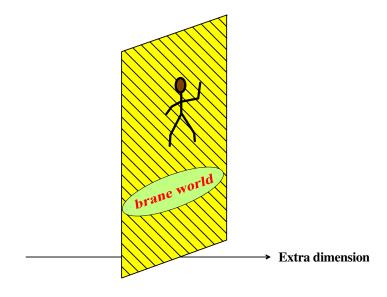
## Theory of Cosmological Perturbations

## Addendum

— a touch on brane-cosmological perturbations —

## § 1. Braneworld scenario

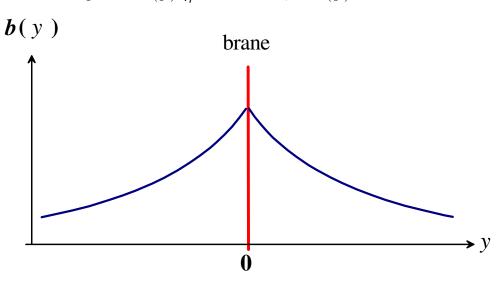
 $\cdot$  Our universe is a (timelike) singular surface (brane) in a higher dimensional spacetime.



- · Standard matter fields are confined on the brane, while gravity propagates in the whole spacetime (called the bulk).
- · One of the most studied models is the Randall-Sundrum model in which the bulk is 5D anti-de Sitter (AdS) space.
- $\cdot$  Cosmology on the brane = brane dynamics in the bulk.

• Randall-Sundrum braneworld metric (AdS<sub>5</sub> + Minkowski brane):

$$ds^2 = dy^2 + b^2(y)\eta_{\mu\nu} dx^{\mu} dx^{\nu}; \quad b(y) = e^{-|y|/\ell}$$



• Generalization to cosmology:

Brane fixed at y = 0 ( $\sim$  view from the brane)

$$ds^{2} = dy^{2} + \left(-n^{2}(y,t)dt^{2} + a^{2}(y,t)d\sigma_{(3)}^{2}\right)$$

Or, a brane moving in the AdS bulk ( $\sim$  view from the bulk)

$$ds^{2} = -A(R)dT^{2} + \frac{dR^{2}}{A(R)} + R^{2}d\Omega_{K}^{2} \quad A(R) = K + \frac{R^{2}}{\ell^{2}} \quad (K = \pm 1, 0)$$

with brane trajectory:  $(T, R) = (T(\tau), R(\tau))$ 

## § 2. Large-scale cosmological perturbations on the brane

· General formalism for braneworld perturbations

· Essentially a 5-dimensional, PDE problem:

$$Q(y, x^{\mu}) = Q_p(y, t)Y_p(x^i) \cdots Q_p(y, t)$$
 is not separable.

· Nevertheless, some simplifications on super-horizon scales.

• Effective gravitational equations on the brane (in AdS<sub>5</sub> bulk background)

$$G_{\mu\nu} + \Lambda_4 \, q_{\mu\nu} = 8\pi G_4 \, T_{\mu\nu} + (8\pi G_5)^2 \, \Pi_{\mu\nu} - E_{\mu\nu}$$
  
 $\equiv 8\pi G_4 \, T_{\mu\nu}^{\text{tot}} \quad (\Pi_{\mu\nu} \sim \text{quadratic in } T_{\mu\nu})$ 

$$E_{\mu\nu} \equiv \overset{(5)}{C}_{a\mu b\nu} n^a n^b;$$
  $\overset{(5)}{C}_{acbd} \cdots$  5D Weyl tensor  $n^a \cdots$  unit normal to the brane

- $\cdot E_{\mu\nu}$  contains all the dynamics of the bulk
- · By definition,  $E^{\mu}_{\mu} = 0$ . In addition,  $D^{\mu}E_{\mu\nu} = 0$  on FLRW background.

"
$$-E_{\mu\nu}$$
": Weyl fluid (or "dark radiation")

The energy momentum tensor of the brane matter:

$$T_{\mu\nu} = \rho \, u_{\mu} u_{\nu} + p \, (u_{\mu} u_{\nu} + g_{\mu\nu}) + \pi_{\mu\nu}$$

For FLRW background, with perturbations of  $O(\epsilon)$ ,

$$\pi_{\mu\nu}$$
: anisotropic stress  $= O(\epsilon)$   $\Rightarrow$   $u^{\mu}D^{\nu}E_{\mu\nu} = O(\epsilon^2)$ 

On superhorizon scales, only the energy conservation law is important. (: momentum cannot be transfered over super-Hubble scales)

Weyl fluid decouples from the brane matter on superhorizon scales.



Standard 4D theory is applicable with slight modifications due to  $\Pi_{\mu\nu}$  (quadratic in  $T_{\mu\nu}$ ).

$$\rho_{\text{eff}} = \rho \left( 1 + \frac{\rho}{2\sigma} \right) \quad (\sigma = \text{brane tension}), \quad \rho_{\mathcal{E}} = \frac{E^0_0}{8\pi G_4}, \quad \text{etc.}$$

What is missing is an equation for anisotropic stress of the Weyl fluid:

$$\pi_{\mathcal{E}}^{ij} = -\frac{E_T^{ij}}{8\pi G_4}; \quad E_{Tj}^i \equiv E_j^i - \frac{1}{3}\delta_j^i E_k^k$$

This must be determined by solving the bulk perturbation equations.

$$\Rightarrow$$
 back to 5D problem

- Some recent ( $\gtrsim 2003$ ) efforts to solve the bulk perturbation equations: (in the context of inflation in RS-type brane cosmology)
  - $\cdot$  Expanding equations around the brane ( $\sim$  gradient expansion).

Koyama ('03), Battye, Van de Bruck & Mennim ('03), · · ·

· Approximating the brane by a de Sitter space.

Minamitsuji, Himemoto & MS ('03), Du, Wang, Abdalla & Su ('04), · · ·

· Analyzing exactly soluble models.

Koyama & Takahashi ('03), Yoshiguchi & Koyama ('04), Kobayashi & Tanaka ('04), . . .

· Focusing on tensor-type perturbations

Easther, Langlois, Maartens & Wands ('03), Hiramatsu, Koyama & Taruya ('03), ...

· Elaborating the perturbation theory

Deffayet ('03), Malik, Rodriguez-Martinez & Langlois ('03), ...

But no rigorous treatment of scalar-type perturbations so far