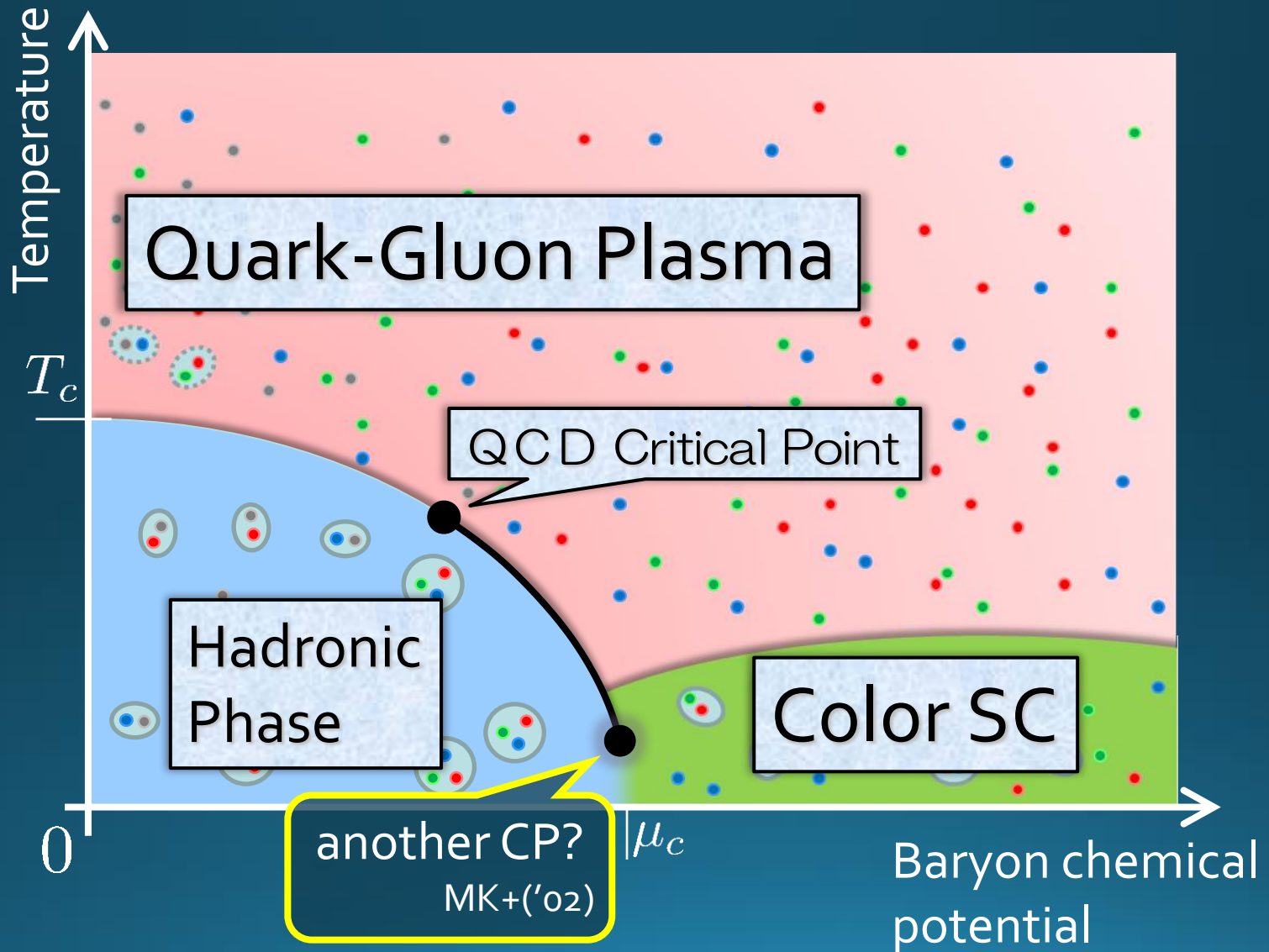


Exploring Phases of Dense QCD in Heavy-Ion Collisions

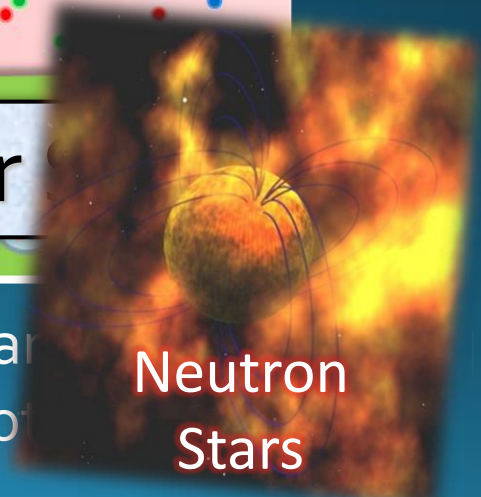
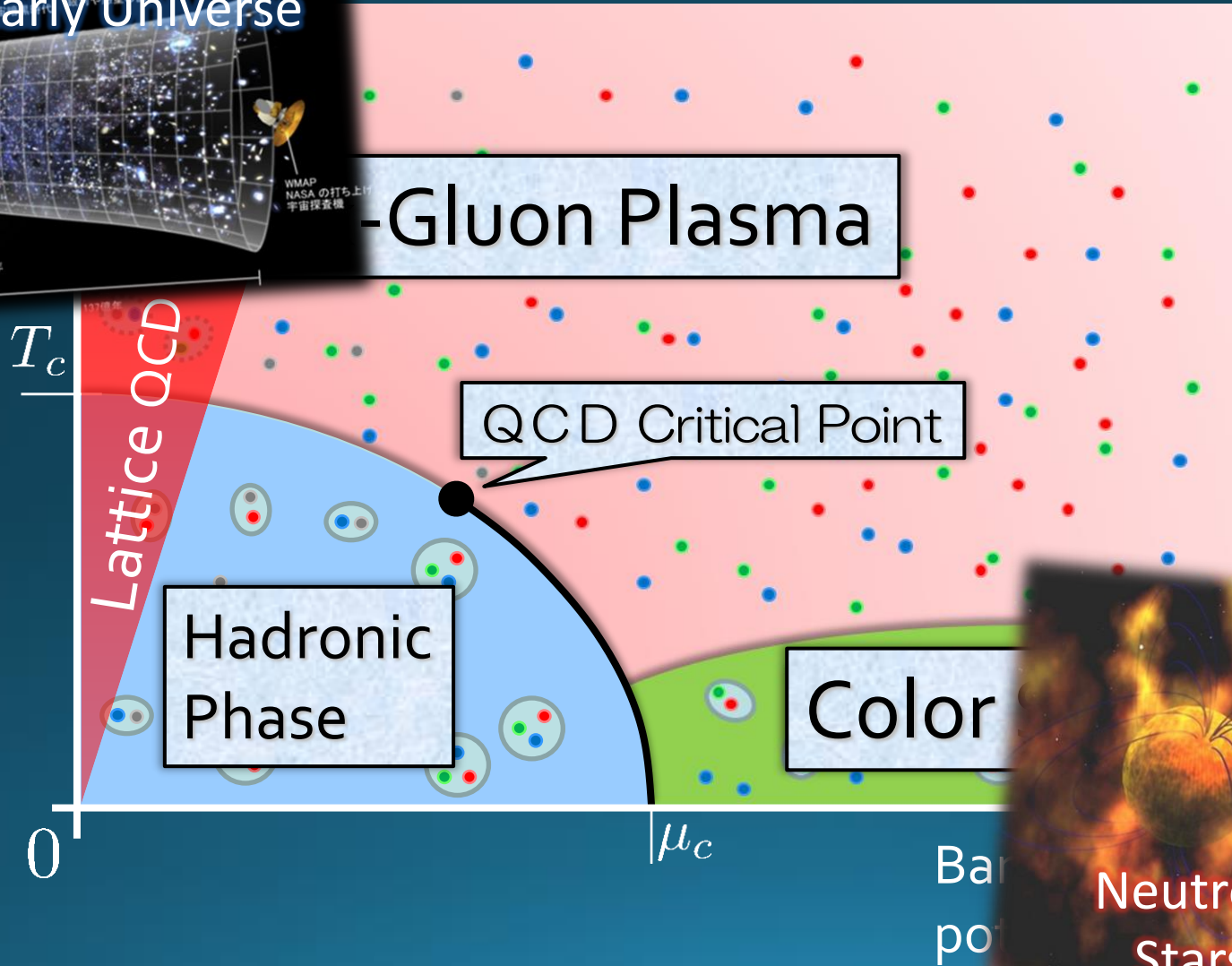
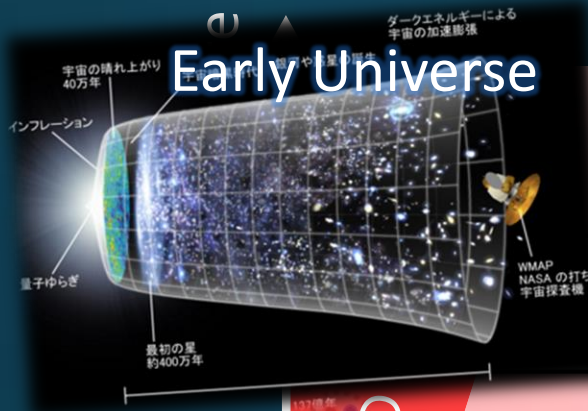
Masakiyo Kitazawa
(Osaka U.)

Reimei Workshop "Hadrons in dense matter at J-PARC"
KEK, Tokai, 2022/Feb./23

QCD Phase Diagram

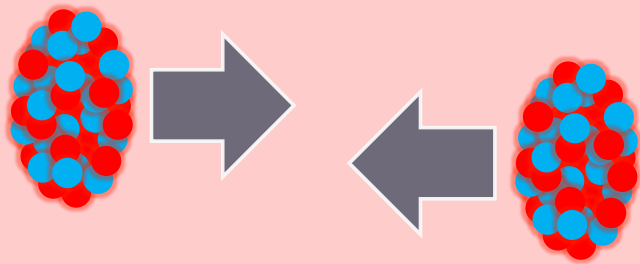


QCD Phase Diagram



Relativistic Heavy-Ion Collisions

Collide 2 heavy nuclei



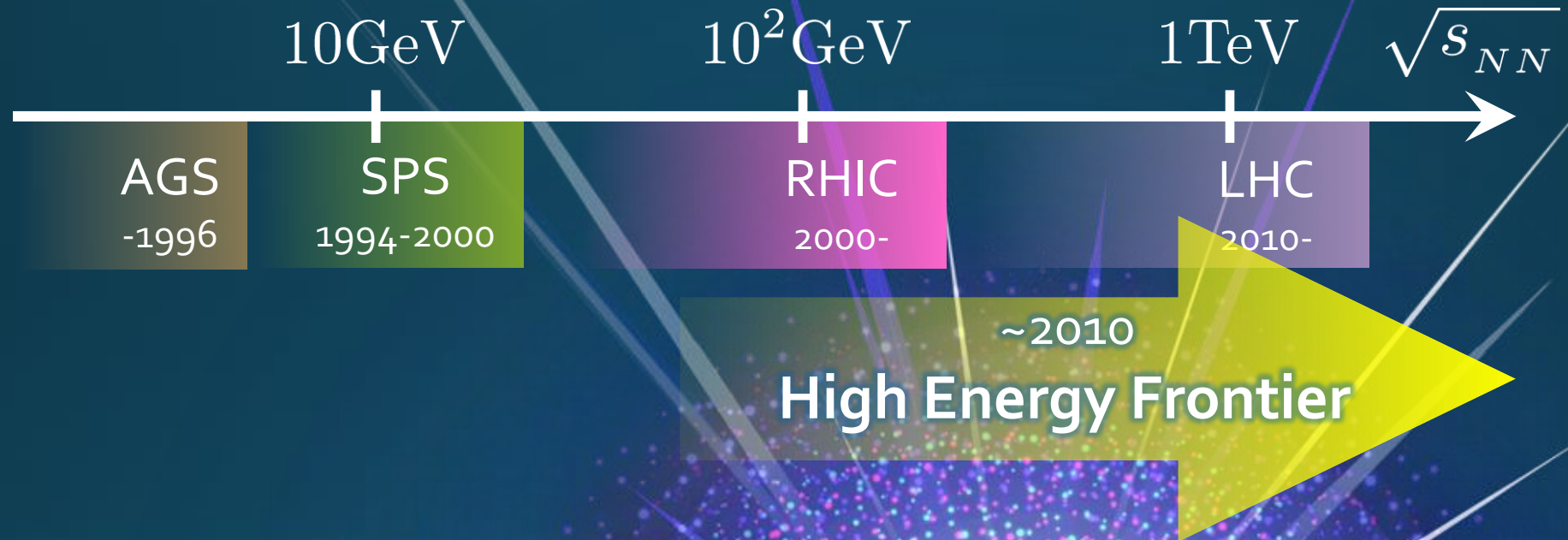
RHIC (2000~)
QGP Formation
Strongly coupled QGP

LHC (2010~)
Precision measurement
of the QGP

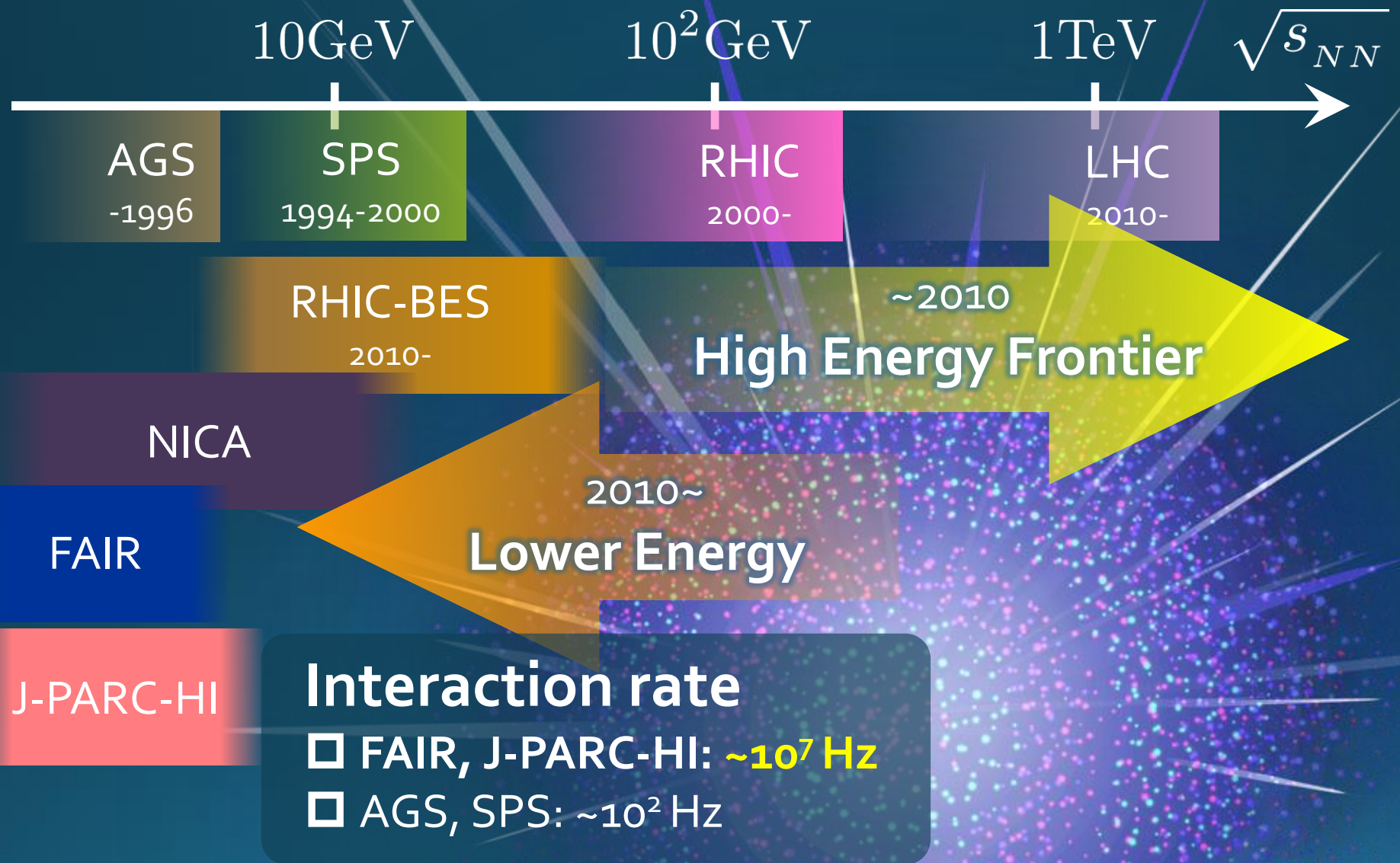
Physics

- ❑ Hot & dense medium
- ❑ Early Universe
- ❑ Quark-gluon plasma
- ❑ QCD phase structure

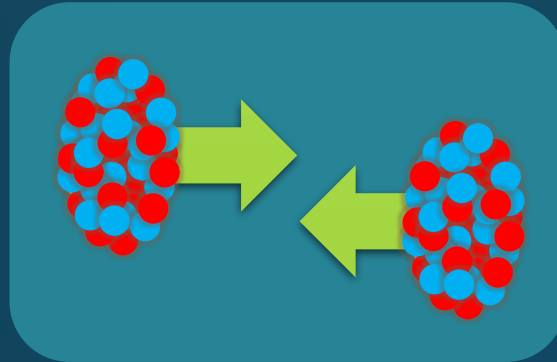
Brief History of Relativistic HIC



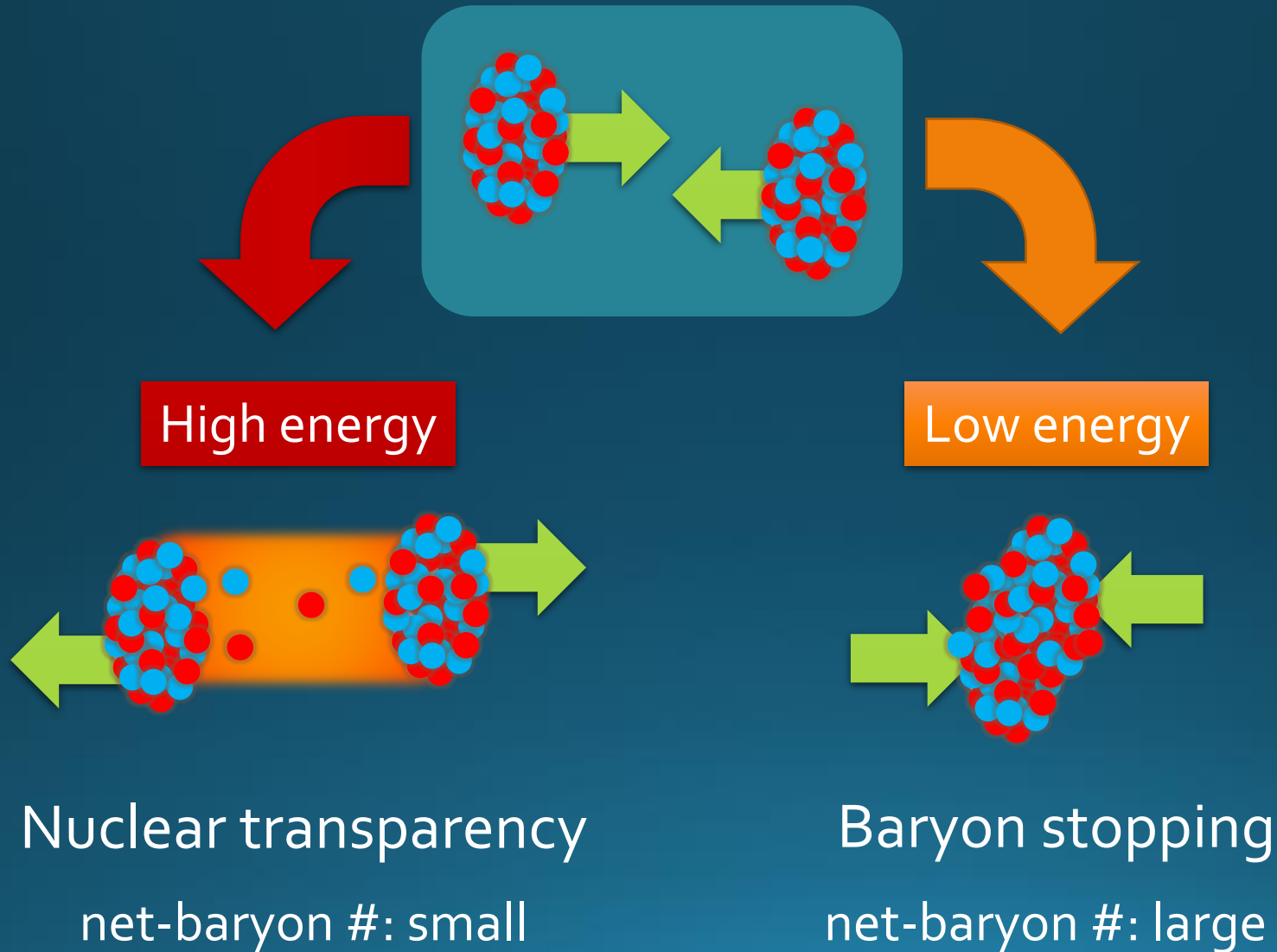
Brief History of Relativistic HIC



Beam-Energy Dependence

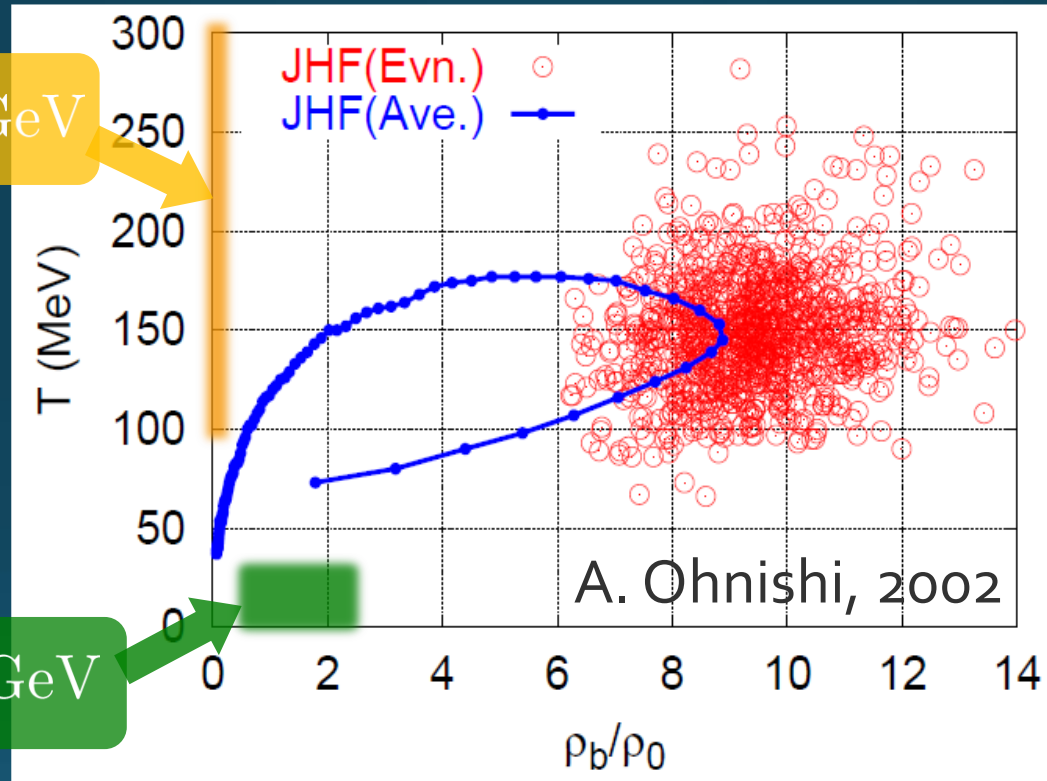


Beam-Energy Dependence



Maximum Density

Time evolution in T - ρ plane by JAM



$$E/A = 20 \text{ GeV}$$

$$\sqrt{s_{NN}} \simeq 6 \text{ GeV}$$

- Maximum density $5 \sim 10\rho_0$ @ $E/A \sim 20 \text{ GeV}$
- Large event-by-event fluctuations?

First-order Transition at Highest Density in the Universe

QGP formation in HIC
Lattice QCD

○ Quark-Gluon
○ Plasma

LHC
RHIC

RHIC
BES - I · II

QCD-CP

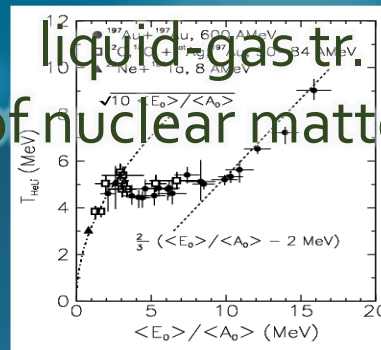
crossover

first-order
transition



boiling water

1g/cm^3



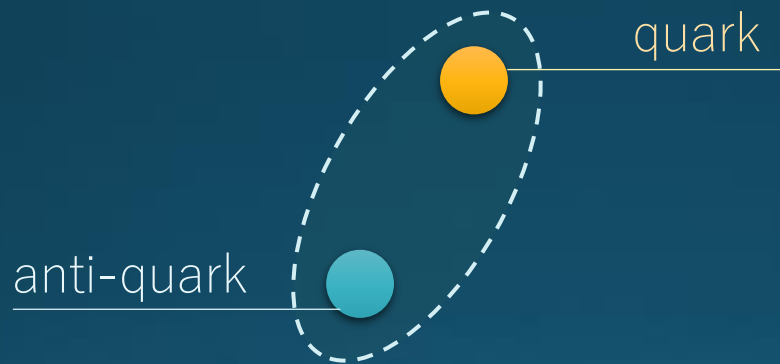
10^{14}g/cm^3

μ_B
Chiral
Transition

10^{15}g/cm^3

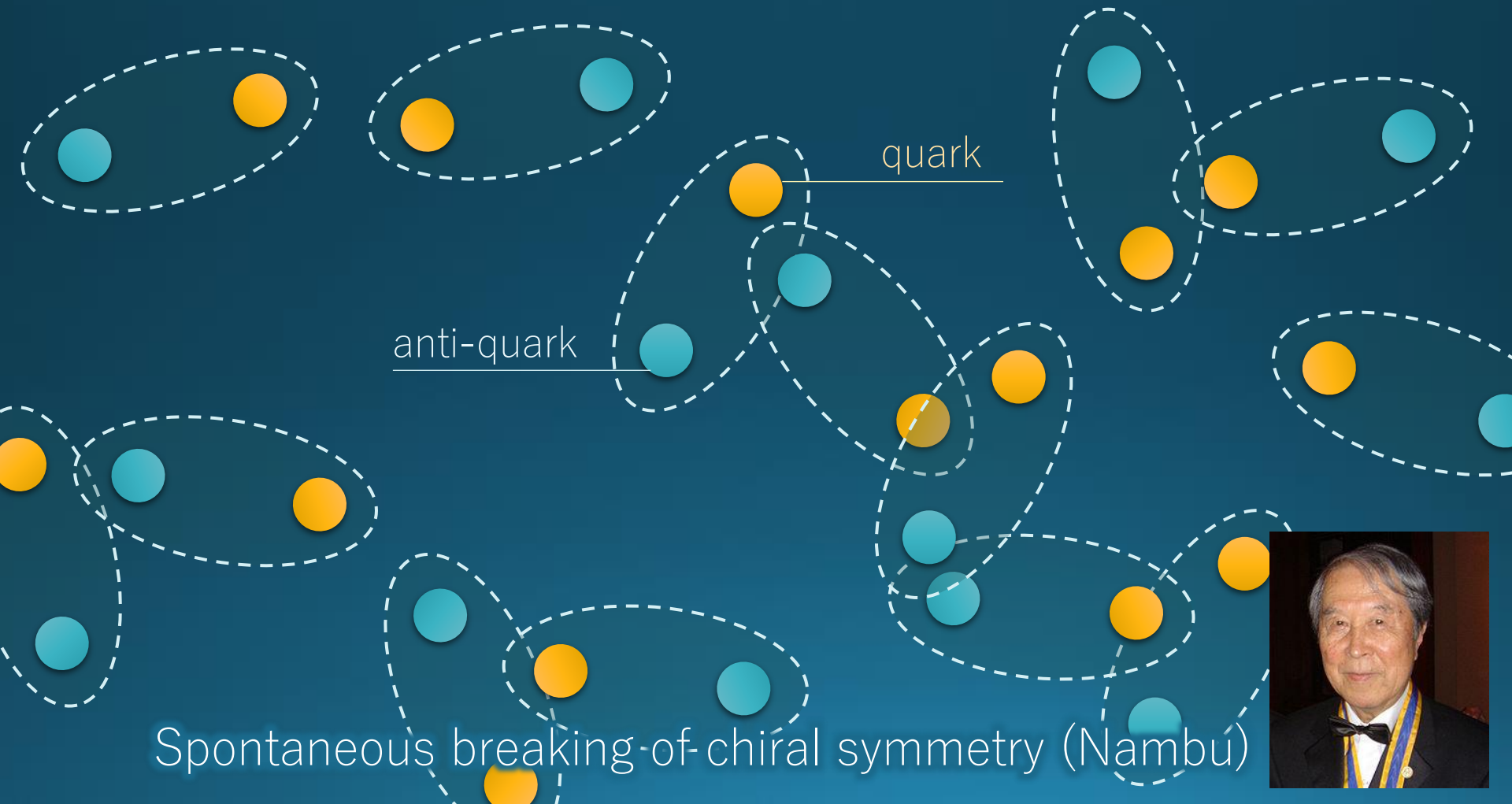
Collapse of Vacuum

Our vacuum is filled with the quark condensate.

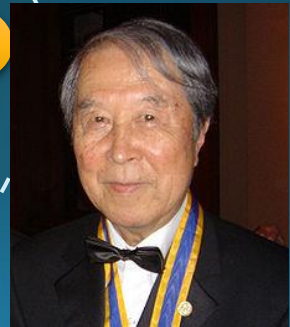


Collapse of Vacuum

Our vacuum is filled with the quark condensate.

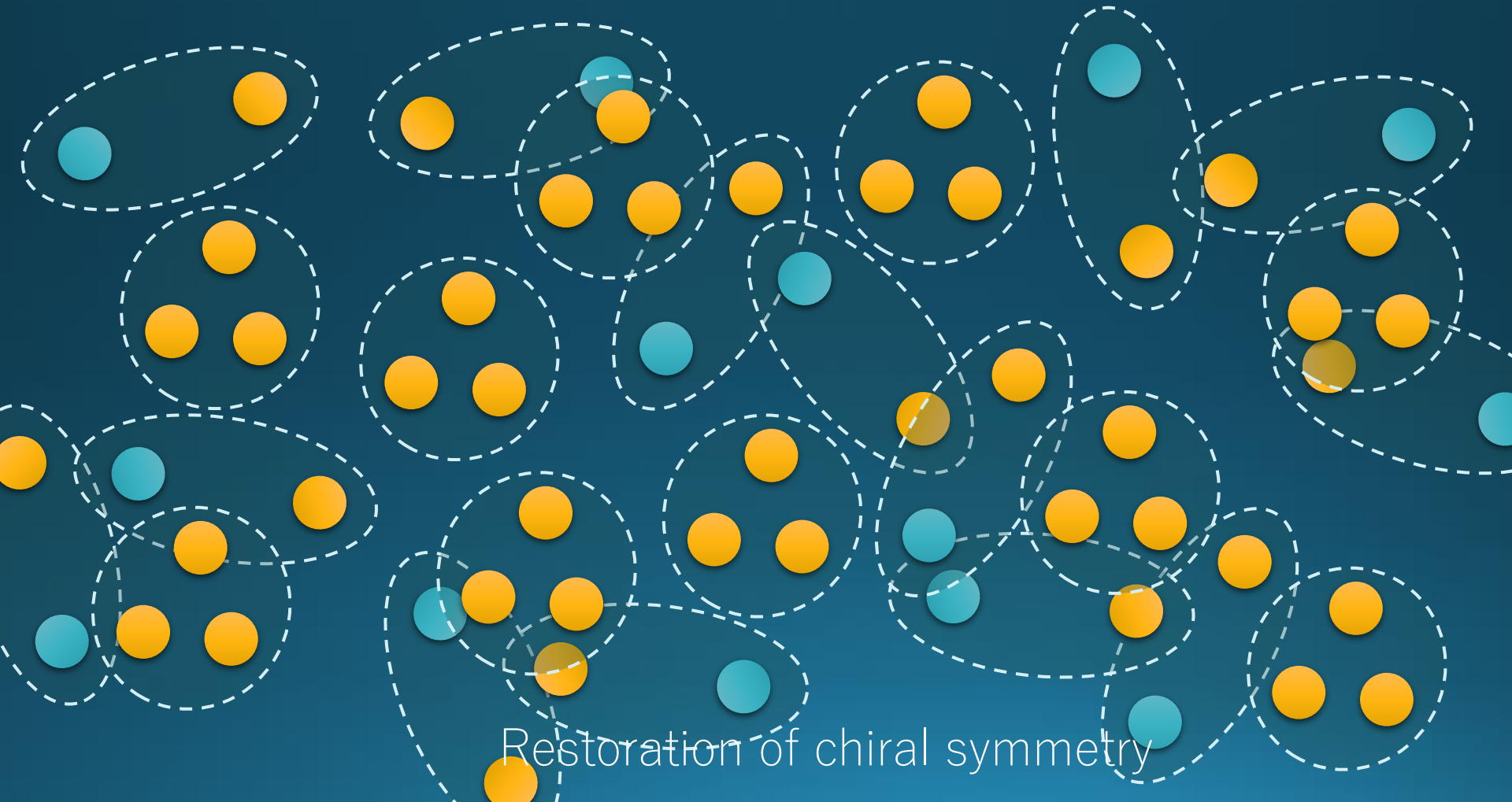


Spontaneous breaking of chiral symmetry (Nambu)



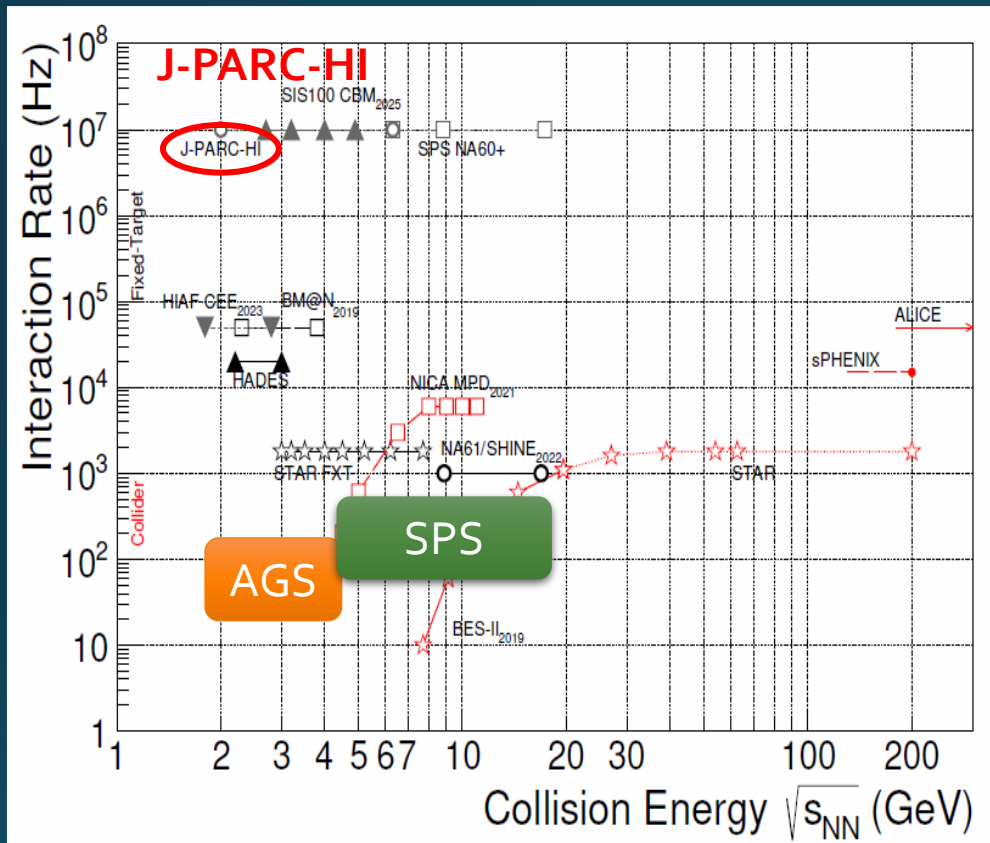
Collapse of Vacuum

Matter modifies the vacuum.



Restoration of chiral symmetry

Collision Rate



Galatyuk, NPA982,163 (2019)

J-PARC-HI:

High-luminosity x Fixed target
 → World highest rate $\sim 10^8$ Hz



5-order higher than AGS, SPS

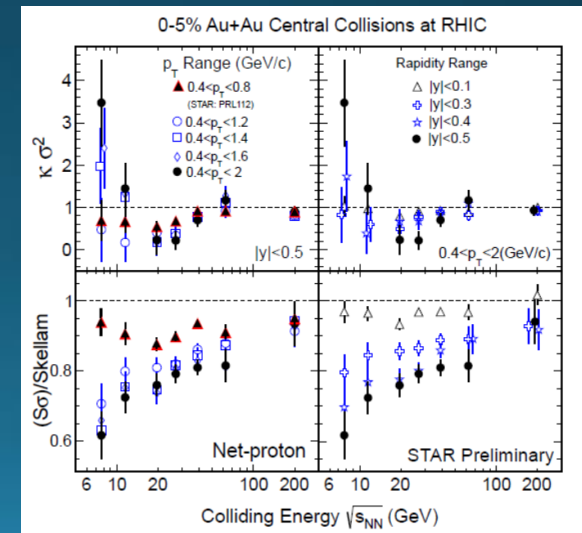
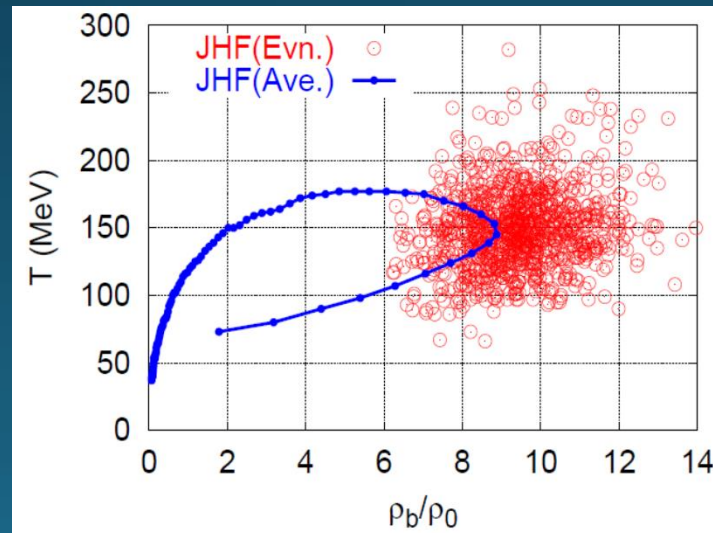
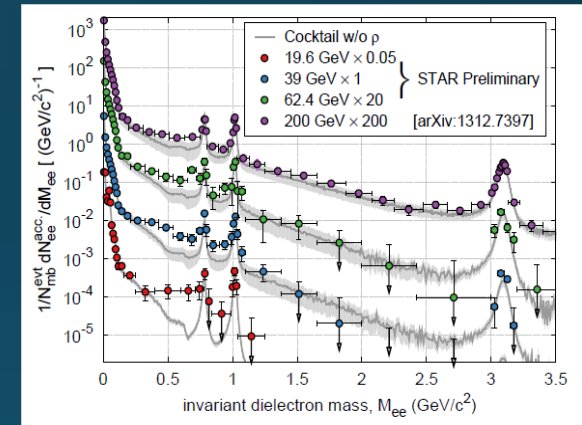
AGS, SPS = J-PARC-HI
 1 year = 5 min.



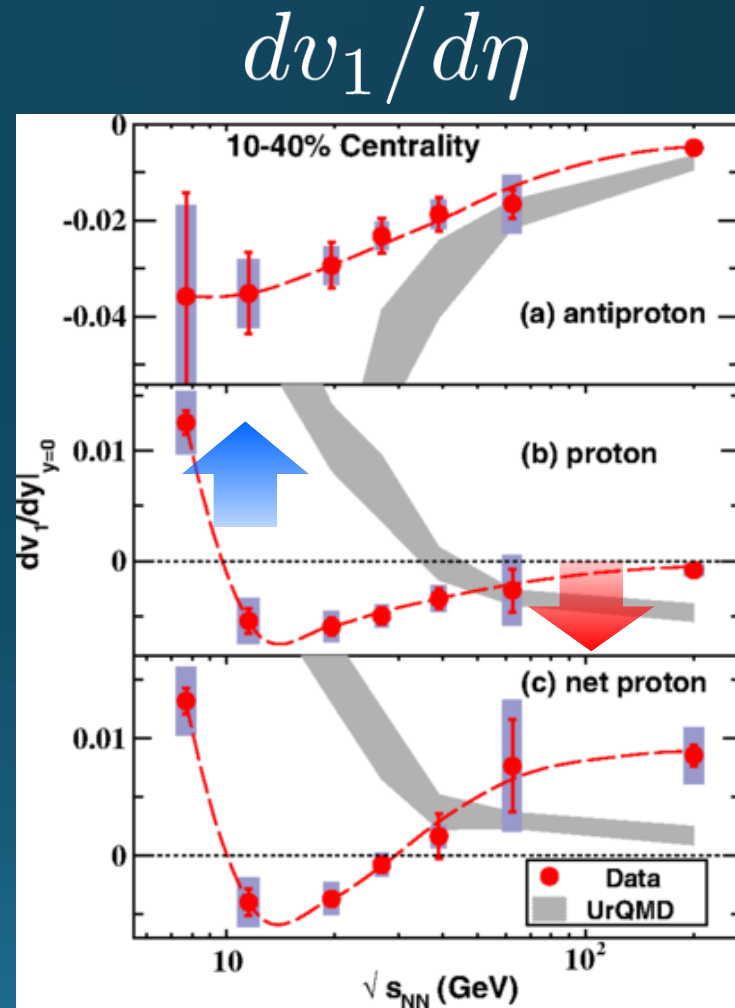
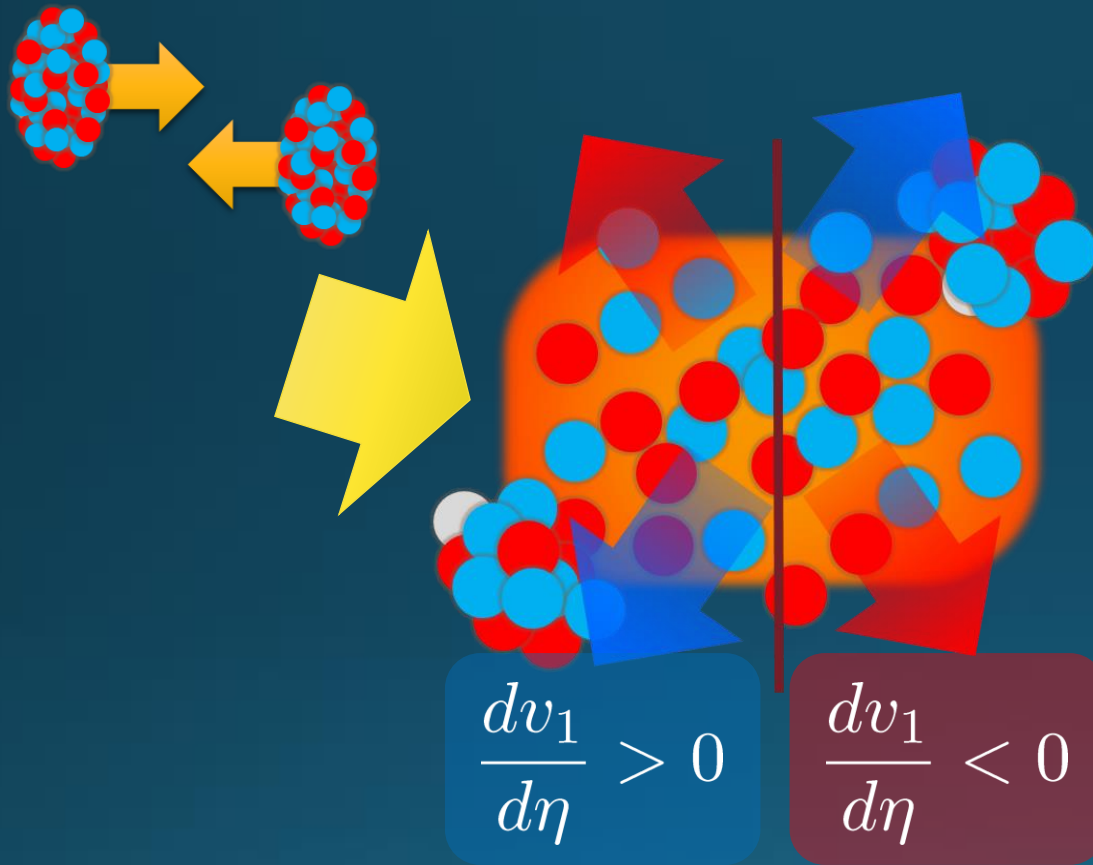
- High-statistical exp.
- various event selections
- higher order correlations
- search for rare events

Various Observables

- Flow
- Dilepton / photon
- Fluctuations, higher-order cumulants
- Ξ, Ω, \dots
- Sophisticated event selections
- Various correlations



Directed Flow

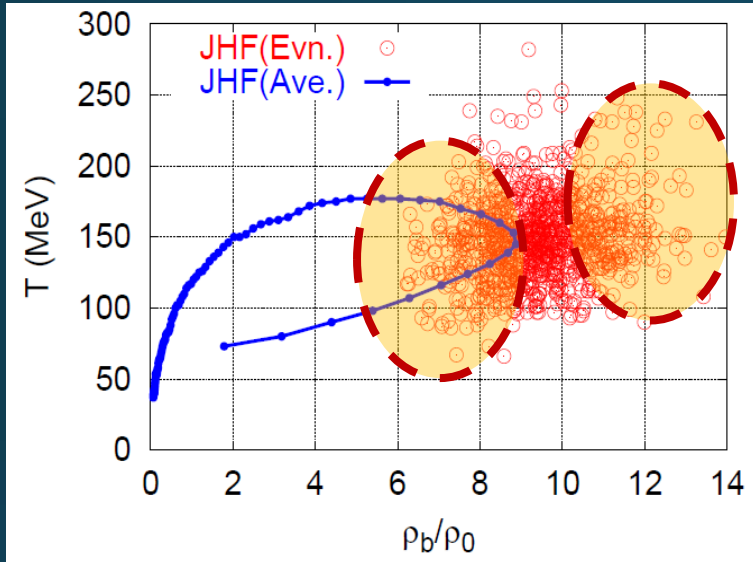


Reproduction of $\sqrt{s_{NN}}$ dependence

Nara, Ohnishi, PRC ('22)

STAR, PRL('14)

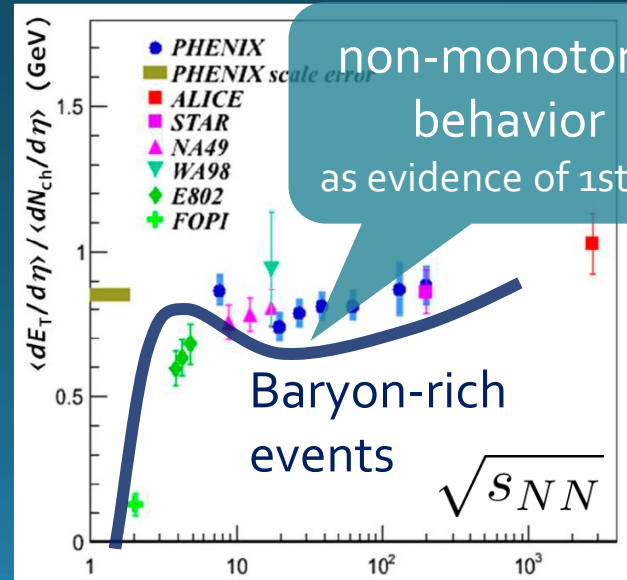
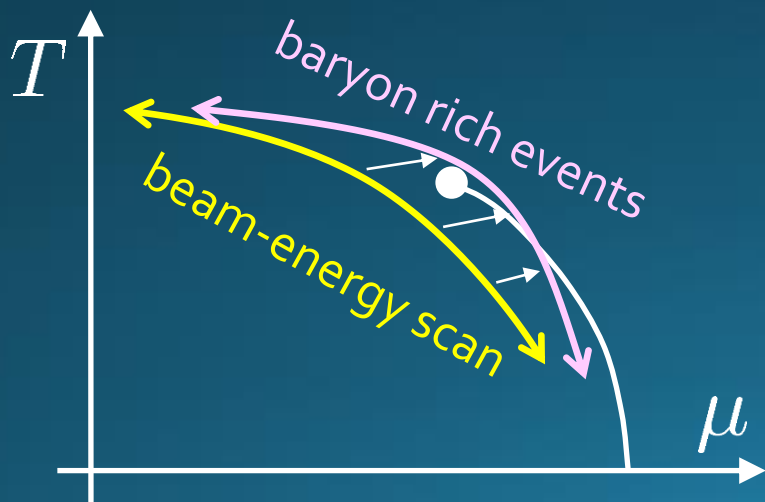
Event Selection



Large fluctuations in highest density

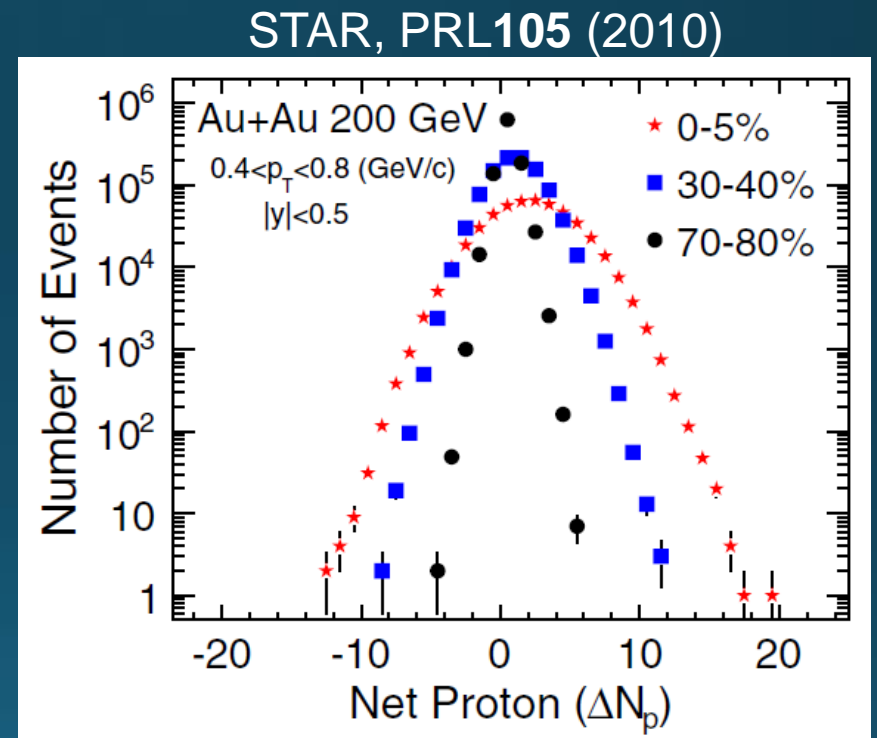
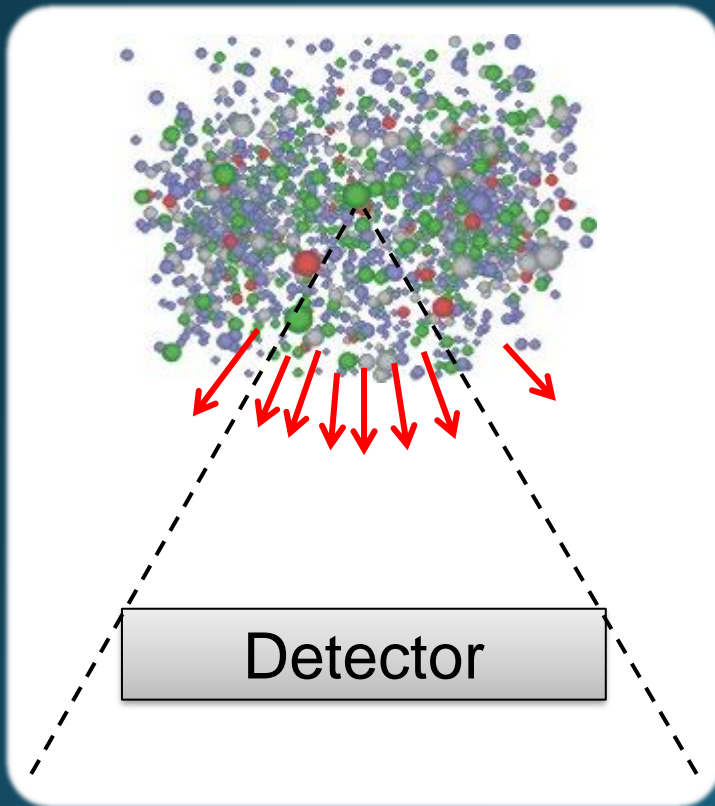
↓

“Density scan” by event selection?



Event-by-Event Fluctuations

Review: Asakawa, MK, PPNP 90 (2016)



Cumulants

$$\langle \delta N_p^2 \rangle, \langle \delta N_p^3 \rangle, \langle \delta N_p^4 \rangle_c$$

Non-Gaussian Cumulants

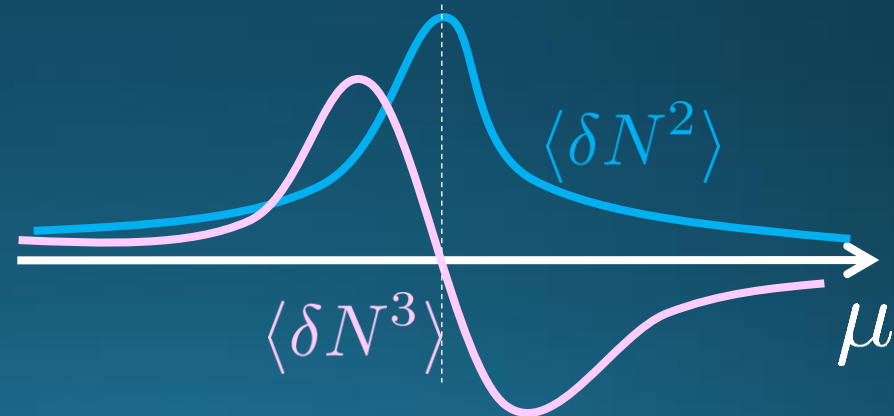
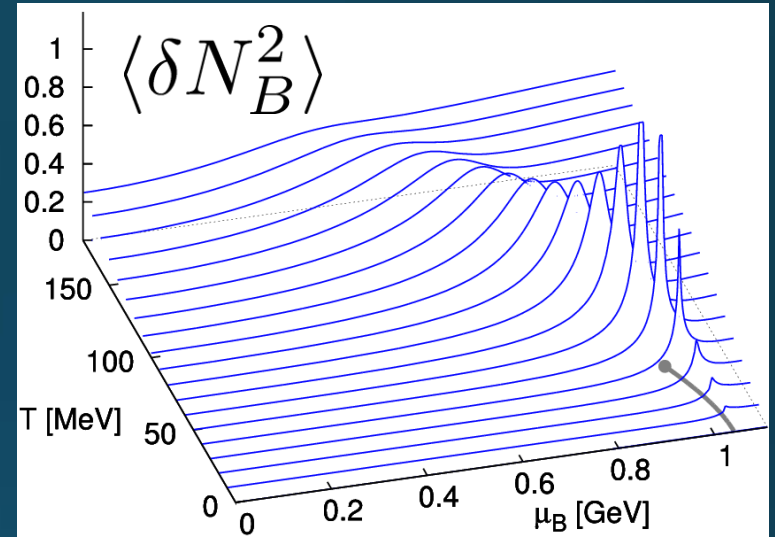
Gaussian fluctuations diverge at the QCD-CP



- Higher order cumulants change sign at the phase boundary

$$\langle \delta N^3 \rangle = T \frac{\partial \langle \delta N^2 \rangle}{\partial \mu}$$

Asakawa, Ejiri, MK, 2009



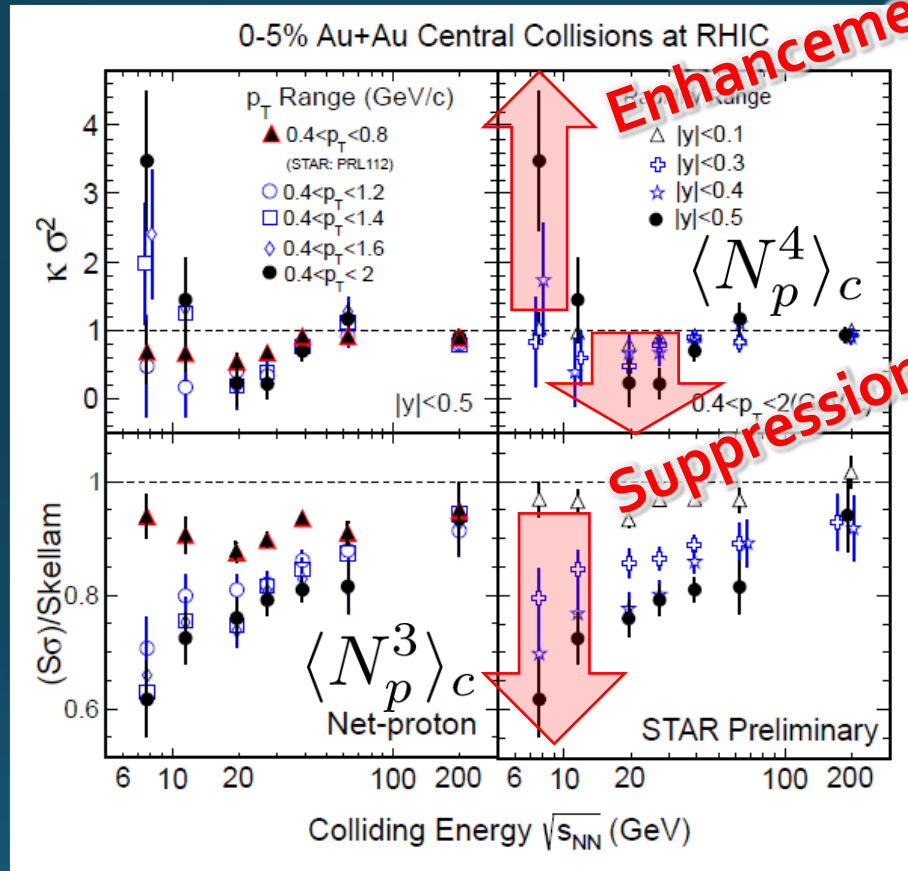
- Steeper divergence for higher-order cumulants

Stephanov, 2009

Experimental Results

$$\frac{\langle N_p^4 \rangle_c}{\langle N_p^2 \rangle_c}$$

$$\frac{\langle N_p^3 \rangle_c}{\langle N_p^2 \rangle_c}$$



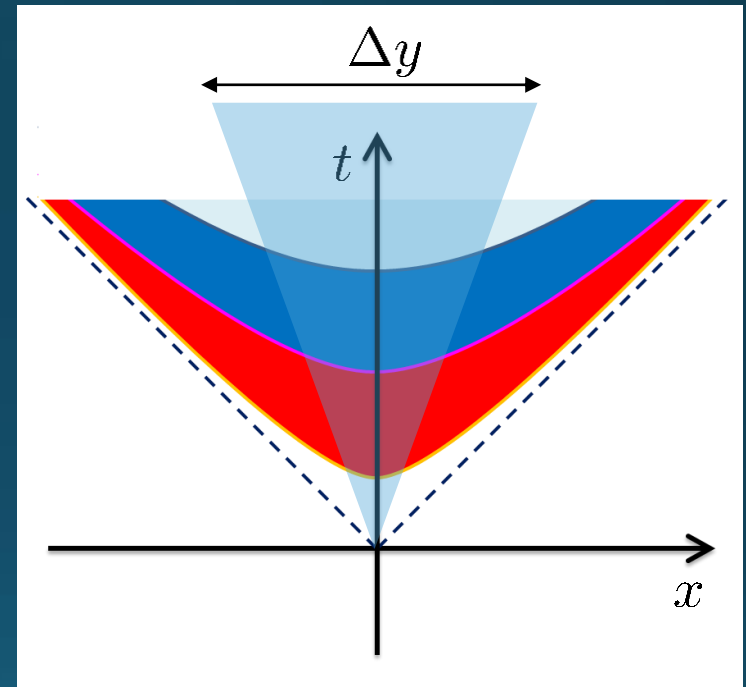
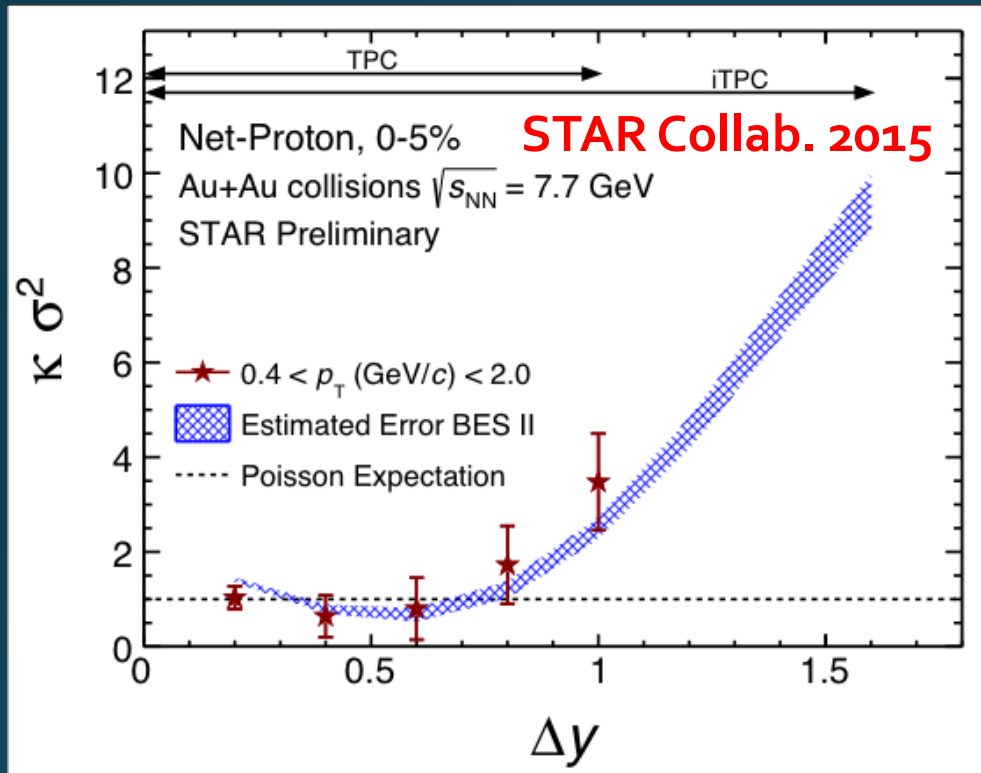
STAR Collab.
2010~

Enhancement & Suppression
of non-Gaussian cumulants!



Have we observed
QCD critical point?

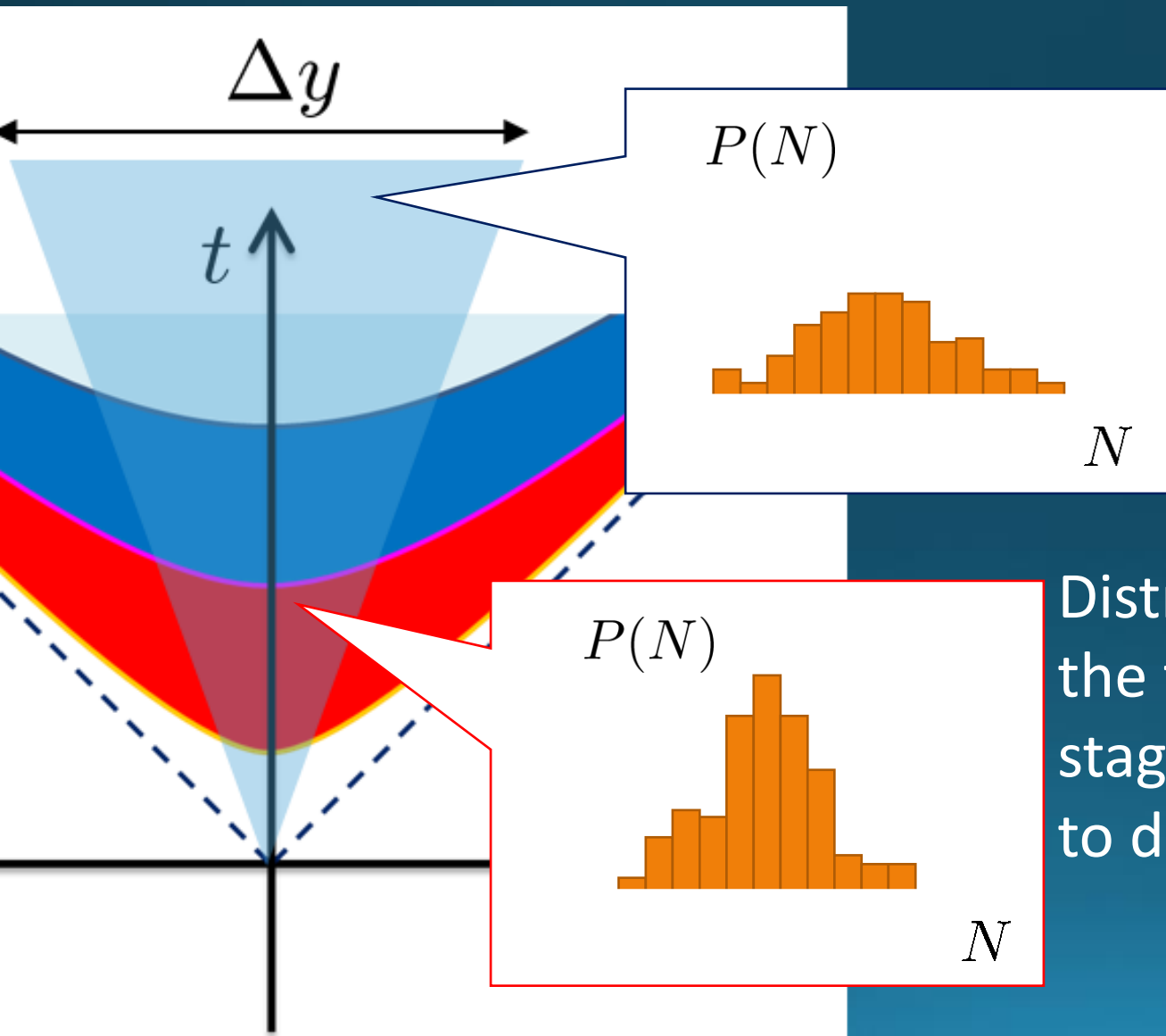
Rapidity Window Dependence



- Non-Gaussian Cumulants have been observed as a function of rapidity window Δy .
- Some results have non-monotonic Δy dependence.

Diffusion of Fluctuations

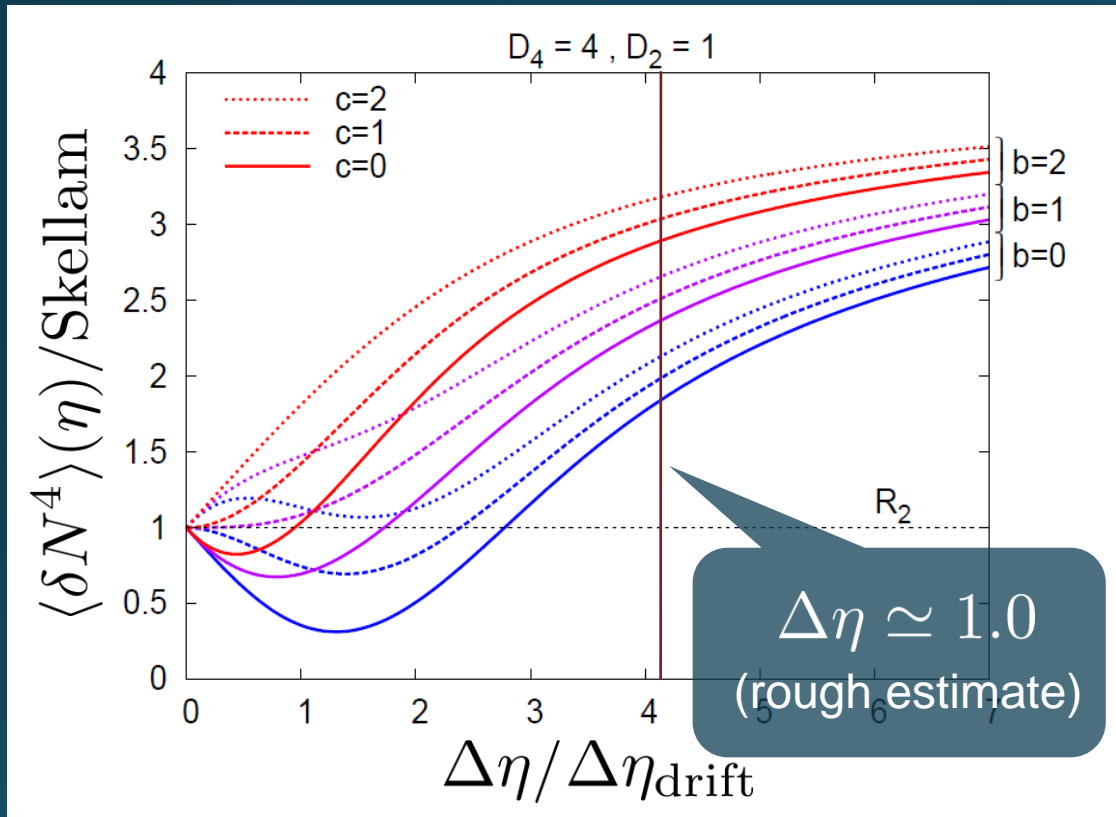
MK, Ohno, Asakawa 2014
MK 2015



Distributions in Δy in the final state and early stage are different due to diffusion.

Rapidity Window dependence as a Result of Diffusion

MK+ (2014); MK (2015)



Parameters

$$D_4 = \frac{\langle Q_{(\text{net})}^4 \rangle_c}{\langle Q_{(\text{tot})} \rangle} = 4$$

$$D_2 = \frac{\langle Q_{(\text{net})}^2 \rangle_c}{\langle Q_{(\text{tot})} \rangle} = 1$$

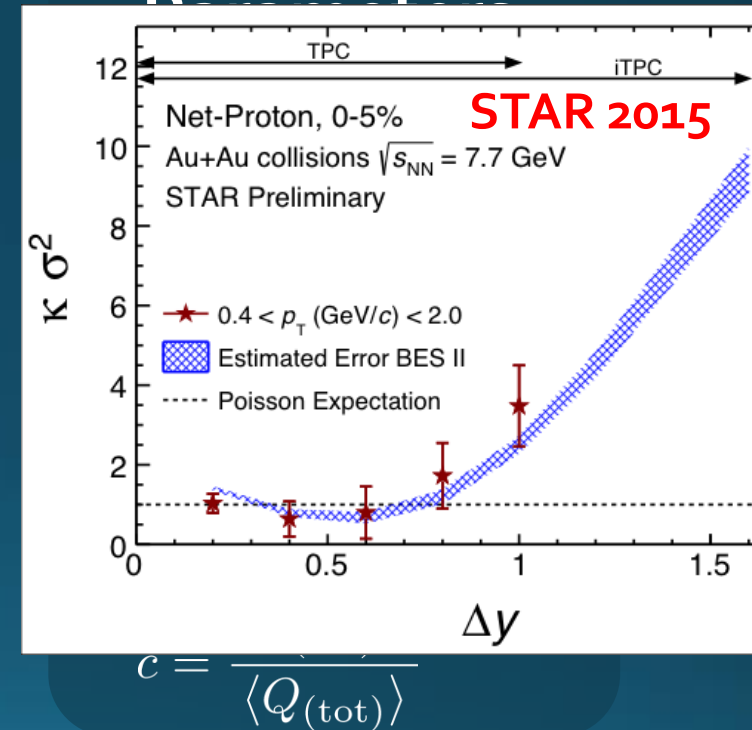
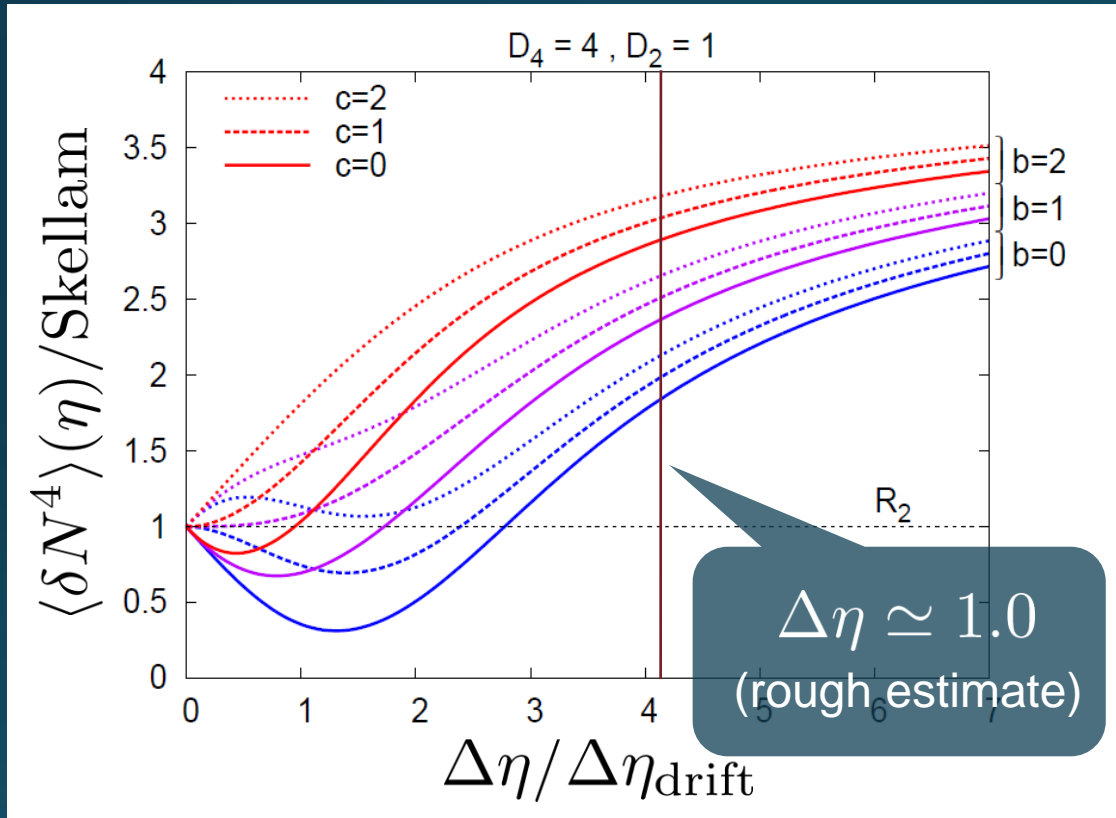
$$b = \frac{\langle Q_{(\text{net})}^2 Q_{(\text{tot})} \rangle_c}{\langle Q_{(\text{net})} \rangle}$$

$$c = \frac{\langle Q_{(\text{tot})}^2 \rangle_c}{\langle Q_{(\text{tot})} \rangle}$$

- Higher order cumulants can behave non-monotonically.
- $\Delta\eta$ dependence encodes history of time evolution.

Rapidity Window dependence as a Result of Diffusion

MK+ (2014); MK (2015)

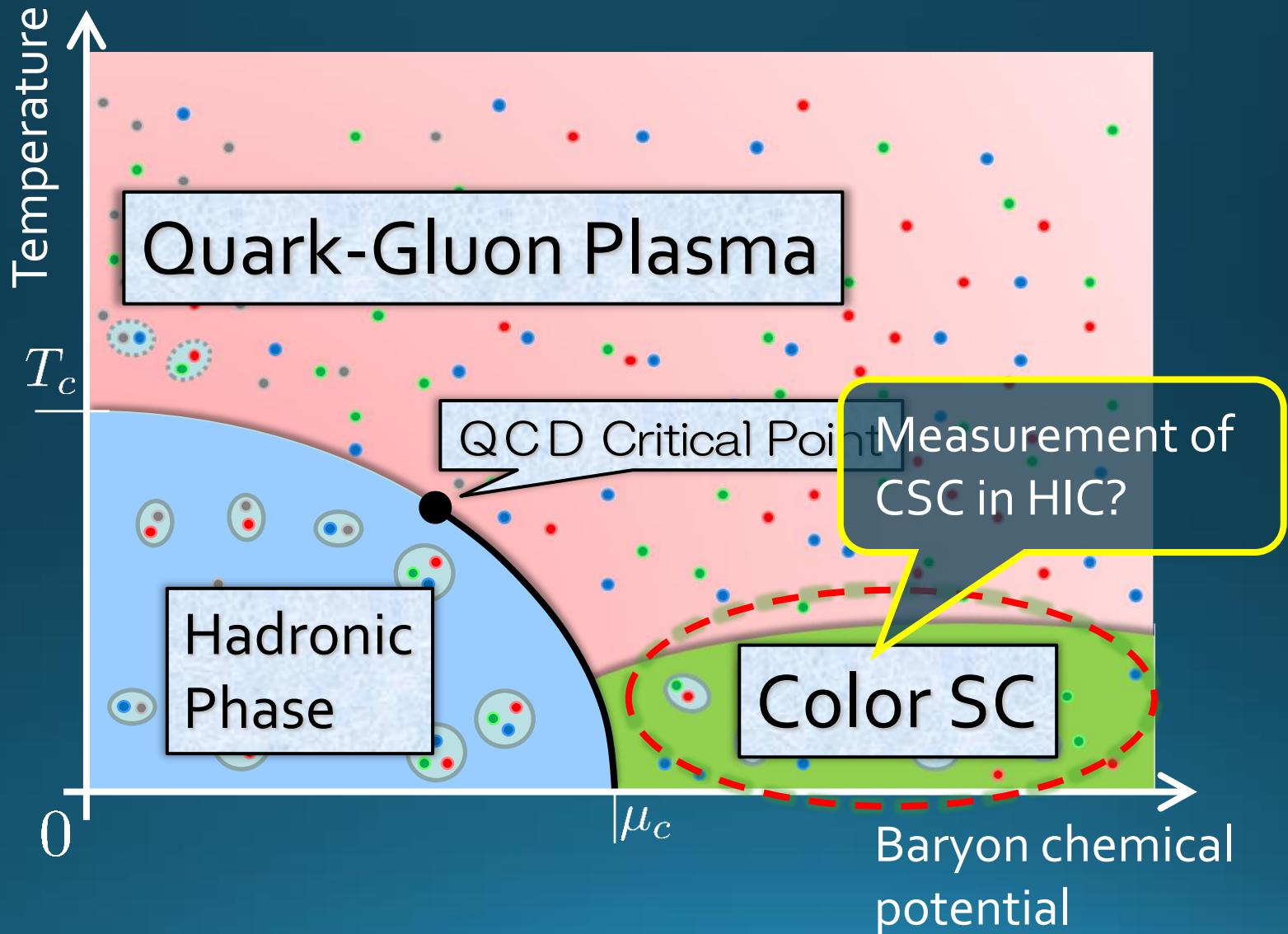


- Higher order cumulants can behave non-monotonically.
- $\Delta\eta$ dependence encodes history of time evolution.

Search for **CSC Phase Transition** using dilepton production rates

Nishimura, MK, Kunihiro, arXiv:2201.01963

Color Superconductivity



Observing CSC in HIC?

Difficulties

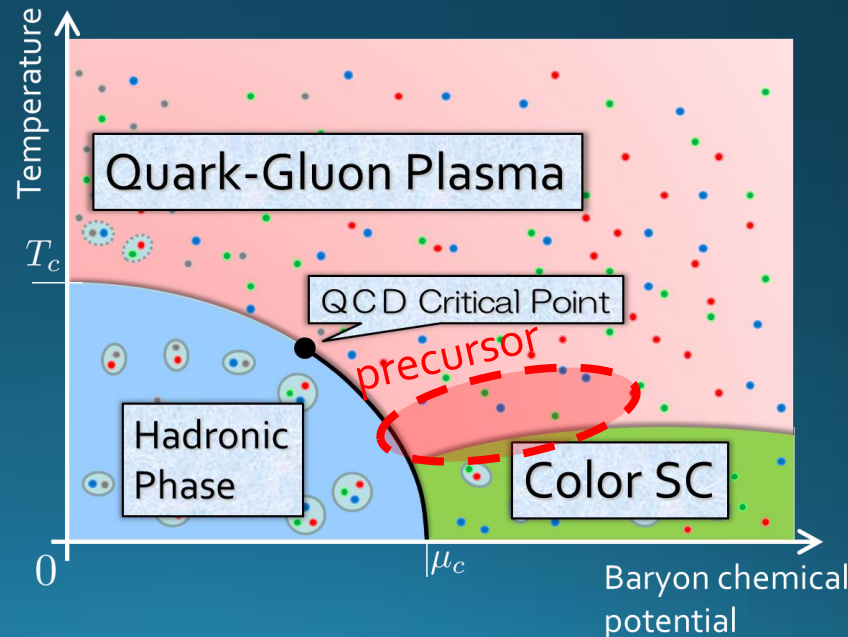
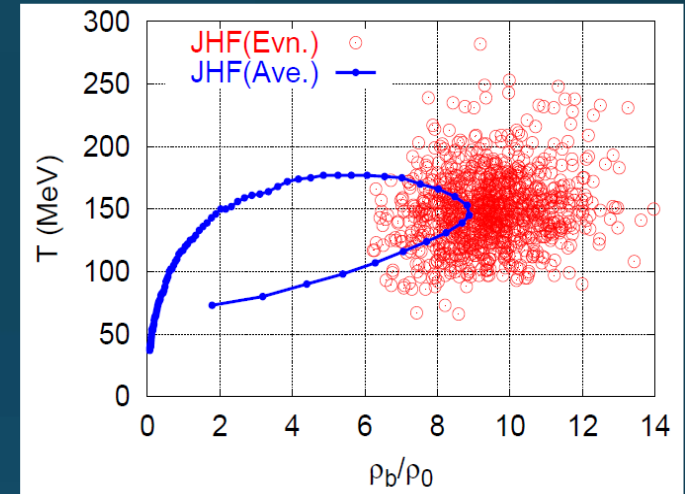
- 1) Creation of CSC itself would be impossible.
- 2) CSC would be realized only in the early stage.



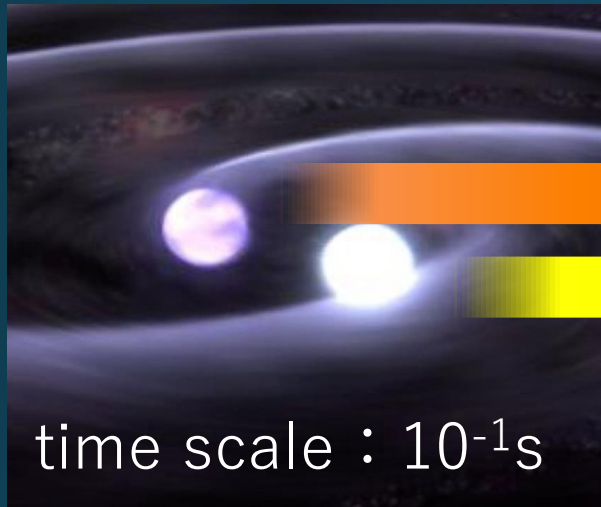
Solution

- 1) Focus on **precursor of CSC**
- 2) Use **dilepton production**

Nishimura, MK, Kunihiro, 2201.01963

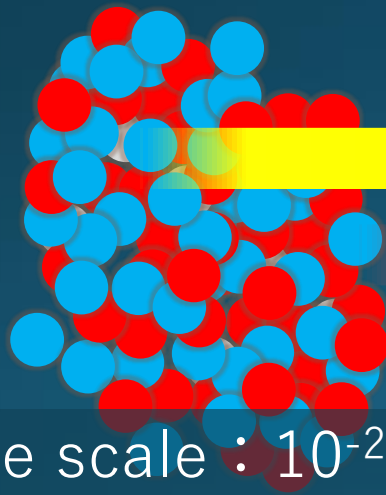


Multi-Messenger Observation



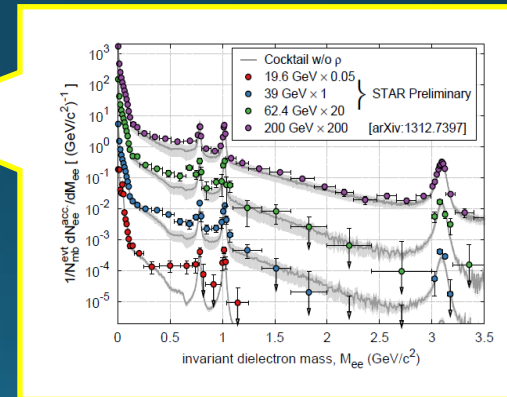
gravitational waves

electromagnetic waves



leptons,
photons

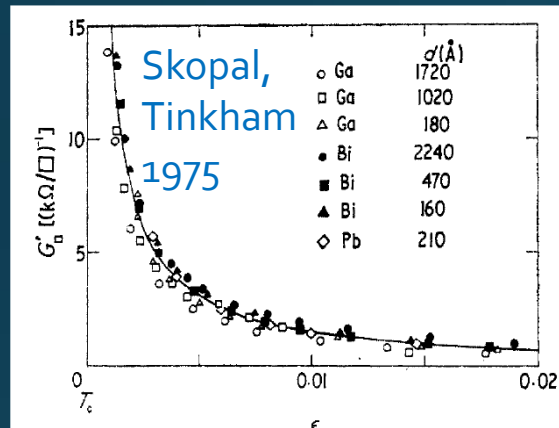
hadronic observables



Precursory Phenomena

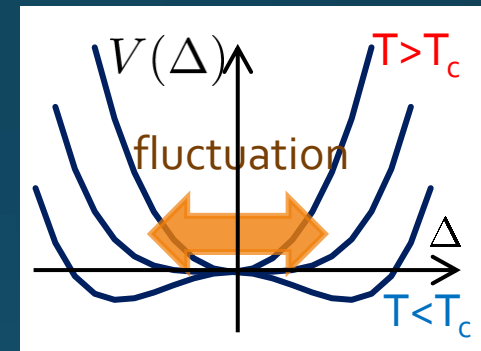
= anomalous behavior of observables near but above T_c

Electric conductivity in metals



$$\epsilon = \frac{T - T_c}{T_c}$$

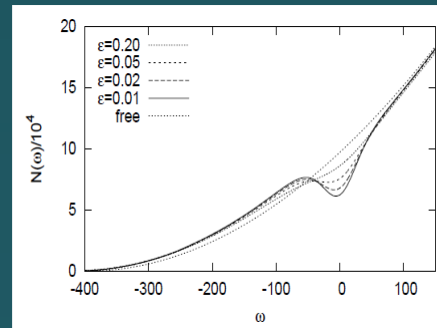
Landau's Free Energy



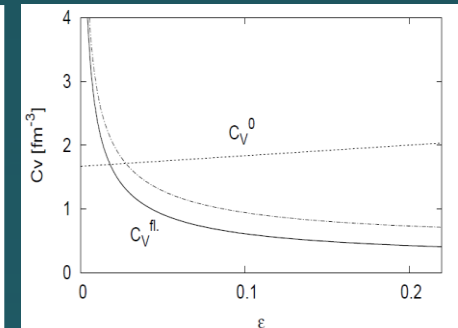
Precursor of CSC

- Pseudogap = Depression in the density of states
- Specific heat
- etc.

Density of States



Specific Heat



MK, Koide, Kunihiro, Nemoto, '03, '05

Model

NJL model (massless 2-flavor)

$$\mathcal{L} = \bar{\psi} i \not{\partial} \psi + \mathcal{L}_S + \mathcal{L}_C$$

$$\mathcal{L}_S = G_S ((\bar{\psi}\psi)^2 + (\bar{\psi} i \gamma_5 \tau \psi)^2)$$

$$\mathcal{L}_C = G_C ((\bar{\psi} i \gamma_5 \tau_A \lambda_A \psi^C)(\text{h.c.}))$$

diquark interaction

$$G_S = 5.01 \text{ GeV}^{-2}, \quad \Lambda = 650 \text{ MeV}$$

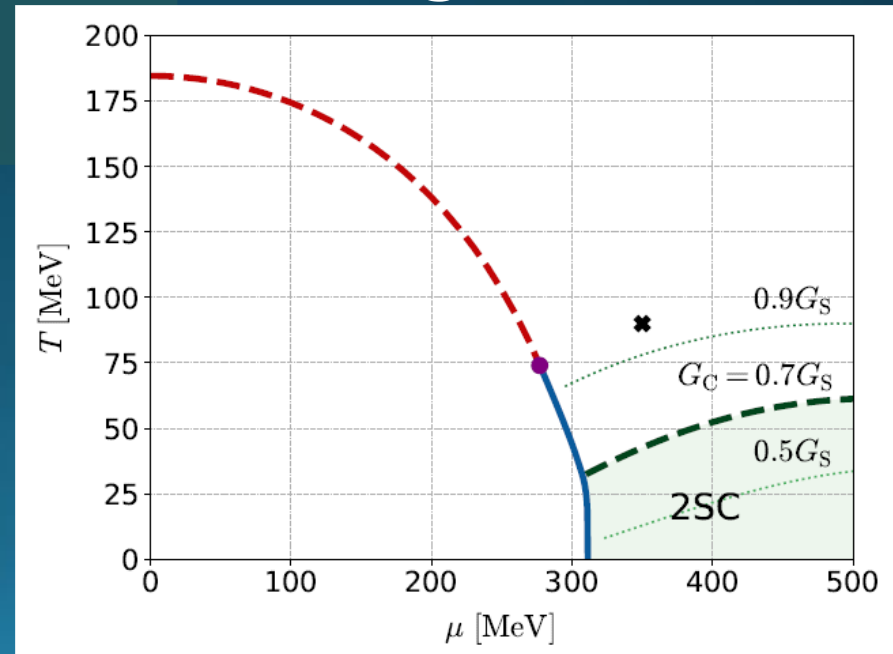
□ G_C : free parameter

Order of CSC phase transition

Matsuura+('04), Giannakis+('04)

Noronha+('06), Fejos, Yamamoto('19)


Phase Diagram in MFA



Diquark Mode

Diquark Propagator (T-matrix)

in random-phase approximation

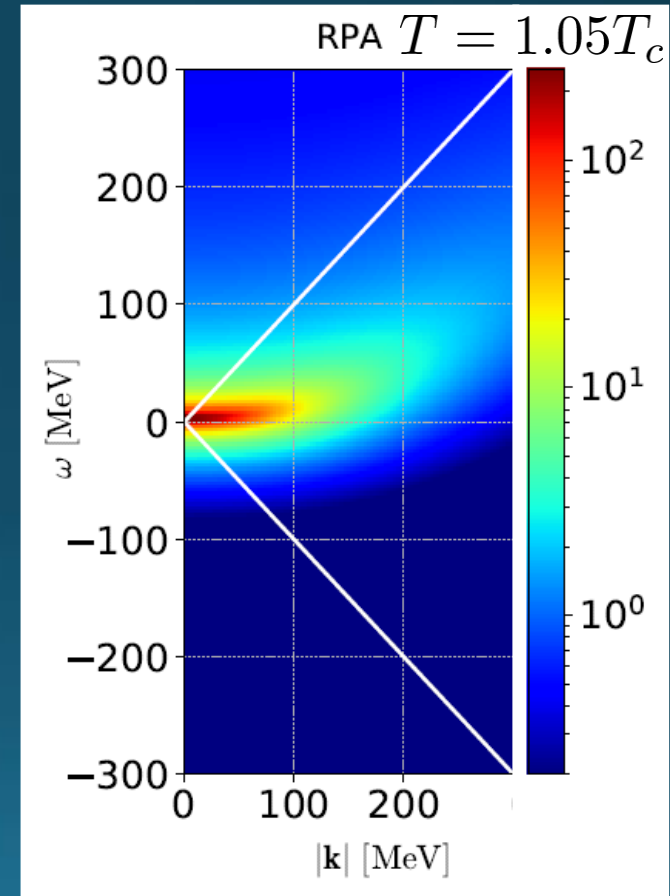
$$\Xi^R(\mathbf{k}, \omega) =$$


The diagram shows a series of bubble diagrams representing the Dyson equation for the diquark propagator. It starts with a single bubble, followed by a plus sign, then two bubbles connected in series, followed by a plus sign and an ellipsis. Below this is the equation:

$$= \frac{Q^R(\mathbf{k}, \omega)}{1 + G_C Q^R(\mathbf{k}, \omega)}$$

- Massless at $T=T_c$ as a soft mode of CSC transition
- Strength in the space-like region

Dynamical Structure Factor

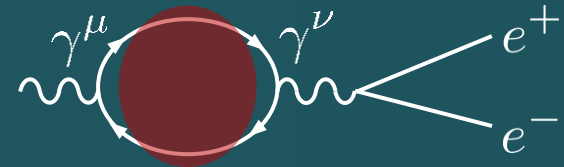


Photon Self-Energy

Dilepton Production Rate

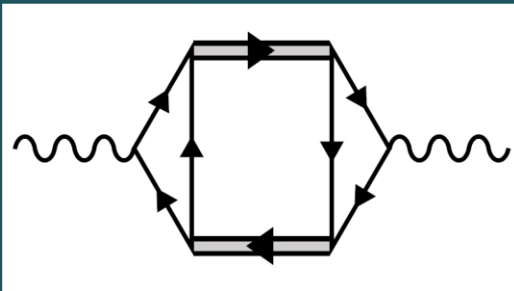
$$\frac{d^4\Gamma}{dk^4} = \frac{\alpha}{12\pi^4} \frac{1}{k^2} \frac{1}{e^{\beta\omega}-1} \text{Im}\Pi_{\mu}^{R\mu}(k)$$

$$\Pi^{\mu\nu}(k) =$$

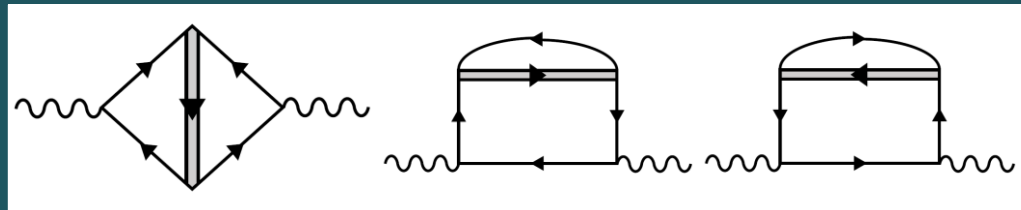


Terms included in $\Pi^{\mu\nu}$

Aslamasov-Larkin



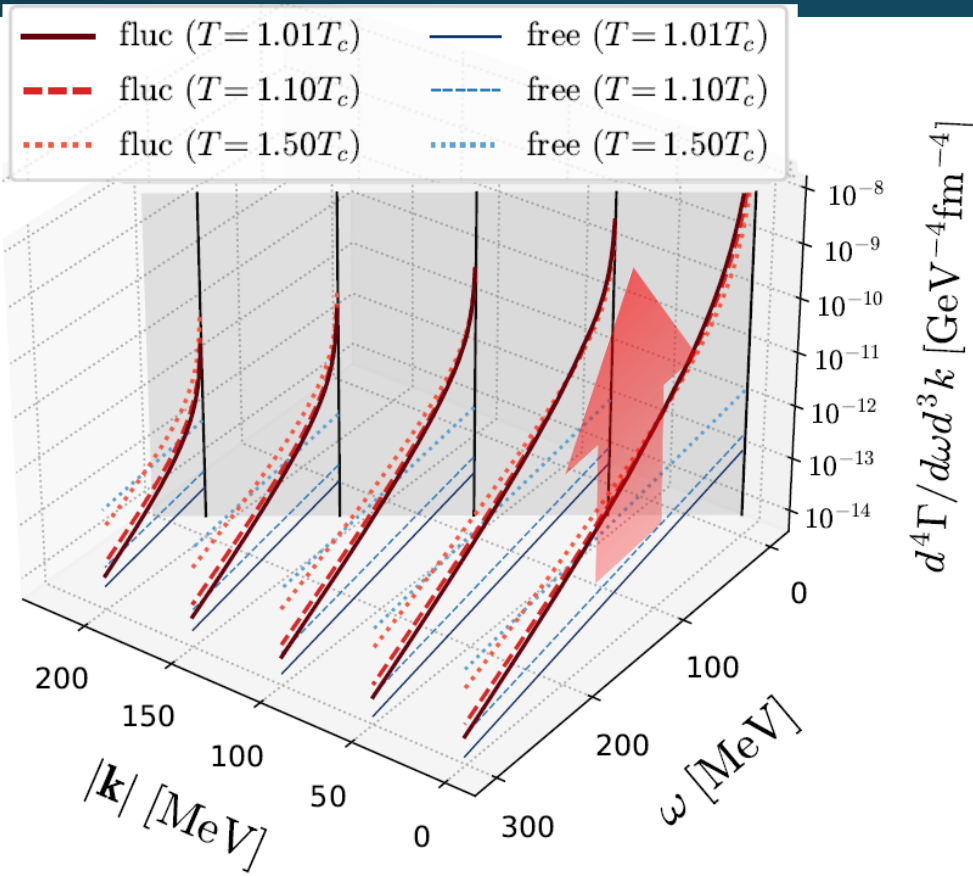
Maki-Thompson, Density of states



\Rightarrow = diquark propagator

- common in metallic superconductors (conductivity)
- time-dependent GL approx. for diquark field
- gauge-invariant construction

Dilepton Production Rates



$$\frac{d^4\Gamma}{dk^4}(\omega, k)$$

Red: fluctuation contribution
Blue: free quarks

$$\mu = 350 \text{ MeV}$$

$$G_C = 0.7G_S, \quad (T_c \simeq 43 \text{ MeV})$$

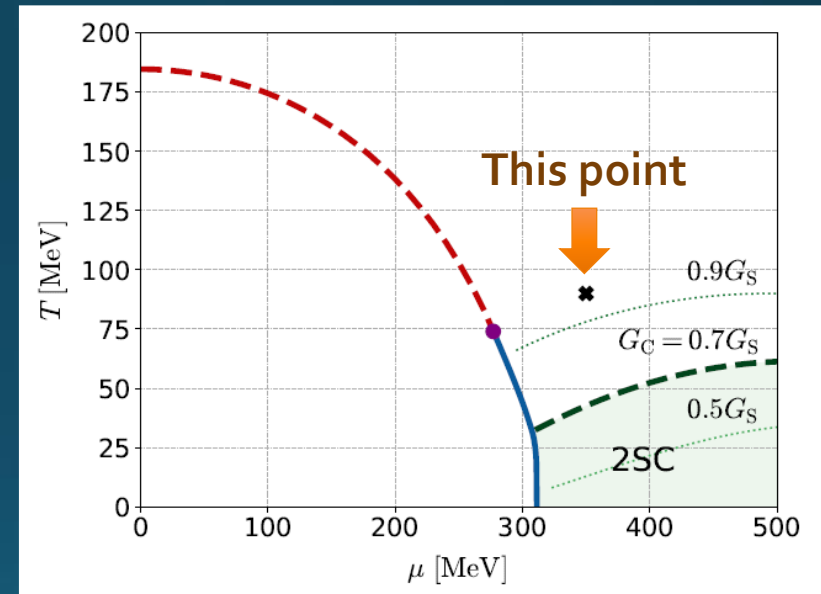
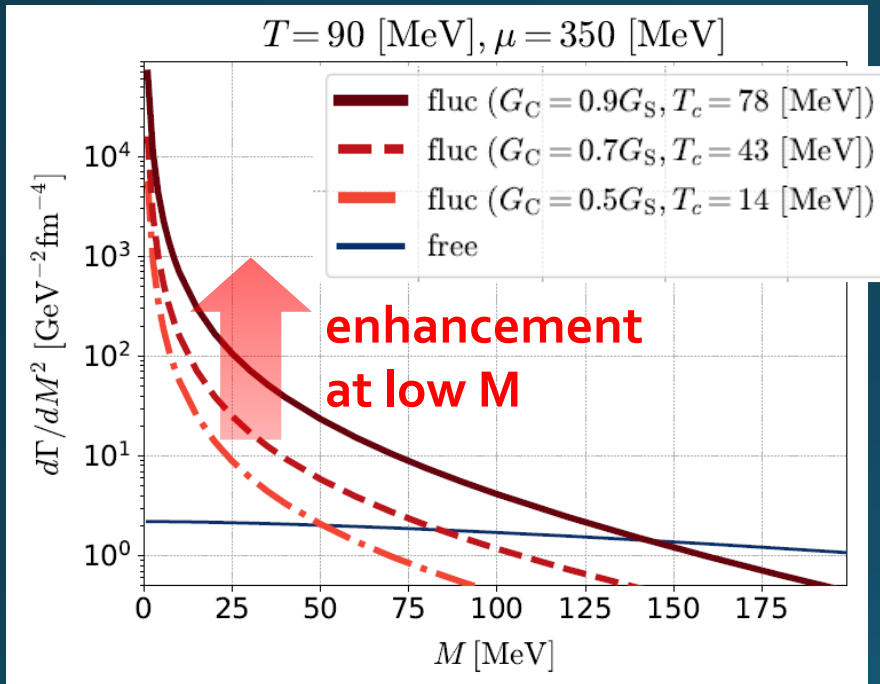
Nishimura, MK, Kunihiro,
 2201.01963

- Diquark fluctuations give rise to anomalous enhancement in the low energy-momentum region for $T < 1.5T_c$.

Invariant-Mass Spectrum

Fixed Temperature

$$\frac{d\Gamma}{dM^2} = \int \frac{d^3k}{2\omega} \frac{d^4\Gamma}{dk^4} \Big|_{\omega=E_k}$$

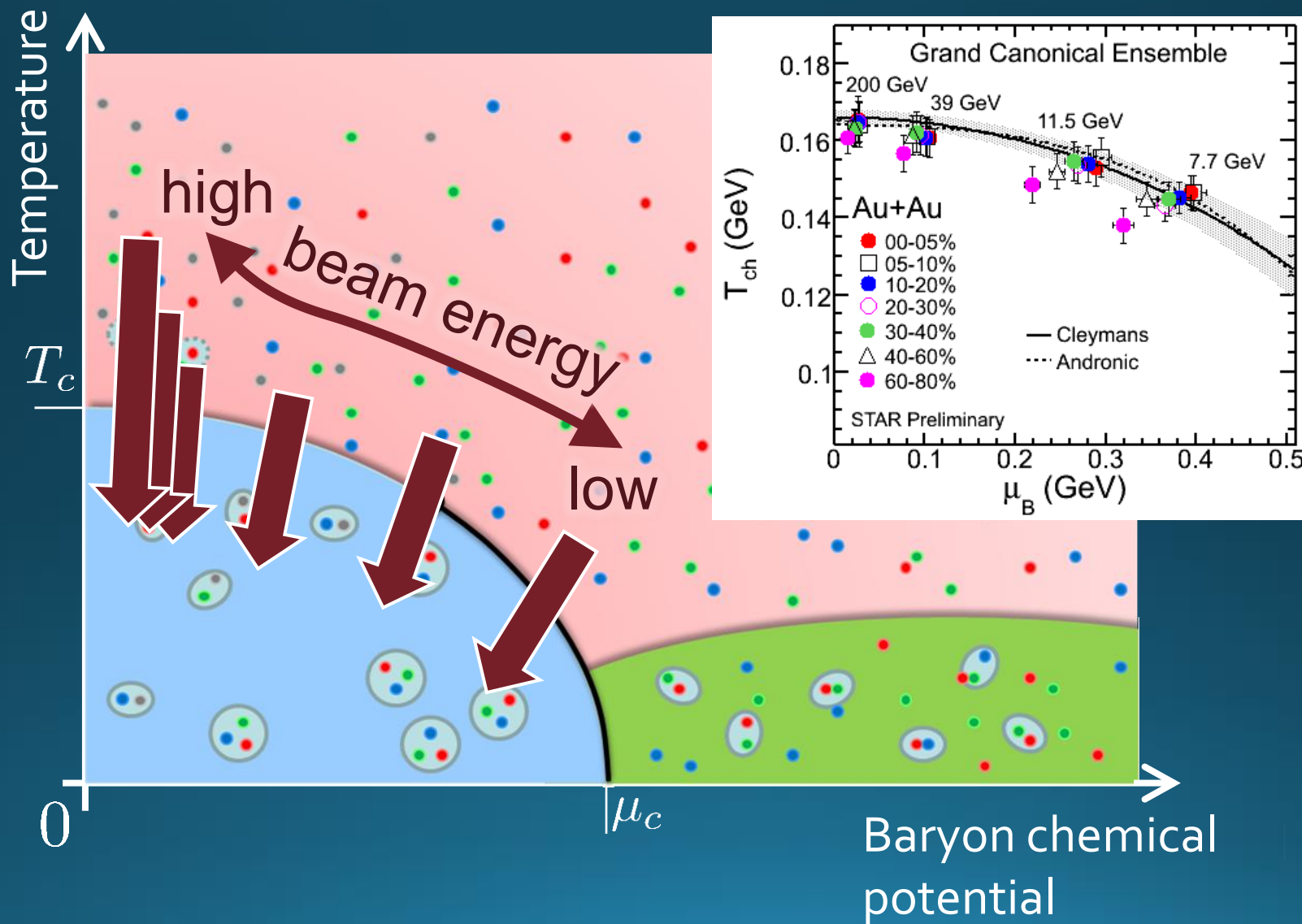


- Strong enhancement at low invariant mass, though the range of M is narrower than the previous results.
- Observable in the HIC?

Summary

- ❑ Exploring dense medium in relativistic heavy-ion collisions is a hot topic in this field. The beam-energy scan is ongoing, and many new experiments will start in near future!
- ❑ Among them, search for
 - ❑ the **QCD critical point** using fluctuation observables
 - ❑ the **color superconductivity** using dileptonsare especially interesting and important subjects.
- ❑ Heavy-ion collisions at J-PARC will play important roles in pursuing these subjects.

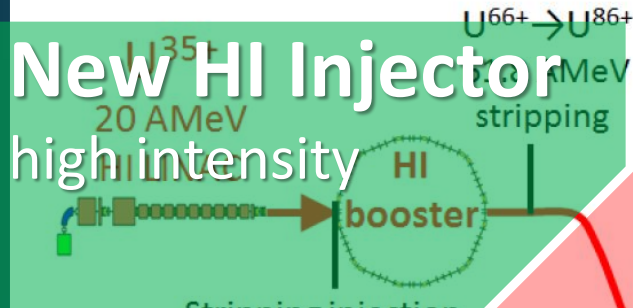
Beam-Energy Scan



HI Acceleration @ J-PARC

New HI Injector

high intensity



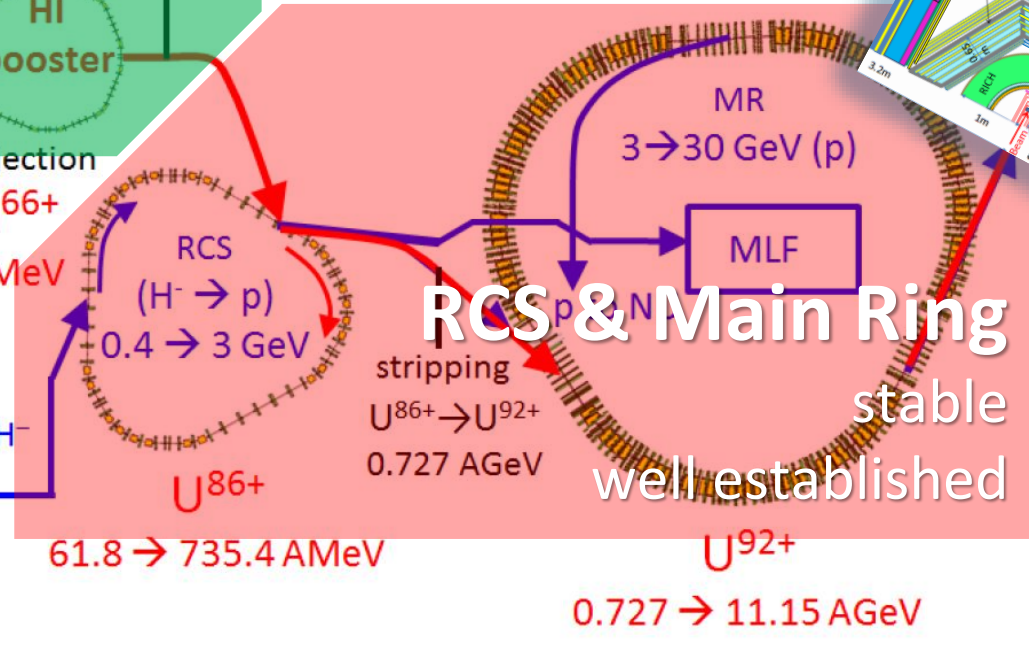
Stripping injection

U³⁵⁺ → U⁶⁶⁺
 20 → 67 AMeV

H⁻ Linac: 0.4 GeV

U⁸⁶⁺
 61.8 → 735.4 AMeV

RCS & Main Ring



J-PARC Heavy Ion Spectrometer

Figures: No

— proton (exist)

— HI (under)



- Use of reliable / high-performance RCS & main ring
- → Reduce cost and time