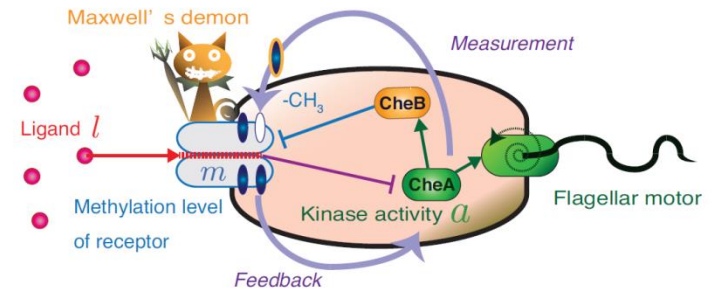
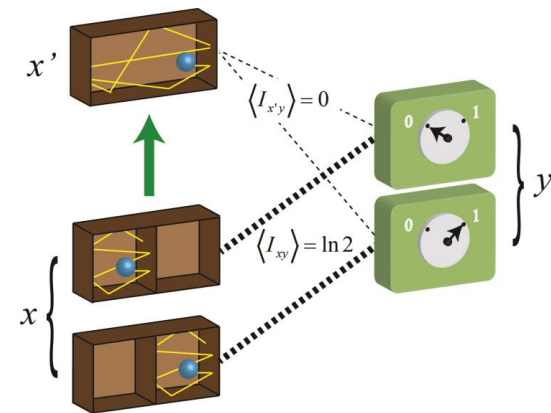


Maxwell's Demon in Biochemical Signal Transduction



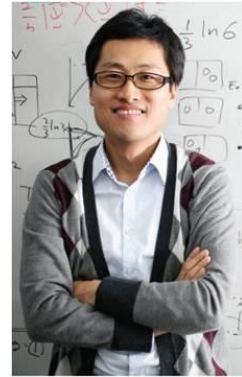
Takahiro Sagawa

Department of Applied Physics, University of Tokyo

New Frontiers in Non-equilibrium Physics 2015
28 July 2015, YITP, Kyoto

Collaborators on Information Thermodynamics

- Masahito Ueda (Univ. Tokyo)
- Shoichi Toyabe (Tohoku Univ.)
- Eiro Muneyuki (Chuo Univ.)
- Masaki Sano (Univ. Tokyo)



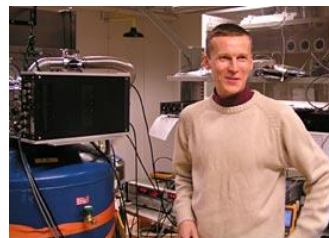
- Sosuke Ito (Titech)

- Naoto Shiraishi (Univ. Tokyo)
- Sang Wook Kim (Pusan National Univ.)
- Jung Jun Park (National Univ. Singapore)
- Kang-Hwan Kim (KAIST)
- Simone De Liberato (Univ. Paris VII)



- Juan M. R. Parrondo (Univ. Madrid)
- Jordan M. Horowitz (MIT)

- Jukka Pekola (Aalto Univ.)
- Jonne Koski (Aalto Univ.)
- Ville Maisi (Aalto Univ.)



Outline

- Introduction
 - Information and entropy
 - Information thermodynamics: a general framework
 - Paradox of Maxwell's demon
- Review of previous results**
- Thermodynamics of autonomous information processing
 - Application to biochemical signal transduction
- Today's main part!**
- Summary

Outline

- **Introduction**
- Information and entropy
- Information thermodynamics: a general framework
- Paradox of Maxwell's demon

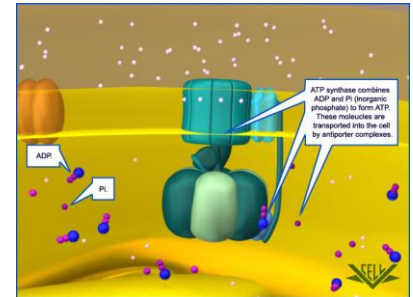
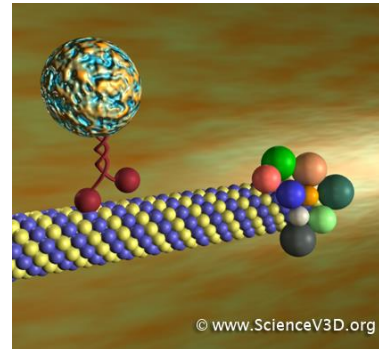
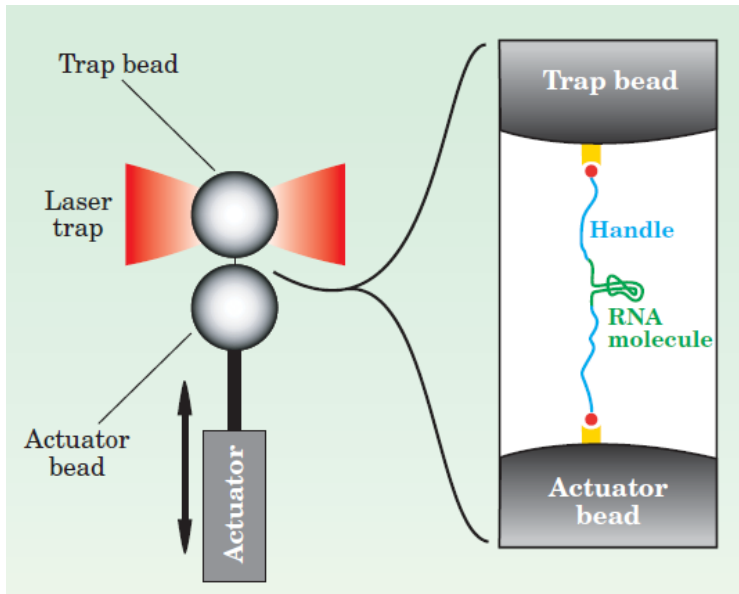
- Thermodynamics of autonomous information processing
- Application to biochemical signal transduction

- Summary

Thermodynamics in the Fluctuating World

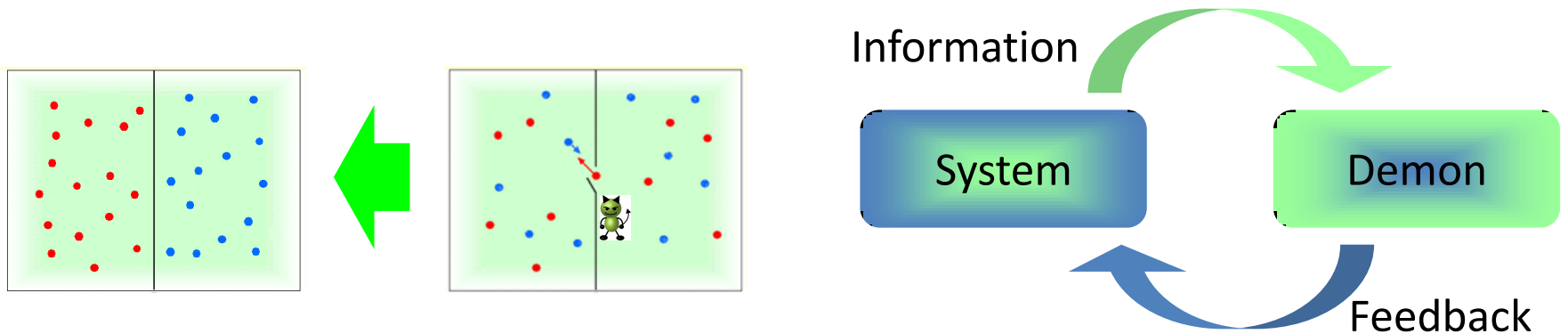
Thermodynamics of small systems with large heat bath(s)

➔ Thermodynamic quantities are fluctuating!



- ✓ Second law $\langle W \rangle \geq \Delta F$
- ✓ Nonlinear & nonequilibrium relations

Information Thermodynamics



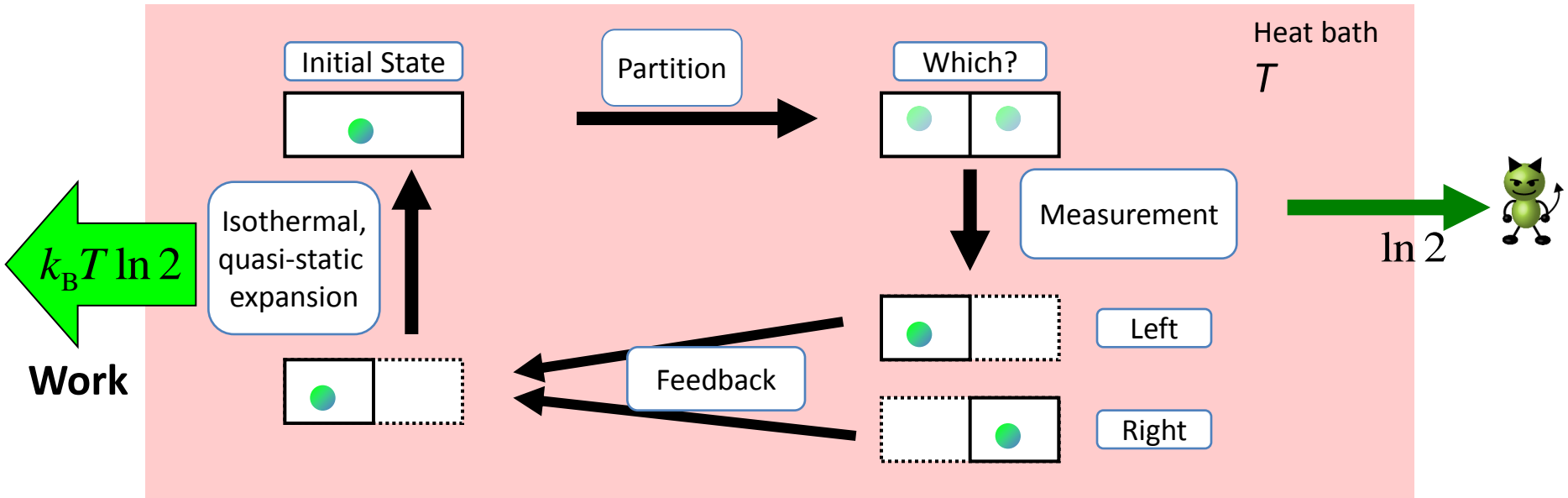
Information processing at the level of thermal fluctuations



- ✓ Foundation of the second law of thermodynamics
- ✓ Application to nanomachines and nanodevices

Review: J. M. R. Parrondo, J. M. Horowitz, & T. Sagawa, *Nature Physics* **11**, 131-139 (2015).

Szilard Engine (1929)



Free energy: $F = E - TS$

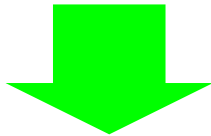
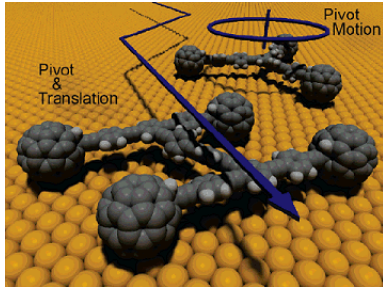
Increase F

Decrease by feedback TS

Can control physical entropy by using information

Information Heat Engine

System
(Working engine)



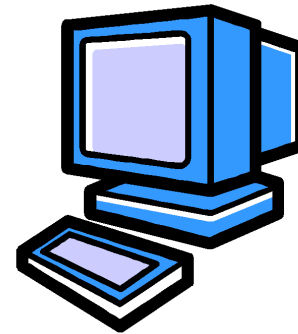
Free energy / work

Information



Feedback

Memory
(Controller)



Entropic cost

- ✓ Can increase the system's free energy even if there is no energy flow between the system and the controller

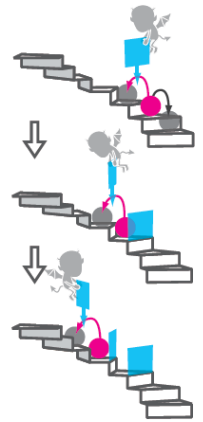
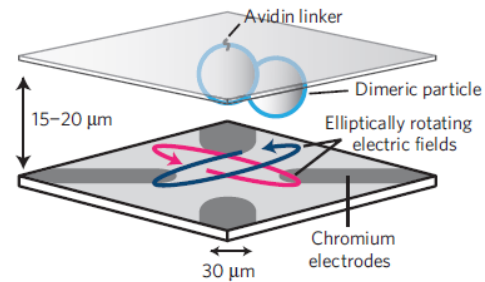
Experimental Realizations

- With a colloidal particle

Toyabe, TS, Ueda, Muneyuki, & Sano, Nature Physics (2010)

Efficiency: 30%

Validation of $\langle e^{-\beta(W-\Delta F)} \rangle = \gamma$

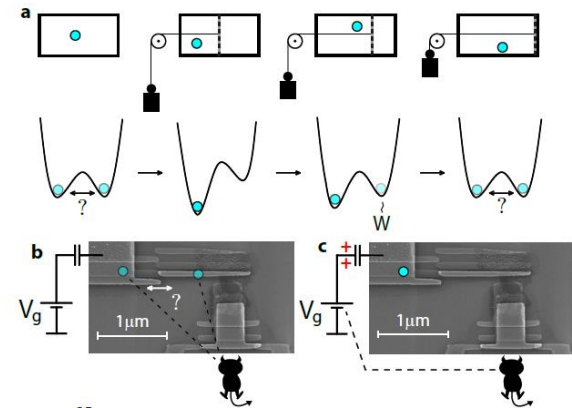


- With a single electron

Koski, Maisi, TS, & Pekola, PRL (2014)

Efficiency: 75%

Validation of $\langle e^{-\beta(W-\Delta F)-I} \rangle = 1$



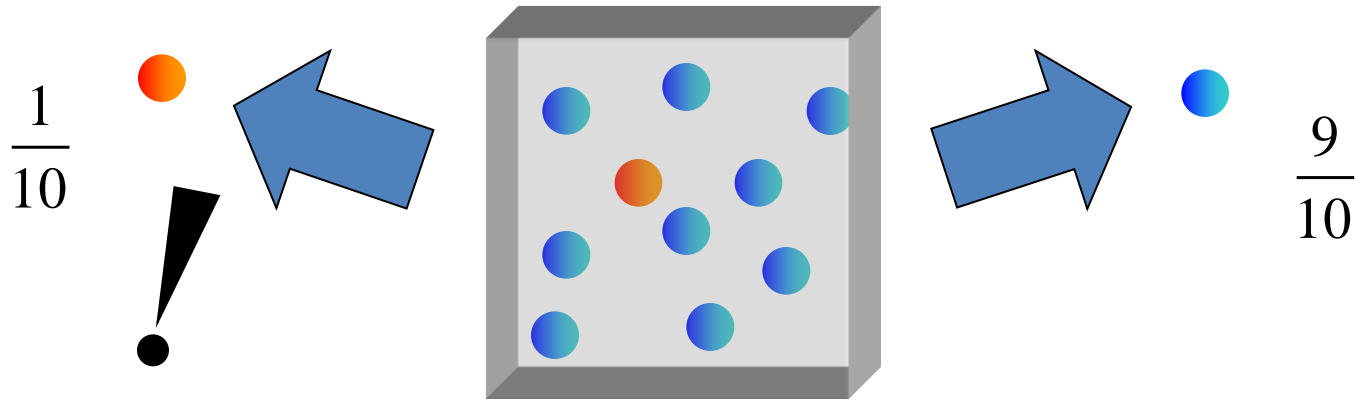
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Shannon Information



Information content with event k : $\ln \frac{1}{p_k}$

Average 

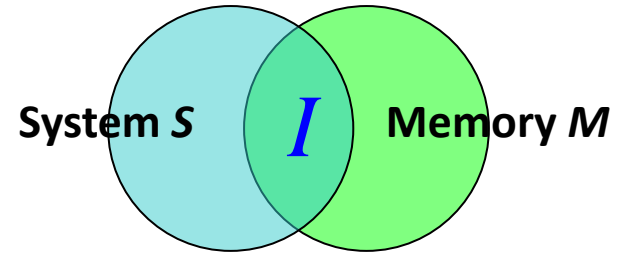
Shannon information: $H = \sum_k p_k \ln \frac{1}{p_k}$

Mutual Information

System S



Memory M
(measurement device)



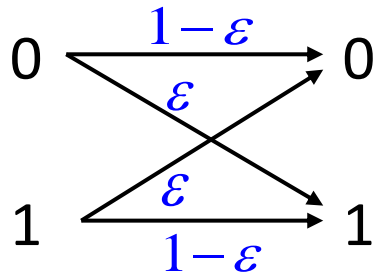
Measurement with stochastic errors

$$I(S : M) \equiv H(S) + H(M) - H(SM)$$

$$0 \leq I \leq H(M)$$

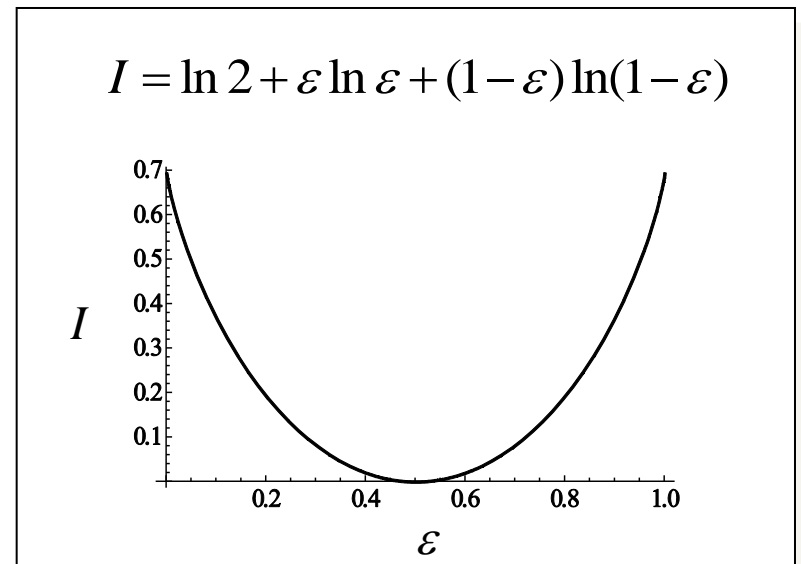
No information

No error



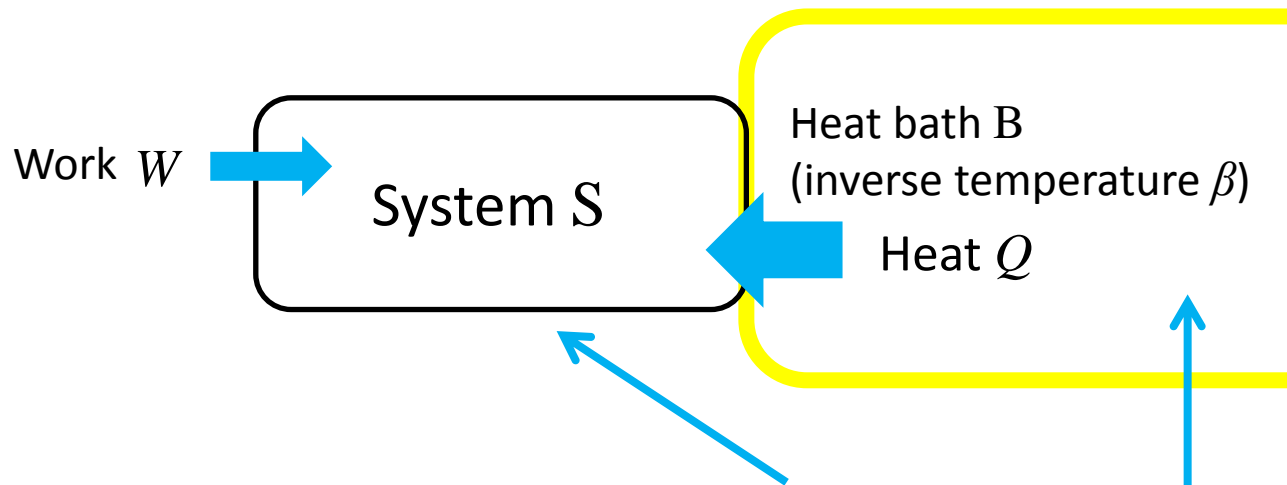
Ex. Binary symmetric channel

Correlation between S and M



Entropy Production

Stochastic dynamics of system S (e.g., Langevin system)



Entropy production
in the total system:

$$\Delta S_{SB} \equiv \Delta S_S - \beta \langle Q \rangle$$

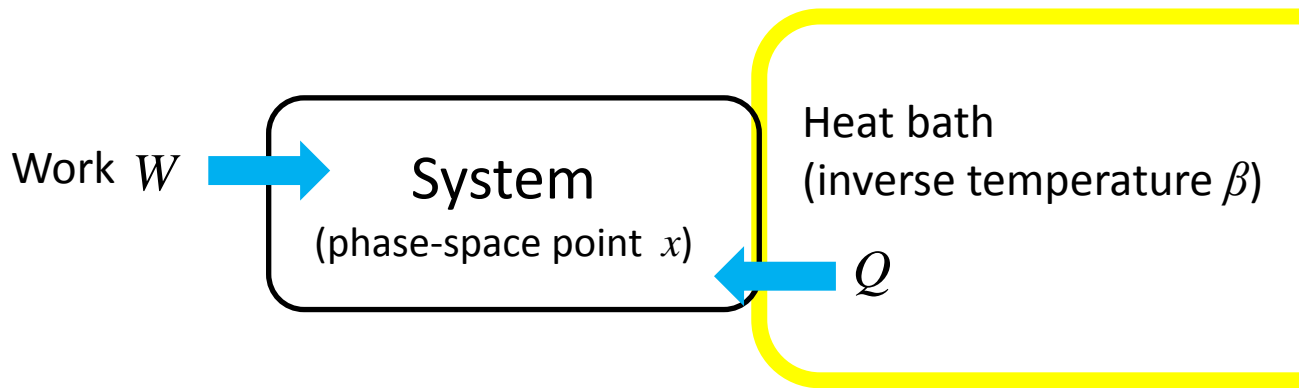
Change in
the Shannon
entropy of S

Averaged heat
absorbed by S

If the initial and the final states are canonical distributions: $\Delta S_{SB} = \beta (\langle W \rangle - \Delta F)$

Free-energy difference

Stochastic Entropy Production



Stochastic entropy production along a trajectory of the system from time 0 to τ

$$\Delta s_{\text{SB}} \equiv \Delta s_{\text{S}} - \beta Q$$

$$\Delta s_{\text{S}} \equiv s_{\text{S}}[x(\tau), \tau] - s_{\text{S}}[x(0), 0] \quad s_{\text{S}}[x, t] \equiv -\ln P[x, t]$$

$$\langle \Delta s_{\text{S}} \rangle = \Delta S_{\text{S}}$$

$P[x, t]$: probability distribution at time t

If the initial and the final states are canonical distributions: $\Delta s_{\text{SB}} = \beta(W - \Delta F)$

Integral Fluctuation Theorem and Jarzynski Equality

Integral fluctuation theorem

$$\left\langle e^{-\Delta S_{\text{SB}}} \right\rangle = 1$$

Seifert, PRL (2005), ...

for any initial and final distributions

Second law can be expressed by an **equality** with full cumulants



The second law of thermodynamics (Clausius inequality)

$$\left\langle \Delta S_{\text{SB}} \right\rangle \geq 0$$



$$\Delta S_s \geq \beta \langle Q \rangle$$

Jarzynski equality

Jarzynski, PRL (1997)

$$\Delta S_{\text{SB}} = \beta(W - \Delta F)$$



$$\left\langle e^{-\beta W} \right\rangle = e^{-\beta \Delta F}$$



$$\langle W \rangle \geq \Delta F$$

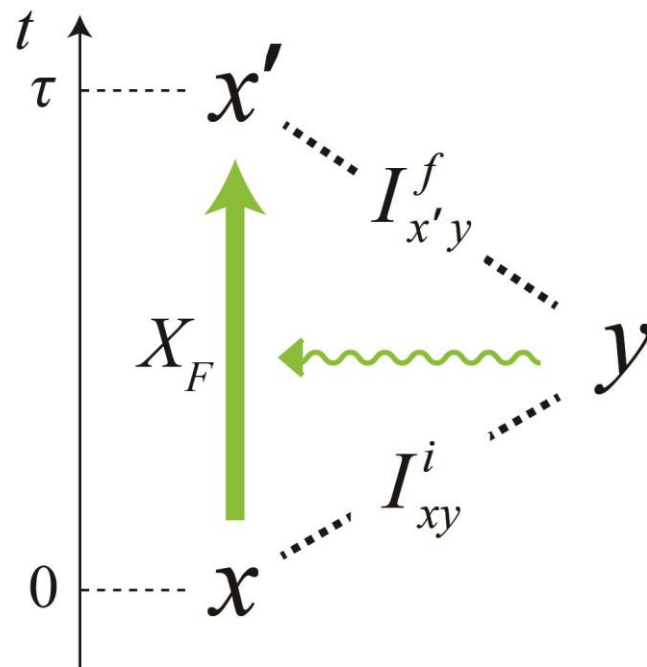
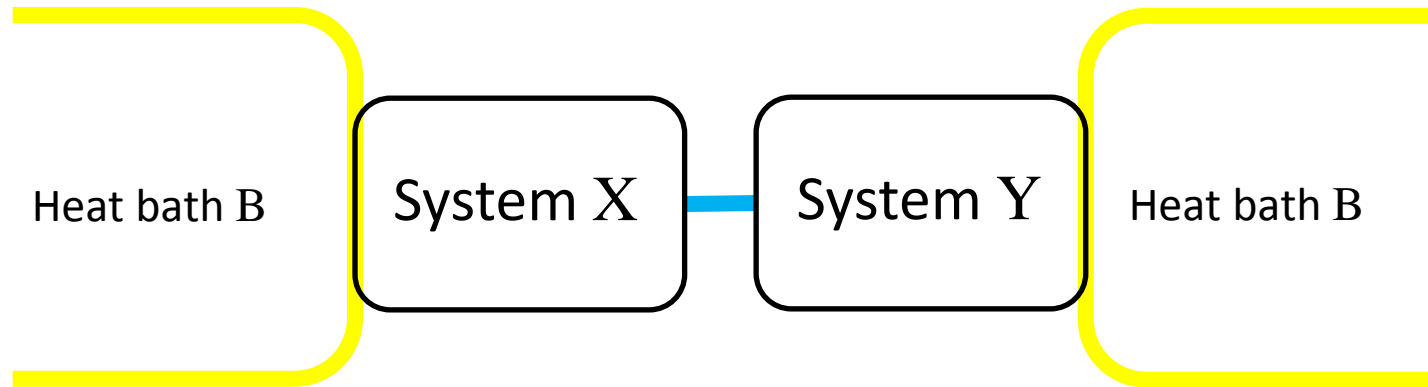
Outline

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Setup

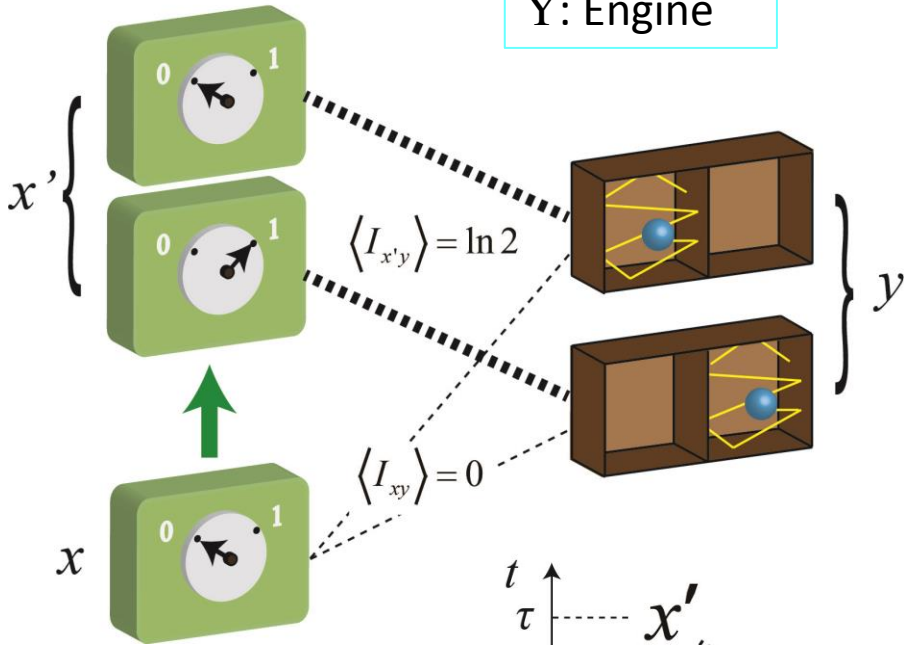


Time evolution of X
under the influence of Y with
initial and final correlations

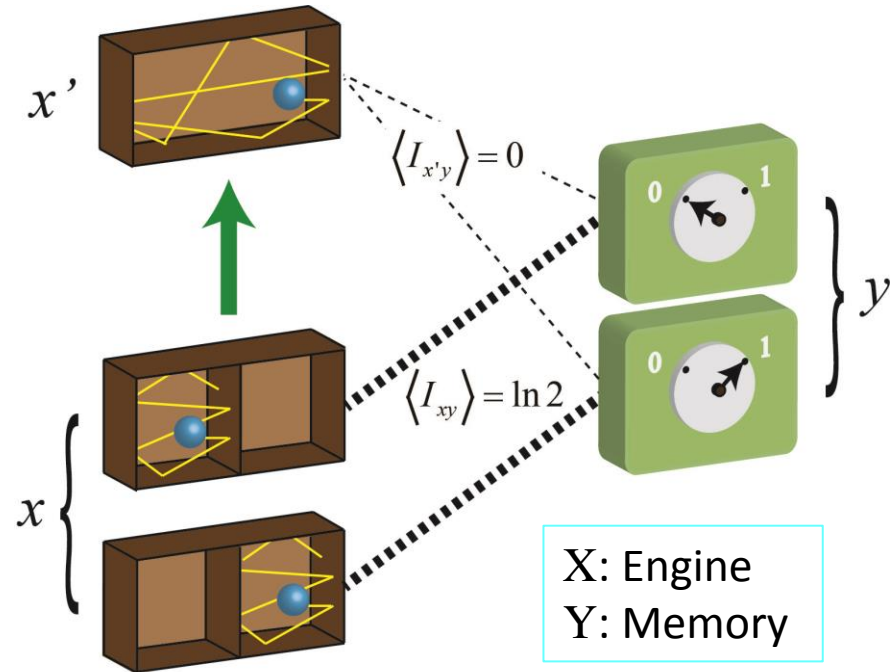
Special Cases: Measurement and Feedback

Measurement

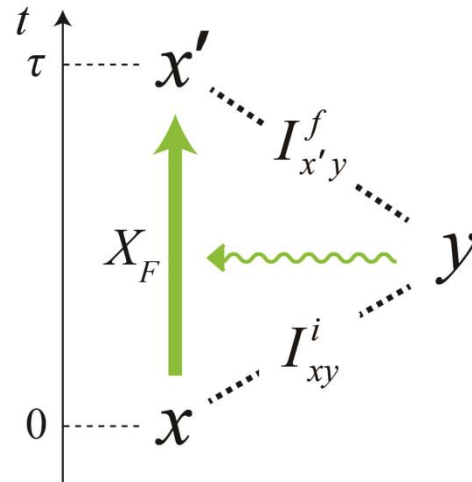
X: Memory
Y: Engine



Feedback



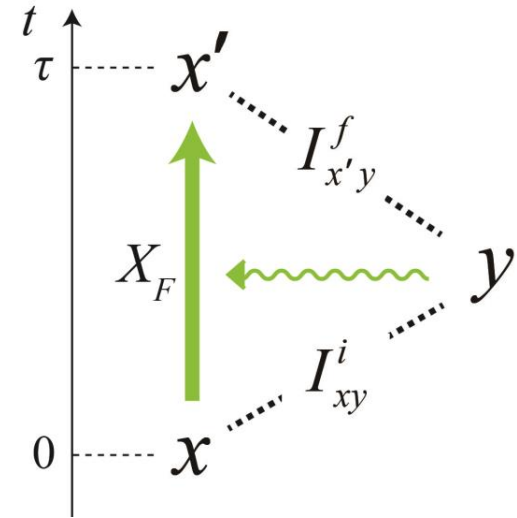
X: Engine
Y: Memory



Stochastic Entropy and Mutual Information

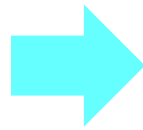
Entropy increase in XB

$$\Delta s_{\text{XB}} \equiv \Delta s_{\text{X}} - \beta Q_{\text{X}}$$



Initial correlation

$$I_{xy}^i \equiv \ln \frac{P[x, y]}{P[x]P[y]}$$



$$\langle I_{xy}^i \rangle = \int dx dy P[x, y] \ln \frac{P[x, y]}{P[x]P[y]}$$

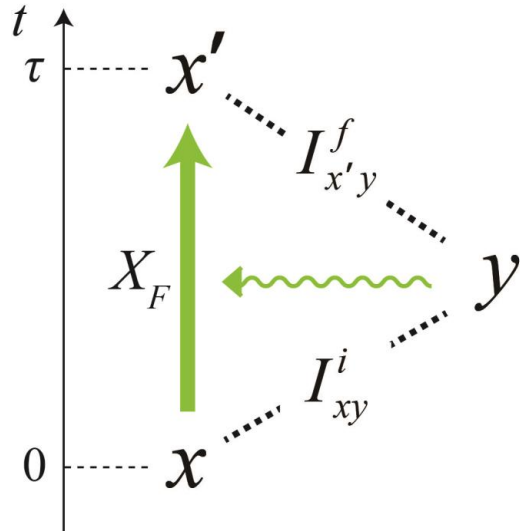
Final correlation

$$I_{x'y}^f \equiv \ln \frac{P[x', y]}{P[x']P[y]}$$



$$\langle I_{x'y}^f \rangle = \int dx' dy P[x', y] \ln \frac{P[x', y]}{P[x']P[y]}$$

Decomposition of Entropy Production



Total entropy production in XYB

$$\Delta S_{XYB} \equiv \Delta S_{XB} - \Delta I$$

$$\Delta S_{XYB} \equiv \Delta S_{XY} - \beta Q_X$$

$$\Delta S_{XB} \equiv \Delta S_X - \beta Q_X$$

$$\Delta I \equiv I_{x'y}^f - I_{xy}^i$$

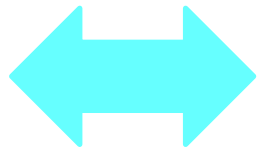
$$\begin{aligned} \Delta S_{XY} &= \Delta S_X + \cancel{\Delta S_Y} - \Delta I \\ &= \Delta S_X - \Delta I \end{aligned}$$

Fluctuation Theorem

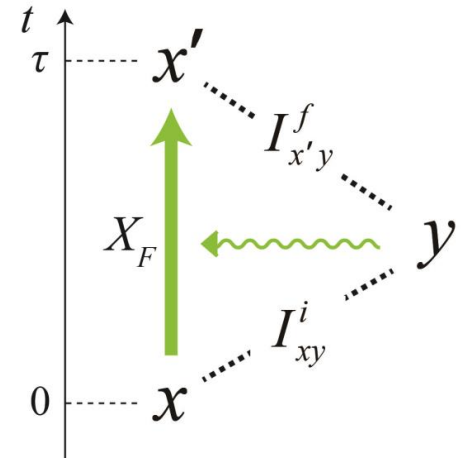
Integral fluctuation theorem

$$\Delta S_{\text{XYB}} = \Delta S_{\text{XB}} - \Delta I$$

$$\left\langle e^{-\Delta S_{\text{XYB}}} \right\rangle = 1$$

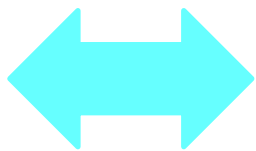


$$\left\langle e^{-\Delta S_{\text{XB}} + \Delta I} \right\rangle = 1$$



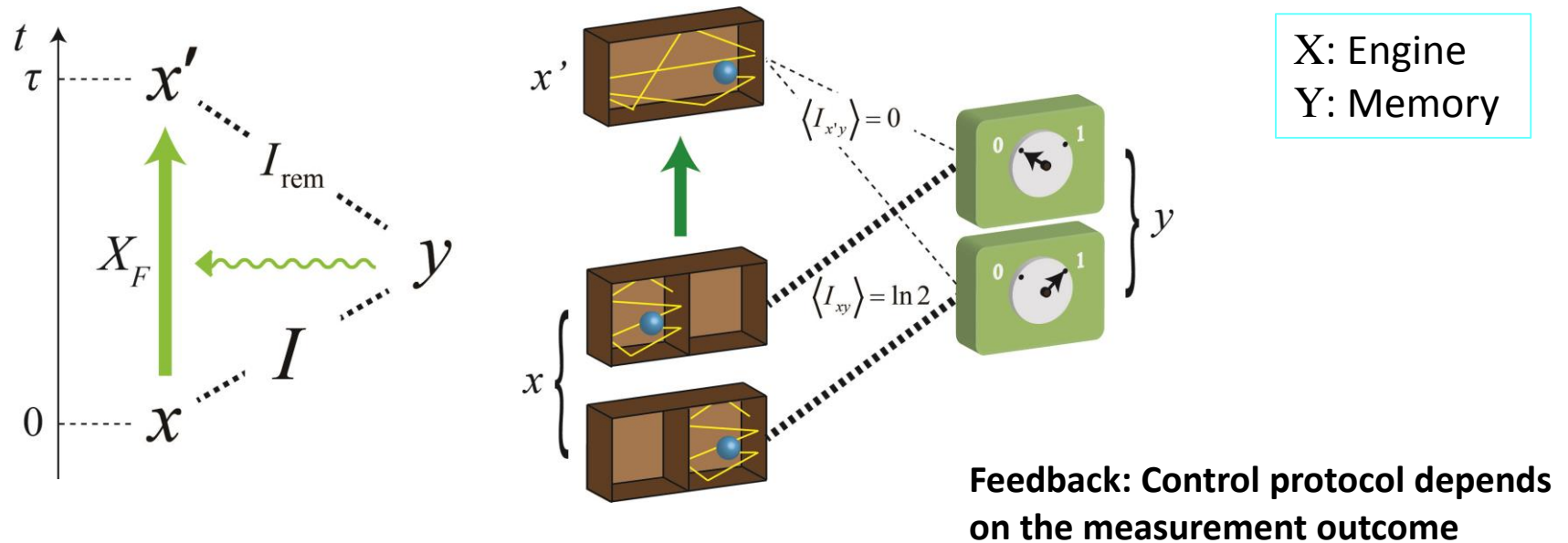
Second law

$$\left\langle \Delta S_{\text{XYB}} \right\rangle \geq 0$$



$$\left\langle \Delta S_{\text{XB}} \right\rangle \geq \left\langle \Delta I \right\rangle$$

Special Case 1: Feedback Control

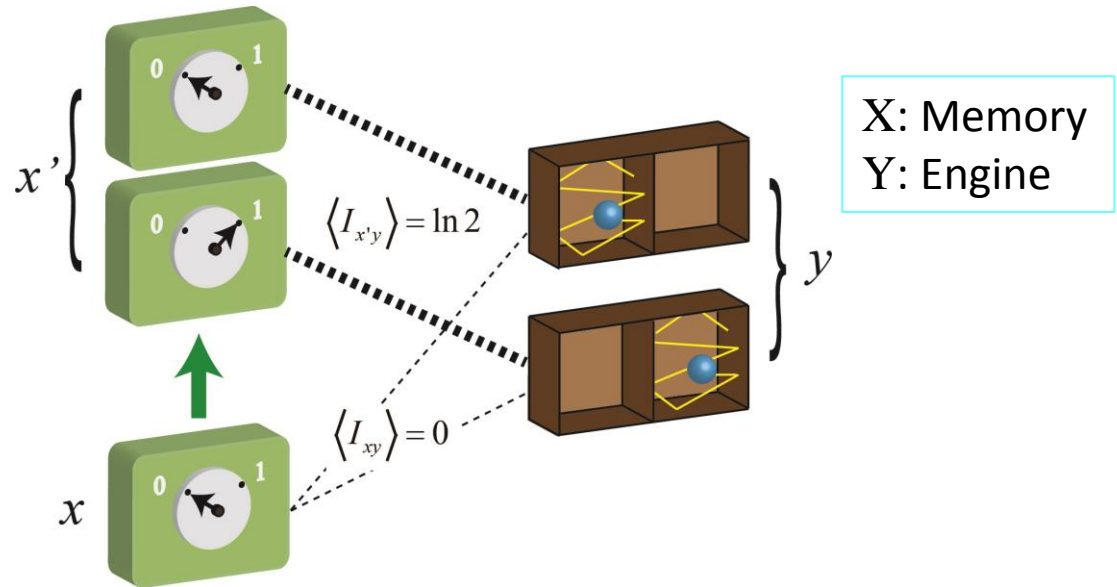
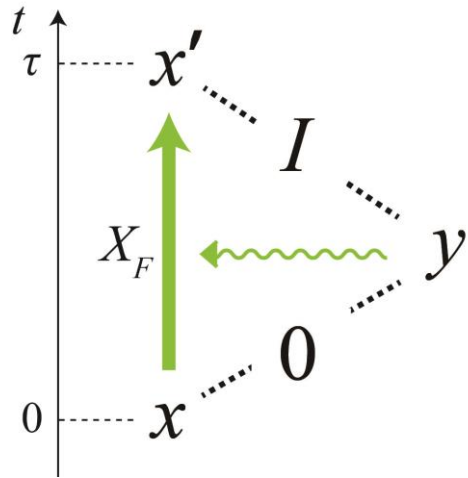


$$\Delta I \equiv I_{\text{rem}} - I$$

$$\left\langle e^{-\Delta S_{\text{XB}} + (I_{\text{rem}} - I)} \right\rangle = 1 \quad \Rightarrow \quad \langle \Delta S_{\text{XB}} \rangle \geq -\langle I - I_{\text{rem}} \rangle$$

$$\Rightarrow \quad W_{\text{ext}} \leq -\Delta F + k_{\text{B}} T \langle I \rangle \quad F : \text{equilibrium free energy}$$

Special Case 2: Measurement

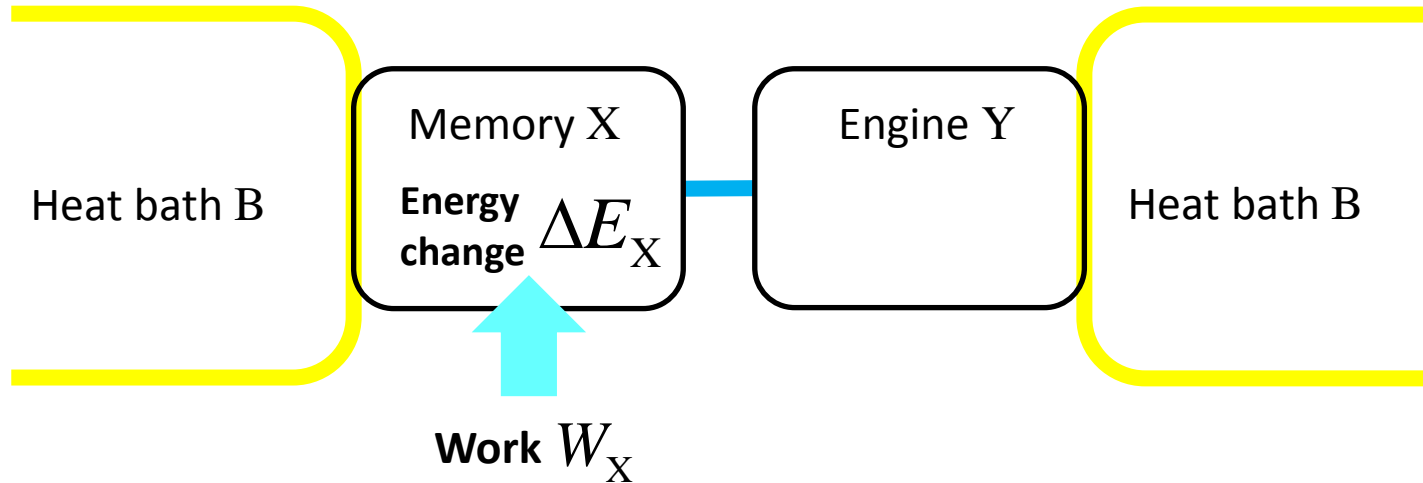


$$\Delta I \equiv I$$

$$\rightarrow \left\langle e^{-\Delta s_{XB} + I} \right\rangle = 1$$

$$\rightarrow \left\langle \Delta s_{XB} \right\rangle \geq \left\langle I \right\rangle$$

Minimal Energy Cost for Measurement



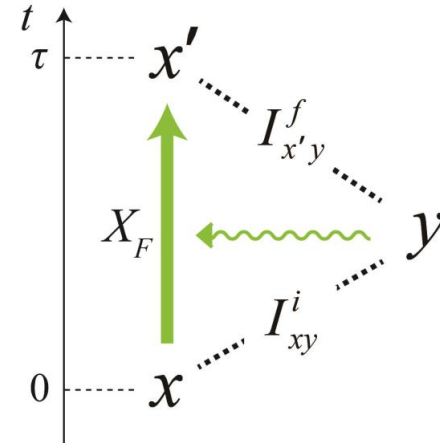
$$\langle \Delta s_{XB} \rangle \geq \langle I \rangle \quad \Rightarrow \quad \beta \langle W_X \rangle \geq \underbrace{\beta \langle \Delta E_X \rangle - \langle \Delta s_X \rangle}_{\text{Change in the nonequilibrium free energy of only X}} + \underbrace{\langle I \rangle}_{\text{Additional energy cost to obtain information}}$$

Information is not free

General Principle of Information Thermodynamics

$$\left\langle e^{-\Delta S_{\text{XB}} + \Delta I} \right\rangle = 1$$

$$\Rightarrow \left\langle \Delta S_{\text{XB}} \right\rangle \geq \left\langle \Delta I \right\rangle$$



Feedback:

$$\left\langle e^{-\Delta S_{\text{XB}} + (I_{\text{rem}} - I)} \right\rangle = 1$$

$$\Rightarrow \left\langle \Delta S_{\text{XB}} \right\rangle \geq -\left\langle I - I_{\text{rem}} \right\rangle$$

Measurement:

$$\left\langle e^{-\Delta S_{\text{XB}} + I} \right\rangle = 1$$

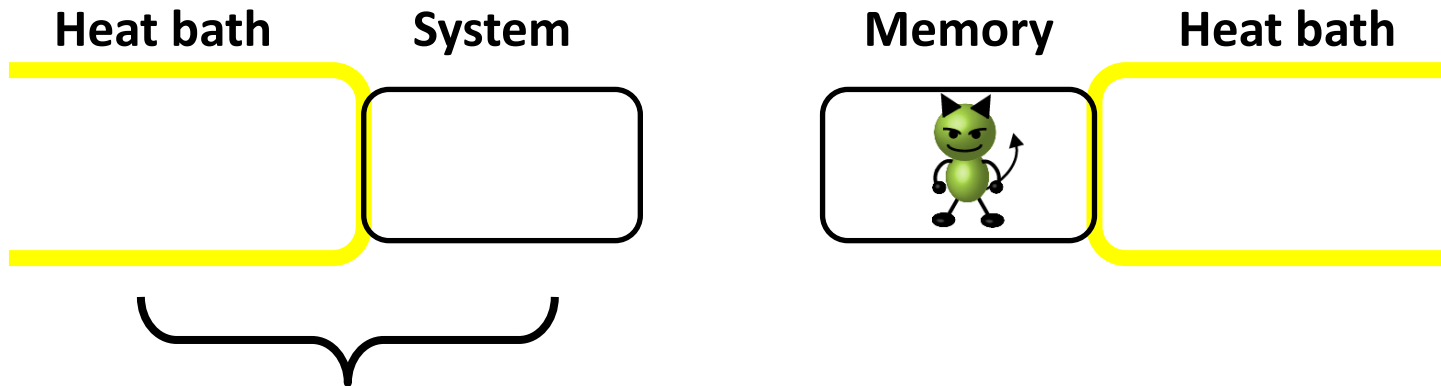
$$\Rightarrow \left\langle \Delta S_{\text{XB}} \right\rangle \geq \left\langle I \right\rangle$$

Unified formulation of measurement and feedback

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Problem



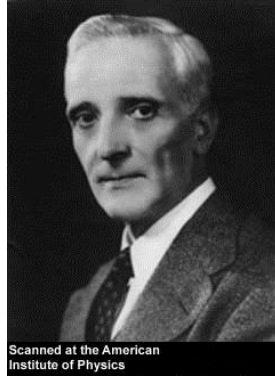
What compensates for the entropy decrease here?

For Szilard engine, $\langle \Delta S_{SB} \rangle = -\ln 2$



Conventional Arguments

Measurement
process



Brillouin

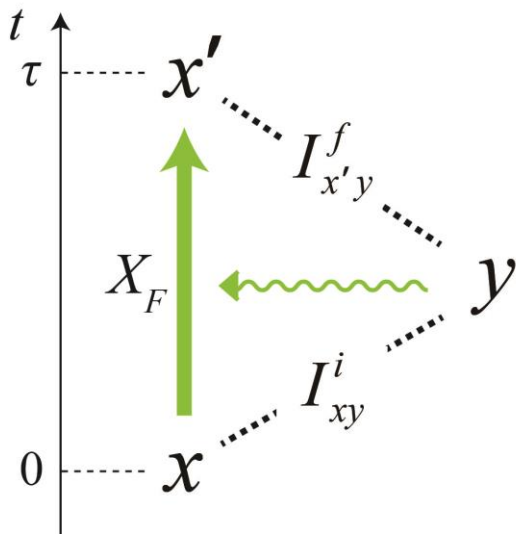
Erasure process
(From Landauer principle)



Bennett
&
Landauer

Widely accepted since 1980's

Total Entropy Production



Total entropy production in XYB

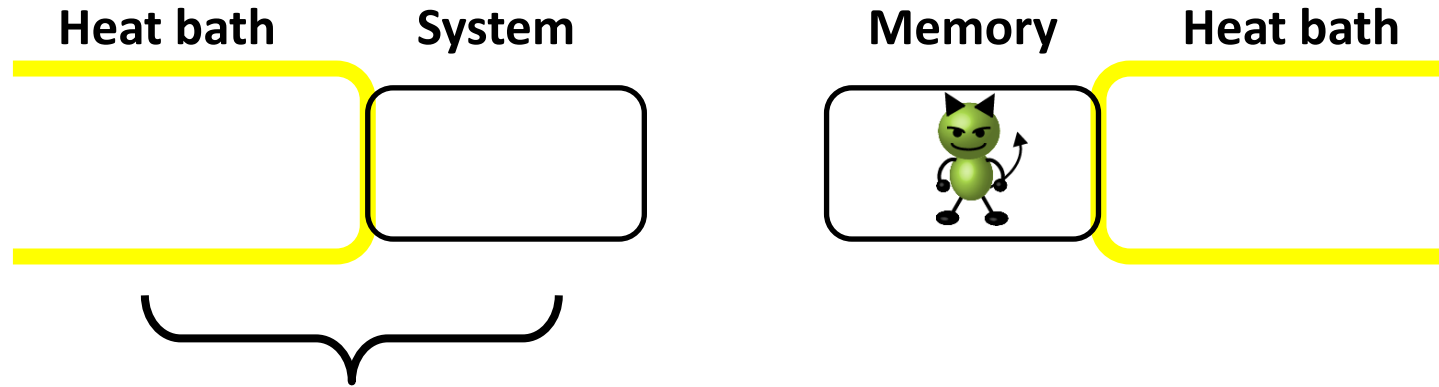
$$\begin{aligned} \Delta S_{\text{XYB}} &\equiv \Delta S_{\text{XY}} - \beta Q_X \\ &= \Delta S_{\text{XB}} - \Delta I \end{aligned}$$

$$\langle \Delta S_{\text{XB}} \rangle \geq \langle \Delta I \rangle \quad \longleftrightarrow \quad \langle \Delta S_{\text{XYB}} \rangle \geq 0$$

Equality: thermodynamically reversible

If the mutual information is taken into account, the total entropy production is always nonnegative for each process of measurement or feedback.

Revisit the Problem



What compensates for the entropy decrease here?

➡ Mutual-information change compensates for it.

For Szilard engine, $\langle \Delta S_{SB} \rangle = -\ln 2$ 

➡ $\langle \Delta S_{SMB} \rangle = \langle \Delta S_{SB} \rangle + \langle I \rangle = -\ln 2 + \ln 2 = 0$

Key to Resolve the Paradox

- Maxwell's demon is consistent with the second law for measurement and feedback processes **individually**
 - The mutual information is the key
- We don't need the Landauer principle to understanding the consistency

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Thermodynamics of Autonomous Information Processing

- Second law & fluctuation theorem**

Allahverdyan, Dominik & Guenter, *J. Stat. Mech.* (2009)

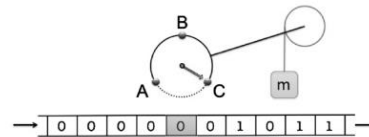
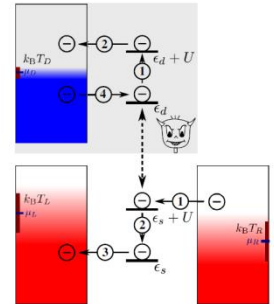
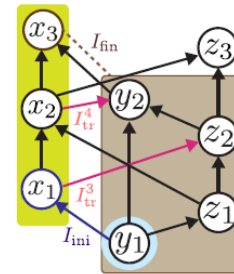
Hartich, Barato, & Seifert, *J. Stat. Mech.* (2014)

Horowitz & Esposito, *Phys. Rev. X* (2014)

Horowitz & Sandberg, *New J. Phys.* (2014)

Shiraishi & Sagawa, *Phys. Rev. E* (2015)

Ito & Sagawa, *Phys. Rev. Lett.* (2013)



- Models of autonomous Maxwell's demons**

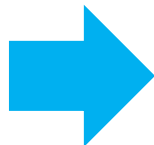
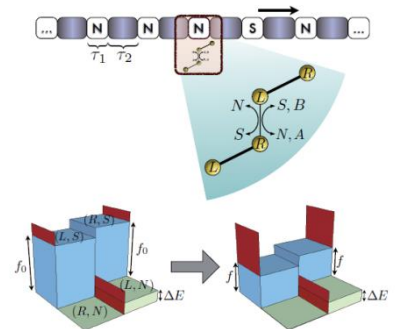
Mandal & Jarzynski, *PNAS* (2012)

Mandal, Quan, & Jarzynski, *Phys. Rev. Lett.* (2013)

Strasberg, Schaller, Brandes, & Esposito *Phys. Rev. Lett.* (2013)

Horowitz, Sagawa, & Parrondo, *Phys. Rev. Lett.* (2013)

Shiraishi, Ito, Kawaguchi & Sagawa, *New J. Phys.* (2015)



Toward deeper understanding of information nanomachines

Two Approaches

- **“Transfer entropy”** approach
 - ✓ Applicable to non-Markovian dynamics
 - ✓ Second law is weaker in Markovian dynamics

Ito & Sagawa, Phys. Rev. Lett. (2013)



But we derived a stronger version! (**Poster by Ito**)

- **“Information flow”** approach
 - ✓ Not applicable to non-Markovian dynamics
 - ✓ Second law is stronger in Markovian dynamics

Second law: Allahverdyan, Dominik & Guenter, J. Stat. Mech. (2009)

Hartich, Barato, & Seifert, J. Stat. Mech. (2014)

Horowitz & Esposito, Phys. Rev. X (2014)

Horowitz & Sandberg, New J. Phys. (2014)

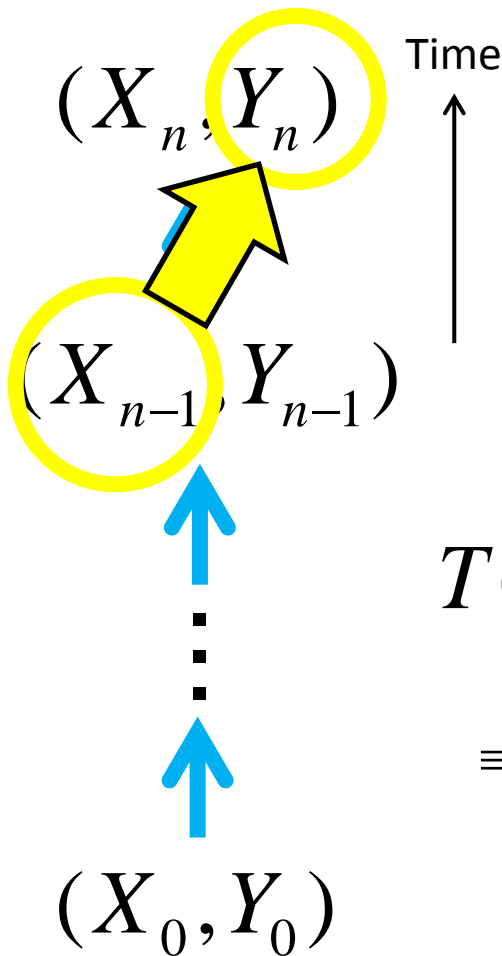
Fluctuation theorem: Shiraishi & Sagawa, PRE (2015)



Poster by Shiraishi

Transfer Entropy

Directional information transfer between two systems



Transfer entropy:

Directional information flow
from X to Y

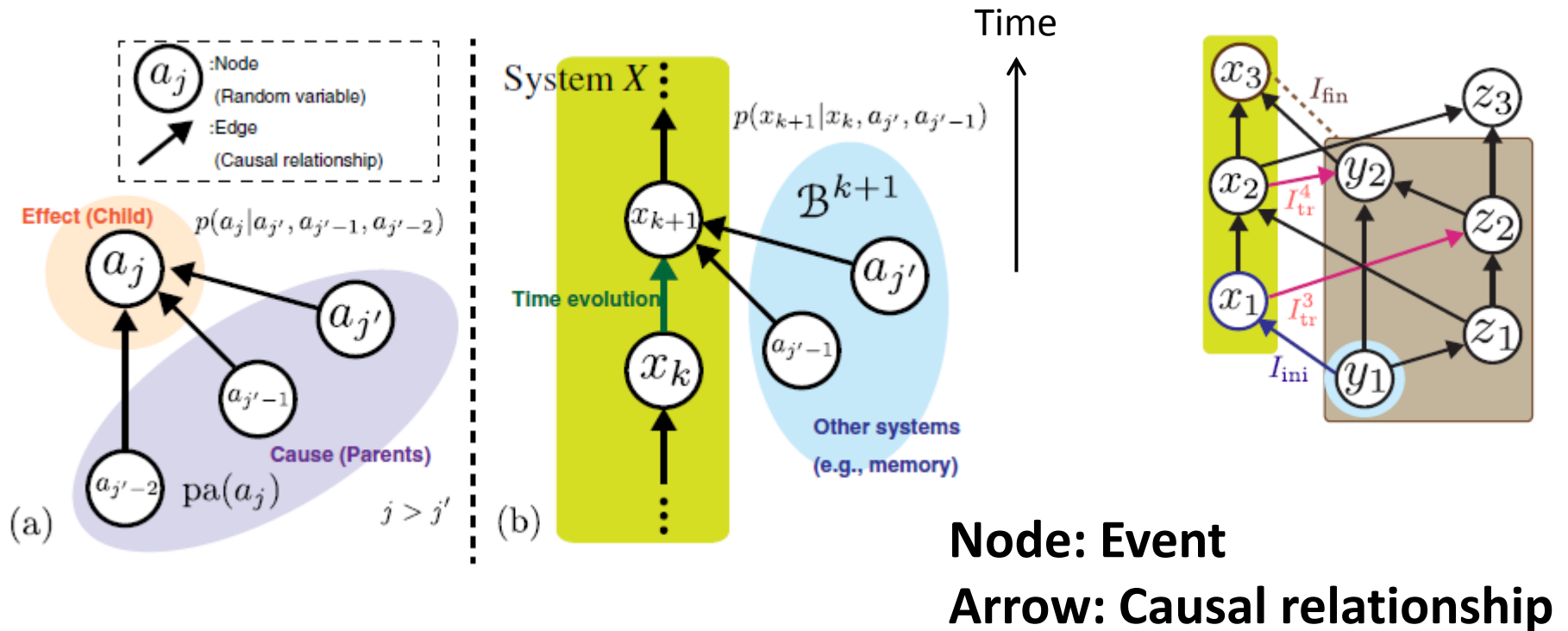
during time n and $n+1$

Conditional mutual information

$$T(X_{n-1} \rightarrow Y_n) \equiv I(X_{n-1} : Y_n | Y_{n-1} \cdots Y_0)$$

$$\equiv \sum_{x_{n-1}, y_0, \dots, y_n} p(x_{n-1}, y_0, \dots, y_n) \ln \frac{p(x_{n-1}, y_n | y_0, \dots, y_{n-1})}{p(x_{n-1} | y_0, \dots, y_{n-1}) p(y_n | y_0, \dots, y_{n-1})}$$

Many-body Systems with Complex Information Flow

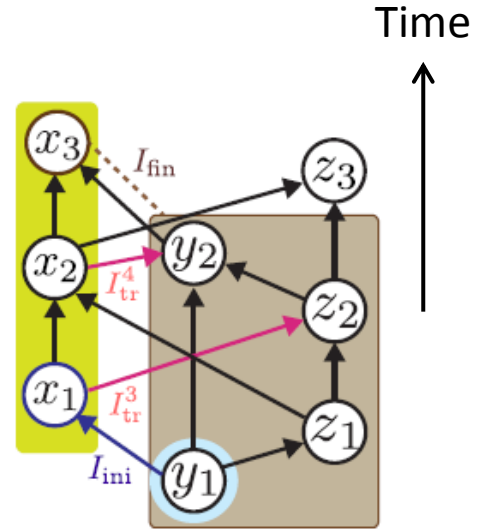


Characterize the dynamics by **Bayesian networks**

Second Law on Bayesian Networks

$$\Delta S_{\text{XB}} \geq \Theta$$

$$\Theta \equiv I_{\text{fin}} - I_{\text{ini}} - \sum_l I_{\text{tr}}^l$$



S. Ito & T. Sagawa, PRL 111, 180603 (2013)

ΔS_{XB} : Entropy production in X and the bath

I_{ini} : Initial mutual information between X and other systems

I_{fin} : Final mutual information between X and other systems

I_{tr}^l : **Transfer entropy** from X to other systems

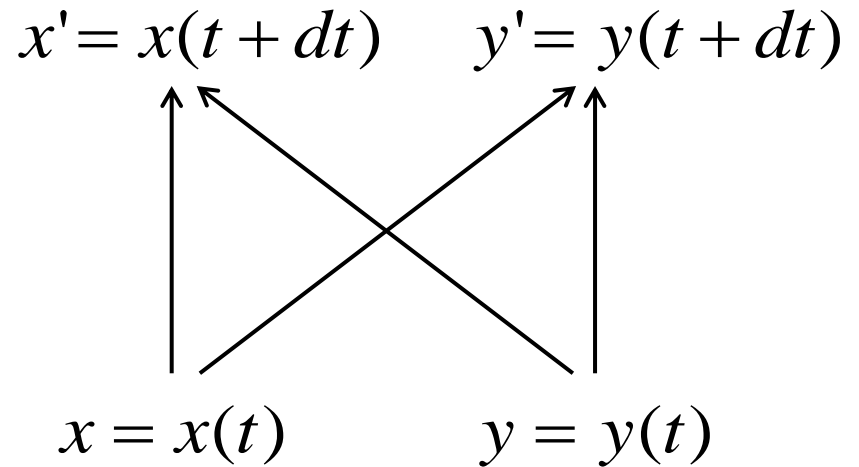
Information Flow VS Transfer Entropy

Infinitesimal transition of coupled Langevin system

$$\dot{x}(t) = f(x(t), y(t)) + \xi_x(t)$$

$$\dot{y}(t) = g(x(t), y(t)) + \xi_y(t)$$

$$\langle \xi_x(t) \xi_y(t) \rangle = 0 \quad : \text{independent noise}$$



Stronger: $\langle s(x') - s(x) - \beta Q \rangle \geq \langle \underline{I(x': y) - I(x : y)} \rangle$

Information flow

Weaker: $\langle s(x') - s(x) - \beta Q \rangle \geq \langle \underline{I(x' : y') - I(x : y) - I(x : y' | y)} \rangle$

Transfer entropy

$\longleftrightarrow \langle s(x' | y') - s(x | y) - \beta Q \rangle \geq -\langle \underline{I(x : y' | y)} \rangle$

Outline

- Introduction
- Information and entropy
- Information thermodynamics: a general framework
- Paradox of Maxwell's demon

- Thermodynamics of autonomous information processing
- **Application to biochemical signal transduction**

- Summary

Toward Biological Information Processing

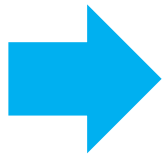
What is the role of information in living systems?

Mutual information is experimentally accessible

ex. Apoptosis path: Cheong *et al.* *Science* (2011).

There is no explicit channel coding inside living cells;

Shannon's second theorem is not straightforwardly applicable



Application of information thermodynamics

Barato, Hartich & Seifert, *New J. Phys.* **16**, 103024 (2014).

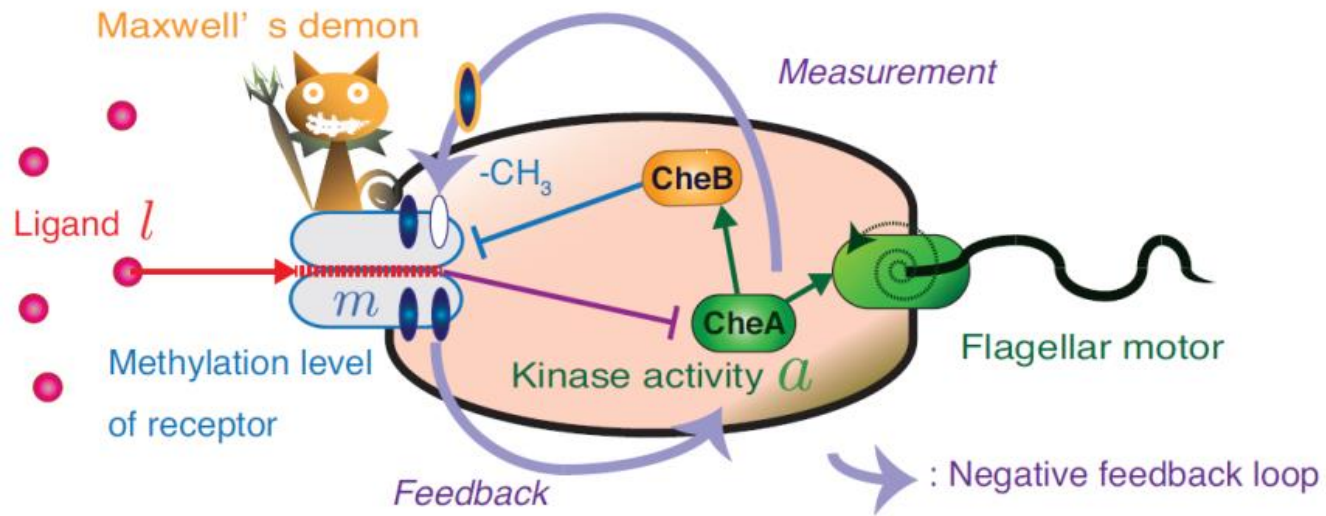
Sartori, Granger, Lee & Horowitz, *PLoS Comput. Biol.* **10**, e1003974 (2014).

Ito & Sagawa, *Nat. Commu.* **6**, 7498 (2015).

Our finding:

Relationship between information and the robustness of adaptation

Signal Transduction of *E. Coli* Chemotaxis



E. Coli moves toward food (ligand)

The information about **ligand density** is transferred to the **methylation level** of the receptor, and used for the feedback to the **kinase activity**.

Adaptation Dynamics

2D Langevin model

Y. Tu *et al.*, *Proc. Natl. Acad. Sci. USA* **105**, 14855 (2008).
F. Tostevin and P. R. ten Wolde, *Phys. Rev. Lett.* **102**, 218101 (2009).
F. G. Lan *et al.*, *Nature Physics* **8**, 422 (2012).

$$\dot{a}_t = -\frac{1}{\tau^a} [a_t - \bar{a}_t(m_t, l_t)] + \xi_t^a$$

$$\dot{m}_t = -\frac{1}{\tau^m} a_t + \xi_t^m$$

$$\langle \xi_t^x \rangle = 0 \quad \langle \xi_t^x \xi_{t'}^{x'} \rangle = 2T_t^x \delta_{xx'} \delta(t - t')$$

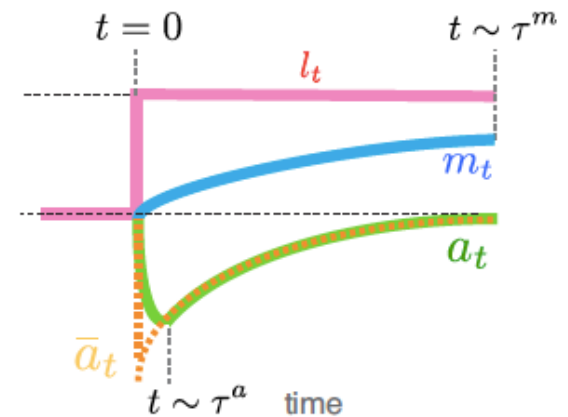
$\bar{a}_t(m_t, l_t) \simeq \alpha m_t - \beta l_t$: stationary value of a_t

$$\alpha, \beta > 0$$

Negative feedback loop:

- ✓ Instantaneous change of a_t in response to l_t
- ✓ Memorize l_t by m_t
- ✓ a_t goes back to the initial value

a_t : kinase activity
 m_t : methylation level
 l_t : average ligand density
 $\tau^m \gg \tau^a > 0$: time constants



Second Law of Information Thermodynamics

$$dI_t^{\text{tr}} + dS_t^{a|m} \geq \frac{J_t^a}{T_t^a} dt$$

(Weaker version with transfer entropy)

$dS_t^{a|m} := \langle \ln p(a_t|m_t) \rangle - \langle \ln p(a_{t+dt}|m_{t+dt}) \rangle$: Change in the conditional Shannon entropy

$dI_t^{\text{tr}} := I(a_t : m_{t+dt}|m_t)$: **Transfer entropy**

$\frac{J_t^a}{T_t^a} = \frac{1}{\tau^a T_t^a} \left[T_t^a - \frac{\langle (a_t - \bar{a}_t)^2 \rangle}{\tau^a} \right]$: **Robustness against the environmental noise**

Upper bound of the robustness is given by the transfer entropy

Stationary State

$$\langle (a_t - \bar{a}_t)^2 \rangle \geq \tau^a T_t^a \left[1 - \frac{dI_t^{\text{tr}}}{dt} \right]$$

Fluctuation (inaccuracy of information transmission) induced by environmental noise

Transfer entropy

Without feedback : $\langle (a_t - \bar{a}_t)^2 \rangle \geq \tau^a T_t^a$

Exact Expression of Transfer Entropy

If the Langevin equation is linear:

$$dI_t^{\text{tr}} = \frac{1}{2} \ln \left(1 + \frac{dP_t}{N_t} \right)$$

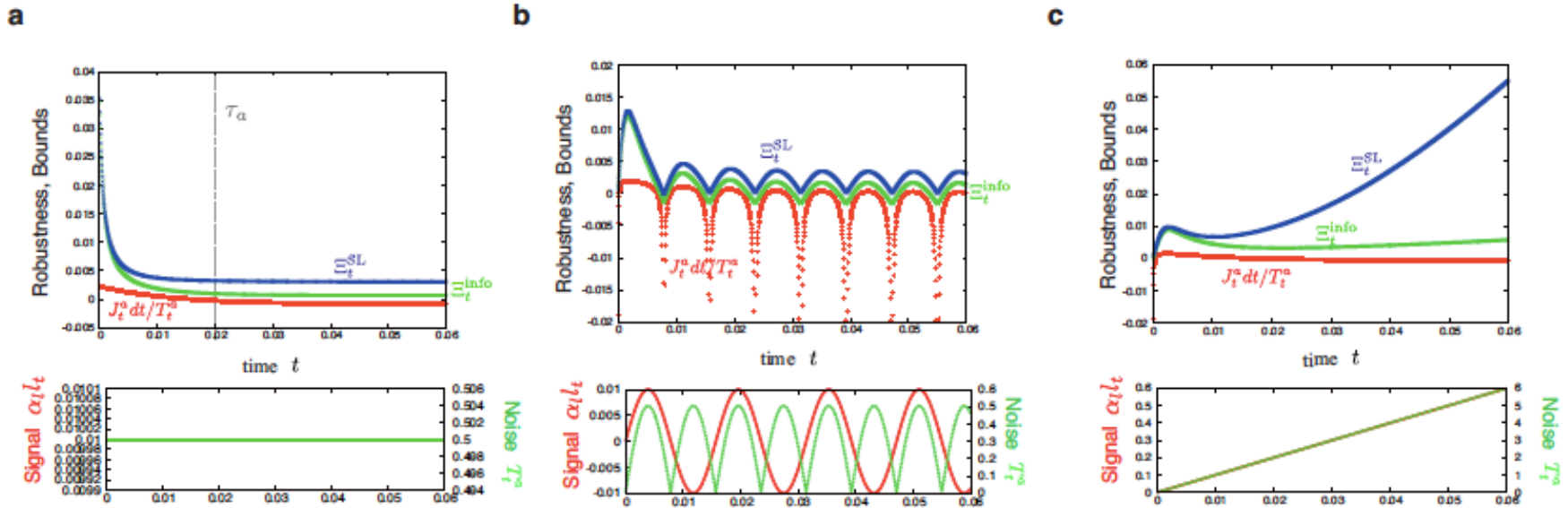
Signal-to-noise ratio

$$dP_t := \frac{(\rho_t^{am})^2 V_t^a}{(\tau^m)^2} dt \quad : \text{power of the signal from } a \text{ to } m$$

$$N_t := 2T_t^m \quad : \text{noise of } m \quad V_t^x := \langle x_t^2 \rangle - \langle x_t \rangle^2 \quad \rho_t^{am} := \frac{\langle a_t m_t \rangle - \langle a_t \rangle \langle m_t \rangle}{\sqrt{V_t^a V_t^m}}$$

Analogous to the Shannon–Hartley theorem

Information-Thermodynamic Efficiency



Input ligand signal: a, step function. b, sinusoidal function. c, linear function.

Numerical simulation:

Red: robustness of adaptation

Green: information-thermodynamic bound

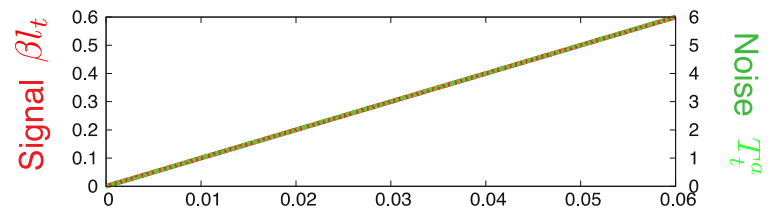
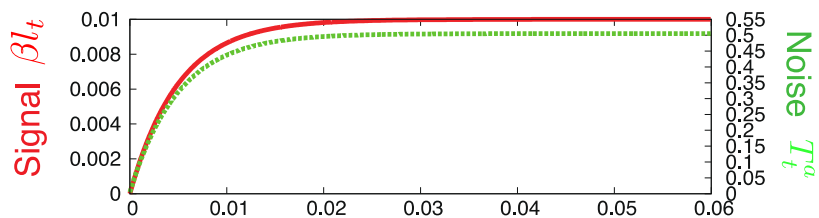
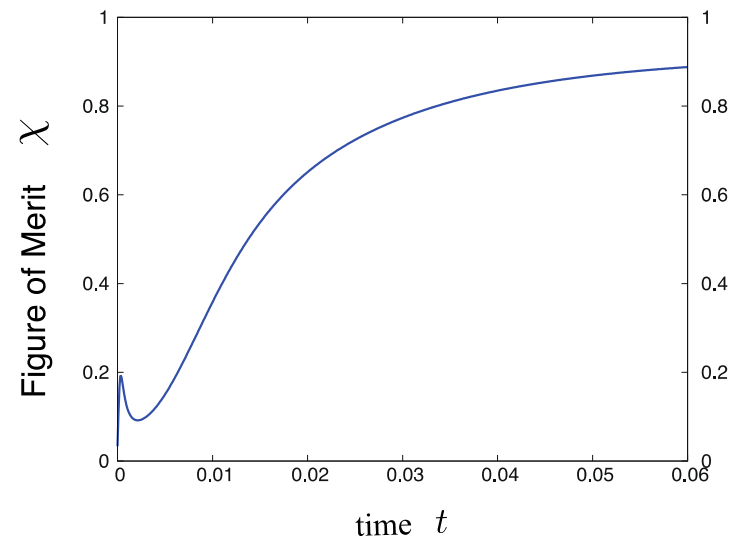
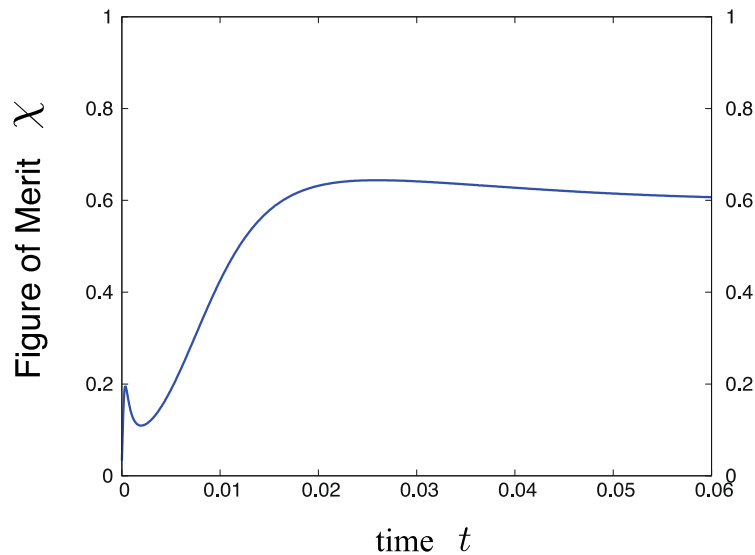
Blue: conventional thermodynamic bound

$\mathbb{I}_t^{\text{info}}$
 \mathbb{I}_t^{SL}

- ✓ Information thermodynamics gives a stronger bound!
- ✓ The adaptation dynamics is inefficient (dissipative) as a conventional thermodynamic engine, but efficient as an information-thermodynamic engine.

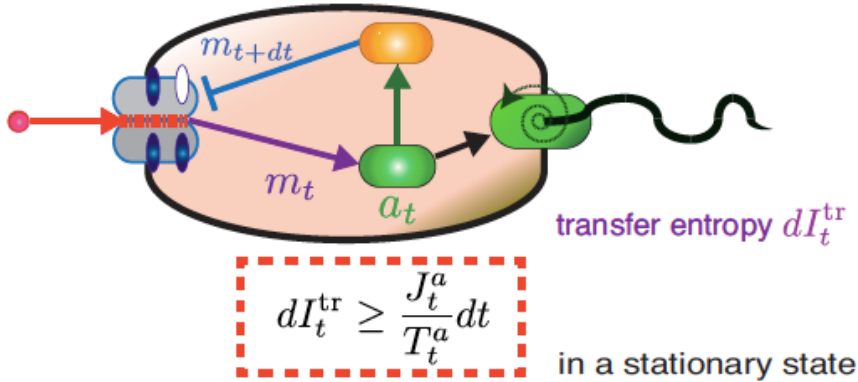
Information-Thermodynamic Figure of Merit

$$\chi := 1 - \frac{\Xi_t^{\text{info}} - \int J_t^a dt / T_t^a}{\Xi_t^{\text{SL}} - \int J_t^a dt / T_t^a}$$



Comparison with Shannon's Information Theory

a Robustness of signal transduction against noise J_t^a

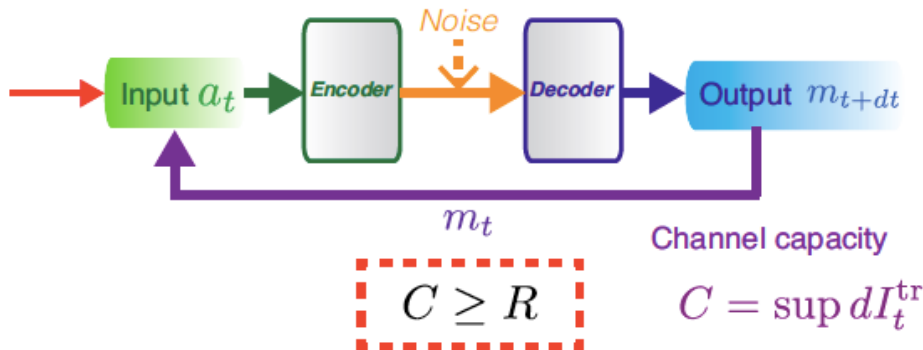


Second law of information thermodynamics

Information dI_t^{tr}
gives the bound of robustness J_t^a

Well-defined in living cells

b Achievable information rate
(Accuracy of information transmission against noise) R



Shannon's second theorem

Information (channel capacity) C
gives the bound of achievable rate R

How to define in living cells??

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Summary

- Unified framework of information thermodynamics

T. Sagawa & M. Ueda, *Phys. Rev. Lett.* **109**, 180602 (2012).

T. Sagawa & M. Ueda, *New J. Phys.* **15**, 125012 (2013).

- Fluctuation theorem for autonomous information processing

S. Ito & T. Sagawa, *Phys. Rev. Lett.* **111**, 180603 (2013).

Review: S. Ito & T. Sagawa, arXiv:1506.08519 (2015).

N. Shiraishi & T. Sagawa, *Phys. Rev. E* **91**, 012130 (2015).

- Information thermodynamics of biochemical signal transduction

S. Ito & T. Sagawa, *Nature Communications* **6**, 7498 (2015).

Review:

J. M. R. Parrondo, J. M. Horowitz, & T. Sagawa, *Nature Physics* **11**, 131-139 (2015).

Thank you for your attention!