

Symmetry breaking patterns in QCD and the study of cold dense matter



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with

Martin Pospiech and Jens Braun

New Frontiers in QCD 2018

YITP, Kyoto University

[J. Braun, ML, M. Pospiech, PRD **96**, 076003 (2017)]

[J. Braun, ML, M. Pospiech, PRD **97**, 076010 (2018)]

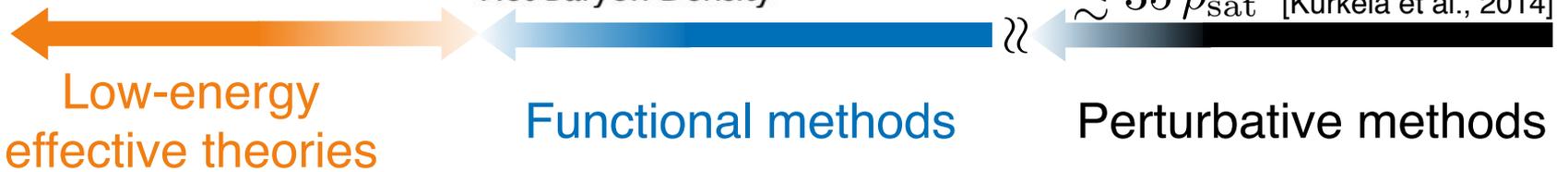
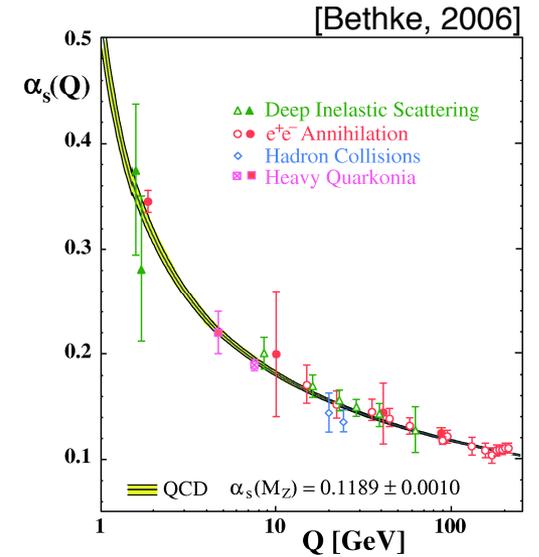
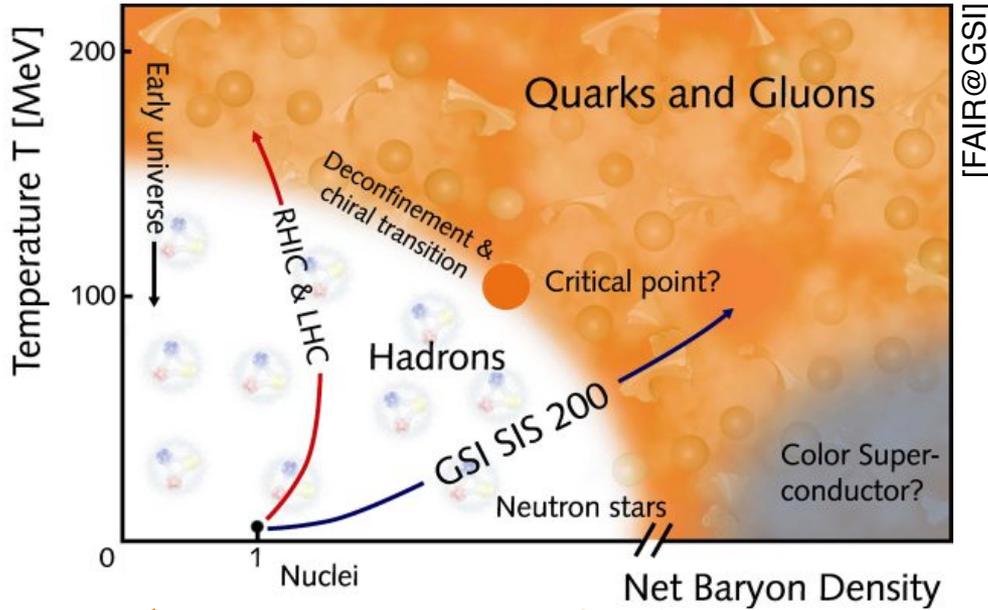
[J. Braun, ML, J. M. Pawłowski, arXiv:1806.04432]

[J. Braun, C. Drischler, K. Hebeler, ML, M. Pospiech,
A. Schwenk, in preparation]



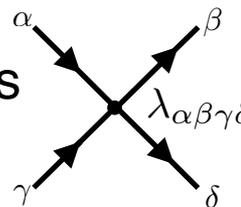
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QCD phase diagram: Neutron stars and the cold dense EoS

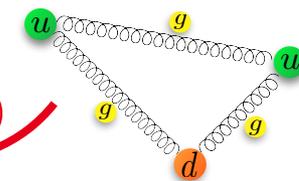


EoS II

Martin's
Talk



I

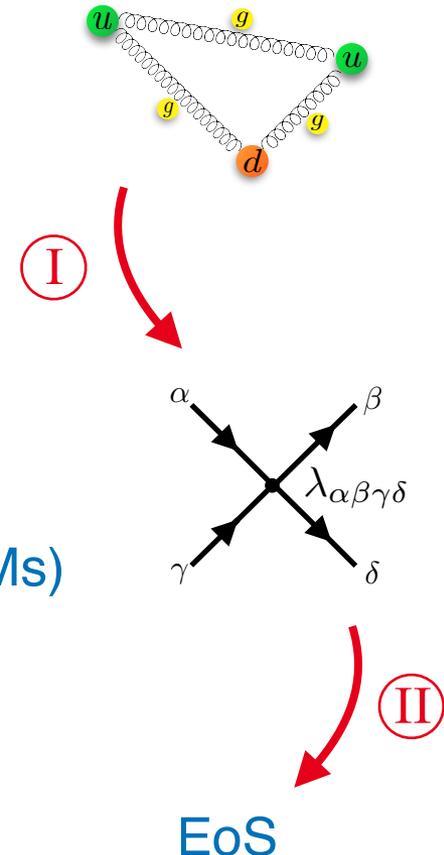


① Including dynamic gauge fields

- Framework and structure of β functions
- Symmetry breaking patterns
- Implications on the phase diagram

② Connecting to low-energy effective models (LEMs)

- Quark-meson-diquark model
- Findings on the equation of state (EoS)

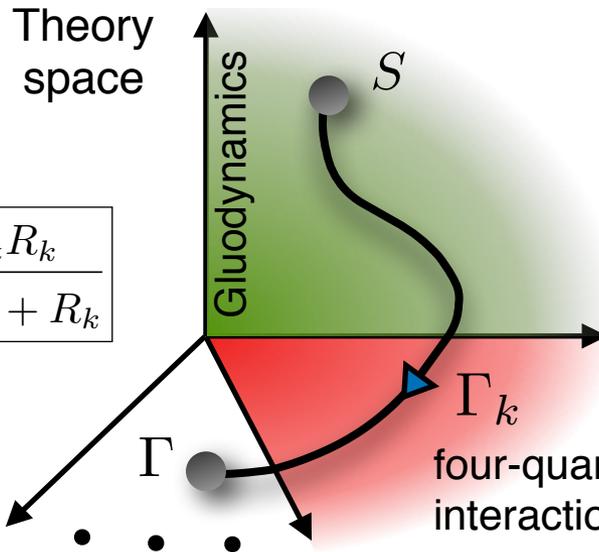
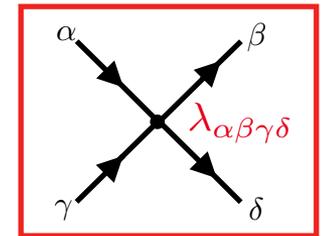
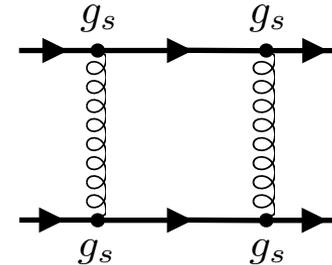
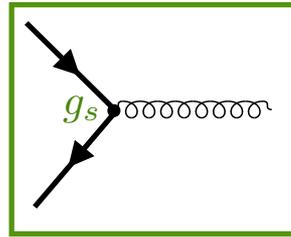
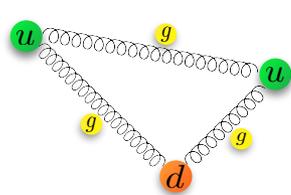


Conclusions and outlook

Functional renormalization group (FRG)

From high to low energies in QCD

RG steps



$$k \frac{\partial}{\partial k} \Gamma_k = \frac{1}{2} \text{STr} \frac{\partial_t R_k}{\Gamma_k^{(2)} + R_k}$$

[C. Wetterich, 1993]

$$[(\bar{\psi} \gamma_\mu \psi)^2]$$

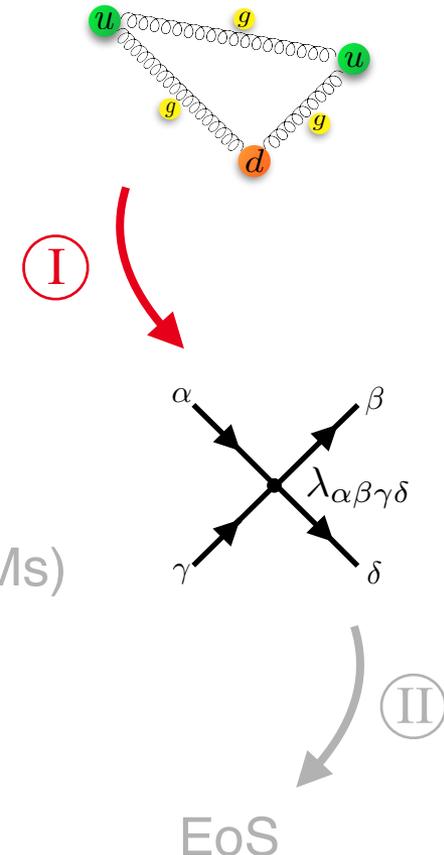
$$[(\bar{\psi} \gamma_\mu \gamma_5 T^a \psi)^2]$$

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Conclusions and outlook

Ansatz for the effective average action

Quark self-interactions and gauge degrees of freedom

$$\Gamma_{\text{LO}}[\bar{\psi}, \psi, A] = \int_0^{\frac{1}{T}} d\tau \int d^3x \left\{ \bar{\psi} (i\gamma_0 \partial_0 + i\gamma_i \partial_i - i\mu\gamma_0) \psi + \frac{1}{2} \sum_{j \in \mathcal{B}} Z_j \bar{\lambda}_j \mathcal{L}_j \right.$$

Symmetries

Lorentz group: $SO(1, 3) \longrightarrow SO(3)$

Flavor space: $SU_L(2) \otimes SU_R(2) \otimes U_V(1)$

Color space: $SU(N_c)$



Fierz-complete basis:
10 channels

$$\mathcal{L}_{(\bar{\psi}\psi)^2}^{(\sigma-\pi)} = (\bar{\psi}\psi)^2 - (\bar{\psi}\gamma_5\vec{\tau}\psi)^2$$

\leftrightarrow formation of
chiral condensate

$$\mathcal{L}_{(\bar{\psi}\psi)^2}^{\text{CSC}} \sim (i\bar{\psi}\gamma_5\tau_A t_c^{A'} \mathcal{C}\bar{\psi}^T)(i\psi^T \mathcal{C}\gamma_5\tau_A t_c^{A'} \psi) \quad J^P = 0^+$$

e.g. [Rapp, Schäfer, Shuryak, Velkovsky, 1998]

\leftrightarrow formation of
diquark condensate

Ansatz for the effective average action

Quark self-interactions and gauge degrees of freedom



$$\Gamma_{\text{LO}}[\bar{\psi}, \psi, A] = \int_0^{\frac{1}{T}} d\tau \int d^3x \left\{ \bar{\psi} (i\gamma_0 \partial_0 + i\gamma_i \partial_i - i\mu\gamma_0) \psi + \frac{1}{2} \sum_{j \in \mathcal{B}} Z_j \bar{\lambda}_j \mathcal{L}_j \right.$$

$$+ \frac{Z_A}{4} (F_{\mu\nu}^a)^2 + \frac{1}{2\xi} (\partial_\mu A_\mu^a)^2 \quad \text{kinetic term for gluons}$$

$$\left. + \bar{g}_s \bar{\psi} \gamma_\mu A_\mu \psi \right\} \quad \text{quark-gluon interaction}$$

RG flow of the gauge coupling: $k \frac{\partial}{\partial k} g_s^2(k) = \eta_{A,k} g_s^2(k)$ [L. F. Abbott, 1980]

$$\eta_{A,k} = \underline{\eta_{A,k}^{\text{YM}}(T)} + \Delta\eta_{A,k}(T, \mu)$$

$$\Delta\eta_{A,k}(T, \mu) = \frac{Z_A^{-1}}{3(N_c^2 - 1)} \left(\frac{\partial^2}{\partial p^2} \Pi^\perp \cdot \text{[diagram]} \right) \Big|_{p=0}$$

[J. Braun and H. Gies, 2007]

[J. Braun, 2009]

[J. Braun et al., 2011]

[J. M. Pawłowski, 2011]

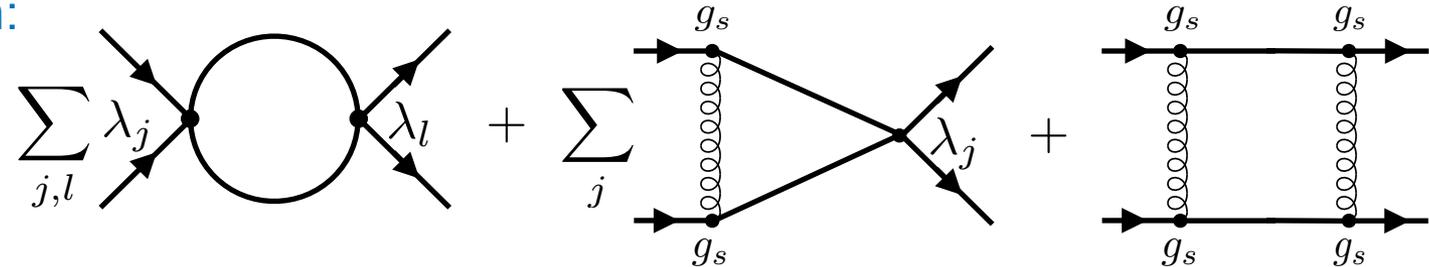
[J. Braun et al., 2014]

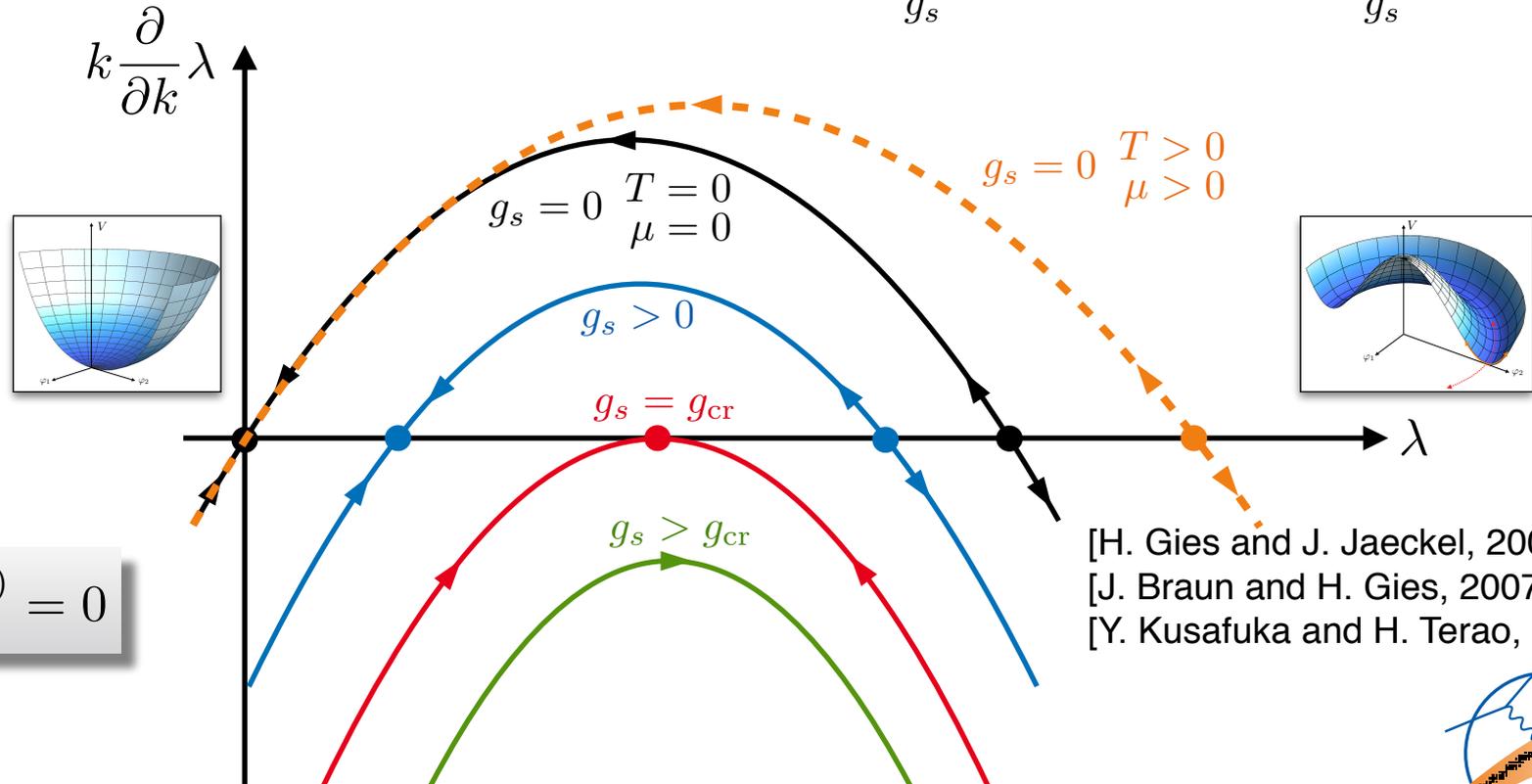


RG flow of four-quark interactions

Characteristics of the β functions and the fixed-point structure

RG flow equation:

$$k \frac{\partial}{\partial k} \lambda_i = 2\lambda_i - \sum_{j,l} \lambda_j \lambda_l + \sum_j g_s \lambda_j + \dots$$




$$\lambda_i^{(UV)} = 0$$

[H. Gies and J. Jaeckel, 2006]
[J. Braun and H. Gies, 2007]
[Y. Kusafuka and H. Terao, 2011]

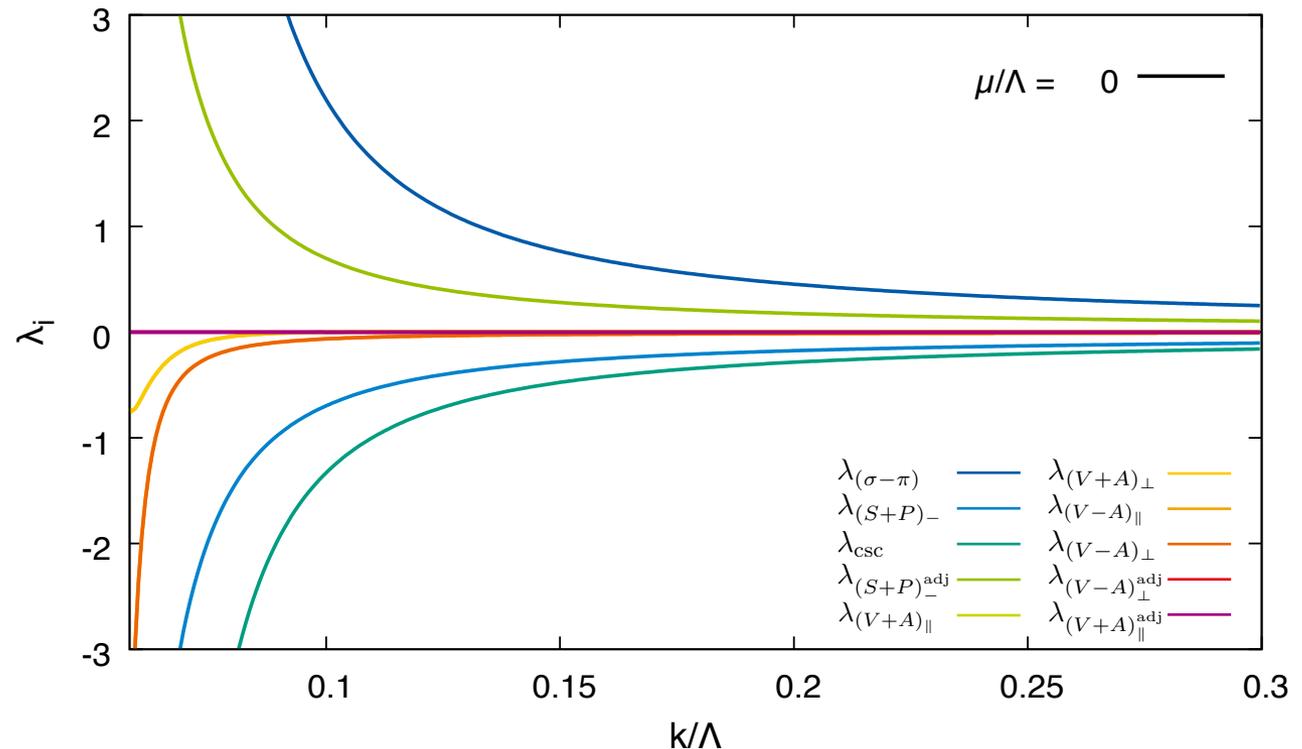
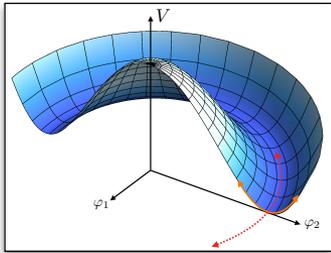
RG flow of four-quark interactions

Fixed-point structure and patterns of symmetry breaking

Illustration of the RG flow of the four-quark couplings
at zero temperature:

$T = 0$

$$\frac{1}{\lambda_i} \Big|_{k_0} \rightarrow 0$$



[J. Braun, 2006]
[M. Mitter, J. M. Pawłowski, N. Strodthoff, 2014]

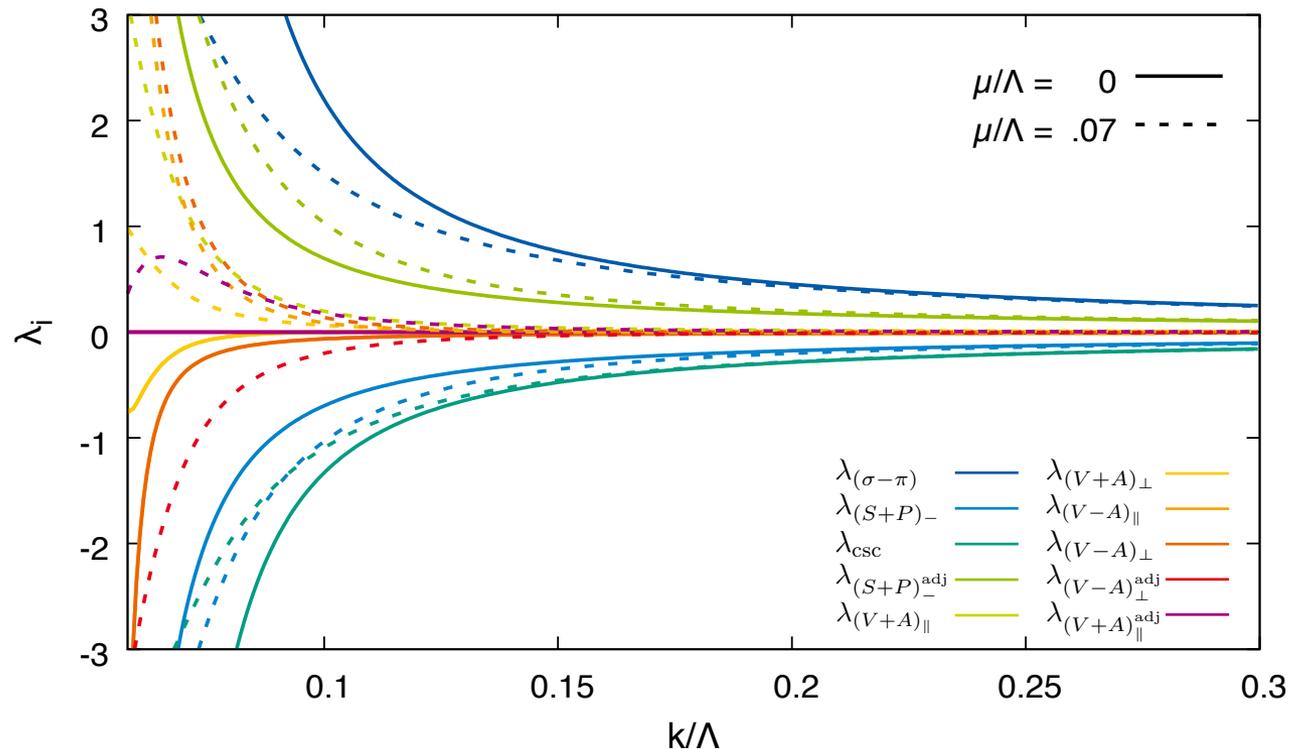
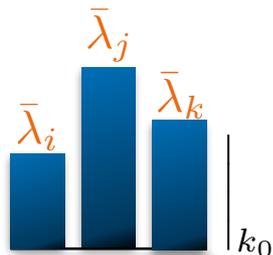
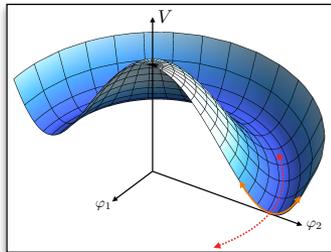
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Assessing relative
interaction strengths



Formation of
specific condensates

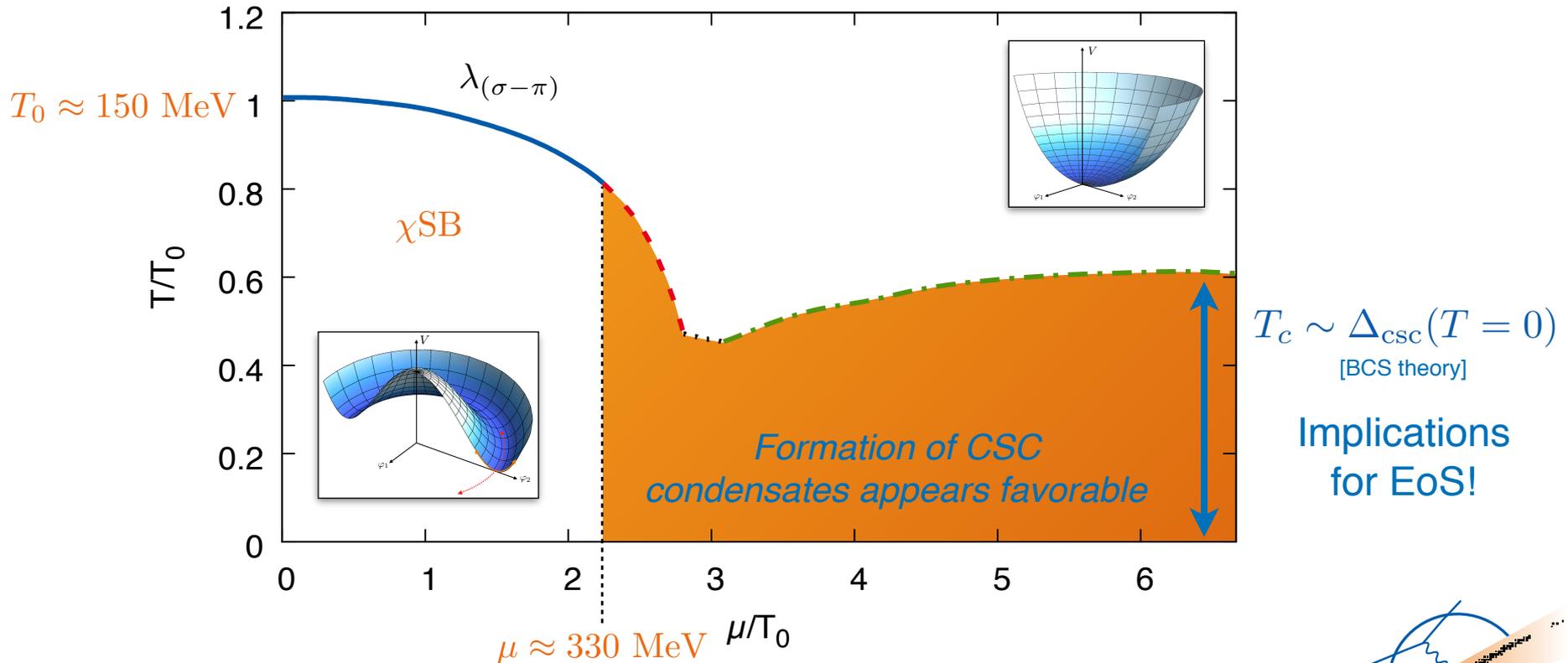
e.g. [J. Braun, H. Gies, L. Janssen, D. Roscher 2014]

Exploring the phase diagram

Fixed-point structure and patterns of symmetry breaking

Scale fixing: $g_s(\Lambda)$ at $\Lambda = 10$ GeV within experimental error [Bethke, 2006]

$$T_{\text{cr}}(\mu=0) \equiv T_0 \approx 150 \text{ MeV}$$

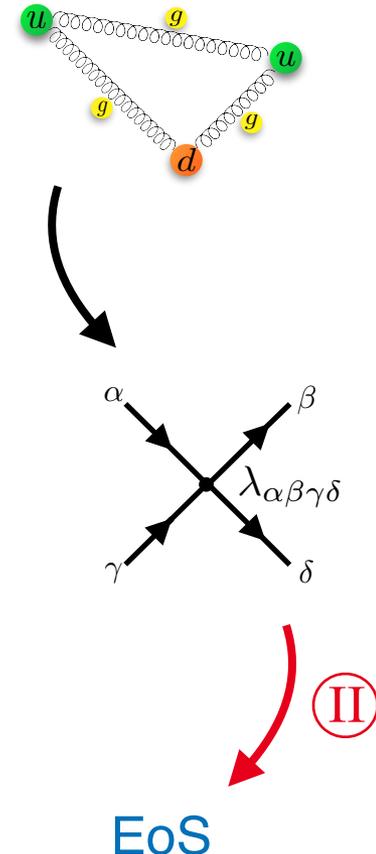


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Conclusions and outlook

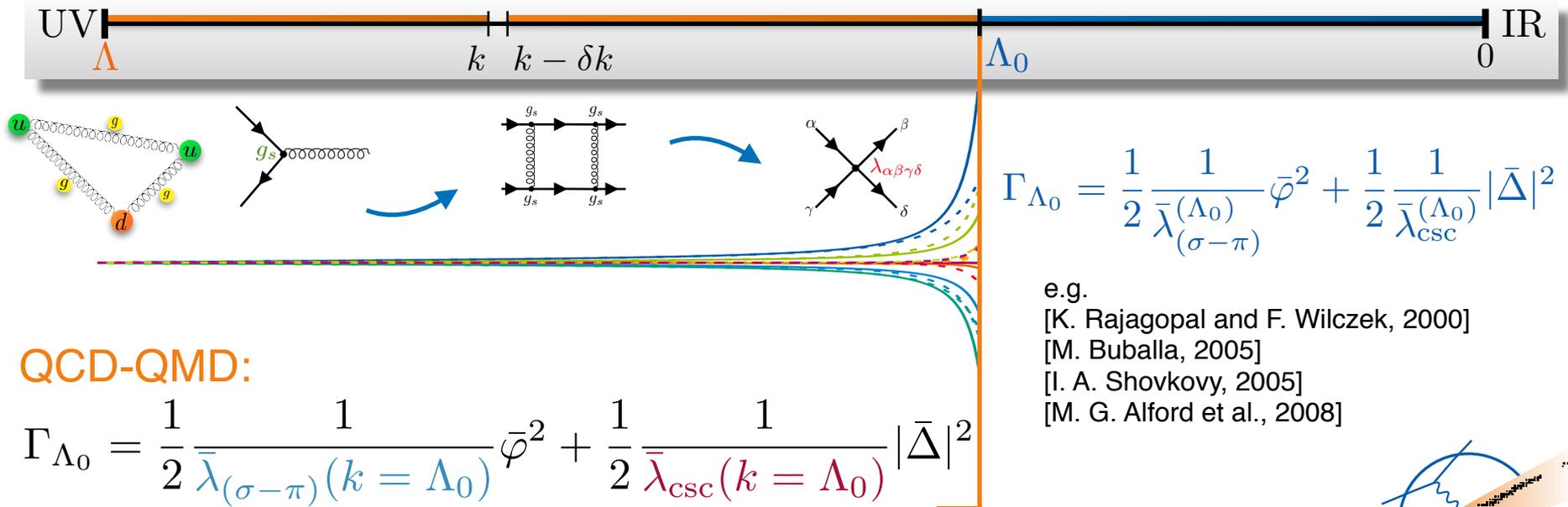
Connecting to low energy effective models

Determining parameters at high densities

Quark-meson-diquark model (QMD)

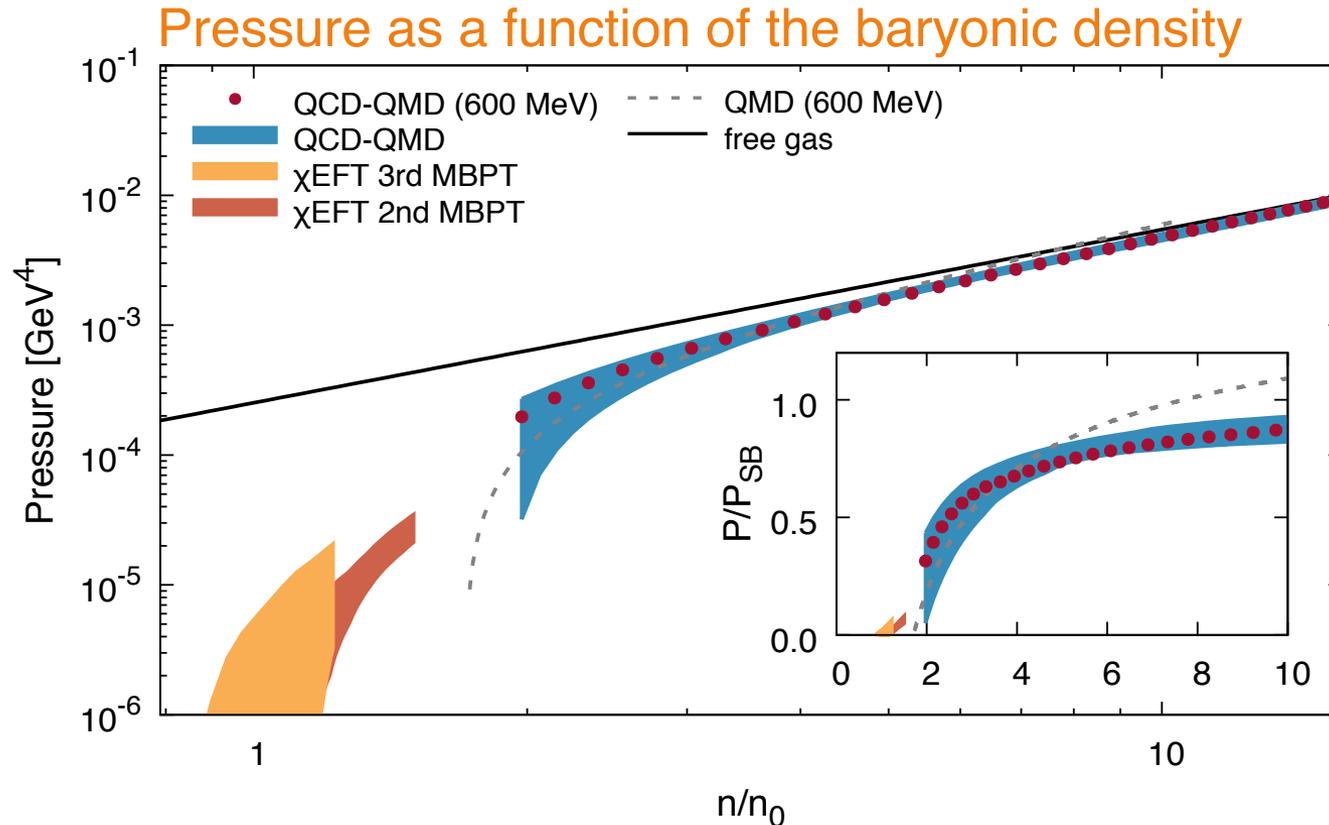
$$S = \int_x \left\{ \bar{q} \left(i\not{\partial} - i\mu\gamma_0 + i(\sigma + i\vec{\tau} \cdot \vec{\pi}\gamma_5) \right) q + \frac{1}{2} \frac{1}{\bar{\lambda}_{(\sigma-\pi)}} \varphi^2 \right\} \text{QM part}$$

$$+ \left\{ \bar{q} \gamma_5 \tau_2 \Delta_A^* T^A C \bar{q}^T - q^T C \gamma_5 \tau_2 \Delta_A T^A C q + \frac{1}{2} \frac{1}{\bar{\lambda}_{\text{csc}}} |\Delta|^2 \right\} \text{diquarks}$$



Implications on the equation of state

Connecting to the quark-meson-diquark model



preliminary

[χ EFT data by C. Drischler, K. Hebeler and A. Schwenk]

- Compatible to pQCD [E. S. Fraga, A. Kurkela, and A. Vuorinen, 2015]
- Degrees of freedom at smaller densities [M. Drews and W. Weise, 2016]

Conclusions and outlook

Dynamic gauge fields

- Dynamical generation of four-quark interactions
- Importance of Fierz completeness at high quark chemical potential and low temperature
- Formation of CSC condensates appears to be favoured at high densities

Connecting to LEMs: quark-meson-diquark model

- RG flow of four-quark couplings to fix parameters at high densities
- Equation of state:
 - Connects to free gas pressure at asymptotically large densities
 - Pressure towards smaller densities lifted to higher values; compatibility with χ EFT results appears to be improved

Outlook

- Improve connection between high- and low-energy regime by employing dynamical hadronization techniques