

Determination of the ω - and η' -nucleus optical potential

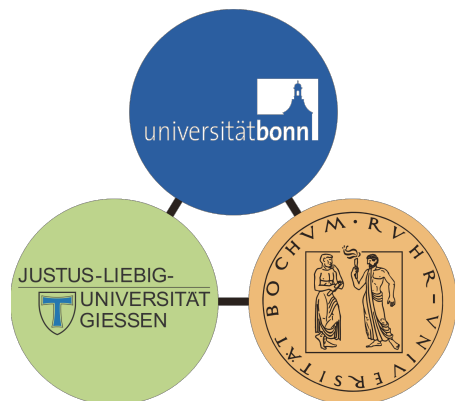
Volker Metag
II. Physikalisches Institut



Outline:

- ◆ theoretical predictions for in-medium modifications of hadron properties
- ◆ exp. approaches and results on the real part of the ω , η' - nucleus potential
- ◆ exp. approaches and results on the imaginary part of the ω , η' - nucleus potential
- ◆ search for meson-nucleus bound states
- ◆ summary & outlook

*funded by the DFG within SFB/TR16



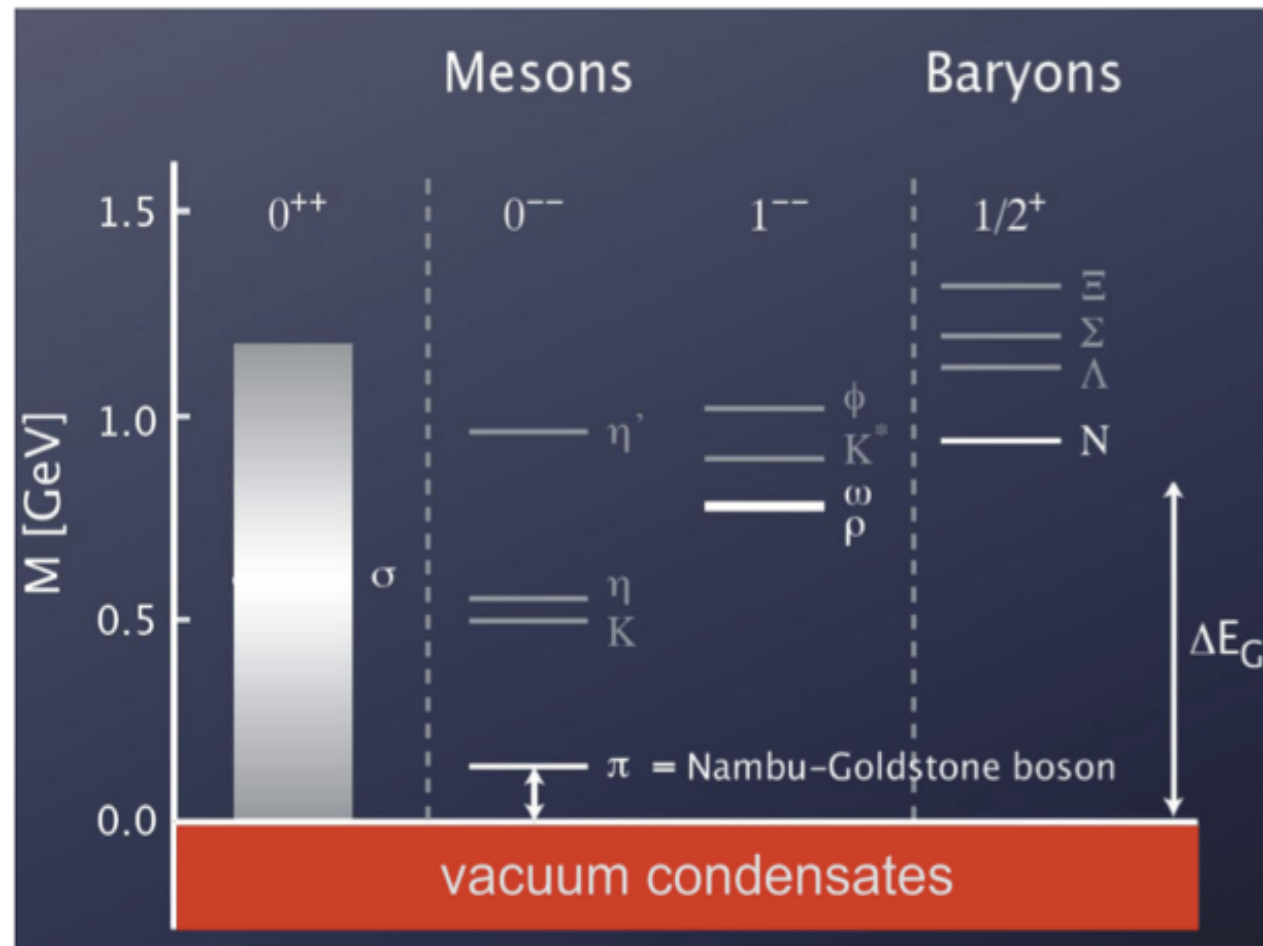
Mesons in Nuclei
International workshop
31.7.-2.8.2016, Kyoto, Japan

HIC | **FAIR**
for
Helmholtz International Center

hadron masses

how do hadron properties (mass, width) change
in a dense nuclear medium ??

J. Wambach



- mesons = excitations of the QCD vacuum
- QCD vacuum: complicated structure characterized by condensates
- in the nuclear medium: condensates are changed

⇒ **change of the hadronic excitation energy spectrum**

V. Bernard and U.-G. Meißner,
NPA 489 (1988) 647

G.E. Brown and M. Rho, $\frac{m^*}{m} \approx \frac{\langle \bar{q}q \rangle^*}{\langle \bar{q}q \rangle_0} \approx 0.8 (\rho \approx \rho_0)$
PRL 66 (1991) 2720

T. Hatsuda and S. Lee, $\frac{m_V^*}{m_V} = (1 - \alpha \frac{\rho}{\rho_0}); \alpha \approx 0.18$
PRC 46 (1992) R34

widespread theoretical and experimental
activities to search for in-medium
modifications of hadrons

symmetry breaking pattern for pseudoscalar mesons

V. Bernard, R.L. Jaffe, U.-G. Meissner, NPB 308 (1988) 753

S. Klimt, M. Lutz, U. Vogel, W. Weise, NPA 516 (1990) 429

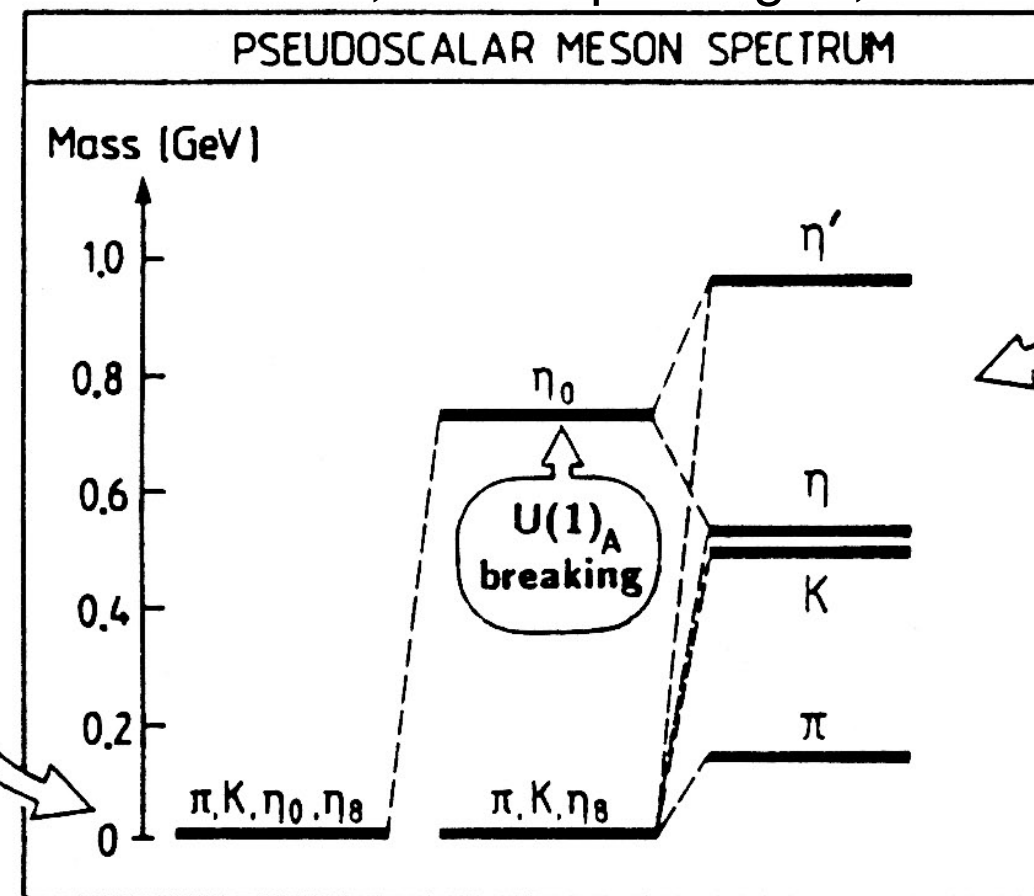
H. Nagahiro et al., PRC 87 (2013) 045201

D. Jido et al., arXiv:1208.0982

mass as a result of
symmetry breaking

SPONTANEOUS
breaking of
 $U(3)_L \otimes U(3)_R$
 $m_u = m_d = m_s = 0$
NINE
PSEUDOSCALAR
GOLDSTONE
BOSONS

W. Weise, NBI Copenhagen, 1993



EXPLICIT breaking
of
 $SU(3)_L \otimes SU(3)_R$
 $m_u = m_d = 5 \text{ MeV}$
 $m_s = 130 \text{ MeV}$

symmetry breaking

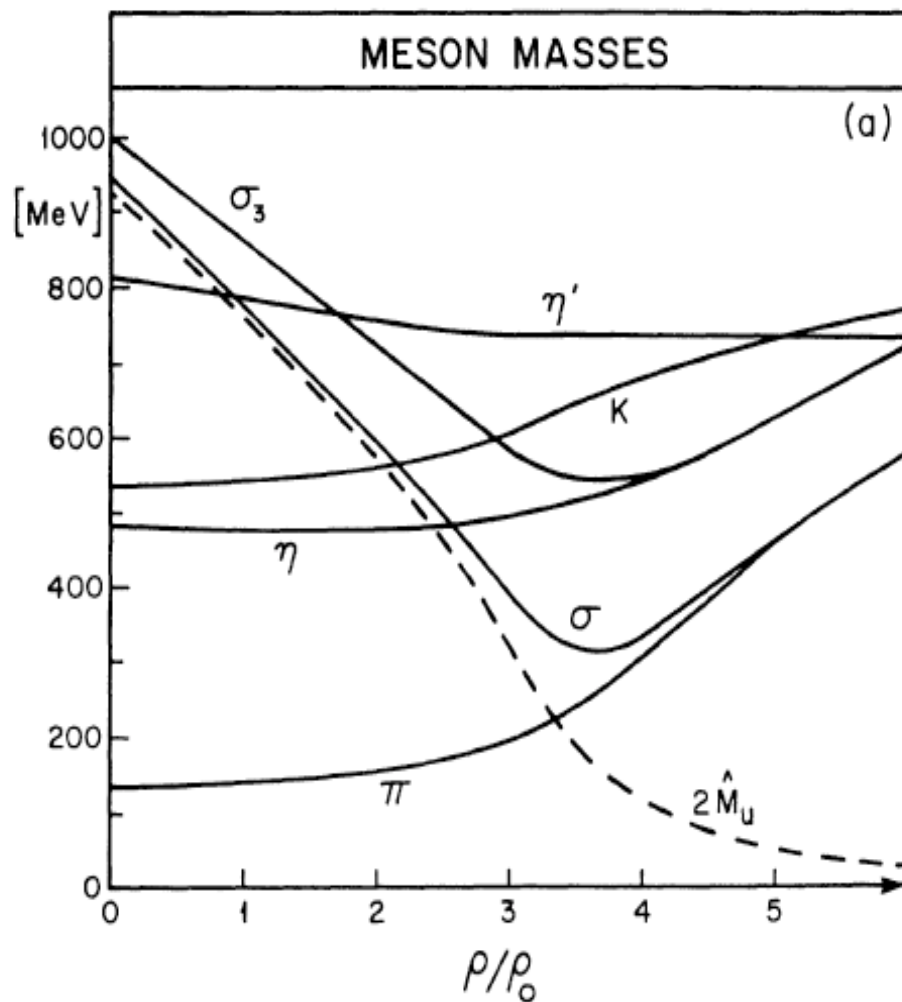
partial symmetry restoration

partial restoration of
chiral symmetry
predicted to occur
in a nucleus $\Rightarrow$ impact
on meson masses ??

hadronic models: predictions for η' in-medium mass

NJL-model

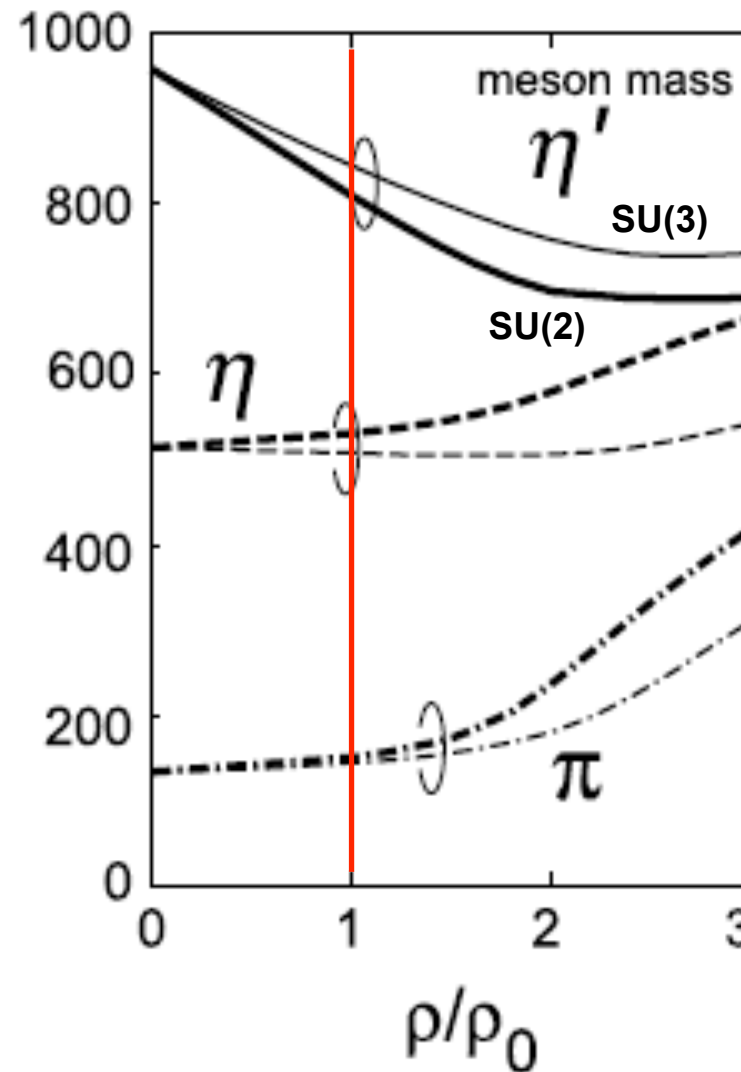
V. Bernard and U.-G. Meissner,
Phys. Rev.D 38 (1988) 1551



almost no dependence of
 η' mass on density

NJL-model

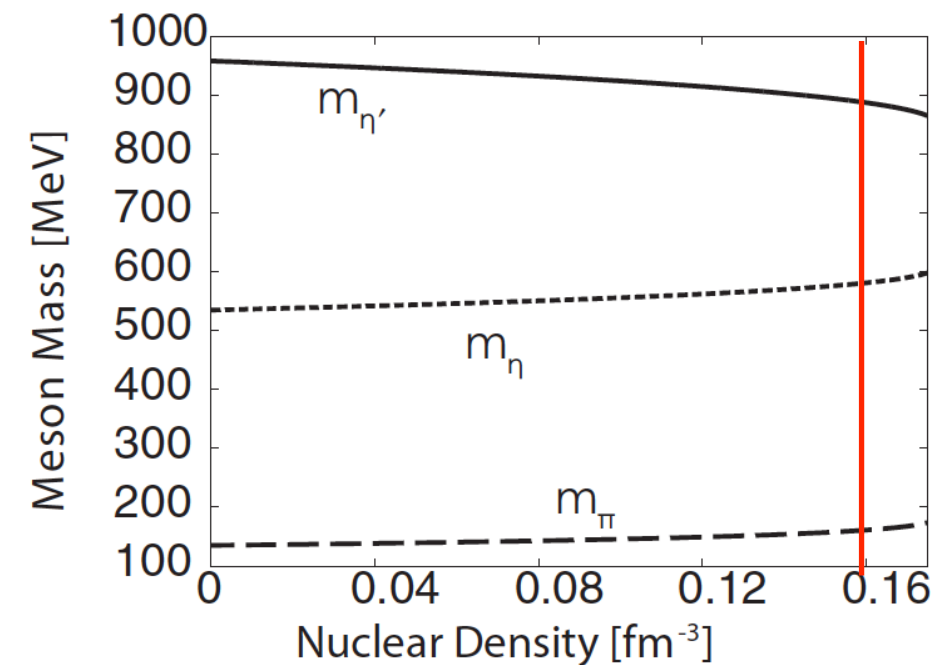
H. Nagahiro et. al,
Phys. Rev. C 74 (2006) 045203



$\Delta m_{\eta'}(\rho_0) \approx -150 \text{ MeV}$
 $\Delta m_{\eta}(\rho_0) \approx +20 \text{ MeV}$

linear σ model

S. Sakai and D. Jido
PRC 88 (2013) 064906



$\Delta m_{\eta'}(\rho_0) \approx -80 \text{ MeV}$

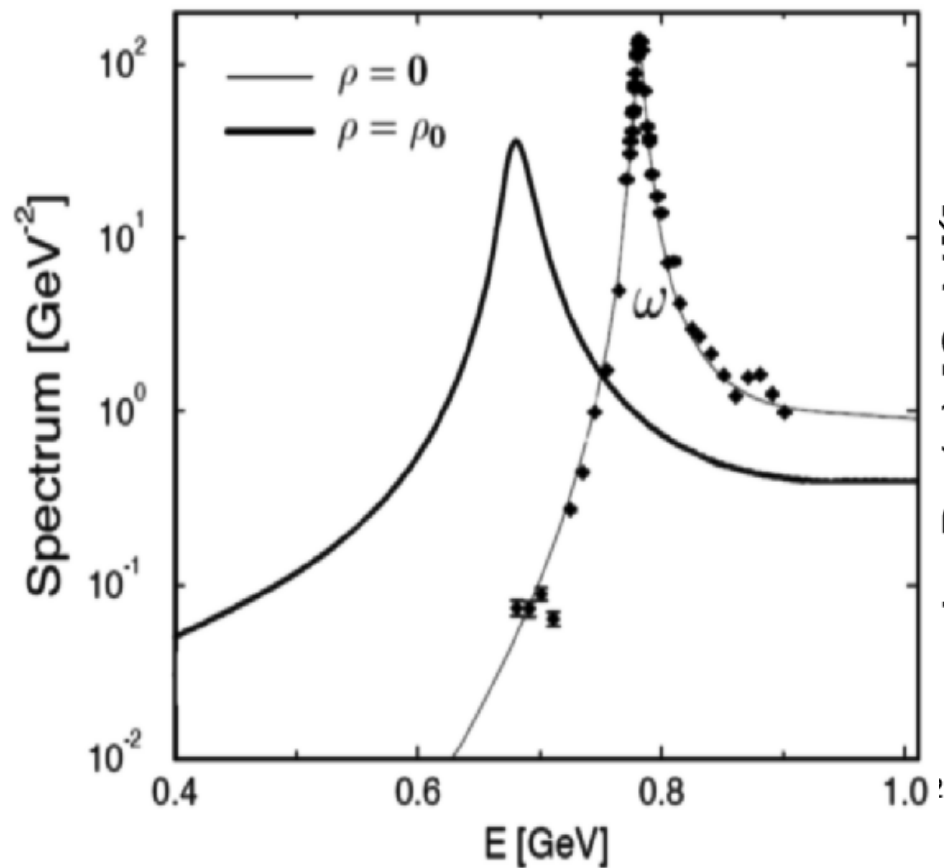
QMC-model

S. Bass and A. Thomas,
PLB 634 (2006) 368

$\Delta m_{\eta'}(\rho_0) \approx -40 \text{ MeV}$
for $\theta_{\eta\eta'} = -20^\circ$

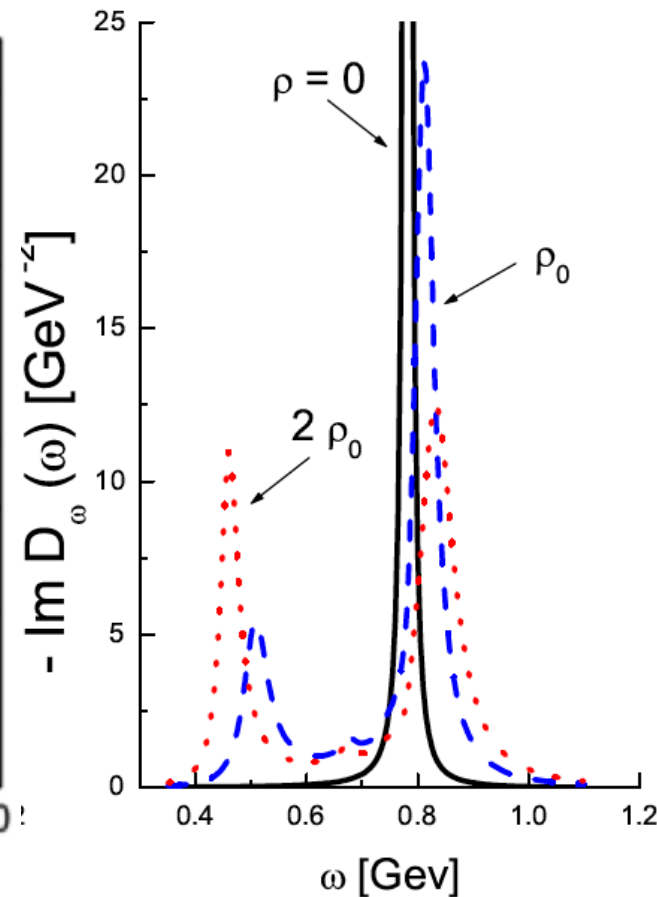
hadronic models: predictions for ω -spectral functions

F. Klingl et al.,
NPA 610 (1997) 297;
NPA 650 (1999) 299



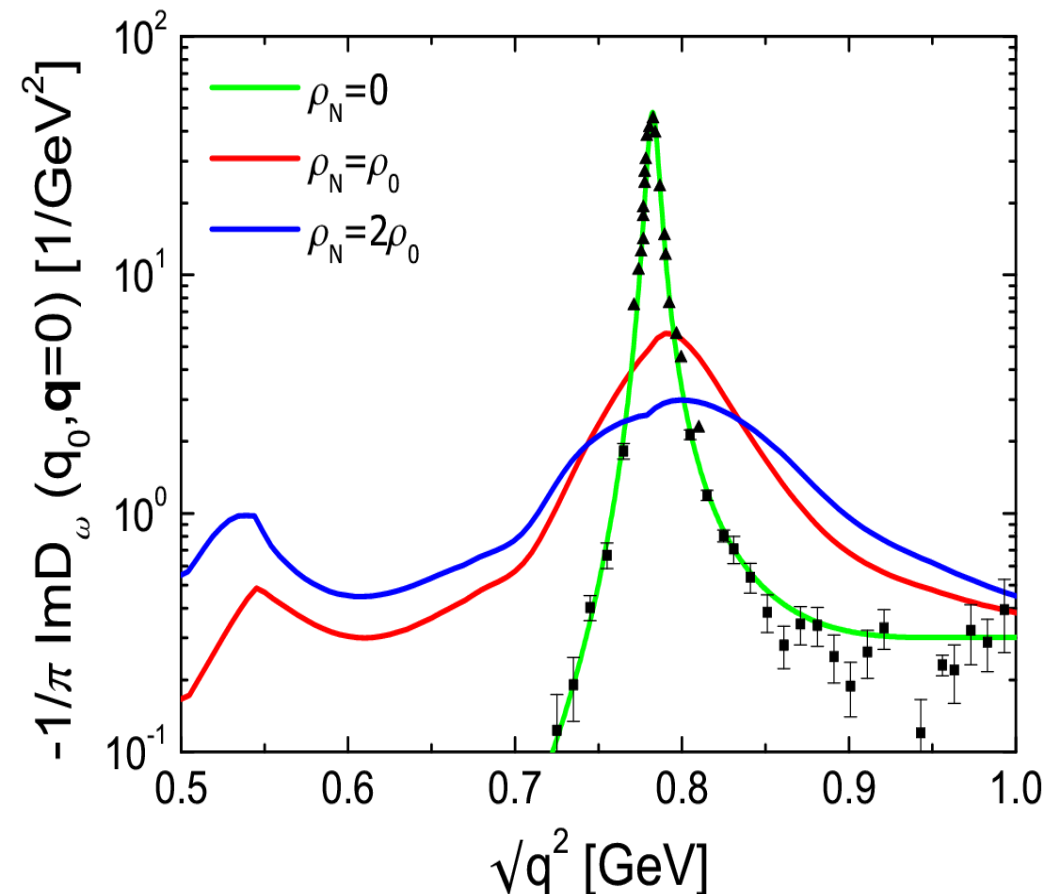
- ◆ lowering of in-medium mass
 - ◆ broadening of resonance with increasing nuclear density
- $\text{Re}(U) \neq 0; \text{Im}(U) \neq 0$

M. Lutz et al.,
NPA 706 (2002) 437



splitting into ω -like and N^*N^{-1} mode due to coupling to nucleon resonances

P. Mühlich et al., NPA 780 (2006) 187



spectral function for ω meson at rest:

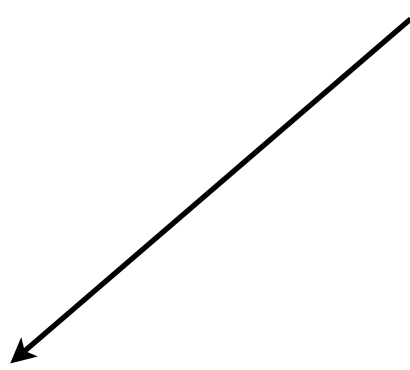
almost no mass shift;
strong in-medium broadening
 $\text{Re}(U) \approx 0; \text{Im}(U) \text{ large}$

experimental task: search for $\left\{ \begin{array}{l} \text{mass shift ?} \\ \text{broadening?} \\ \text{structures?} \end{array} \right\}$ of hadronic spectral functions

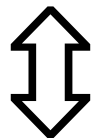
meson-nucleus optical potential

H. Nagahiro and S. Hirenzaki,
PRL 94 (2005) 232503

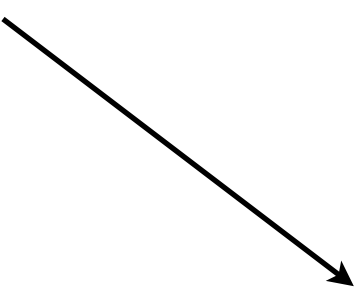
$$U(r) = V(r) + iW(r)$$


$$V(r) = \Delta m(\rho_0) \cdot \frac{\rho(r)}{\rho_0}$$

real part



in-medium mass modification


$$W(r) = -\Gamma_0/2 \cdot \frac{\rho(r)}{\rho_0}$$
$$= -\frac{1}{2} \cdot \hbar c \cdot \rho(r) \cdot \sigma_{inel} \cdot \beta$$

imaginary part



lifetime shortened

in-medium width, absorption
inelastic cross section

mass and lifetime (width) may be changed in the medium

experimental approaches to determine the meson-nucleus optical potential

$$U(r) = V(r) + iW(r)$$

real part

$$V(r) = \Delta m(\rho_0) \cdot \frac{\rho(r)}{\rho_0}$$

- ◆ line shape analysis
- ◆ excitation function
- ◆ momentum distribution
- ◆ meson-nucleus bound states

imaginary part

$$\begin{aligned} W(r) &= -\Gamma_0/2 \cdot \frac{\rho(r)}{\rho_0} \\ &= -\frac{1}{2} \cdot \hbar c \cdot \rho(r) \cdot \sigma_{inel} \cdot \beta \end{aligned}$$

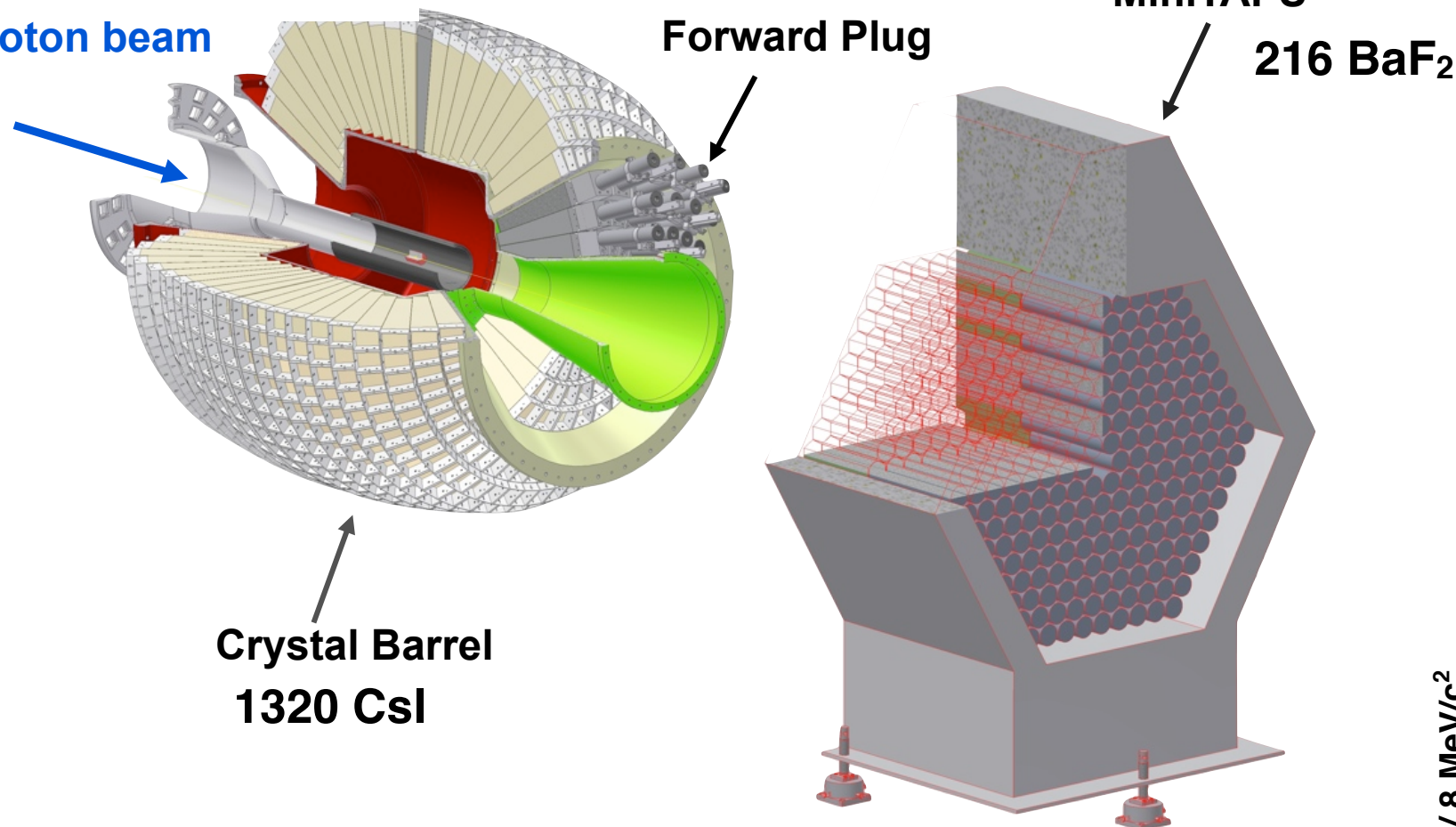
- ◆ transparency ratio measurement

$$T_A = \frac{\sigma_{\gamma A \rightarrow \eta' X}}{A \cdot \sigma_{\gamma N \rightarrow \eta' X}}$$

D. Cabrera et al., NPA 733 (2004)130

CBELSA/TAPS experiment

$E_\gamma = 0.7 - 3.1 \text{ GeV}$
photon beam

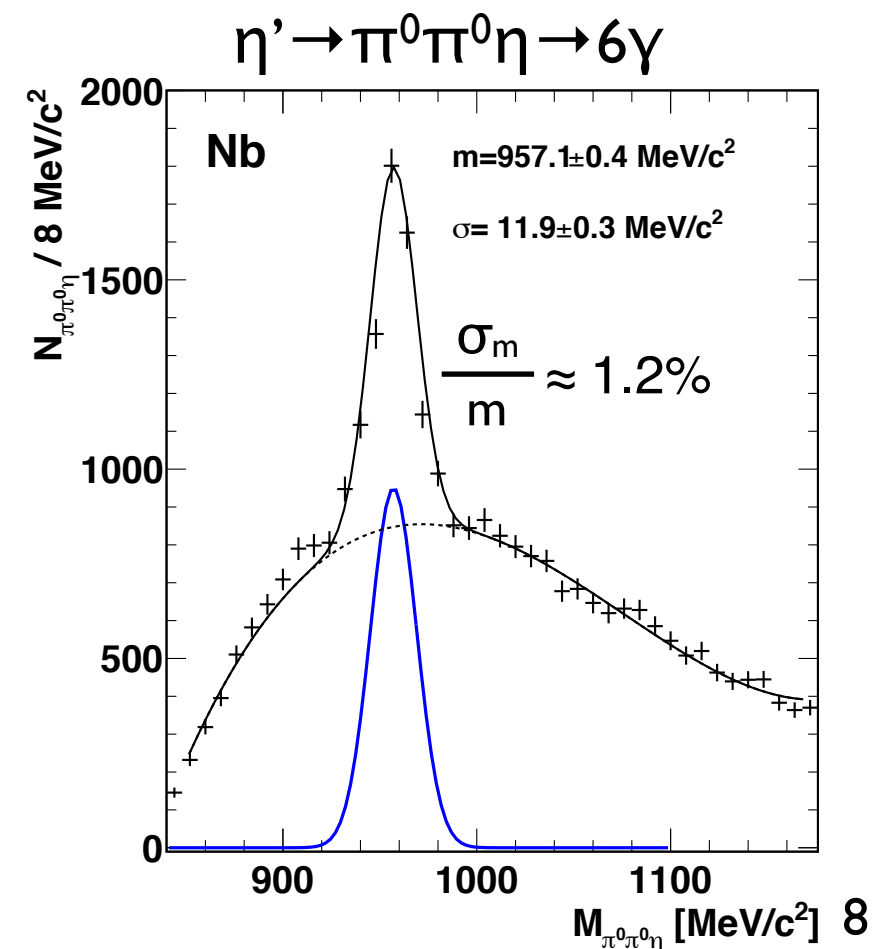
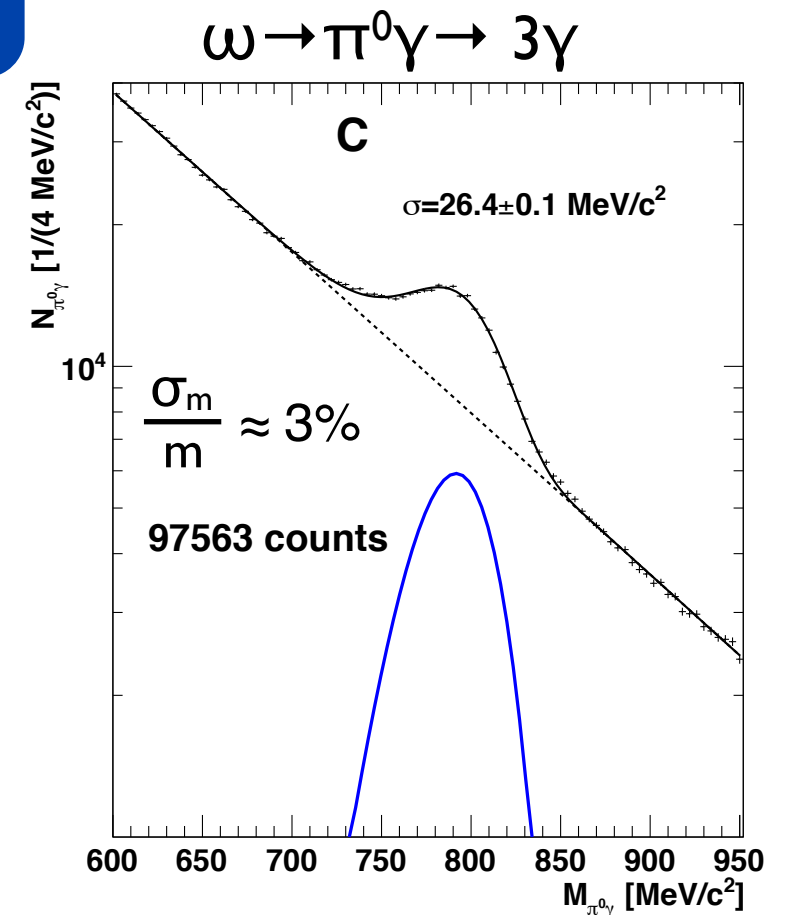


solid target: ^{12}C and ^{93}Nb

4π photon detector: ideally suited for
identification of multi-photon final states

$\omega \rightarrow \pi^0 \gamma \rightarrow 3\gamma$ BR 8.2%

$\eta' \rightarrow \pi^0 \pi^0 \eta \rightarrow 6\gamma$ BR 8.5%



The real part of the meson-nucleus
optical potential

the real part of the ω -nucleus potential

J.Weil, U. Mosel and V. Metag, PLB 723 (2013) 120 $\omega \rightarrow \pi^0 \gamma$

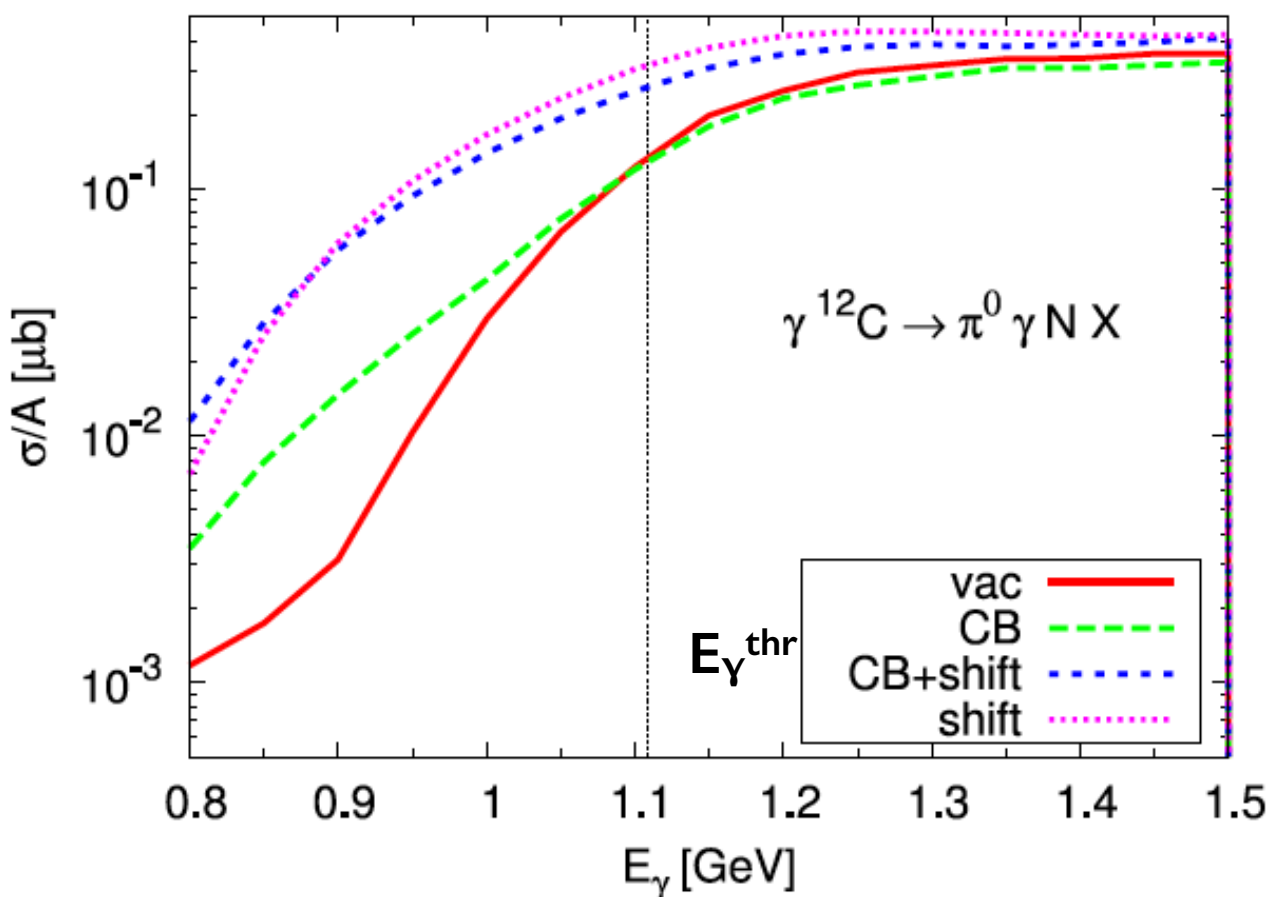
sensitive to nuclear density at **production** point and not at **decay** point

◆ measurement of the excitation function of the meson

in case of dropping mass -
higher meson yield for given $\sqrt{s}$
because of increased phase space
due to lowering of the production threshold

⇒ cross section enhancement

$\pi^0 \gamma$ excitation function

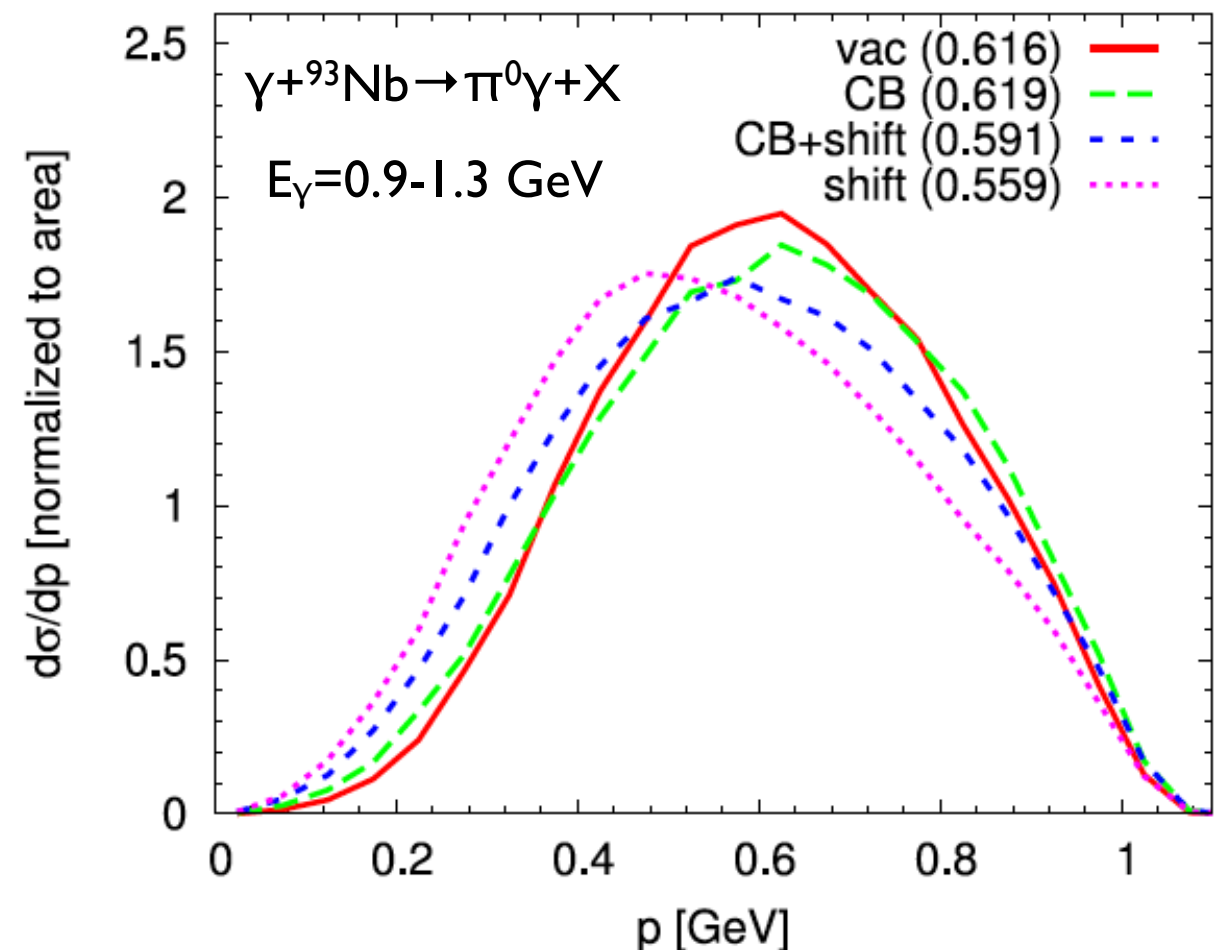


◆ momentum distribution of the meson:

in case of dropping mass - when leaving the nucleus hadron has to become on-shell;
mass generated at the expense of kinetic energy

⇒ downward shift of momentum distribution

$\pi^0 \gamma$ momentum distribution

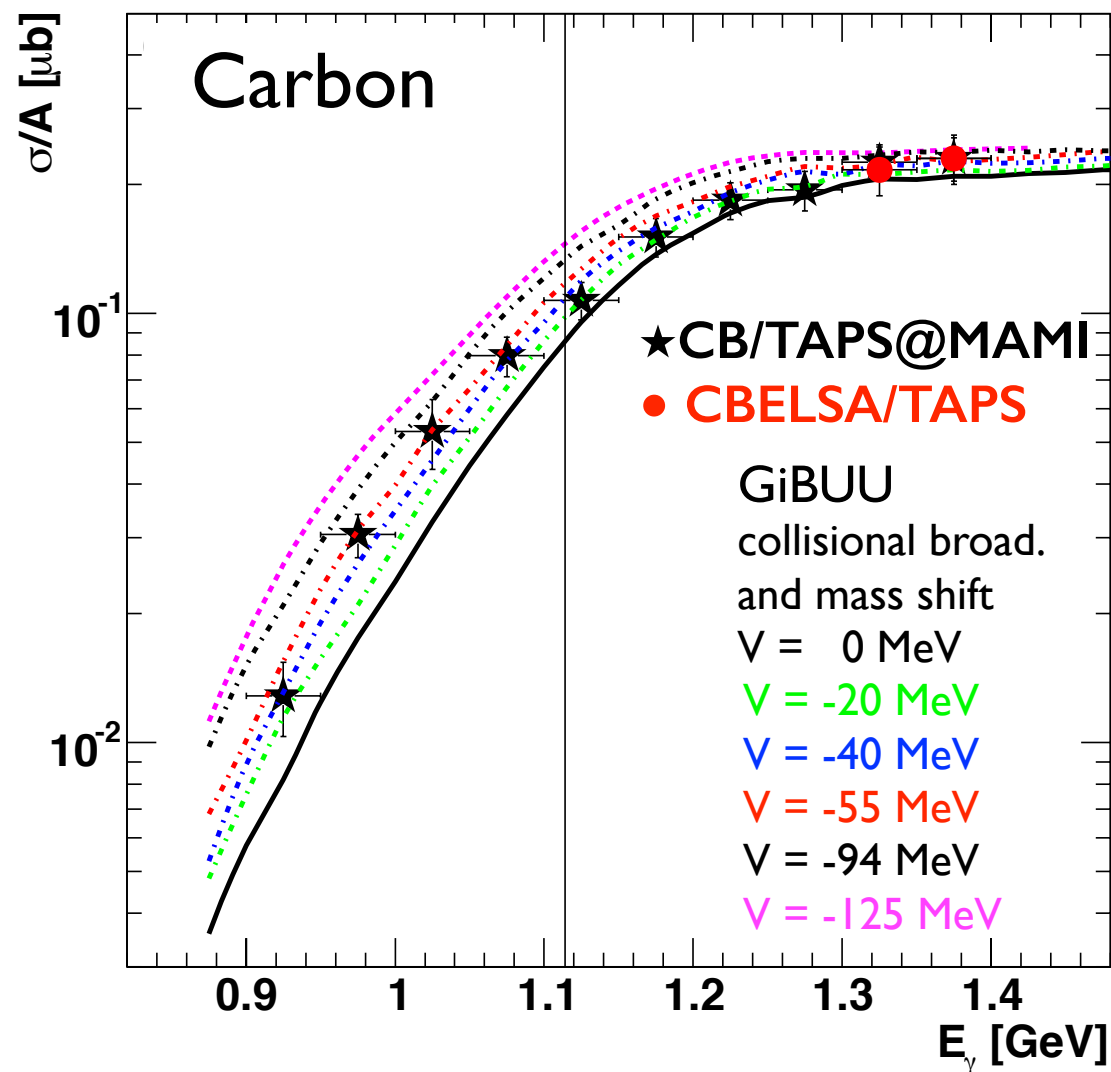


excitation function for ω photoproduction off C comparison with GiBUU calculation

CB/TAPS @ MAMI

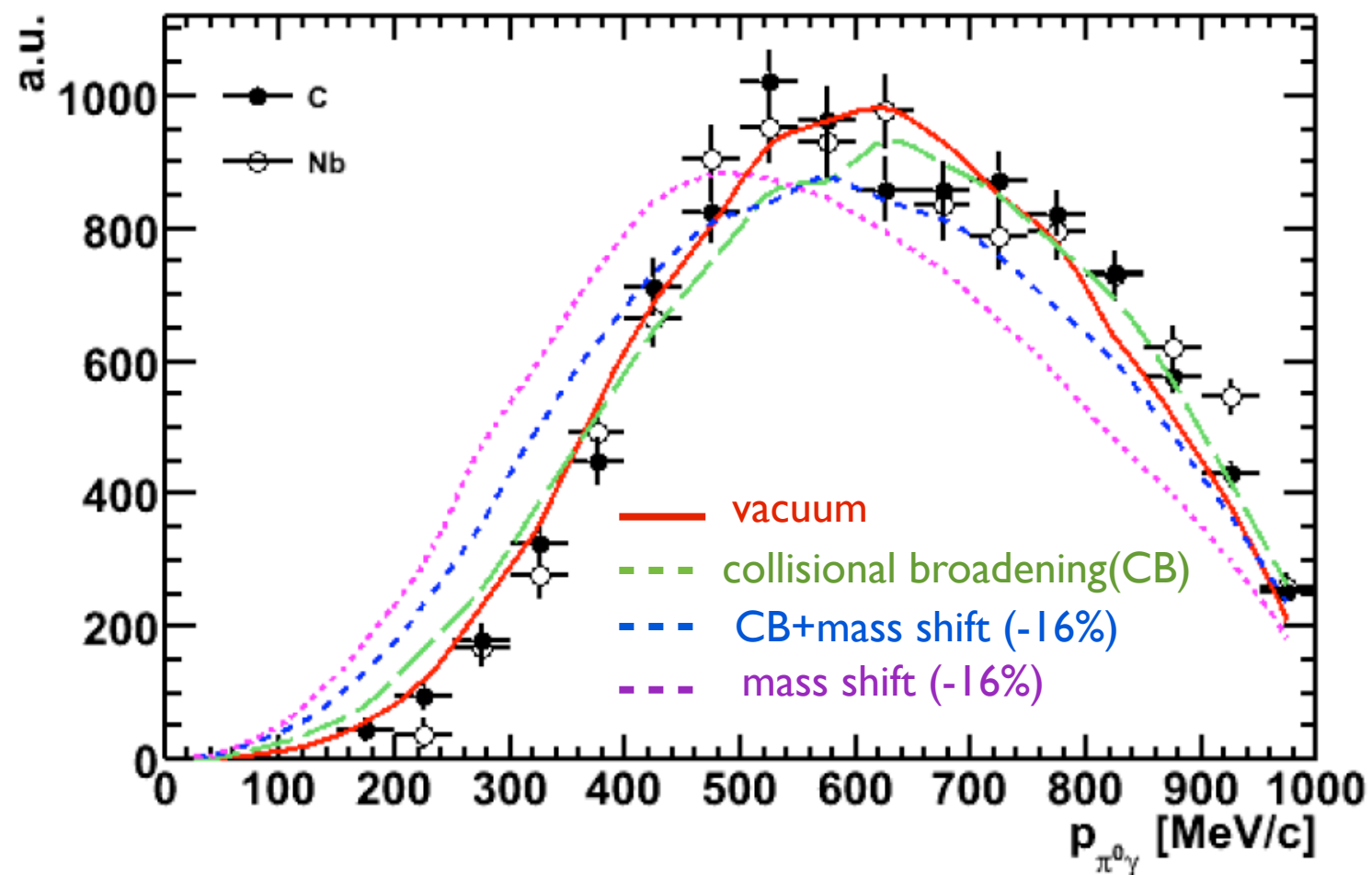
V. Metag et al., PPNP, 67 (2012) 530

excitation function



M.Thiel et al., EPJA 49 (2013) 132

momentum distribution



$$V(\rho=\rho_0) = -(42 \pm 17(\text{stat}) \pm 20(\text{syst})) \text{ MeV}$$

data not consistent with strong mass shift scenario ($\Delta m/m \approx -16\%$)

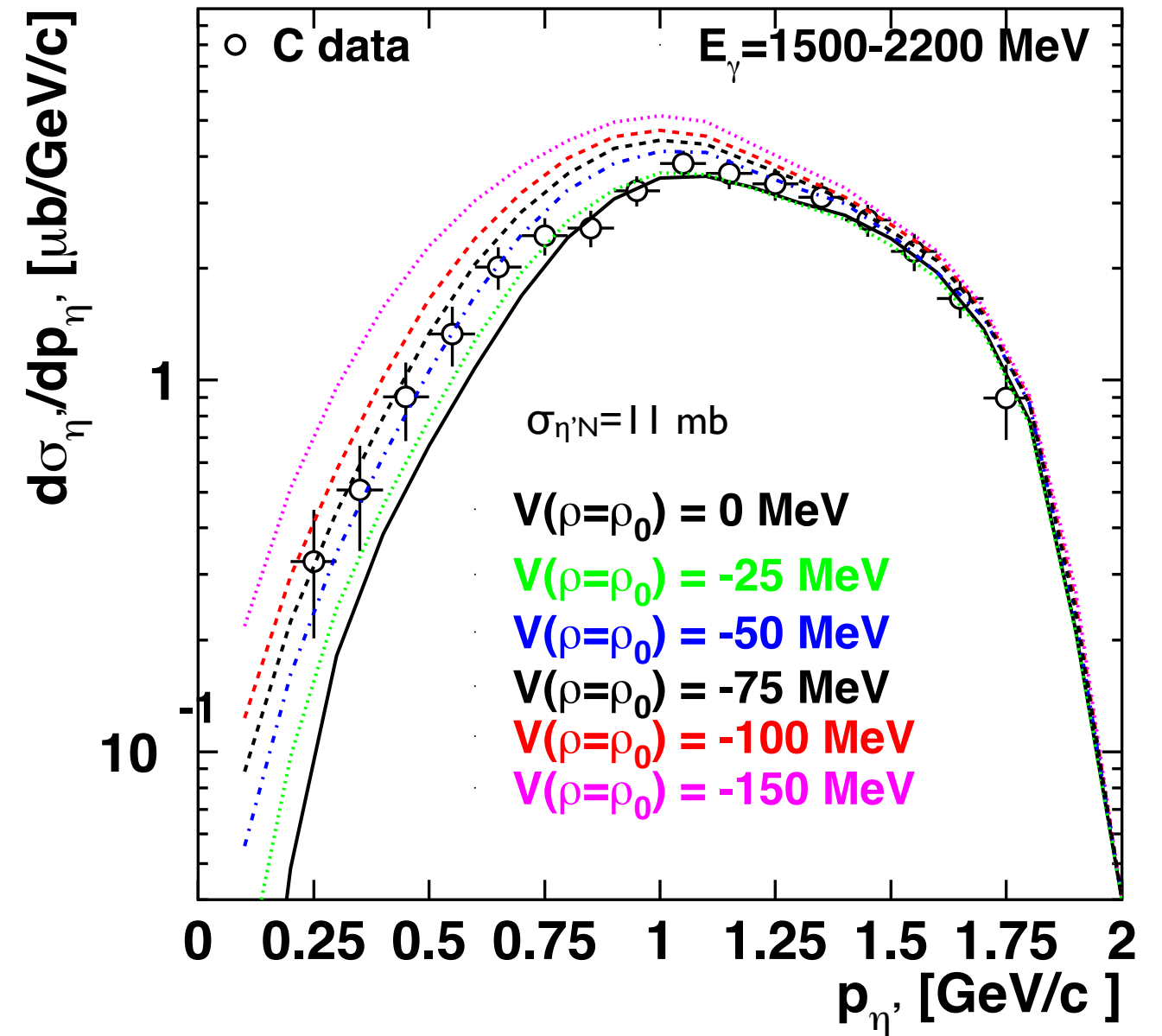
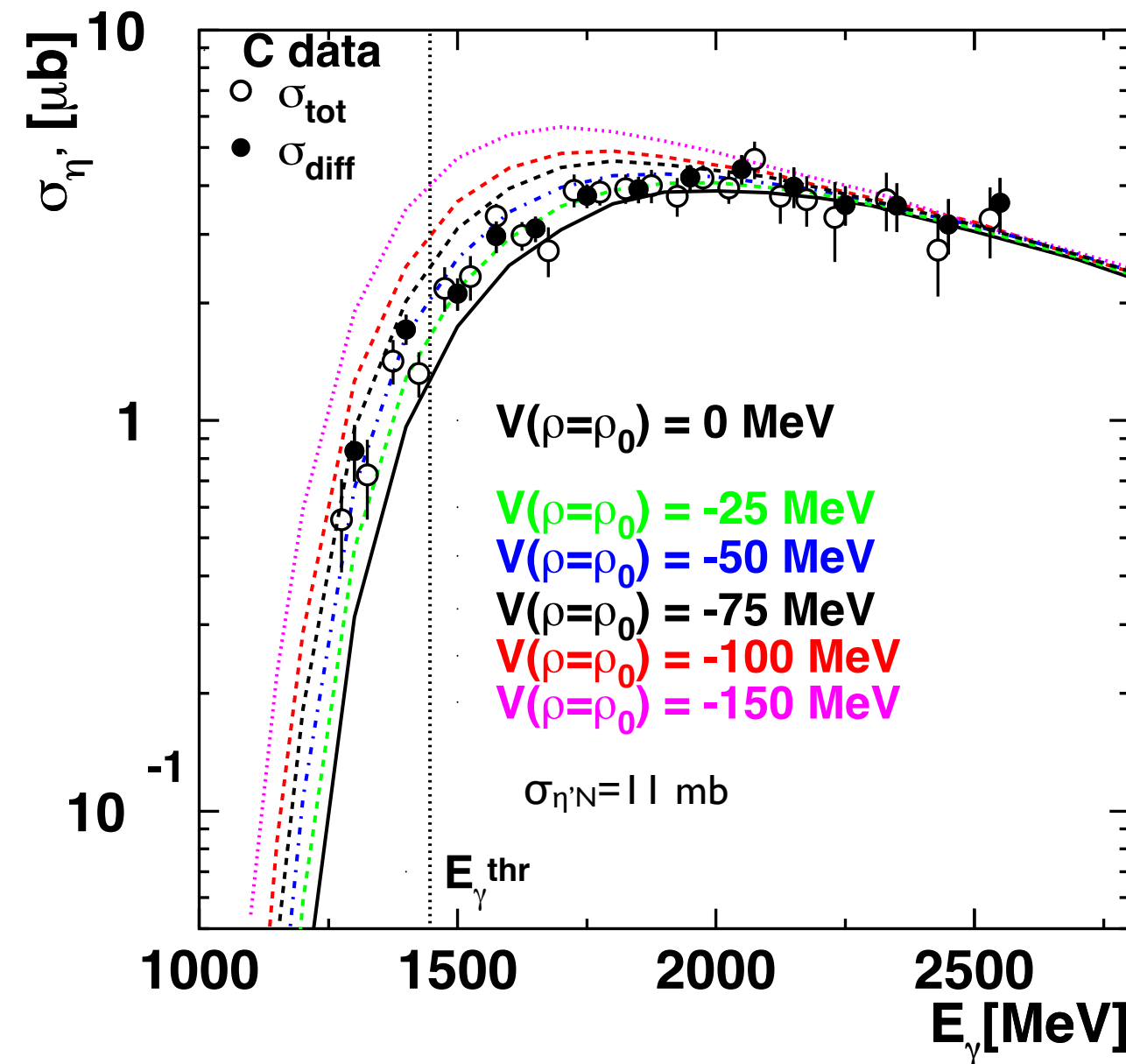
excitation function and momentum distribution for η' photoproduction off C

CBELSA/TAPS @ ELSA

data: M. Nanova et al., PLB 727 (2013) 417

$\gamma C \rightarrow \eta' X$

calc.: E. Paryev, J. Phys. G 40 (2013) 025201



$$V_{\eta'}(\rho=\rho_0) = -(40 \pm 6) \text{ MeV}$$

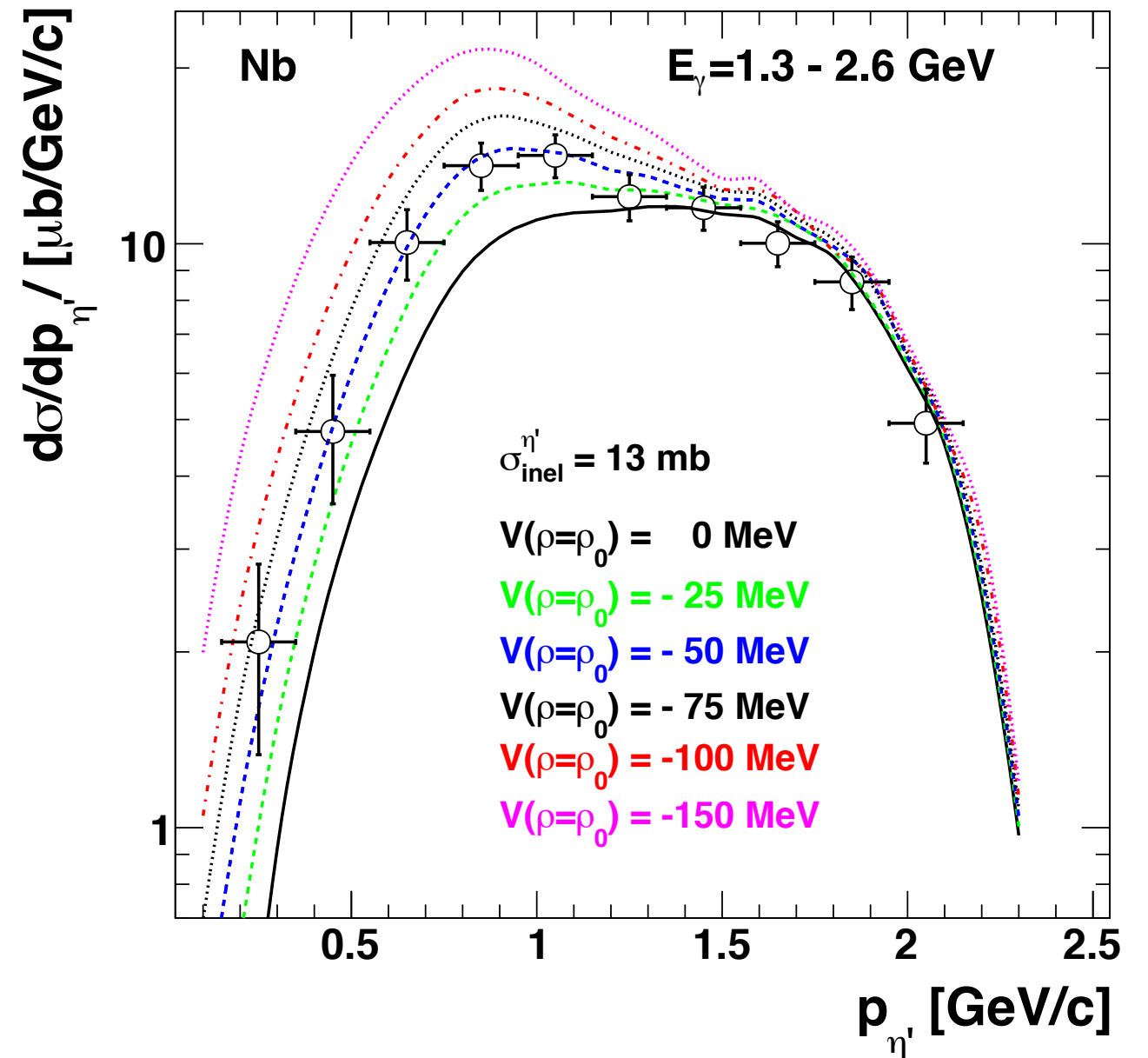
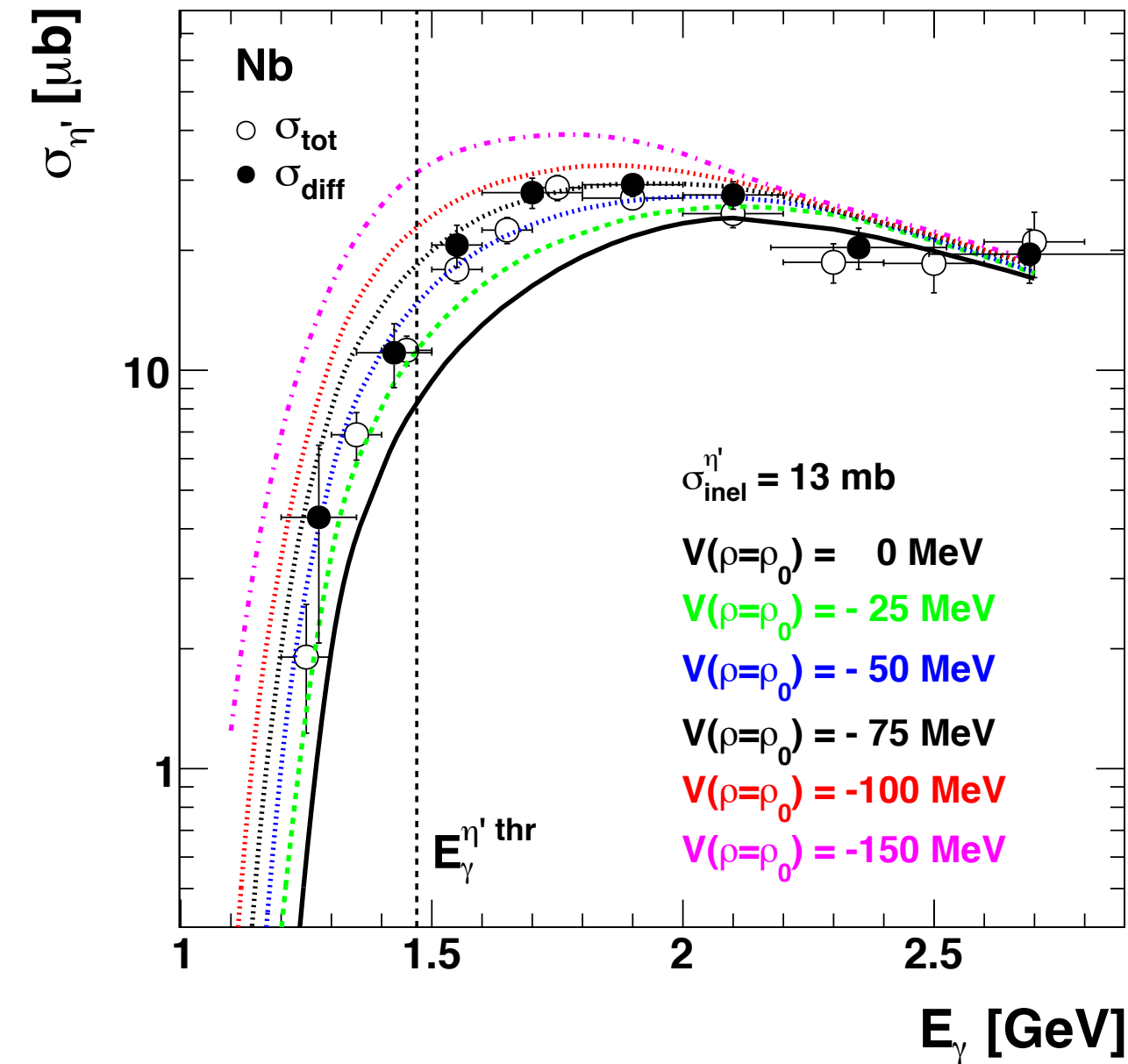
$$V_{\eta'}(\langle p_{\eta'} \rangle \approx 1.1 \text{ GeV}/c; \rho=\rho_0) = -(32 \pm 11) \text{ MeV}$$

data disfavour strong mass shifts

excitation function and momentum distribution for η' photoproduction off Nb

CBELSA/TAPS @ ELSA

$\gamma \text{ Nb} \rightarrow \eta' X$ M. Nanova et al., arXiv:1607.07228; accepted for publication in PRC



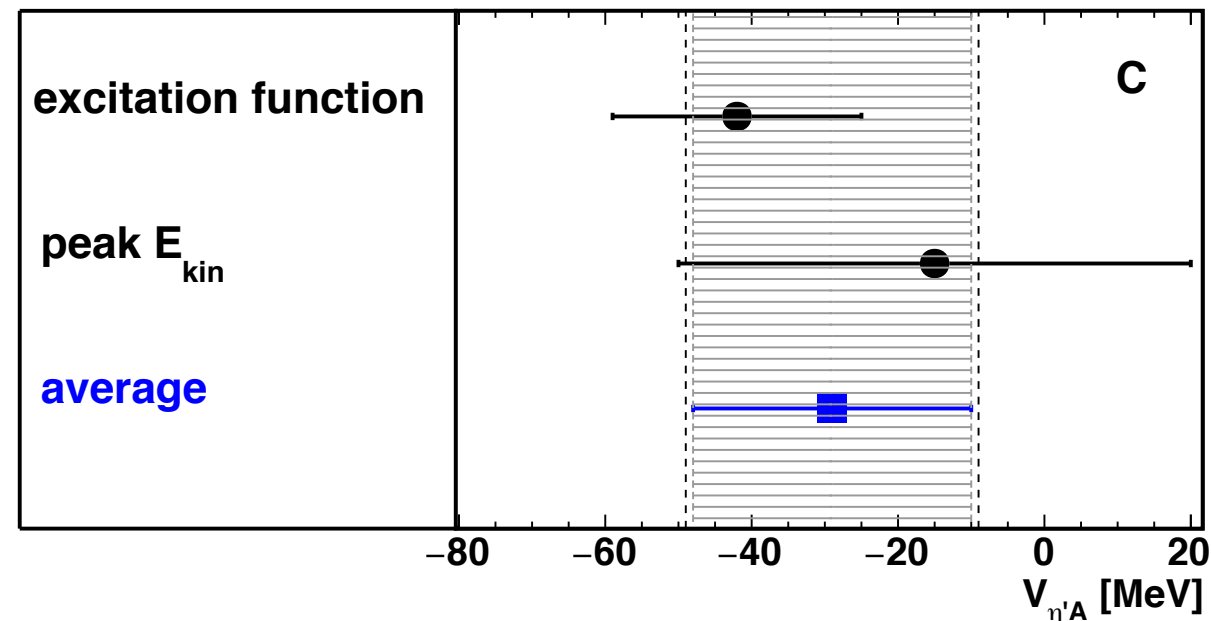
$$V_{\eta'}(\rho=\rho_0) = -(40 \pm 12) \text{ MeV}$$

$$V_{\eta'}(\langle p_{\eta'} \rangle \approx 1.14 \text{ GeV}/c; \rho=\rho_0) = -(41 \pm 10) \text{ MeV}$$

data disfavour strong mass shifts

compilation of results for the real part of the ω - and η' -nucleus optical potential

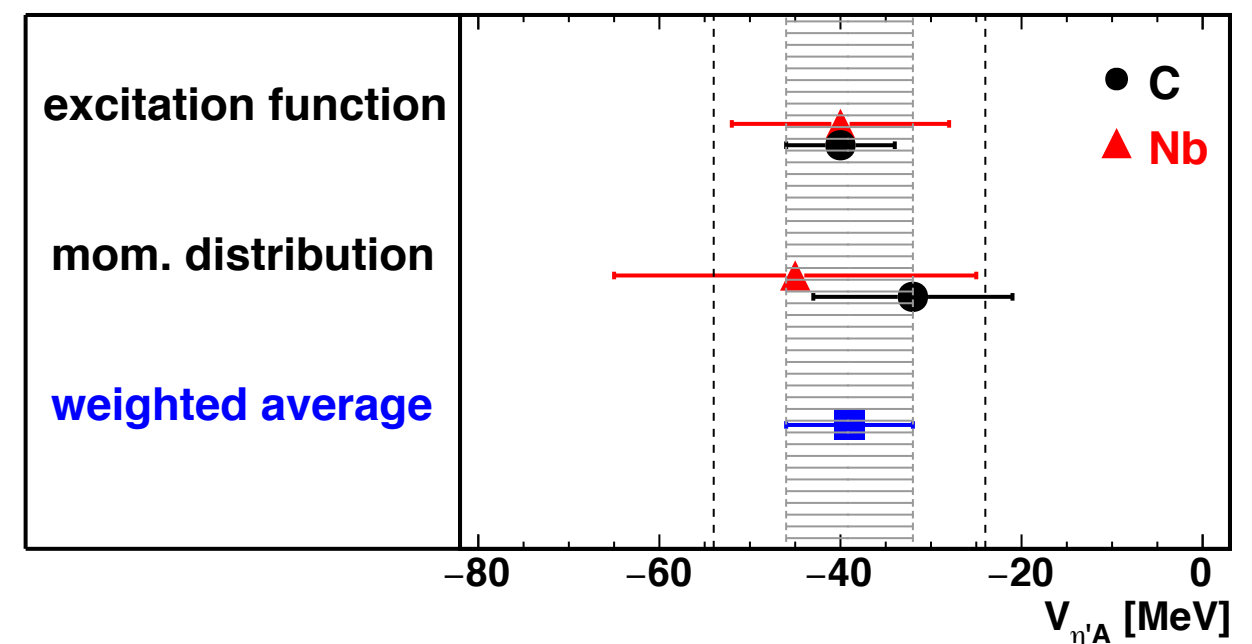
ω



$$V_{\omega A}(\rho=\rho_0) = -(29 \pm 19(\text{stat}) \pm 20(\text{syst})) \text{ MeV}$$

M. Nanova et al., arXiv:1607.07228;
 accepted for publication in PRC

η'



$$V_{\eta' A}(\rho=\rho_0) = -(40 \pm 8(\text{stat}) \pm 15(\text{syst})) \text{ MeV}$$

first (indirect) observation of in-medium mass shift of η' at $\rho=\rho_0$ and $T=0$
 in good agreement with QMC model predictions (S. Bass et al., PLB 634 (2006) 368)

The imaginary part of the meson-nucleus
optical potential: momentum dependence

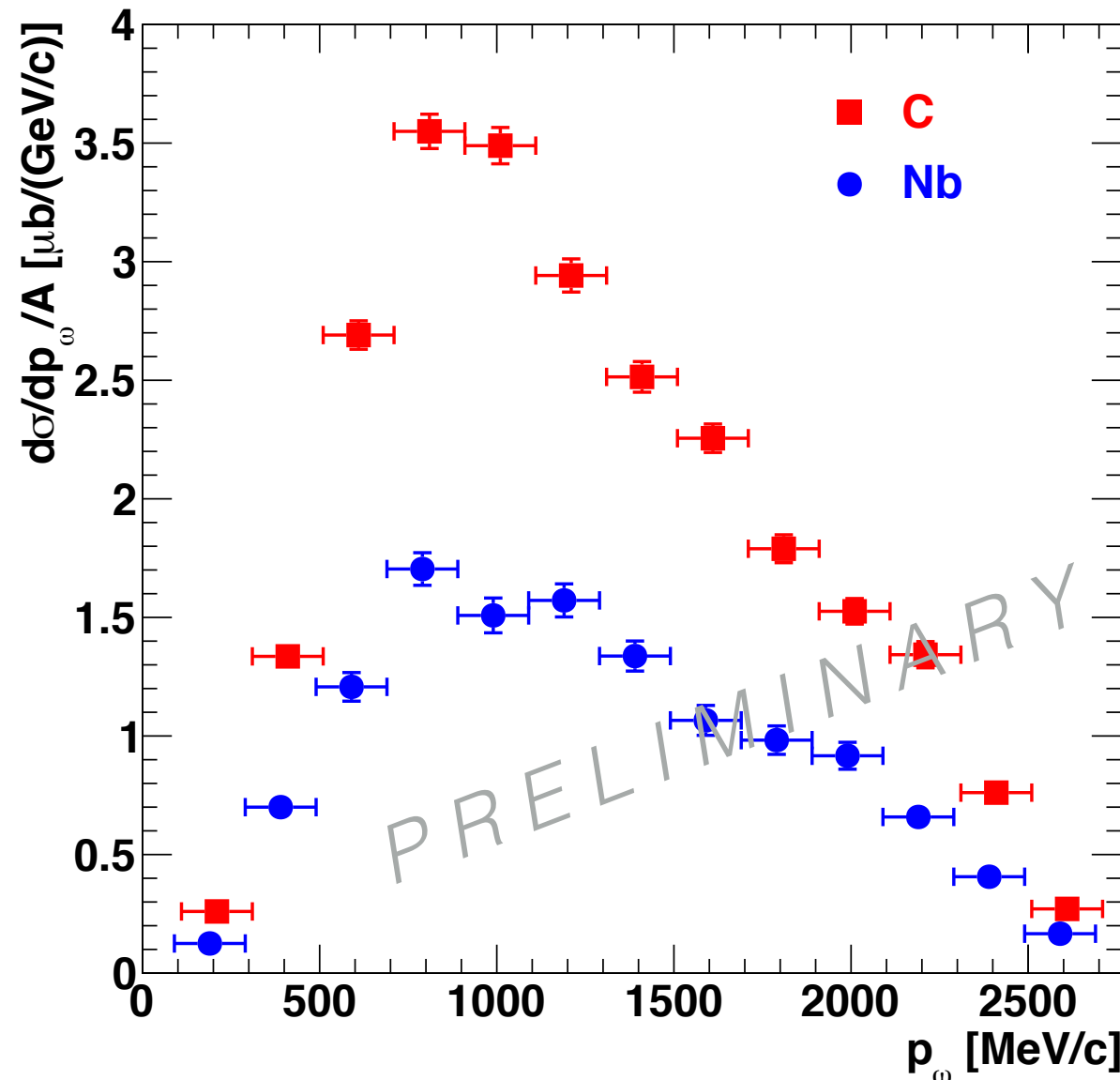
momentum differential cross section for ω, η' produced off C, Nb

ω

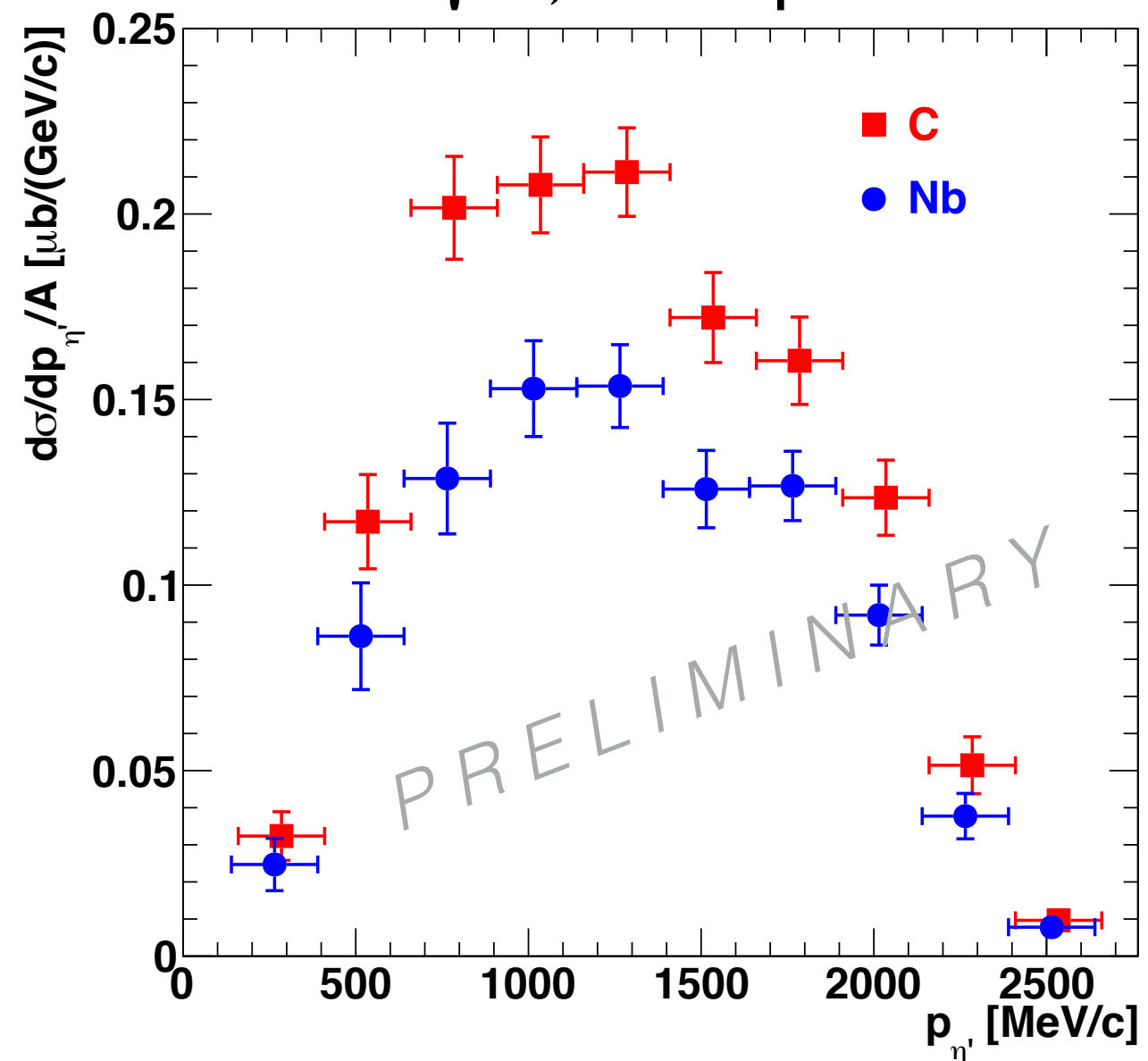
$E_\gamma = 1.2 - 2.9 \text{ GeV}$

η'

$\gamma \text{ C, Nb} \rightarrow \omega X$



$\gamma \text{ C, Nb} \rightarrow \eta' X$



momentum differential cross sections $\Rightarrow T_{\text{Nb/C}}^m(p_m) = \frac{12 \cdot \sigma_{\gamma \text{Nb} \rightarrow mX}(p_m)}{93 \cdot \sigma_{\gamma \text{C} \rightarrow mX}(p_m)}$

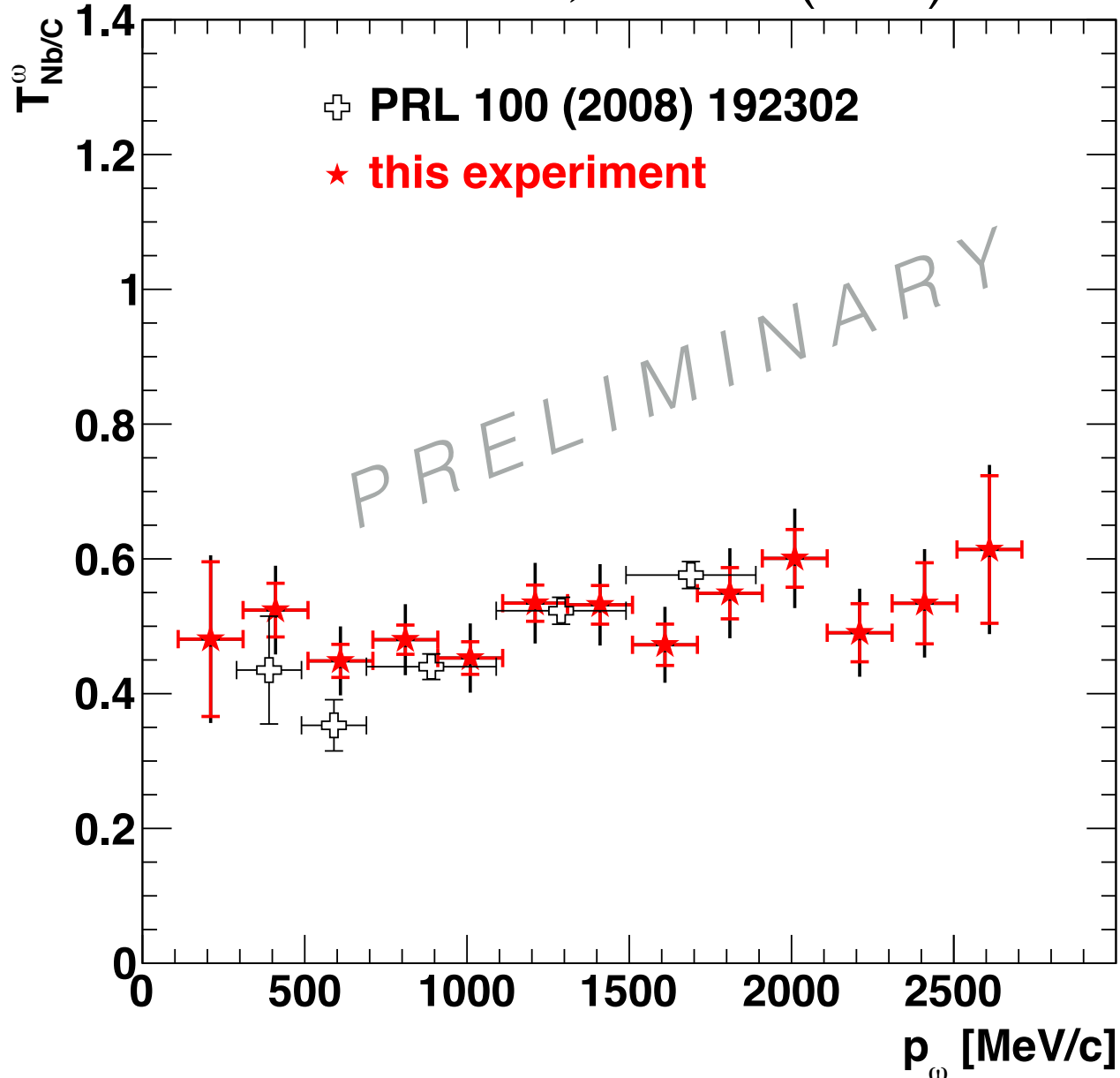
momentum dependence of transparency ratio for ω , η'

$$T_{\text{Nb/C}}^m(p_m) = \frac{12 \cdot \sigma_{\gamma\text{Nb} \rightarrow mX}(p_m)}{93 \cdot \sigma_{\gamma\text{C} \rightarrow mX}(p_m)}$$

ω

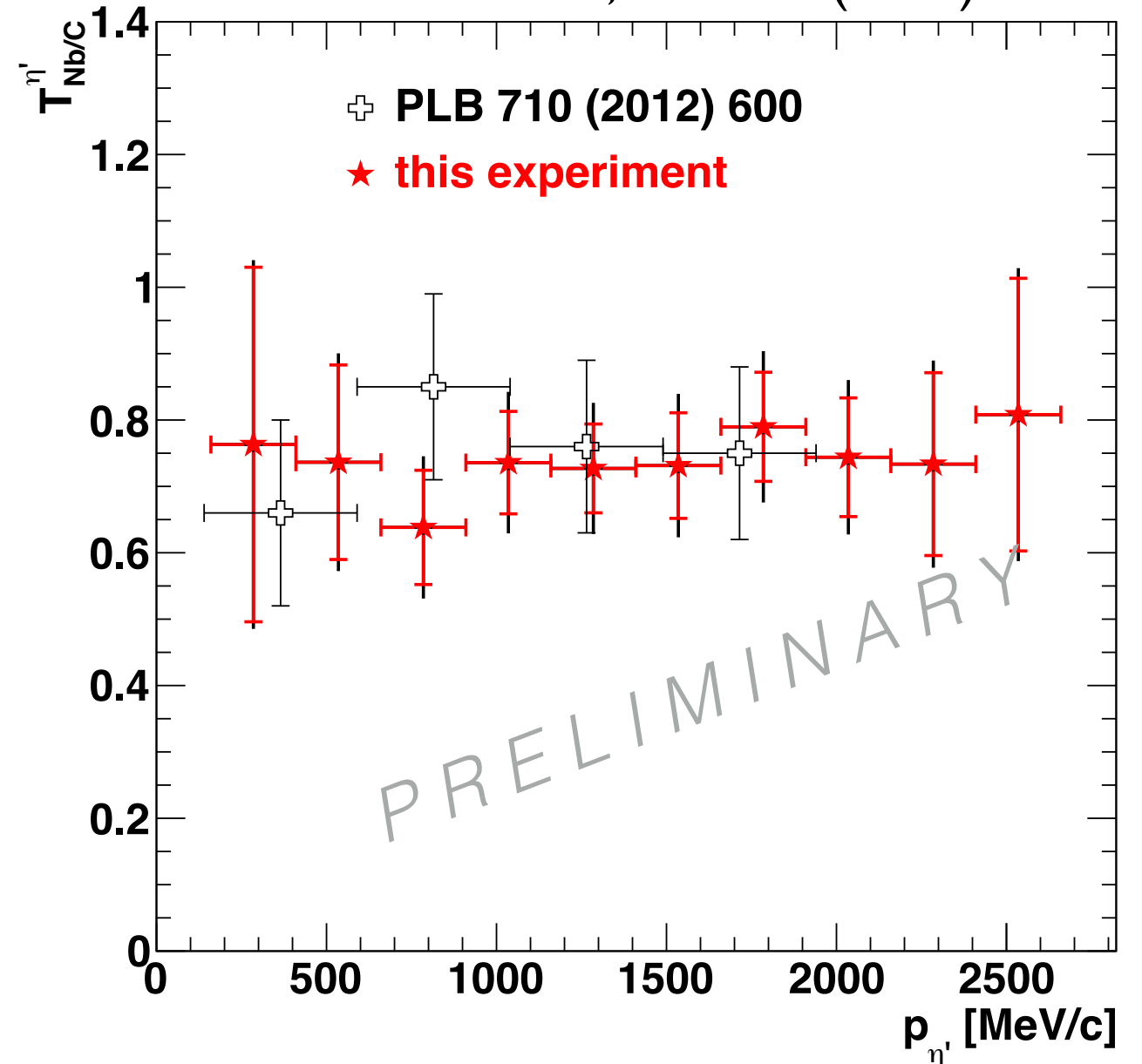
η'

+ M. Kotulla et al., PRL 100 (2008) 192302



$$T_{\text{Nb/C}}^\omega \approx 0.4-0.6$$

+ M. Nanova et al., PLB 710 (2012) 600



$$T_{\text{Nb/C}}^{\eta'} \approx 0.7-0.8$$

absorption of η' mesons much weaker than for ω mesons !!

momentum dependence of the ω , η' in-medium width

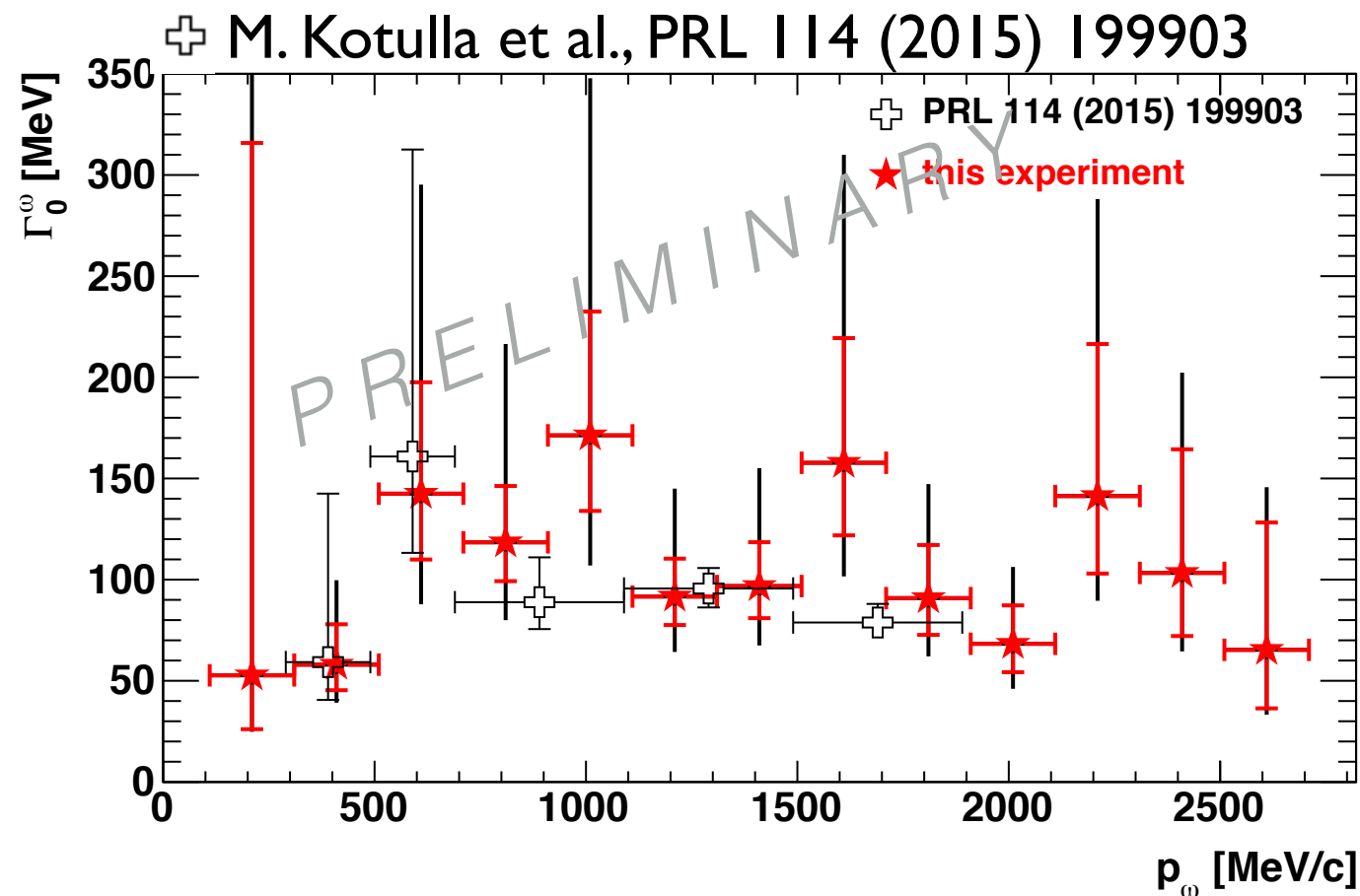
Glauber model: high energy Eikonal approximation

$$T_{\text{Nb/C}}^m(p_m) \Rightarrow \Gamma_0^m(\rho=\rho_0)(p_m)$$

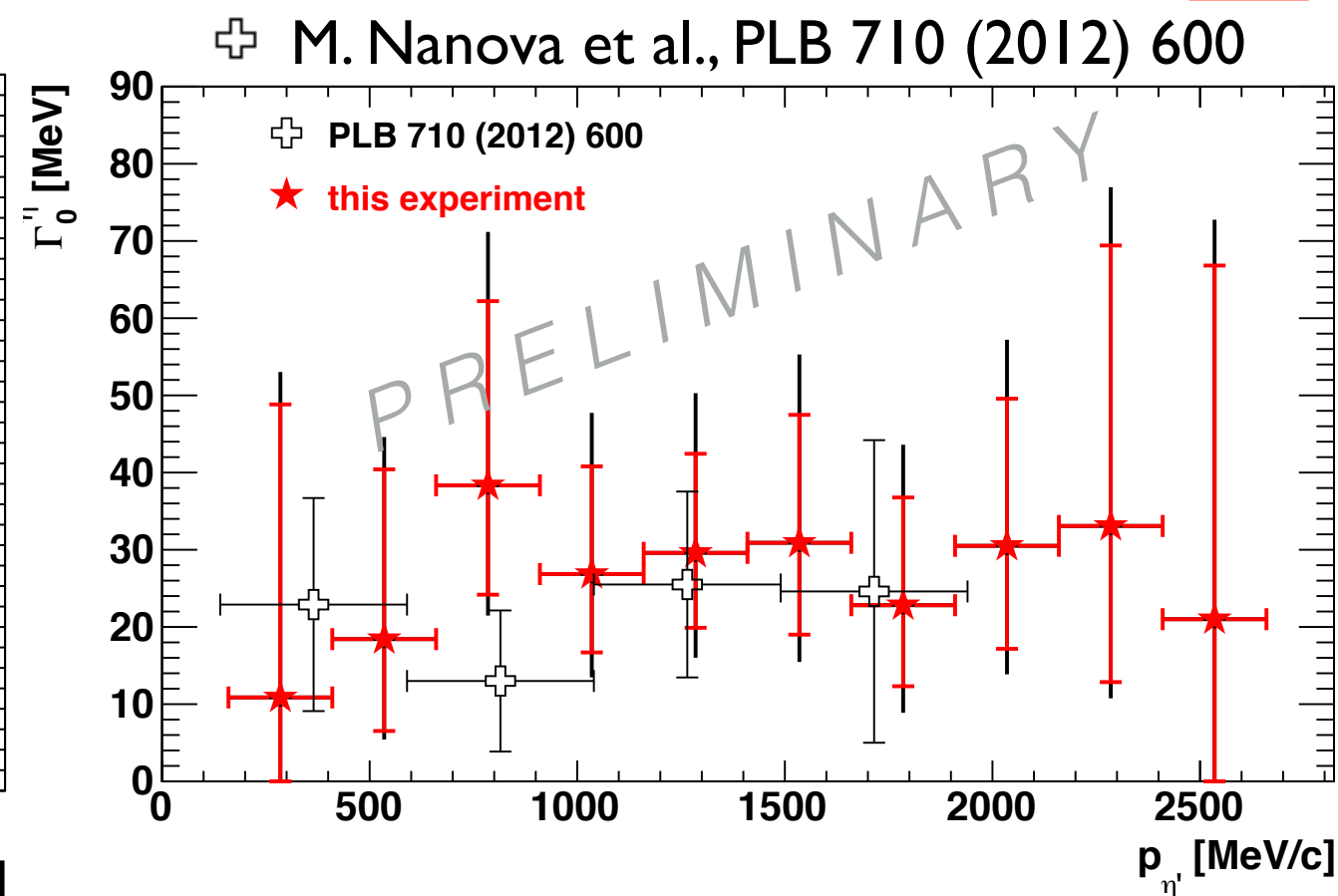
S. Friedrich et al., submitted to EPJA

ω

η'



ω in-medium width ≈ 50 -150 MeV



η' in-medium width ≈ 10 -30 MeV

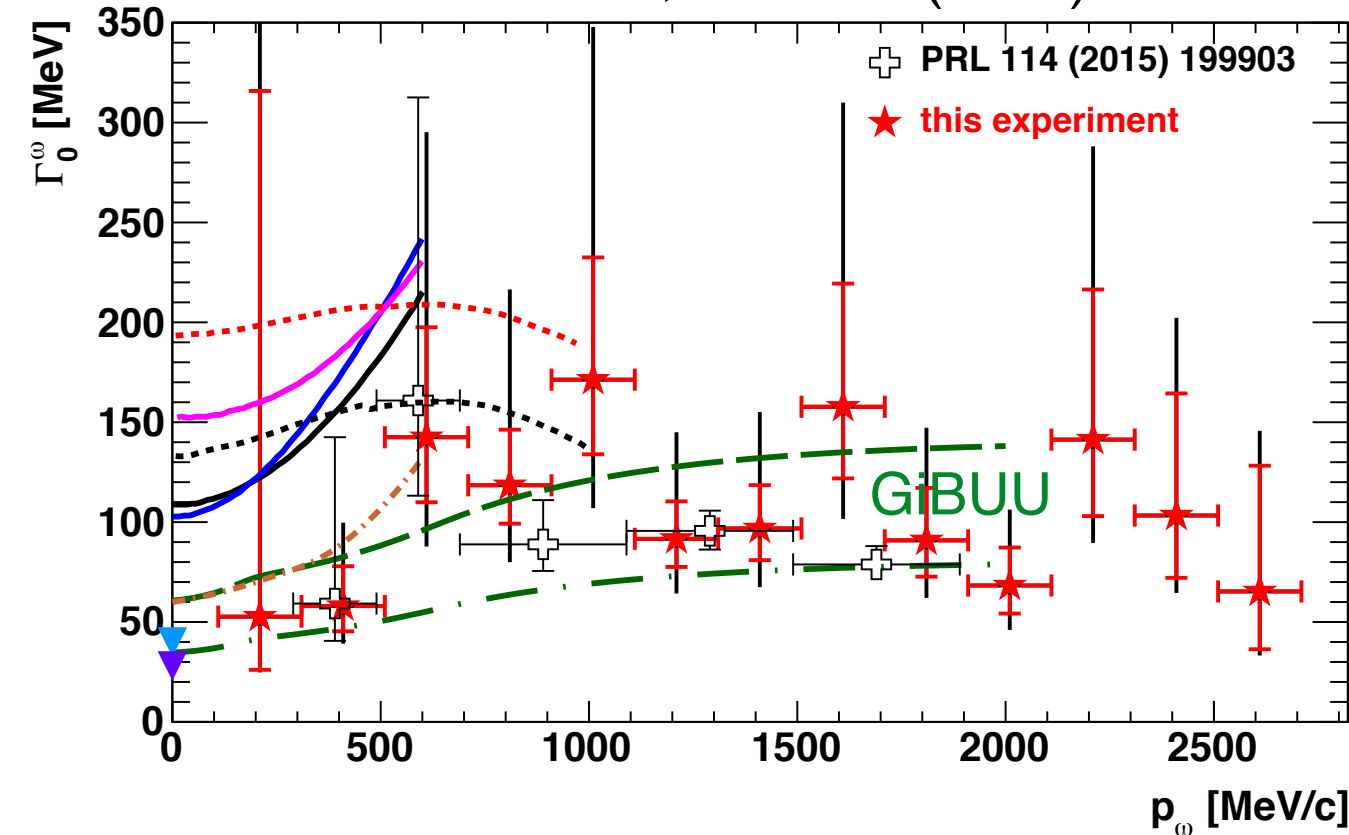
ω in-medium width in comparison to theoretical predictions

ω

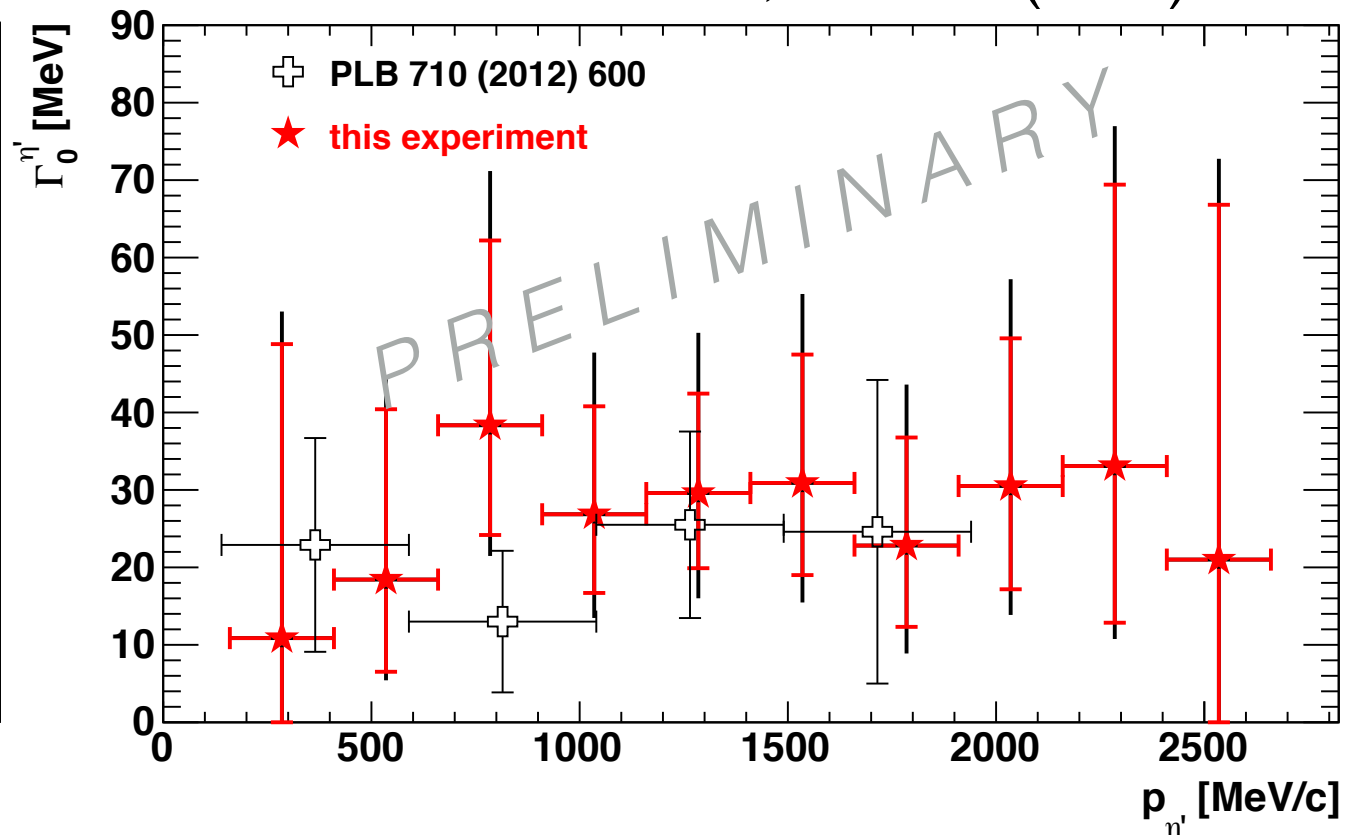
S. Friedrich et al., submitted to EPJA

η'

+ M. Kotulla et al., PRL 114 (2015) 199903



+ M. Nanova et al., PLB 710 (2012) 600



two groups of calculations:

$\omega \rightarrow \rho\pi$: --- D. Cabrera and R. Rapp, PLB 729 (2014) 67
 — A. Ramos et al., EPJA 49(2013) 148

NN* ω coupl.:
 ▼ F. Klingl et al., NPA 650 (1999) 299
 ▼ B. Friman et al., arXiv: 9811040
 ▼ M. Lutz et al., NPA 706 (2002) 437
 P. Mühlich et al., NPA 780 (2006) 187

imaginary part of the potential for ω , η'

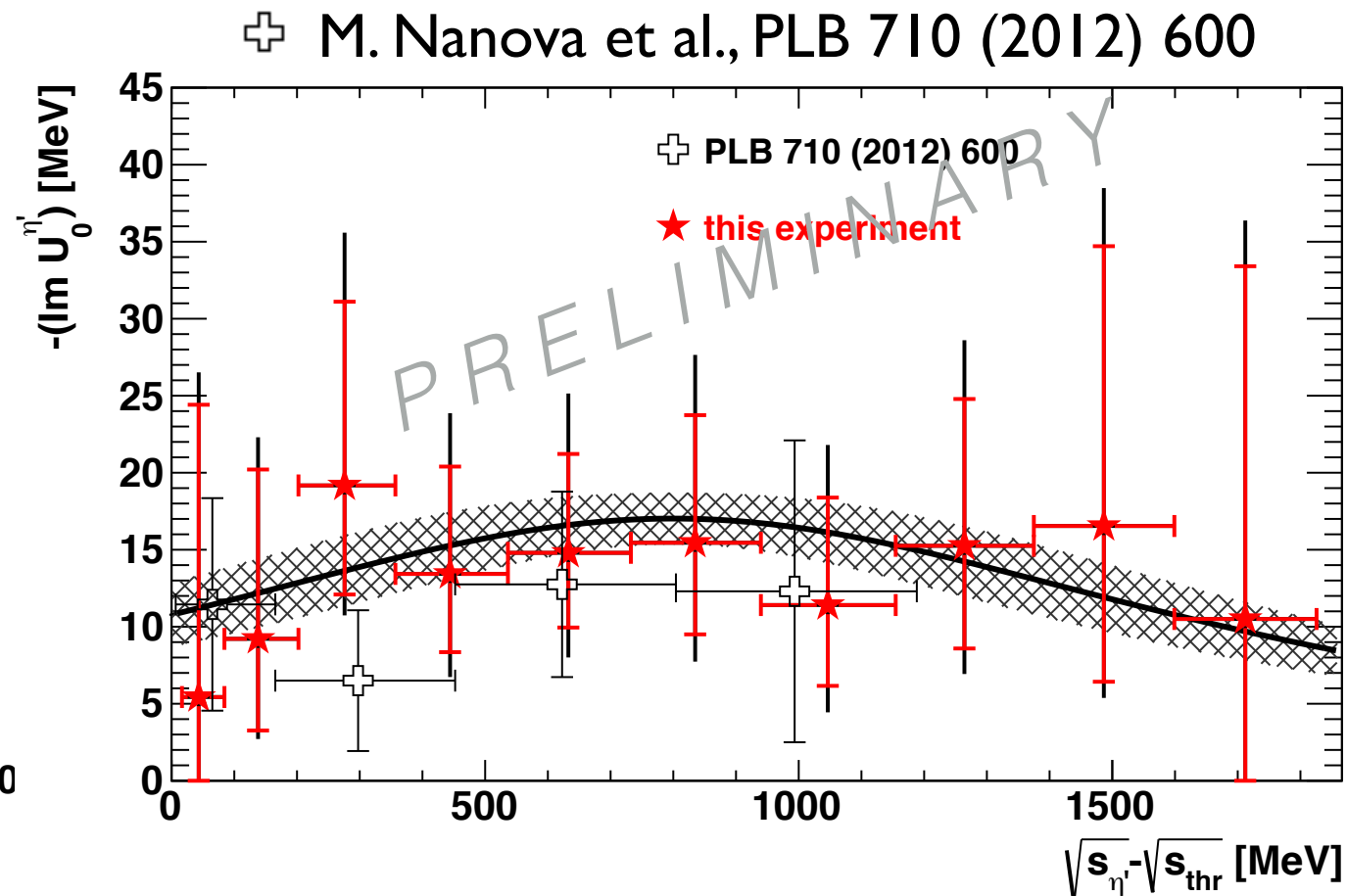
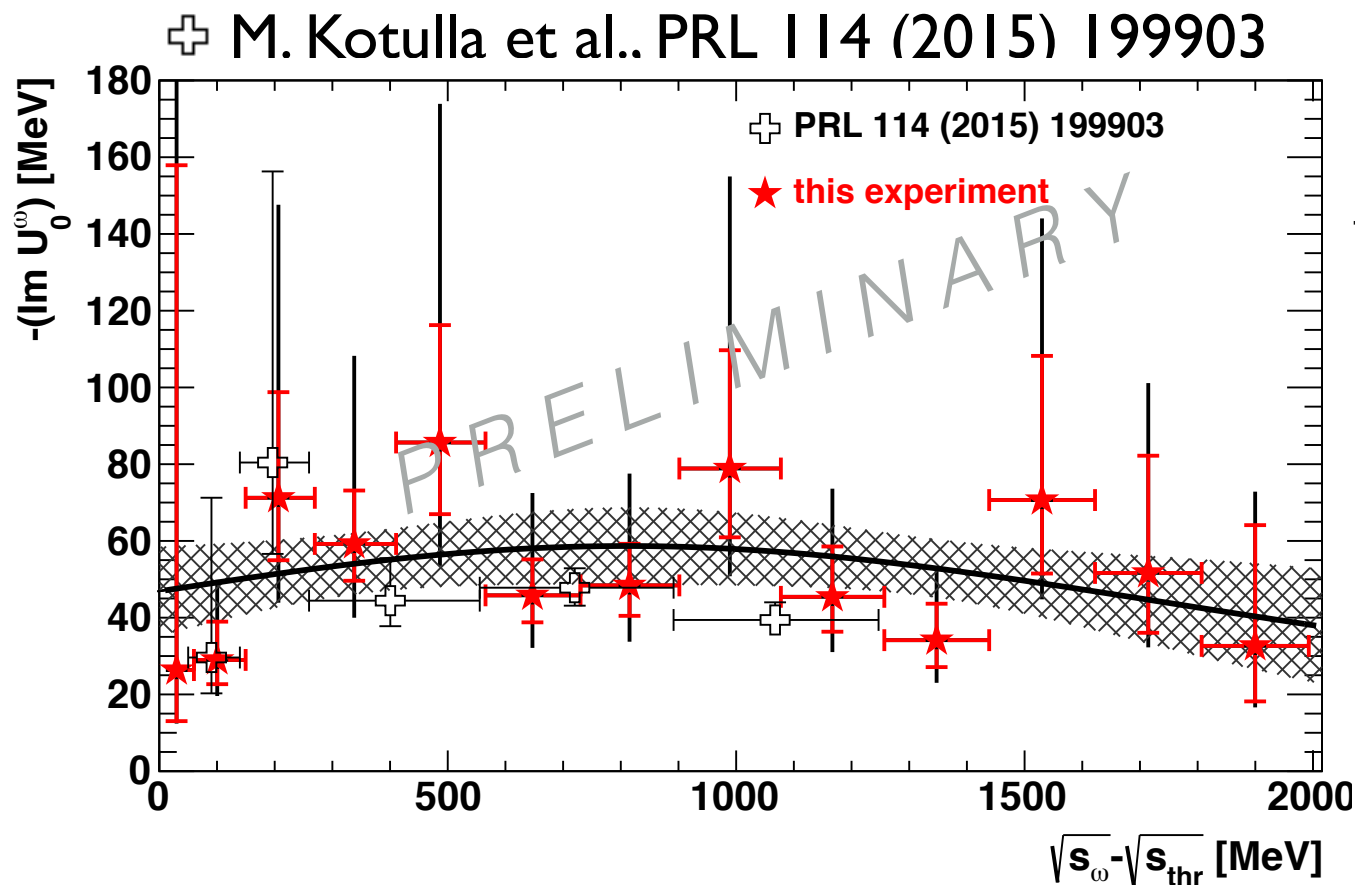
Glauber model: high energy Eikonal approximation

$$T_{\text{Nb/C}}^m(p_m) \Rightarrow \Gamma_0^m(\rho=\rho_0)(p_m) = -2 \text{Im } U_0^m(p_m)$$

S. Friedrich et al., submitted to EPJA

ω

η'



- ◆ extrapolation to production threshold:

$$\text{Im } U_0^\omega(\rho=\rho_0, p_\omega=0) = -(30 \pm 10) \text{ MeV}$$

$$\text{Im } U_0^{\eta'}(\rho=\rho_0, p_{\eta'}=0) = -(10 \pm 3) \text{ MeV}$$

- ◆ extension to higher energies allows for dispersion relation analysis, providing link between real and imaginary part of potential

compilation of results for real and imaginary part of the ω, η' -nucleus optical potential

ω

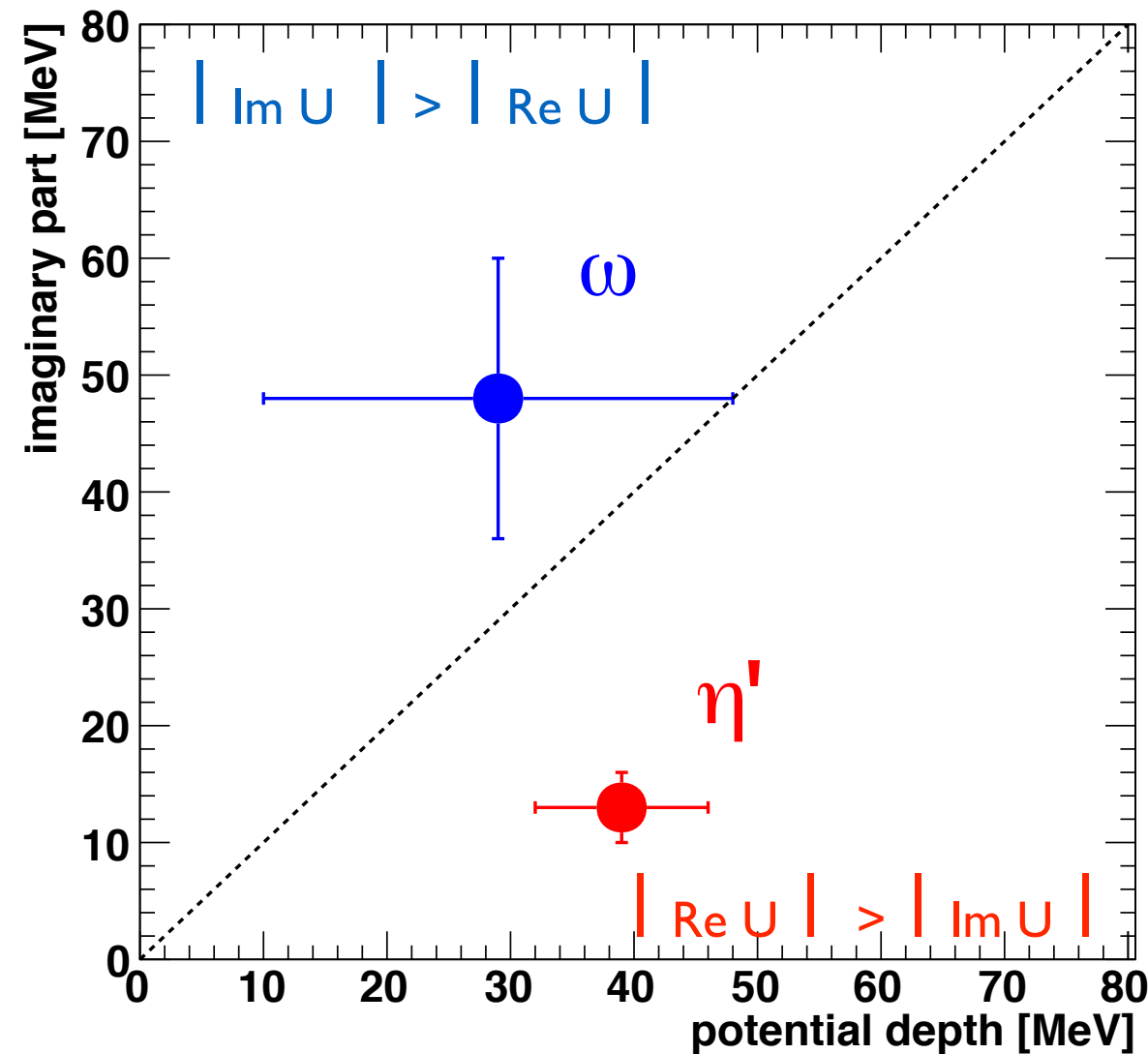
$$U_{\omega A}(\rho=\rho_0)=$$

$$-((29 \pm 19(\text{stat}) \pm 20(\text{syst}) + i(48 \pm 10 \pm 9)) \text{ MeV}$$

η'

$$U_{\eta' A}(\rho=\rho_0)=$$

$$-((40 \pm 8(\text{stat}) \pm 15(\text{syst}) + i(13 \pm 3 \pm 3)) \text{ MeV}$$



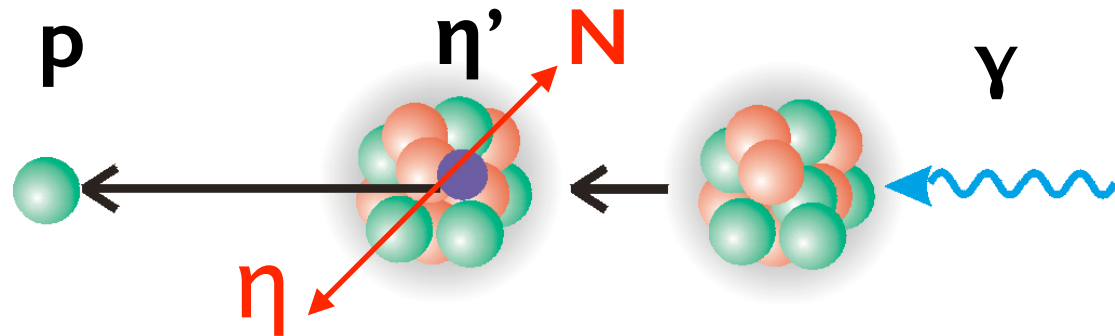
$|Im U| \gtrsim |Re U| ; \Rightarrow \omega$ not a good candidate
 to search for meson-nucleus bound states!

$|Re U| \gg |Im U| ; \Rightarrow \eta'$ promising
candidate to search for mesic states

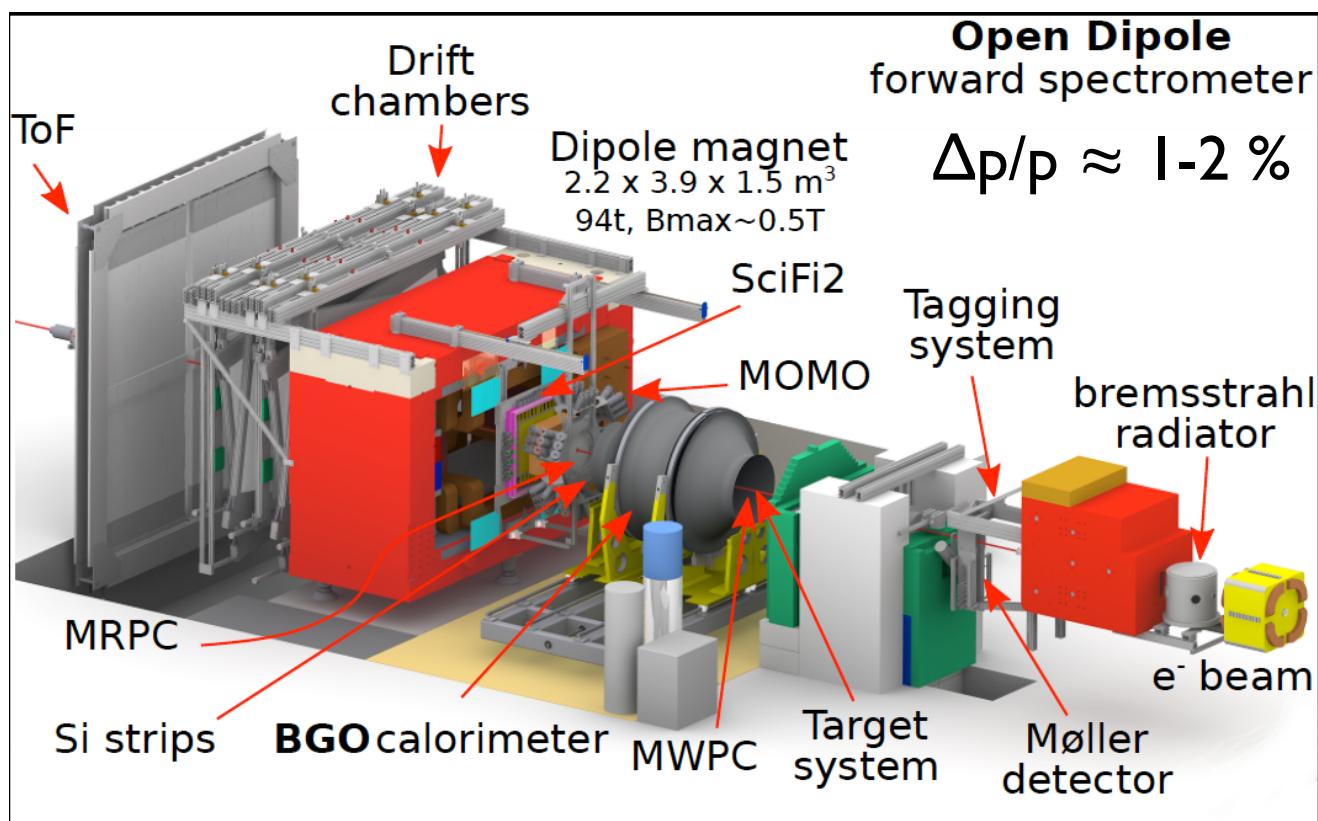
outlook: search for η' -mesic states in photo-nuclear reactions

B1: BGO-OD@ELSA

$^{12}\text{C}(\gamma, p) \eta' X @ 1.5\text{-}2.8 \text{ GeV}$



formation and decay of η' -mesic state

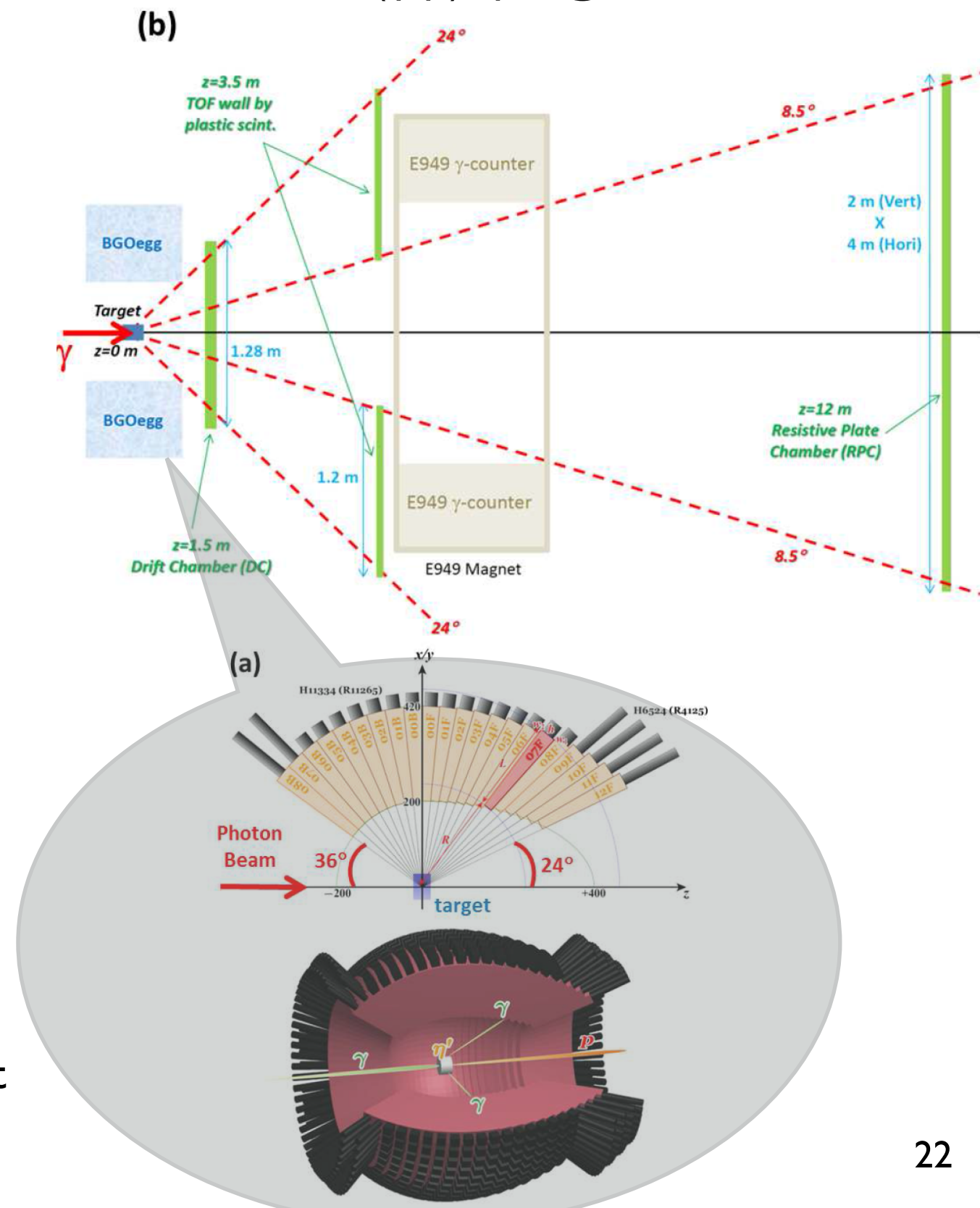


BGO-OD ideally suited for exclusive measurement

approved proposal: ELSA/3-2012-BGO

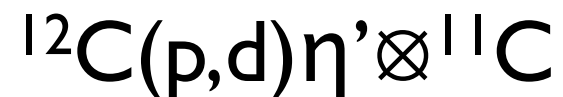
LEPS2@SPRing-8

$^{12}\text{C}(\gamma, p) \eta' X @ 1.5\text{-}2.4 \text{ GeV}$



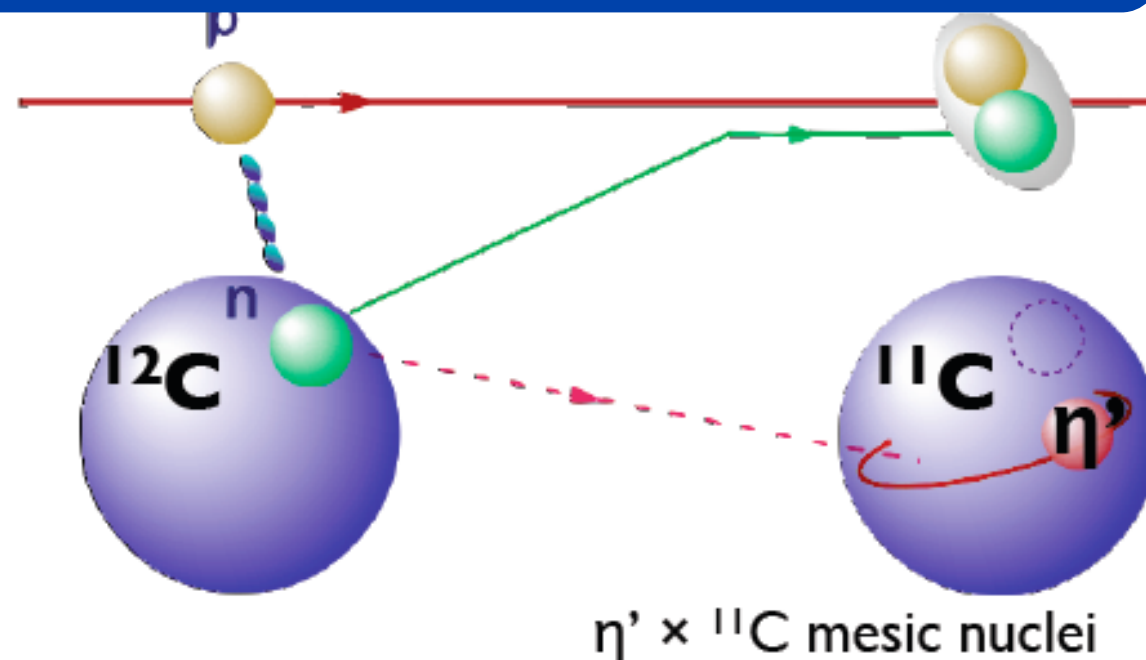
search for η' -mesic states in hadronic reactions

FRS@GSI: PRIME



K. Itahashi et al., PETP 128 (2012) 601

H. Nagahiro et al., PRC 87 (2013) 045201

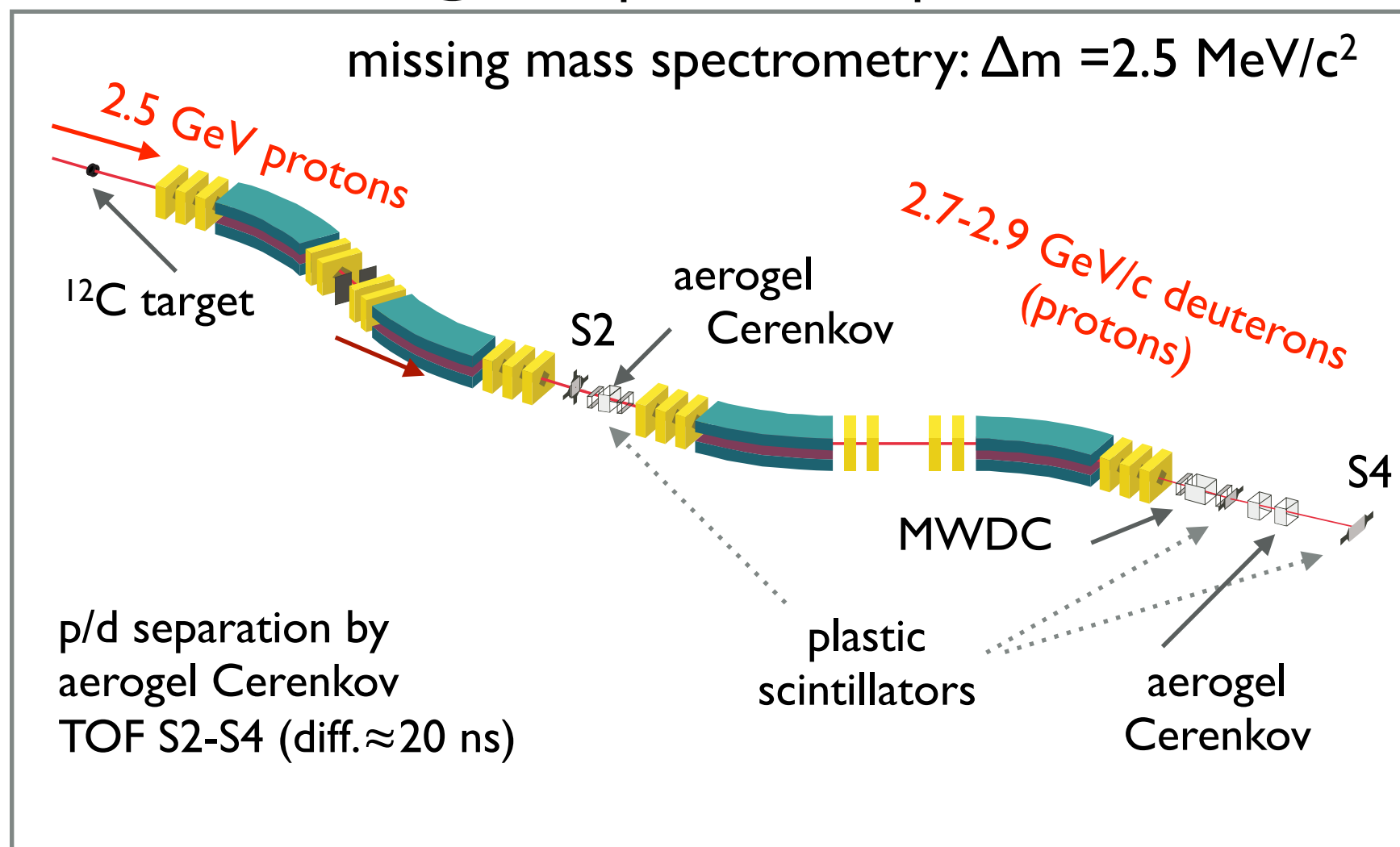


particle identification
by time-of-flight

no structures in
bound state region
observed;
deep potentials with
 $|V_0| > 100 \text{ MeV}$
excluded

FRS@GSI operated as spectrometer

missing mass spectrometry: $\Delta m = 2.5 \text{ MeV}/c^2$



Summary & Outlook

how do the hadron properties (mass, width) change
in a dense nuclear medium ??

meson properties do change in a strongly interacting medium !!

- ◆ all mesons are broadened; their lifetime is shortened through inelastic collisions
 $\Gamma_{\omega}(\rho=\rho_0; p=0) \approx 90 \text{ MeV}$; $\Gamma_{\eta'}(\rho=\rho_0; p=0) \approx 25 \text{ MeV}$;
- ◆ large mass modifications $|\Delta m| > 100 \text{ MeV}$ (as predicted by some calculations) have not been observed
- ◆ for the η' meson an in-medium mass drop of $\Delta m(\rho=\rho_0) \approx -40 \text{ MeV}$ has been determined
- ◆ in-medium effects described within meson-nucleus optical
- ◆ the η' meson is a good candidate for forming meson-nucleus bound states since $|\text{Im } U| \ll |\text{Re } U|$
- ◆ search for η' mesic states ongoing

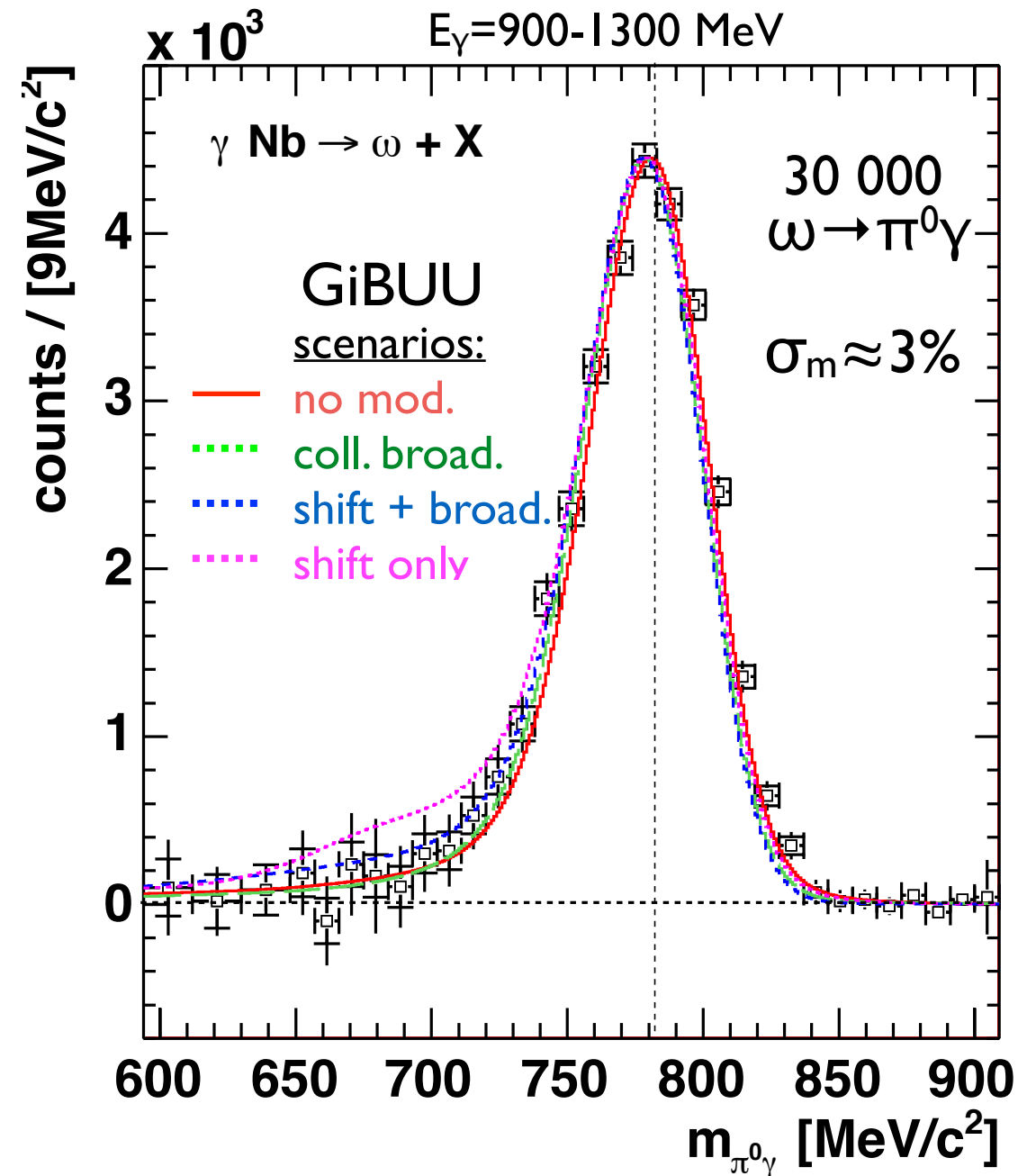
BACKUP

ω line shape from $\omega \rightarrow \pi^0 \gamma$ in photo-nuclear reaction

M.Thiel et al., EPJA 49 (2013) 132

CB/TAPS@MAMI

$\gamma \text{ Nb} \rightarrow \omega + X \rightarrow \pi^0 \gamma + X$ at $E_\gamma = 0.9\text{--}1.3$ GeV



different in-medium modification scenarios almost indistinguishable

ω meson

- $c\tau = 22$ fm/c
for $\langle p_\omega \rangle \approx 600$ MeV/c: $\beta\gamma c\tau = 17$ fm
fraction of ω decays in Nb: $\approx 36\%$
- only mass shifts $\gg \sigma_m/m = 3\%$ observable
- ω signal smeared out due to in-medium broadening ($\Gamma \approx 140$ MeV)
- due to π^0 absorption (π^0 -FSI) $\omega \rightarrow \pi^0 \gamma$ decays in the center of the nucleus are suppressed

η' meson

- $c\tau = 1000$ fm/c
for $\langle p_{\eta'} \rangle \approx 1000$ MeV/c: $\beta\gamma c\tau = 1000$ fm
fraction of η' decays in Nb: $\approx 0.5\%$

line shape analysis very difficult or even impossible