicable!
 $\sqrt{8t} \ll \Lambda^{-1}, T^{-1}$

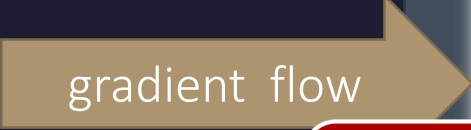
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$T R_{\mu\nu}$
 continuum theory (with dim. reg.)

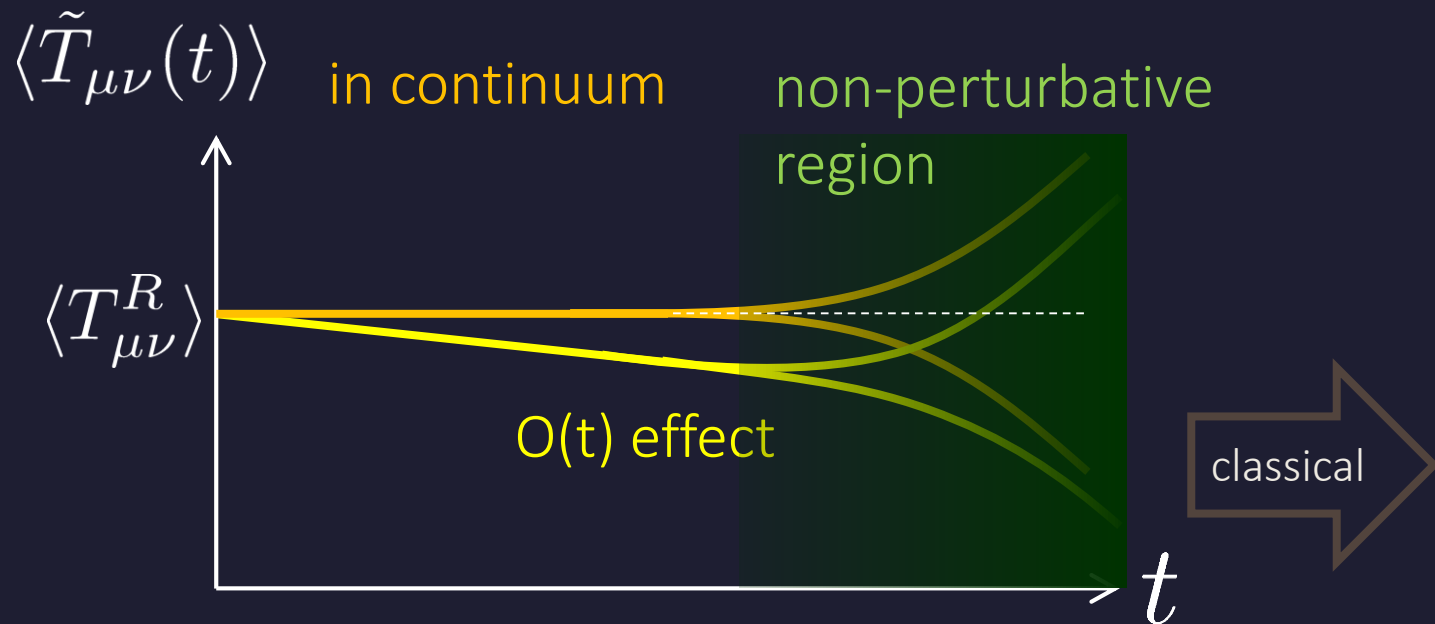
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$U_{\mu\nu}, E$
 continuum theory (with dim. reg.)

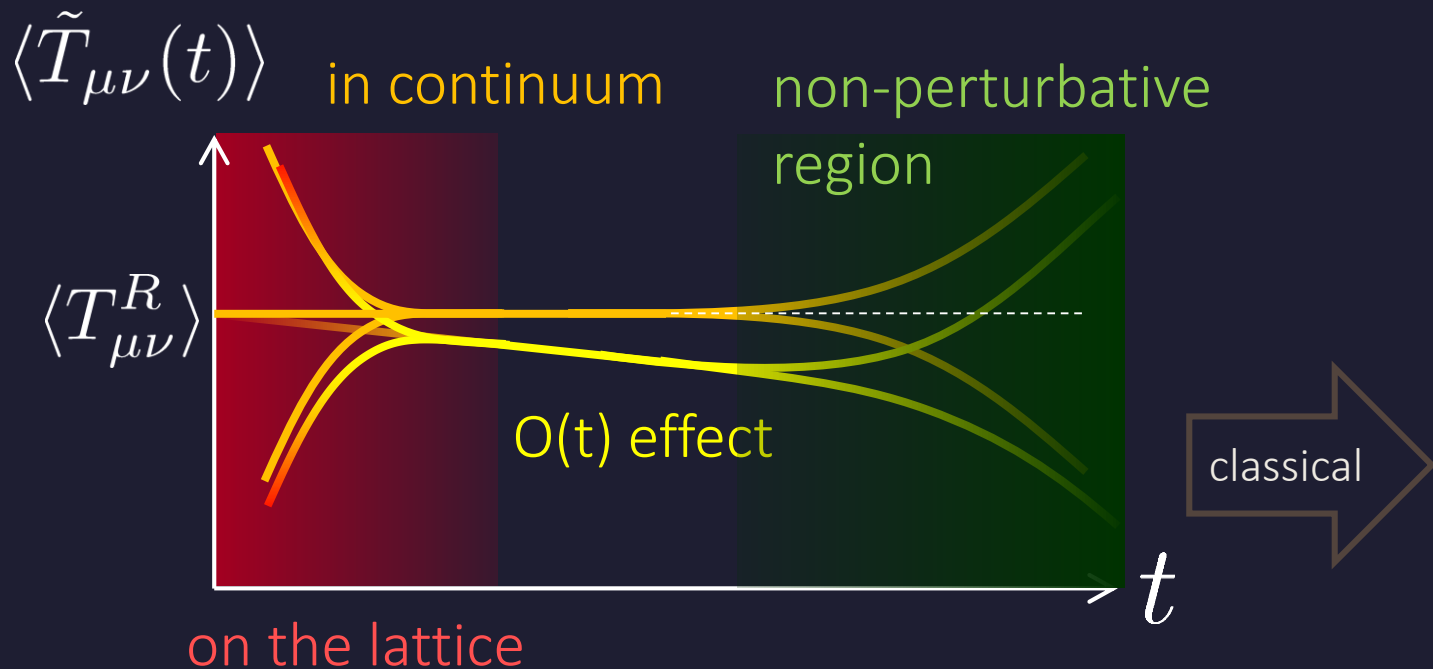


$a \ll \sqrt{8t} \ll \Lambda^{-1}$

$$\tilde{T}_{\mu\nu}(t) = \frac{1}{\alpha_U(t)} U_{\mu\nu}(t) + \frac{\delta_{\mu\nu}}{4\alpha_E(t)} E(t)_{\text{subt.}} \quad T_{\mu\nu}^R = \lim_{t \rightarrow 0} \tilde{T}_{\mu\nu}(t)$$



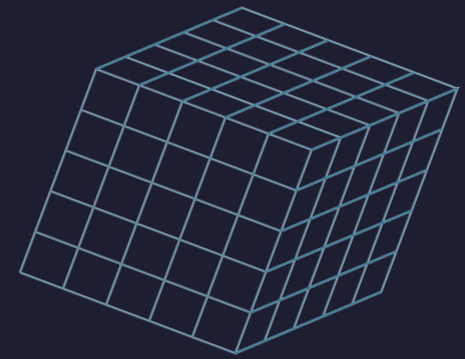
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□ $t \rightarrow 0$ limit with keeping $t \gg a^2$

Numerical Simulation

- SU(3) YM theory
- Wilson gauge action



twice finer lattice!

Simulation 1

(arXiv:1312.7492)

- lattice size: $32^3 \times N_t$
- $N_t = 6, 8, 10$
- $\beta = 5.89 - 6.56$
- ~ 300 configurations

using SX8 @ RCNP
SR16000 @ KEK



Simulation 2

(*new*, preliminary)

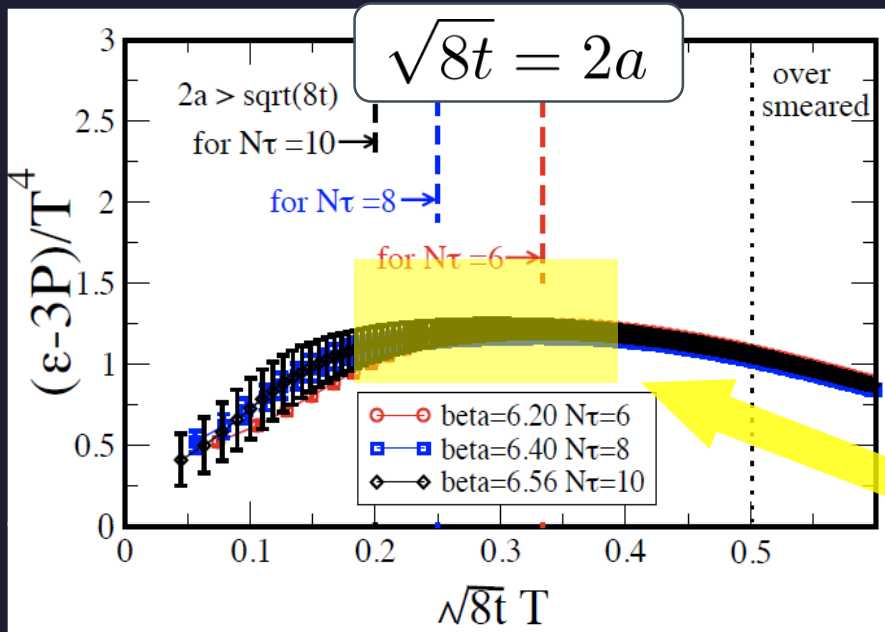
- lattice size: $64^3 \times N_t$
- $N_t = 10, 12, 14, 16$
- $\beta = 6.4 - 7.4$
- ~ 2000 configurations

using BlueGeneQ @ KEK
efficiency $\sim 40\%$

ε -3p at $T=1.65T_c$

$$\tilde{T}_{\mu\nu}(t) = \frac{1}{\alpha_U(t)} U_{\mu\nu}(t) + \frac{\delta_{\mu\nu}}{4\alpha_E(t)} E(t)_{\text{subt.}}$$

$$T_{\mu\nu}^R = \lim_{t \rightarrow 0} \tilde{T}_{\mu\nu}(t)$$



Emergent plateau!

$$2a \lesssim \sqrt{8t} \lesssim 0.4T^{-1}$$

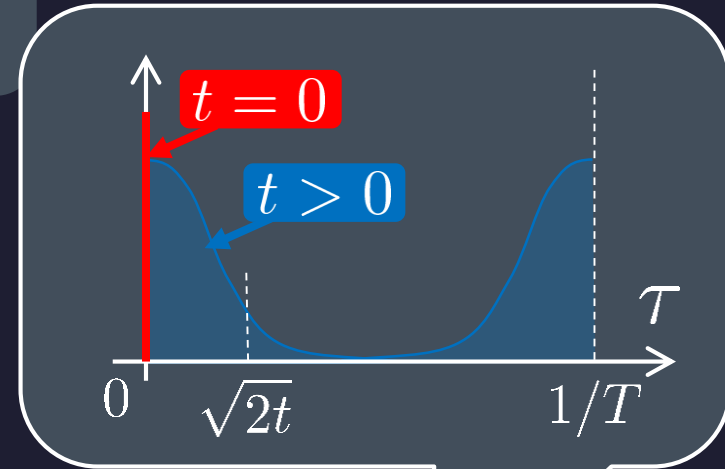
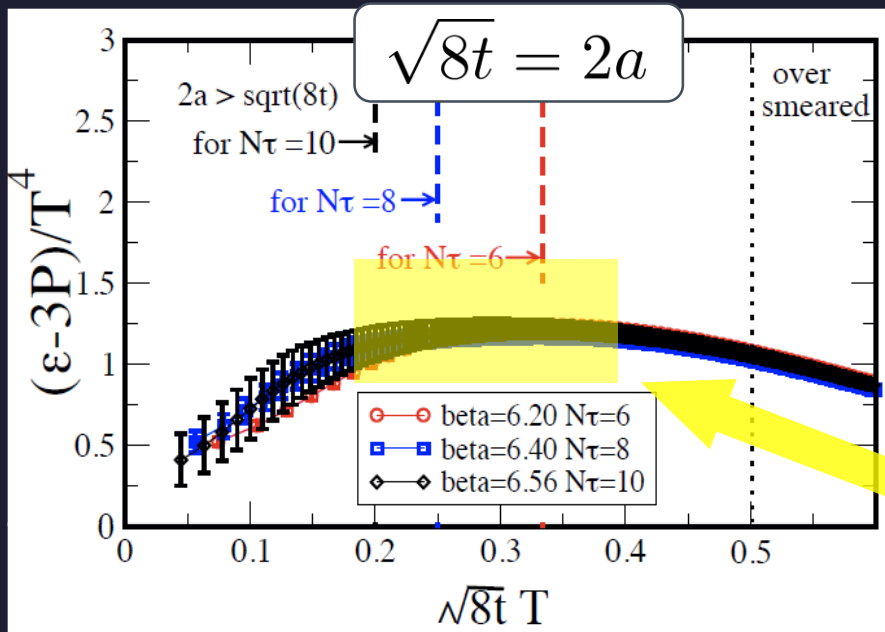
$Nt=6,8,10$
 ~ 300 confs.

the range of t where the EMT formula is successfully used!

ϵ -3p at $T=1.65T_c$

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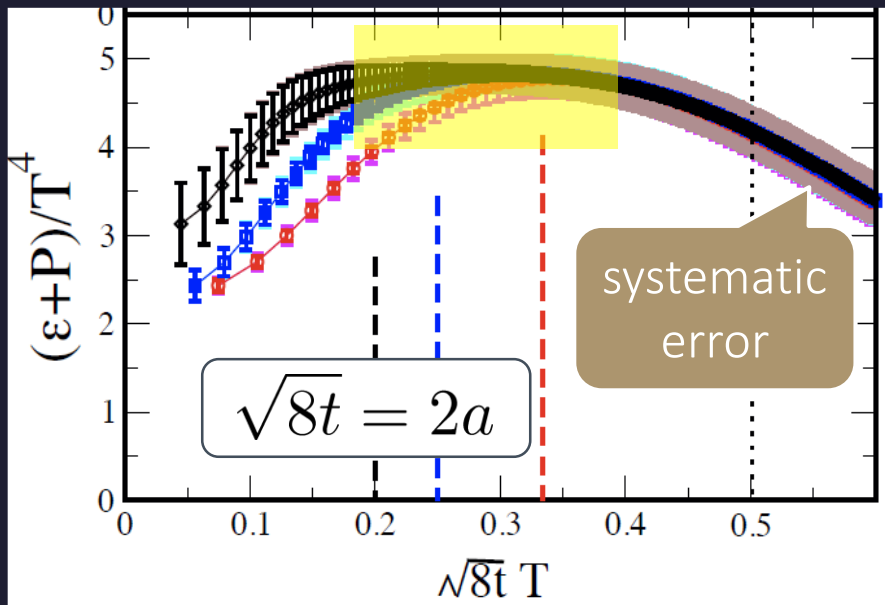
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the range of t where the EMT formula is successfully used!

Entropy Density at $T=1.65T_c$



Emergent plateau!

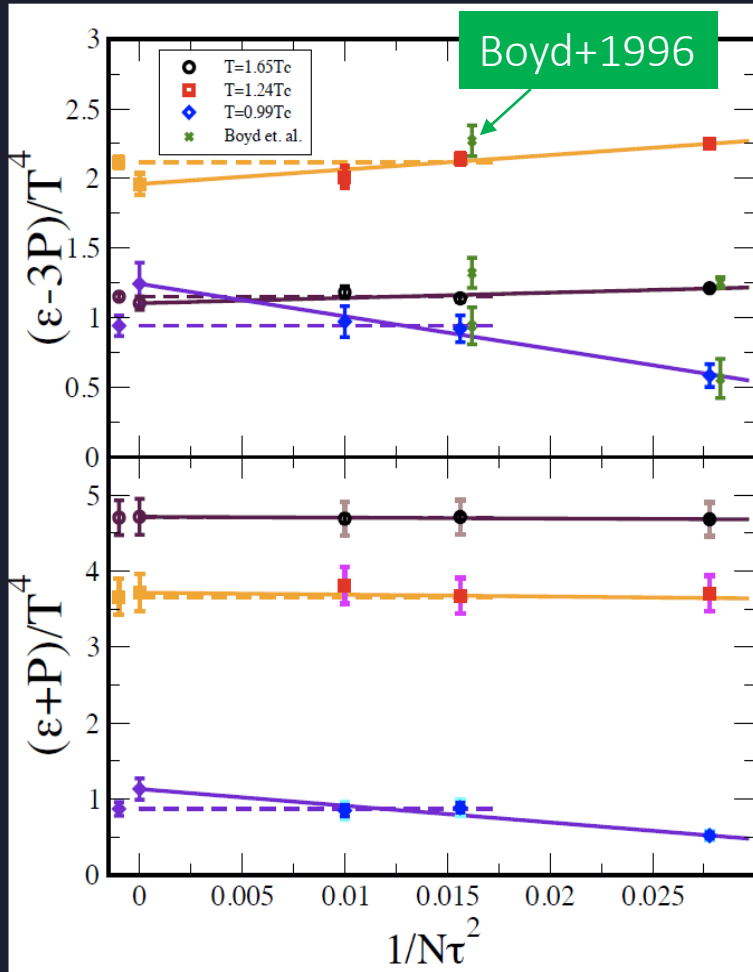
$$2a \lesssim \sqrt{8t} \lesssim 0.4T^{-1}$$

$N_t=6,8,10$
 ~ 300 confs.

Direct measurement of $e+p$ on a given T !

NO integral / **NO** vacuum subtraction

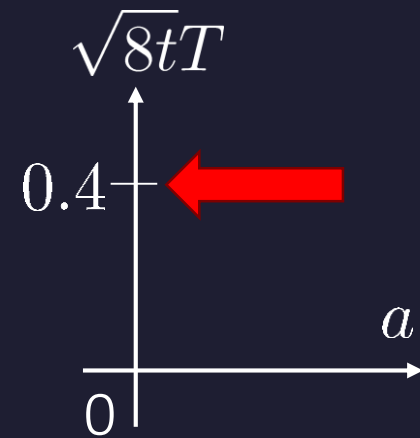
Continuum Limit



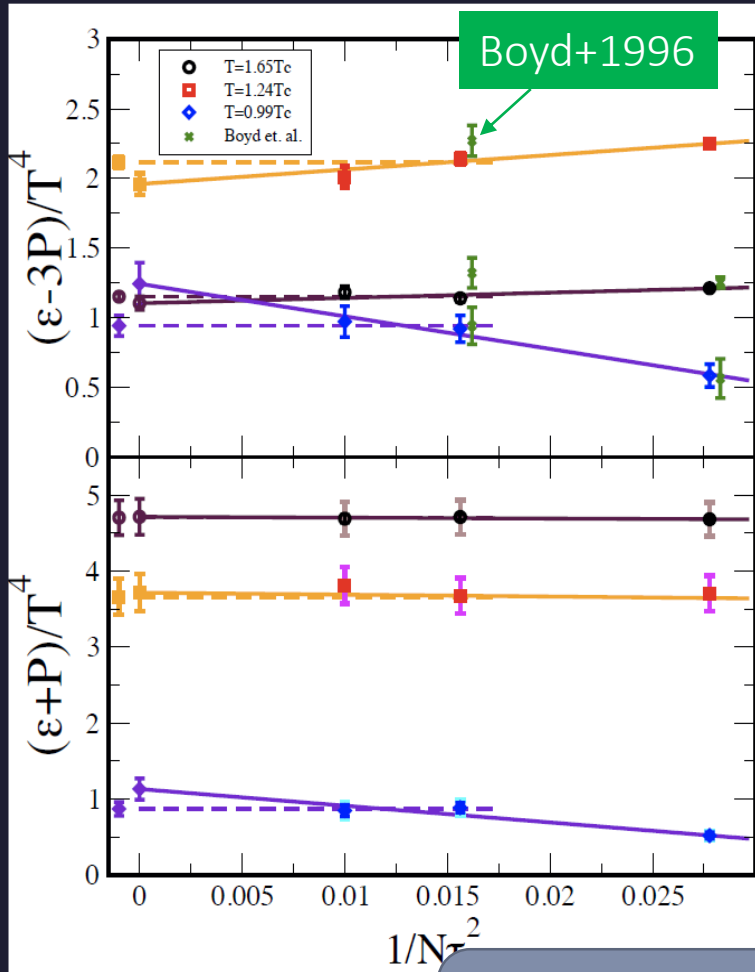
$32^3 \times N_t$

$N_t = 6, 8, 10$

$T/T_c = 0.99, 1.24, 1.65$



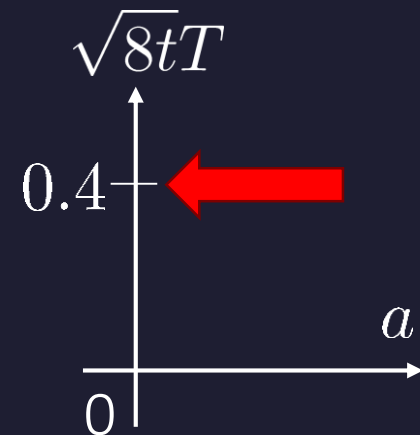
Continuum Limit



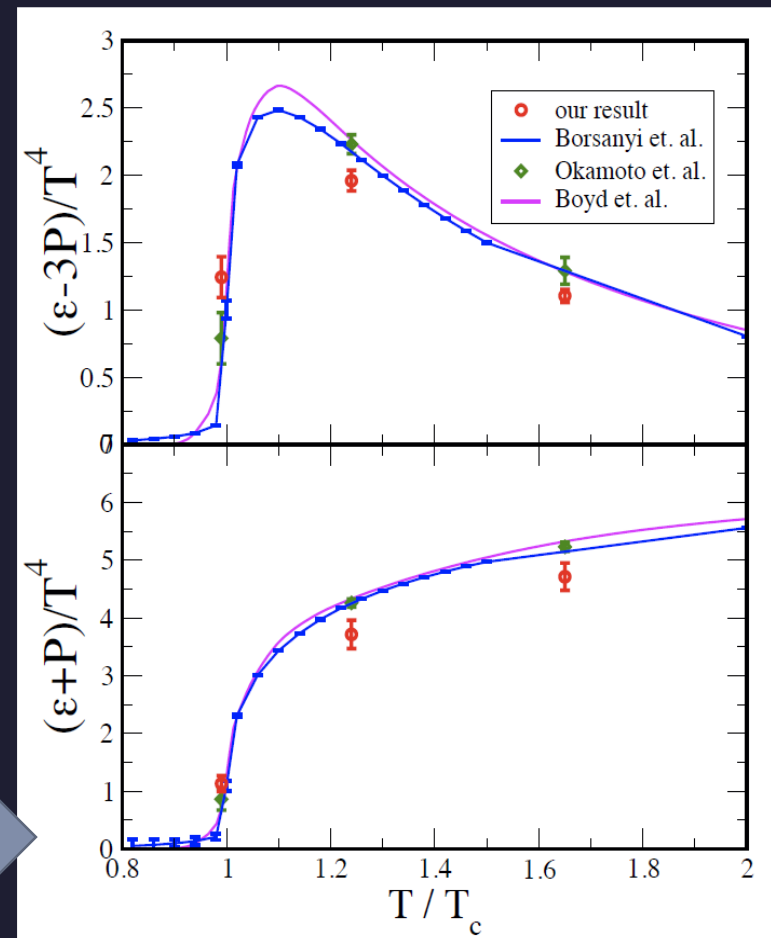
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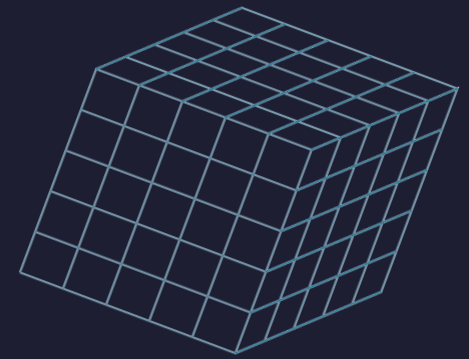
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Comparison with previous studies



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Simulation 2

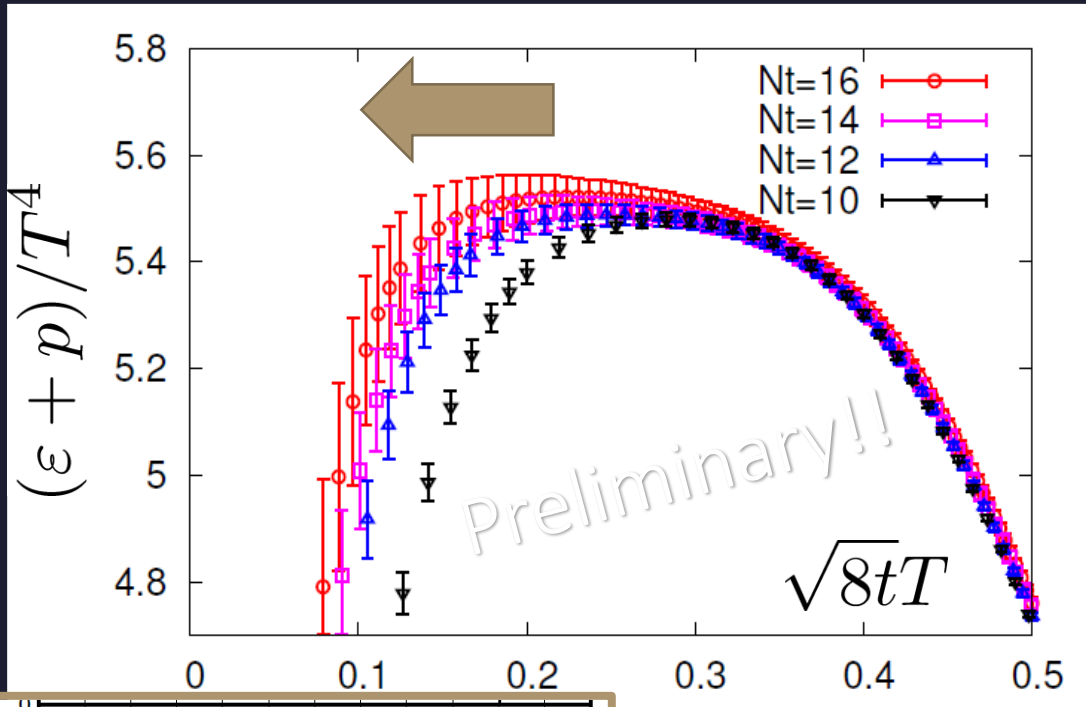
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efficiency $\sim 40\%$

twice finer lattice!

Entropy Density on Finer Lattices

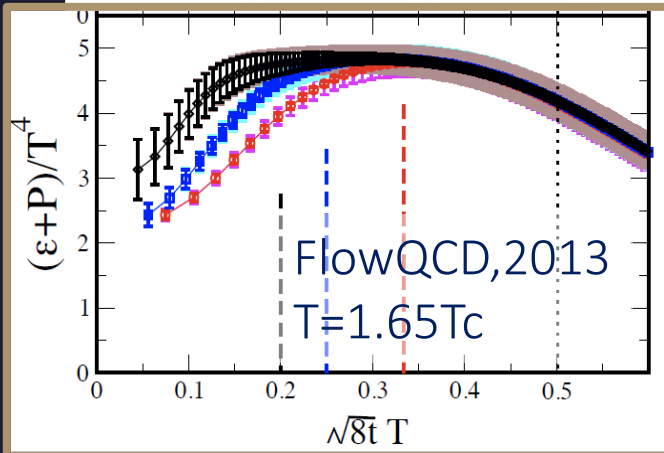


$T = 2.31T_c$

$64^3 \times N_t$

$N_t = 10, 12, 14, 16$

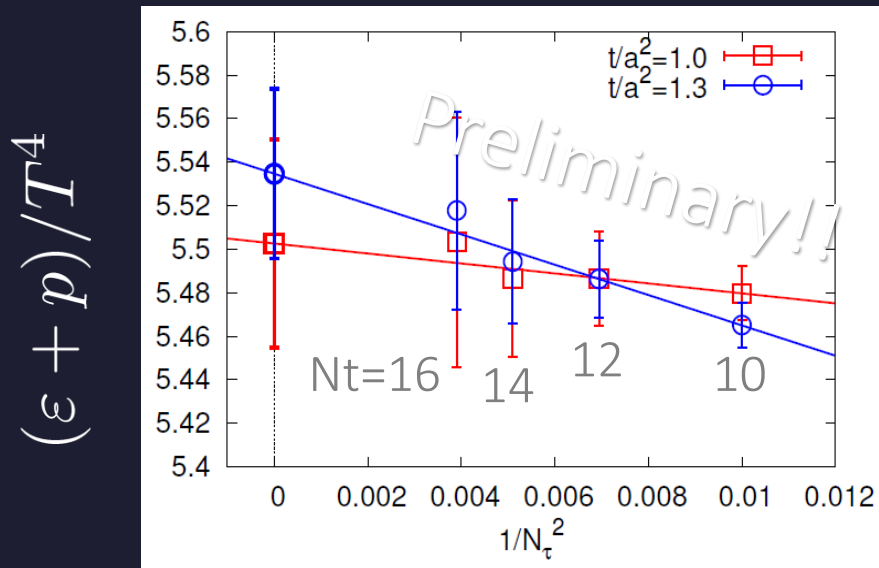
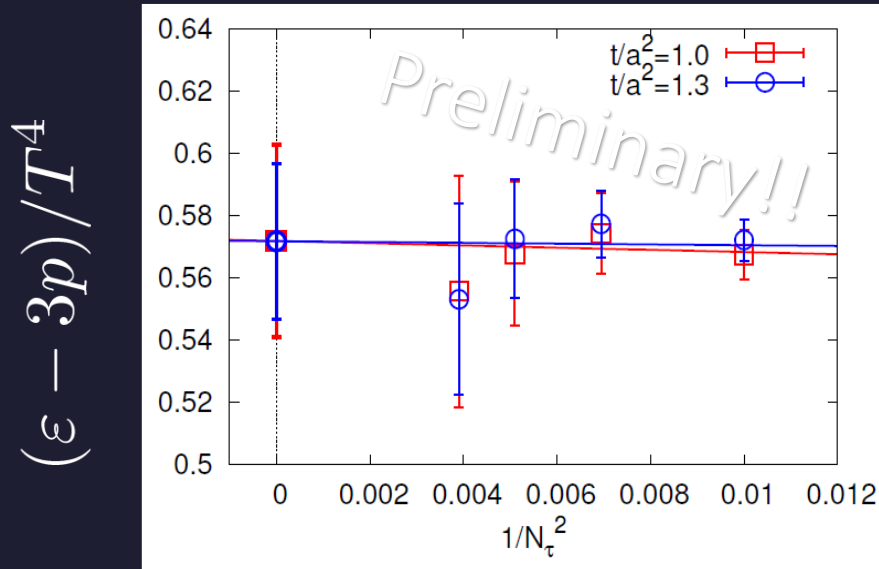
2000 confs.



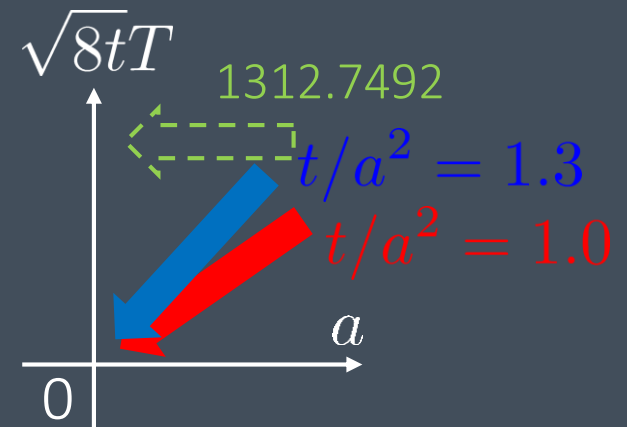
- The wider plateau on the finer lattices
- Plateau may have a nonzero slope

Continuum Extrapolation

- $T=2.31T_c$
- 2000 confs
- $Nt = 10 \sim 16$



$a \rightarrow 0$ limit with fixed t/a^2



Continuum extrapolation
is stable

Numerical Analysis: EMT Correlators

Fluctuations and Correlations

viscosity, specific heat, ...

$$\eta = \int_0^\infty dt \langle T_{12}; T_{12} \rangle$$

$$c_V \sim \langle \delta T_{00}^2 \rangle$$

EMT Correlator

□ Kubo Formula: T_{12} correlator \leftrightarrow shear viscosity

$$\eta = \int_0^\infty dt \int_0^{1/T} d\tau \int d^3x \langle T_{12}(x, -i\tau) T_{12}(0, t) \rangle$$

➤ Hydrodynamics describes long range behavior of $T_{\mu\nu}$

□ Energy fluctuation \leftrightarrow specific heat

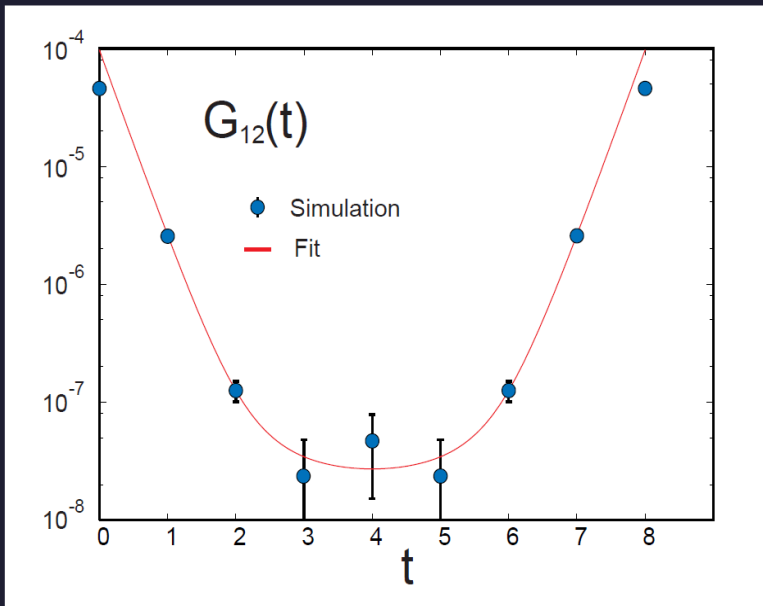
$$c_V = \frac{\langle \delta E^2 \rangle}{VT^2}$$

EMT Correlator : Noisy...

With naïve EMT operators

$$\langle T_{12}(\tau) T_{12}(0) \rangle$$

$$\langle T_{\mu\nu}(\tau) T_{\mu\nu}(0) \rangle$$

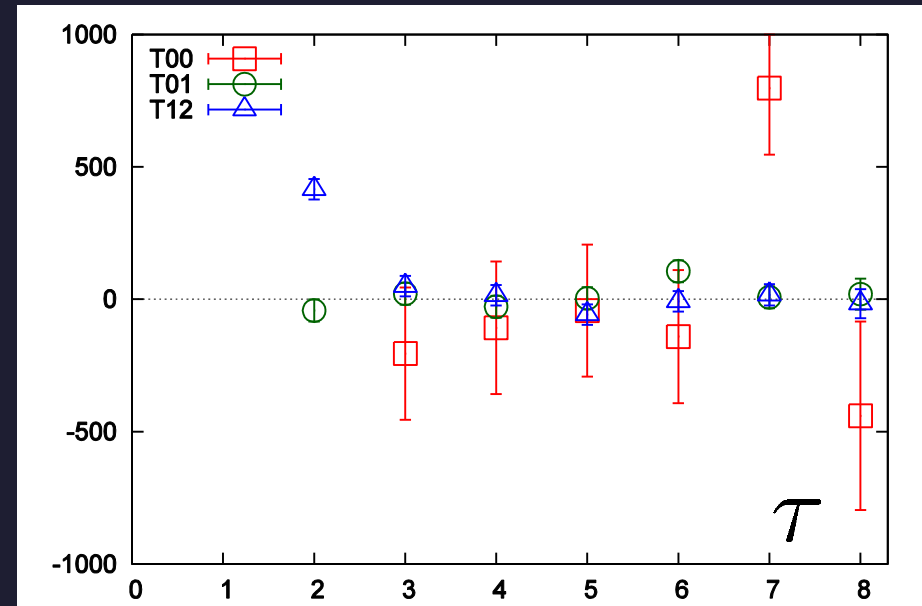


Nakamura, Sakai, PRL, 2005

$N_t=8$

improved action

$\sim 10^6$ configurations



$N_t=16$

standard action

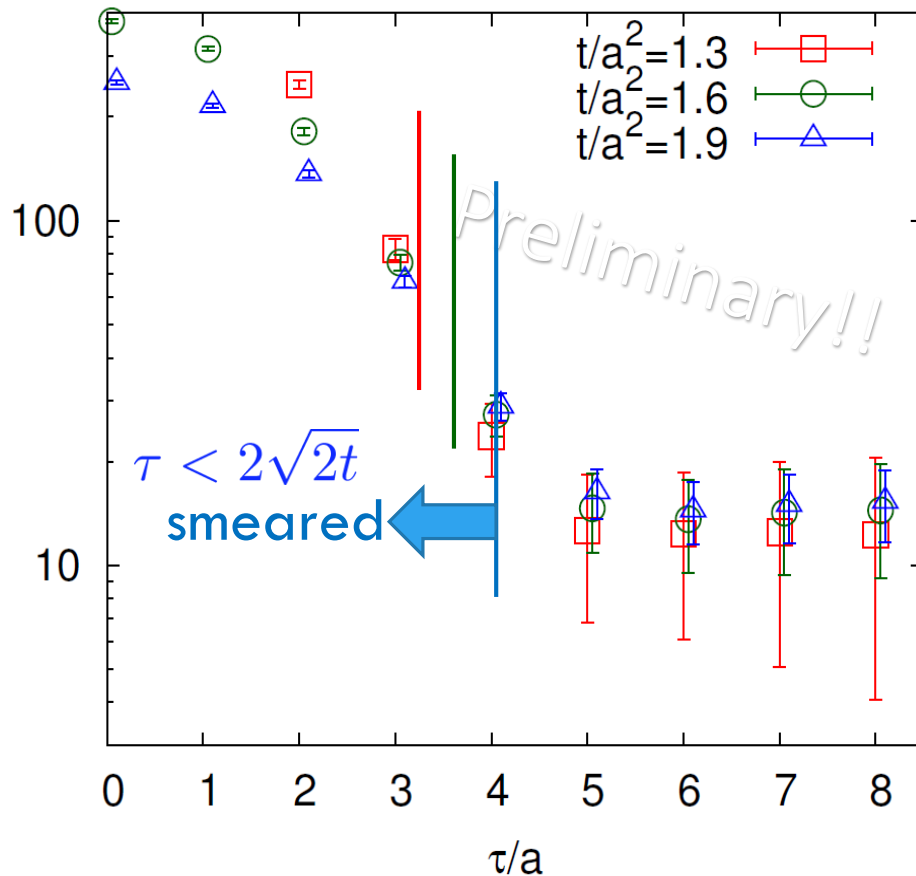
5×10^4 configurations

... no signal

Energy Correlation Function

$$\langle T_{00}(\tau)T_{00}(0) \rangle / T^5$$

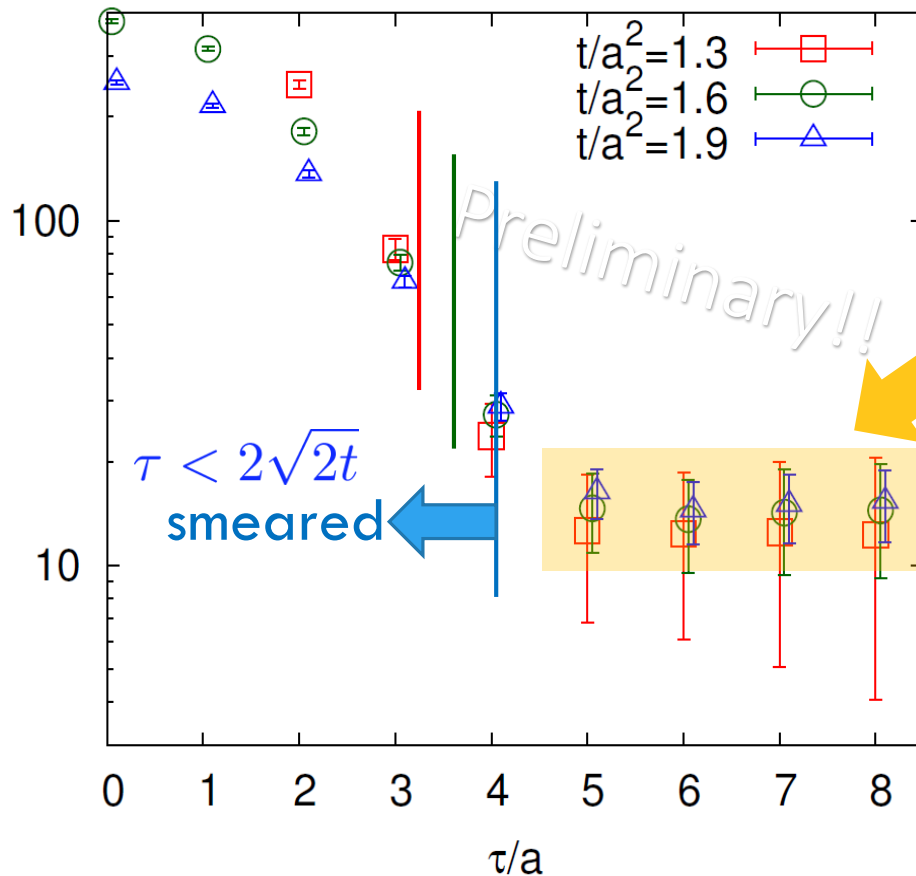
$T=2.31T_c$
 $b=7.2, Nt=16$
2000 confs
 $p=0$ correlator



Energy Correlation Function

$$\langle T_{00}(\tau)T_{00}(0) \rangle / T^5$$

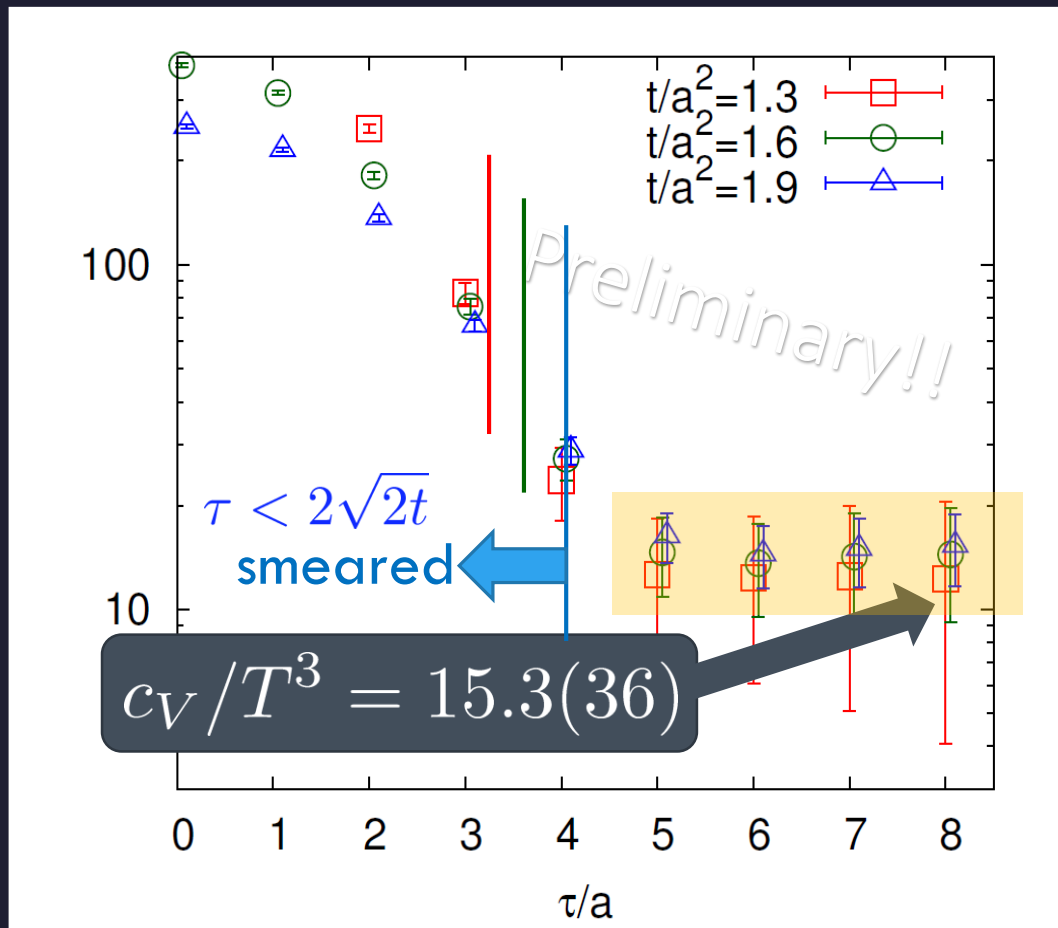
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Energy Correlation Function

$$\langle T_{00}(\tau) T_{00}(0) \rangle / T^5$$

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□ specific heat

$$c_V = \frac{\langle \delta E^2 \rangle}{VT^2}$$

→ Novel approach to
 measure specific heat!

Gavai, Gupta, Mukherjee, 2005

$$c_V/T^3 = 15(1) \quad T/T_c = 2$$

$$= 18(2) \quad T/T_c = 3$$

differential method / cont lim.

Summary

$$T_{\mu\nu}^R(x)$$

Summary

EMT formula from gradient flow

$$T_{\mu\nu}^R(x) = \lim_{t \rightarrow 0} \left[\frac{1}{\alpha_U(t)} U_{\mu\nu}(t, x) + \frac{\delta_{\mu\nu}}{4\alpha_E(t)} E(t, x)_{\text{subt.}} \right]$$

This formula can successfully define and calculate the EMT on the lattice

It provides us with novel approaches to measure various observables on the lattice!

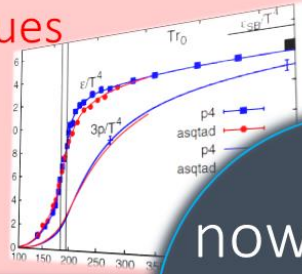
This method is direct, intuitive and less noisy

Many Future Studies!!

Thermodynamics

direct measurement of expectation values

$$\langle T_{00} \rangle, \langle T_{ii} \rangle$$



Fluctuations and Correlations

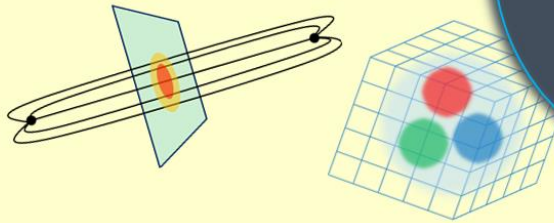
viscosity, specific heat, ...

$$c_V \sim \langle \delta T_{00}^2 \rangle$$

$$\eta = \langle T_{12}; T_{12} \rangle$$

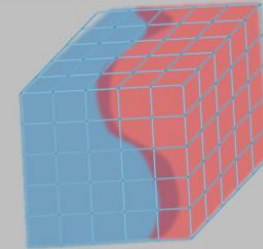
now we have

$$T_{\mu\nu}$$



- confinement string
- EM distribution in hadrons

Hadron Structure



- vacuum configuration
- mixed state on 1st transition

Vacuum Structure

Other observables

full QCD Makino, Suzuki, 2014

non-pert. improvement Patella 7E(Thu)

O(a) improvement

Nogradi, 7E(Thu); Sint, 7E(Thu)

Monahan, 7E(Thu)

and etc.

Correlation Function

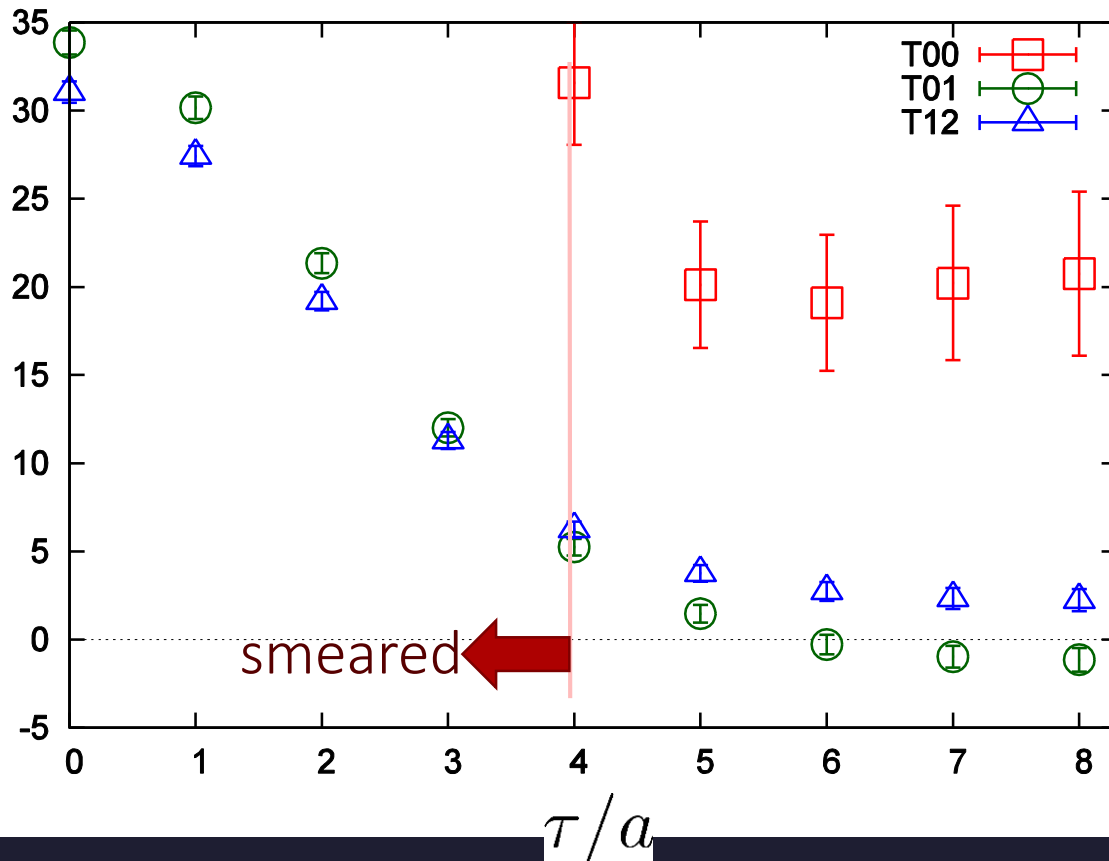
$$C_{\mu\nu}(\tau) = \int d^3x \langle T_{\mu\nu}(x, \tau) T_{\mu\nu}(0, 0) \rangle$$

$64^3 \times 16$

$\beta = 7.2$ ($T \sim 2.3 T_c$)

1200 confs

$t/a^2 = 1.9$



$C_{44}(\tau)$: constant
 ← conservation law!

$$\partial_\tau \langle \delta E(\tau) \delta E(0) \rangle = 0$$

(for $\tau \neq 0$)

$C_{12}(\tau)$

$C_{41}(\tau)$

negative ← $i^2 = -1$