

Lee-Yang Zeros and Critical Points

Masakiyo Kitazawa

In collab. with **Tatsuya Wada**, Kazuyuki Kanaya

Wada, MK, Kanaya, PRL ('25), arXiv:2508.20422

Lee-Yang Zeros and

PHYSICAL REVIEW LETTERS **134**, 162302 (2025)

Editors' Suggestion

Locating Critical Points Using Ratios of Lee-Yang

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In collab. with **Tatsuya**

Wada, MK, Kanaya, PR

京都大学



KYOTO UNIVERSITY



ホーム > 最新の研究成果を知る > 多様な系に現れる臨界点の性質を系統的に決定する新手法を開発—リーヤンゼロを用いた一般的手法—

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リーヤンゼロを用いた一般的手法—

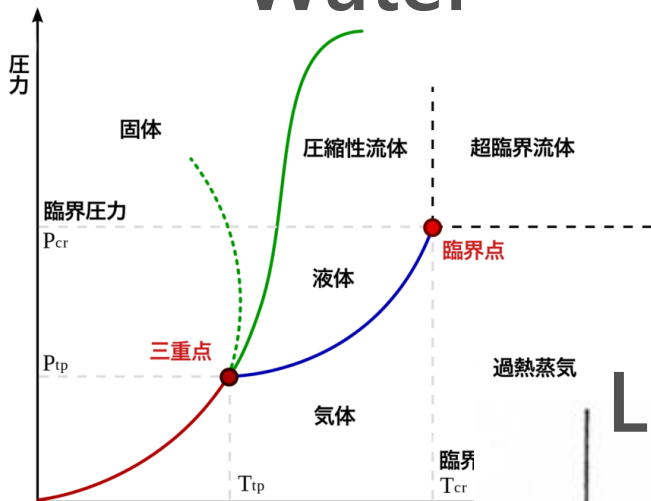
企業・研究者の方

公開日: 2025年05月02日

臨界点は、宇宙から素粒子まで様々なスケールの相転移現象に現れる普遍的な概念です。臨界点近傍で

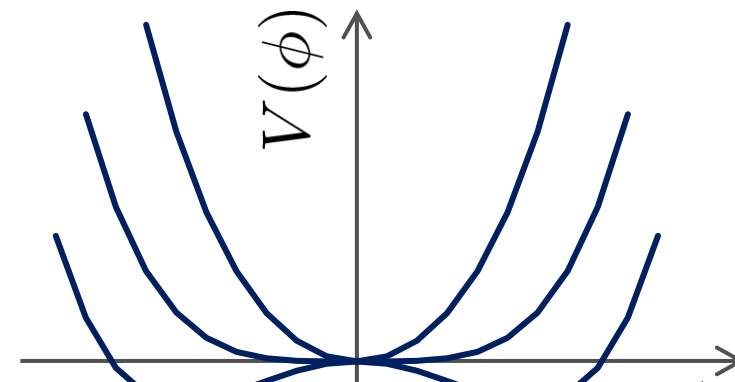
Critical Points Everywhere!

Water

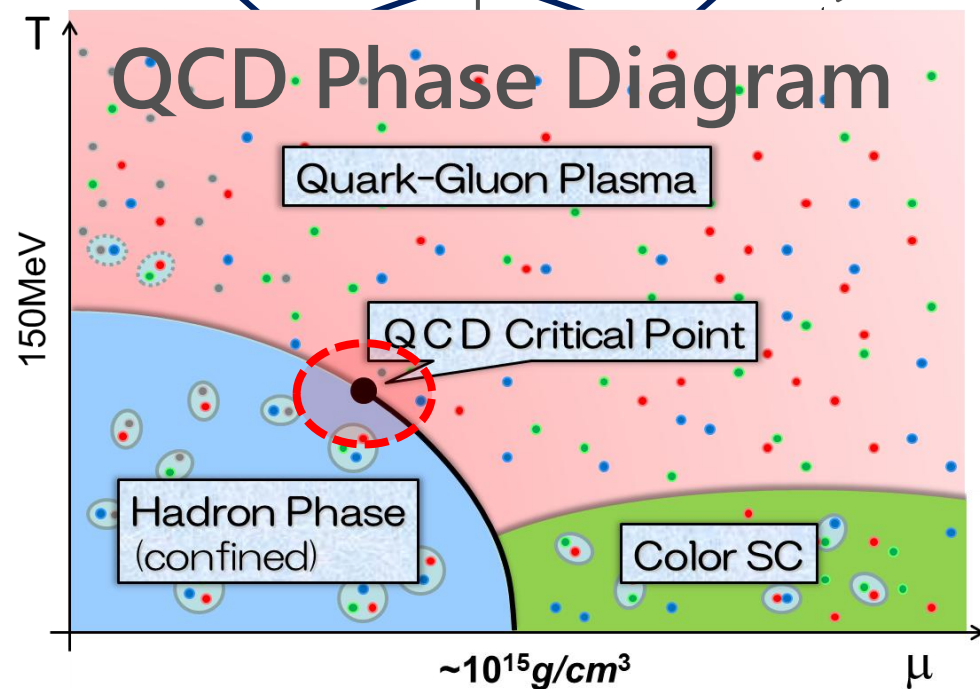
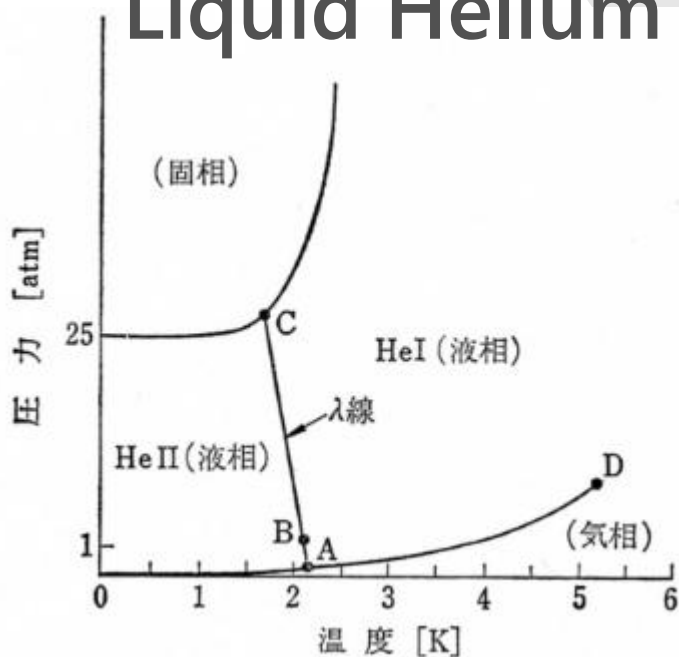


Ising

$$V(\phi) = -a\phi^2 + b\phi^4$$

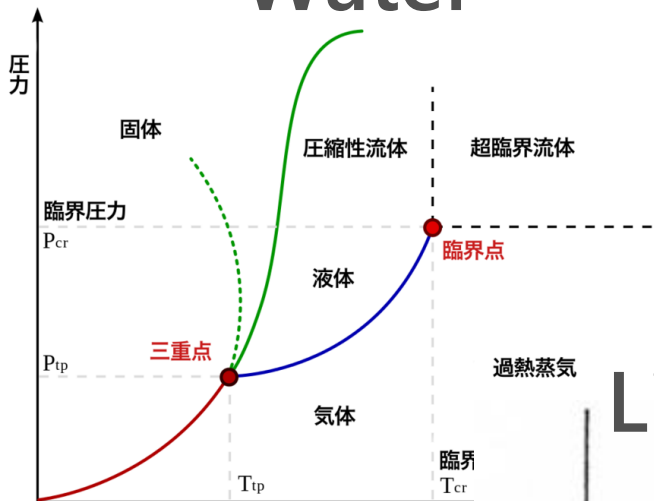


Liquid Helium



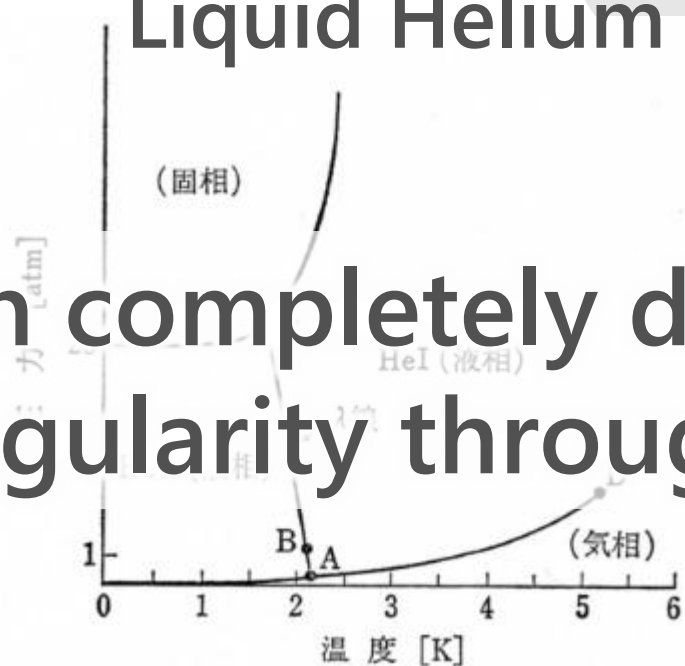
Critical Points Everywhere!

Water

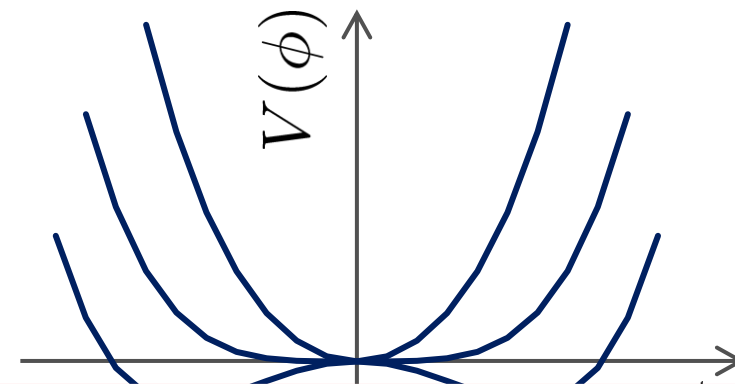


Ising

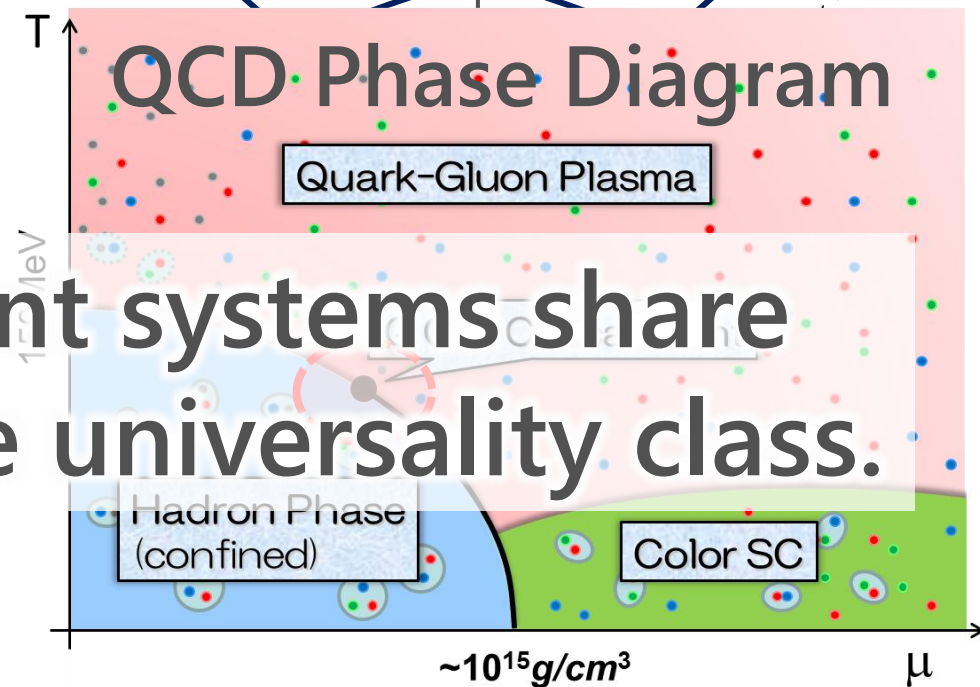
Liquid Helium



$$V(\phi) = -a\phi^2 + b\phi^4$$



QCD Phase Diagram



CPs in completely different systems share the singularity through the universality class.

Search for QCD-CP with Lee-Yang Zeros

in Lattice QCD Simulations

...;

Dimopoulos+, PRD ('22);

Clarke+, 2405.10196;

Adam+, 2507.13254; ...

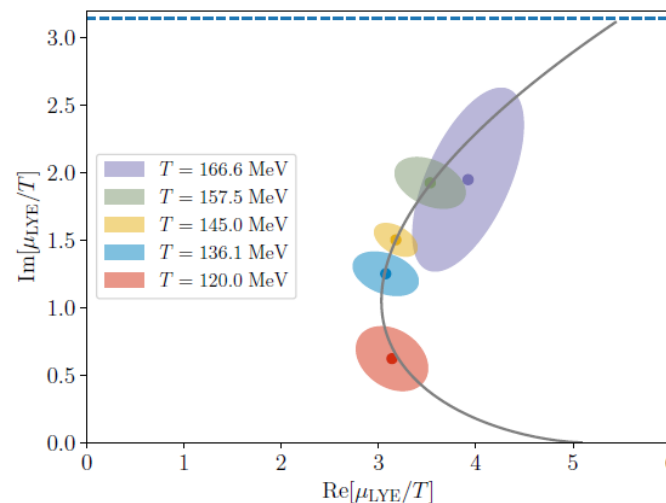
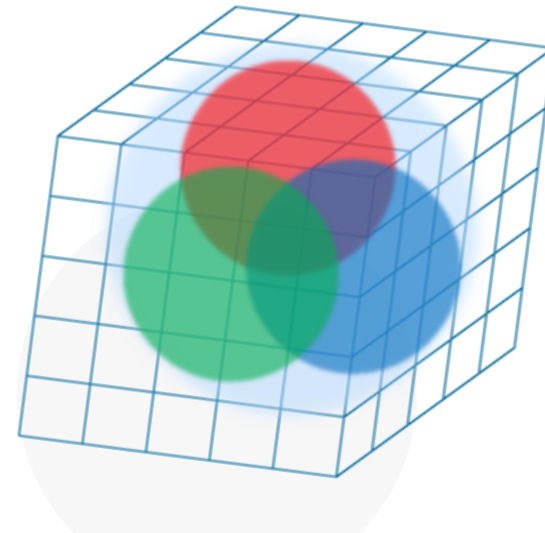
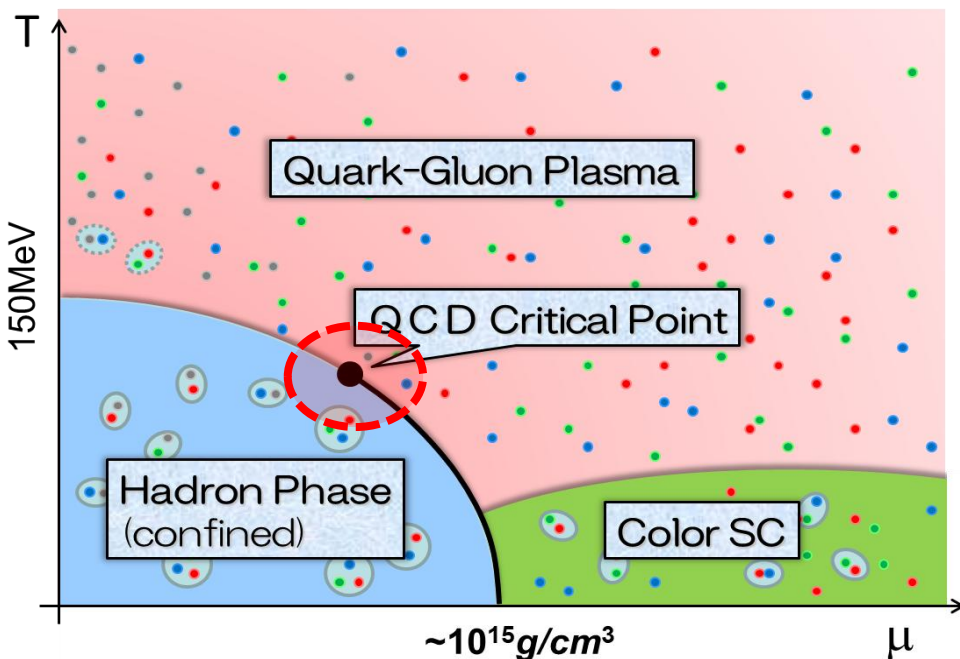
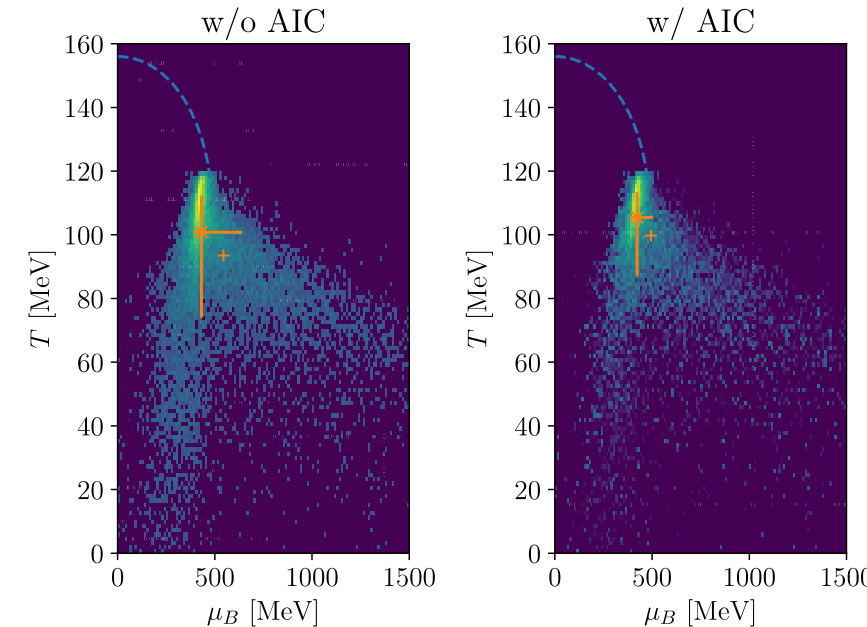


FIG. 3. Singularities at $T = 166.6, 157.5, 145.0, 136.1$ and 120.0 MeV. The dashed line lies at $\hat{\mu}_B = i\pi$.

arXiv:2405.10196v1 [hep-lat] 16 May 2024



$$\begin{cases} \mu^{\text{CEP}} = 422^{+80}_{-35} \text{ MeV} \\ T^{\text{CEP}} = 105^{+8}_{-18} \text{ MeV} \end{cases}$$

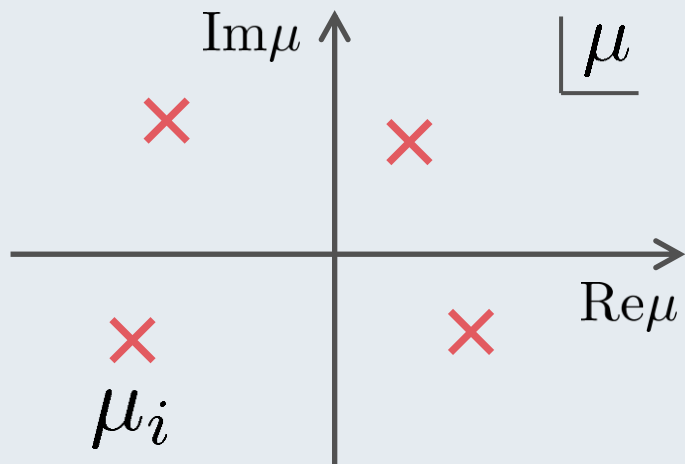
Lee-Yang Zeros

Yang, Lee; Lee, Yang ('52)

Partition Function $Z(T, \mu)$

Finite V \rightarrow Polynomial of μ (or T)

$$Z(T, \mu) = \prod_i (\mu - \mu_i)$$



\rightarrow zeros on the complex plane
= Lee-Yang Zeros

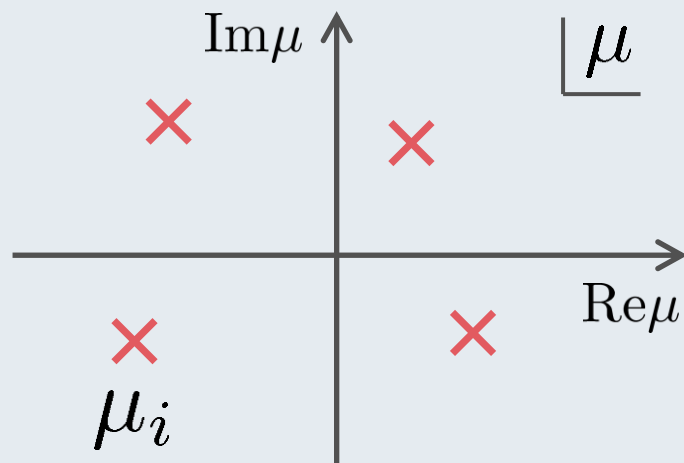
Lee-Yang Zeros

Yang, Lee; Lee, Yang ('52)

Partition Function $Z(T, \mu)$

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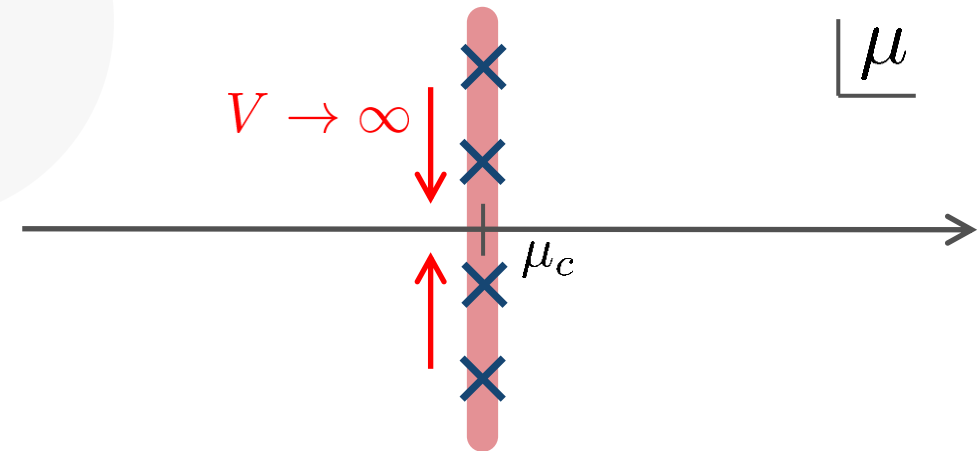
$$Z(T, \mu) = \prod_i (\mu - \mu_i)$$



\rightarrow zeros on the complex plane
= Lee-Yang Zeros

Phase Transition & LYZ

First-order transition
at $\mu = \mu_c$



— For $V \rightarrow \infty$, LYZs gather on the line
crossing the real axis at $\mu = \mu_c$.

LYZ around a Critical Point in Ising Model

t

$$t = \frac{T - T_c}{T_c}$$

1st-transition

singularity on the real h axis

Crossover

no singularity on the real axis

Note:

LYZ in complex- h plane are purely imaginary.

Lee-Yang, 1952

h

LYZ around a Critical Point in Ising Model

t

$$t = \frac{T - T_c}{T_c}$$

1st-transition

singularity on the real h axis

Crossover

no singularity on the real axis

Note:

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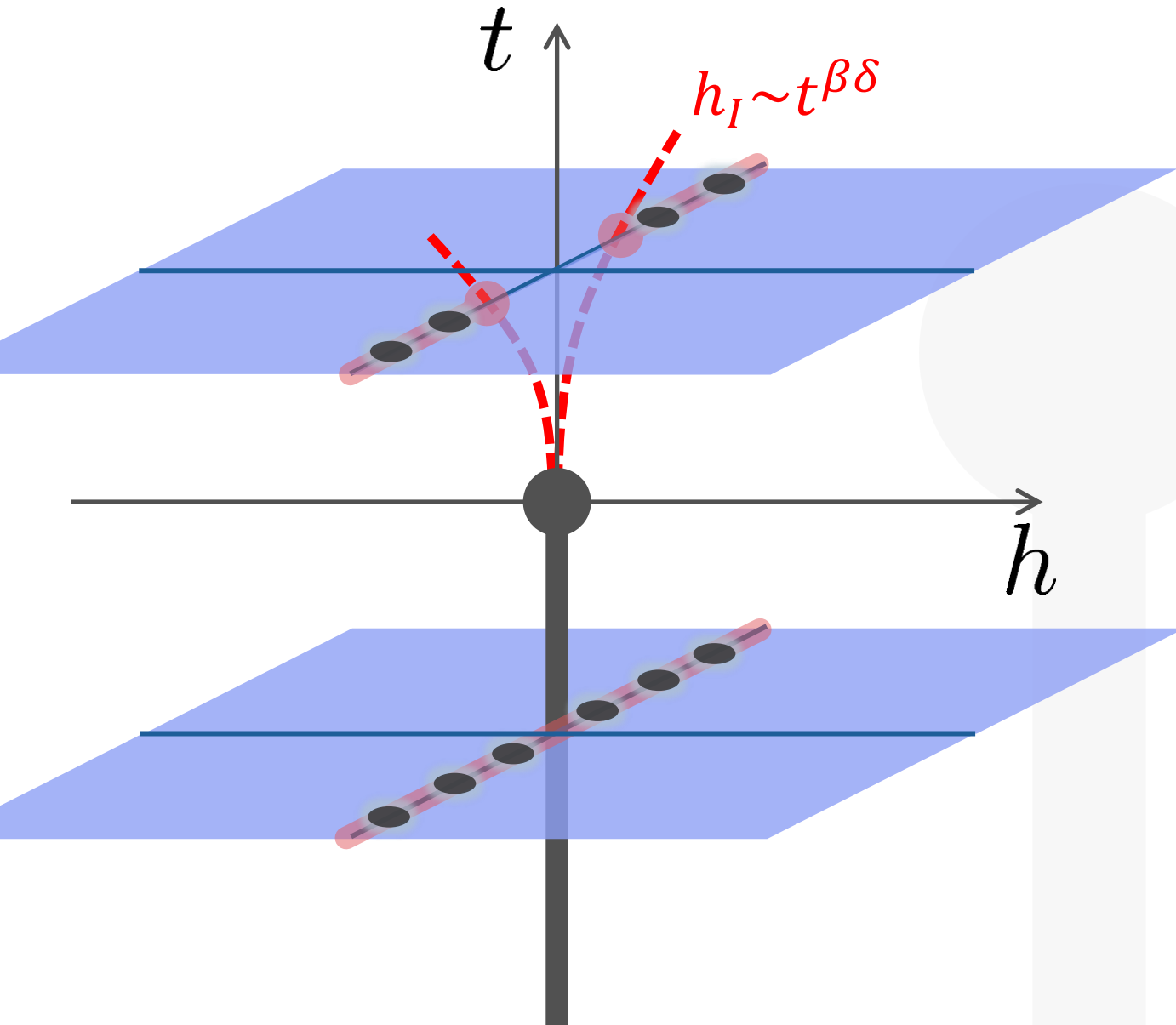
Lee-Yang, 1952

h



The diagram shows a complex plane with a vertical axis labeled t and a horizontal axis labeled h . A black dot marks the origin. A blue shaded region represents a branch cut along the real axis for $h < 0$. A red dashed line with black dots represents the Lee-Yang zeros, which are purely imaginary and located on the imaginary axis. A vertical black line passes through the origin, and a horizontal blue line is drawn below the real axis.

LYZ around a Critical Point in Ising Model



1st-transition

singularity on the real h axis

Crossover

no singularity on the real axis

LY edge singularity

Starting from the CP

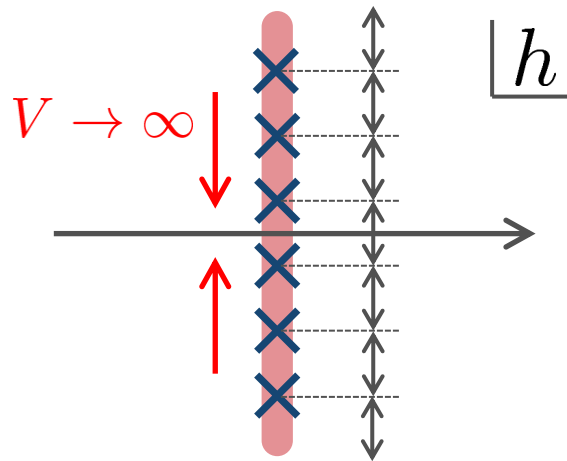
Its behavior is governed by the the scaling function.

$$h_I \sim t^{\beta\delta}$$

Lee-Yang Zero Ratios

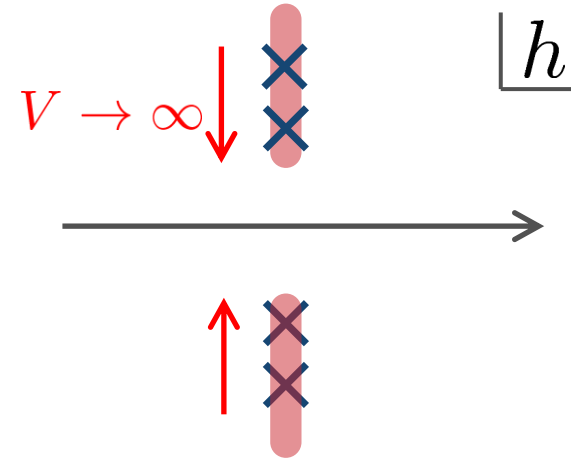
$$R_{nm}(t) = \frac{h^{(n)}(t)}{h^{(m)}(t)}$$

First-Order Side ($t < 0$)



$$R_{nm}(t) \xrightarrow{V \rightarrow \infty} \frac{2n - 1}{2m - 1}$$

Crossover Side ($t > 0$)



$$R_{nm}(t) \xrightarrow{V \rightarrow \infty} 1$$

Finite-Size Scaling

Scaling Hypothesis

$$F_{\text{sing}}(t, h, L^{-1}) = \tilde{F}_{\text{sing}}(L^{y_t} t, L^{y_h} h)$$

$$Z_{\text{sing}}(t, h, L^{-1}) = \tilde{Z}_{\text{sing}}(L^{y_t} t, L^{y_h} h)$$

$$F = F_{\text{sing}} + F_{\text{reg}}$$

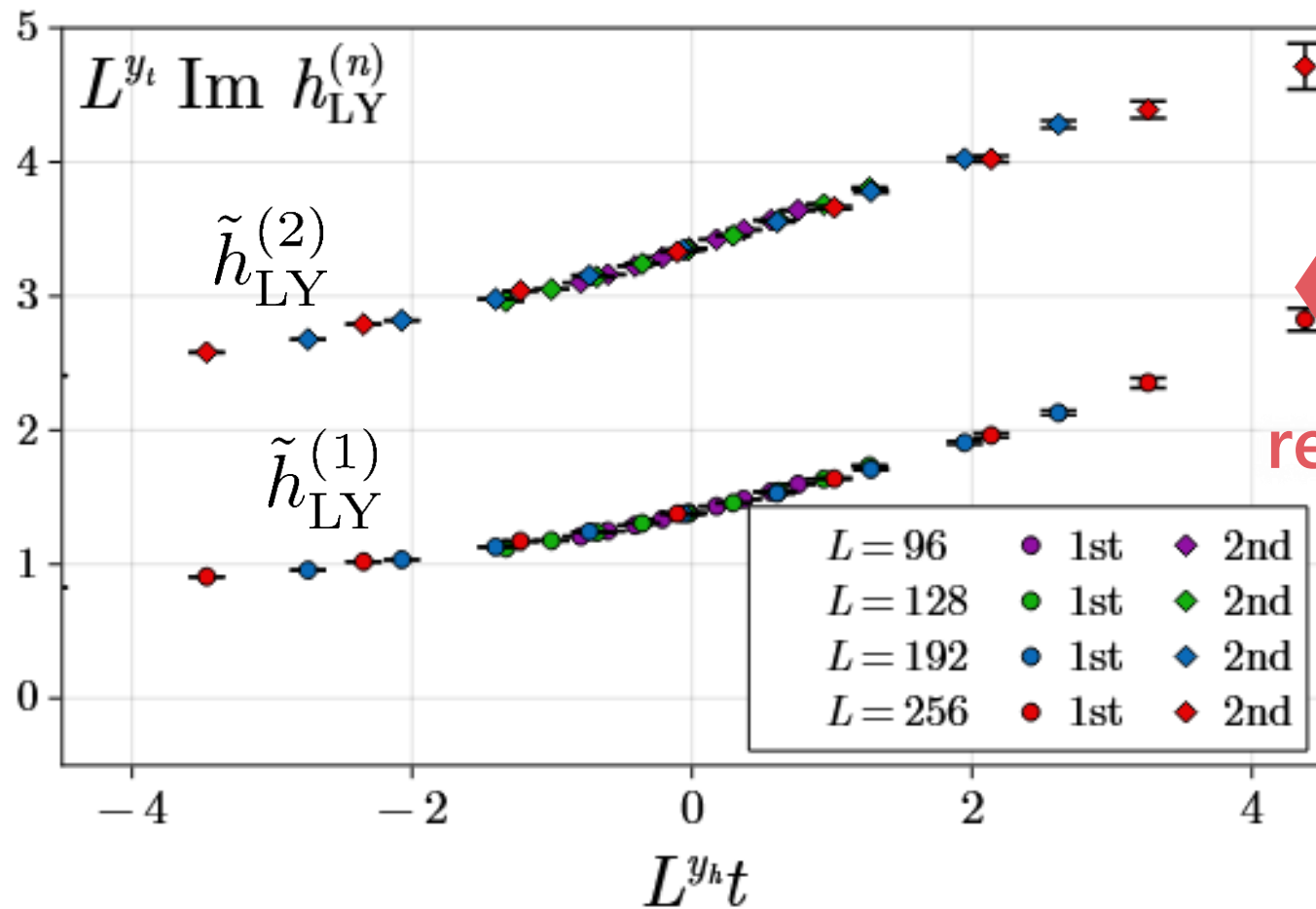
$$Z = Z_{\text{sing}} \times Z_{\text{reg}}$$

LYZ in the scaling region on finite volume

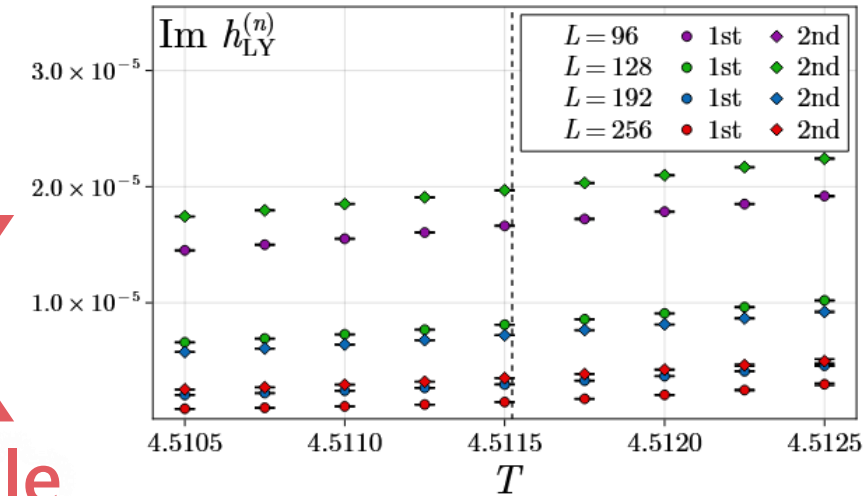
$$Z(t, h, L^{-1}) \sim \tilde{Z}_{\text{sing}}(L^{y_t} t, L^{y_h} h) = 0$$



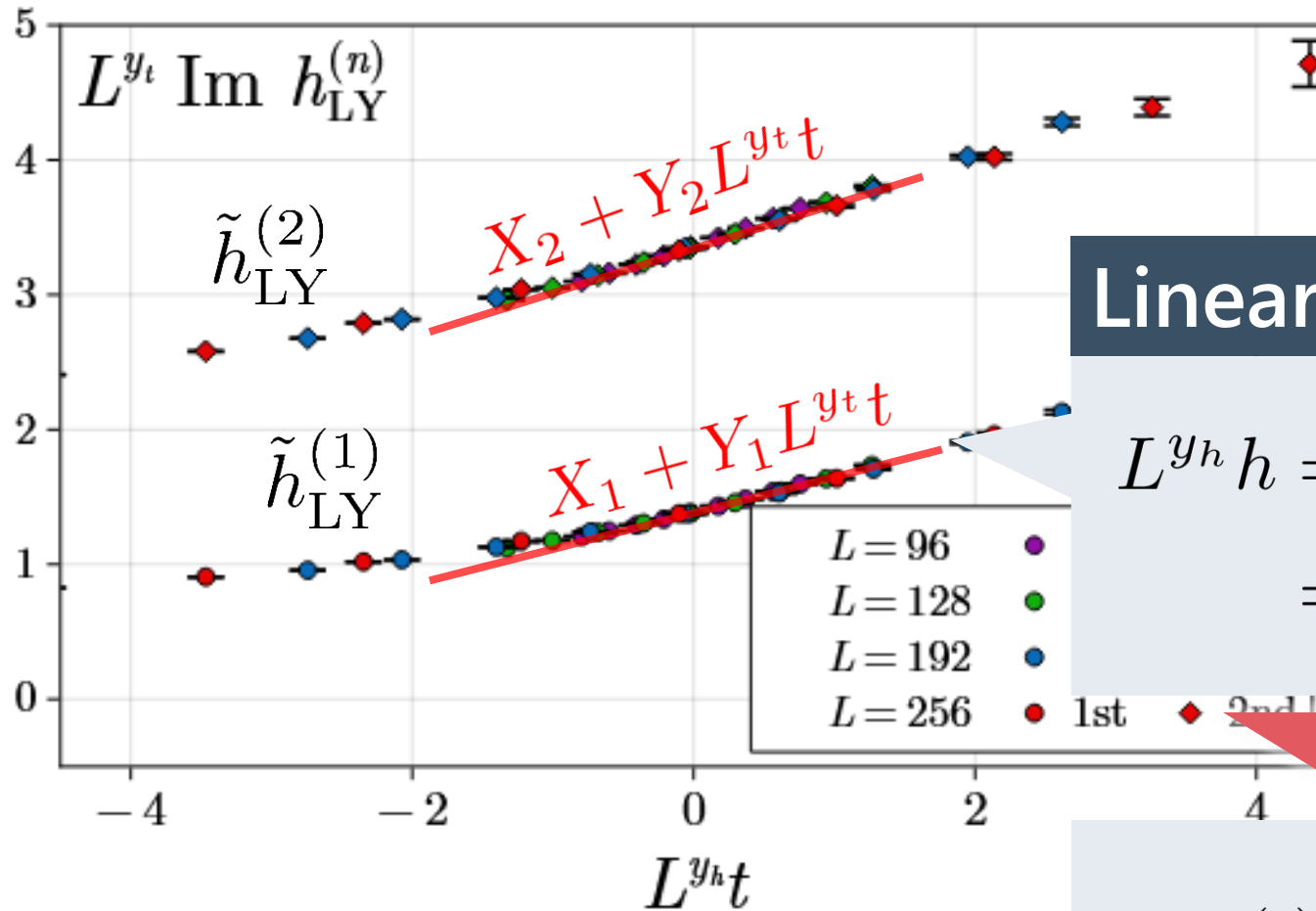
$$L^{y_h} h_{\text{LY}}^{(n)}(t) = \tilde{h}_{\text{LY}}^{(n)}(L^{y_t} t)$$

LYZ near $t = 0$ 

rescale



LYZ near $t = 0$



Linear Approx. at $t = 0$

$$L^{y_h} h = \tilde{h}_{\text{LY}}^{(i)}(L^{y_t} t)$$

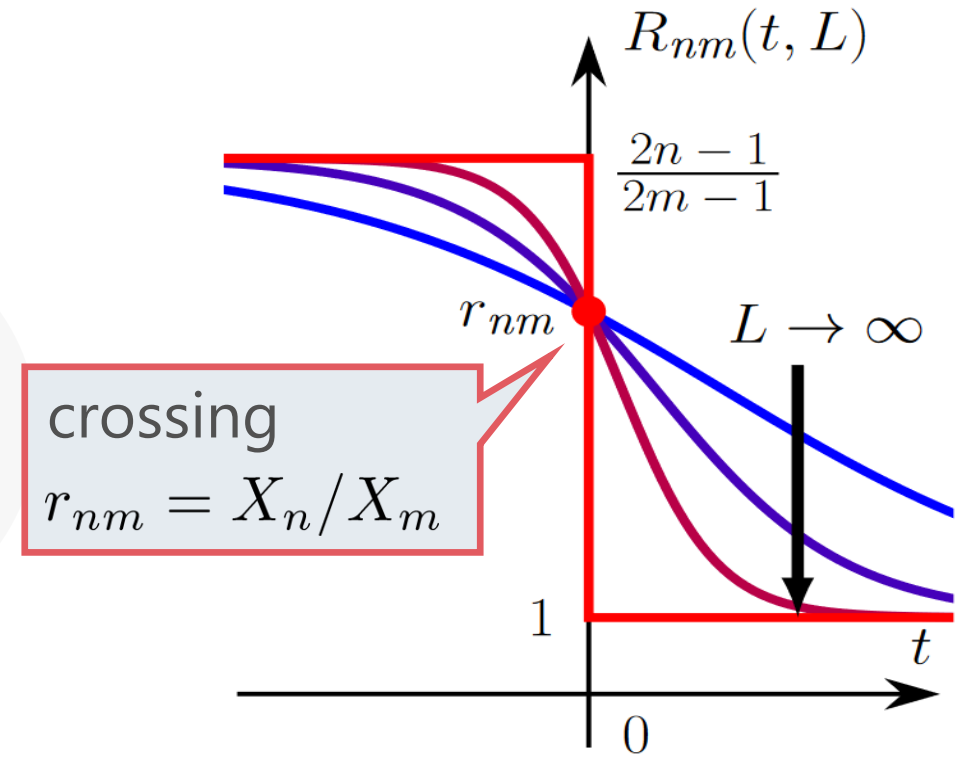
$$= X_i + Y_i L^{y_t} t + \mathcal{O}(t^2)$$

$$R_{nm}(t) = \frac{X_n}{X_m} \left(1 + C_{nm} t L^{y_t} + \mathcal{O}(t^2) \right)$$

LYZ Ratios $R_{nm}(t) = \frac{h^{(n)}(t)}{h^{(m)}(t)}$

$$R_{n1}(t) \xrightarrow{V \rightarrow \infty} \begin{cases} 2n - 1 & t < 0 \text{ (1st order)} \\ 1 & t > 0 \text{ (crossover)} \end{cases}$$

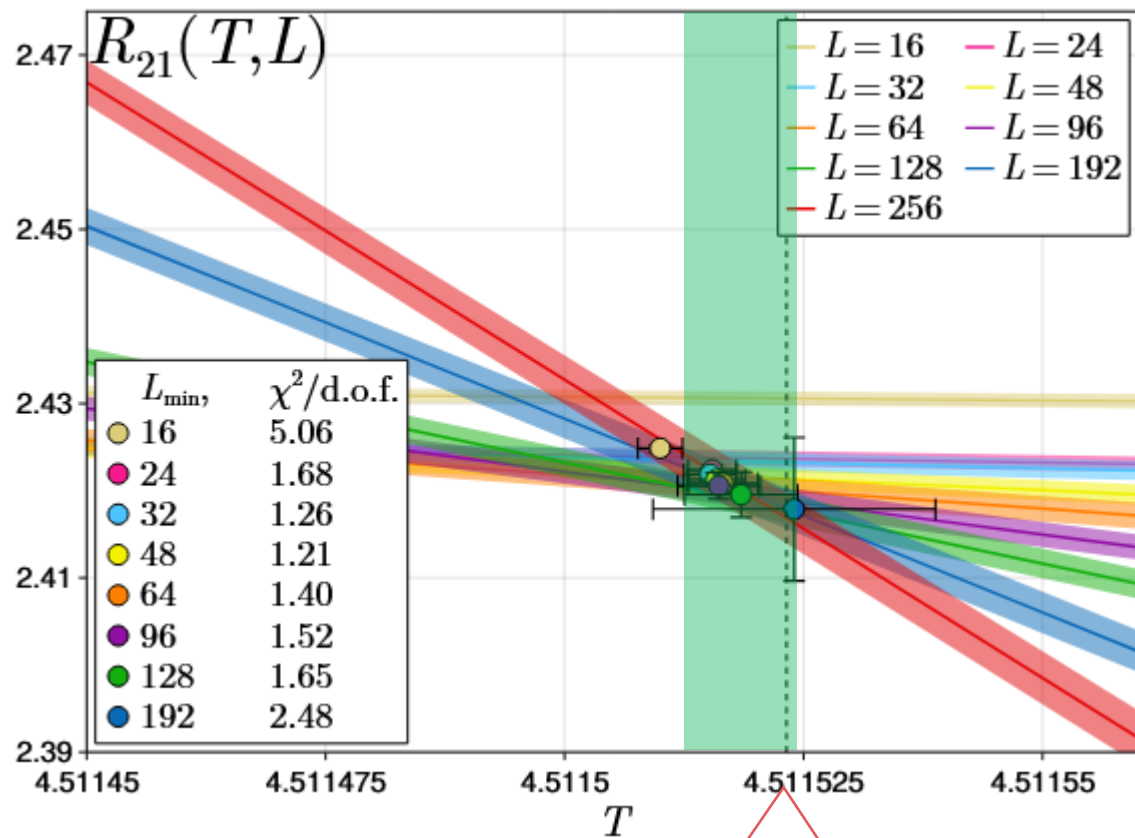
$$R_{nm}(t) = r_{nm} \left(1 + C_{nm} t L^{y_t} + \mathcal{O}(t^2) \right) \text{ near } t = 0$$



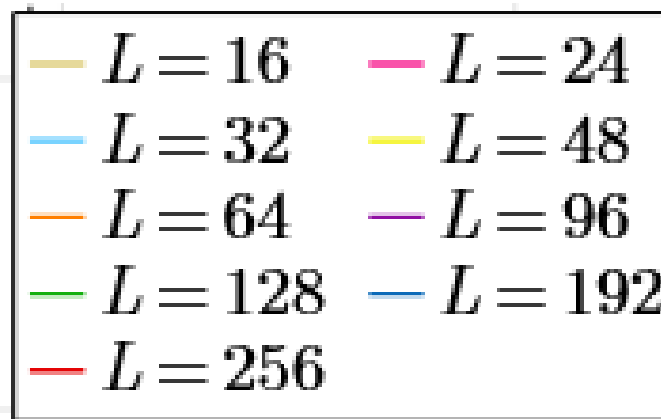
- $R(0)$ is L independent, the universal value.
- Intersection point of various L gives the CP.
- Reminiscent of Binder-cumulant analysis

LYZR in 3d-Ising

LYZR

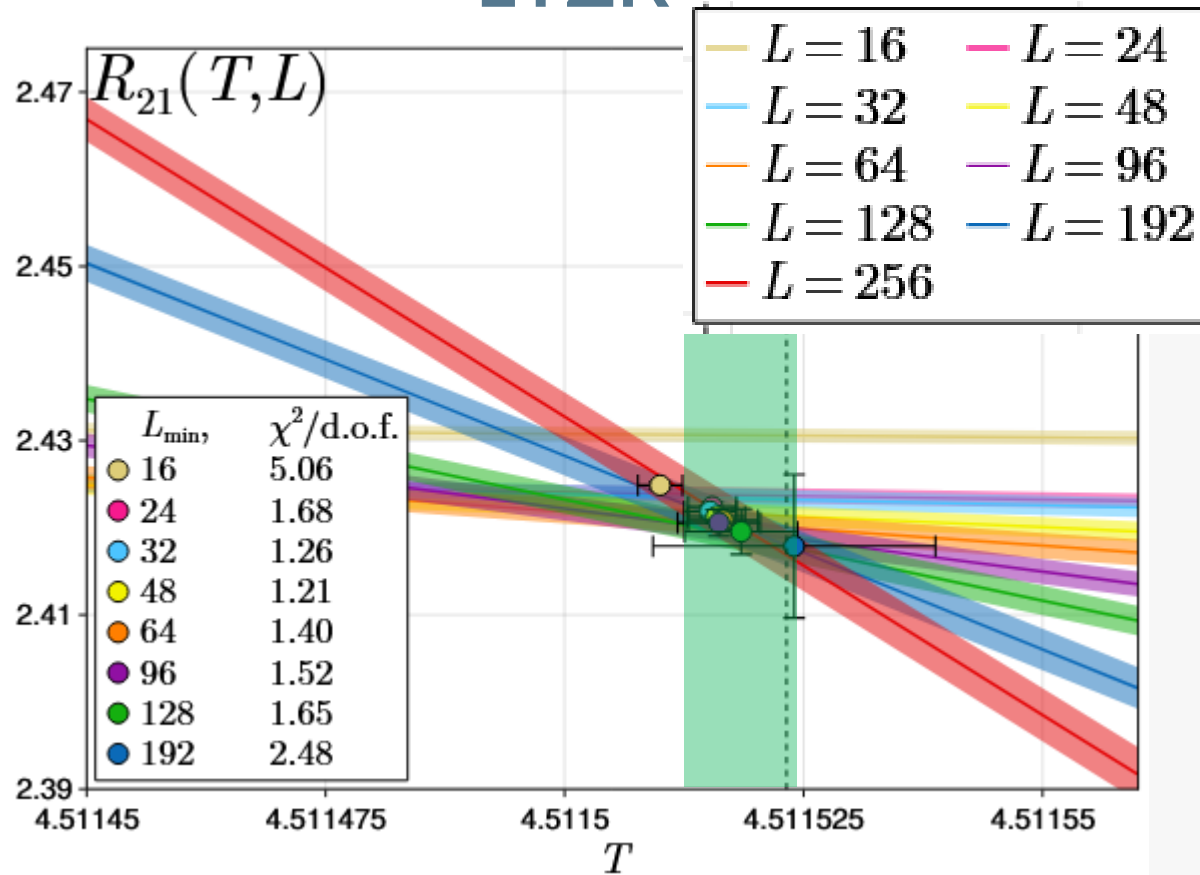


T_C Ferrenberg ('18)

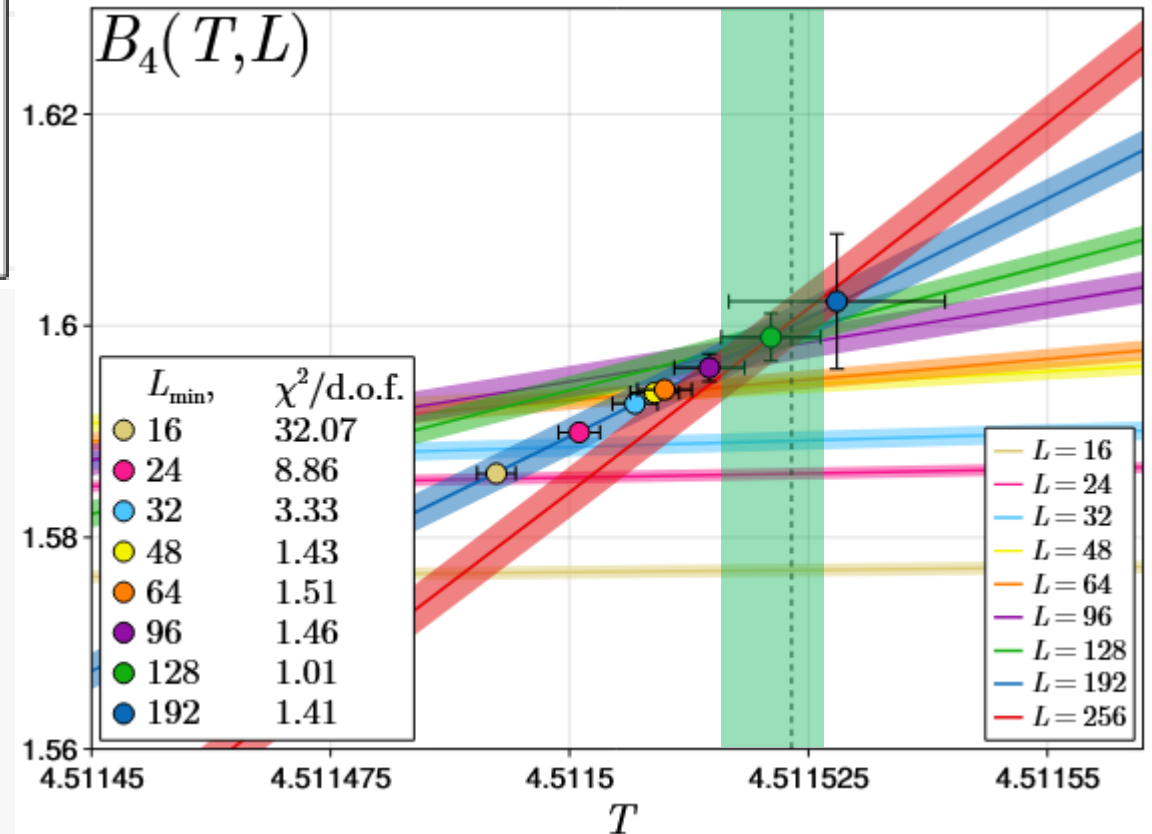


LYZR vs Binder Cumulant in 3d-Ising

LYZR



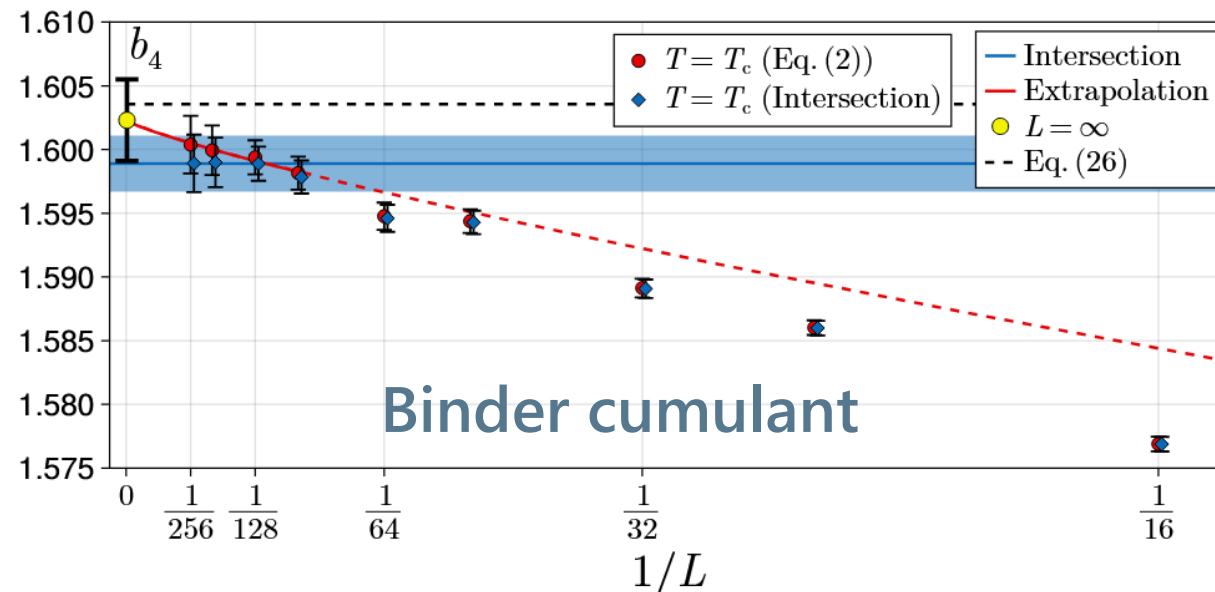
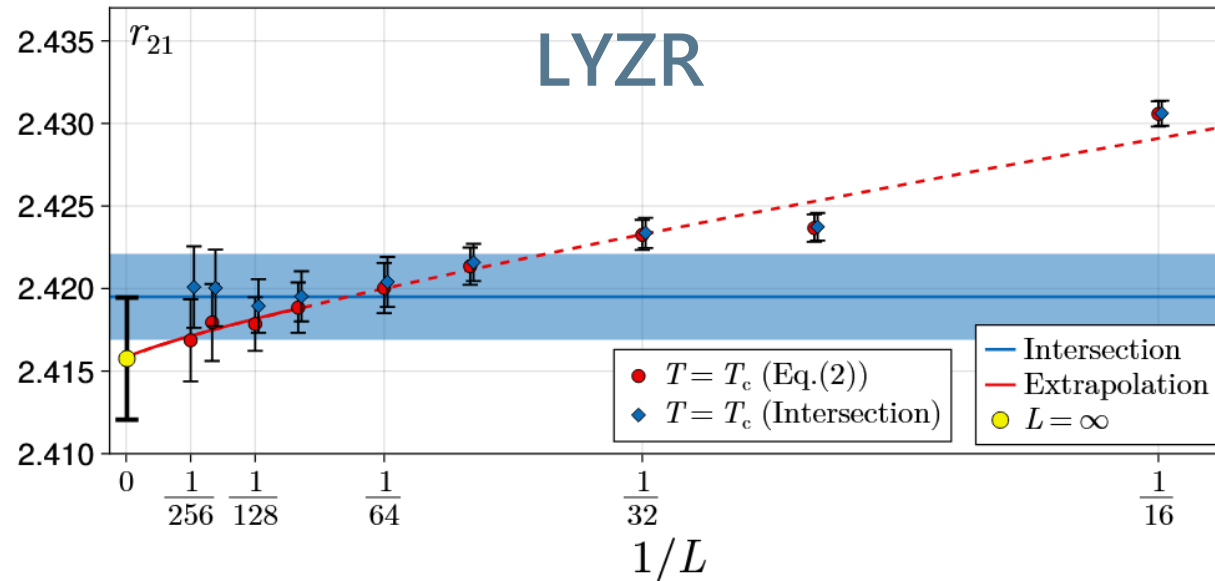
Binder cumulant



Faster convergence of the violation of FSS in LYZ?

Convergence at $T = T_c$

Wada, MK, Kanaya, arXiv:2508. 20422



Red: T_c of Ferrenberg ('18)
Blue: T_c of intersection point

$$R_{21}(0, L) = r_{21}(1 + cL^{-\omega})$$

In 3d-Ising, violation of FSS is more quickly suppressed in the LYZR than the Binder cumulant for $L \rightarrow \infty$.

$$Z = Z_{\text{sing}} \times Z_{\text{reg}}$$

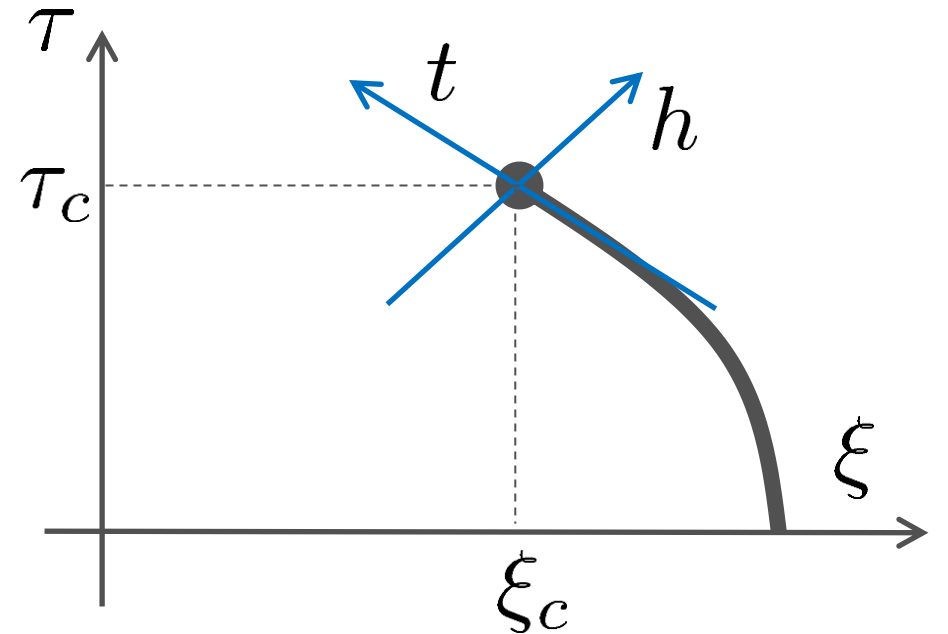
CP in a General System

- CP on a $\tau - \xi$ plane
- **LYZ on the complex ξ plane**

$$\begin{pmatrix} t \\ h \end{pmatrix} = \begin{pmatrix} a_{11} & a_{12} \\ a_{21} & a_{22} \end{pmatrix} \begin{pmatrix} \tau - \tau_c \\ \xi - \xi_c \end{pmatrix} = A \begin{pmatrix} \delta\tau \\ \delta\xi \end{pmatrix}$$

$$L^{y_h} h^{(n)}(t) \simeq X_n + Y_n L^{y_t} t$$

$$\bar{y} = y_t - y_h = -0.894$$



$$Z_{\text{sing}}(\tau, \xi) = \tilde{Z}_{\text{Ising}}(t(\tau, \xi), \xi(\tau, \xi))$$

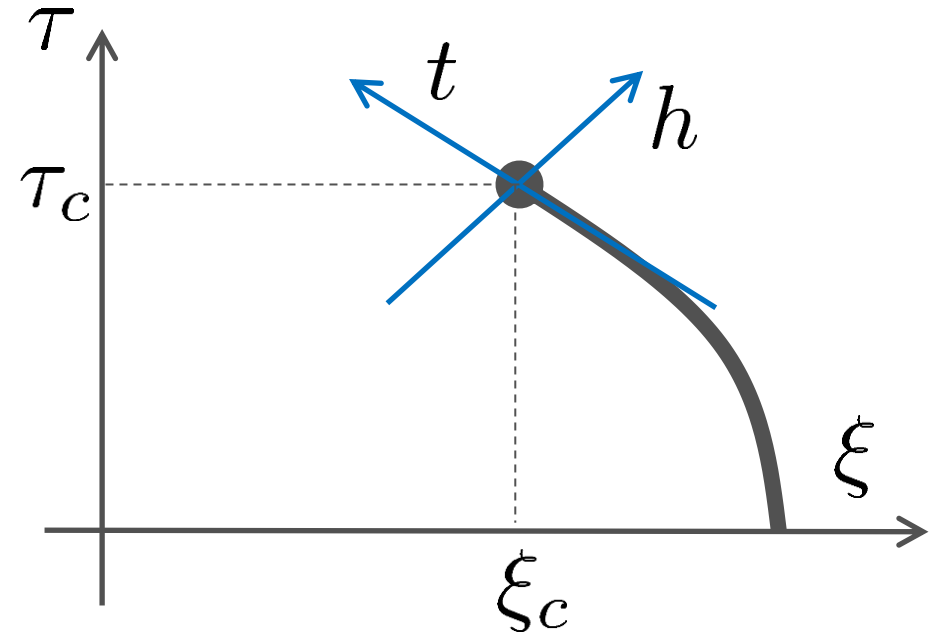
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CP in a General System


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$$L^{y_h} h^{(n)}(t) \simeq X_n + Y_n L^{y_t} t$$



$$\begin{cases} \xi_{\text{R}}^{(n)} = \xi_c - \frac{a_{21}}{a_{22}} \delta\tau + \mathcal{O}(L^{2\bar{y}}) \\ \xi_{\text{I}}^{(n)} = \frac{X_n}{a_{22}} L^{-y_h} + \frac{\det AY_n}{a_{22}^2} \delta\tau L^{\bar{y}} + \mathcal{O}(L^{2\bar{y}}) \end{cases}$$

$L \rightarrow \infty$

**generalization
to finite V**

LY Edge Singularity

$$\begin{cases} \text{Re}\xi_{\text{LYES}} \simeq c_1 \tau \\ \text{Im}\xi_{\text{LYES}} \simeq c_2 \tau^{\beta\delta} \end{cases}$$

Stephanov, 2006

LYZ Ratios for General CP

$$\bar{y} = y_t - y_h = -0.894$$

LYZ Ratio

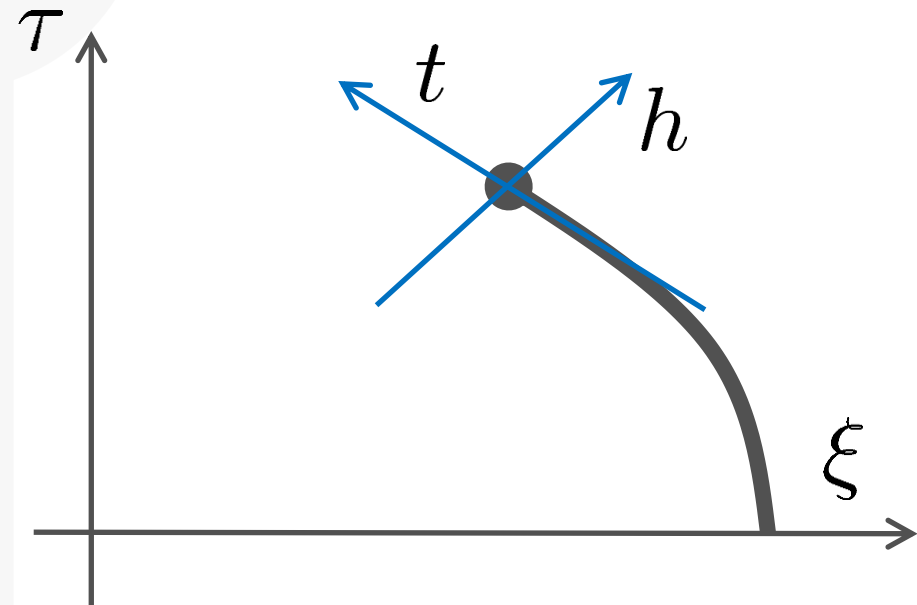
$$R_{nm}(\tau) = \frac{\xi_I^{(n)}(\tau)}{\xi_I^{(m)}(\tau)} = r_{nm} \left(1 + CL^{y_t} \delta\tau + \mathcal{O}(\delta\tau^2) \right) \left(1 + DL^{2\bar{y}} + \mathcal{O}(L^{4\bar{y}}) \right)$$

nonzero for $a_{12} \neq 0$

$$r_{nm} = \frac{X_n}{X_m}, \quad C = \frac{\det A}{a_{22}} \left(\frac{Y_2}{X_2} - \frac{Y_1}{X_1} \right), \quad D = \frac{a_{12}^2}{a_{22}^2} (Y_1^2 - Y_2^2)$$

$$\begin{pmatrix} t \\ h \end{pmatrix} = \begin{pmatrix} a_{11} & a_{12} \\ a_{21} & a_{22} \end{pmatrix} \begin{pmatrix} \tau - \tau_c \\ \xi - \xi_c \end{pmatrix} = A \begin{pmatrix} \delta\tau \\ \delta\xi \end{pmatrix}$$

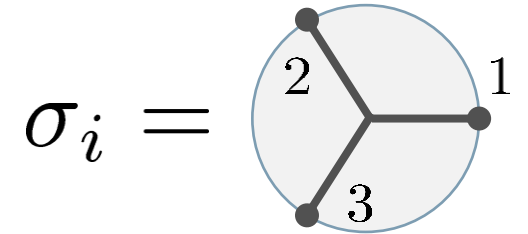
r_{nm} are universal constants specific to universality classes.



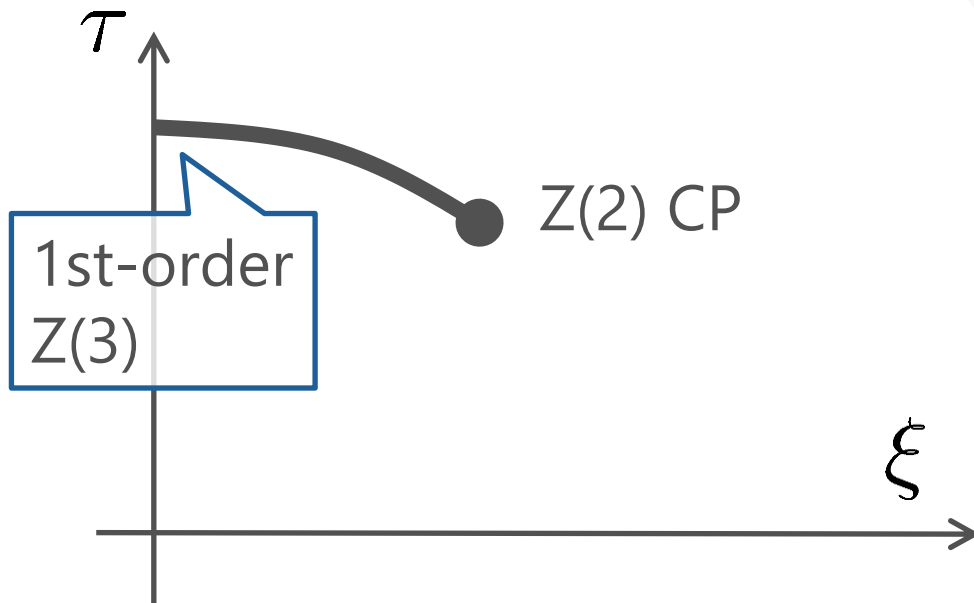
Numerical Analysis: 3d 3-State Potts Model

$$H = -\tau \sum_{\langle i,j \rangle} \delta_{\sigma_i, \sigma_j} - \xi \sum_i \delta_{\sigma_i, 1} \quad \sigma_i = 1, 2, 3$$

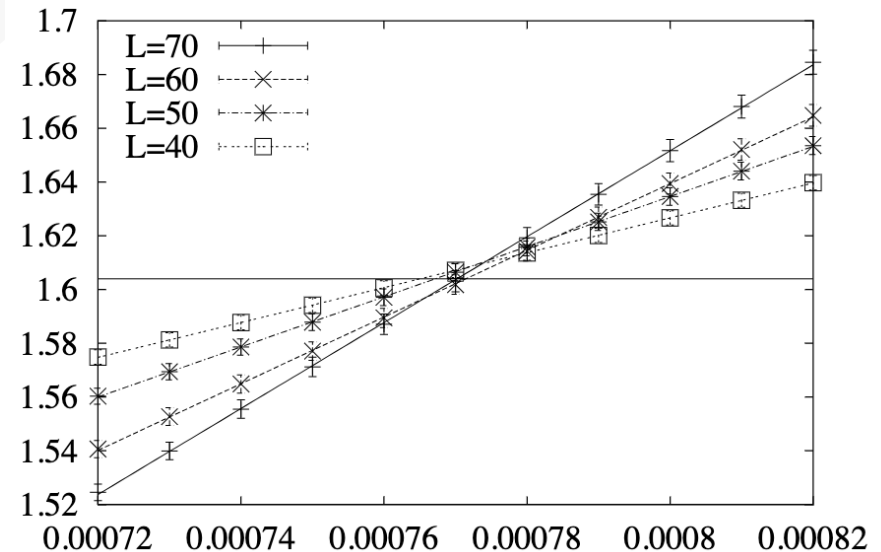
Monte-Carlo + reweighting



Phase Diagram



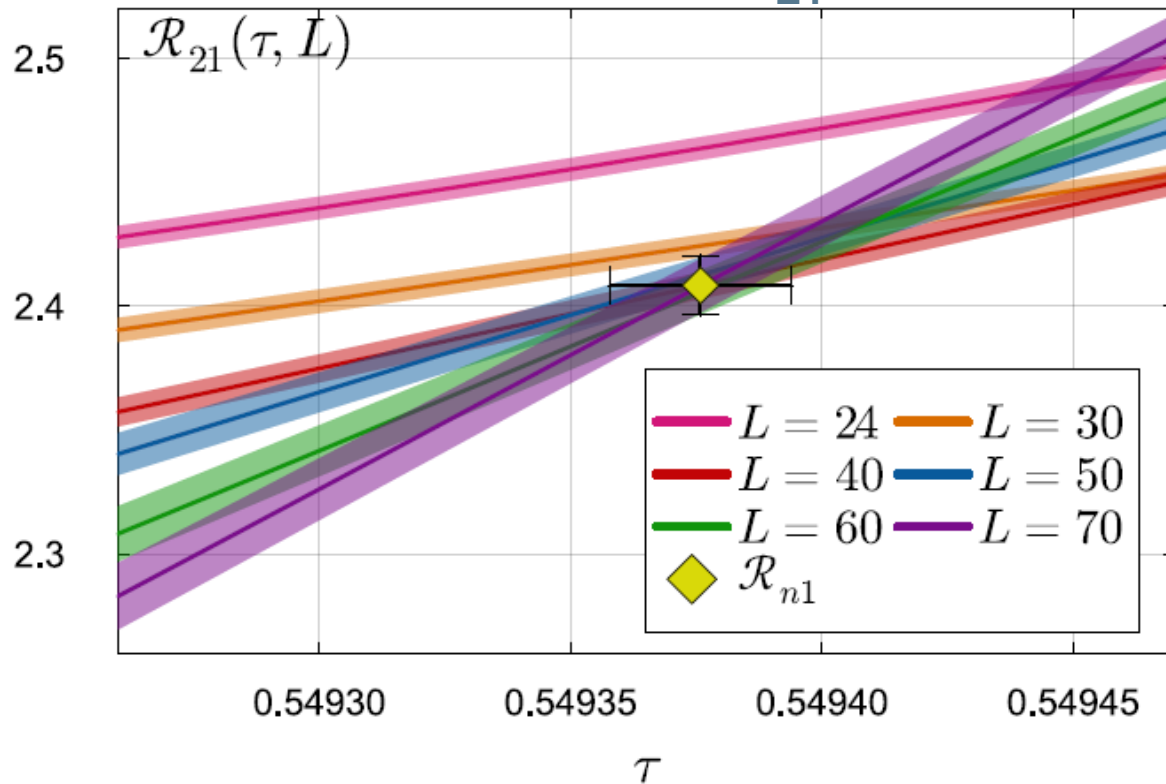
Binder-Cumulant Analysis



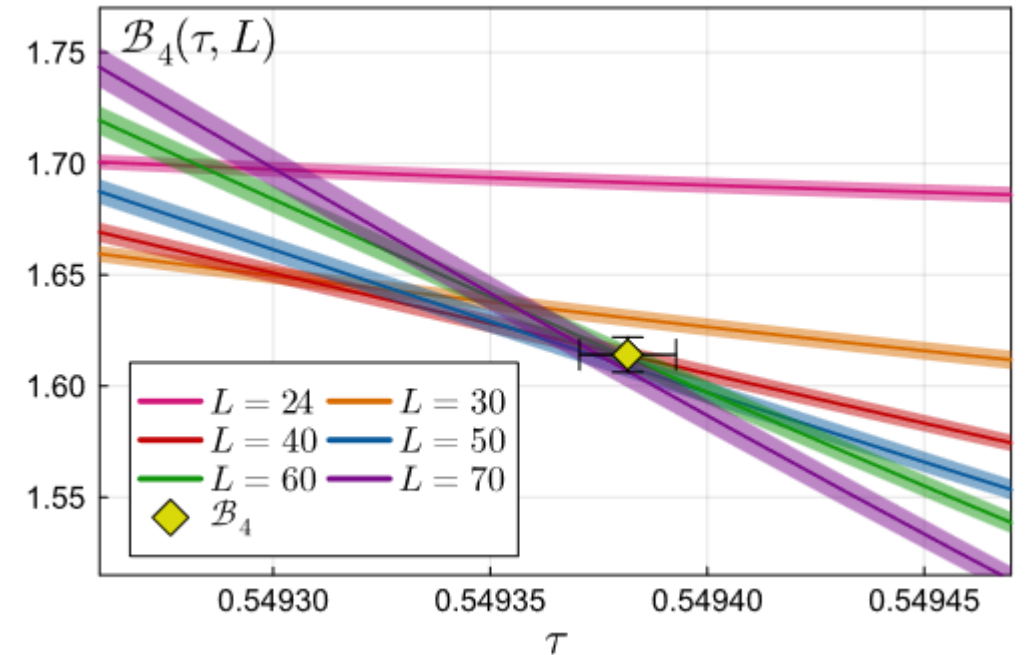
Karsch, Stickan, 2000

3d 3-State Potts Model: LYZ Ratio

LYZ Ratio (\mathcal{R}_{21})



Binder Cumulant



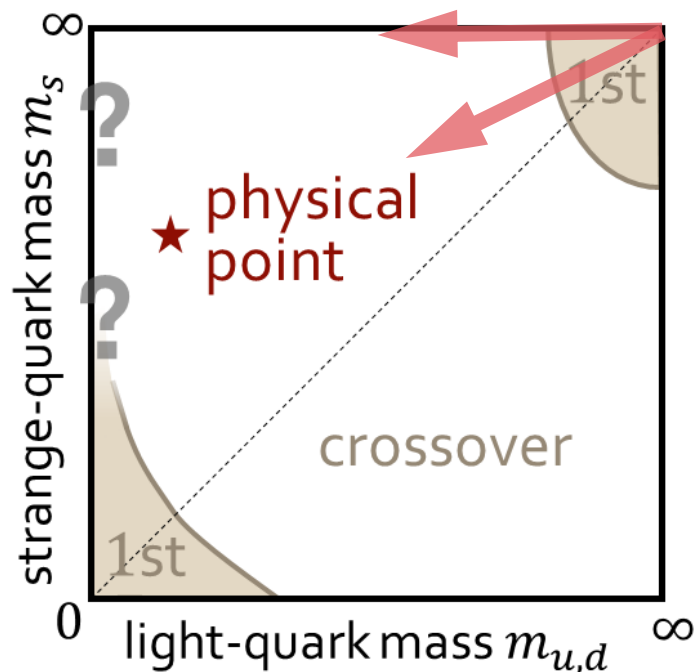
Wada, MK, Kanaya, PRL, '25

fit data	τ_c	y_t	r_{n1} or b_4	χ^2/dof
\mathcal{R}_{21}	0.549375(18)	1.53(19)	2.408(12)	0.38
\mathcal{B}_4	0.549382(11)	1.63(13)	1.614(8)	0.69

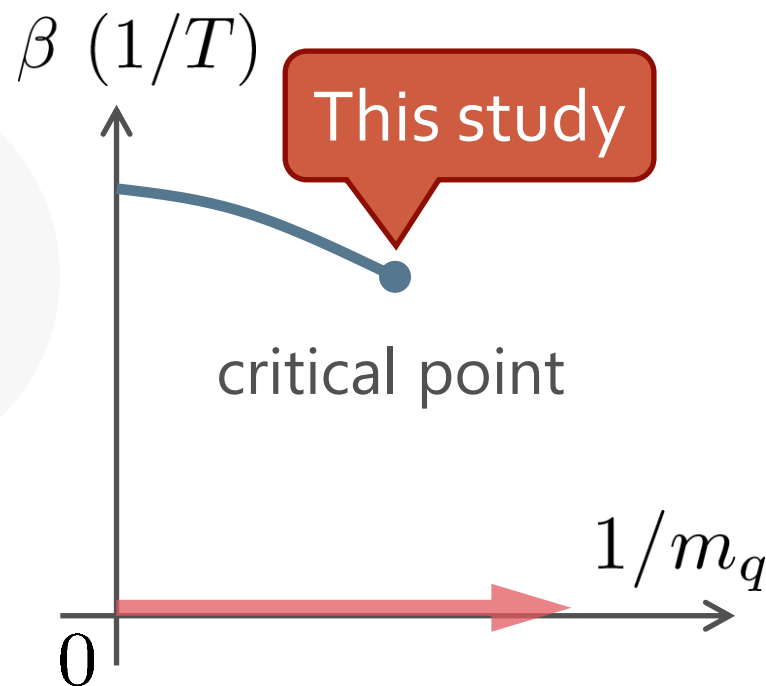
\mathcal{R}_{n1} and \mathcal{B}_4 give the same value of τ_c within statistics.

CP in Heavy-Quark QCD

Columbia Plot



Phase Diagram



CP in heavy-quark QCD

– $\mu_q = 0$ & large m_q

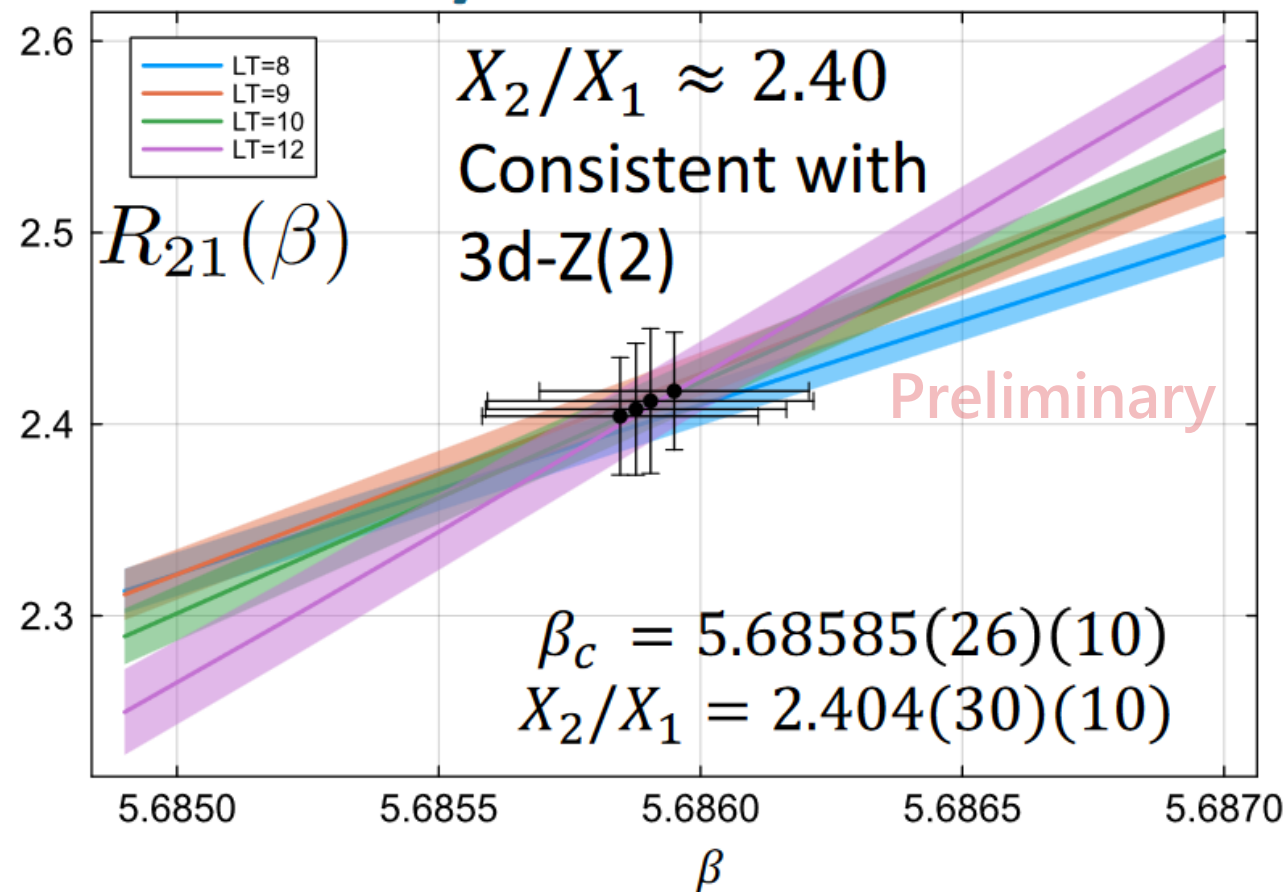
➤ Easy to handle in lattice simulations!

➤ We study the LYZ around the HQ-QCD-CP.

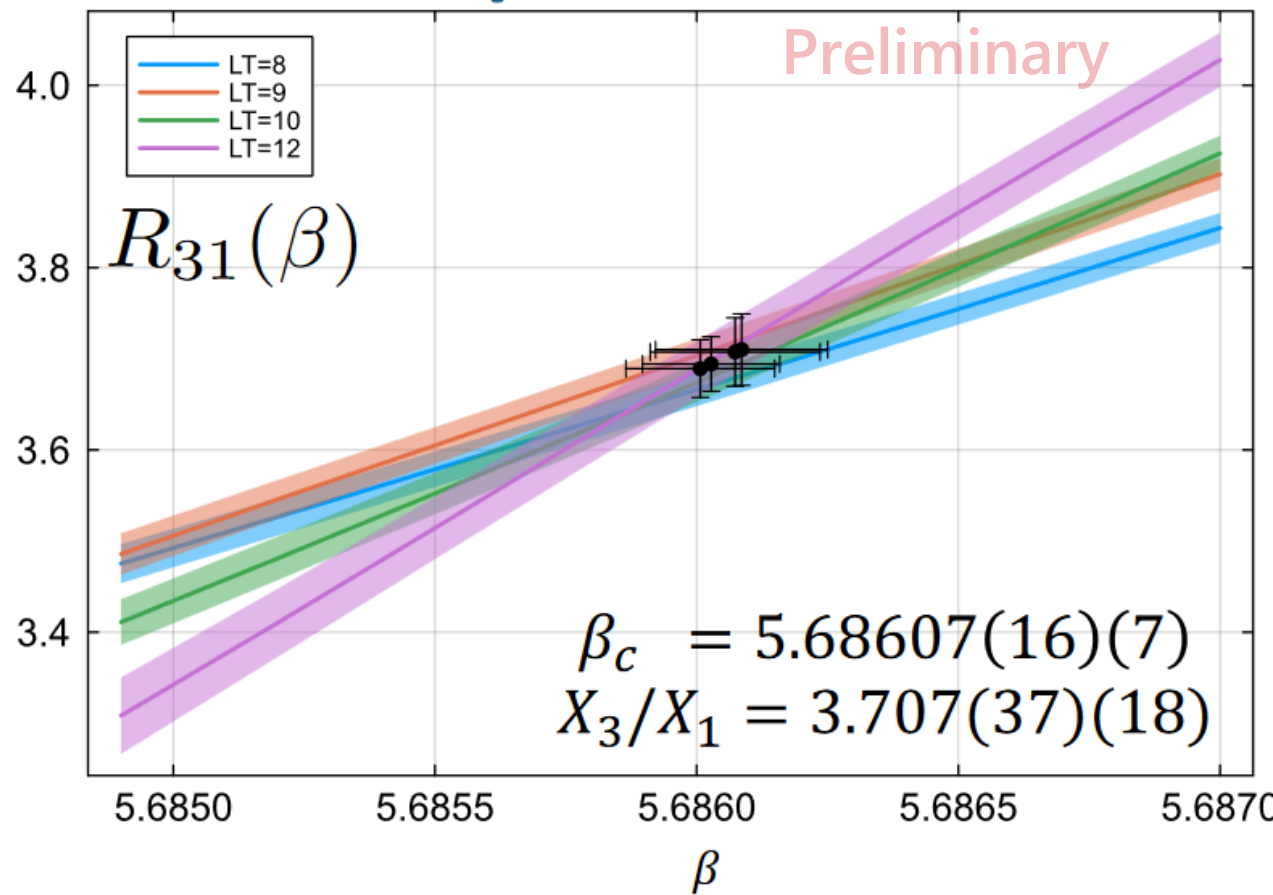
LYZ Ratio

$N_t = 4$, HPE-NLO

2nd/1st LYZ Ratio



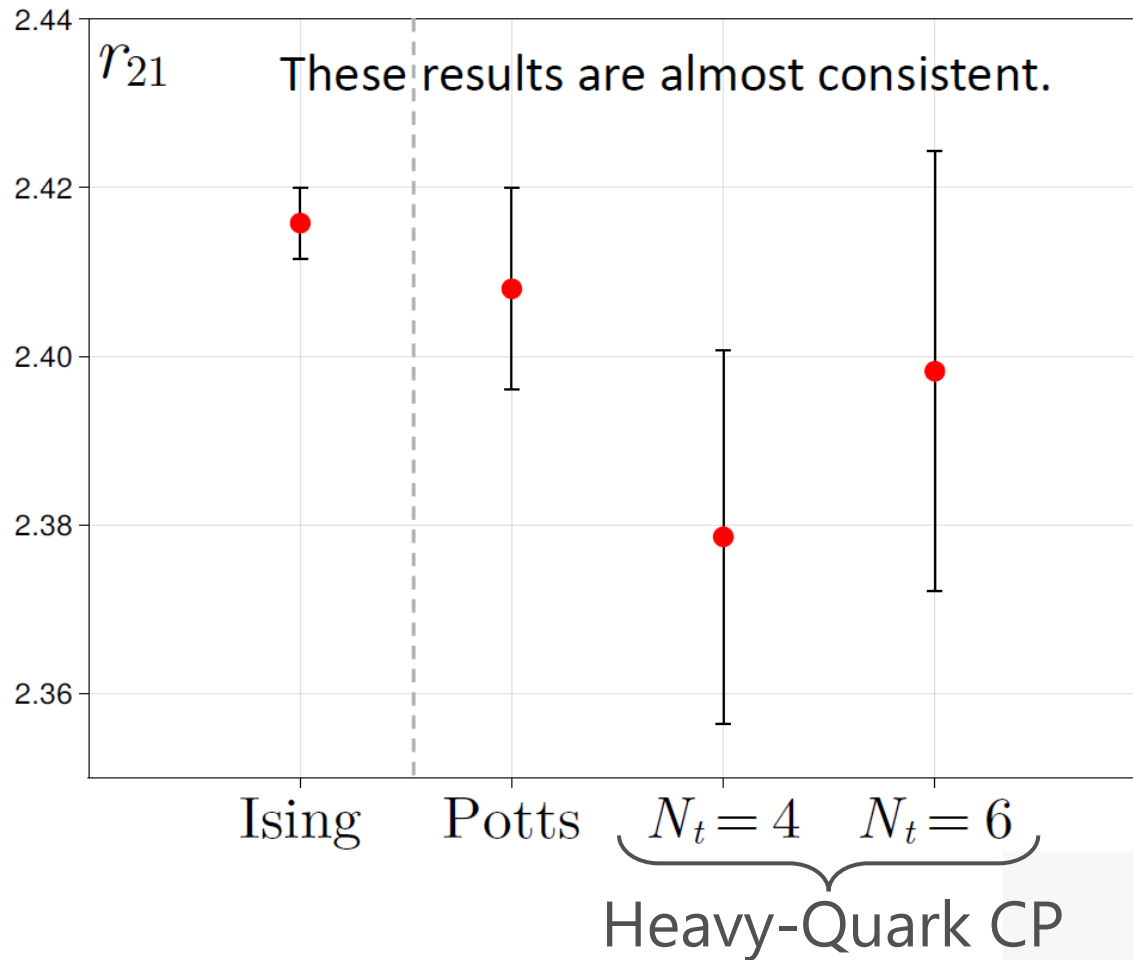
3rd/1st LYZ Ratio



Consistent with the Binder-cumulant analysis $\beta_c = 5.68578(22)$.

LYZR in Various Models

LYZR at the intersection point



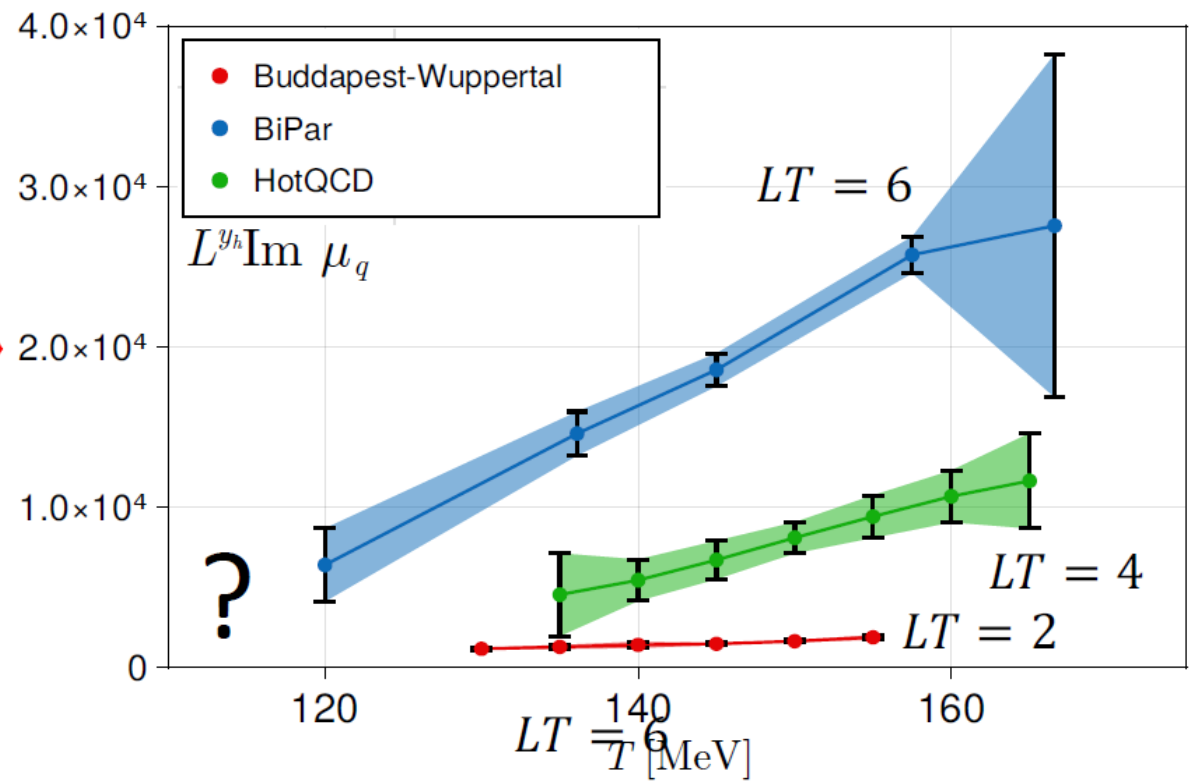
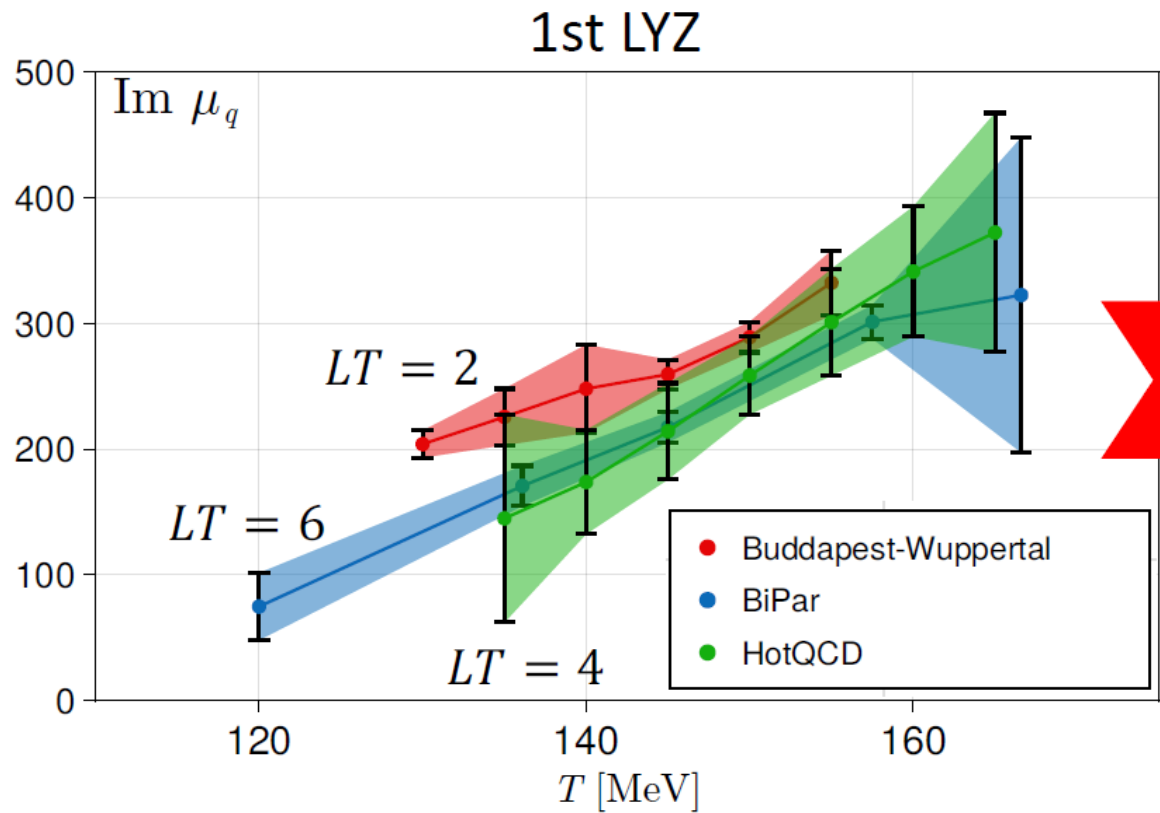
CPs in 3d-Z(2) universality class

- Ising: 2508.20422
- Potts: PRL, '25
- HQ-CP: 2501.18904

Consistent within statistics ➔ Confirmation of universality

Application to QCD Critical Point?

Single-LYZ analysis for 2+1 flavor QCD



Intersection around $T \simeq 100$ MeV?

Yet lower T data are mandatory.

Summary

Lee-Yang-zero ratios

A new tool for **the CP searches** utilizing the **finite-size scaling of Lee-Yang zeros**.

$$R_{nm}(t) = \frac{h^{(n)}(t)}{h^{(m)}(t)}$$

Outlook

- Determination of r_{nm} in each universality class
- More sophisticated utilization of $R_{nm}(t, L)$
- Other quantities: ξ_c , mixing matrix A , etc.
- Application to **QCD-CP** and other CPs



backup