



Australian  
National  
University

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# MASSIVE STAR FORMATION

(A PRIMER FOR PEOPLE WHO DON'T CARE ABOUT THE ISM)

## MOTIVATION FOR THIS TALK

- ▶ Massive stars are the ultimate sources of many of the multi-messenger signals of interest to most at this meeting
- ▶ The demographics of massive stars are therefore a crucial input to efforts to model source populations; conversely, these demographics can be constrained by observations
- ▶ Many demographic features matter: IMF, binary statistics, spin; we care about how these depend on environment
- ▶ My goal: tell you what is known theoretically (and a little observationally) about these demographics

# OUTLINE

- ▶ Observational background on massive star formation
- ▶ Key physical processes for massive star formation
  - ▶ Fragmentation
  - ▶ Feedback and the upper mass limit
  - ▶ Disks: fragmentation and braking
- ▶ Variations: what changes with environment, and why?
  - ▶ The initial mass function
  - ▶ Binarity / multiplicity
  - ▶ Rotation
- ▶ Implications and questions



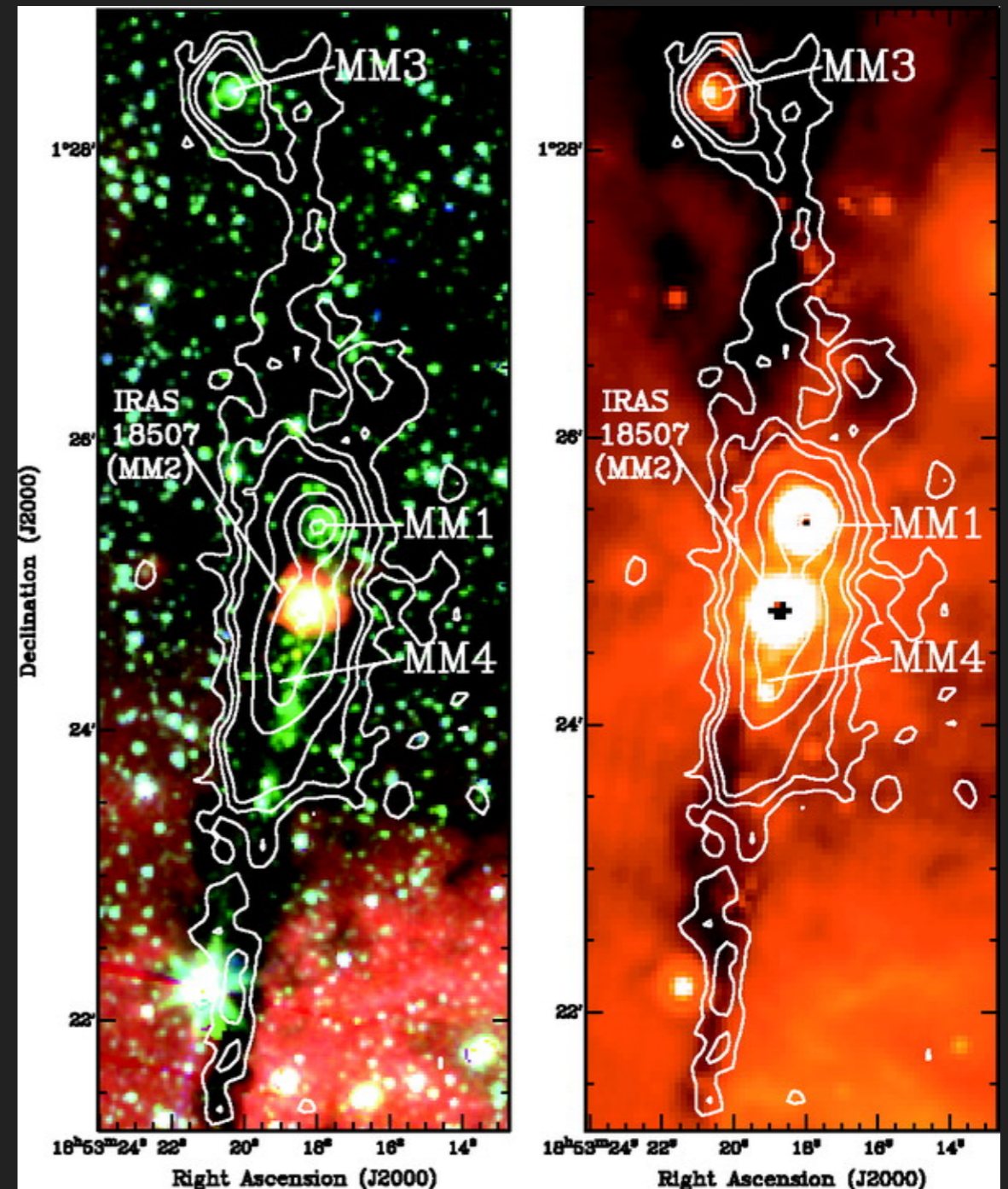
Action Press/REX/Shutterstock

# OBSERVATIONAL BACKGROUND

For no reason whatsoever, here is a baby wombat

## SITES OF MASSIVE STAR FORMATION

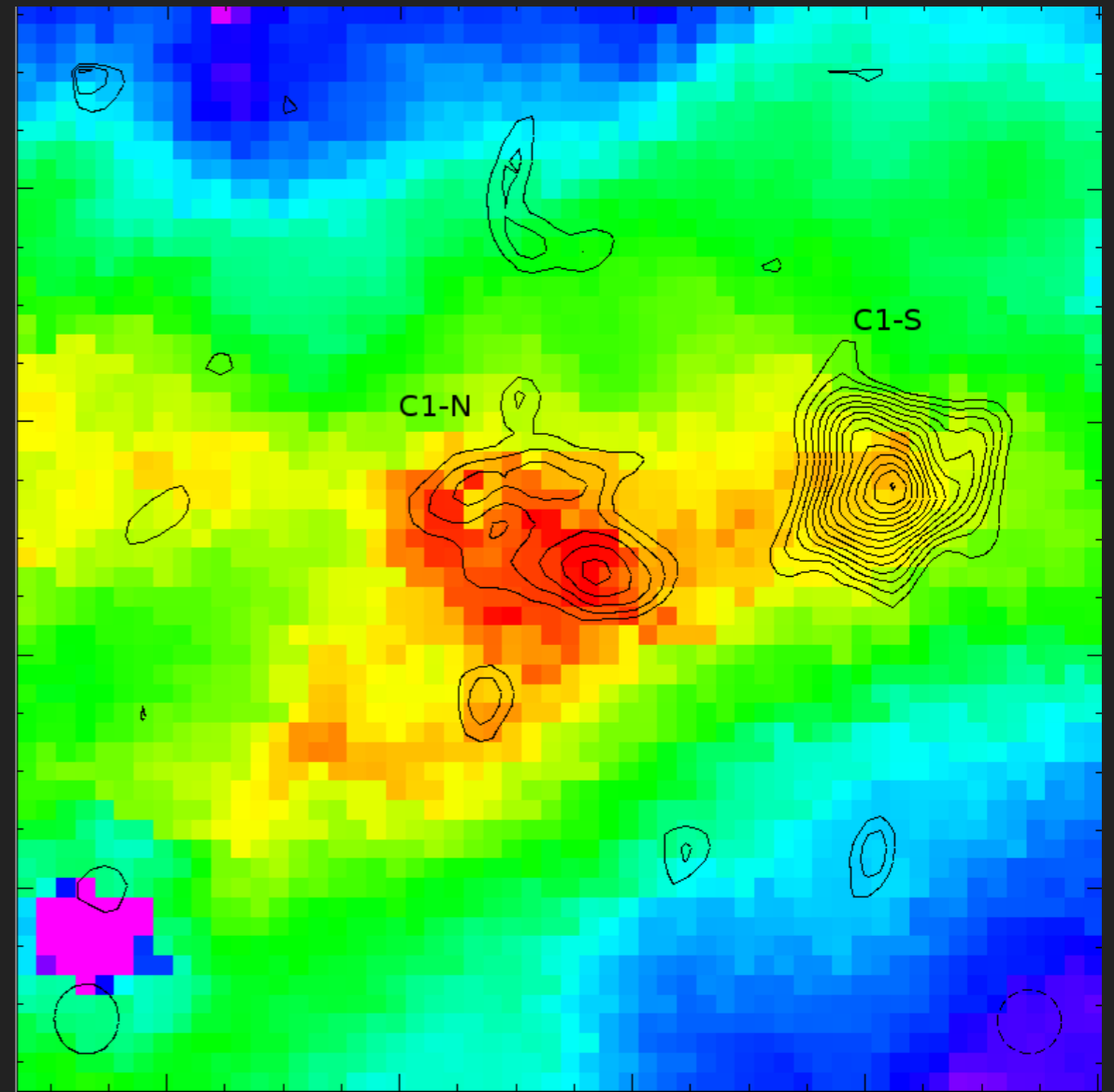
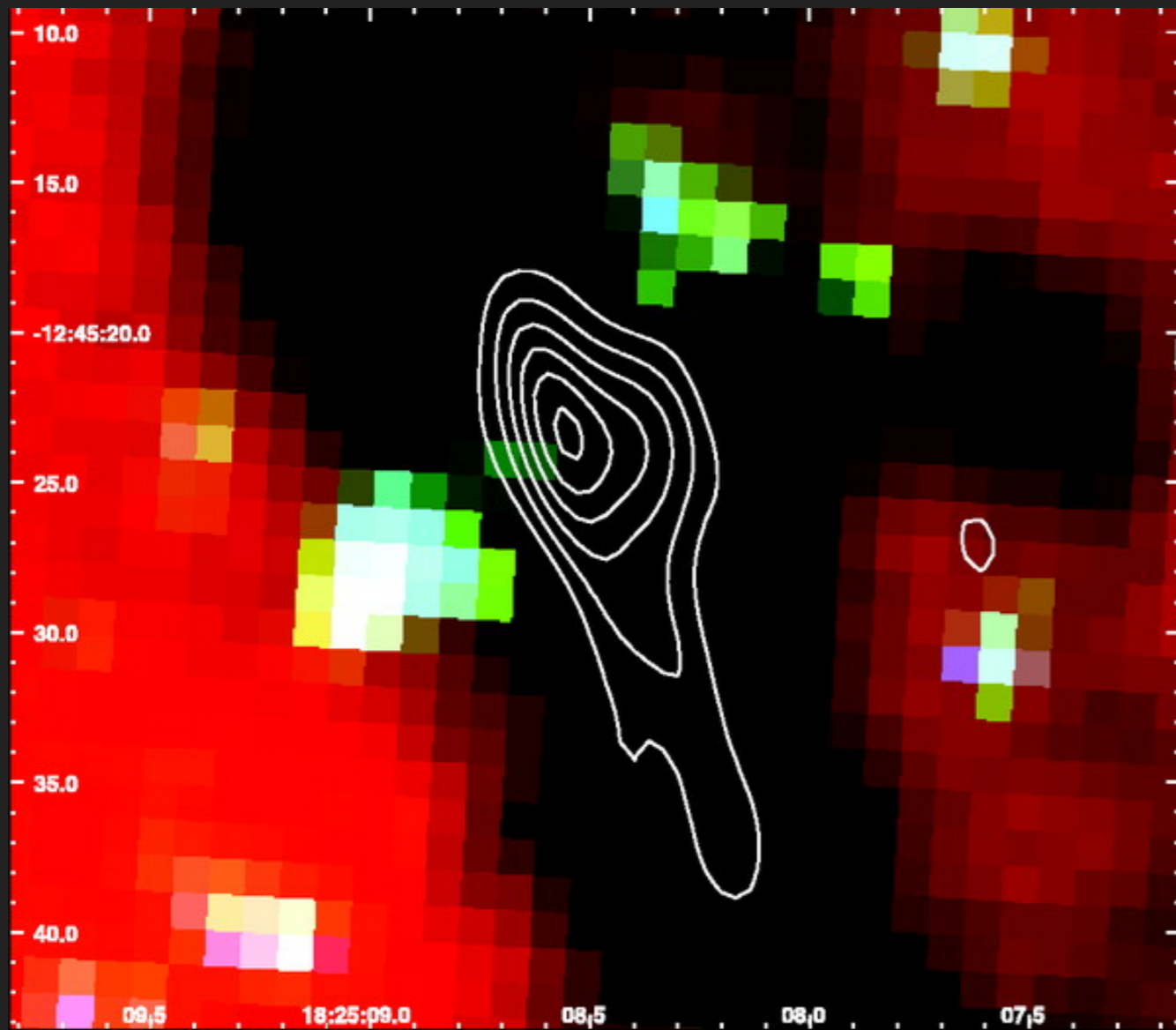
- ▶ Massive stars form in cold, dense, dusty interstellar clouds
- ▶ Detect massive protostars in clouds by mid-IR emission
- ▶ Detect clouds by near- or mid-IR absorption, FIR or mm emission
- ▶ Typical surface density  $\Sigma \sim 0.1 - 1 \text{ g cm}^{-2}$ , temperature  $T \sim 10 \text{ K}$
- ▶ Contain MIR-dark cores with  $M \sim 100 M_{\odot}$ ,  $\Sigma \sim 1 \text{ g cm}^{-2}$ ,  $\sigma \sim 1 \text{ km s}^{-1}$



Rathborne+ 2005

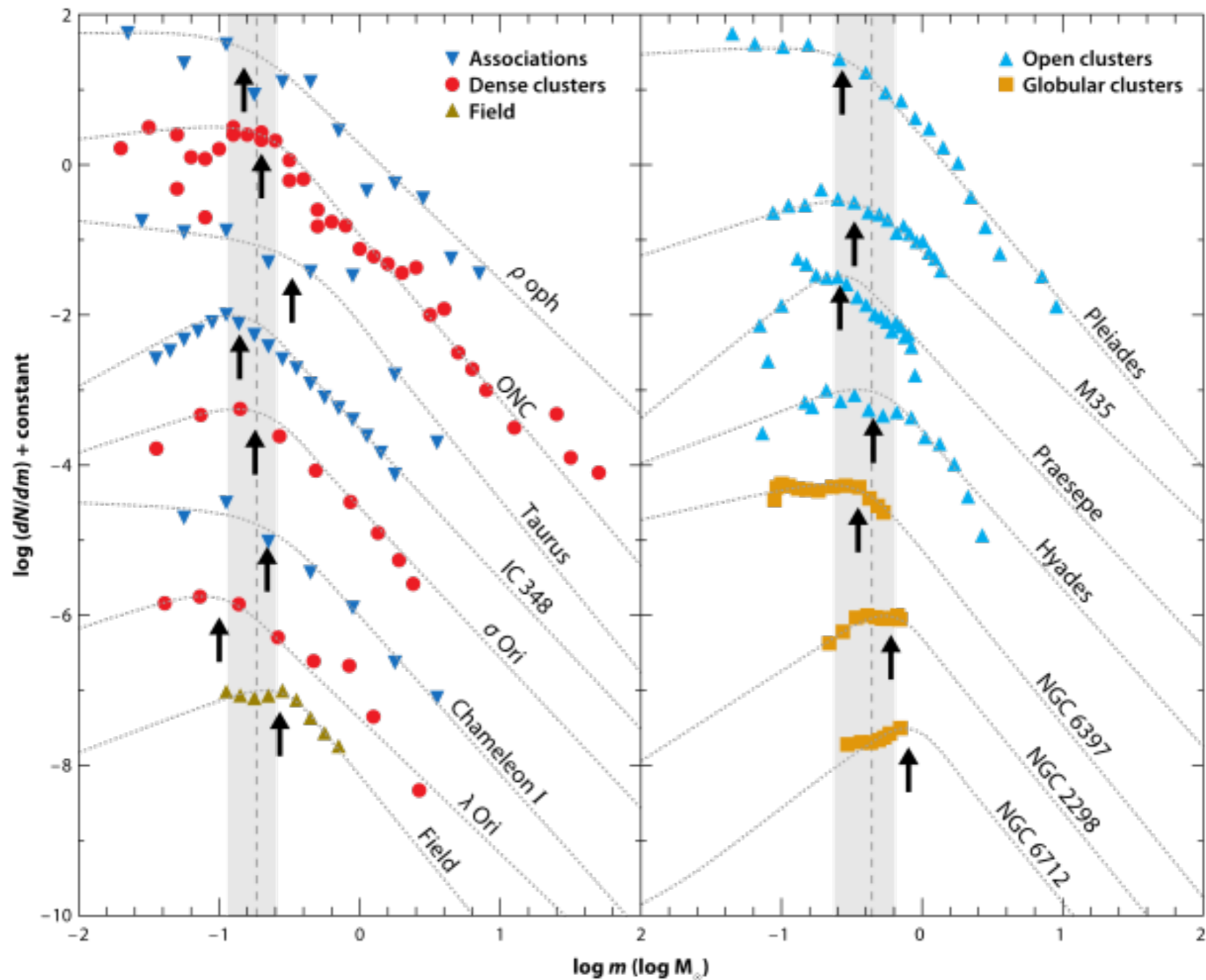
Left: Spitzer NIR + IRAM mm

Right: Spitzer MIR + IRAM mm

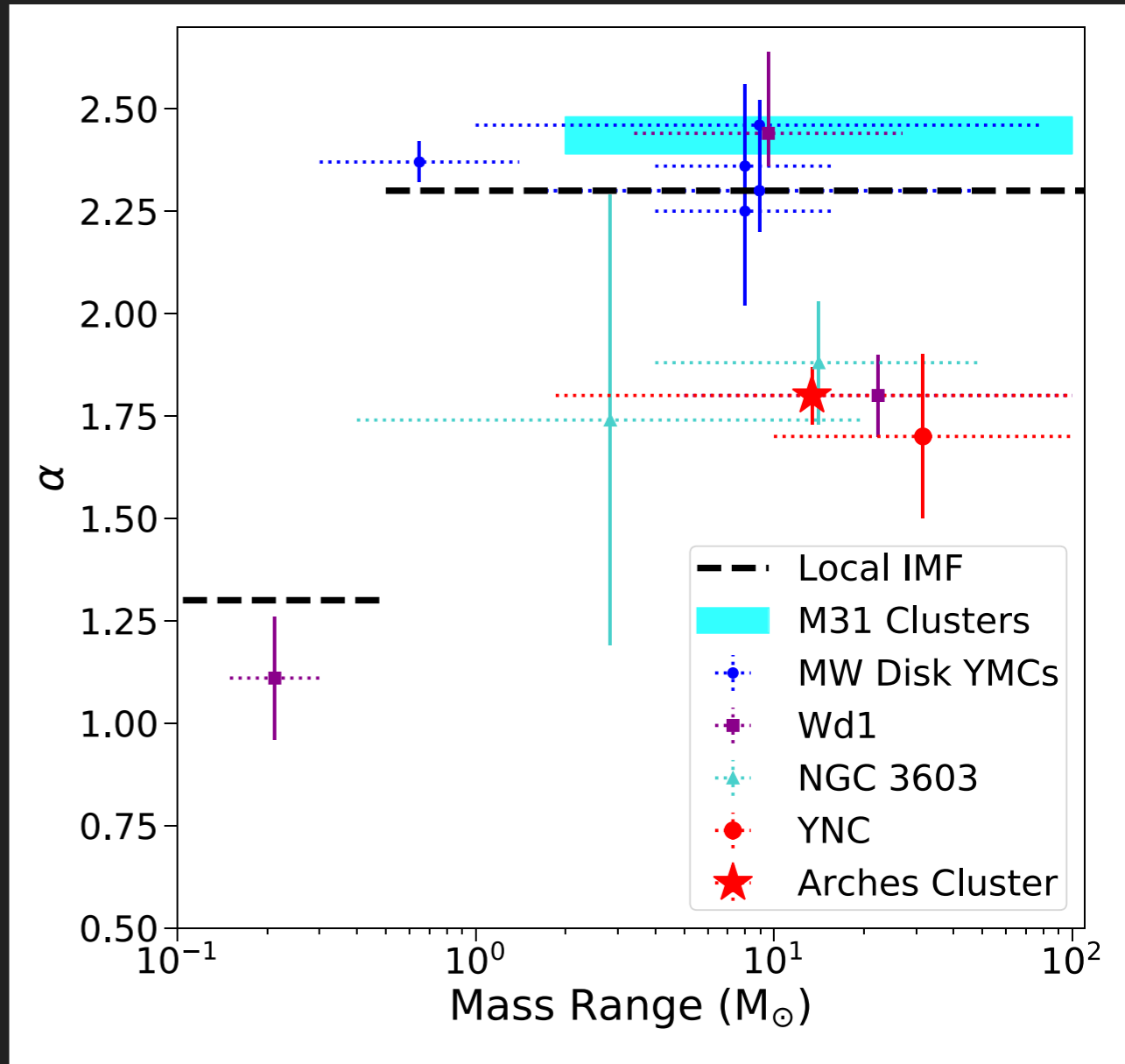
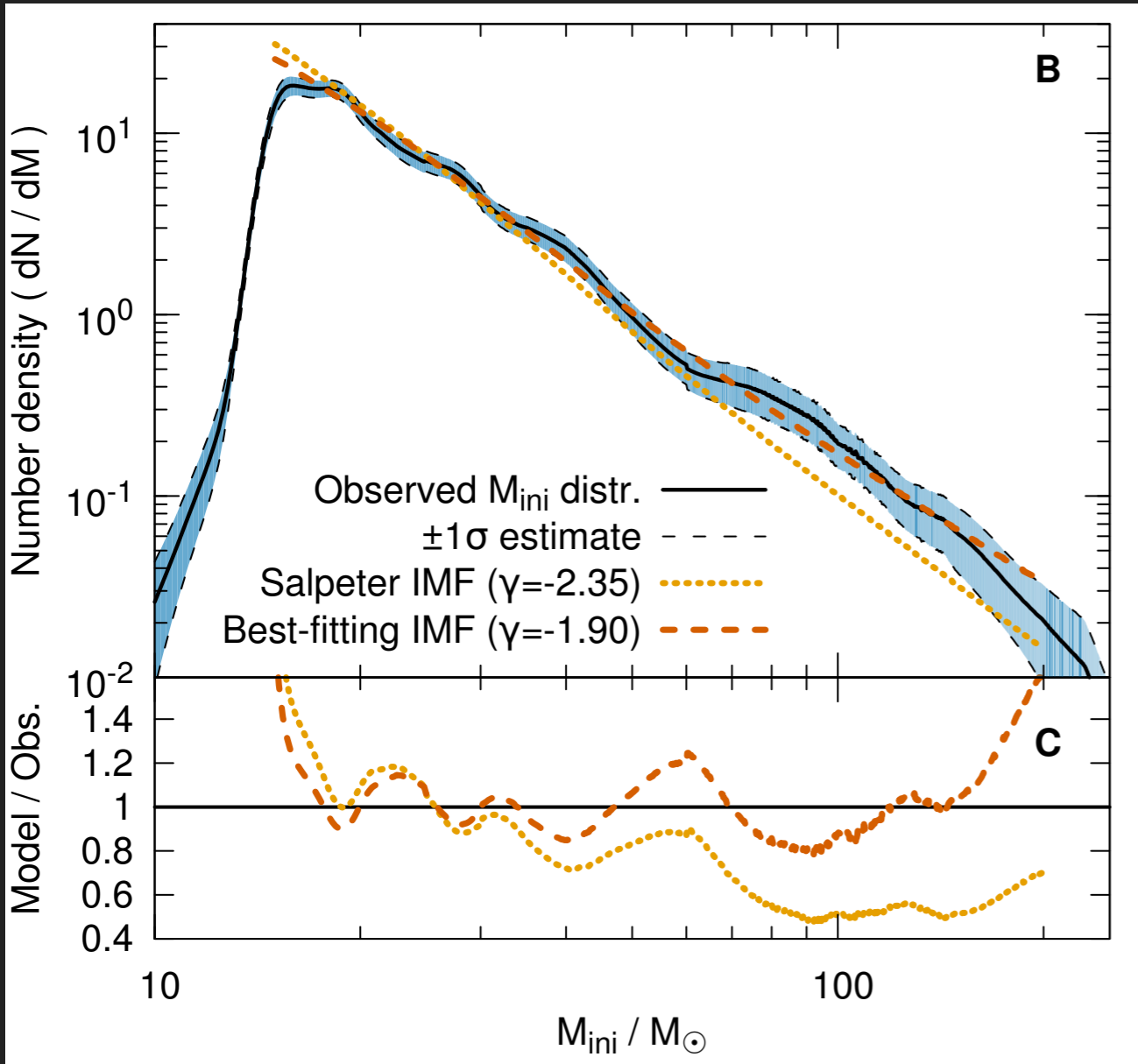


# MASSIVE CORES

Left: Beuther+ 2005, Spitzer MIR + IRAM mm  
Right: Tan+ 2014, Herschel MIR + ALMA mm



# THE OBSERVED IMF



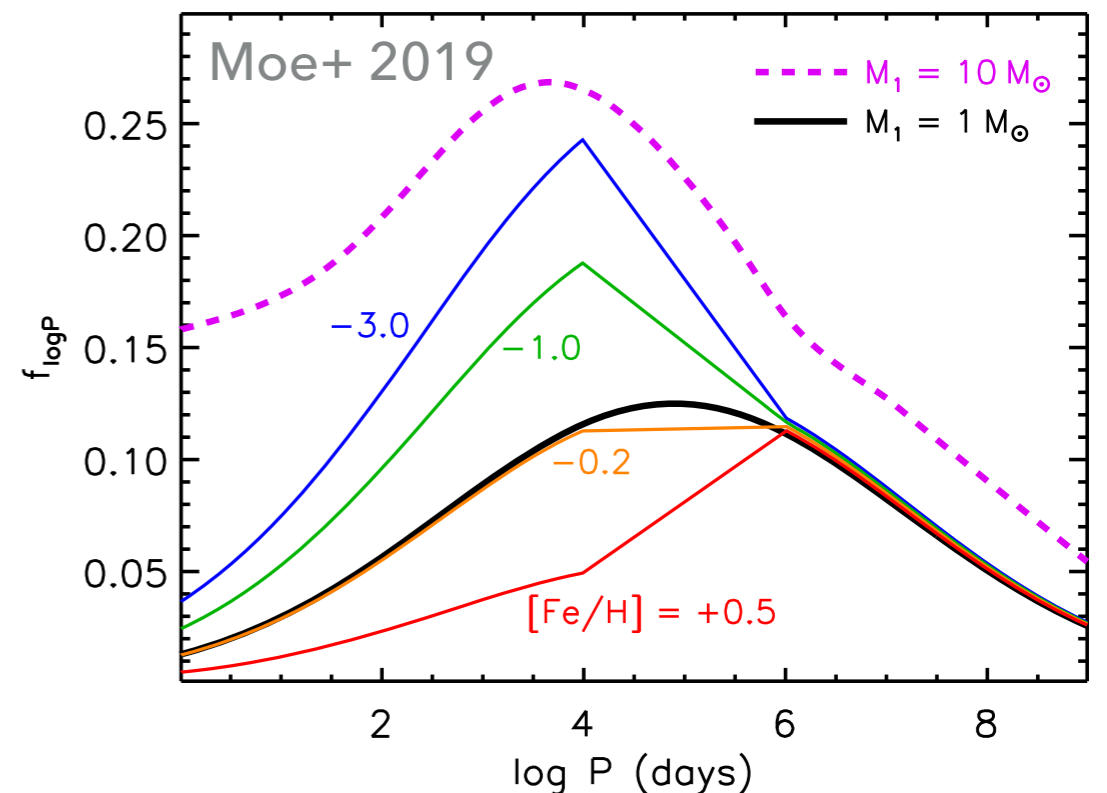
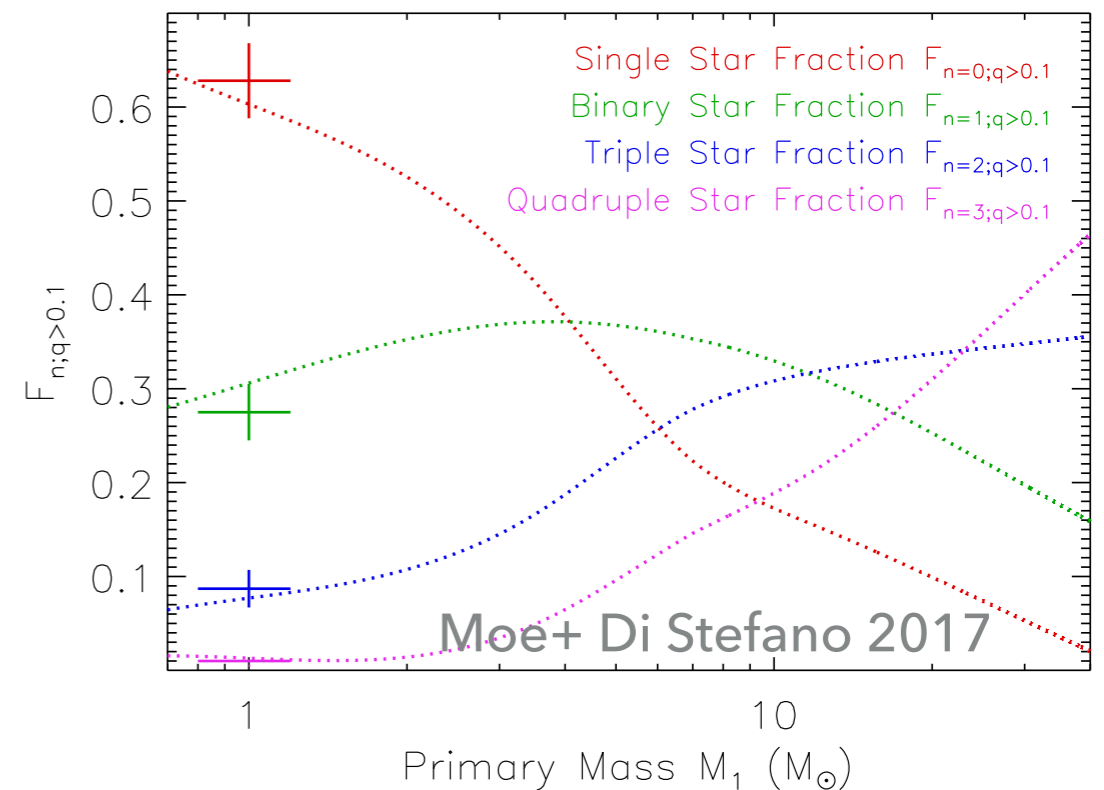
# IMF VARIATIONS

Left: IMF in 30 Doradus, Schneider+ 2018  
 Right: IMF in MW center clusters, Hosek+ 2019



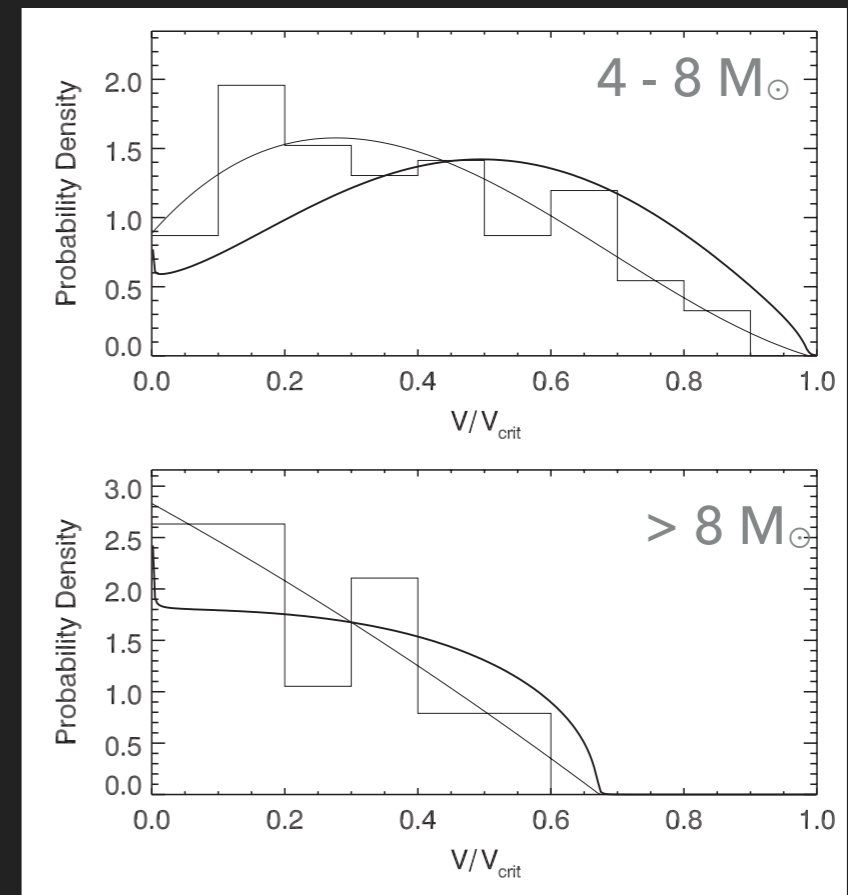
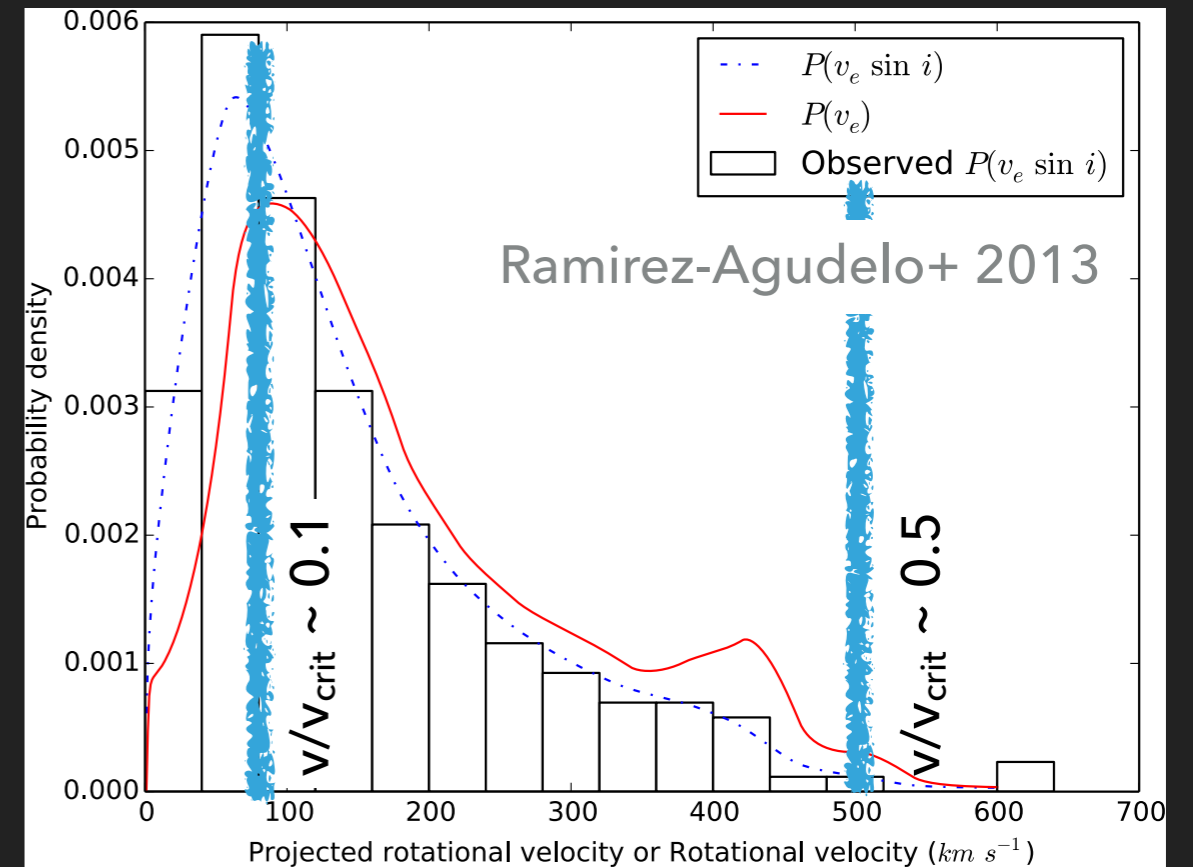
# MASSIVE STAR MULTIPLICITY

- ▶ Essentially all massive stars form in multiple systems, many with multiple companions
- ▶ Compared to lower mass stars at Solar metallicity, significant excess of close companions
- ▶ Mass ratios mostly moderate at most separations,  $q \sim 0.3 - 0.5$ , with 10% excess of "twin" binaries ( $q > 0.95$ ) at small separation



# MASSIVE STAR ROTATION

- ▶ Observed massive stars mostly rotating at  $\sim 10\text{-}20\%$  of breakup, but with tail to up  $\sim 50\%$  of breakup
- ▶ Rotation speed at birth very likely higher, close to  $\sim 50\%$  – stars spin down due to magnetised winds
- ▶ Evidence for wind spin-down: B stars with weak winds typically at  $\sim 50\%$  of breakup





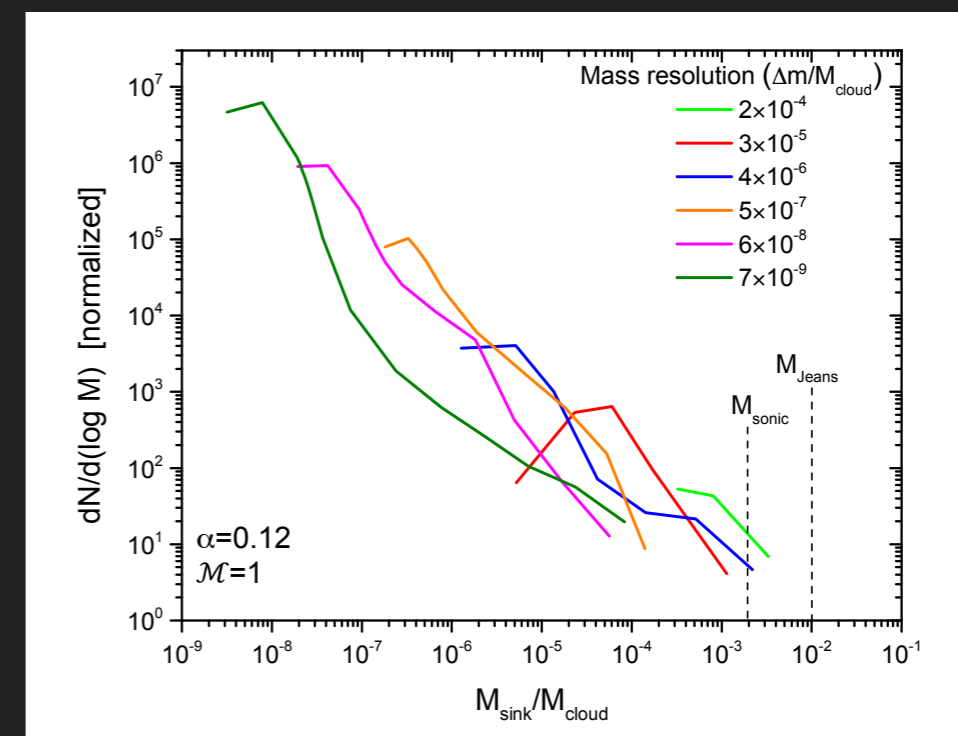
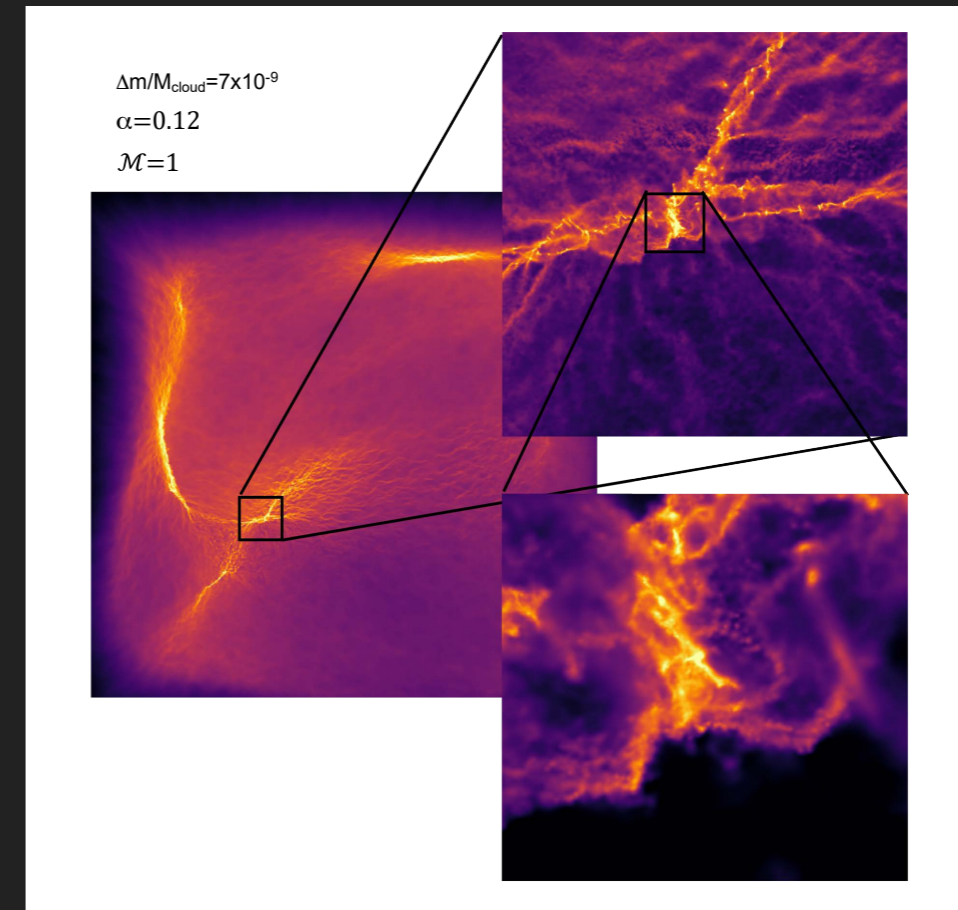
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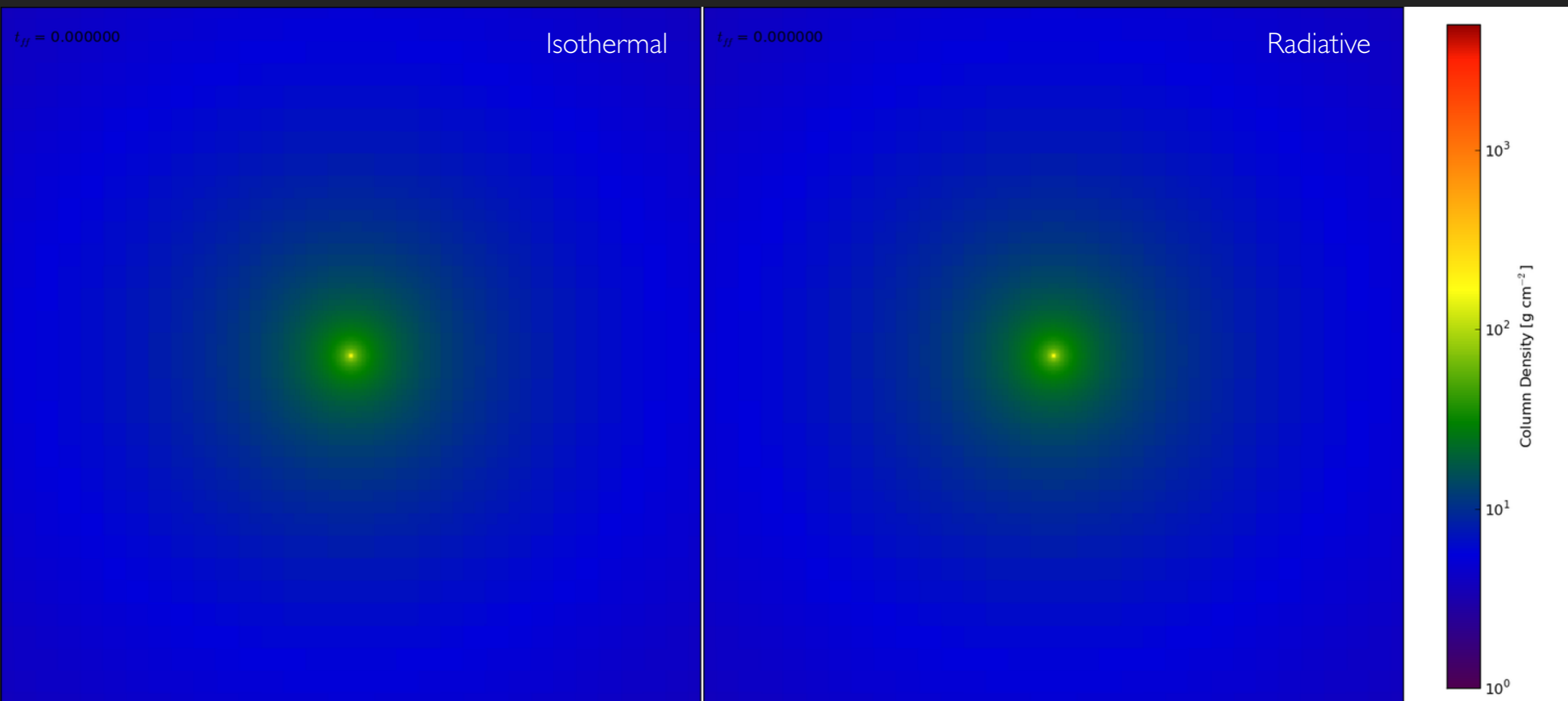
# MASSIVE STAR FORMATION PHYSICS

For no reason whatsoever, here is  
a baby quokka

# ISOTHERMAL FRAGMENTATION

- ▶ Jeans mass  $M_J \sim \rho^{-1/2}$ , so as collapse occurs, mass that is able to fragment goes to zero
- ▶ Numerical experiments show that this produces fragmentation to infinitely small scales
- ▶ To form a massive star, the fragmentation cascade must be halted
- ▶ Likely agent: radiative feedback



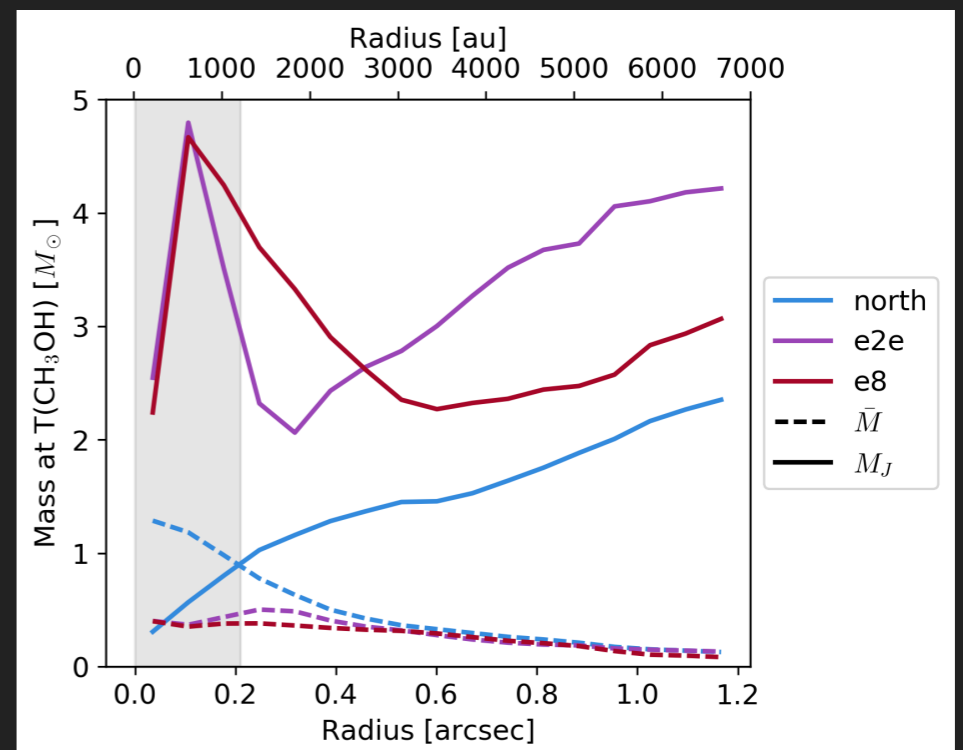
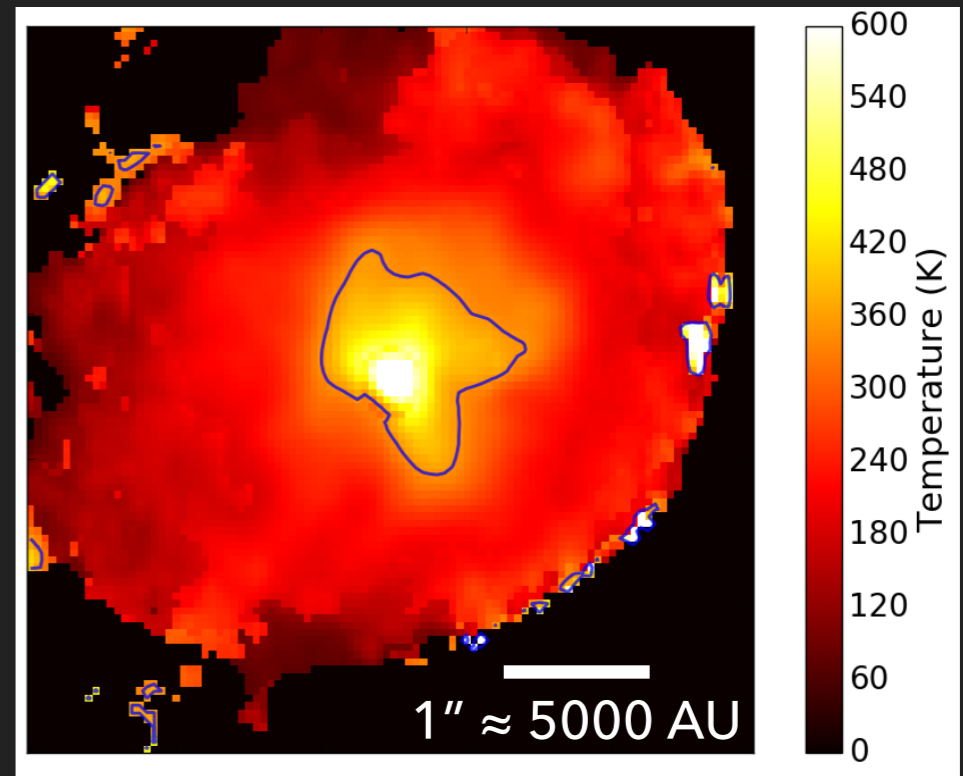


# RMHD SIMULATIONS

Myers+ 2013  
200  $M_{\odot}$  centrally-condensed core

## OBSERVATIONAL EVIDENCE

- ▶ Observations of temperature structure around massive protostars shows warm gas
- ▶ Observed heating sufficient to suppress fragmentation on  $>1000$  AU scales
- ▶ Supports the idea that radiative feedback is key to allowing massive star formation

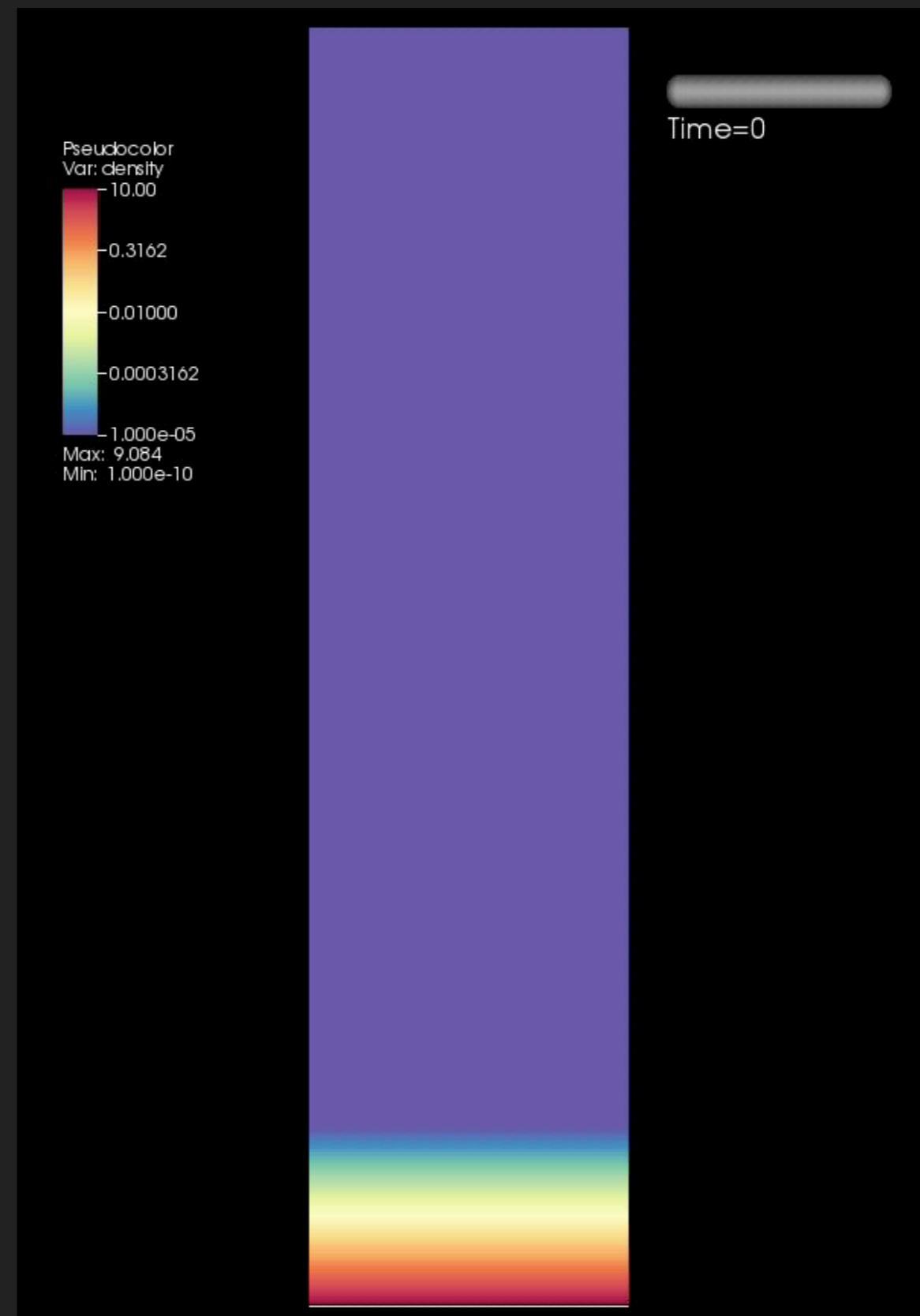


## LIMITING STELLAR MASSES: WINDS AND PHOTOIONIZATION

- ▶ Photoionization feedback mostly ineffective because  $dM/dt \sim 10^{-4}$  sufficient to keep ionized region trapped near star (Walmsley 1995, Keto+ 2002, 2003, 2007)
- ▶ Main sequence winds can only become important at masses above  $\sim 40 M_{\odot}$  – otherwise star is bloated and has  $T_{\text{eff}}$  too low to drive wind
- ▶ Winds conceivably important after that, but only if they become trapped; otherwise too little momentum

## DIRECT RADIATION PRESSURE

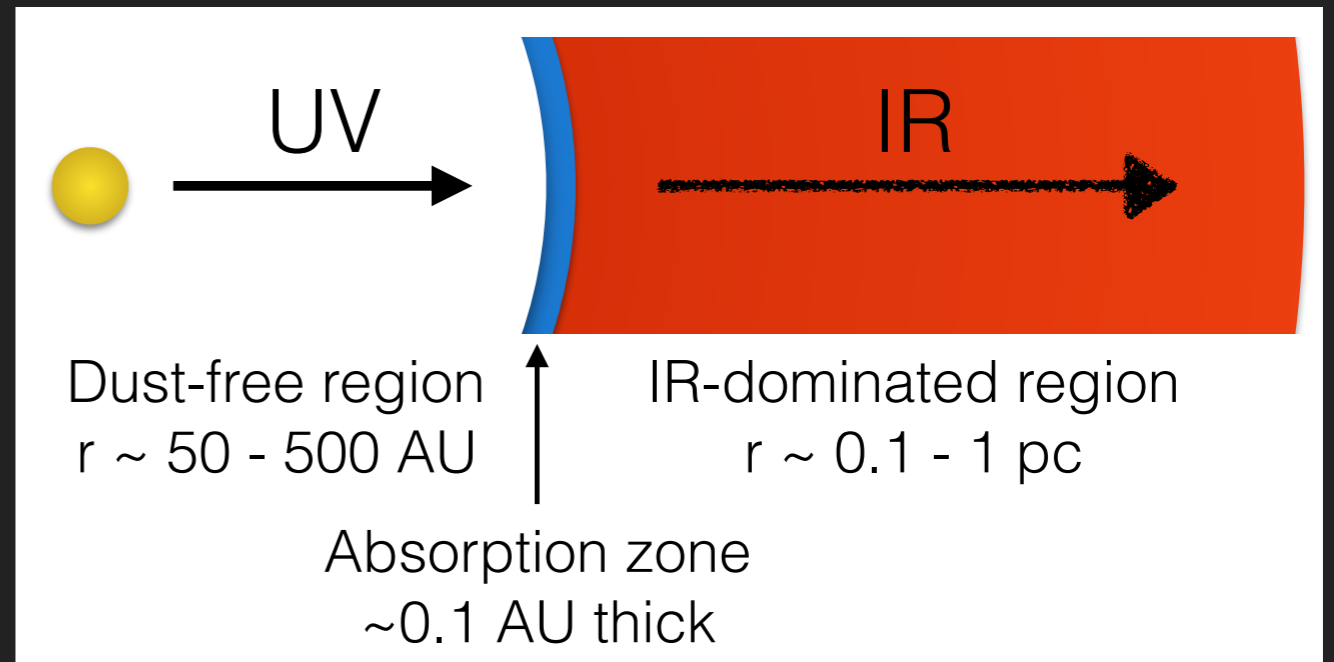
- ▶ Radiation force  $>$  gravitational force on any gas column with  $\Sigma < \Sigma_{\text{crit}} = (L/M) / 4\pi Gc \sim 300 M_{\odot} \text{ pc}^{-2}$  (Fall, Krumholz, & Matzner 2010)
- ▶ In a turbulent medium with a PDF of  $\Sigma$ 's, low  $\Sigma$  regions unstable to ejection even if mean  $\Sigma > \Sigma_{\text{crit}}$  (Thompson & Krumholz 2016)
- ▶ However, most massive stars have  $L/M \approx 10^4 L_{\odot}/M_{\odot} \rightarrow \Sigma_{\text{crit}} \approx 0.8 \text{ g cm}^{-2}$ ; direct radiation pressure cannot set a mass limit in cores of higher  $\Sigma$

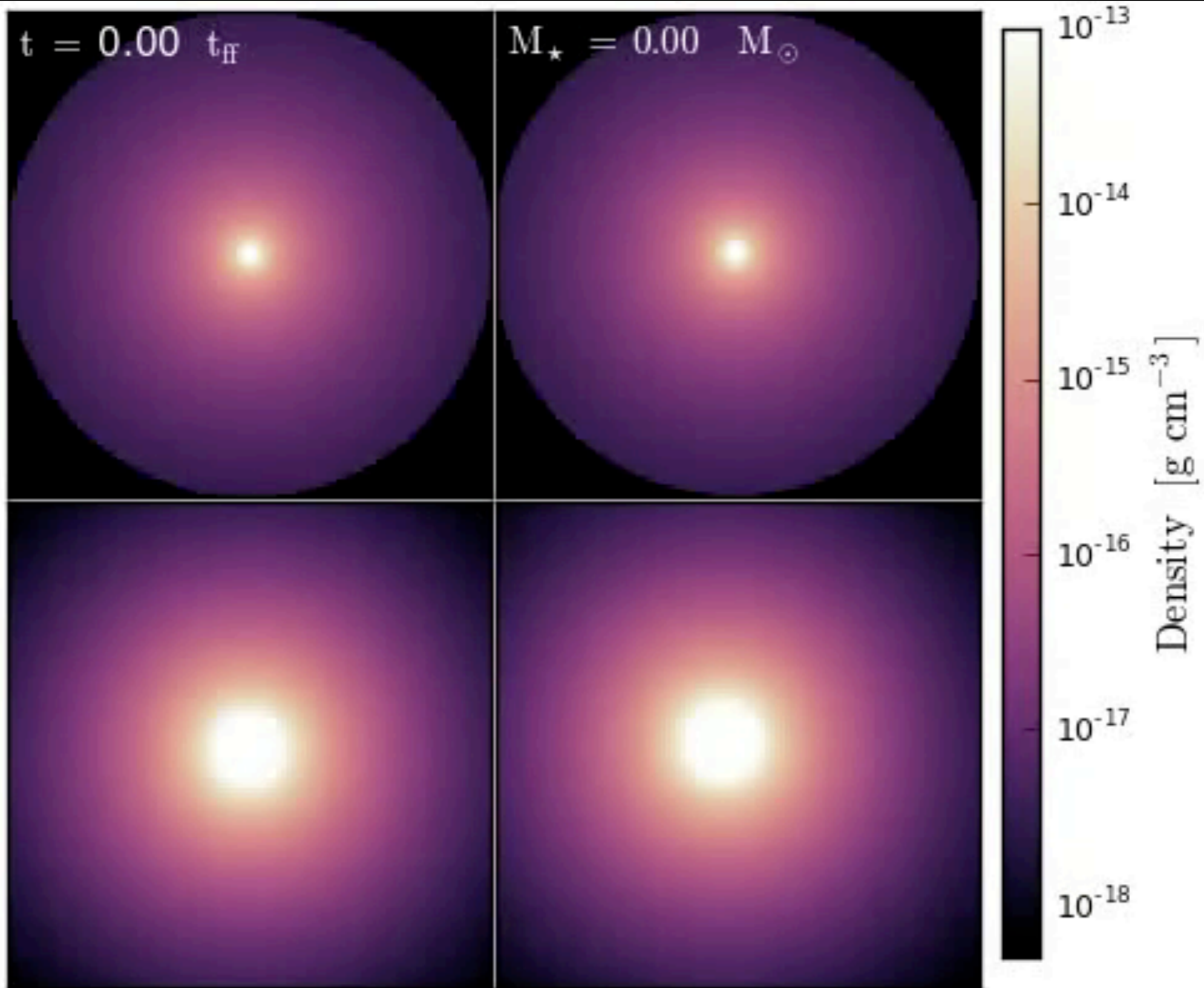




## INDIRECT RADIATION PRESSURE

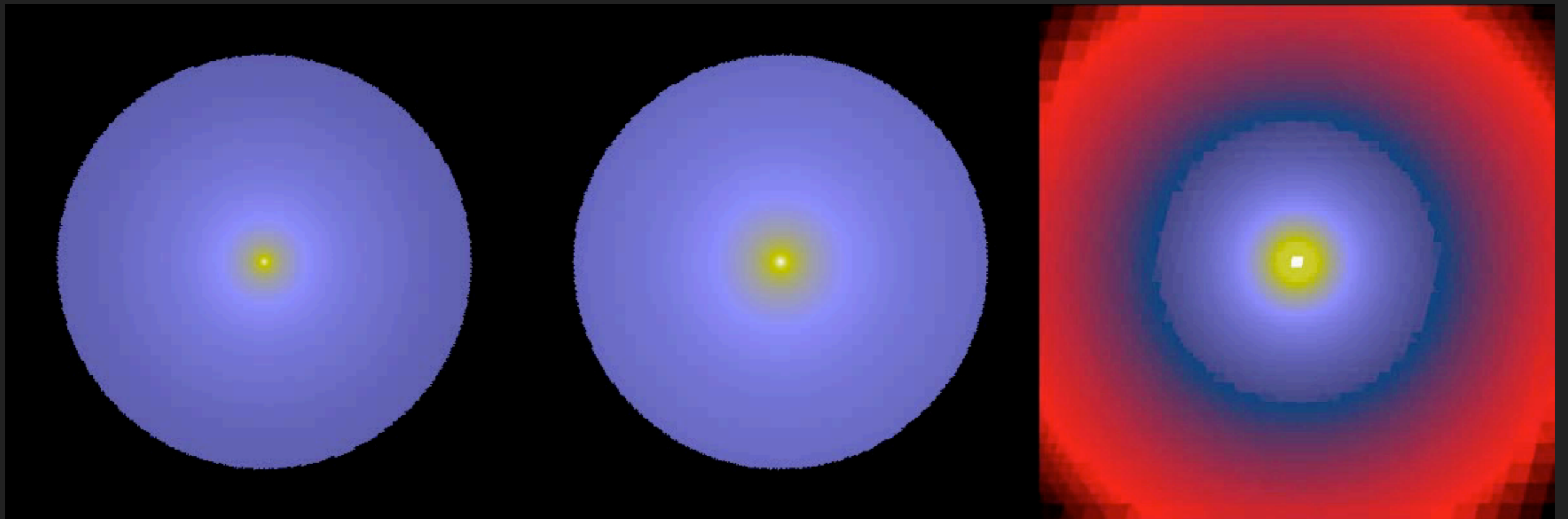
- ▶ Near massive star, radiation creates a dust-free zone with low opacity
- ▶ Radiation free-streams to dust destruction front, is reprocessed into IR
- ▶ In IR-dominated region, Eddington ratio is for isotropic radiation flux is  $f_{\text{Edd}} = \kappa_{\text{IR}} L / 4\pi G M c \approx 8 (\kappa_{\text{IR}} / 10 \text{ cm}^2 \text{ g}^{-1})$
- ▶ Thus accretion is possible only if some mechanism makes the radiation flux anisotropic





# BEATING RP: THE RRT INSTABILITY

Rosen+ 2016, 2017



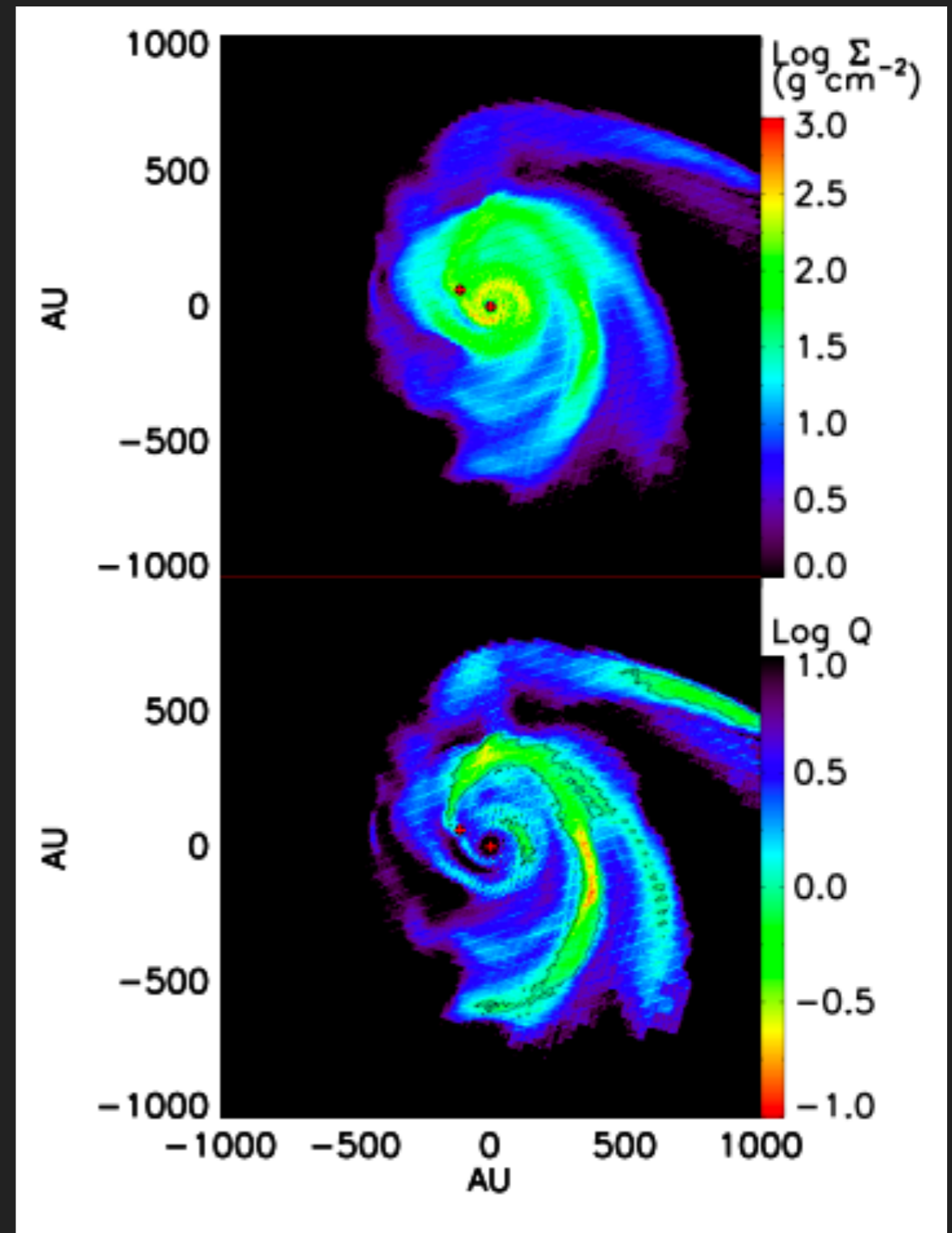
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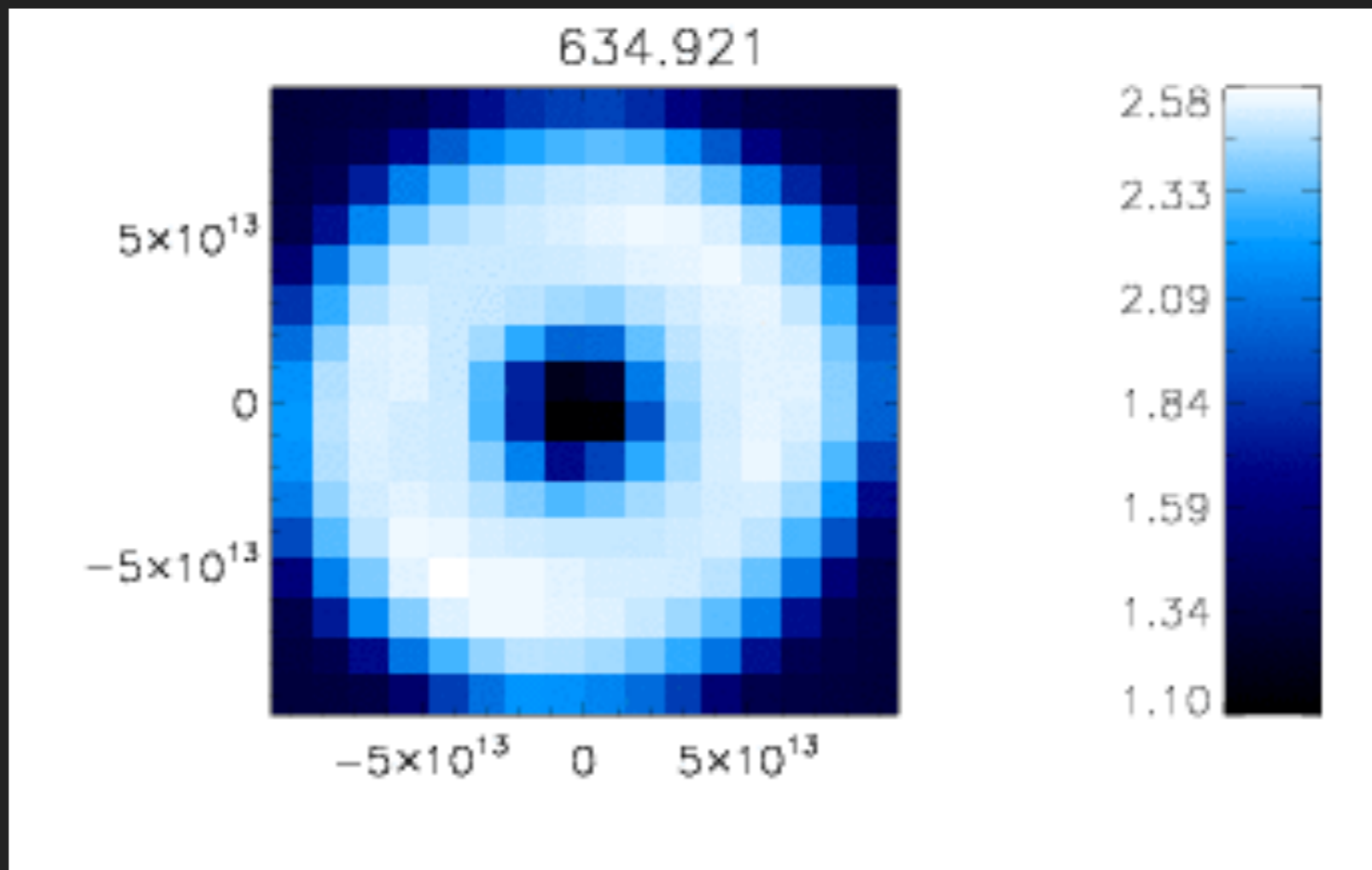
# BEATING RP: OUTFLOWS

Cunningham+ 2011

## MASSIVE STAR DISKS

- ▶ Accretion rate onto disk  $\dot{M} \sim \sigma^3 / G$   
 $\sim \text{few} \times 10^{-4} M_{\odot} \text{ yr}^{-1}$
- ▶ Disk accretion rate  $\dot{M} = 3\alpha c_s^3 / GQ$   
 $= 1.5 \times 10^{-4} T_2^{3/2} (\alpha/Q) M_{\odot} \text{ yr}^{-1}$
- ▶ Implication: disk can only deposit material on star as quickly as it accretes if  $\alpha \approx 1$  AND  $Q \approx 1$
- ▶ Disk likely to be gravitationally unstable (Kratter & Matzner 2006; Kratter+ 2008, 2010)





# DISK FRAGMENTATION

## MAGNETIC BRAKING BY DISKS

- ▶ Bipolar stellar B field:

$$B_z = B_* (r/R_*)^{-3}$$

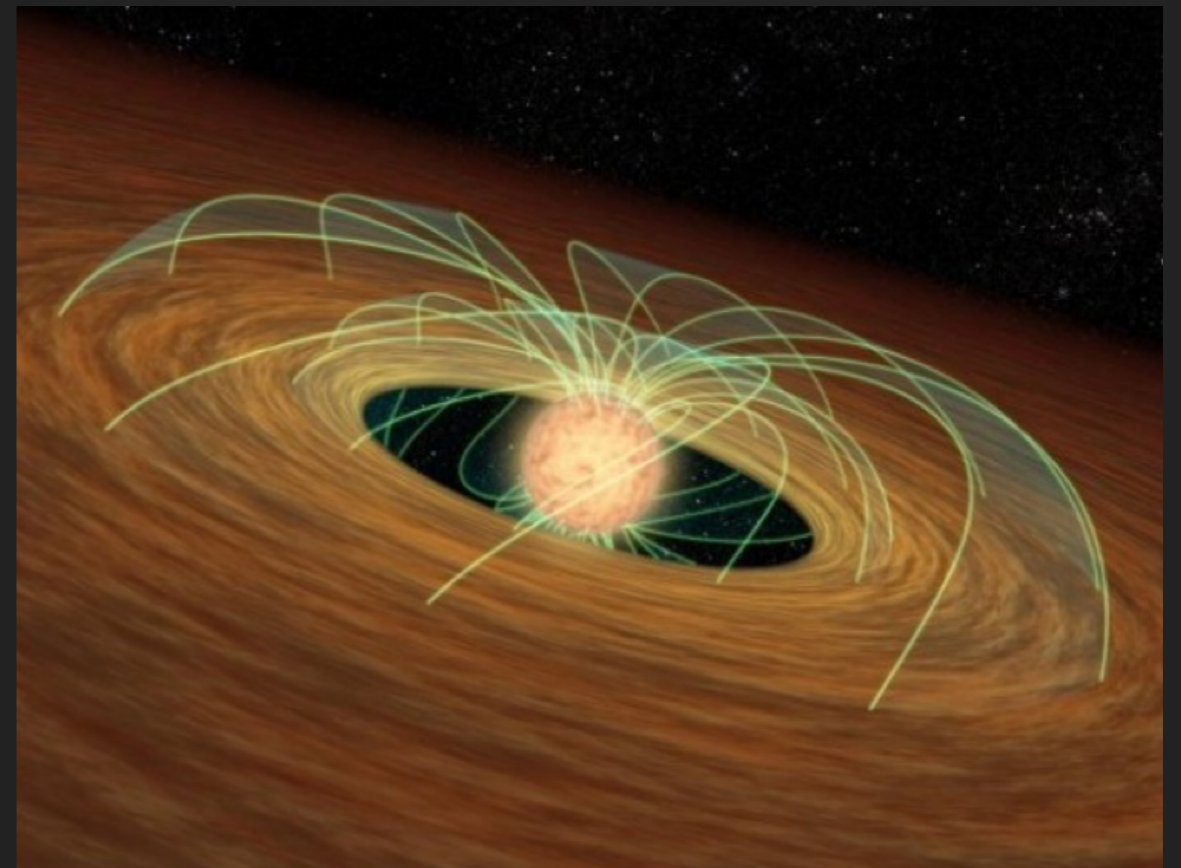
- ▶ Magnetic truncation of disk:

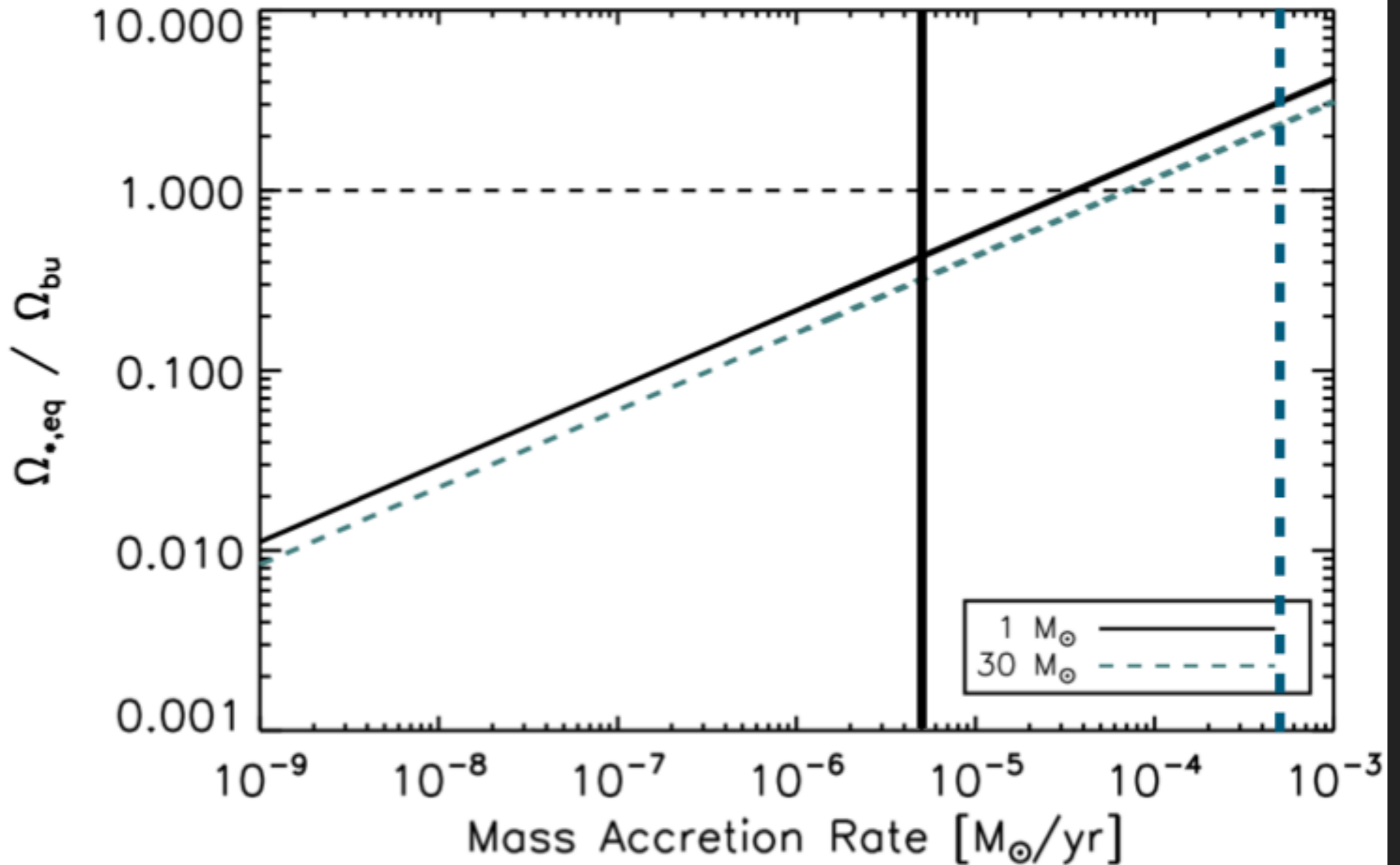
$$B^2 / 8\pi = \rho v^2$$

- ▶ Combining the above:  $R_A/R_* \sim (B_*^4 R_*^5 / \dot{M}^2 M_*)^{1/7}$

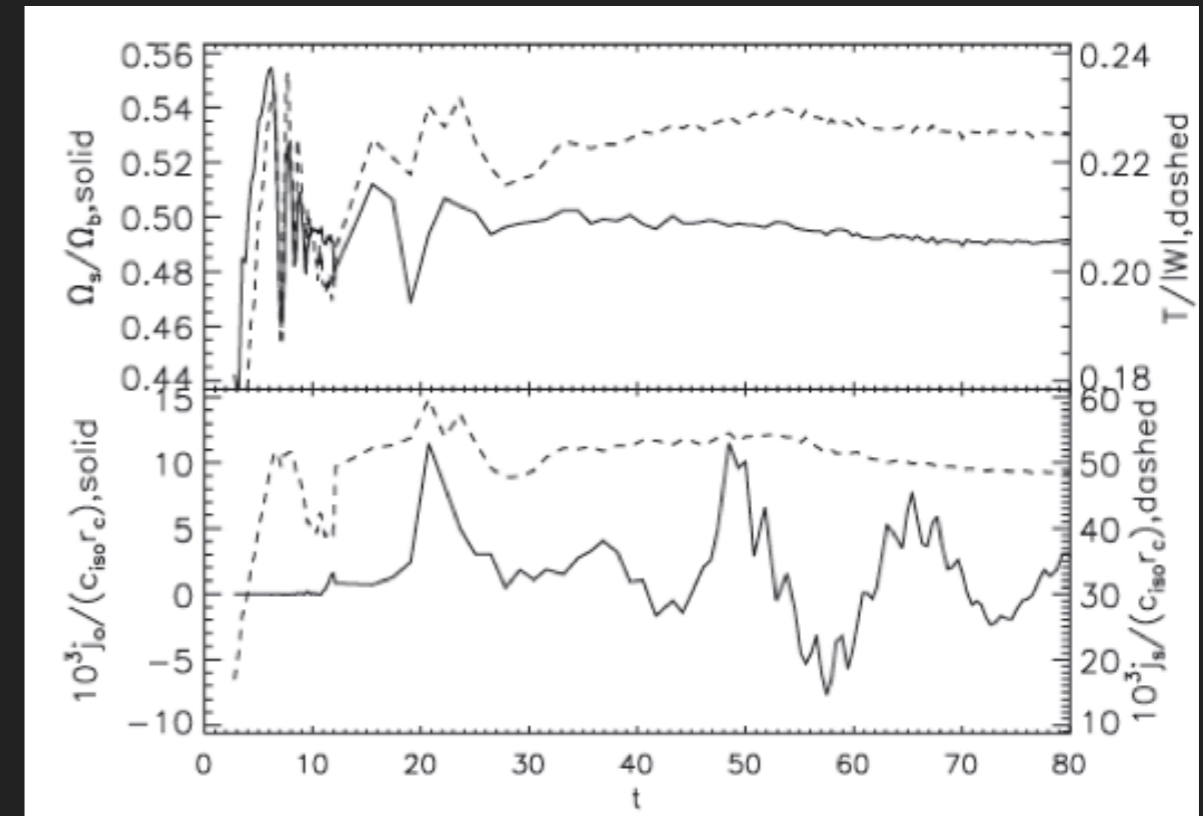
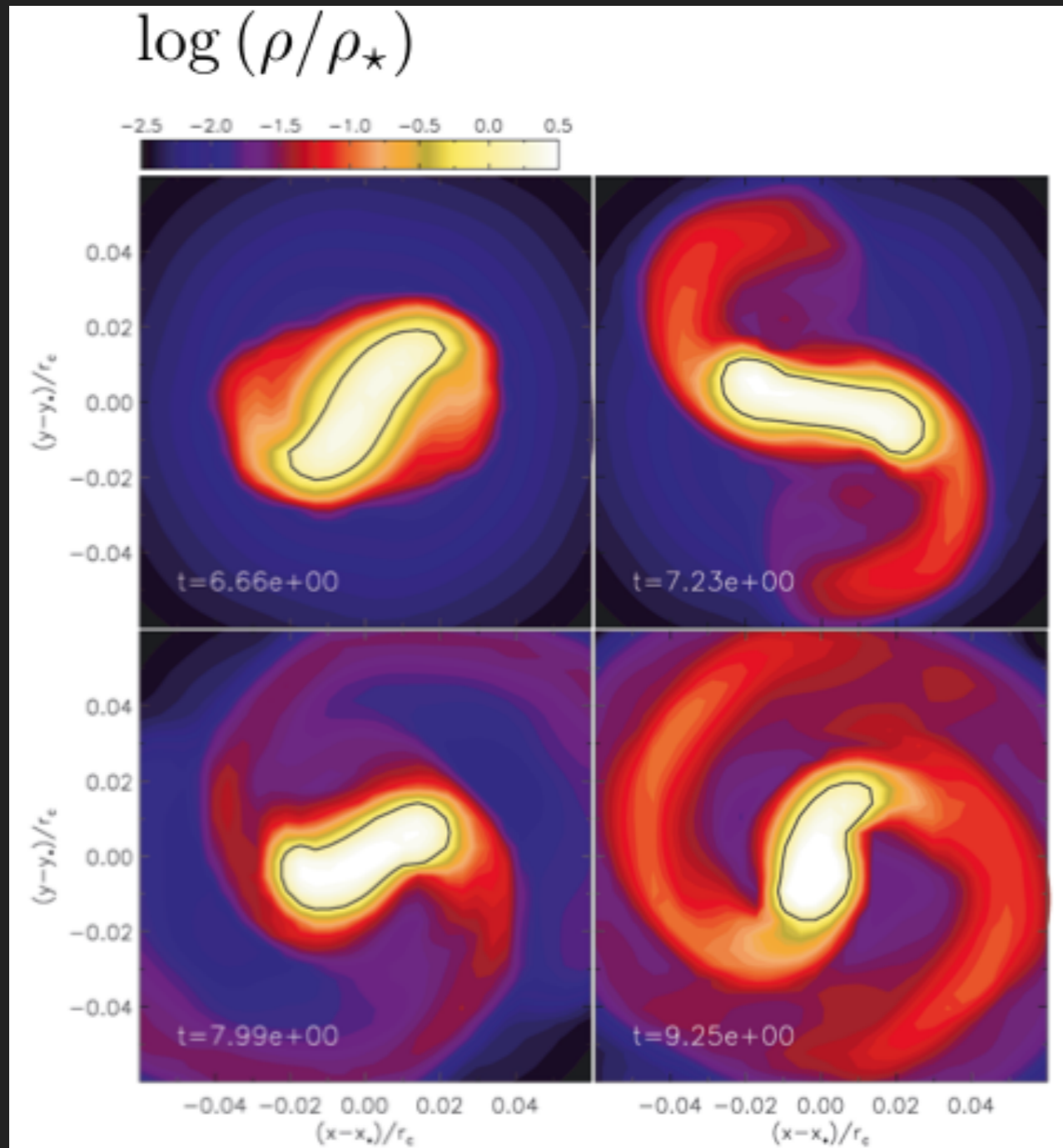
- ▶ Accretion torque  $\tau_A = \dot{M} (GM_* R_A)^{1/2}$

- ▶ Magnetic torque  $\tau_M = (1/3) B_*^2 R_*^6 [R_A^{-3} - 2(R_{CO} R_A)^{-3/2}]$



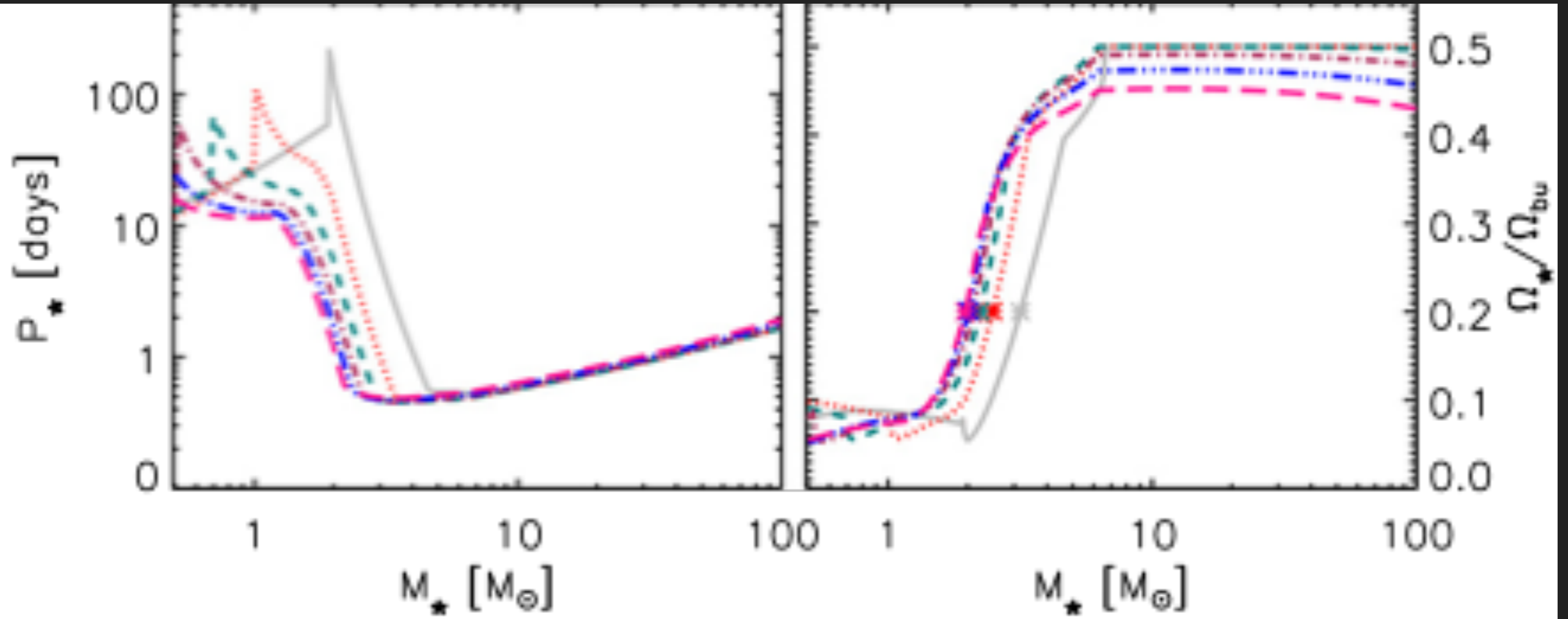


**MAGNETIC BRAKING FAILS**



# GRAVITATIONAL BRAKING





# MAGNETIC + GRAVITATIONAL BRAKING

Rosen+ 2012



# VARIATION WITH ENVIRONMENT

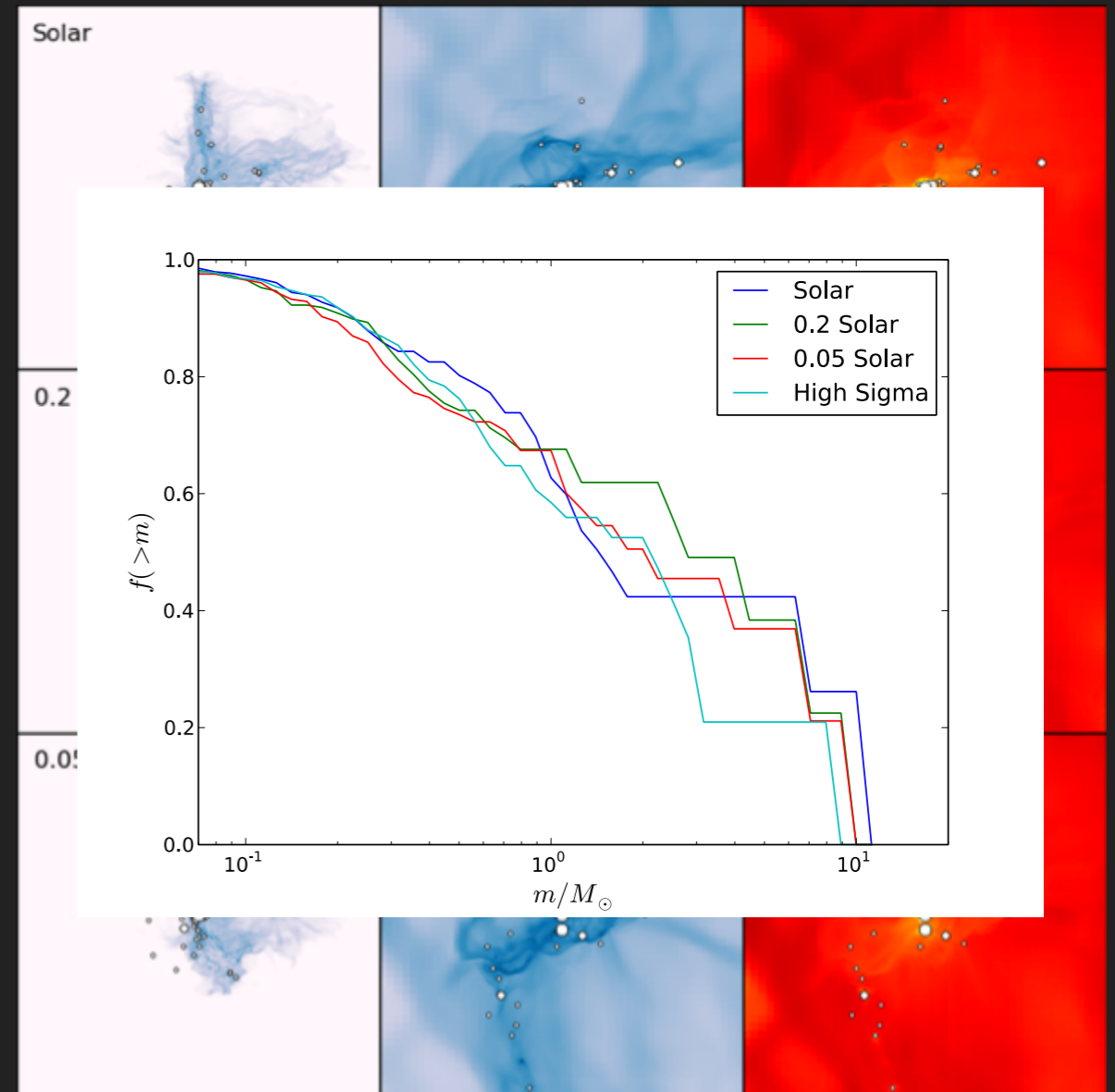
For no reason whatsoever, here are baby platypuses

## POSSIBLE SOURCES OF IMF VARIATION

- ▶ IMF affected by two main factors:
  - ▶ Upper mass limit shaped by radiation pressure
  - ▶ Slope of upper IMF affected by fragmentation
- ▶ Radiation pressure problem appears to be overcome by RRT instability and outflows – seems unlikely that this depends on environment
- ▶ Fragmentation more likely to vary, since this depends on how effectively radiation is able to heat the gas

## DEPENDENCE ON $\Sigma$ AND $Z$

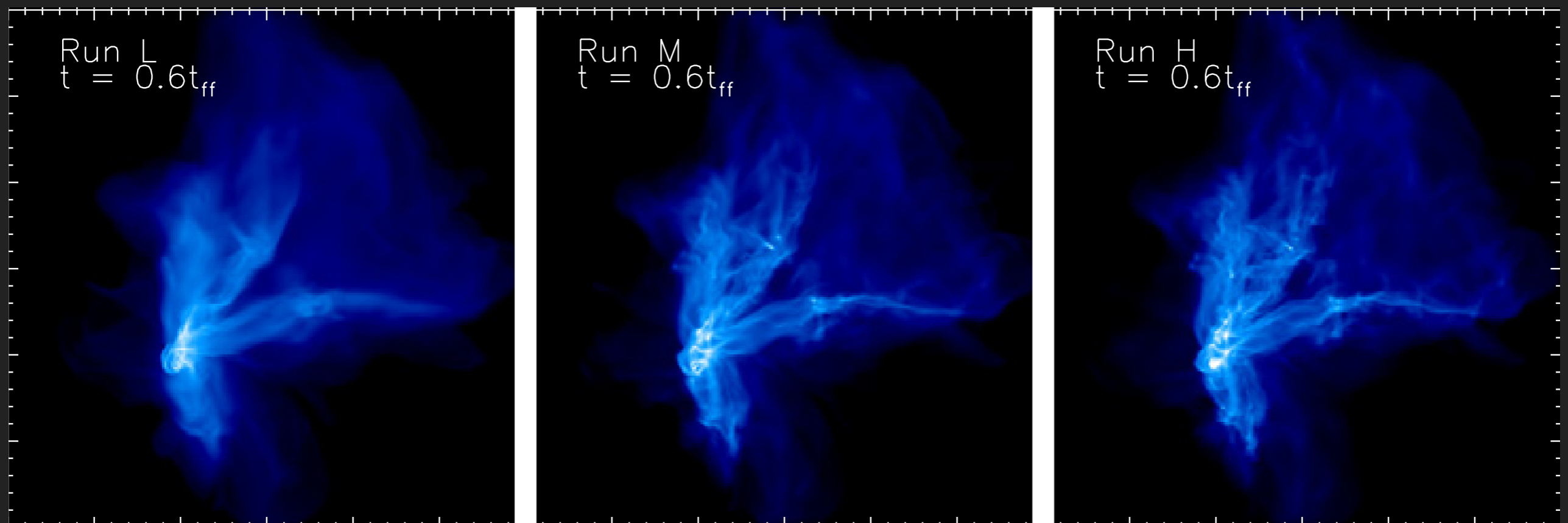
- ▶ Radiation coupled to gas by dust, so metallicity might matter
- ▶ Turns out it doesn't, because at  $\Sigma \sim 1 \text{ g cm}^{-2}$ , even opacity 1% of Milky Way is sufficient to render gas optically thick to stellar photons
- ▶ However,  $\Sigma$  needs to be high enough to trap the radiation
  - ▶ For no B fields, "high enough" is  $\Sigma \sim 1 \text{ g cm}^{-2}$  (Krumholz & McKee 2008)
  - ▶ Value with B fields unknown



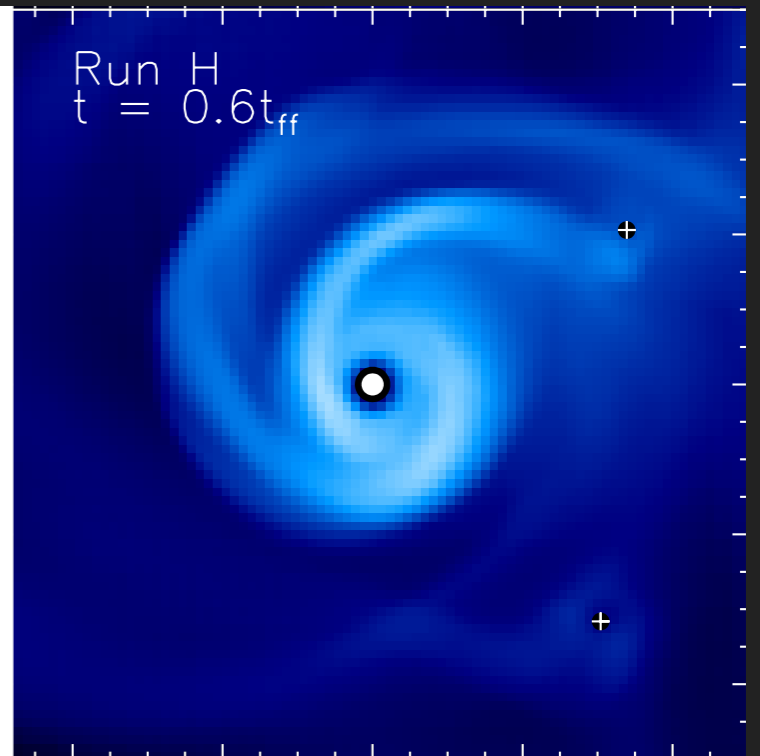
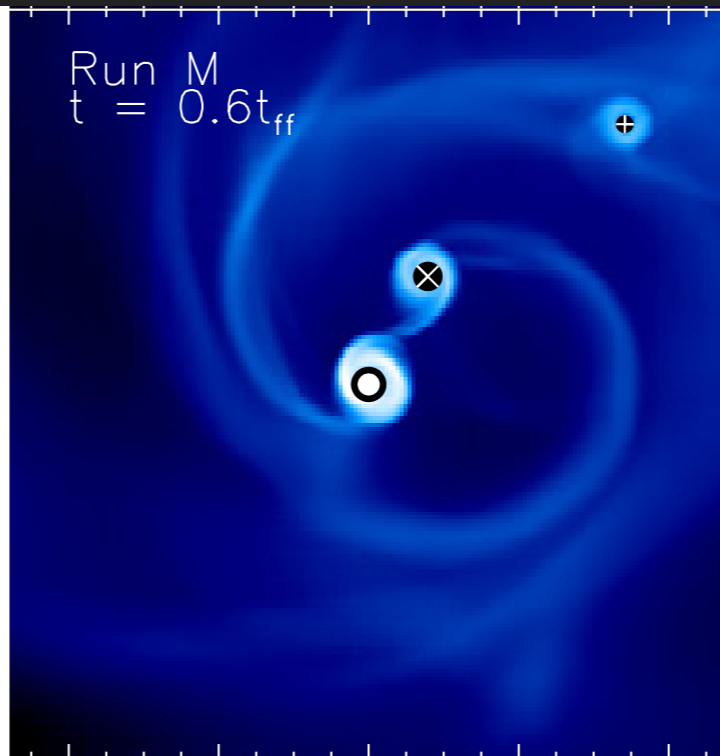
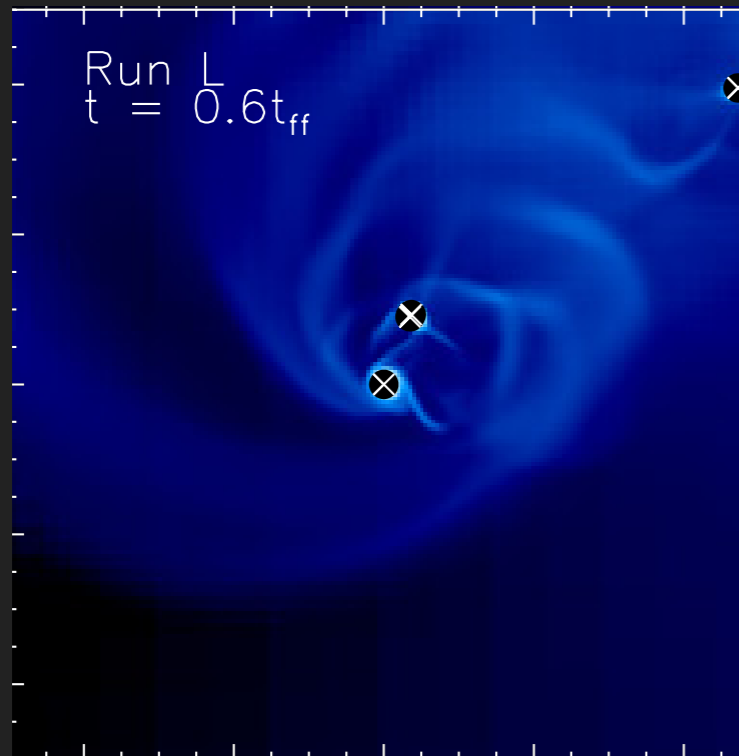
Myers+ 2011

## NUMERICAL EXPERIMENTS: VARIATION WITH $\Sigma$

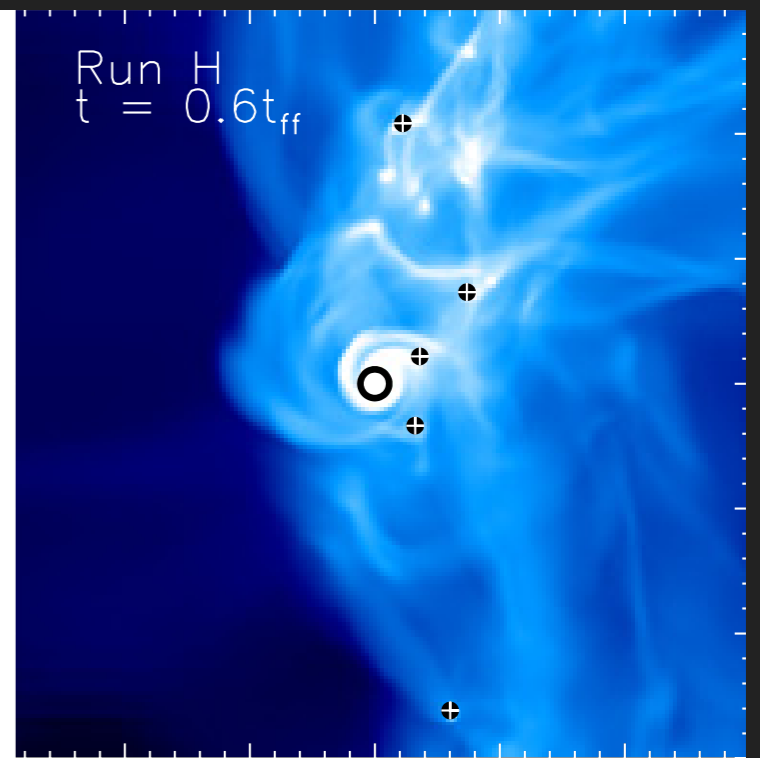
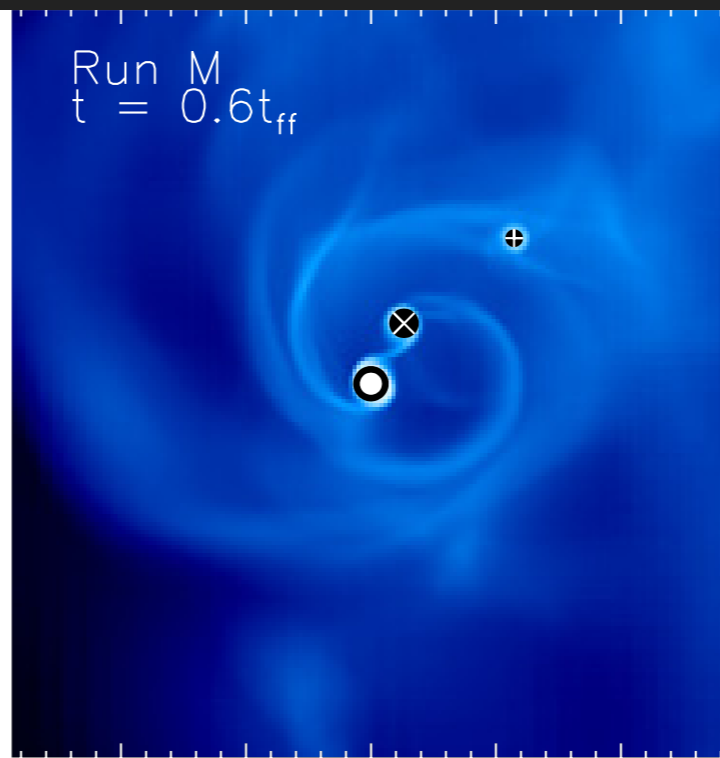
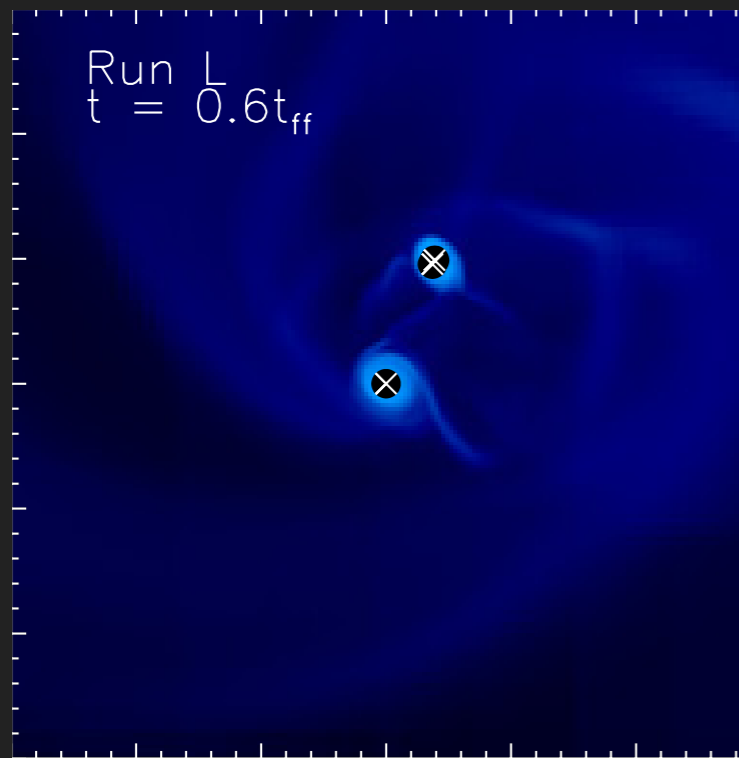
- ▶ Three simulations with identical mass, virial ratio, resolution, velocity field shape, but different column density (Krumholz+ 2010)



Relative scale

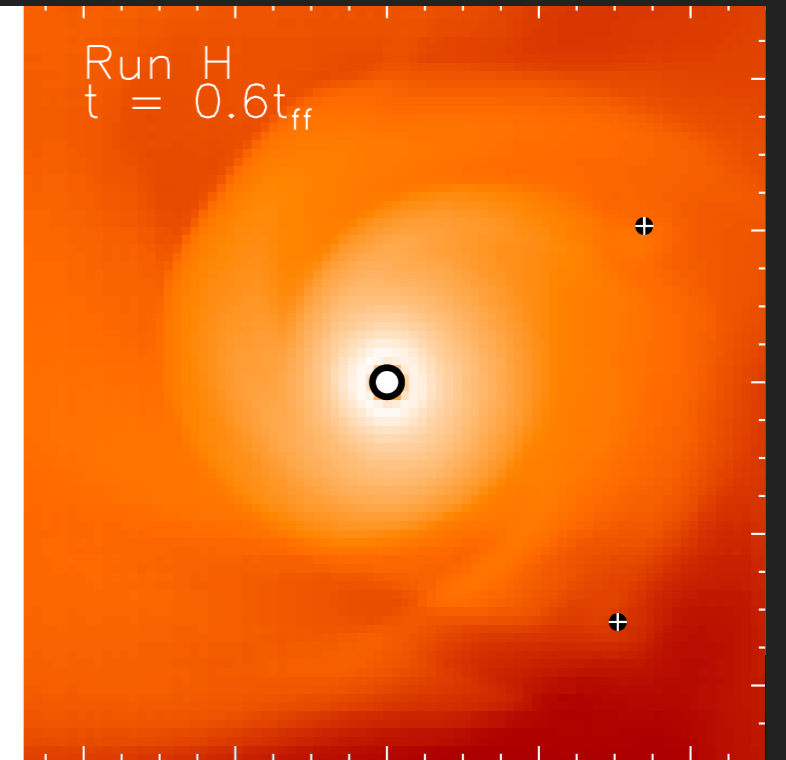
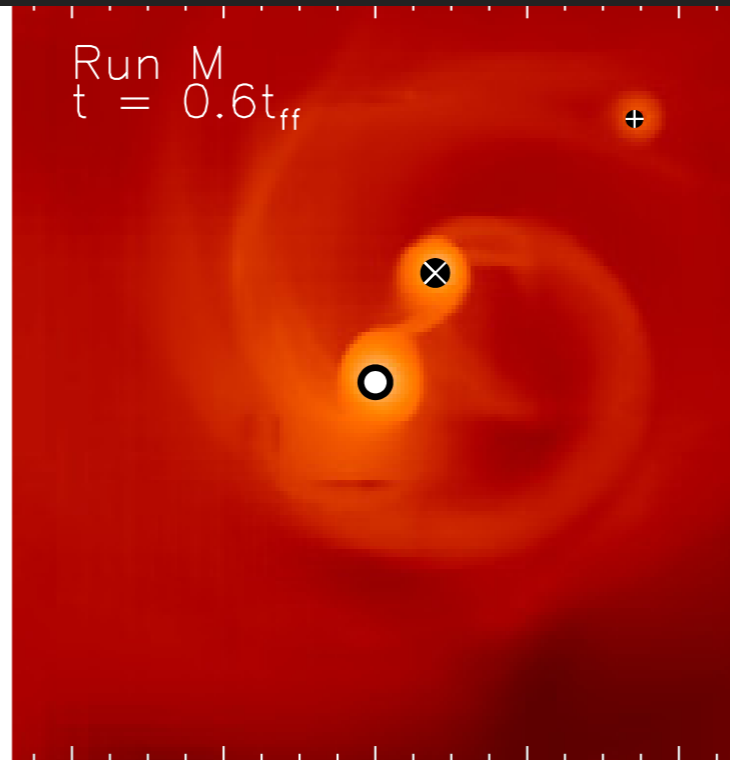
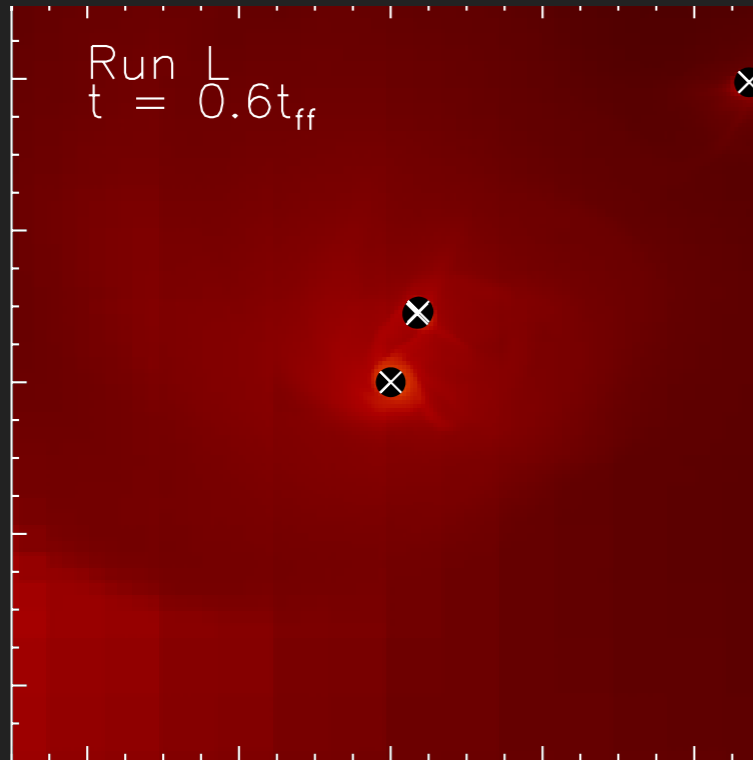


Physical scale

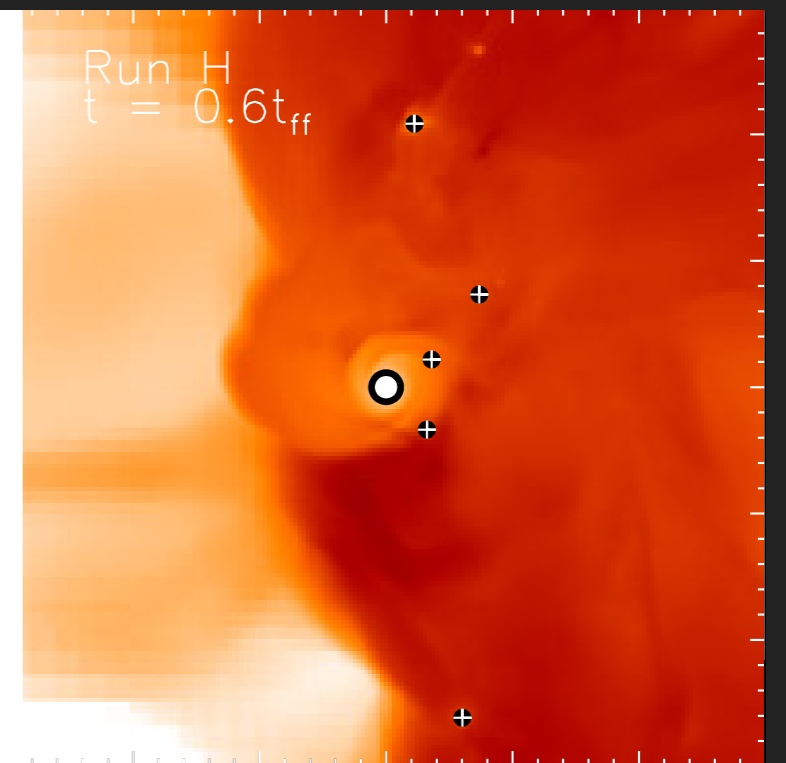
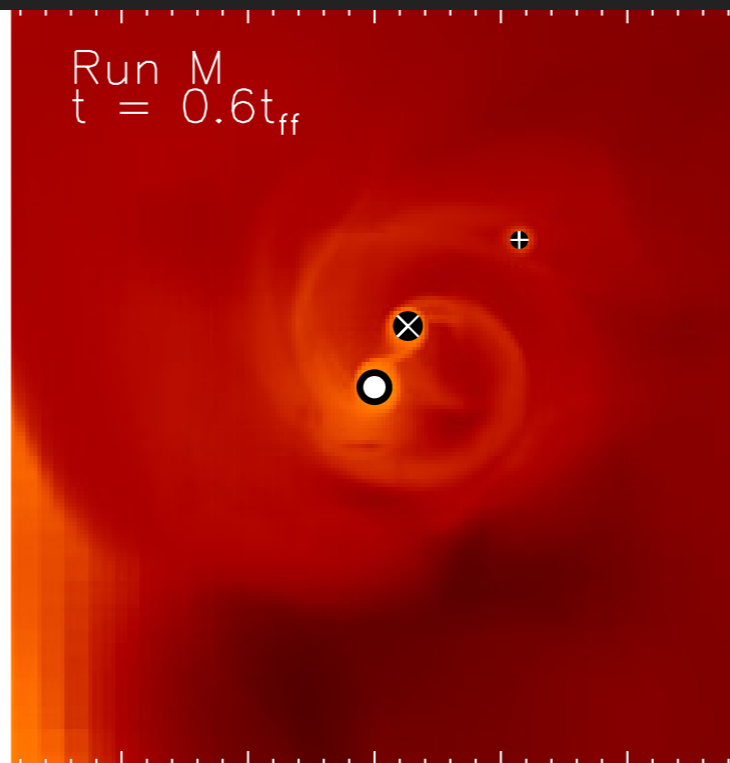
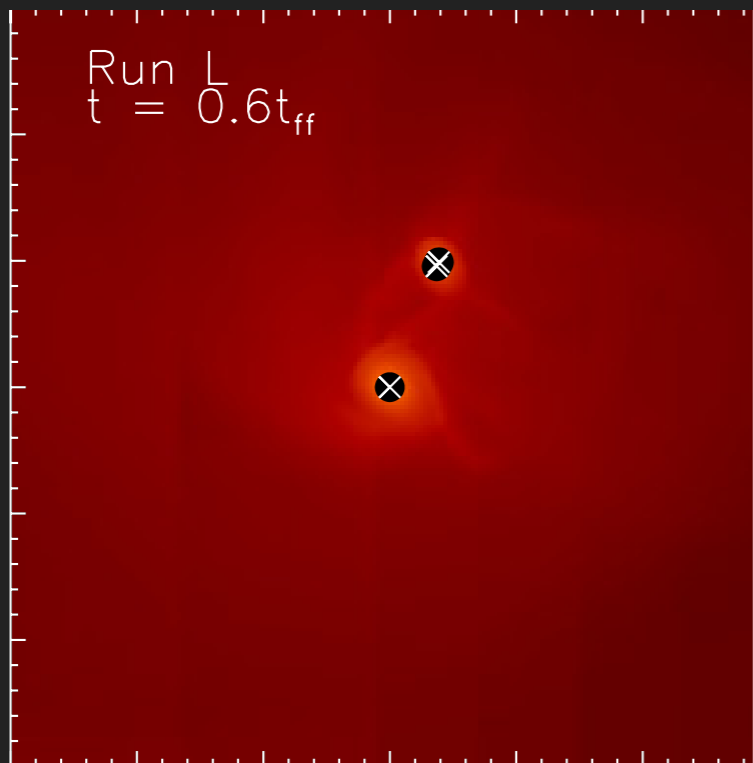


# NUMERICAL EXPERIMENTS WITH $\Sigma$

Relative scale

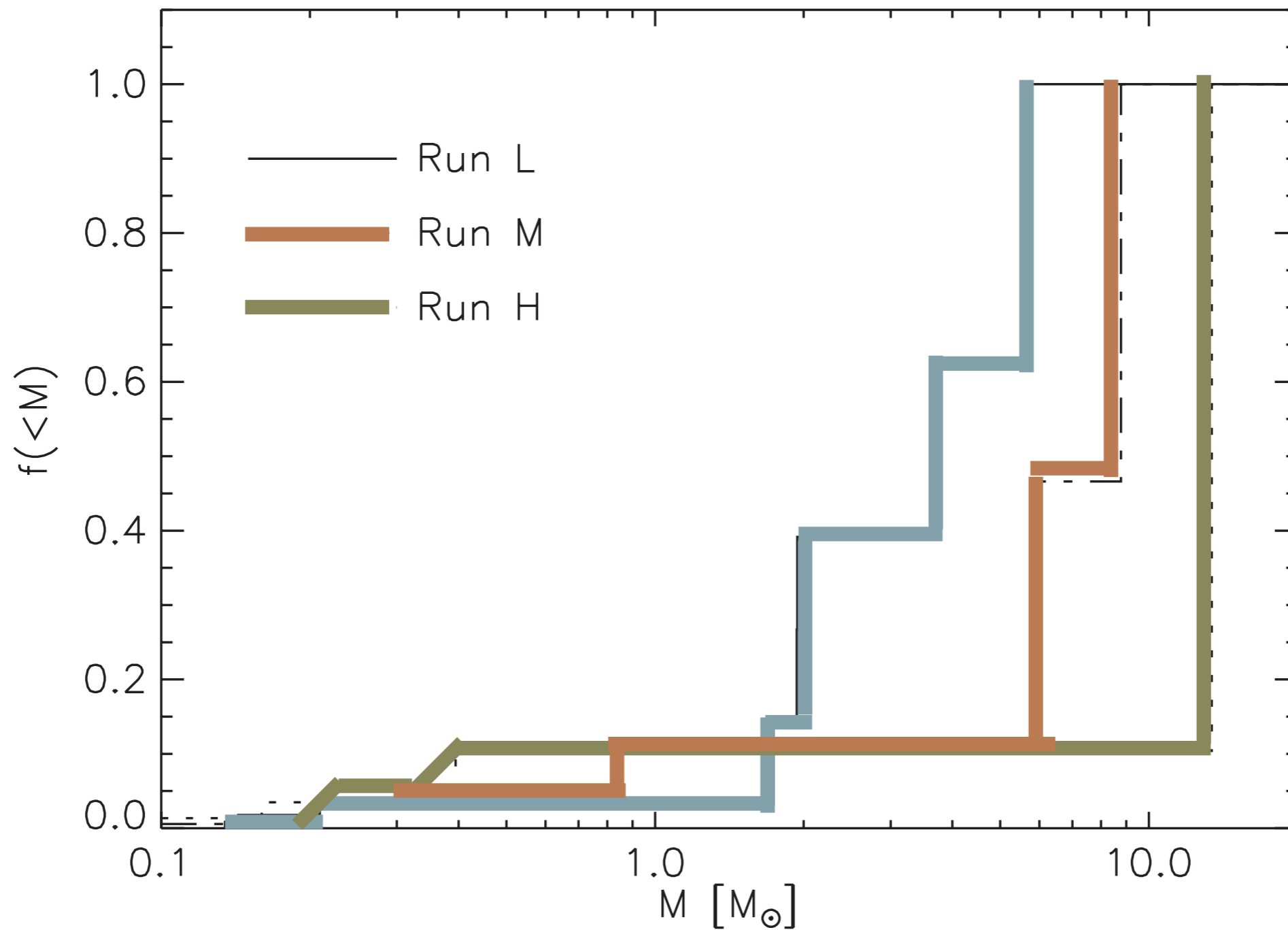


Physical scale



**NUMERICAL EXPERIMENTS WITH  $\Sigma$**

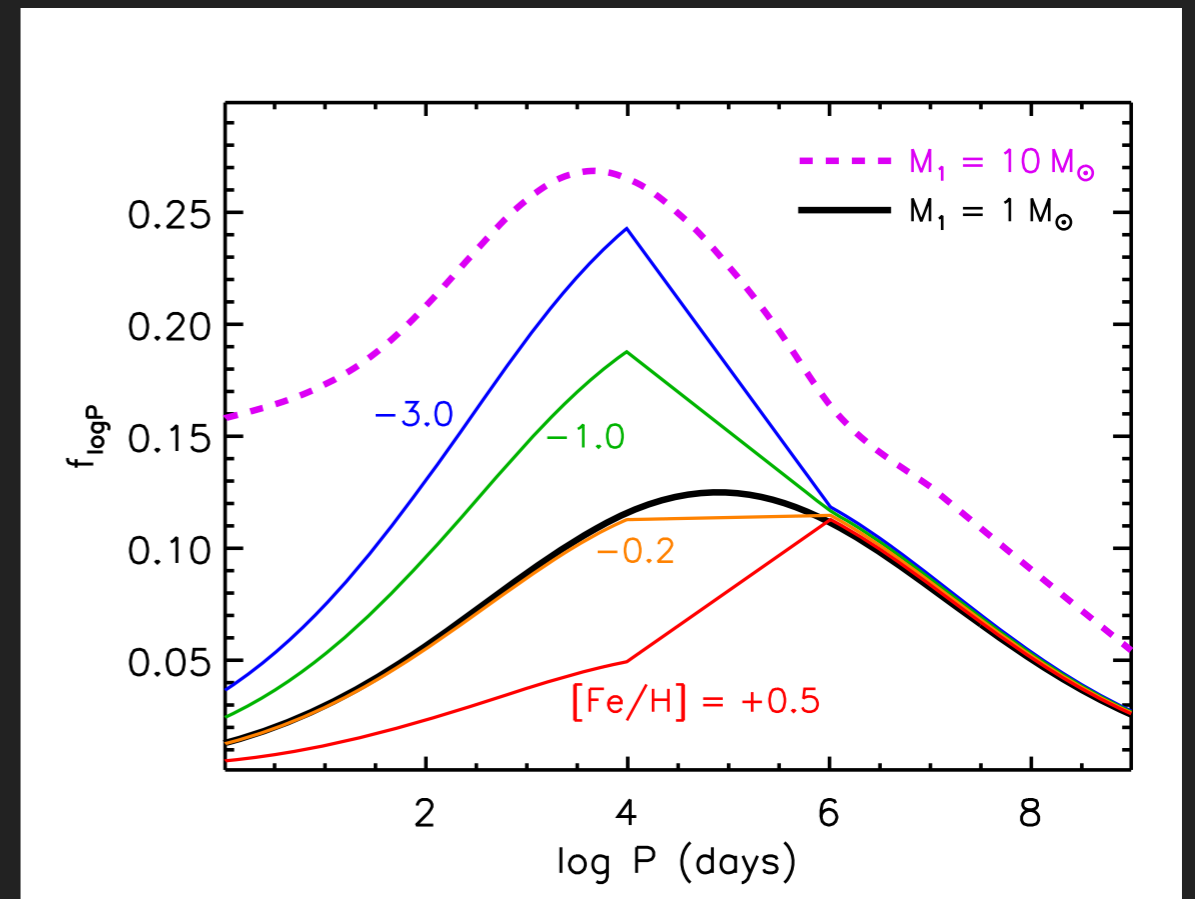
Krumholz+ 2010





# MULTIPLICITY VARIATION

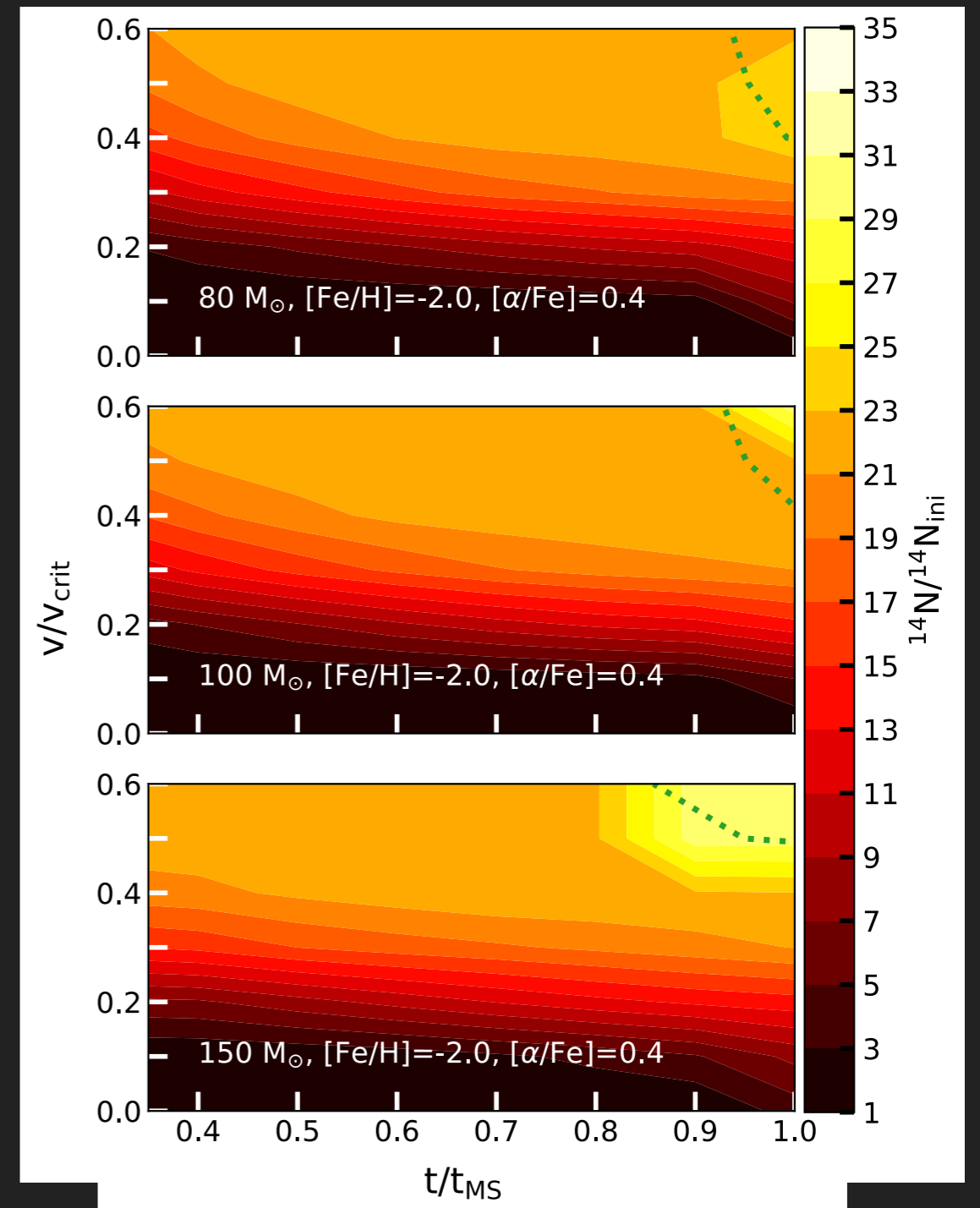
- ▶ Metallicity affects disk ability to absorb starlight and cool; big effect for solar-type stars
- ▶ For massive stars with high accretion rate, disk fragmentation unavailable even at  $Z = Z_{\odot}$ , so probably no big change with  $Z$
- ▶ Changing  $\Sigma$  may affect accretion rate and thus fragmentation, but there have been no studies to date



Moe+ 2019

## VARIATION IN ROTATION

- ▶ Gravitational torques independent of environment, so birth rotation distribution probably constant
- ▶ Post-birth spin down via winds depends on wind strength and therefore on metallicity
- ▶ Main effect: longer rapid rotation for low  $Z$  stars
- ▶ Probably affects WNL star frequency, maybe other things





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# IMPLICATIONS AND QUESTIONS

For no reason whatsoever, here is  
a baby sugar glider

### WHAT WE UNDERSTAND (MAYBE)

- ▶ Formation of massive stars controlled mostly by fragmentation; feedback probably ineffective
  - ▶ Fragmentation does not depend on  $Z$ , but does depend  $\Sigma$  – possibly less fragmentation at high  $\Sigma$
- ▶ All massive stars are multiples due to disk fragmentation; no obvious reason this should vary with environment
- ▶ Massive stars born rotating fast in all environments, but spin-down weaker at lower metallicity

### WHAT WE DEFINITELY DON'T UNDERSTAND

- ▶ How do any potential variations in demographics at birth interact with subsequent evolution (particularly in binaries) to affect rate of compact object production?
- ▶ How does variation with microscopic environment (for example local surface density) change the demographics averaged over galactic or cosmological scales? For example, do we have more high- $\Sigma$  environments, and thus more massive stars, at higher redshift?