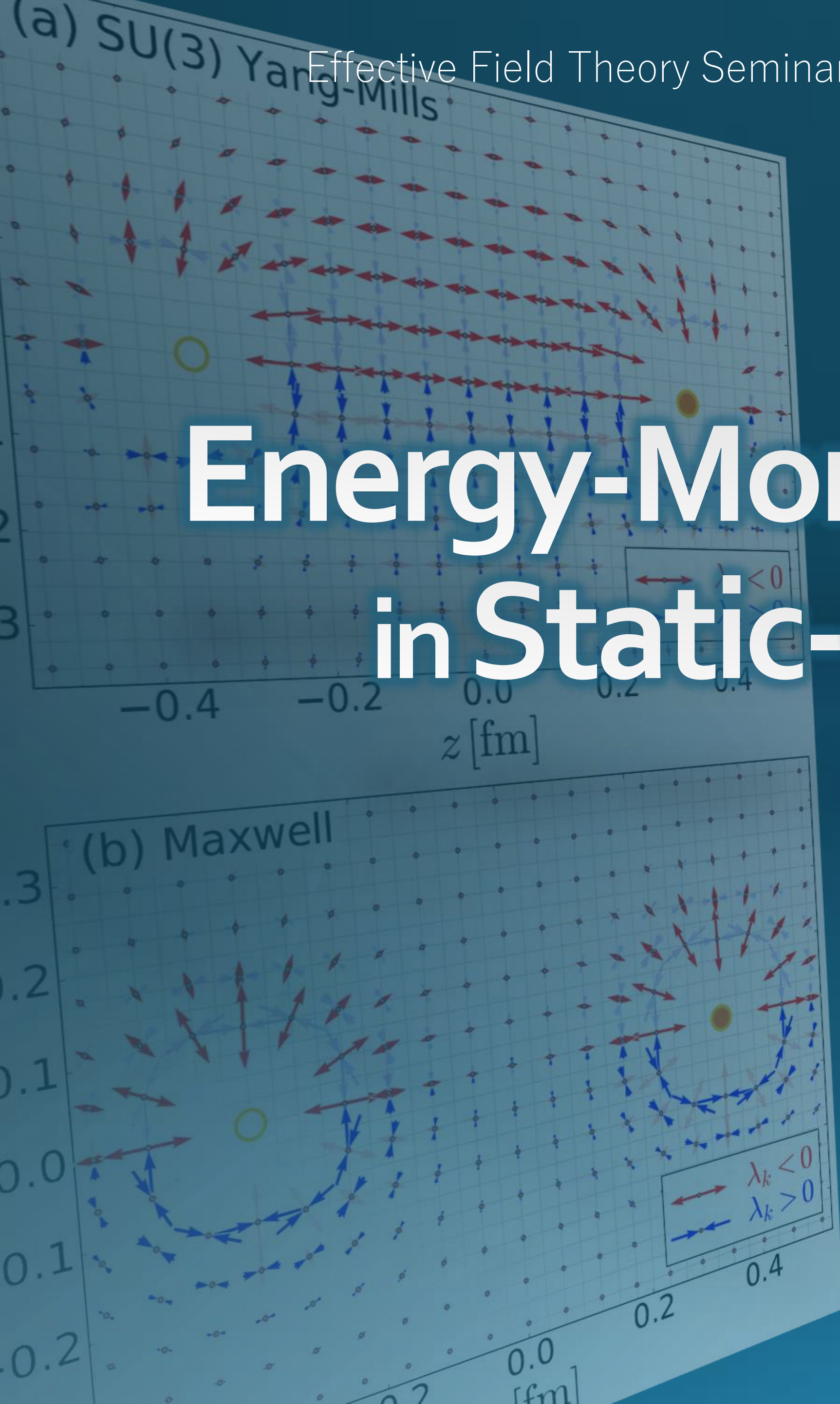
Effective Field Theory Seminar, TUM, München (Online), 2022/Jan./26



Distribution of Energy-Momentum Tensor in Static-Quark Systems

Masakiyo Kitazawa
(Osaka U.)

FlowQCD, PLB **789**, 210 (2019)
Yanagihara, MK, PTEP**2019**, 093B02 (2019)
FlowQCD, PRD **102**, 114522 (2020)

Energy-Momentum Tensor

energy — momentum

$$T_{\mu\nu} = \begin{bmatrix} T_{00} & T_{01} & T_{02} & T_{03} \\ T_{10} & T_{11} & T_{12} & T_{13} \\ T_{20} & T_{21} & T_{22} & T_{23} \\ T_{30} & T_{31} & T_{32} & T_{33} \end{bmatrix}$$

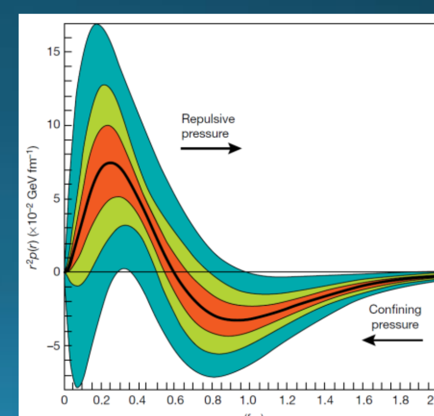
pressure

stress

- ☐ The most fundamental quantity in physics.
- ☐ All components are important quantities.

- ☐ How does EMT behave inside hadrons?

Pressure inside a proton
Nature, 557, 396 (2018)



Burkert+, Nature (2018)

Static-Quark Systems

- ☐ Fundamental probe to study field theories
- ☐ Numerical simulations on the lattice is straightforward!

☐ $Q\bar{Q}$

- Flux tube formation

☐ Single Q

- Heavy-light meson @ $T<T_c$
- Debye screening @ $T>T_c$

EMT in Static Q systems

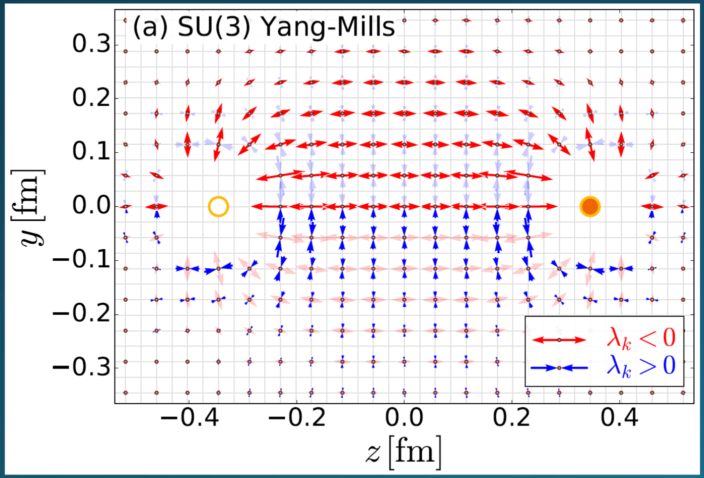
Combine 2 fundamental tools!

Static-Quark Systems

- ☐ Fundamental probe to study field theories
- ☐ Numerical simulations on the lattice is straightforward!

☐ $Q\bar{Q}$

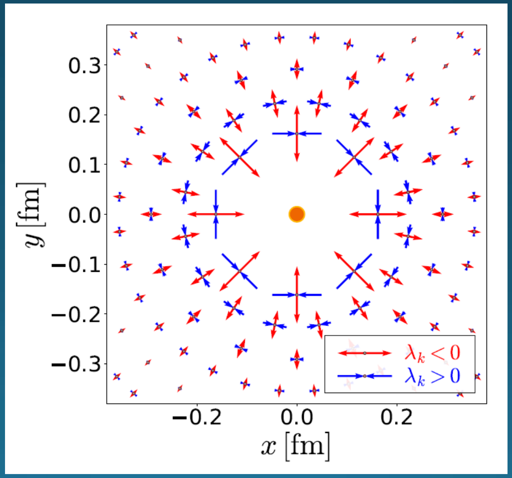
- Flux tube formation



FlowQCD, PLB **789**, 210 (2019)

☐ Single Q

- Heavy-light meson @ $T<T_c$
- Debye screening @ $T>T_c$

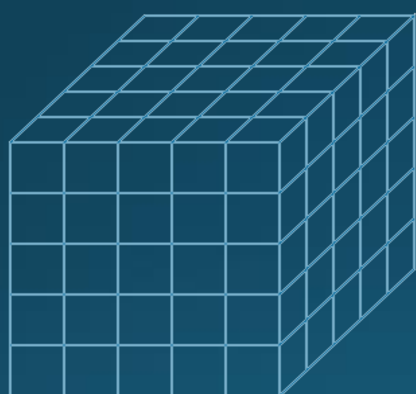


FlowQCD, PRD **102**, 114522 (2020)

$T_{\mu\nu}$: nontrivial observable on the lattice

① Definition of the operator is nontrivial

because of the explicit breaking of translational invariance



ex: $T_{\mu\nu} = F_{\mu\rho}F_{\nu\rho} - \frac{1}{4}\delta_{\mu\nu}FF$

$F_{\mu\nu} =$ 

② Its measurement is noisy

due to high dimensionality and etc.

SF*t*X Method and Gradient Flow

SF*t*X = Small Flow-*t*ime eXpansion

Yang-Mills Gradient Flow

Luscher 2010
Narayanan, Neuberger, 2006
Luscher, Weiss, 2011

$$\frac{\partial}{\partial t} A_\mu(t,x) = -\frac{\partial S_{\rm YM}}{\partial A_\mu}$$

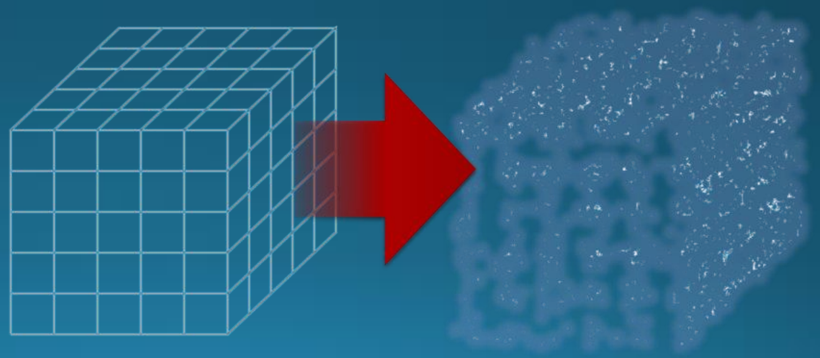
$$A_\mu(0,x) = A_\mu(x)$$

t: "flow time"
dim:[length²]

leading

$$\partial_t A_\mu = D_\nu G_{\mu\nu} = \partial_\nu \partial_\nu A_\mu + \cdots$$

- diffusion equation in 4-dim space
- diffusion distance $d \sim \sqrt{8t}$
- "continuous" cooling/smearing
- No UV divergence at $t > 0$



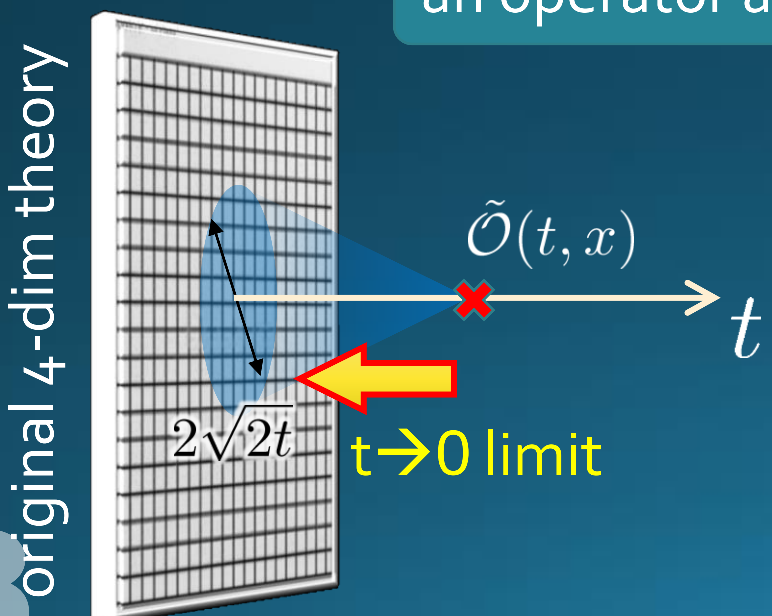
Small Flow-Time Expansion

Luescher, Weisz, 2011
Suzuki, 2013

$$\tilde{\mathcal{O}}(t,x) \xrightarrow[t\to 0]{} \sum_i c_i(t)\mathcal{O}_i^R(x)$$

an operator at t>0

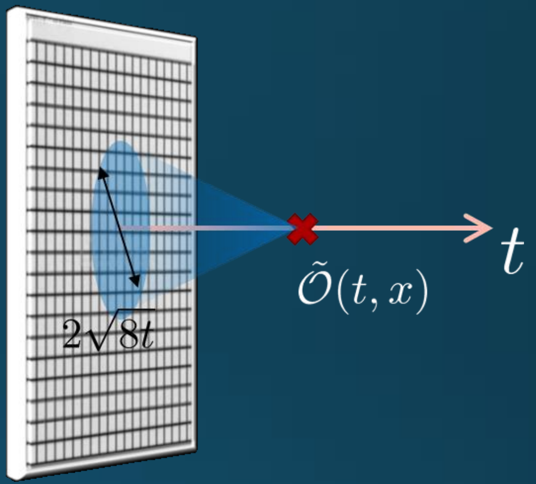
remormalized operators of original theory



Constructing EMT 1

Suzuki, 2013

$$\tilde{O}(t,x) \underset{t\to 0}{\longrightarrow} \sum_i c_i(t)\mathcal{O}_i^R(x)$$



□ **Gauge-invariant dimension 4 operators**

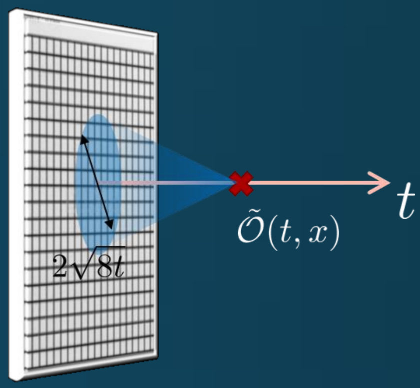
$$\begin{cases} U_{\mu\nu}(t,x) = G_{\mu\rho}(t,x)G_{\nu\rho}(t,x) - \frac{1}{4}\delta_{\mu\nu}G_{\rho\sigma}(t,x)G_{\rho\sigma}(t,x) \\ E(t,x) = \frac{1}{4}G_{\rho\sigma}(t,x)G_{\rho\sigma}(t,x) \end{cases}$$

Constructing EMT

Suzuki, 2013

$$U_{\mu\nu}(t,x) = \alpha_U(t)\left[T^R_{\mu\nu}(x) - \frac{1}{4}\delta_{\mu\nu}T^R_{\rho\rho}(x)\right] + \mathcal{O}(t)$$

$$E(t,x) = \langle E(t,x)\rangle + \alpha_E(t) T^R_{\rho\rho}(x) + \mathcal{O}(t)$$



vacuum subtr.

Remormalized EMT

$$T^R_{\mu\nu}(x) = \lim_{t\to 0}\left[c_1(t)U_{\mu\nu}(t,x) + \delta_{\mu\nu}c_2(t)E(t,x)_{\rm subt.}\right]$$

"SF*t*X method" (Small Flow *t*ime eXpansion)

Perturbative Coefficients

$$T_{\mu\nu}(t) = c_1(t) U_{\mu\nu}(t) + \delta_{\mu\nu} c_2(t) E(t)$$

	LO	1-loop	2-loop	3-loop
$c_1(t)$	○	○	○	
$c_2(t)$	× zero	○	○	○

Suzuki (2013) Harlander+(2018)

Iritani, MK, Suzuki, Takaura, 2019

Suzuki, PTEP 2013, 083B03
Harlander+, 1808.09837
Iritani, MK, Suzuki, Takaura, PTEP 2019

□ Choice of the scale of g²

$$c_1(t) = c_1\left(g^2(\mu(t))\right)$$

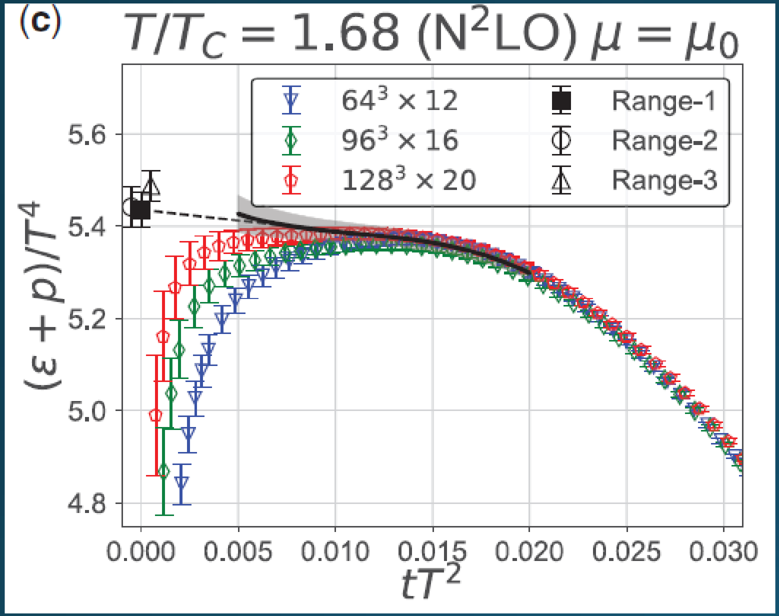
Previous: $\mu_d(t) = 1/\sqrt{8t}$

Improved: $\mu_0(t) = 1/\sqrt{2e^{\gamma_E} t}$

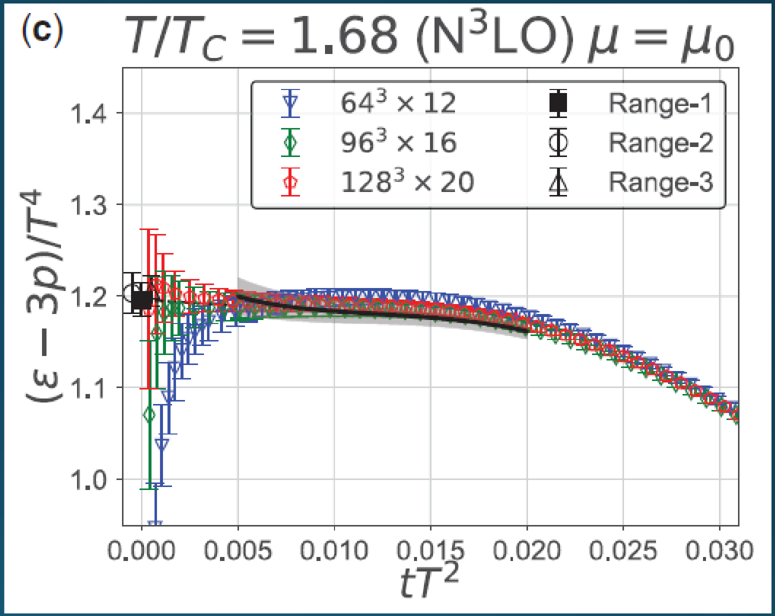
Harlander+ (2018)

t Dependence

$\langle T_{00}\rangle + \langle T_{11}\rangle$



$\langle T_{00}\rangle - 3\langle T_{11}\rangle$



Iritani, MK, Suzuki, Takaura, PTEP 2019

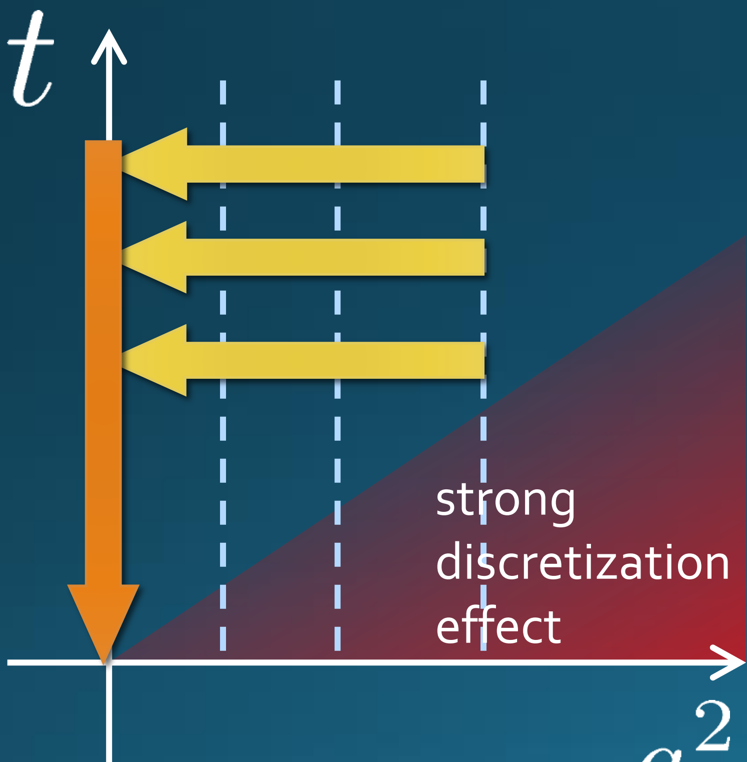
- ☐ Existence of "linear window" at intermediate t

Double Extrapolation

t→0, a→0

$$\langle T_{\mu\nu}(t)\rangle_{\rm latt} = \langle T_{\mu\nu}(t)\rangle_{\rm phys} + C_{\mu\nu}t + D_{\mu\nu}(t)\frac{a^2}{t}$$

O(t) terms in SFTE lattice discretization



Continuum extrapolation

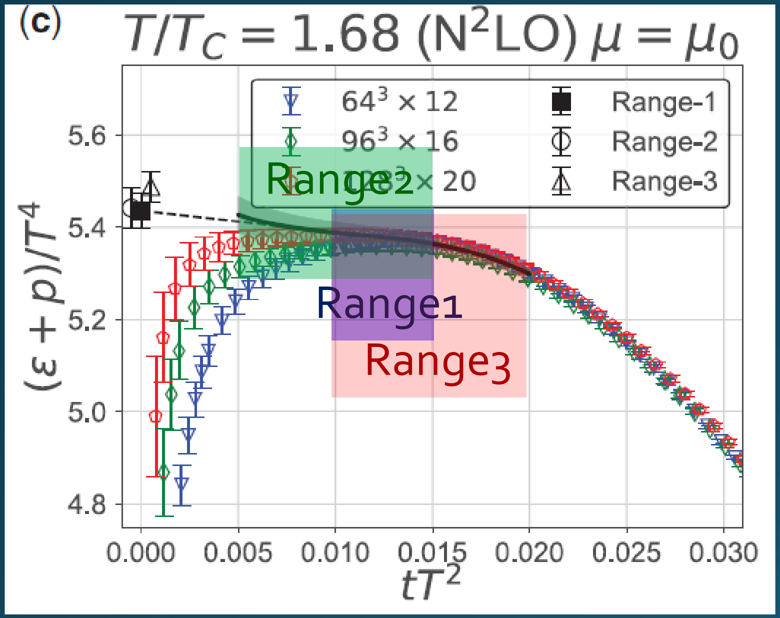
$$\langle T_{\mu\nu}(t)\rangle_{\rm cont} = \langle T_{\mu\nu}(t)\rangle_{\rm lat} + C(t)a^2$$

Small t extrapolation

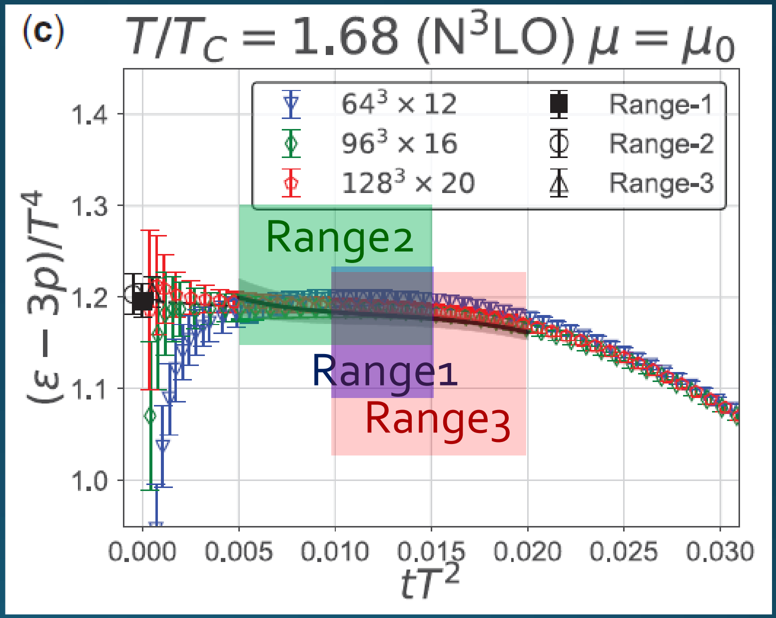
$$\langle T_{\mu\nu}\rangle = \langle T_{\mu\nu}(t)\rangle + C't$$

Thermodynamics: ε+p & ε-3p

$\langle T_{00}\rangle + \langle T_{11}\rangle$



$\langle T_{00}\rangle - 3\langle T_{11}\rangle$

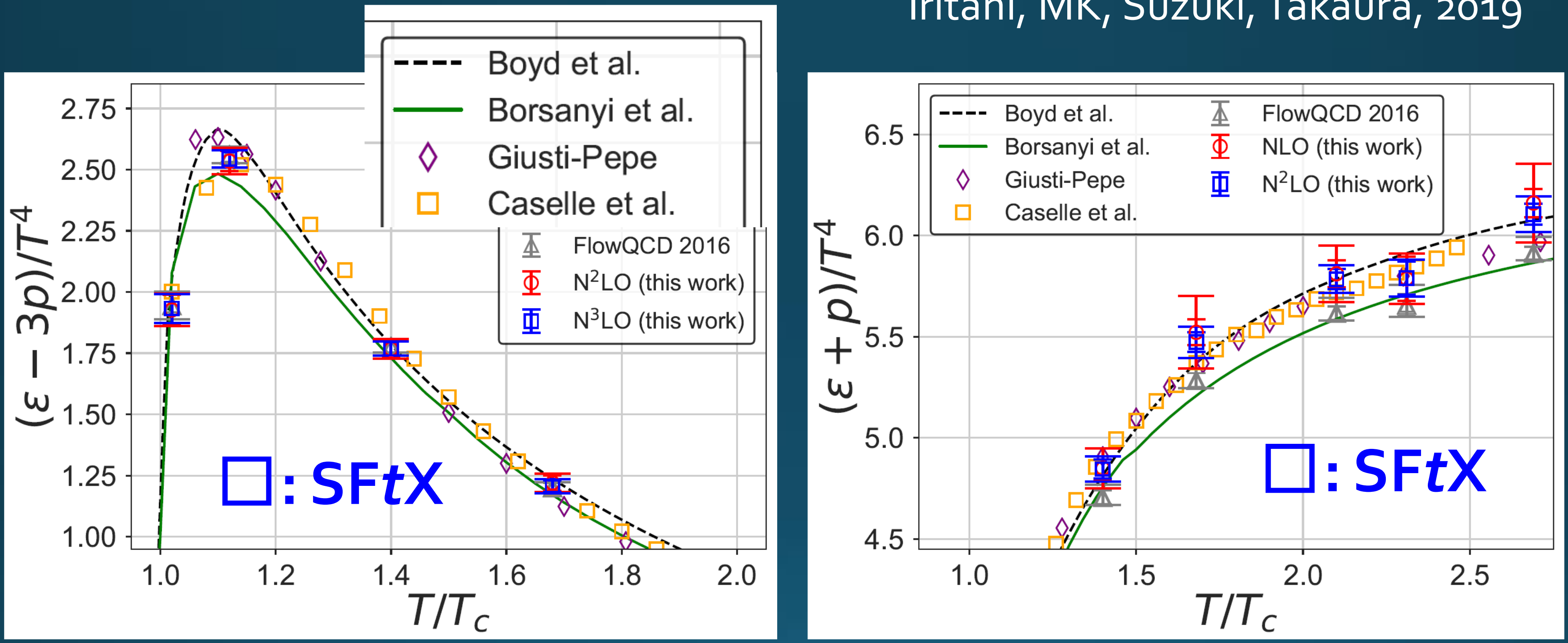


Iritani, MK, Suzuki, Takaura, PTEP 2019

- Existence of "linear window" at intermediate t
- Stable t⟶0 extrapolation
- Systematic errors: fit range, uncertaintyof Λ ($\pm$3%), …

Thermodynamics: $\varepsilon = \langle T_{00} \rangle,\ p = \langle T_{11} \rangle$

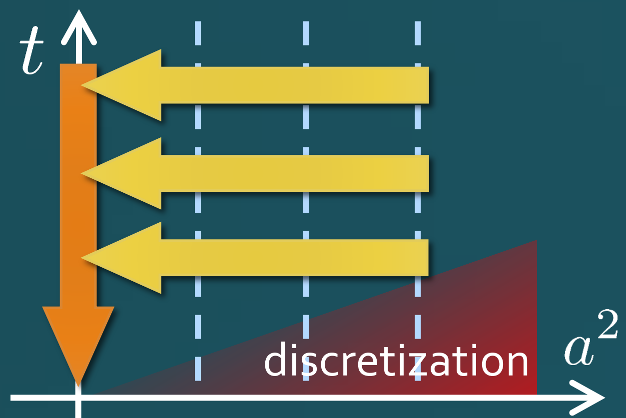
Iritani, MK, Suzuki, Takaura, 2019



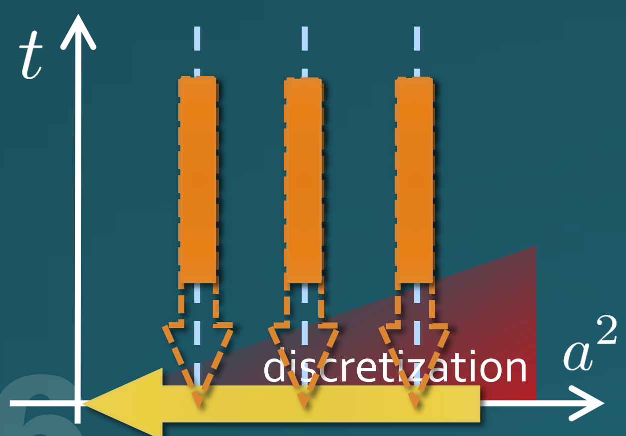
- ☐ Agreement with other methods within 1% level!
- ☐ Smaller statistics thanks to smearing by the flow

Alternative Extrapolation

Method A: $a \to 0,\ t \to 0$

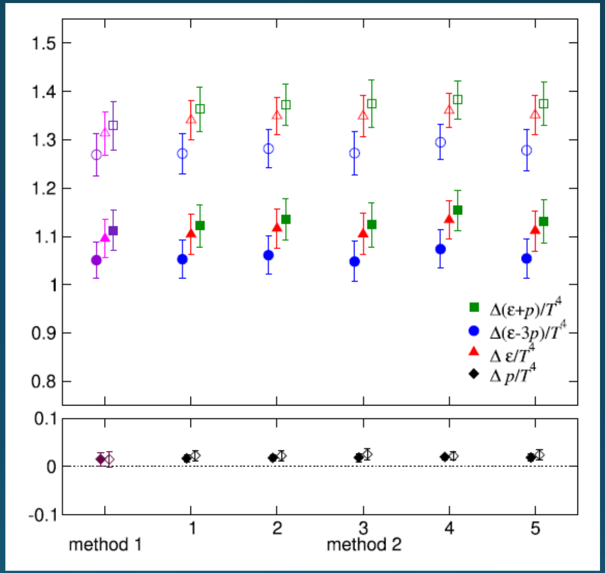


Method B: $t \to 0,\ a \to 0$



□ Consistency check

- □ latent heat & pressure gap



Shirogane+, 2021 (WHOT-QCD)

B | Method A

Consistent result for two methods

Gradient Flow for Fermions

$$\partial_t \psi(t,x) = D_\mu D_\mu \psi(t,x)$$

$$\partial_t \bar{\psi}(t,x) = \psi(t,x) \overleftarrow{D}_\mu \overleftarrow{D}_\mu$$

$$D_\mu = \partial_\mu + A_\mu(t,x)$$

Luscher, 2013
Makino, Suzuki, 2014
Taniguchi+ (WHOT) 2016~

- Not "gradient" flow but a "diffusion"-type equation.

- Divergence in field renormalization of fermions.
- All observables are finite at t>0 once Z(t) is fixed.

$$\tilde{\psi}(t,x) = Z(t)\psi(t,x)$$

- Energy-momentum tensor from SF*t*X Makino, Suzuki, 2014

EMT in QCD

Makino, Suzuki (2014)
Harlander+ (2018)

$$T_{\mu\nu}(t,x) = c_1(t)U_{\mu\nu}(t,x) + c_2(t)\delta_{\mu\nu}\big(E(t,x) - \langle E\rangle_0\big)$$
$$+ c_3(t)\big(O_{3\mu\nu}(t,x) - 2O_{4\mu\nu}(t,x) - \mathrm{VEV}\big)$$
$$+ c_4(t)\big(O_{4\mu\nu}(t,x) - \mathrm{VEV}\big) + c_5(t)\big(O_{5\mu\nu}(t,x) - \mathrm{VEV}\big)$$

$$T_{\mu\nu}(x) = \lim_{t\to 0} T_{\mu\nu}(t,x)$$

$$\tilde{O}^f_{3\mu\nu}(t,x) \equiv \varphi_f(t)\bar{\chi}_f(t,x)\left(\gamma_\mu \overleftrightarrow{D}_\nu + \gamma_\nu \overleftrightarrow{D}_\mu\right)\chi_f(t,x),$$
$$\tilde{O}^f_{4\mu\nu}(t,x) \equiv \varphi_f(t)\delta_{\mu\nu}\bar{\chi}_f(t,x)\overleftrightarrow{\not{D}}\chi_f(t,x),$$
$$\tilde{O}^f_{5\mu\nu}(t,x) \equiv \varphi_f(t)\delta_{\mu\nu}\bar{\chi}_f(t,x)\chi_f(t,x),$$

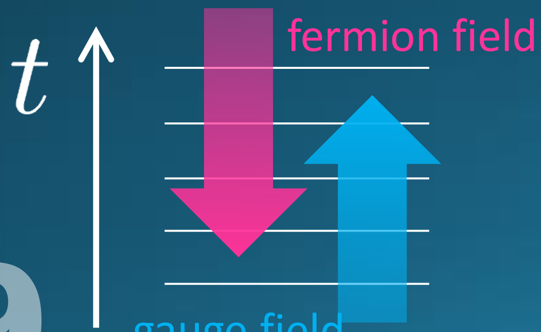
$$\varphi_f(t) \equiv \frac{-6}{(4\pi)^2 t^2 \left\langle \bar{\chi}_f(t,x)\overleftrightarrow{\not{D}}\chi_f(t,x)\right\rangle_0}.$$

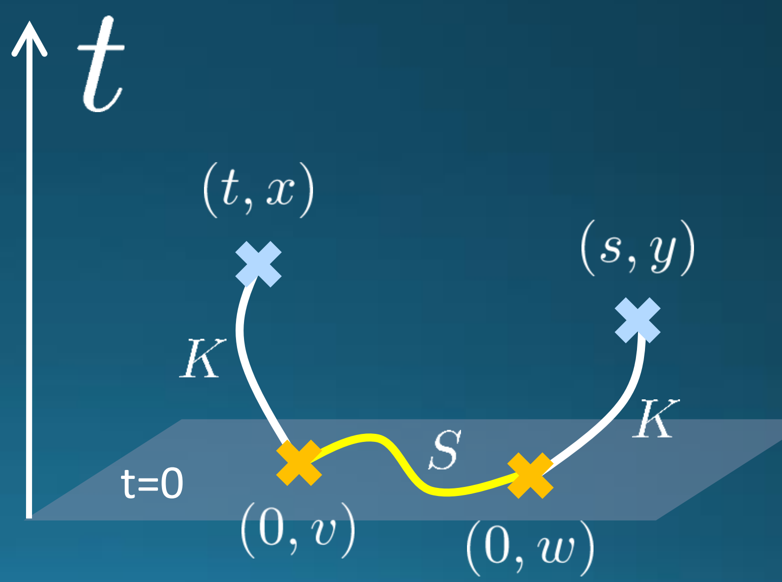
Fermion Propagator

$$S(t,x;s,y) = \langle \chi(t,x)\bar{\chi}(s,y)\rangle$$
$$= \sum_{v,w} K(t,x;0,v)S(v,w)K(s,y;0,w)^{\dagger}$$

$$\Big(\partial_t - D_\mu D_\mu\Big)K(t,x) = 0$$

- propagator of flow equation
- Inverse propagator is needed

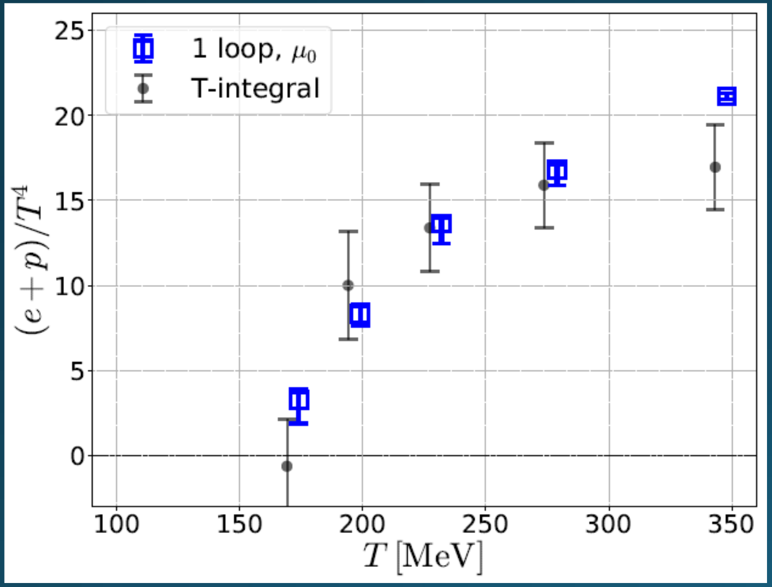




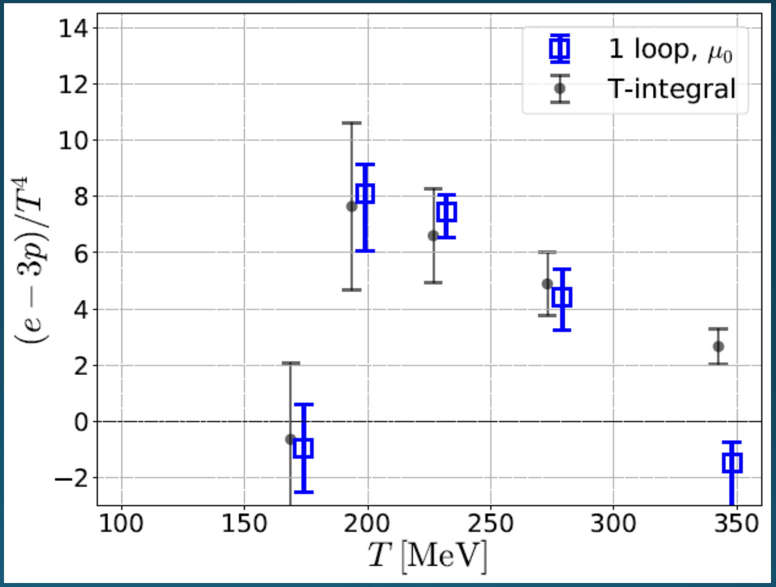
2+1 QCD EoS from Gradient Flow

WHOT-QCD, PR**D96** (2017); PR**D102** (2020)

$$\langle T_{00}\rangle + \langle T_{11}\rangle$$



$$\langle T_{00}\rangle - 3\langle T_{11}\rangle$$



- ☐ Agreement with integral method
- ☐ Substantial suppression of statistical errors

$m_{PS}/m_V \approx 0.63$

Physical mass: Kanaya+ (WHOT-QCD), 1910.13036

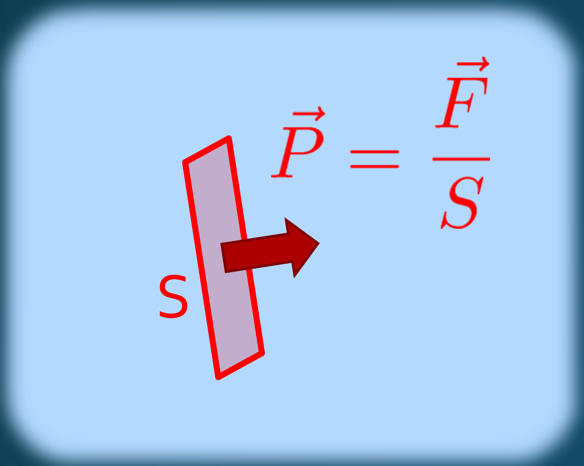
Flux-Tube Formation in $Q\bar{Q}$ System

FlowQCD, PLB **789**, 210 (2019)
Yanagihara, MK, PTEP**2019**, 093B02 (2019)

Stress = Force per Unit Area

Stress = Force per Unit Area

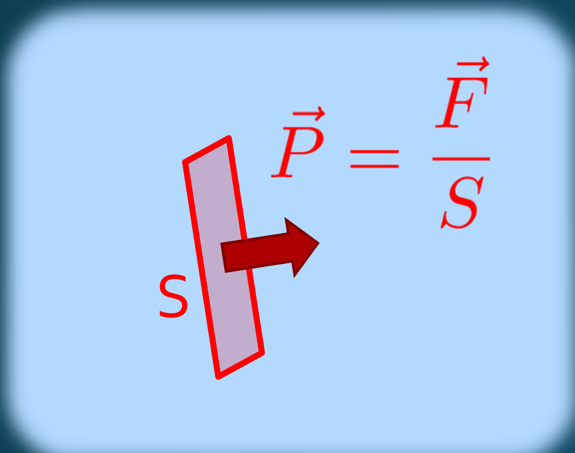
Pressure



$$\vec{P} = P\vec{n}$$

Stress = Force per Unit Area

Pressure

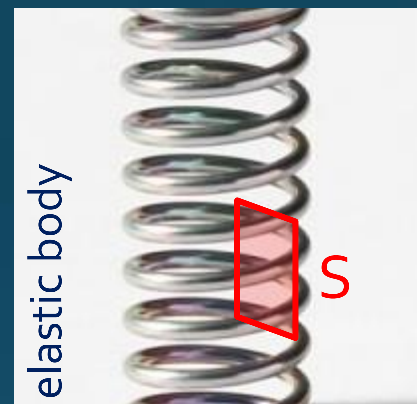


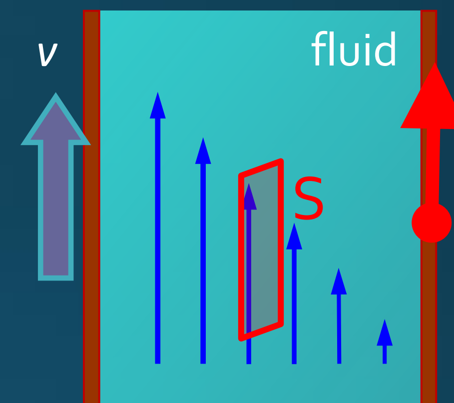
$$\vec{P} = P\vec{n}$$

In thermal medium

$$T_{ij} = P\delta_{ij}$$

Generally, F and n are not parallel





$$\frac{F_i}{S} = \sigma_{ij} n_j$$

Stress Tensor

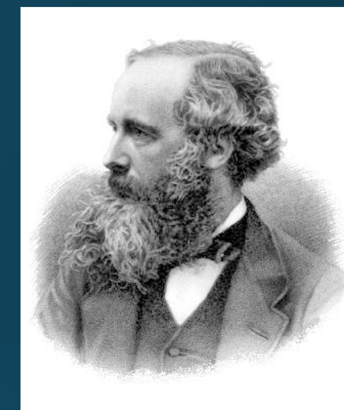
$$\sigma_{ij} = -T_{ij}$$

Landau
Lifshitz

Maxwell Stress

(in Maxwell Theory)

$$\sigma_{ij} = \varepsilon_0 E_i E_j + \frac{1}{\mu_0} B_i B_j - \frac{1}{2}\delta_{ij}\left(\varepsilon_0 E^2 + \frac{1}{\mu_0} B^2\right)$$

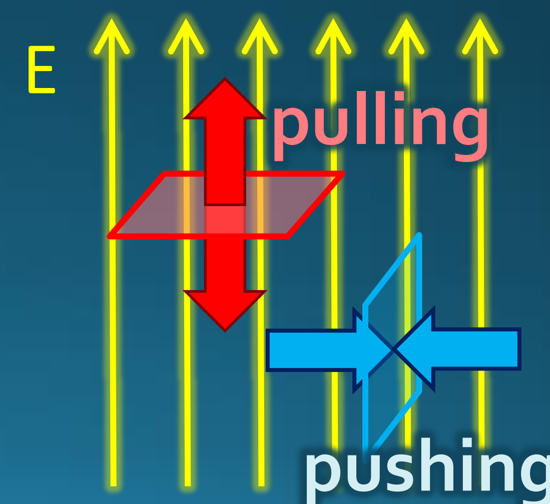


Maxwell

$$\vec{E} = (E, 0, 0)$$

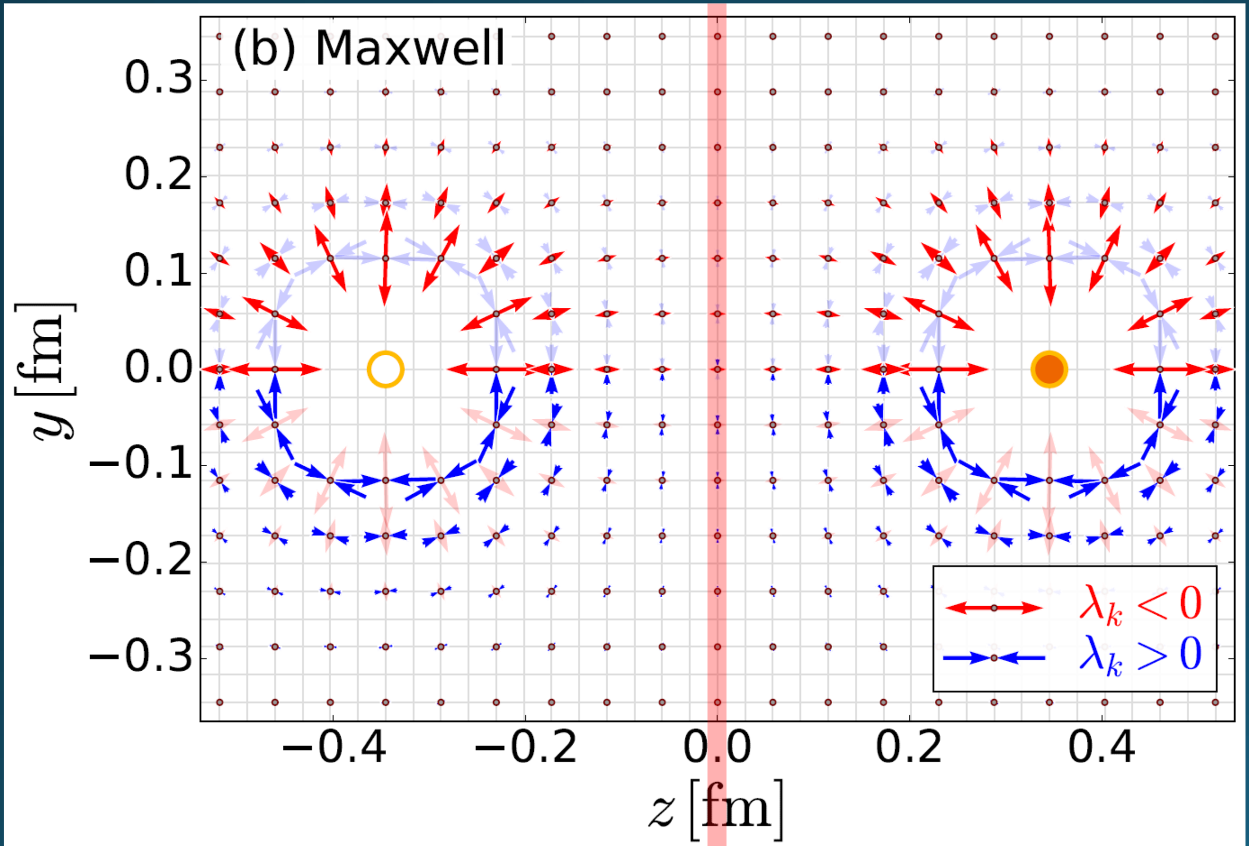
$$T_{ij} = \begin{pmatrix} -E^2 & 0 & 0 \\ 0 & E^2 & 0 \\ 0 & 0 & E^2 \end{pmatrix}$$

- Parallel to field: **Pulling**
- Vertical to field: **Pushing**



Maxwell Stress

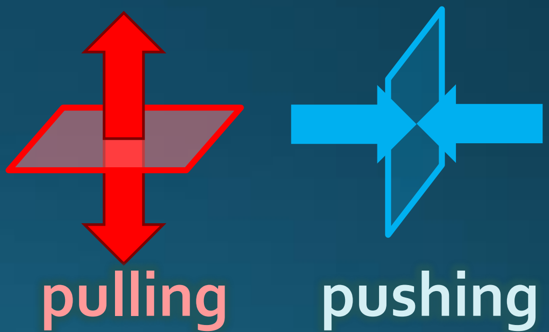
(in Maxwell Theory)



$$T_{ij} v_j^{(k)} = \lambda_k v_i^{(k)}$$

$$(k = 1, 2, 3)$$

length: $\sqrt{|\lambda_k|}$

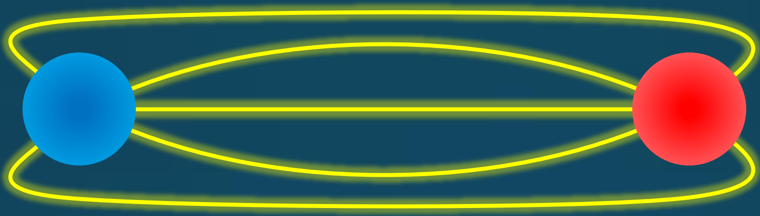


Definite physical meaning

- ☐ Distortion of field, line of the field
- ☐ Propagation of the force as local interaction

Quark-Anti-quark System

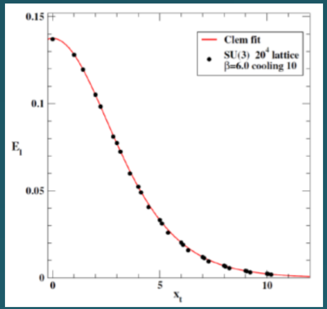
Formation of the flux tube → confinement



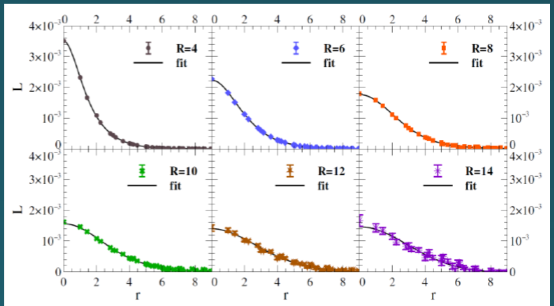
Previous Studies on Flux Tube

- ☐ Potential
- ☐ Action density
- ☐ Color-electric field

so many studies…



Cea+ (2012)



Cardoso+ (2013)

Lattice Setup

FlowQCD, PLB (2019)

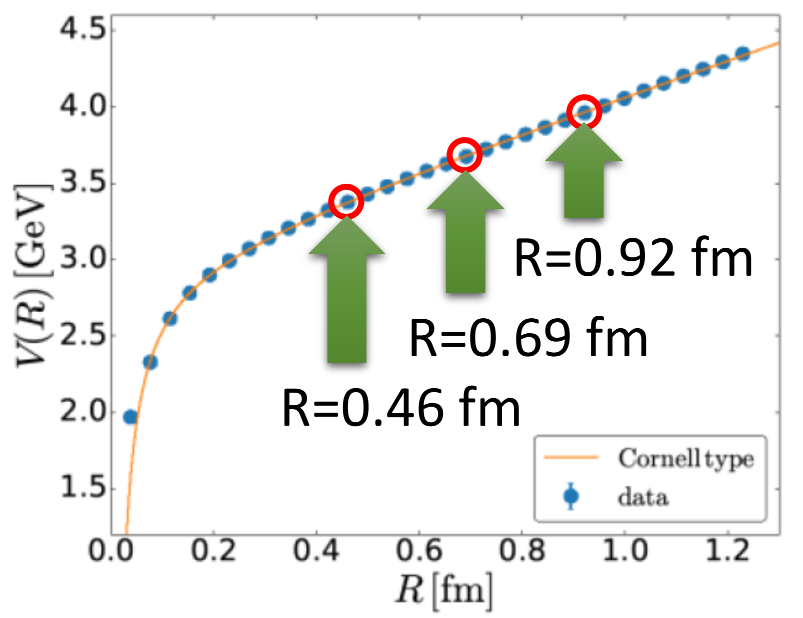
- SU(3) Yang-Mills (Quenched)
- Wilson gauge action
- Clover operator

- EMT around Wilson Loop
- APE smearing / multi-hit

- fine lattices (a=0.029-0.06 fm)
- **continuum extrapolation**

- Simulation: bluegene/Q@KEK

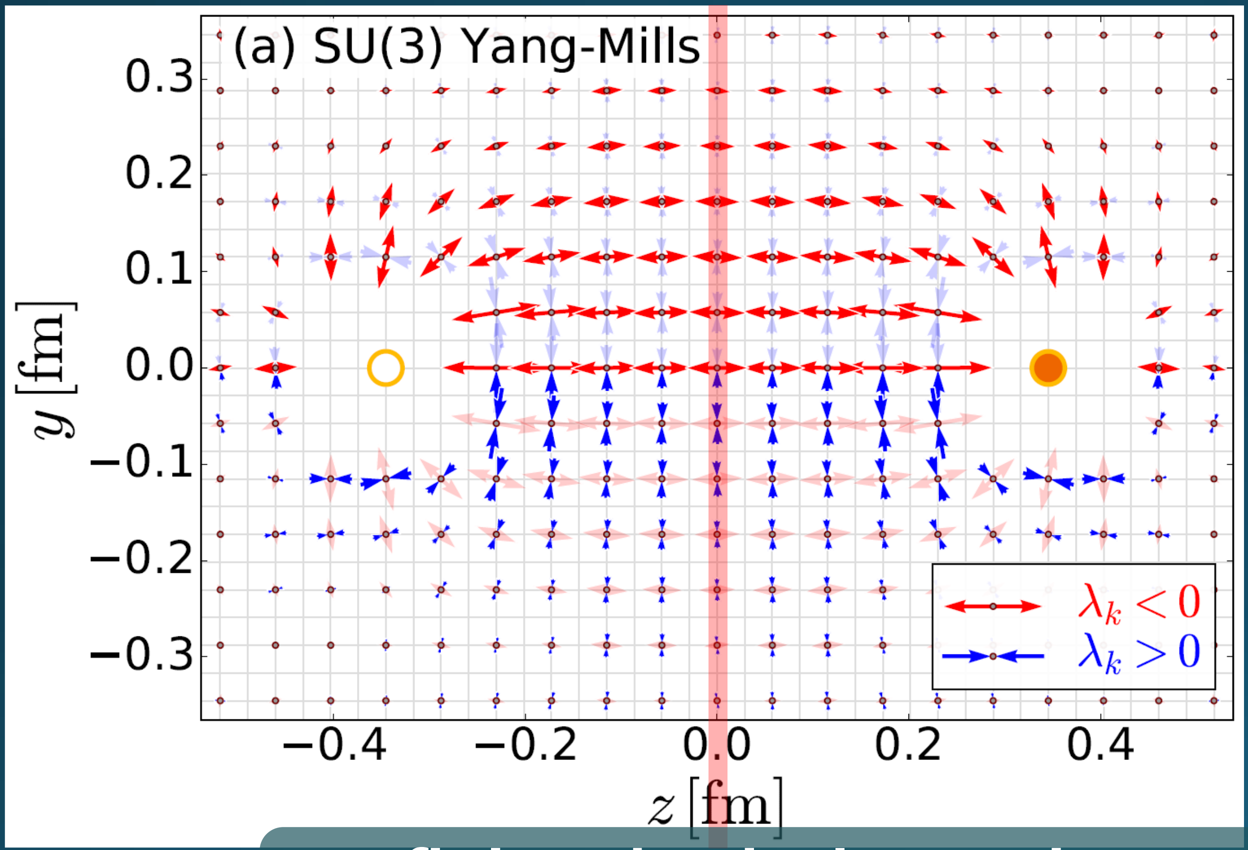
β	a [fm]	$N_{\rm size}^4$	$N_{\rm conf}$	R/a		
6.304	0.058	48^4	140	8	12	16
6.465	0.046	48^4	440	10	–	20
6.513	0.043	48^4	600	–	16	–
6.600	0.038	48^4	1,500	12	18	24
6.819	0.029	64^4	1,000	16	24	32
			R [fm]	0.46	0.69	0.92



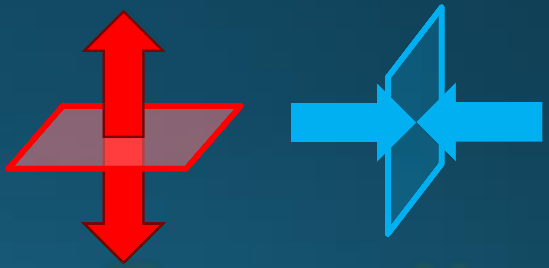
$$\langle O(x)\rangle_{\rm Q\bar{Q}} = \lim_{T\to\infty} \frac{\langle \delta O(x)\delta W(R,T)\rangle}{\langle W(R,T)\rangle}$$

Stress Tensor in $Q\bar{Q}$ System

FlowQCD, PLB (2019)



Lattice simulation
SU(3) Yang-Mills
a=0.029 fm
R=0.69 fm
t/a^2=2.0

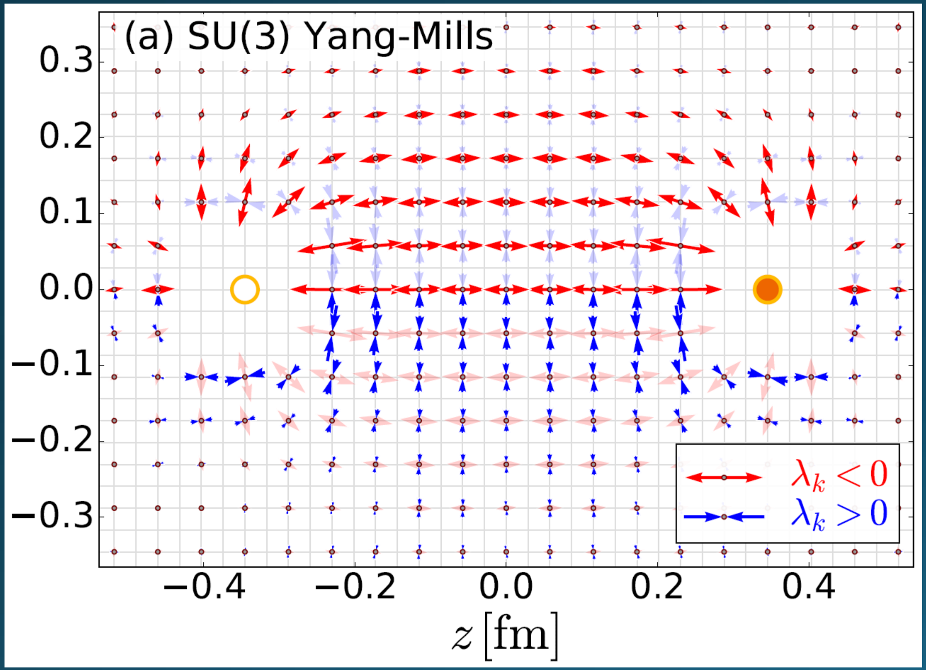


Definite physical meaning

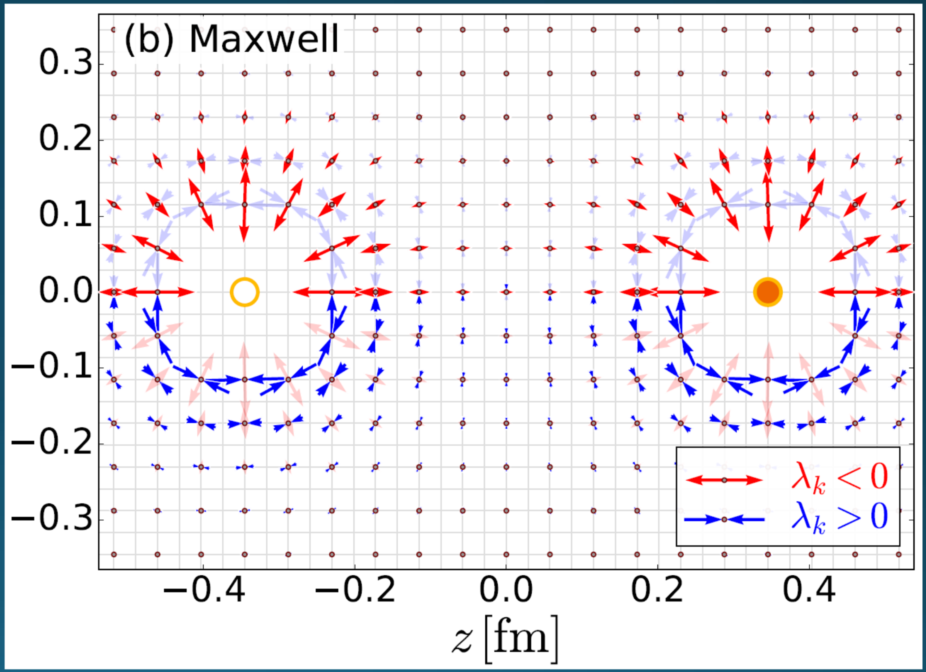
- ☐ Distortion of field, line of field
- ☐ Propagation of the force as local interaction
- ☐ Manifestly gauge invariant

SU(3) YM vs Maxwell

SU(3) Yang-Mills
(quantum)



Maxwell
(classical)



Propagation of the force is clearly different in YM and Maxwell theories!

Stress Distribution on Mid-Plane

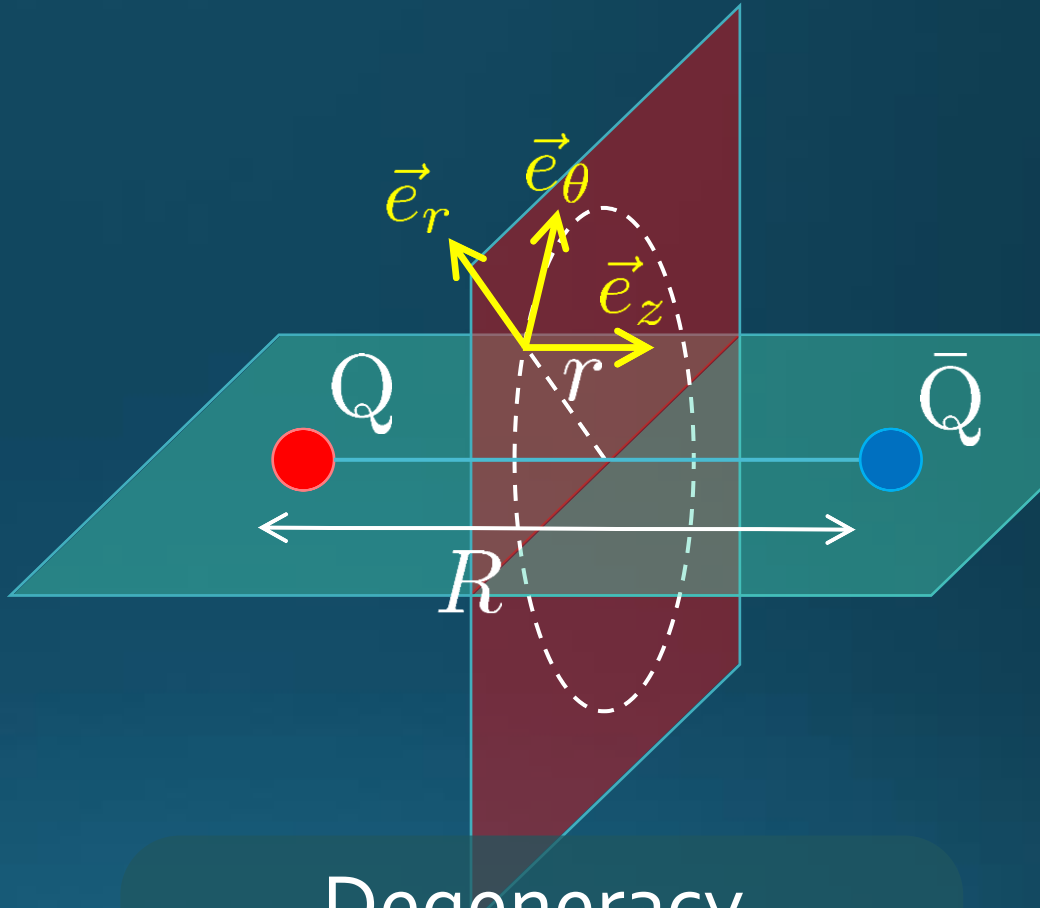
From rotational symm. & parity

EMT is diagonalized
in Cylindrical Coordinates

$$T_{cc'}(r) = \begin{pmatrix} T_{rr} & & & \\ & T_{\theta\theta} & & \\ & & T_{zz} & \\ & & & T_{44} \end{pmatrix}$$

$$T_{rr} = \vec{e}_r^{\,T} T \vec{e}_r$$

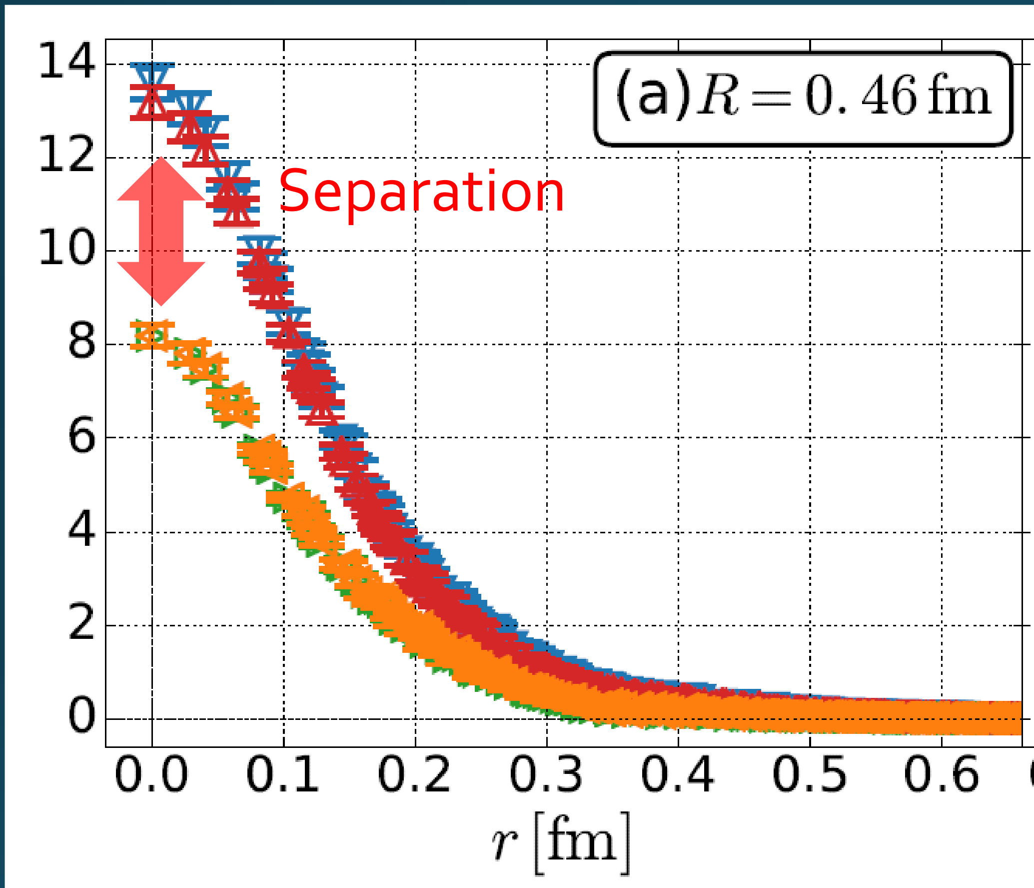
$$T_{\theta\theta} = \vec{e}_\theta^{\,T} T \vec{e}_\theta$$

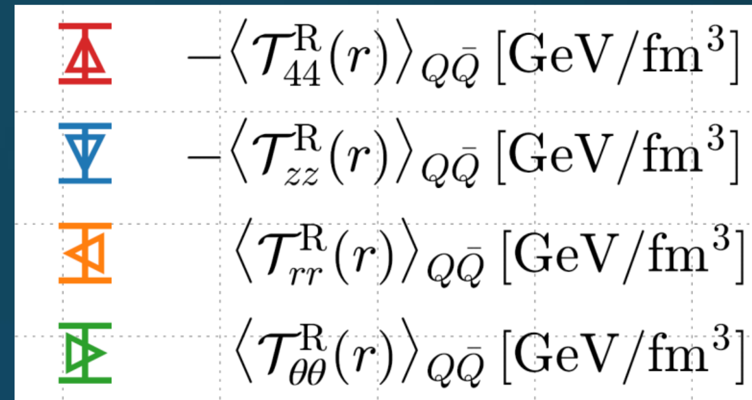


Degeneracy
in Maxwell theory

$$T_{rr} = T_{\theta\theta} = -T_{zz} = -T_{44}$$

Mid-Plane





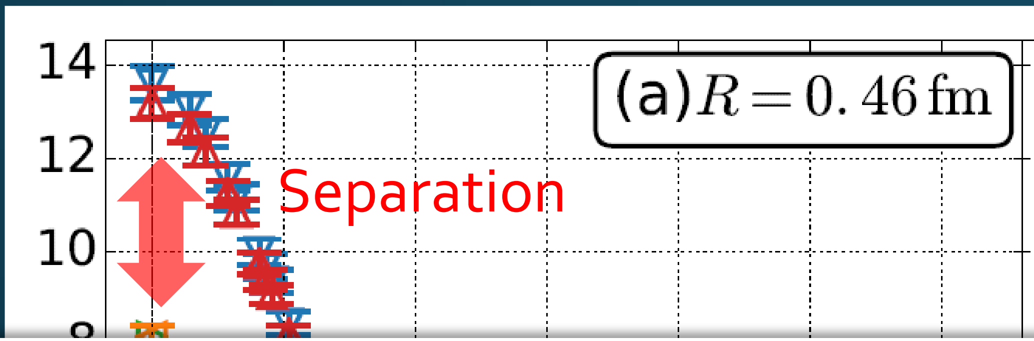
Continuum Extrapolated!

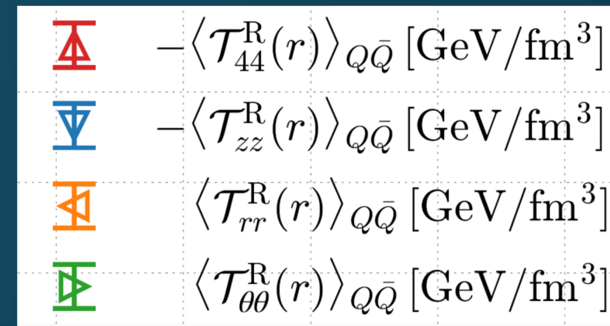
In Maxwell theory

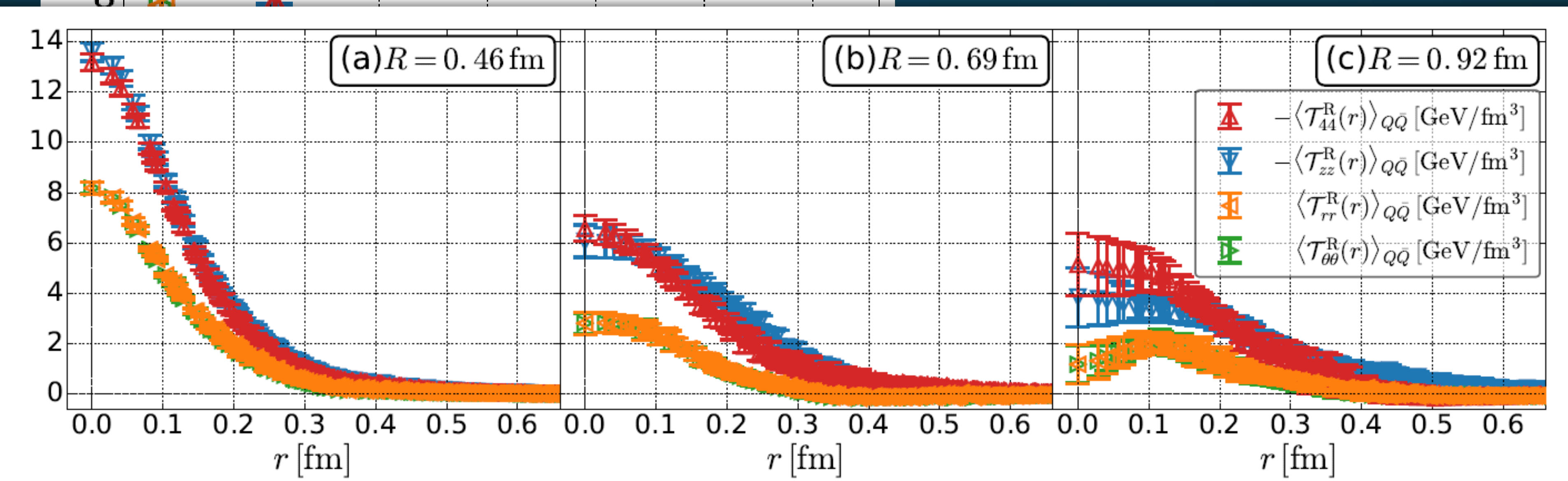
$$T_{rr} = T_{\theta\theta} = -T_{zz} = -T_{44}$$

- ☐ Degeneracy: $T_{44} \simeq T_{zz}, \quad T_{rr} \simeq T_{\theta\theta}$
- ☐ Separation: $T_{zz} \neq T_{rr}$
- ☐ Nonzero trace anomaly $\sum T_{cc} \neq 0$

Mid-Plane

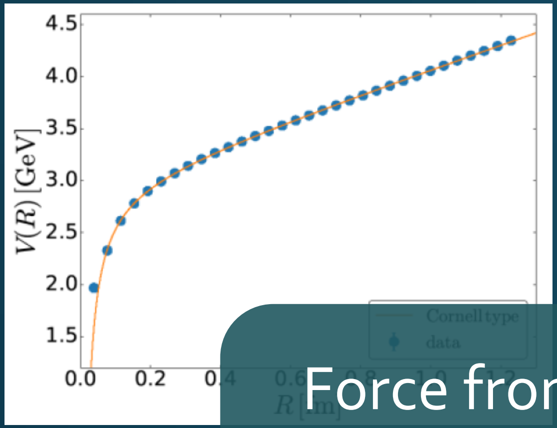


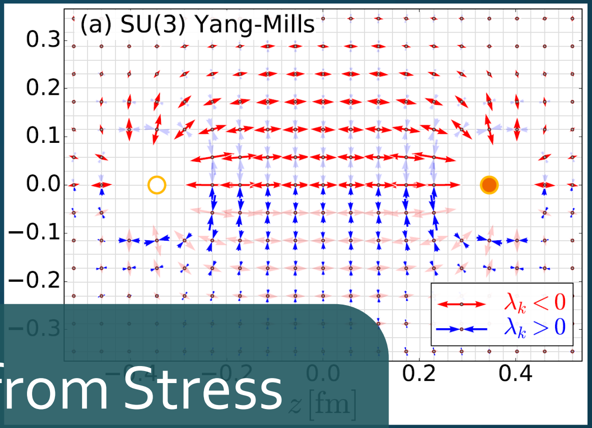




- Degeneracy: $T_{44} \simeq T_{zz}, \quad T_{rr} \simeq T_{\theta\theta}$
- Separation: $T_{zz} \neq T_{rr}$
- Nonzero trace anomaly $\sum T_{cc} \neq 0$

Force





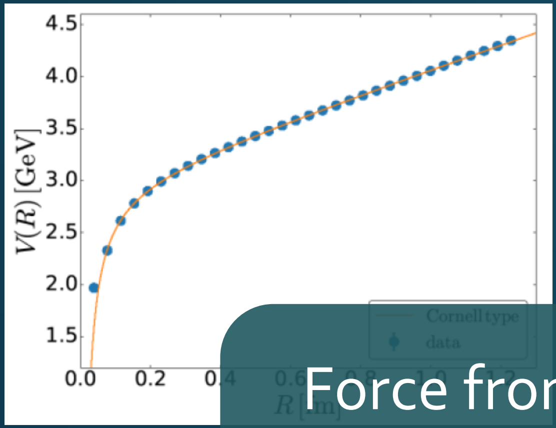
Force from Potential

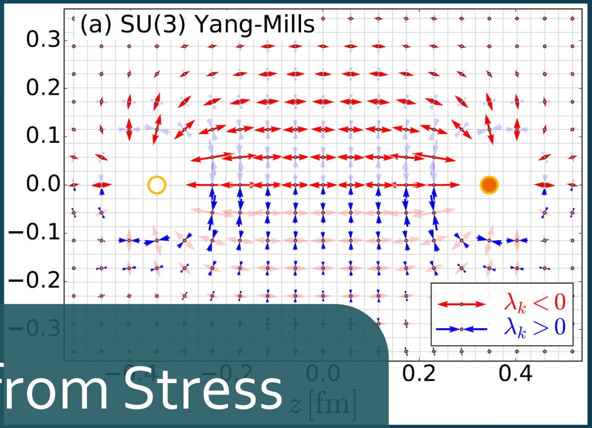
$$F_{\text{pot}} = -\frac{dV}{dR}$$

Force from Stress

$$F_{\text{stress}} = \int_{\text{mid.}} d^2x T_{zz}(x)$$

Force





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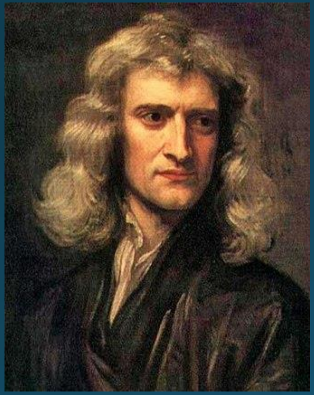
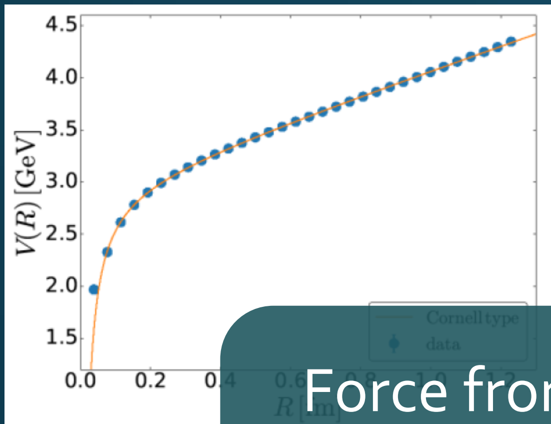
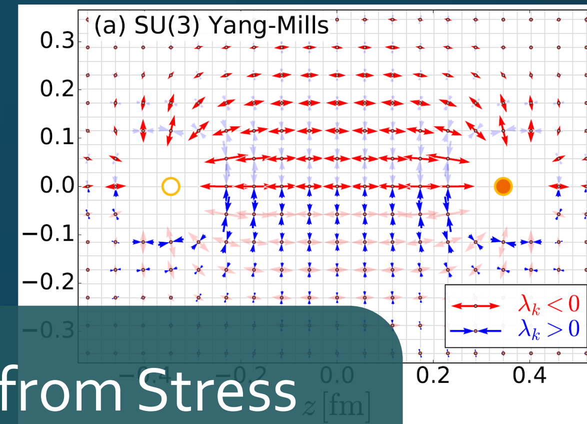


Newton
1687

Faraday
1839

Force





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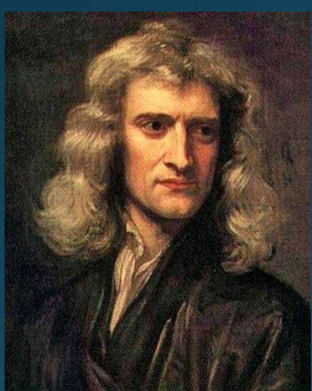
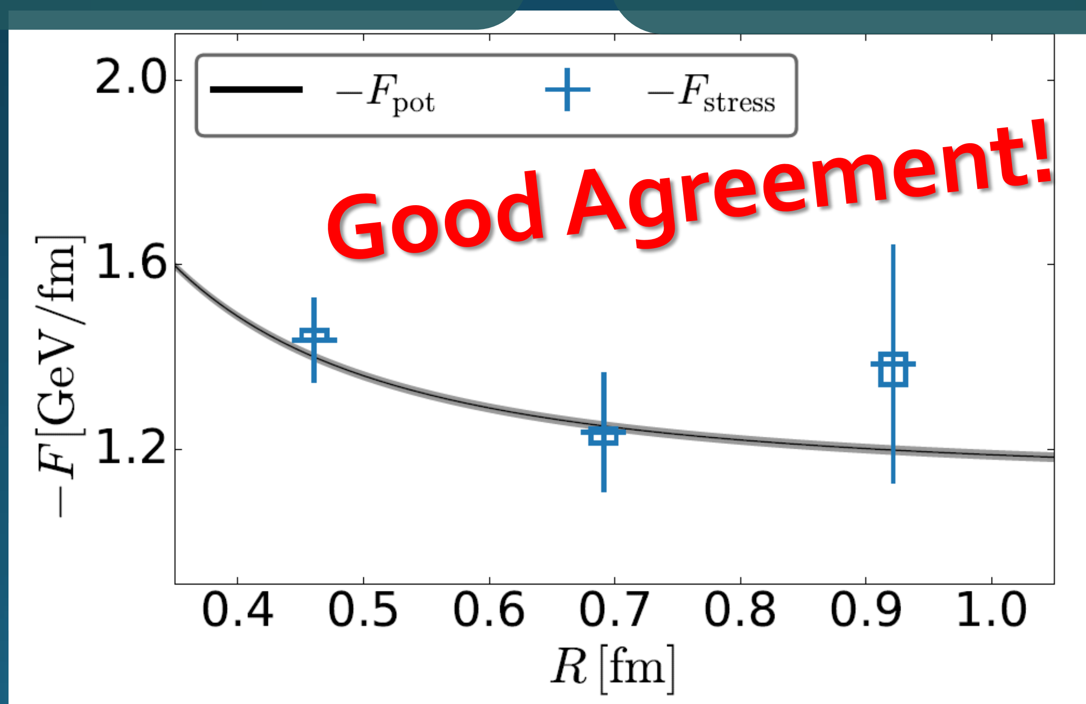


Newton
1687



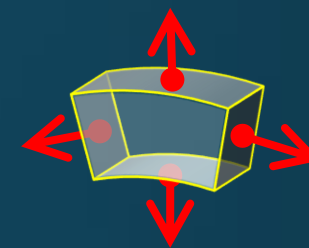
Faraday
1839

Momentum Conservation

Yanagihara, MK, PTEP2019

- In cylindrical coordinats,

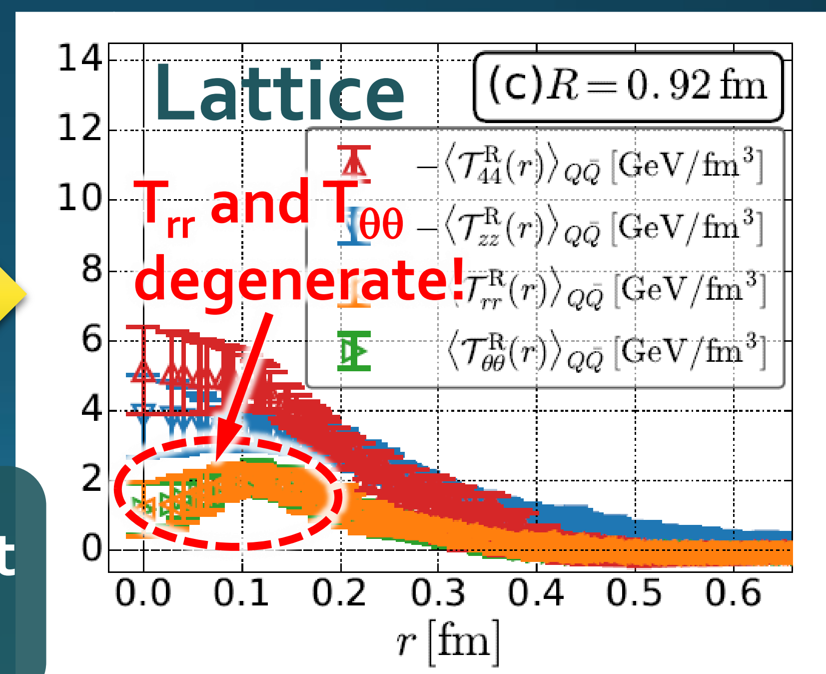
$$\partial_i T_{ij} = 0 \Rightarrow \partial_r (r T_{rr}) = T_{\theta\theta} - r \partial_z T_{rz}$$



- For infinitely-long flux tube

$$\partial_r (r T_{rr}) = T_{\theta\theta}$$

⇒ T_{rr} and $T_{\theta\theta}$ must separate!
$T_{\theta\theta}$ must change sign!



Effect of boundaries is important for the flux tube at R=0.92fm

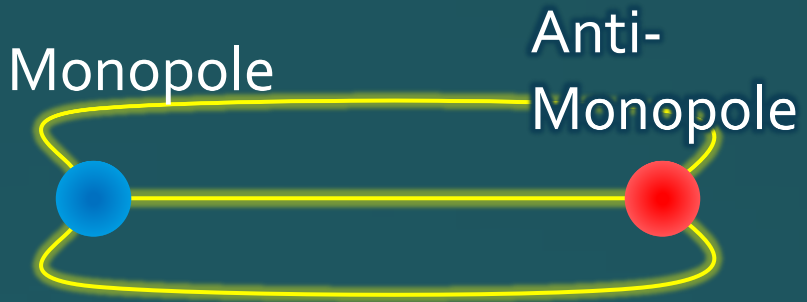
Dual Superconductor Picture

Nambu, 1970
Nielsen, Olesen, 1973
t 'Hooft, 1981
...

QCD Vacuum

Flux Tube

Superconductor



Magnetic Vortex

Dual (E⇔B)

Abelian-Higgs Model

Yanagihara, MK, 2019

Abelian-Higgs Model

$$\mathcal{L}_{\rm AH} = -\frac{1}{4}F_{\mu\nu}^2 + |(\partial_\mu + igA_\mu)\phi|^2 - \lambda(\phi^2 - v^2)^2$$

GL parameter: $\kappa = \sqrt{\lambda}/g$

- type-I : $\kappa < 1/\sqrt{2}$
- type-II : $\kappa > 1/\sqrt{2}$
- Bogomol'nyi bound : $\kappa = 1/\sqrt{2}$

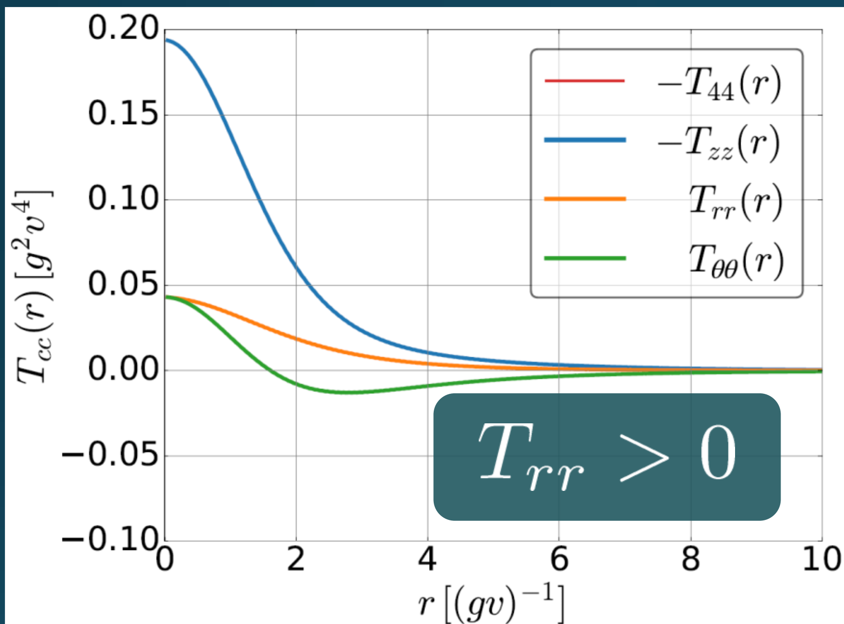
Infinitely long tube

- degeneracy
  $T_{zz}(r) = T_{44}(r)$ Luscher, 1981
- momentum conservation
  $$\frac{d}{dr}(rT_{rr}) = T_{\theta\theta}$$

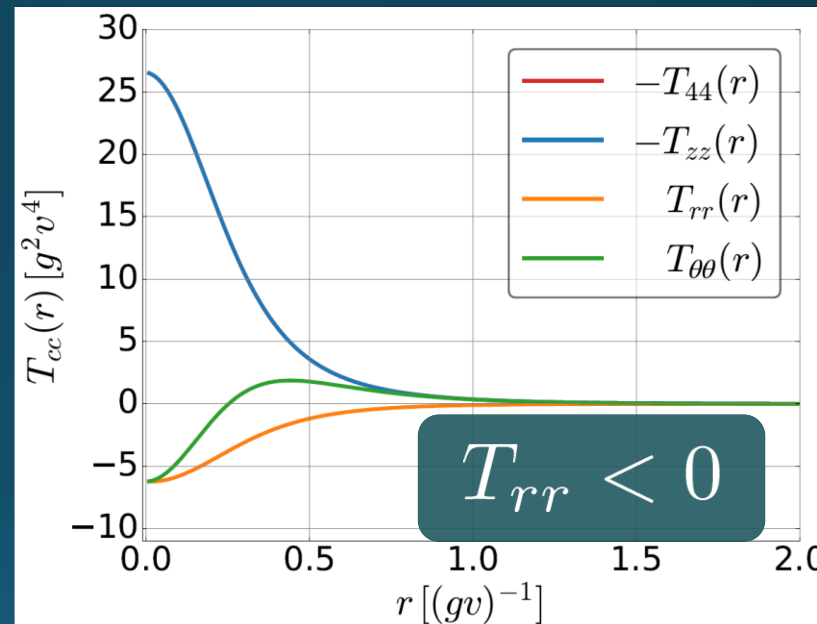
Stress Tensor in AH Model

infinitely-long flux tube

Type-I $\kappa = 0.1$



Type-II $\kappa = 3.0$



- ☐ No degeneracy bw T_{rr} & $T_{\theta\theta}$
- ☐ $T_{\theta\theta}$ changes sign

conservation law

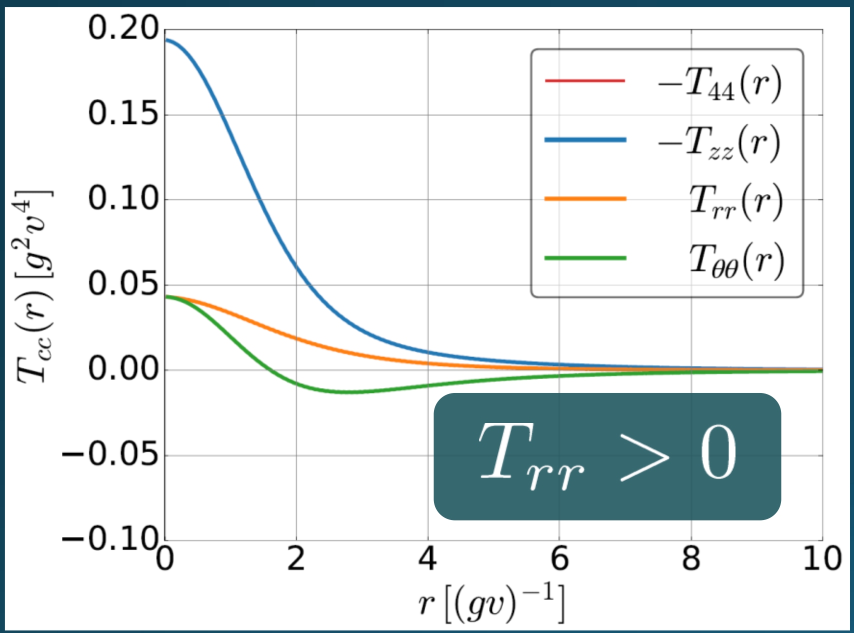
$$\frac{d}{dr}(rT_{rr}) = T_{\theta\theta}$$

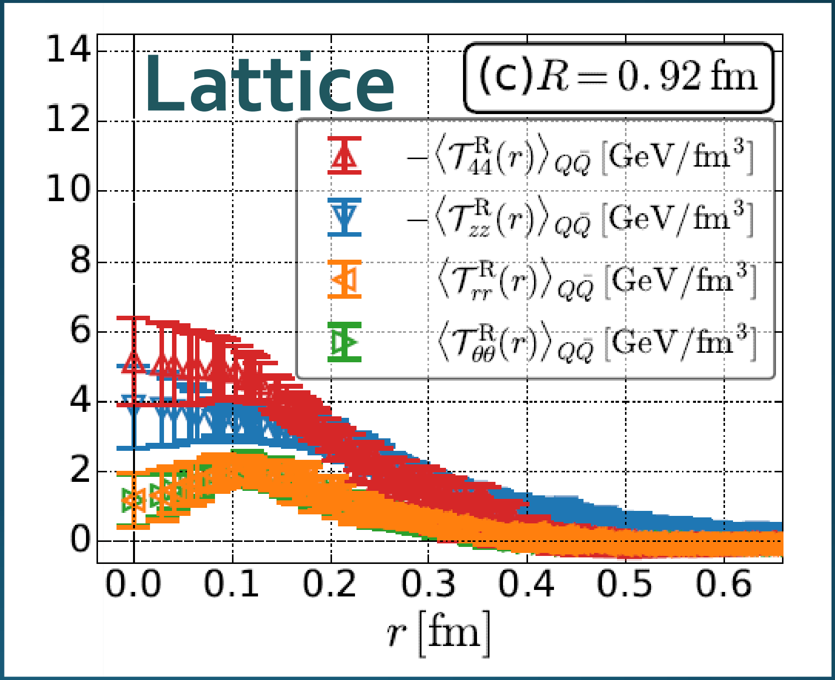
Stress Tensor in AH Model

infinitely-long flux tube

Type-I

$\kappa = 0.1$





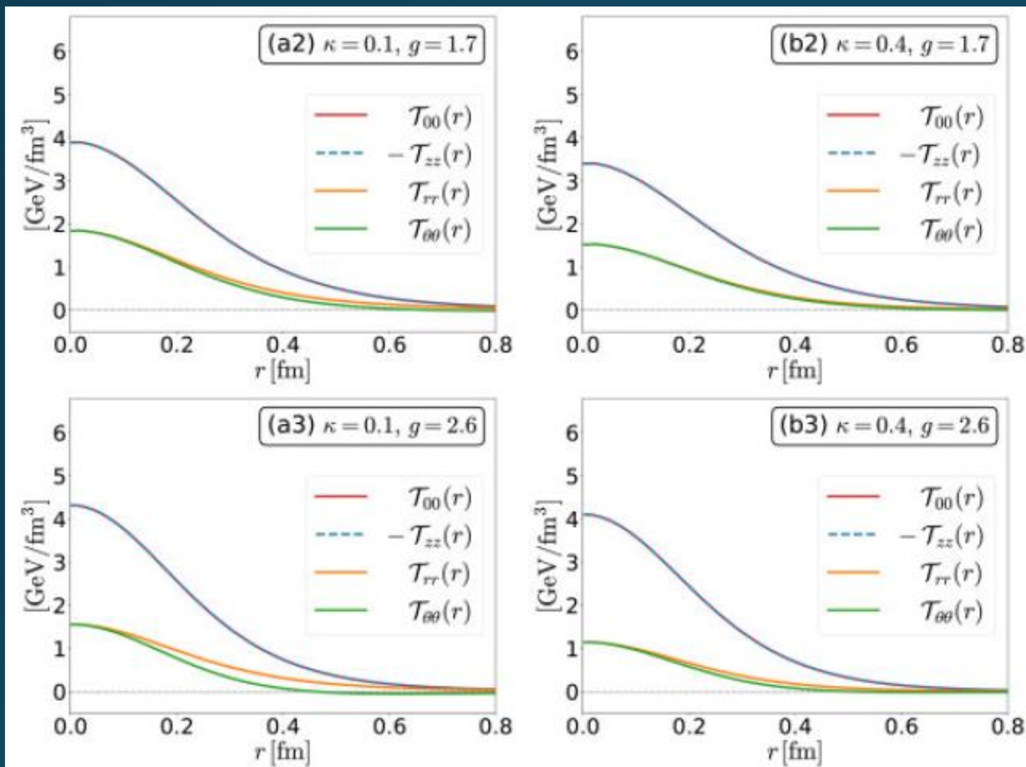
- No degeneracy bw T_{rr} & $T_{\theta\theta}$
- $T_{\theta\theta}$ changes sign

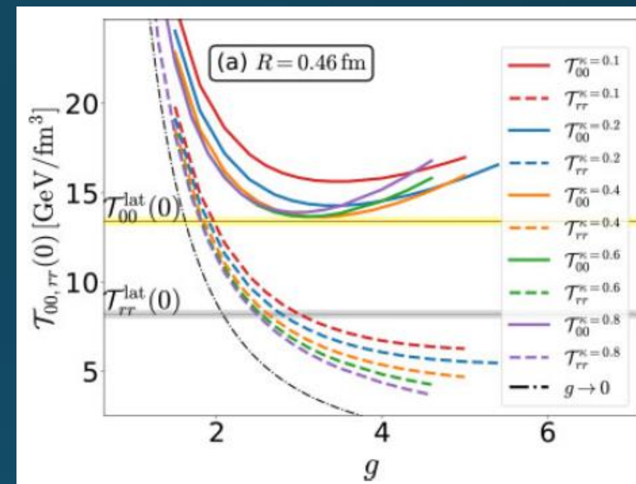
Inconsistent with lattice result

$$T_{rr} \simeq T_{\theta\theta}$$

Flux Tube with Finite Length

Yanagihara, MK (2019)





- ☐ AH model can reproduce lattice results **qualitatively** by tuning parameters.
- ☐ But, **quantitatively** all parameters are rejected.

EMT Distr. in Simple Systems

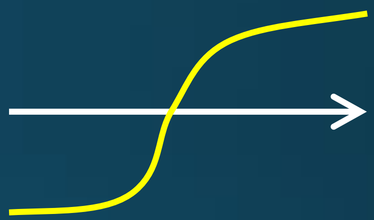
Ito, MK, in prep.

ϕ^4 Theory in 1+1d

$$\mathcal{L} = \frac{1}{2}(\partial_\mu \phi)^2 - \frac{\lambda}{4}\left(\phi^2 - \frac{m^2}{\lambda}\right)^2$$

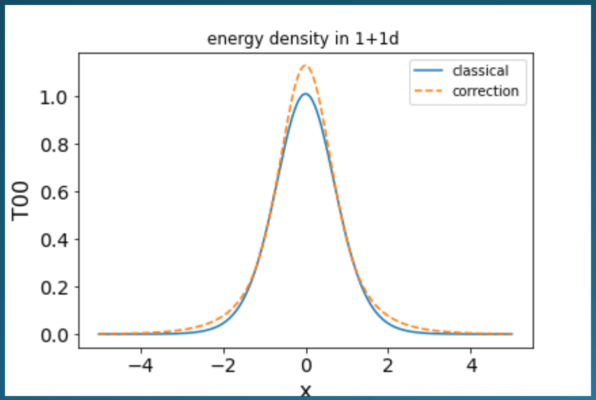
- **Soliton (kink)**

$$\phi(x) = \frac{m}{\sqrt{\lambda}} \tanh \frac{mx}{\sqrt{2}}$$



- **Quantum effect on EMT at 1-loop order**

$T_{00}(x)$



Confirmation of EMT conservation

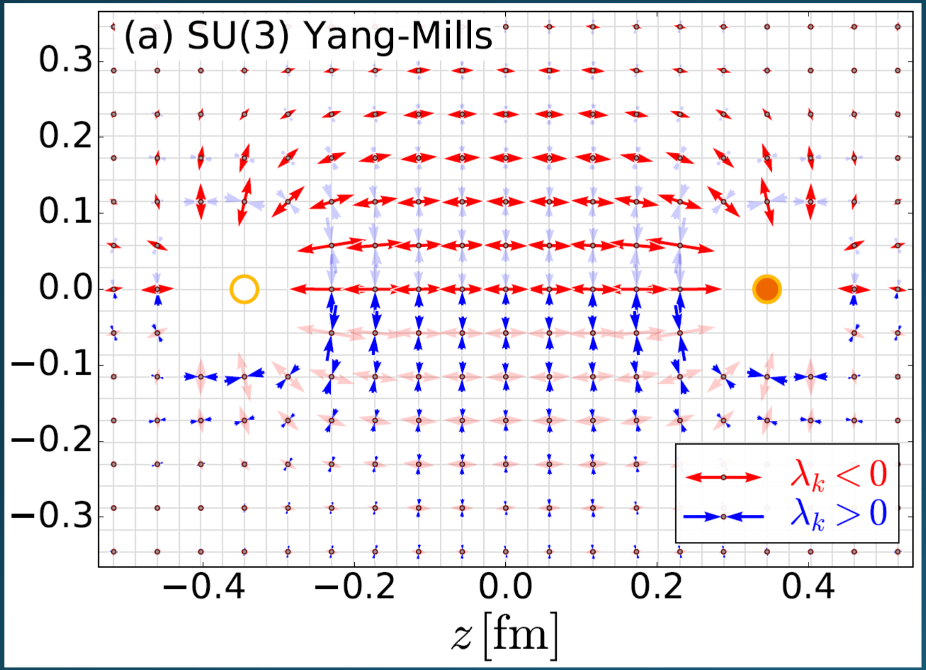
$$\partial_x T_{11}(x) = 0$$

Single Q System
in the Deconfined Phase

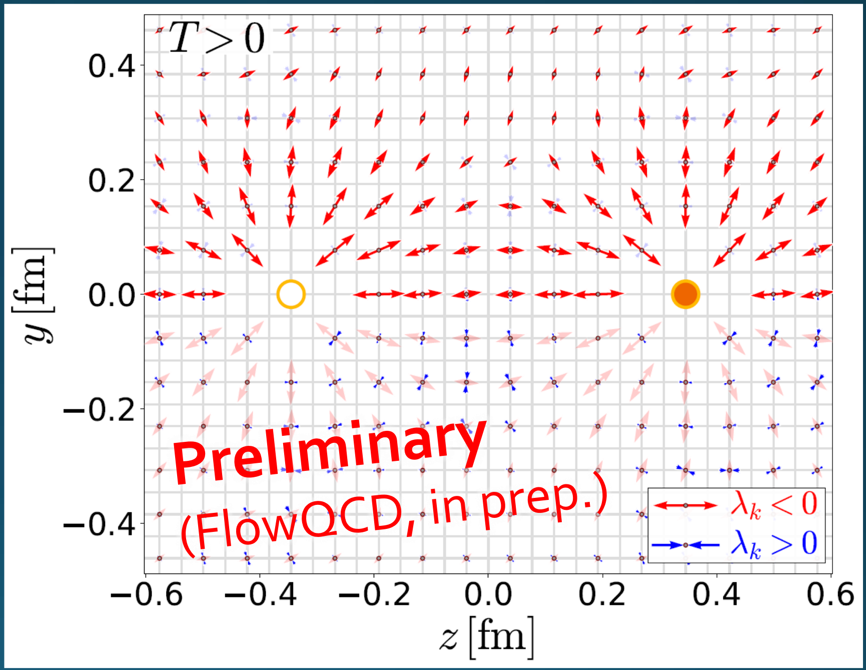
FlowQCD, PRD **102**, 114522 (2020)

Flux Tube at Nonzero T

Vacuum



$T = 1.42T_c$



- Dissociation of the flux tube at $T>T_c$.

Motivations

□ $T < T_c$: Heavy-light meson

- EMT distribution in the meson

□ $T > T_c$: Single charge

- Screening
- Running coupling

□ $T \approx T_c$

- Confinement transition

This study:

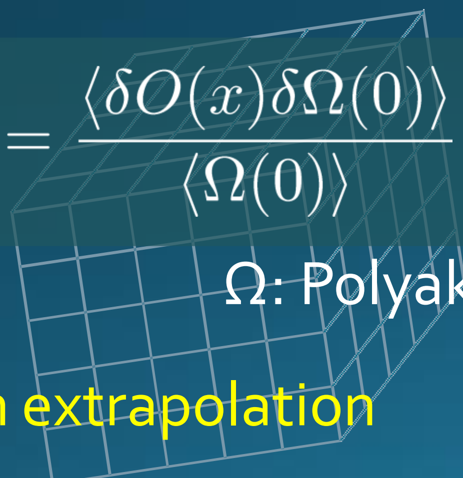
$T > T_c$ in pure YM

Lattice Setup

- SU(3) Yang-Mills (Quenched)
- Wilson gauge action
- Clover operator

- Analysis above Tc
- Simulation on a Z_3 minimum
- EMT around a Polyakov loop

$$\langle O(x) \rangle_{\mathrm{Q}} = \frac{\langle \delta O(x) \delta \Omega(0) \rangle}{\langle \Omega(0) \rangle}$$



Ω: Polyakov loop

- continuum extrapolation

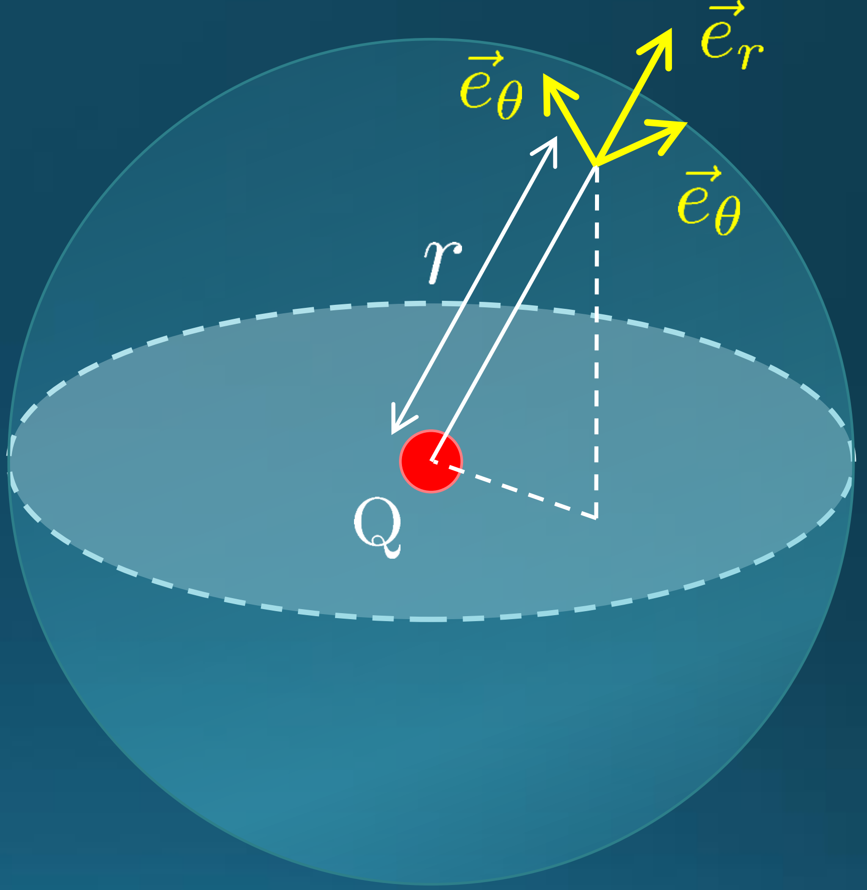
T/T_c	N_s	N_τ	β	a [fm]	N_{conf}
1.20	40	10	6.336	0.0551	500
	48	12	6.467	0.0460	650
	56	14	6.581	0.0394	840
	64	16	6.682	0.0344	1,000
	72	18	6.771	0.0306	1,000
1.44	40	10	6.465	0.0461	500
	48	12	6.600	0.0384	650
	56	14	6.716	0.0329	840
	64	16	6.819	0.0288	1,000
	72	18	6.910	0.0256	1,000
2.00	40	10	6.712	0.0331	500
	48	12	6.853	0.0275	650
	56	14	6.973	0.0236	840
	64	16	7.079	0.0207	1,000
	72	18	7.173	0.0184	1,000
2.60	40	10	6.914	0.0255	500
	48	12	7.058	0.0212	650
	56	14	7.182	0.0182	840
	64	16	7.290	0.0159	1,000
	72	18	7.387	0.0141	1,000

FlowQCD, PRD **102**, 114522 (2020)

Spherical Coordinates

EMT is diagonalized
in Spherical Coordinates

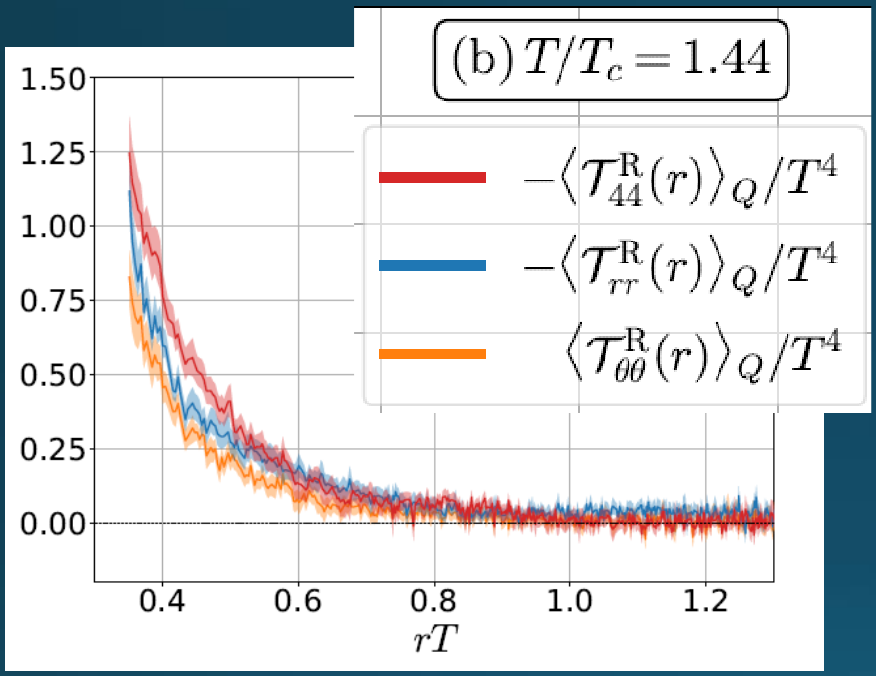
$$T_{cc'}(r) = \begin{pmatrix} T_{rr} & & & \\ & T_{\theta\theta} & & \\ & & T_{\theta\theta} & \\ & & & T_{44} \end{pmatrix}$$



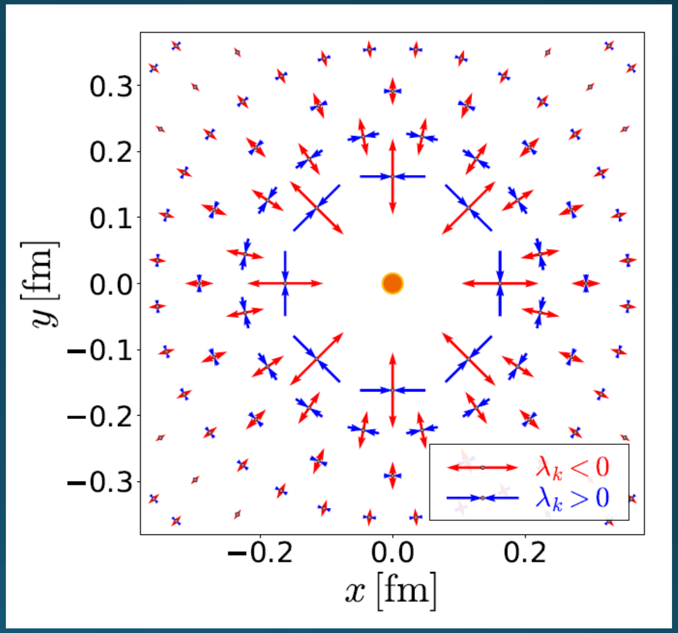
- Maxwell theory

$$T_{44} = T_{rr} = -T_{\theta\theta} = -\frac{|\boldsymbol{E}|^2}{2} = -\frac{\alpha}{8\pi}\frac{1}{r^4}$$

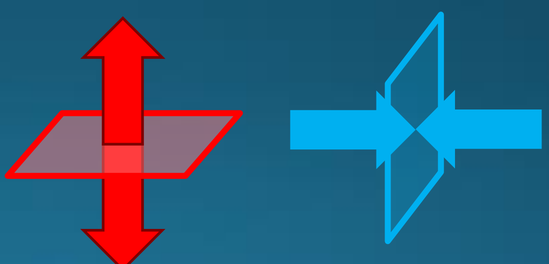
Stress Tensor Around a Quark



$T=1.44T_c$

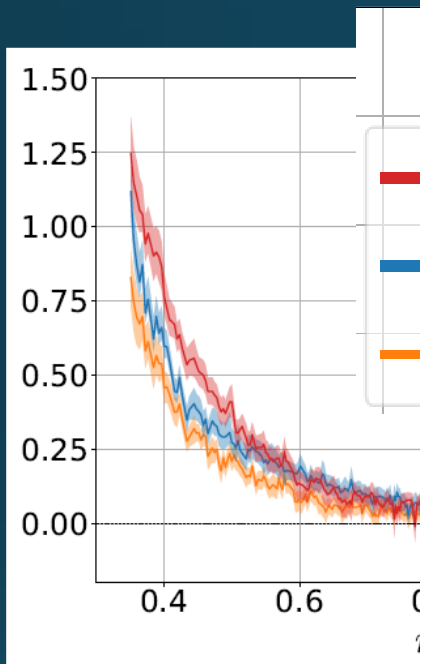


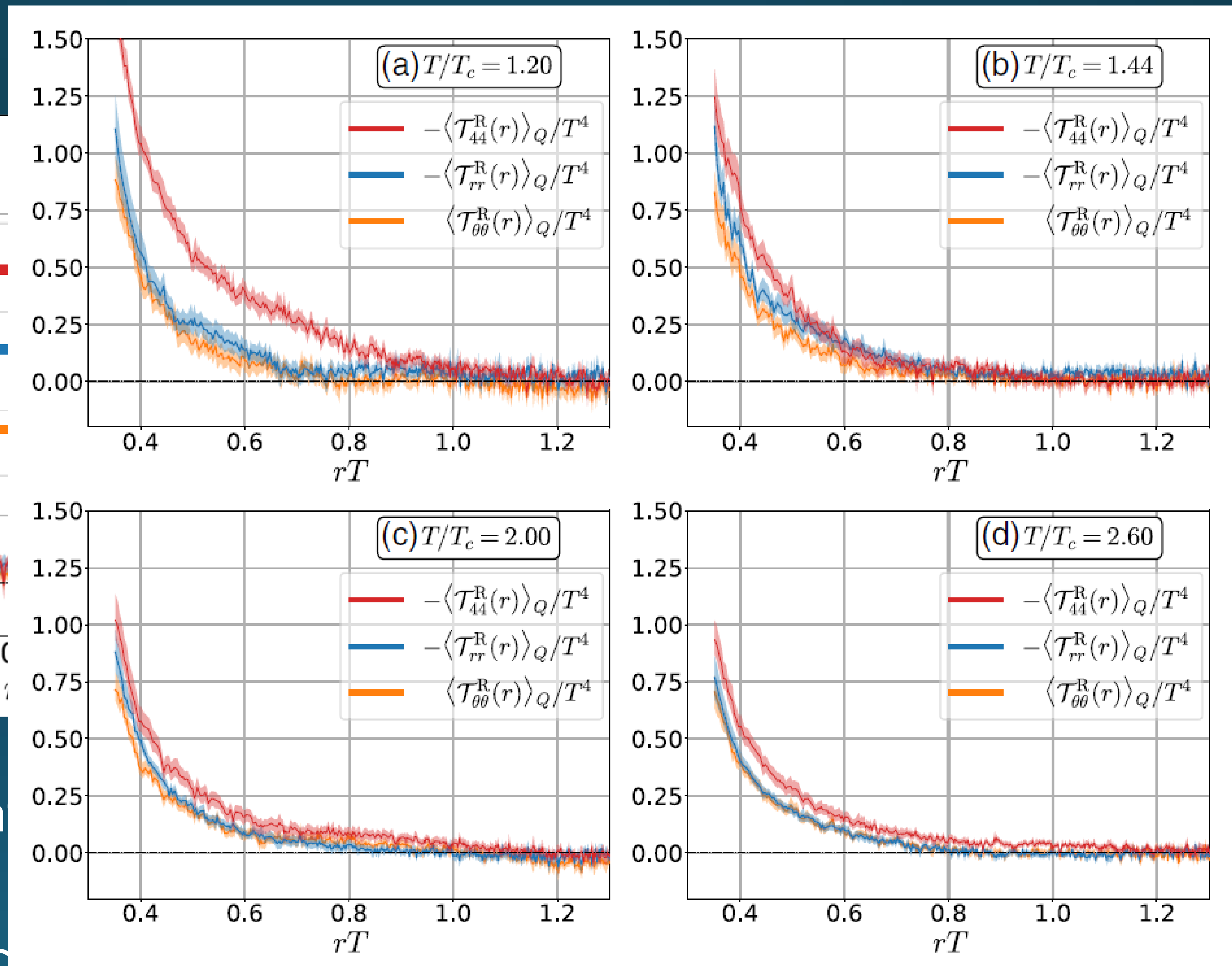
- Suppression at large distance
- Separation of different channels
- $|T_{44}| > |T_{rr}| \sim |T_{\theta\theta}|$



pulling **pushing**

Stress Tensor Around a Quark





- Suppression a
- Separation of
- $|T_{44}| > |T_{rr}| \sim$
- Clearer separation for lower T

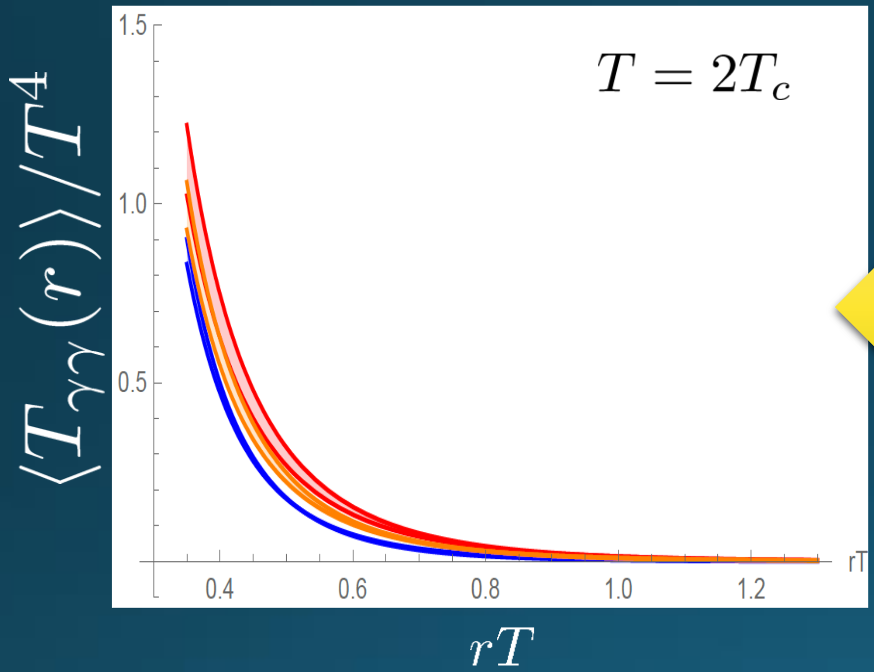
pulling

pushing

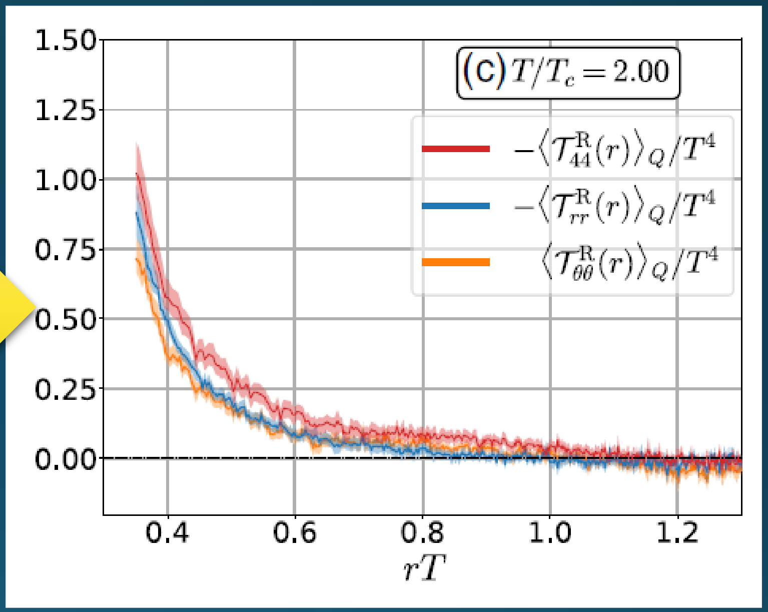
Perturbative Analysis

M. **Berwein**, private comm.

Perturbation



Lattice

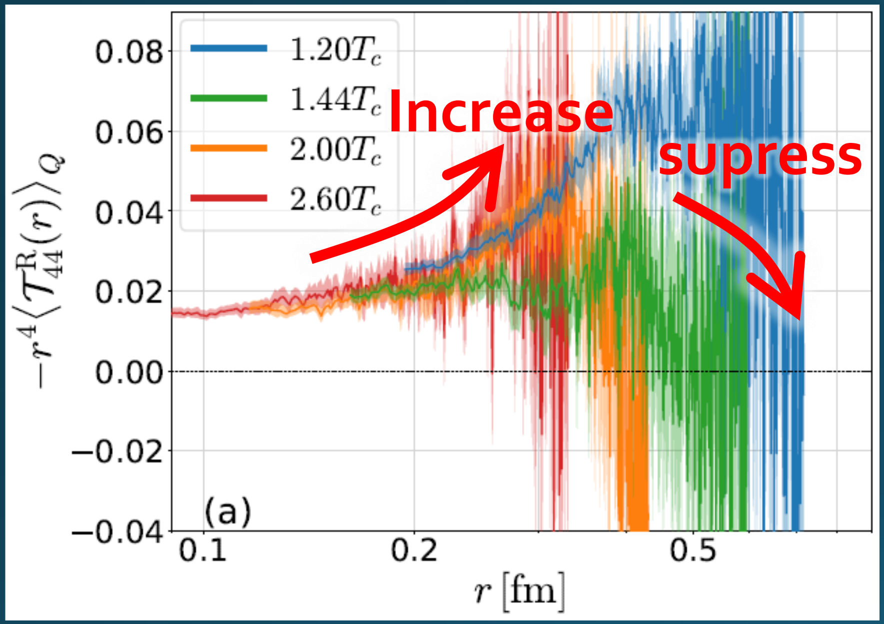


Perturbation:
Combination of
NLO pert. + NLO EQCD

- $|T_{44}| > |T_{rr}|$ is reproduced by perturbation.
- Hierarchy of T_{rr} , $T_{\theta\theta}$ does not match?

r Dependence

$r^4\langle T_{00}(r)\rangle$



Leading order perturbation

$$\langle \mathcal{T}_{44}(r)\rangle = \langle \mathcal{T}_{rr}(r)\rangle = -\langle \mathcal{T}_{\theta\theta}(r)\rangle$$

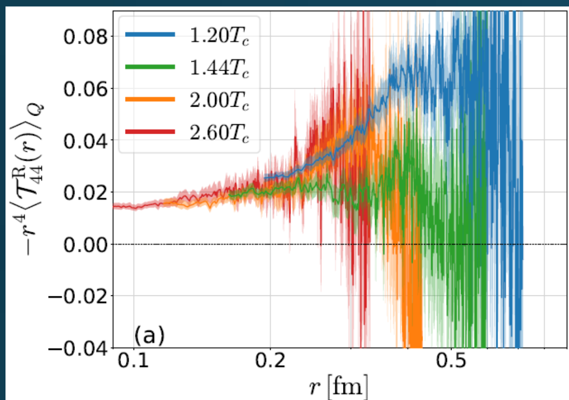
$$= -\frac{C_F}{8\pi}\alpha_s\frac{(m_{\mathrm{D}}r+1)^2}{r^4}e^{-2m_{\mathrm{D}}r}$$

Higher order terms:
M. Berwein, in progress

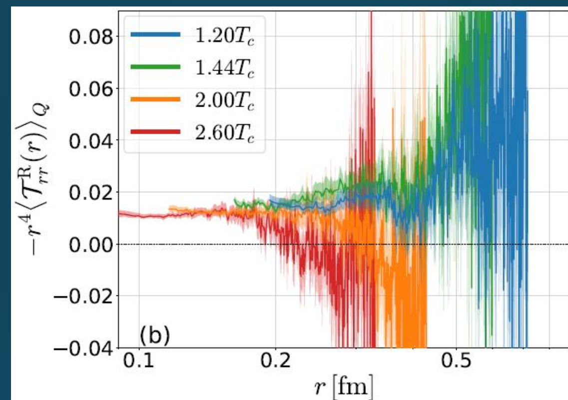
- Increase at short r / suppression at larger r
- T dependence is suppressed at $r < 1/T$
- Too noisy at large r for extracting screening mass m_{D}

Channel Dependence

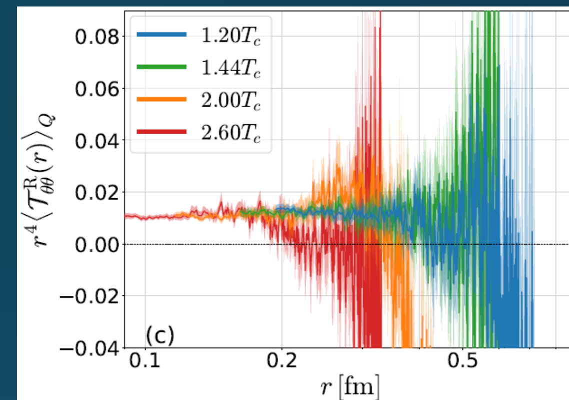
$r^4\langle T_{00}(r)\rangle$



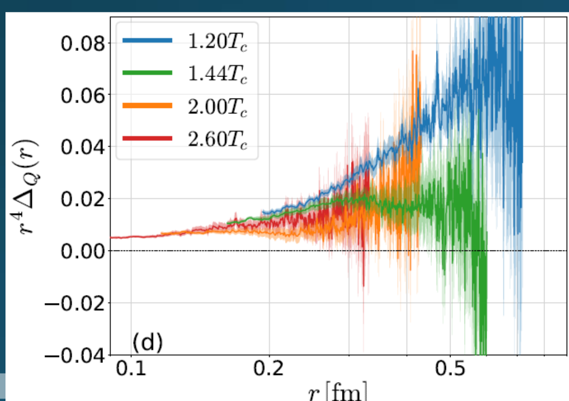
$-r^4\langle T_{rr}(r)\rangle$



$r^4\langle T_{\theta\theta}(r)\rangle$



$r^4\Delta(r) = -r^4\langle T_{\mu\mu}\rangle$



- Separation b/w channels becomes clearer for smaller *T*

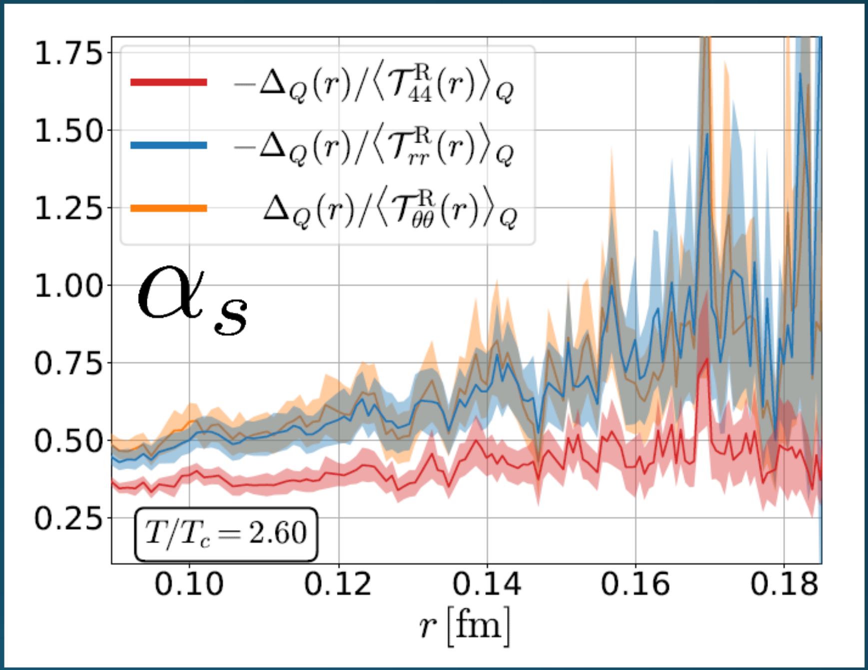
FlowQCD, PRD **102**, 114522 (2020)

Running Coupling

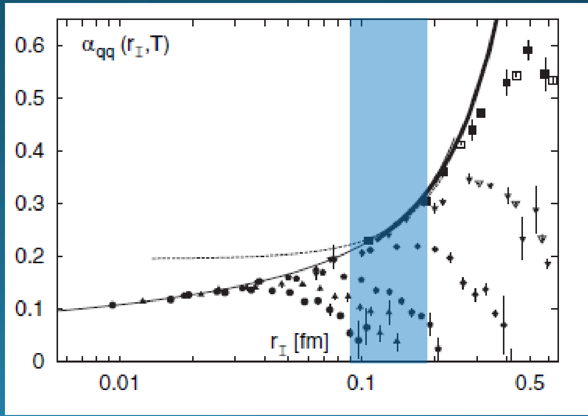
- **Estimate of α_s**

$$\left|\frac{\langle T_{\mu\mu}\rangle}{\langle \mathcal{T}_{44,rr,\theta\theta}(r)\rangle}\right| = \frac{11}{2\pi}\alpha_s + \mathcal{O}(g^3),$$

  - by the formula at the leading-order perturbation theory
  - channel dependent



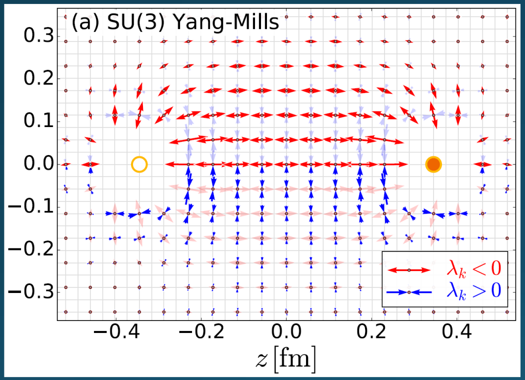
- Consistent with the estimate from $Q\bar{Q}$ potential

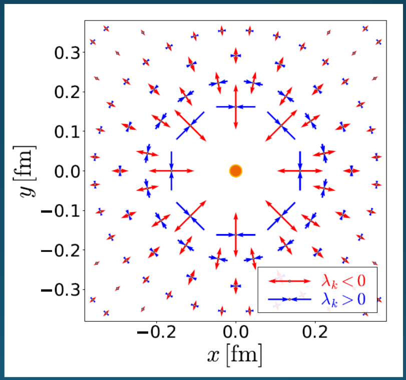


Kaczmarek, Karsch, Zantow, 2004

Summary

- Static charges are fundamental but convenient tools for studying YM gauge theories.
- Now, lattice simulations of EMT in static-quark systems are available thanks to SF*t*X (gradient flow) method.





So many future studies

- Single Q in full QCD @ $T<T_c$ = heavy-light meson
- QQQ, QQ, etc. / T dependence
- Hadrons

backup

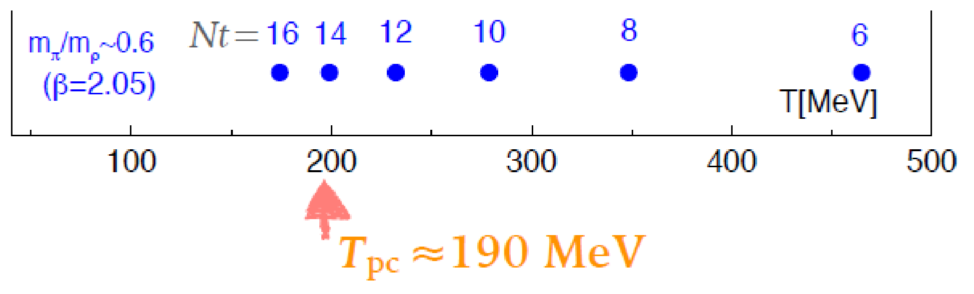
N_f=2+1 QCD Thermodynamics

Taniguchi+ (WHOT-QCD), PRD**96**, 014509 (2017)

- N_f=2+1 QCD, Iwasaki gauge + NP-clover
- m_{PS}/m_V ≈0.63 / almost physical s quark mass

- T=0: CP-PACS+JLQCD (ß=2.05, 28^3x56, a≈0.07fm)
- T>0: 32^3xN_t, N_t = 4, 6, ... , 14, 16):
- T≈174-697MeV

- t→0 extrapolation only (No continuum limit)



EMT on the Lattice: Conventional

Lattice EMT Operator Caracciolo+, 1990

$$T_{\mu\nu} = Z_6 T^{[6]}_{\mu\nu} + Z_3 T^{[3]}_{\mu\nu} + Z_1 \left(T^{[1]}_{\mu\nu} - \langle T^{[1]}_{\mu\nu} \rangle \right)$$

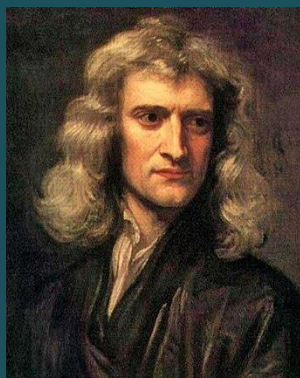
$$T^{[6]}_{\mu\nu} = (1 - \delta_{\mu\nu}) F^a_{\mu\rho} F^a_{\nu\rho}, \;\; T^{[3]}_{\mu\nu} = \delta_{\mu\nu} \left(F^a_{\mu\rho} F^a_{\nu\rho} - \frac{1}{4} F^a_{\rho\sigma} F^a_{\rho\sigma} \right), \;\; T^{[1]}_{\mu\nu} = \delta_{\mu\nu} F^a_{\rho\sigma} F^a_{\rho\sigma}$$

- Determination of Zs are necessary.
- Non-pert. Determination of Zs
  - Shifted-boundary method
  - Full QCD with fermions

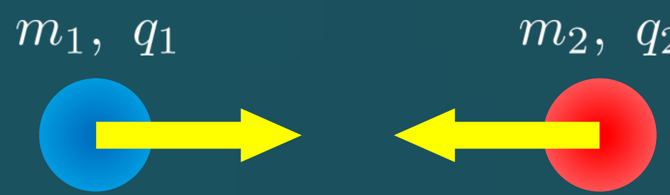
Giusti, Pepe, 2014~; Borsanyi+, 2018
Brida, Giusti, Pepe, 2020

Force

Action-at-a-distance



Newton
1687



$$F = -G\frac{m_1 m_2}{r^2} \qquad F = -\frac{1}{4\pi\epsilon_0}\frac{q_1 q_2}{r^2}$$

Local interaction

Faraday
1839