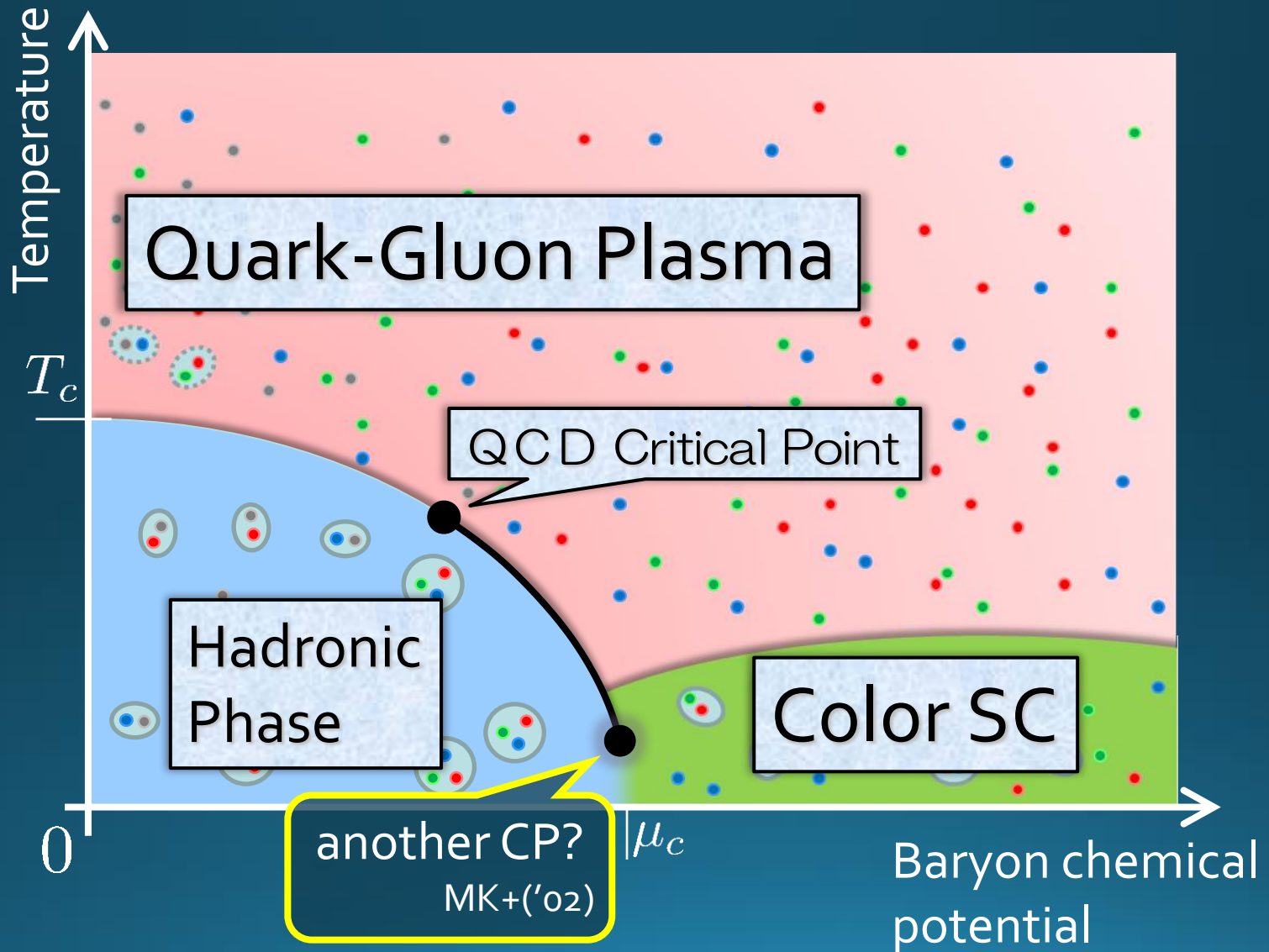


# High-Density Physics at J-PARC

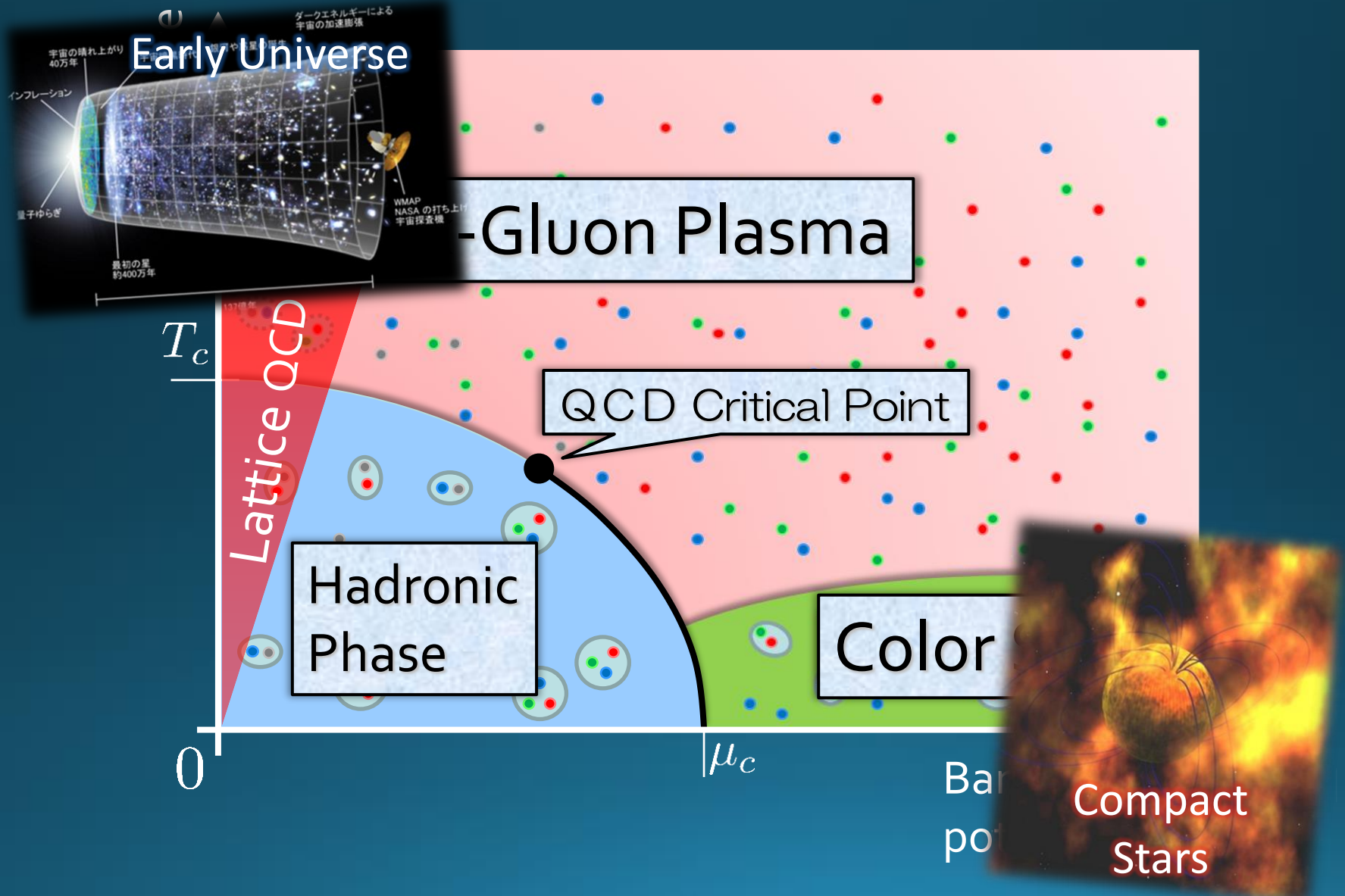
Masakiyo Kitazawa  
(Osaka U.)

2nd J-PARC HEF-ex WS, Online, 2022/Feb./16

# QCD Phase Diagram

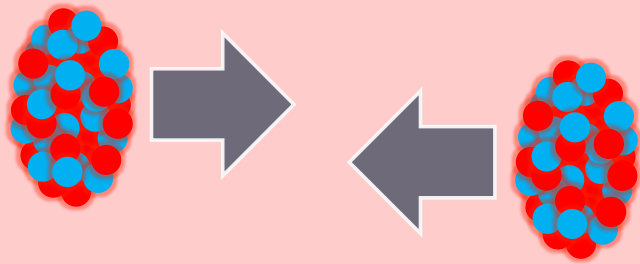


# QCD Phase Diagram



# Relativistic Heavy-Ion Collisions

**Collide 2 heavy nuclei**



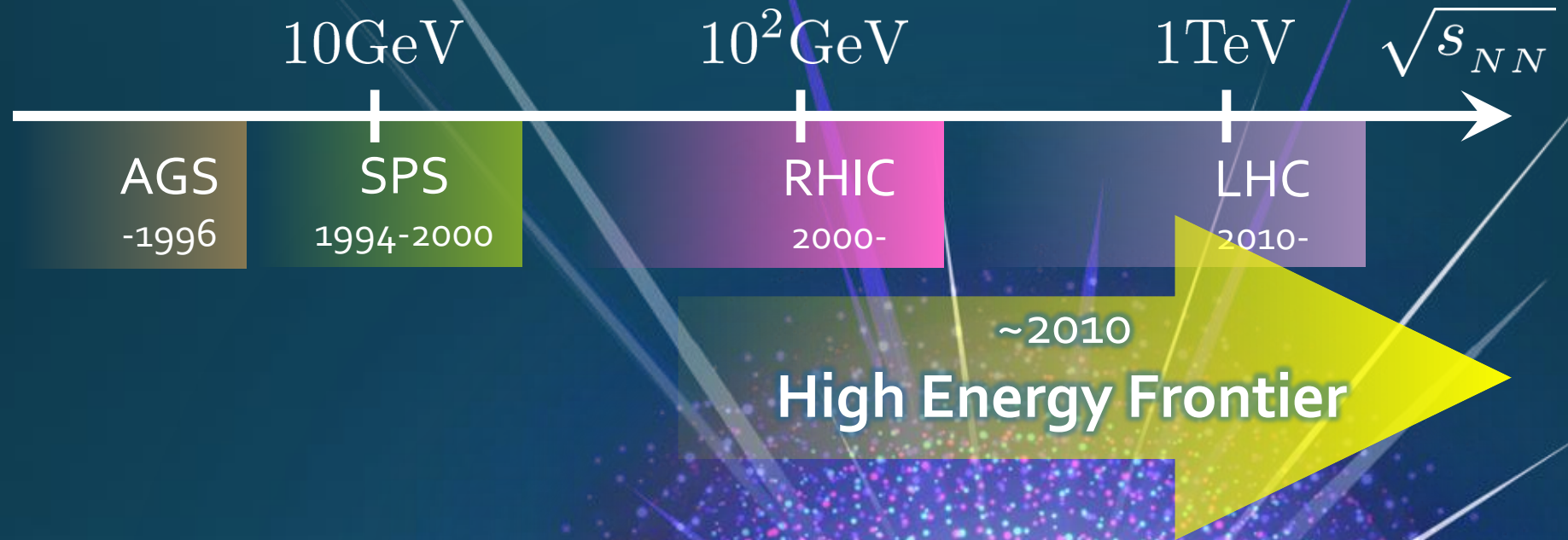
**RHIC** (2000~)  
QGP Formation  
Strongly coupled QGP

**LHC** (2010~)  
Precision measurement  
of the QGP

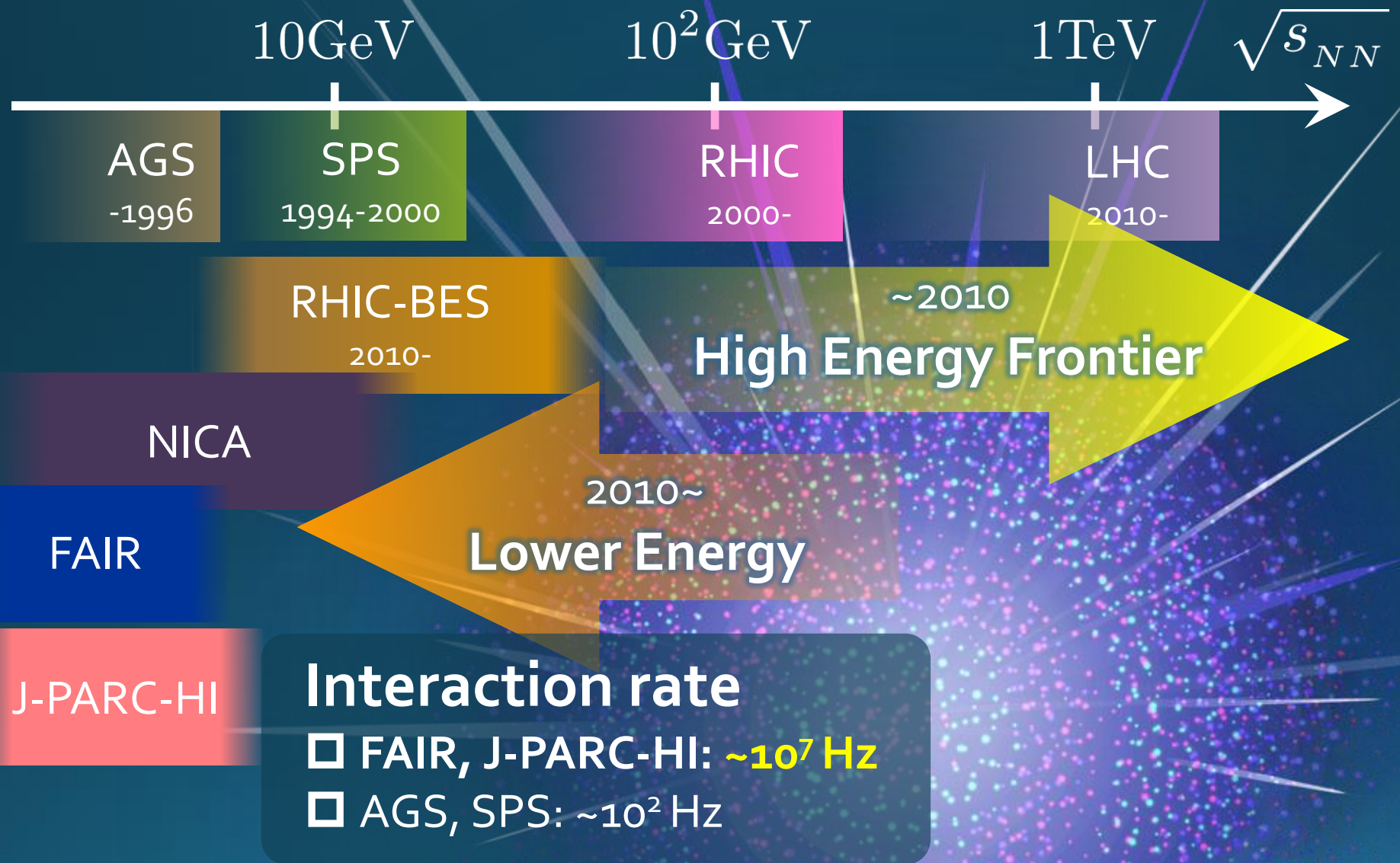
## Physics

- Hot & dense medium
- Early Universe
- Quark-gluon plasma
- QCD phase structure

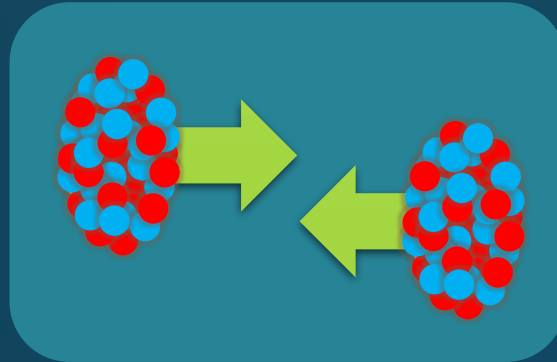
# Brief History of Relativistic HIC



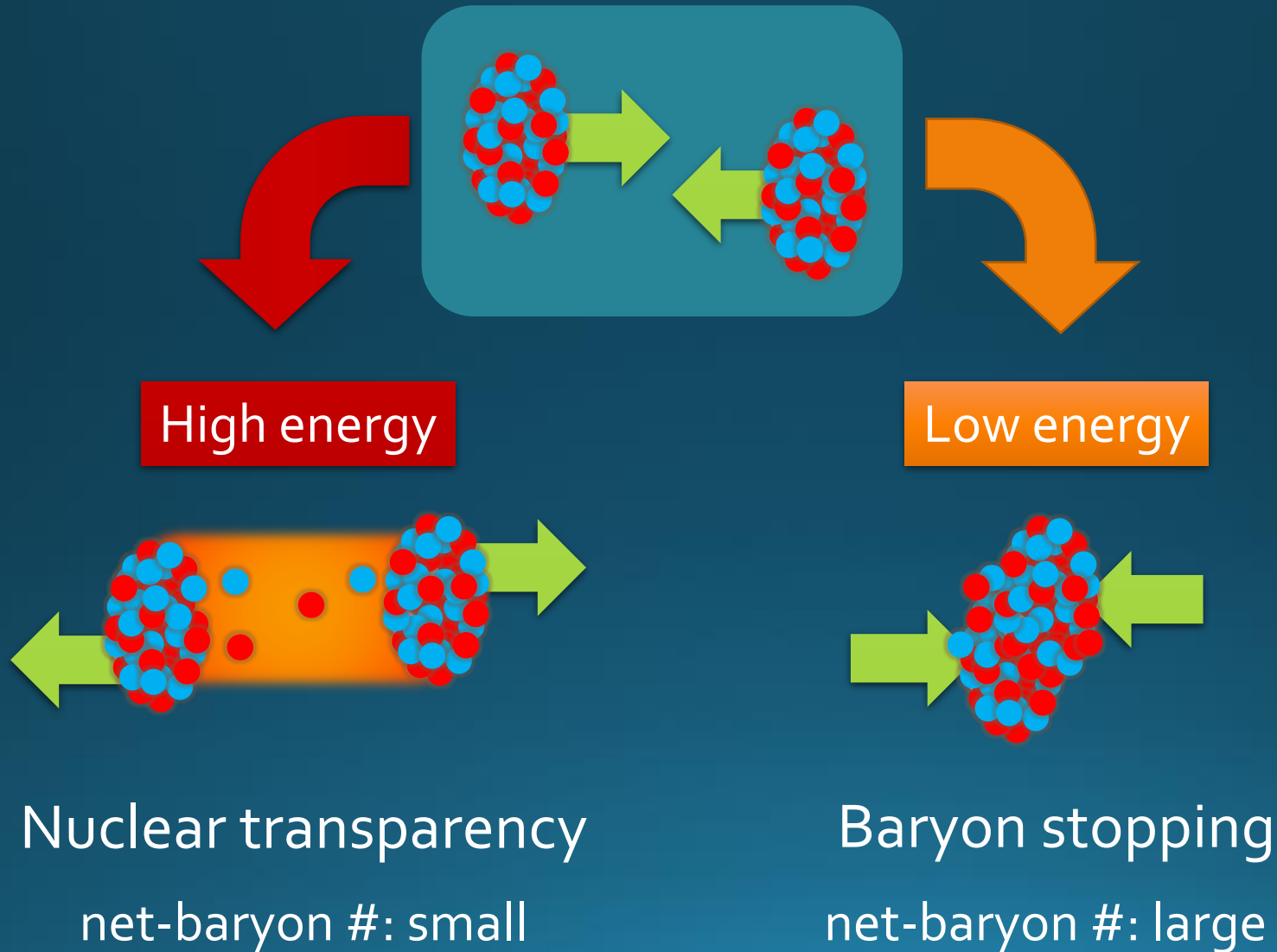
# Brief History of Relativistic HIC



# Beam-Energy Dependence



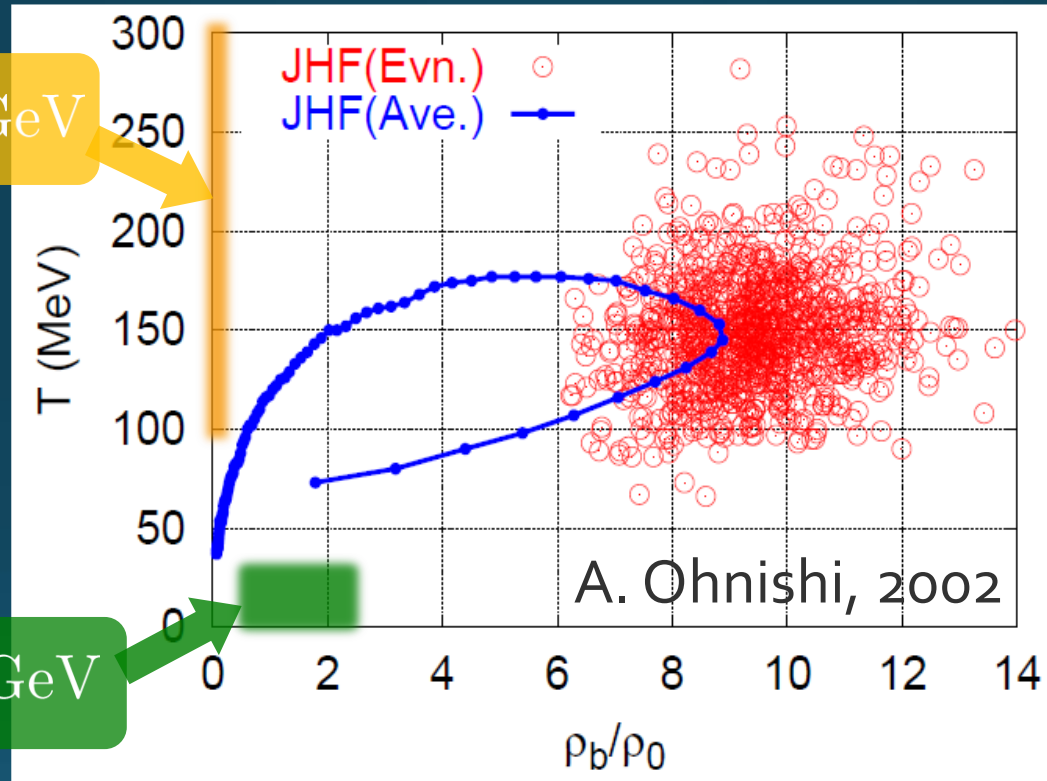
# Beam-Energy Dependence





# Maximum Density

Time evolution in  $T$ - $\rho$  plane by JAM



$\sqrt{s_{NN}} > 100 \text{ GeV}$

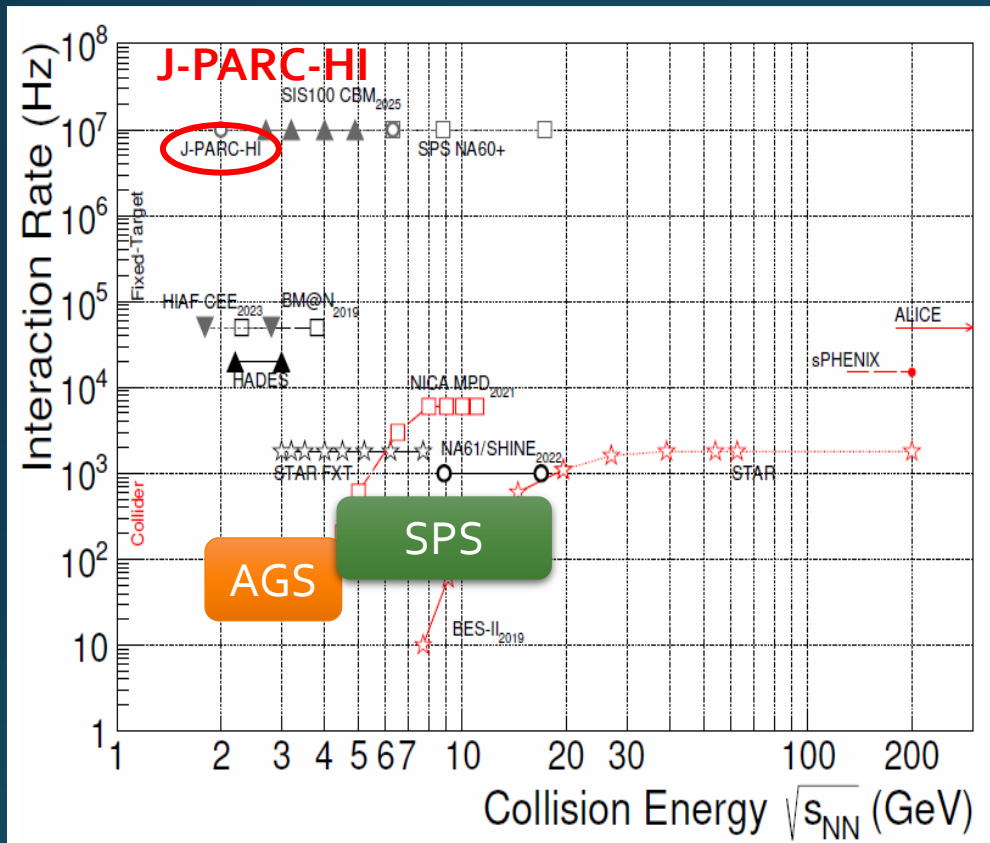
$E/A = 20 \text{ GeV}$

$\sqrt{s_{NN}} \simeq 6 \text{ GeV}$

$E/A < 1 \text{ GeV}$

- Maximum density  $5 \sim 10\rho_0$  @  $E/A \sim 20 \text{ GeV}$
- Large event-by-event fluctuations?

# Collision Rate



Galatyuk, NPA982,163 (2019)

## J-PARC-HI:

High-luminosity X Fixed target  
 → World highest rate  $\sim 10^8$  Hz



5-order higher than AGS, SPS

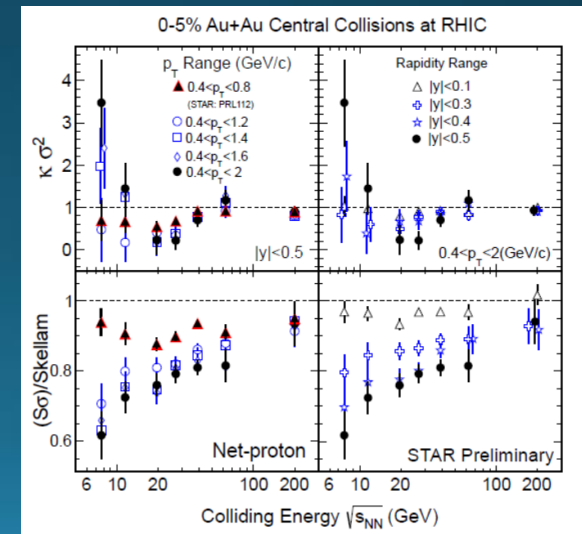
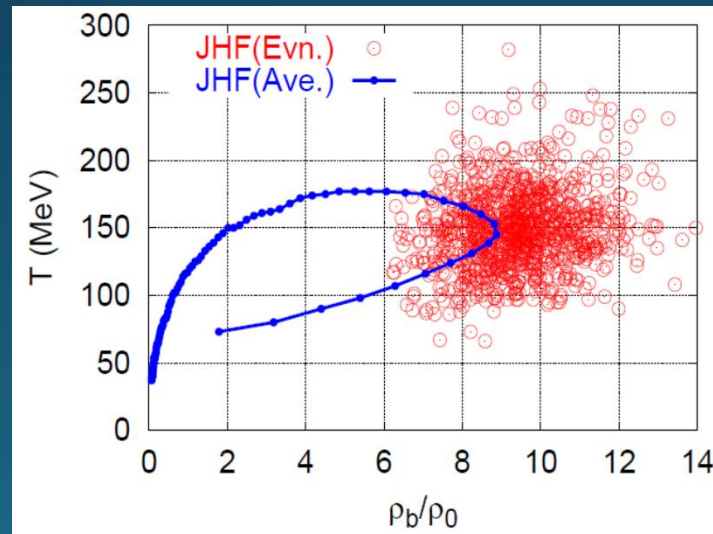
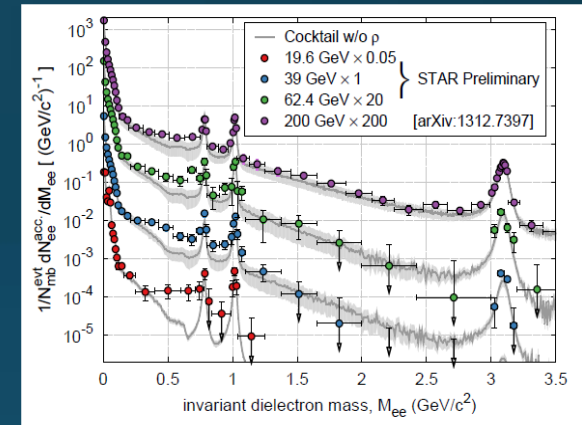
AGS, SPS = J-PARC-HI  
 1 year = 5 min.



- High-statistical exp.
- various event selections
- higher order correlations
- search of rare events

# Various Observables

- Flow
- Dilepton / photon
- Fluctuations, higher-order cumulants
- $\Xi, \Omega, \dots$
- Sophisticated event selections
- Various correlations



# Contents

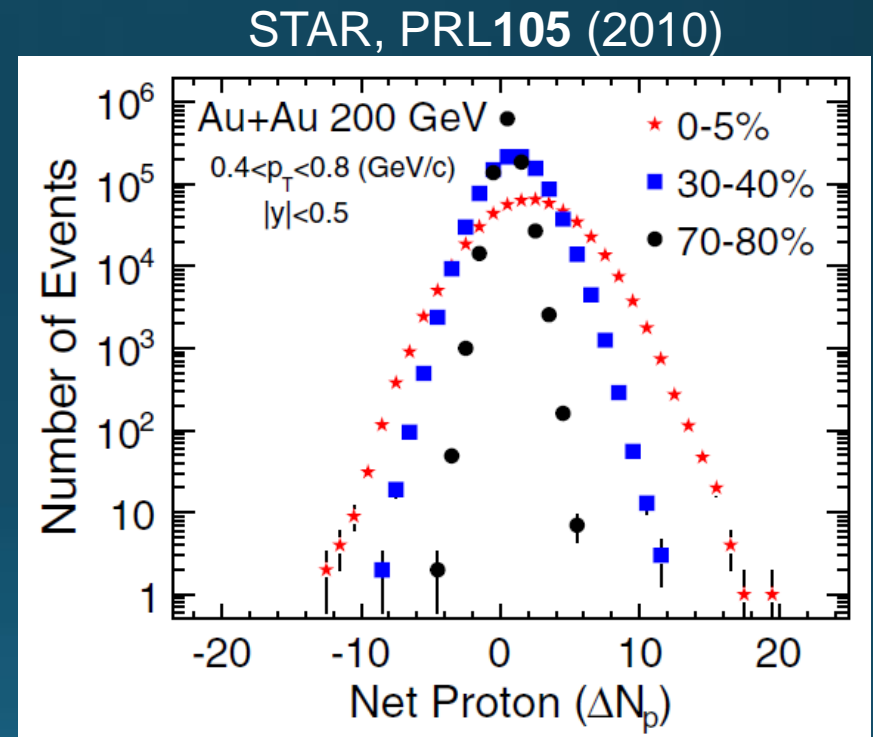
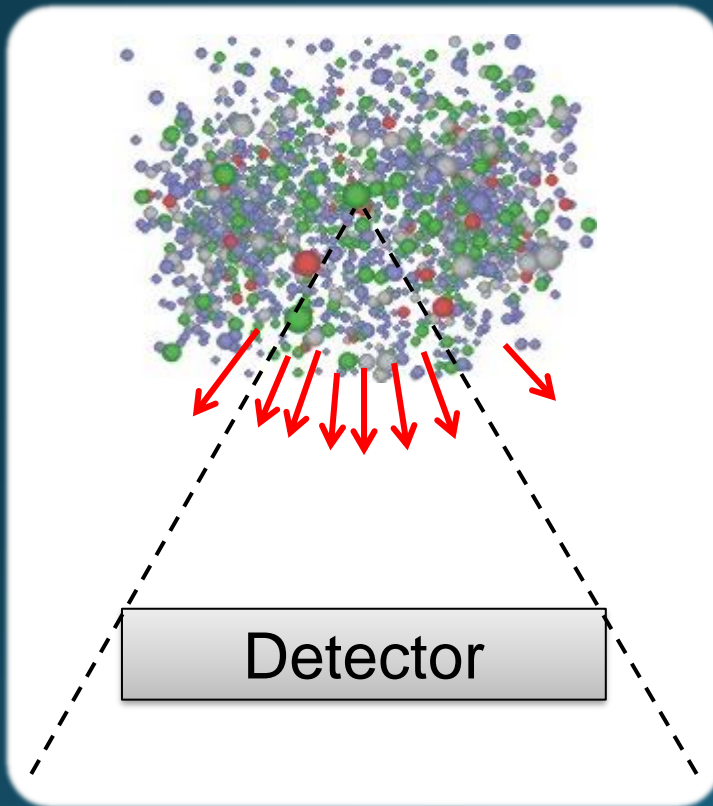
1. Search for **QCD Critical Point**  
using fluctuation observables

2. Search for **CSC Phase Transition**  
using dilepton production rates

Nishimura, MK, Kunihiro, arXiv:2201.01963

# Event-by-Event Fluctuations

Review: Asakawa, MK, PPNP 90 (2016)



Cumulants

$$\langle \delta N_p^2 \rangle, \langle \delta N_p^3 \rangle, \langle \delta N_p^4 \rangle_c$$

# Non-Gaussian Cumulants

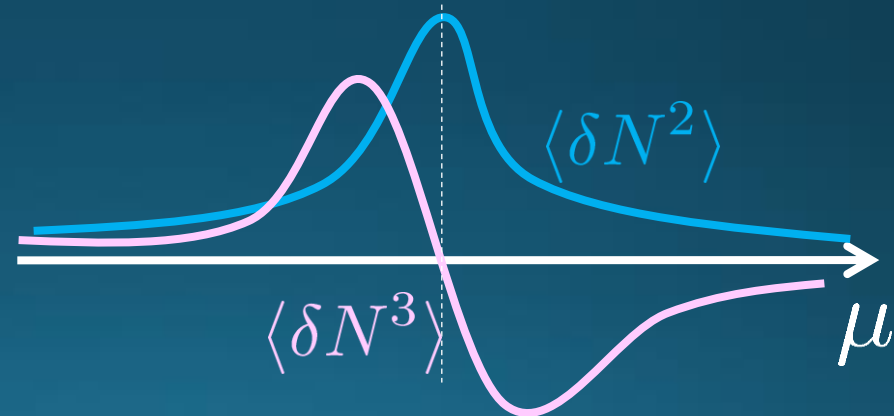
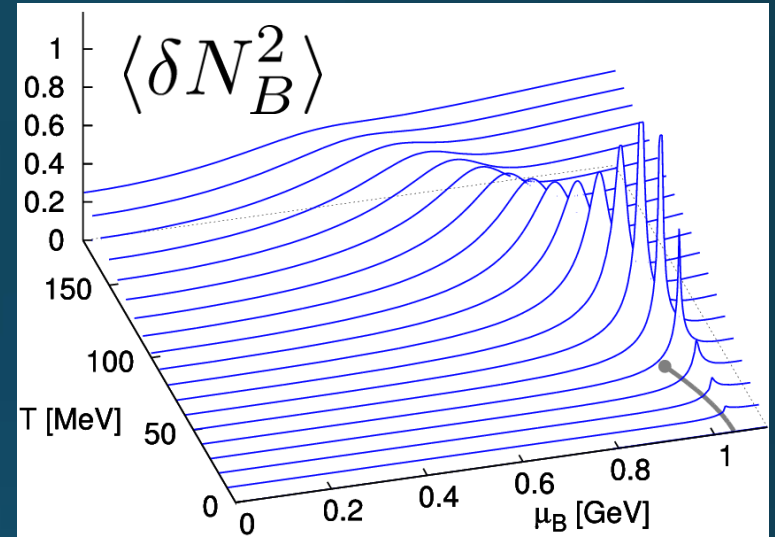
Gaussian fluctuations diverge at the QCD-CP



- Higher order cumulants change sign at the phase boundary

$$\langle \delta N^3 \rangle = T \frac{\partial \langle \delta N^2 \rangle}{\partial \mu}$$

Asakawa, Ejiri, MK, 2009



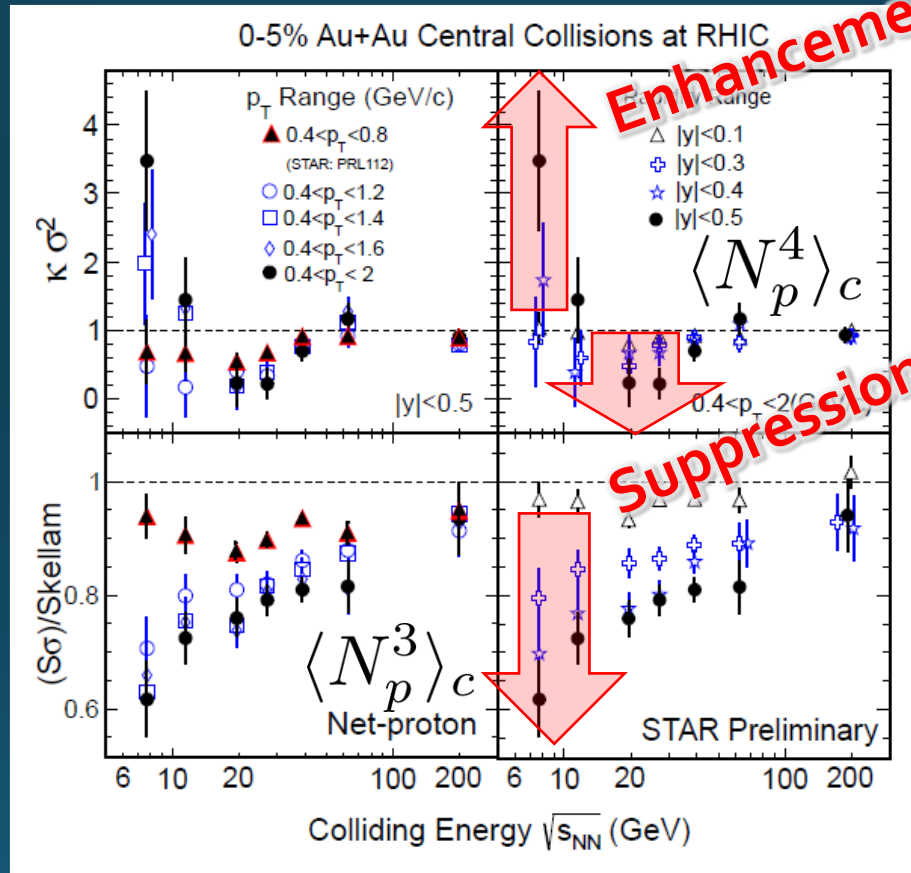
- Steeper divergence for higher-order cumulants

Stephanov, 2009

# Experimental Results

$$\frac{\langle N_p^4 \rangle_c}{\langle N_p^2 \rangle_c}$$

$$\frac{\langle N_p^3 \rangle_c}{\langle N_p^2 \rangle_c}$$



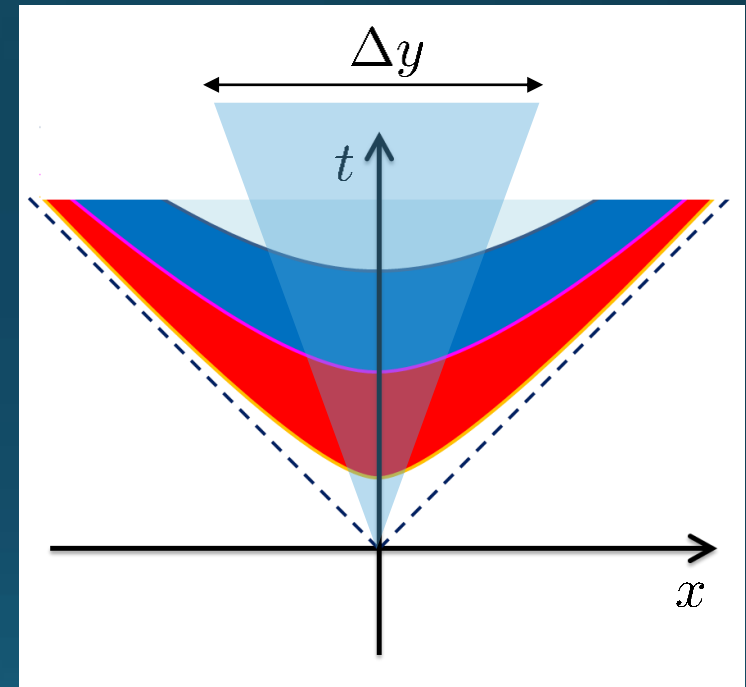
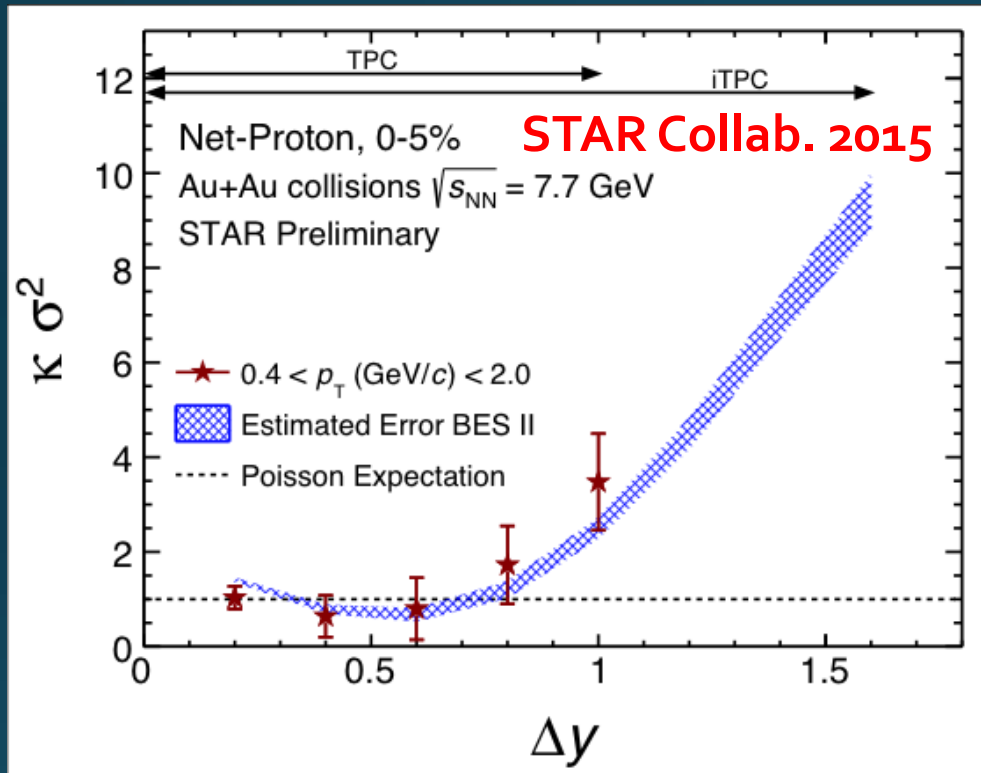
STAR Collab.  
2010~

**Enhancement & Suppression**  
of non-Gaussian cumulants!



Have we observed  
QCD critical point?

# Rapidity Window Dependence

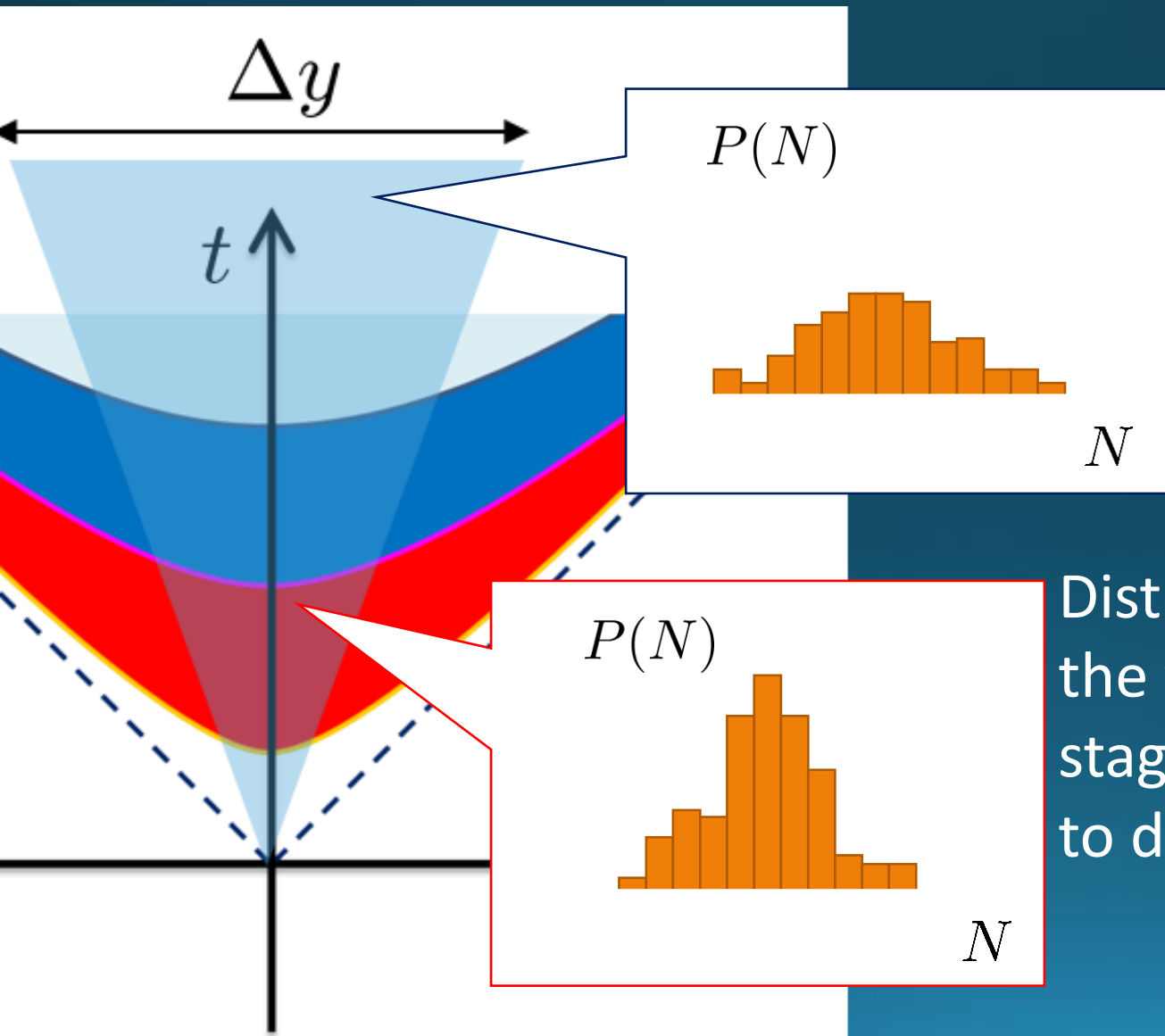


- Non-Gaussian Cumulants have been observed as a function of rapidity window  $\Delta y$ .
- Some results have non-monotonic  $\Delta y$  dependence.



# Diffusion of Fluctuations

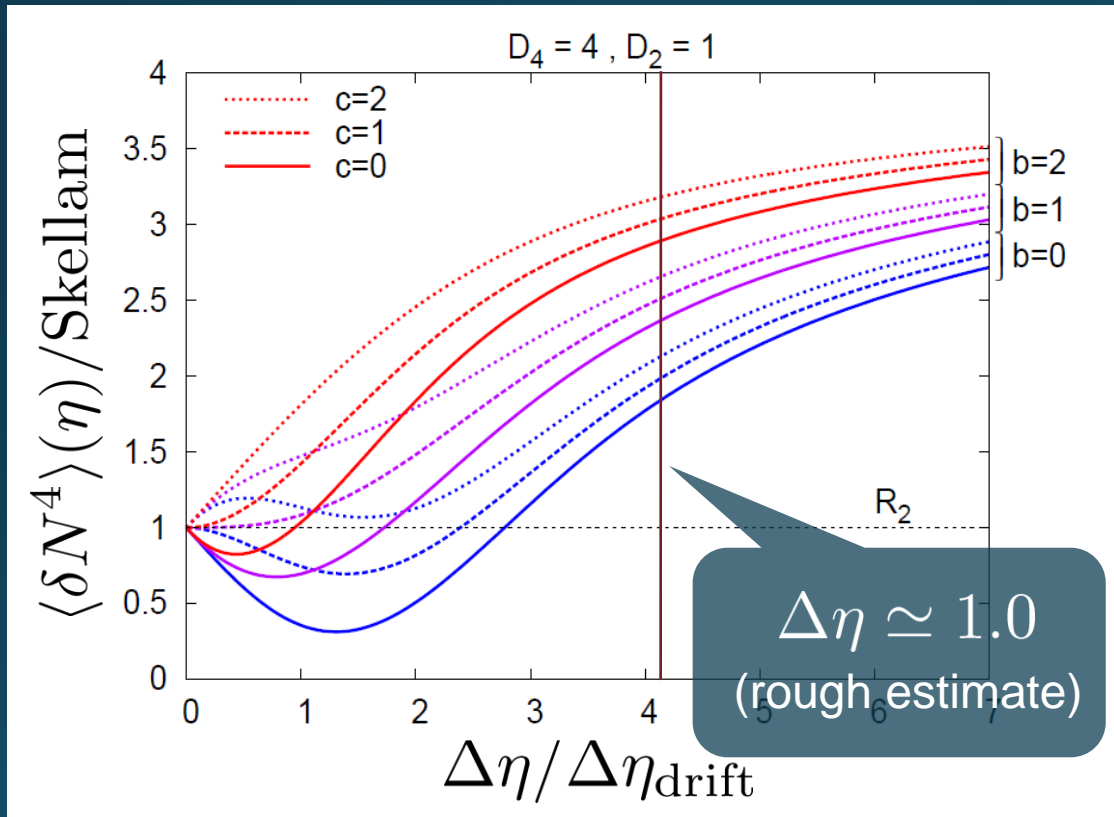
MK, Ohno, Asakawa 2014  
MK 2015



Distributions in  $\Delta y$  in the final state and early stage are different due to diffusion.

# Rapidity Window dependence as a Result of Diffusion

MK+ (2014); MK (2015)



## Parameters

$$D_4 = \frac{\langle Q_{(\text{net})}^4 \rangle_c}{\langle Q_{(\text{tot})} \rangle} = 4$$

$$D_2 = \frac{\langle Q_{(\text{net})}^2 \rangle_c}{\langle Q_{(\text{tot})} \rangle} = 1$$

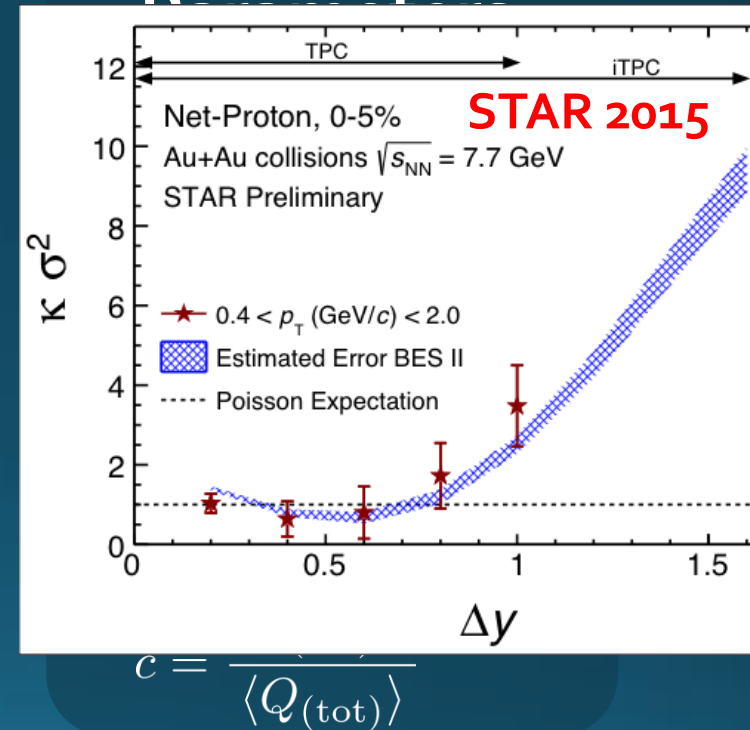
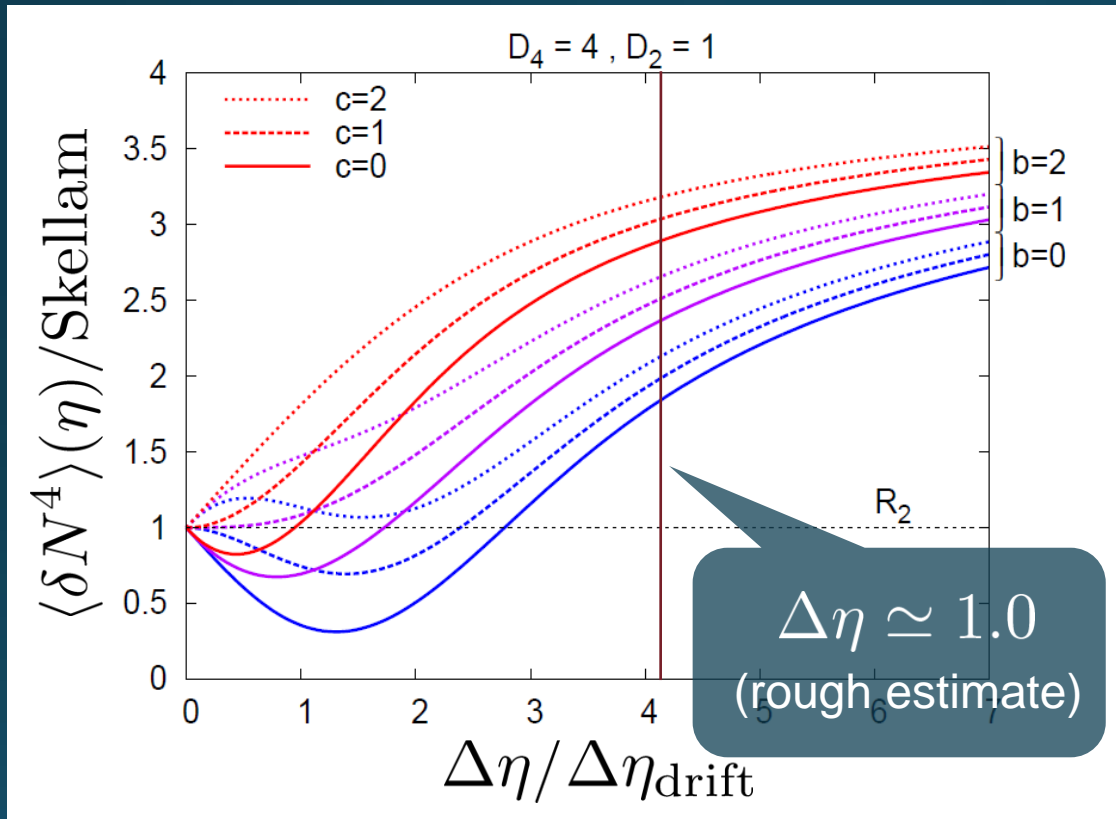
$$b = \frac{\langle Q_{(\text{net})}^2 Q_{(\text{tot})} \rangle_c}{\langle Q_{(\text{net})} \rangle}$$

$$c = \frac{\langle Q_{(\text{tot})}^2 \rangle_c}{\langle Q_{(\text{tot})} \rangle}$$

- Higher order cumulants can behave non-monotonically.
- $\Delta\eta$  dependence encodes history of time evolution.

# Rapidity Window dependence as a Result of Diffusion

MK+ (2014); MK (2015)



- Higher order cumulants can behave non-monotonically.
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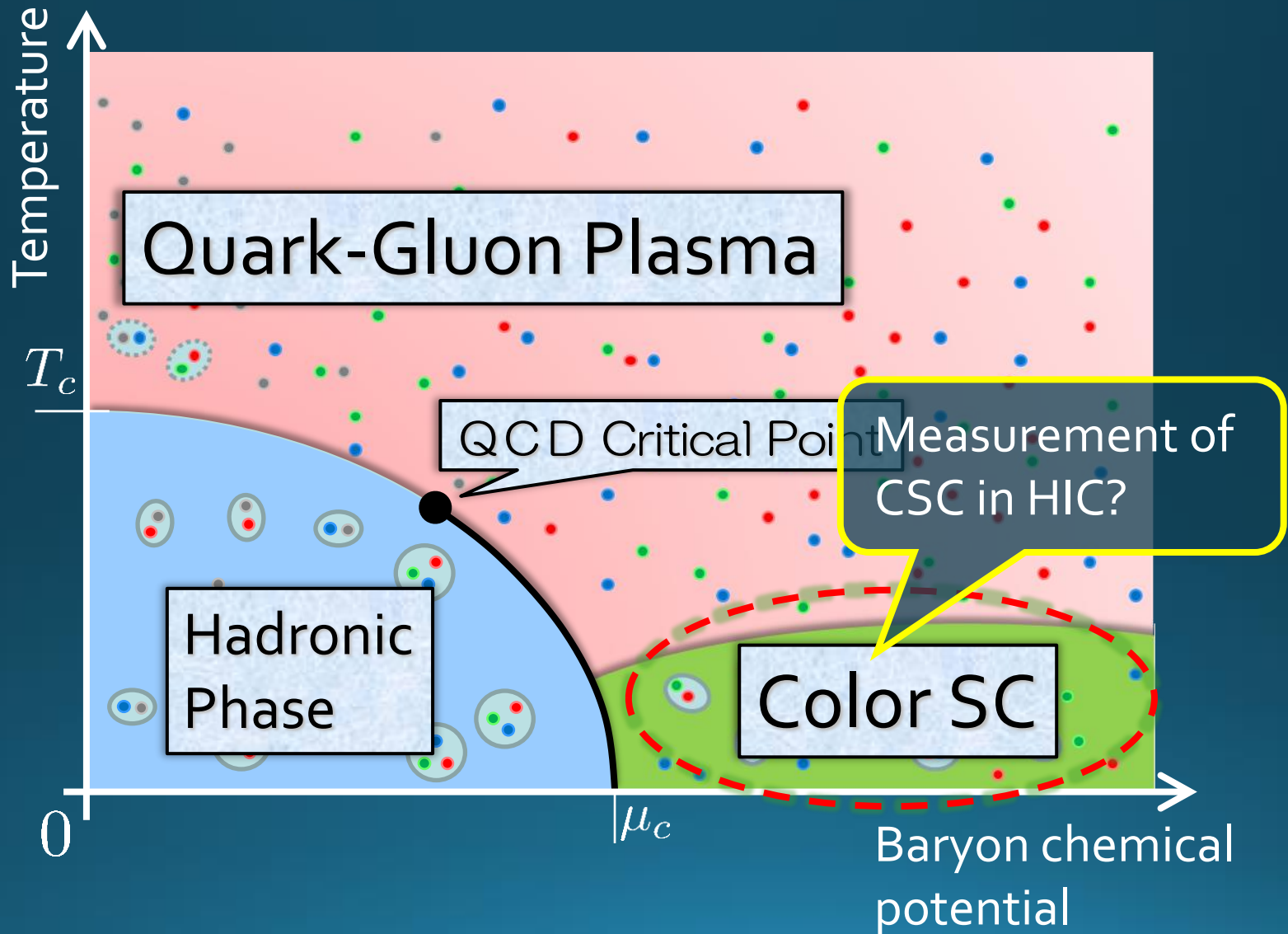
# Contents

1. Search for **QCD Critical Point**  
using fluctuation observables

2. Search for **CSC Phase Transition**  
using dilepton production rates

Nishimura, MK, Kunihiro, arXiv:2201.01963

# Color Superconductivity



# Observing CSC in HIC?

## Difficulties

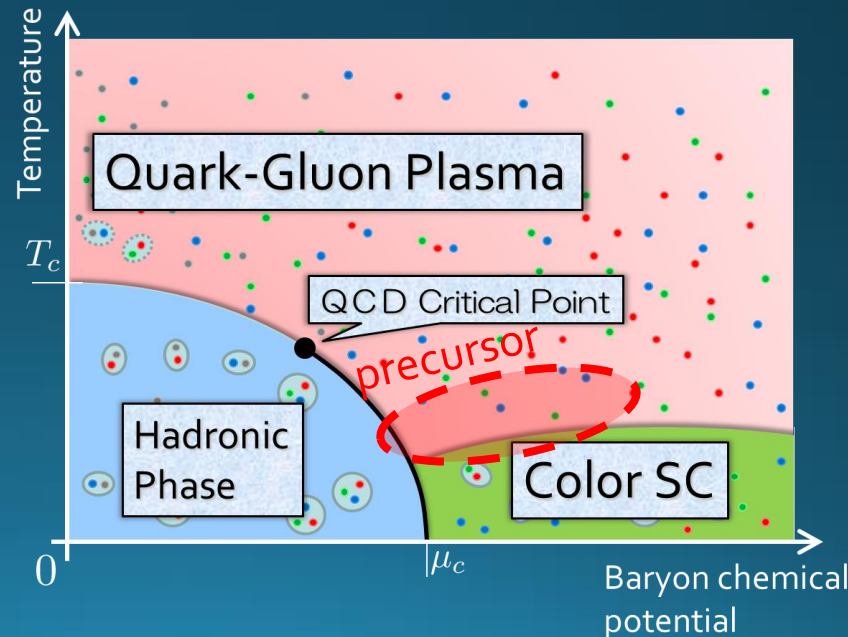
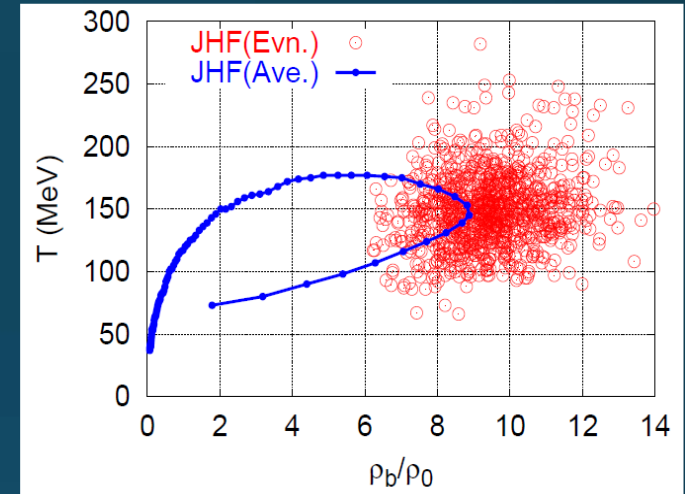
- 1) Creation of CSC itself would be impossible.
- 2) CSC would be realized only in the early stage.



## Solution

- 1) Focus on **precursor of CSC**
- 2) Use **dilepton production**

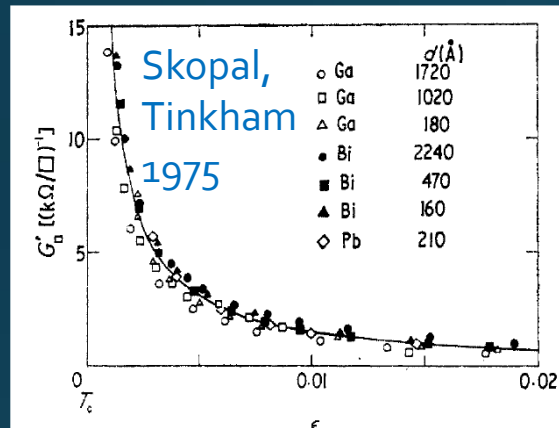
Nishimura, MK, Kunihiro, 2201.01963



# Precursory Phenomena

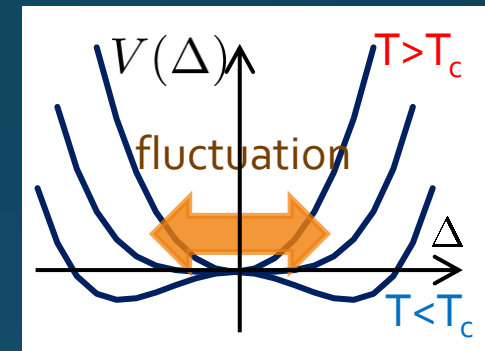
= anomalous behavior of observables near but above  $T_c$

## Electric conductivity in metals



$$\epsilon = \frac{T - T_c}{T_c}$$

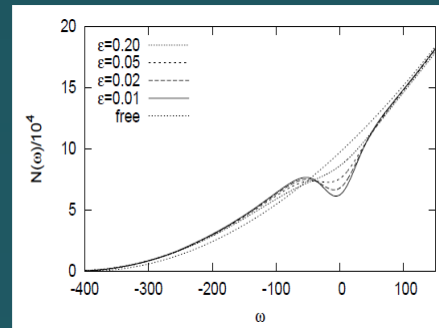
## Landau's Free Energy



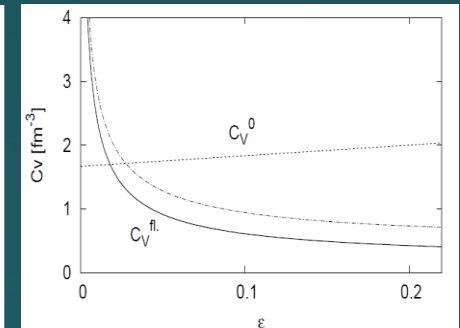
## Precursor of CSC

- Pseudogap = Depression in the density of states
- Specific heat
- etc.

## Density of States



## Specific Heat



MK, Koide, Kunihiro, Nemoto, '03, '05

# Model

## NJL model (massless 2-flavor)

$$\mathcal{L} = \bar{\psi} i \not{\partial} \psi + \mathcal{L}_S + \mathcal{L}_C$$

$$\mathcal{L}_S = G_S ((\bar{\psi}\psi)^2 + (\bar{\psi}i\gamma_5\tau\psi)^2)$$

$$\mathcal{L}_C = G_C ((\bar{\psi}i\gamma_5\tau_A\lambda_A\psi^C)(\text{h.c.}))$$

diquark interaction

$$G_S = 5.01 \text{ GeV}^{-2}, \quad \Lambda = 650 \text{ MeV}$$

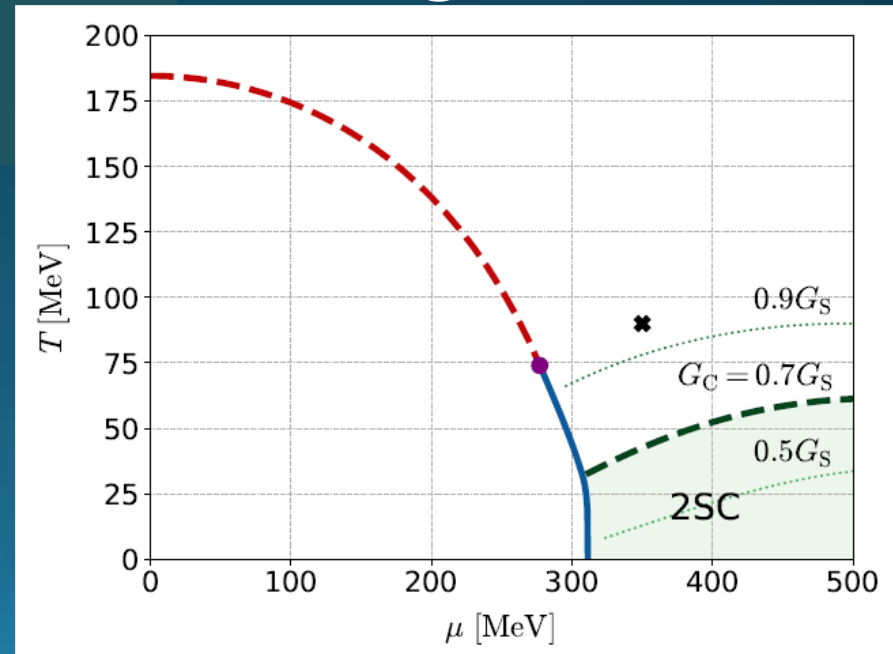
□  $G_C$ : free parameter

Order of CSC phase transition

Matsuura+'04), Giannakis+'04)

Noronha+'06), Fejos, Yamamoto('19)

## Phase Diagram in MFA






# Diquark Mode

## Diquark Propagator (T-matrix)

in random-phase approximation

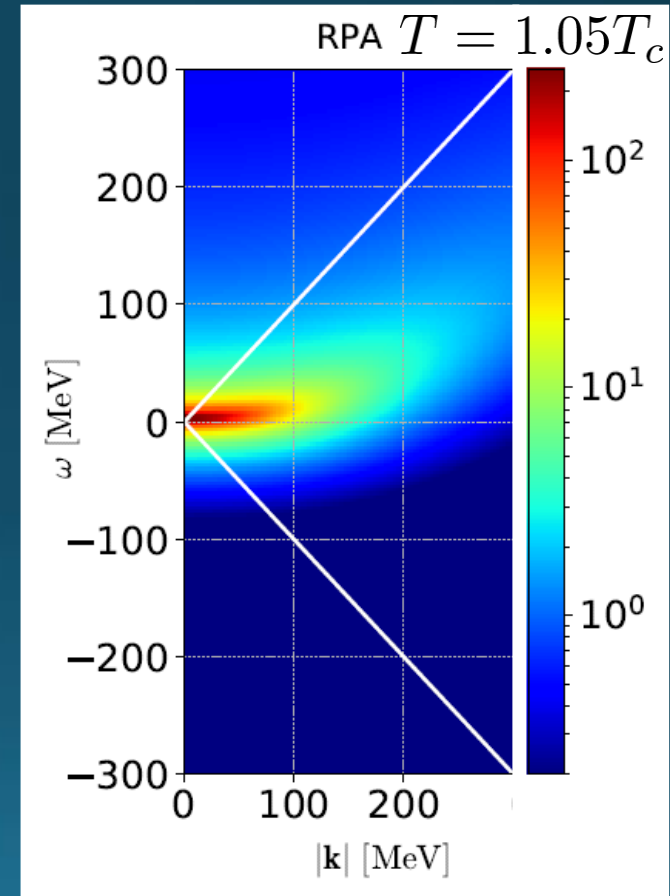
$$\Xi^R(\mathbf{k}, \omega) =$$


The diagram shows a series of bubble diagrams representing the Dyson equation for the diquark propagator. The first term is a single bubble. The second term is two bubbles connected in series. The third term is three bubbles connected in series, followed by an ellipsis indicating the continuation of the series.

$$= \frac{Q^R(\mathbf{k}, \omega)}{1 + G_C Q^R(\mathbf{k}, \omega)}$$

- Massless at  $T=T_c$  as a soft mode of CSC transition
- Strength in the space-like region

## Dynamical Structure Factor

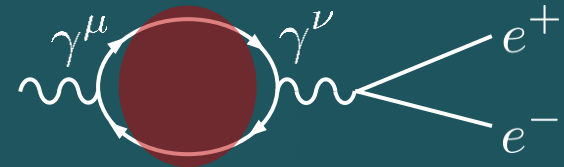


# Photon Self-Energy

## Dilepton Production Rate

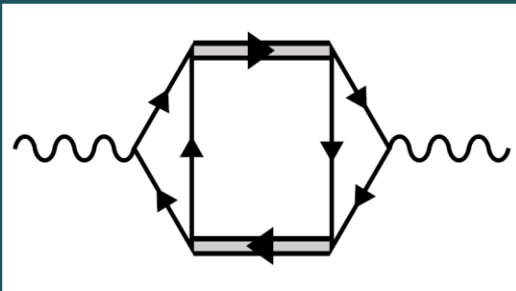
$$\frac{d^4\Gamma}{dk^4} = \frac{\alpha}{12\pi^4} \frac{1}{k^2} \frac{1}{e^{\beta\omega}-1} \text{Im}\Pi_{\mu}^{R\mu}(k)$$

$$\Pi^{\mu\nu}(k) =$$

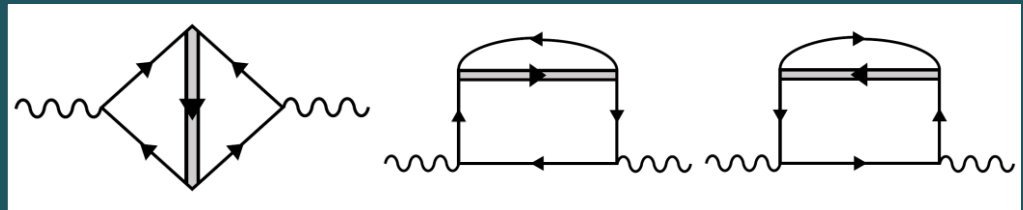


## Terms included in $\Pi^{\mu\nu}$

Aslamasov-Larkin



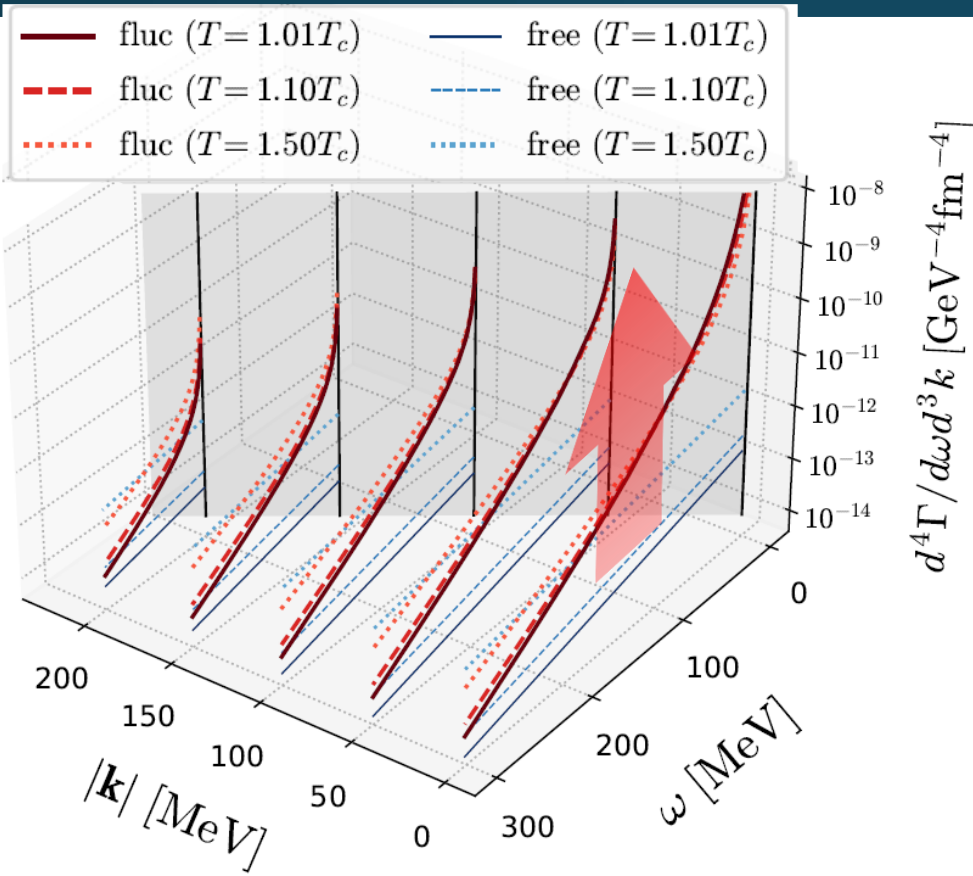
Maki-Thompson, Density of states



$\Rightarrow$  = diquark field

- common in metallic superconductors (conductivity)
- time-dependent GL approx. for diquark field
- gauge-invariant construction

# Dilepton Production Rates



$$\frac{d^4\Gamma}{dk^4}(\omega, k)$$

**Red:** fluctuation contribution  
**Blue:** free quarks

$$\mu = 350 \text{ MeV}$$

$$G_C = 0.7G_S, \quad (T_c \simeq 43 \text{ MeV})$$

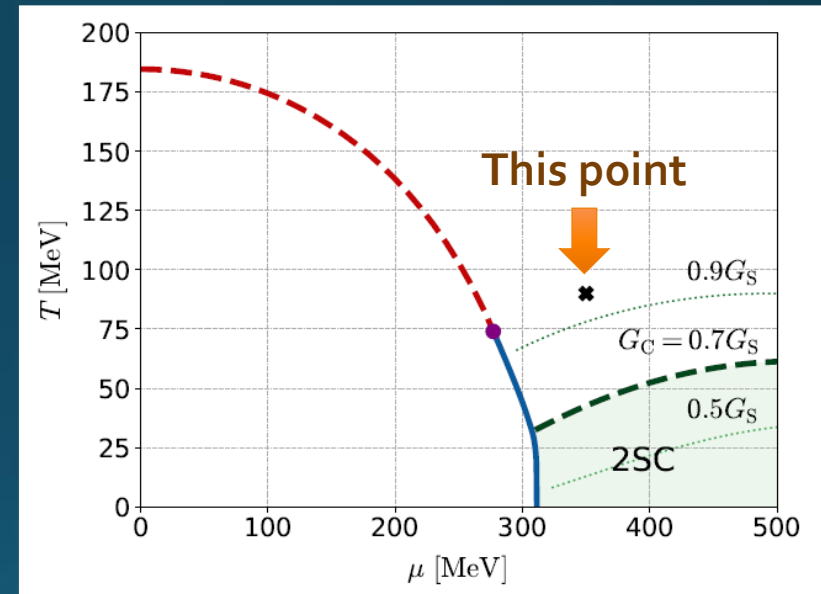
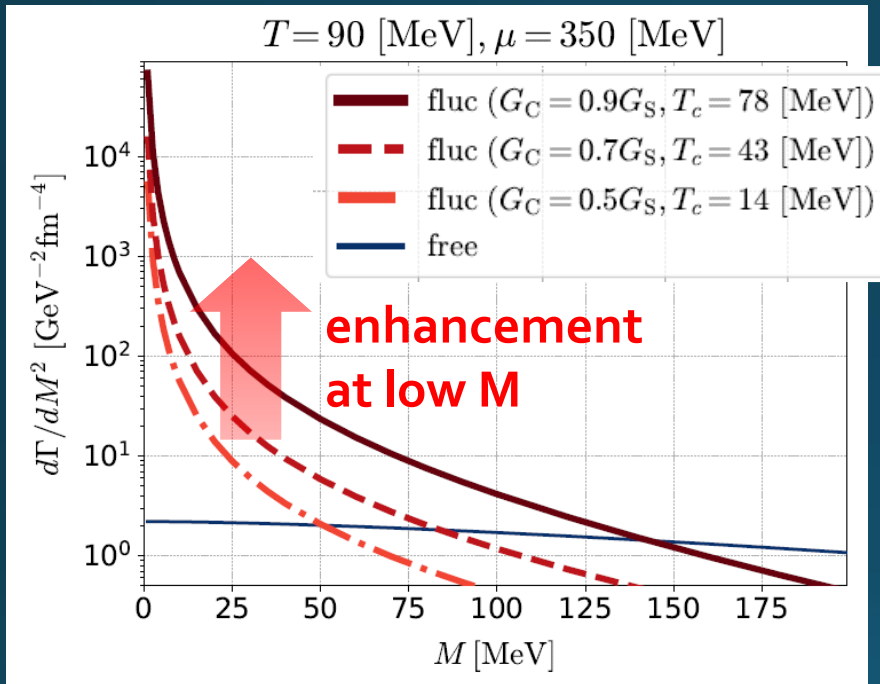
Nishimura, MK, Kunihiro,  
 2201.01963

- Diquark fluctuations give rise to anomalous enhancement in the low energy-momentum region for  $T < 1.5T_c$ .

# Invariant-Mass Spectrum

## Fixed Temperature

$$\frac{d\Gamma}{dM^2} = \int \frac{d^3k}{2\omega} \frac{d^4\Gamma}{dk^4} \Big|_{\omega=E_k}$$

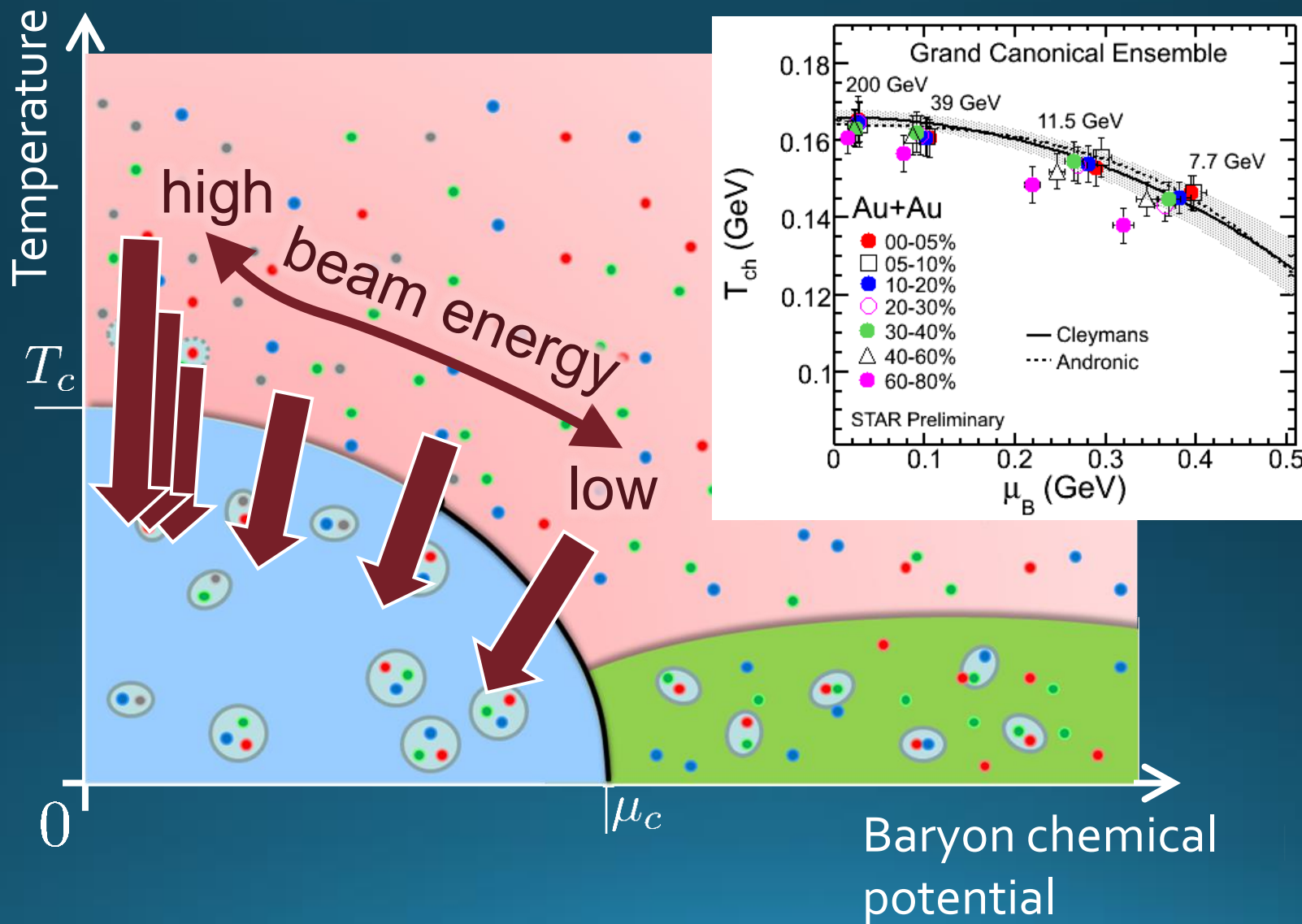


- Strong enhancement at low invariant mass, though the range of  $M$  is narrower than the previous results.
- Observable in the HIC?

# Summary

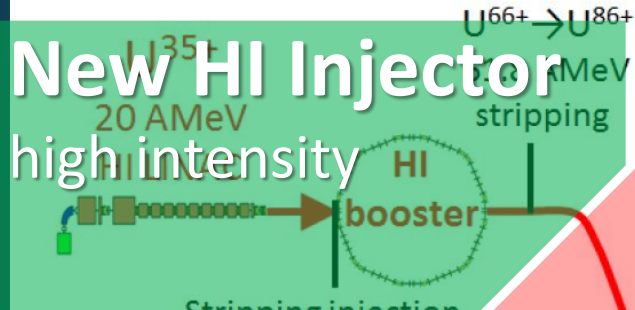
- ❑ Exploring dense medium in relativistic heavy-ion collisions is one of the hottest topics in this field. The beam-energy scan is ongoing, and many new experiments will start in near future!
- ❑ Among them, search for
  - ❑ the **QCD critical point** using fluctuation observables
  - ❑ the **color superconductivity** using dileptonsare especially interesting and important subjects.
- ❑ Heavy-ion collisions at J-PARC will play important roles in pursuing these subjects.

# Beam-Energy Scan



# HI Acceleration @ J-PARC

## New HI Injector



Stripping injection  
 U<sup>35+</sup> → U<sup>66+</sup>  
 20 → 67 A MeV

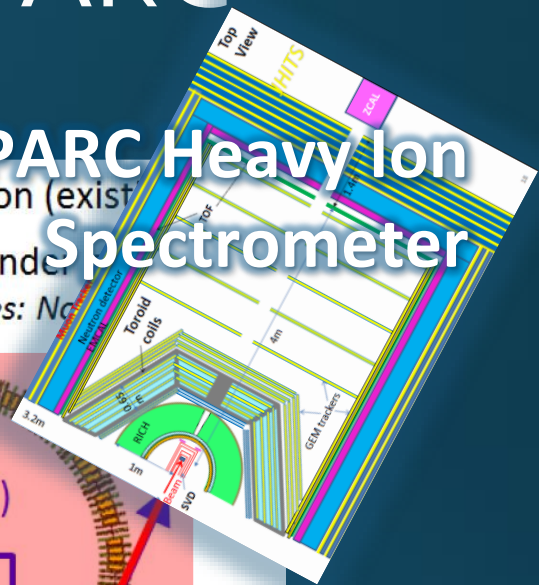
H<sup>-</sup> Linac: 0.4 GeV

61.8 → 735.4 A MeV

RCS & Main Ring  
 stable well established

0.727 → 11.15 A GeV

J-PARC Heavy Ion Spectrometer  
 Figures: No



- proton (exist)
- HI (under)

- Use of reliable / high-performance RCS & main ring
- → Reduce cost and time