

Rotating/magnetized fermionic matter confined in a finite-size system


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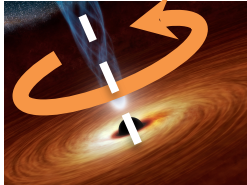
HL. Chen, K. Fukushima, XG. Huang, KM, PRD 93, 104052 (2016)

S. Ebihara, K. Fukushima, KM, PLB 764, 94 (2017)

HL. Chen, K. Fukushima, XG. Huang, KM PRD 96, 054032 (2017)

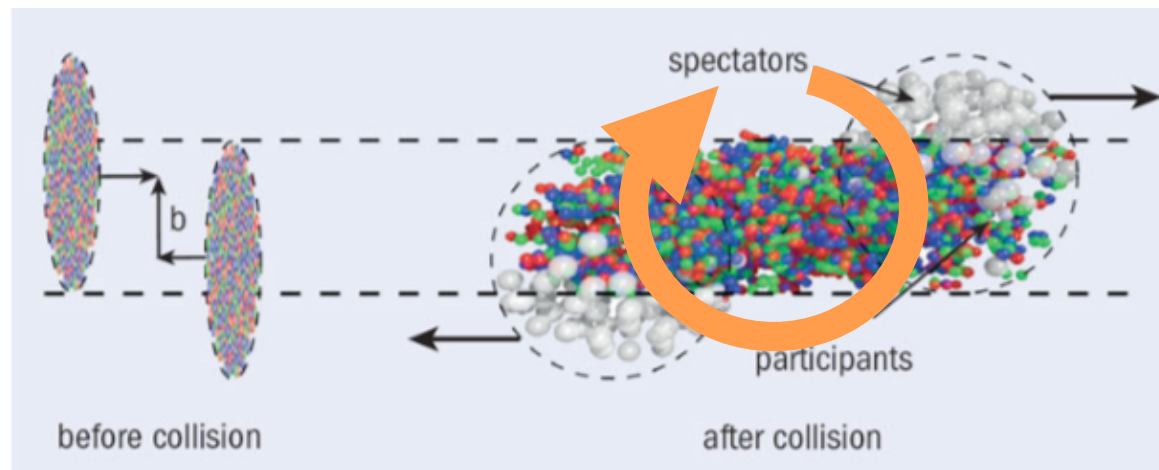
Rotating Relativistic Systems

• binary star merger 

• black hole 

• neutron star 

• noncentral heavy-ion collision

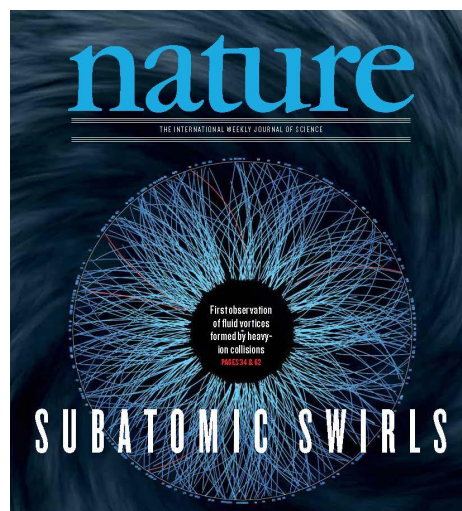


Angular momentum of heavy-ions

→ **Vorticity** (or **Rotation**)

First evidence:
global Λ -polarization (2017)

$$\Omega \gtrsim 10^{21} \text{ Hz}$$



#38



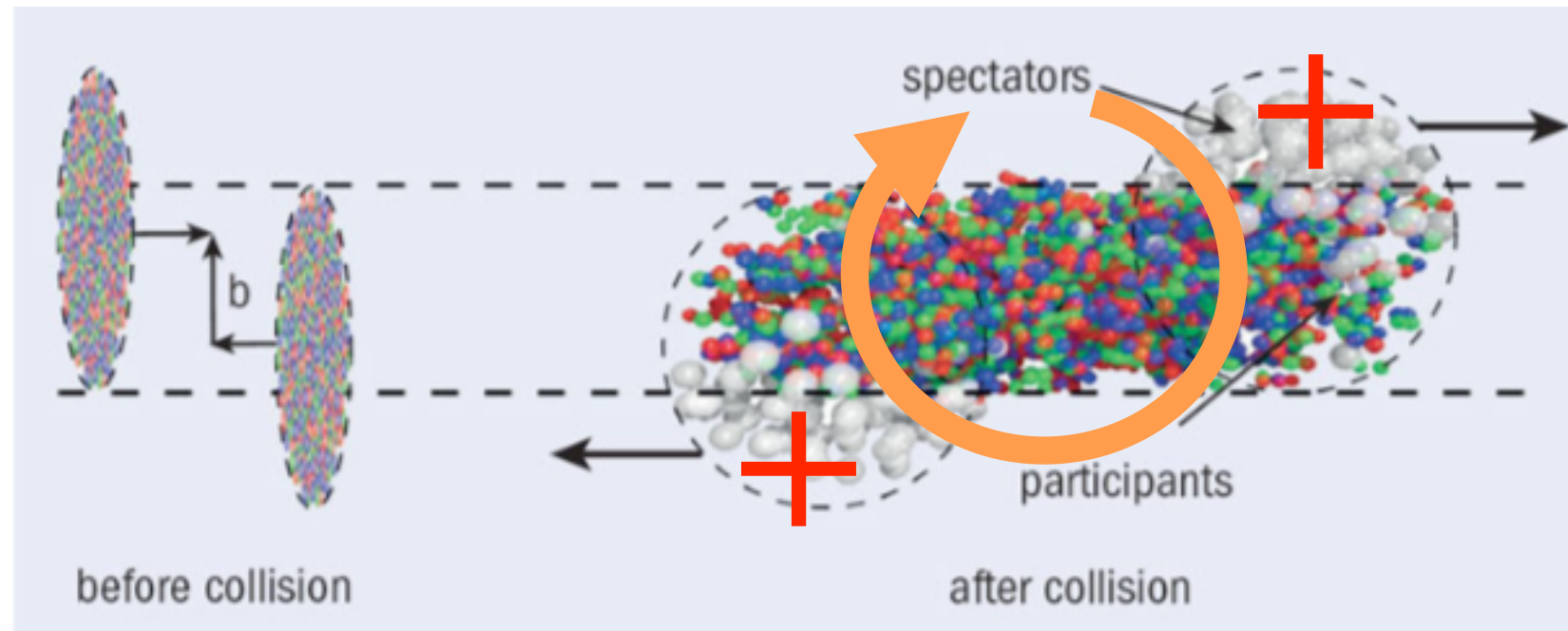
The Fastest Fluid

by Sylvia Morrow

Superhot material spins at an incredible rate.

Necessary to know **Rotational effect on QFT**

Rotating Magnetized System



angular momentum conservation

$$\Omega \sim 10^{22} \text{ Hz} \simeq 10 \text{ MeV} \quad \text{STAR (2017)}$$

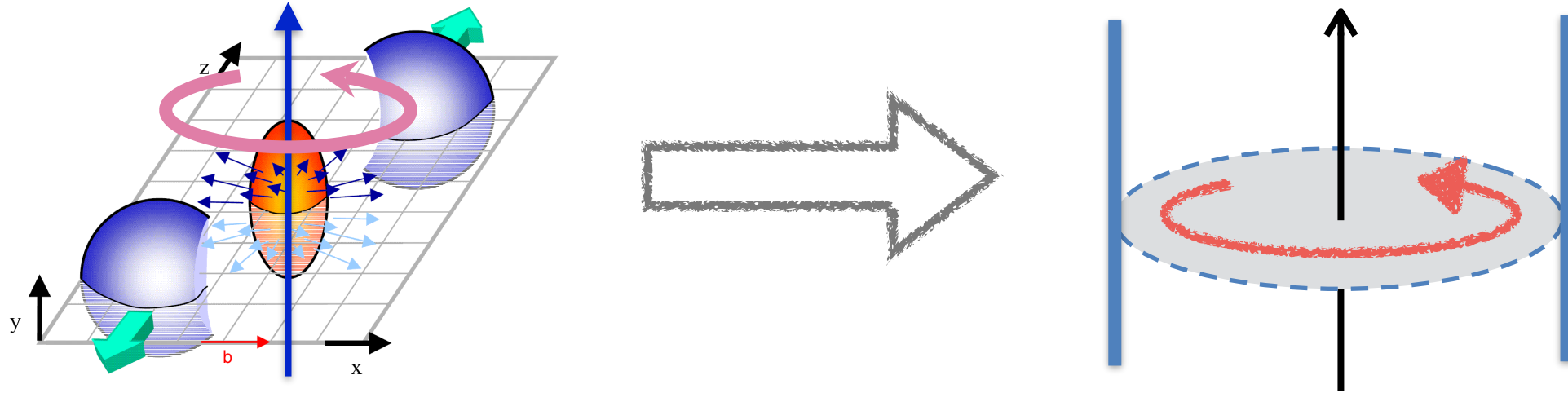
Ampere's law

$$eB \sim 10^{18} \text{ G} \simeq m_{\pi}^2 \quad \text{Skokov, Illarionov, Toneev (2009)}$$

How are the **magnetic** properties modified by **rotation**?

e.g. anomalous transport and chiral symmetry breaking

Rigidly Rotating System

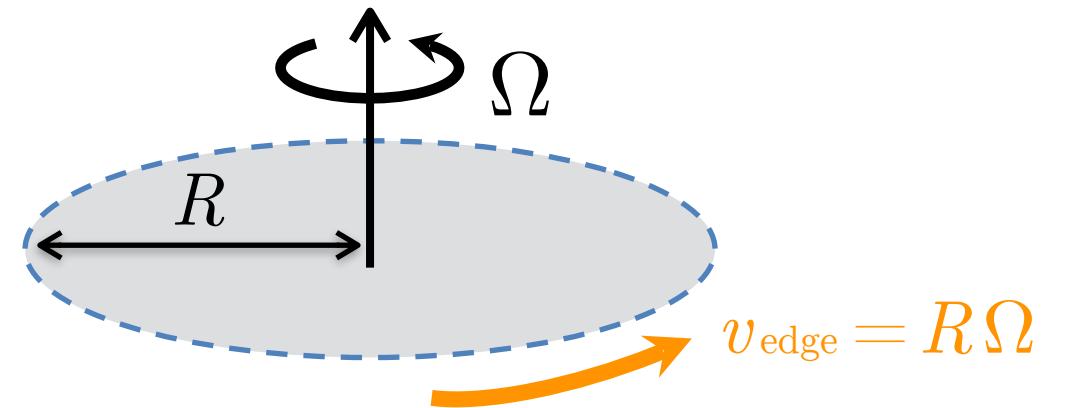


causality constraint

$$v_{\text{edge}} = \Omega R \leq 1$$



$$R \leq 1/\Omega < \infty$$



Rotating systems must be finite-sized

coordinate transformation

$$x_{\text{lab}}^{\mu} \longrightarrow x_{\text{rot}}^{\mu}$$

Contents

1. Finite-size system with Ω

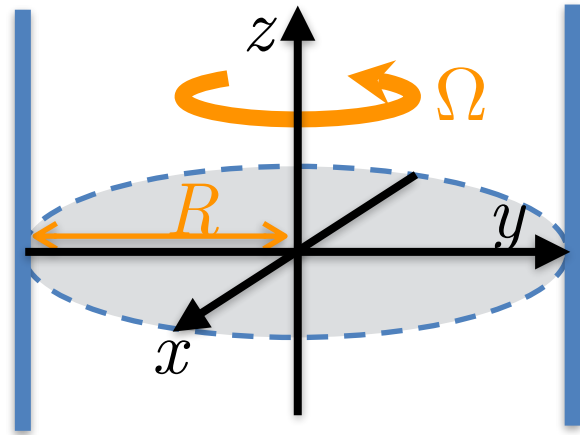
2. Finite-size system with eB

3. Finite-size system with Ω and eB

1. Finite-size system with Ω

Rotating Fermions

$$[i\gamma^\mu (\partial_\mu + \Gamma_\mu) - m]\psi = 0$$



$$g_{\mu\nu} = \begin{pmatrix} 1 - r^2\Omega^2 & y\Omega & -x\Omega & 0 \\ y\Omega & -1 & 0 & 0 \\ -x\Omega & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix}$$

$$\longrightarrow \left[i\gamma^0 \left\{ \partial_0 + \Omega \left(-x\partial_2 + y\partial_1 - \frac{i}{2}\sigma^{12} \right) \right\} + i\gamma^1\partial_1 + i\gamma^2\partial_2 + i\gamma^3\partial_3 - m \right] \psi = 0$$

rotational energy (effective chemical potential)

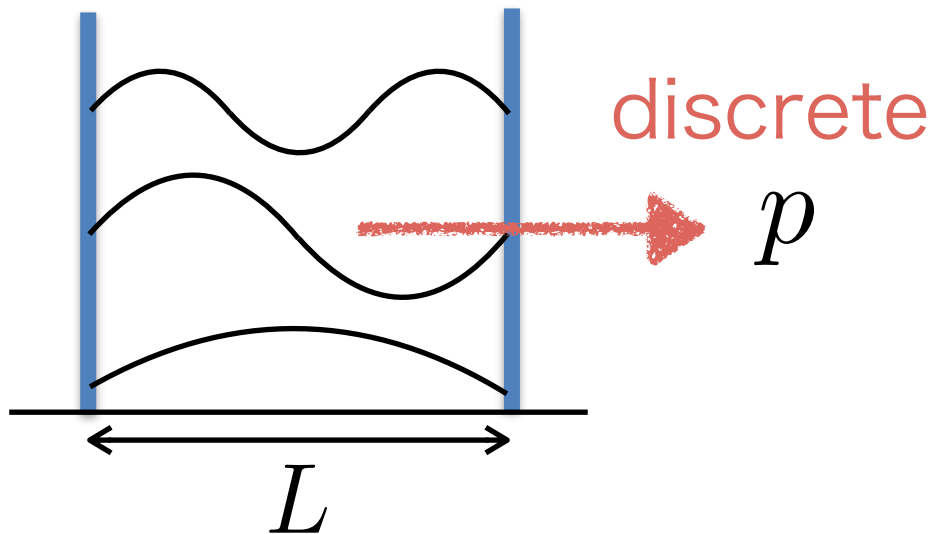
Fermi distribution

$$f(\varepsilon) = \frac{1}{e^{\beta\{\varepsilon - \Omega(l+1/2)\}} + 1}$$

eigenvalue of $\hat{L}_z = -i(x\partial_y - y\partial_x)$

Momentum Discretization

Bosons confined in a well



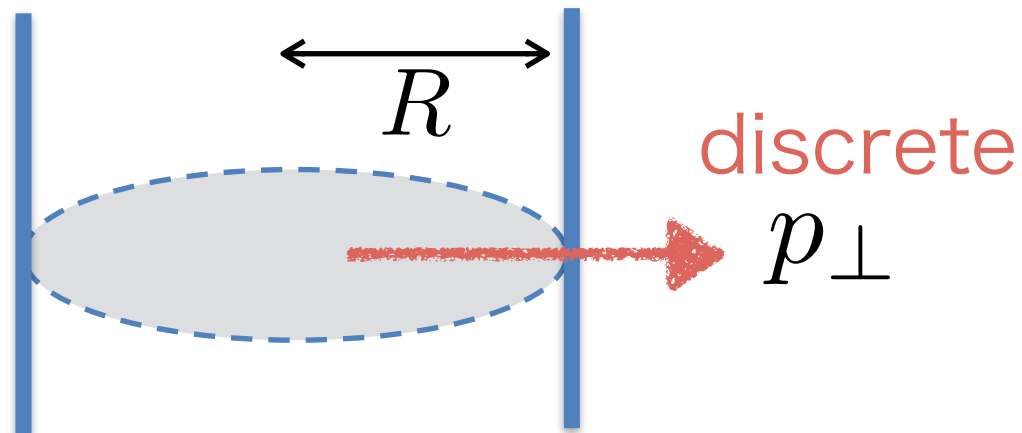
Dirichlet type

$$\sin(px)|_{x=L} = 0$$

$$p = \frac{n\pi}{L} \geq \frac{\pi}{L}$$

IR gap

Fermions confined in a cylinder



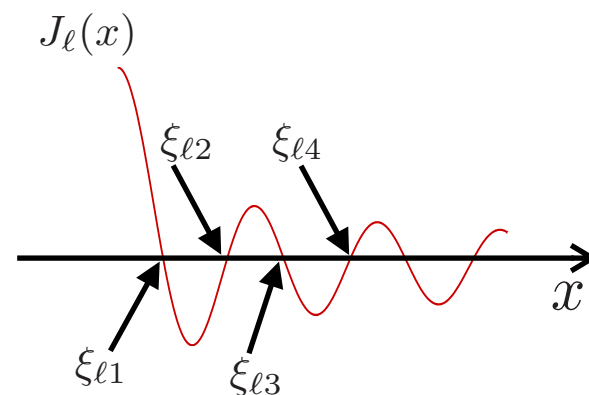
NO incoming current

$$\int \bar{\psi} \gamma^r \psi \Big|_{r=R} = 0$$

$$p_{\perp} = \frac{\xi_{l,k}}{R} \geq \frac{\xi_{l,1}}{R}$$

IR gap

$\xi_{l,k}$: the k th root of $J_l(x)$



Rotational Effect at $T = 0$

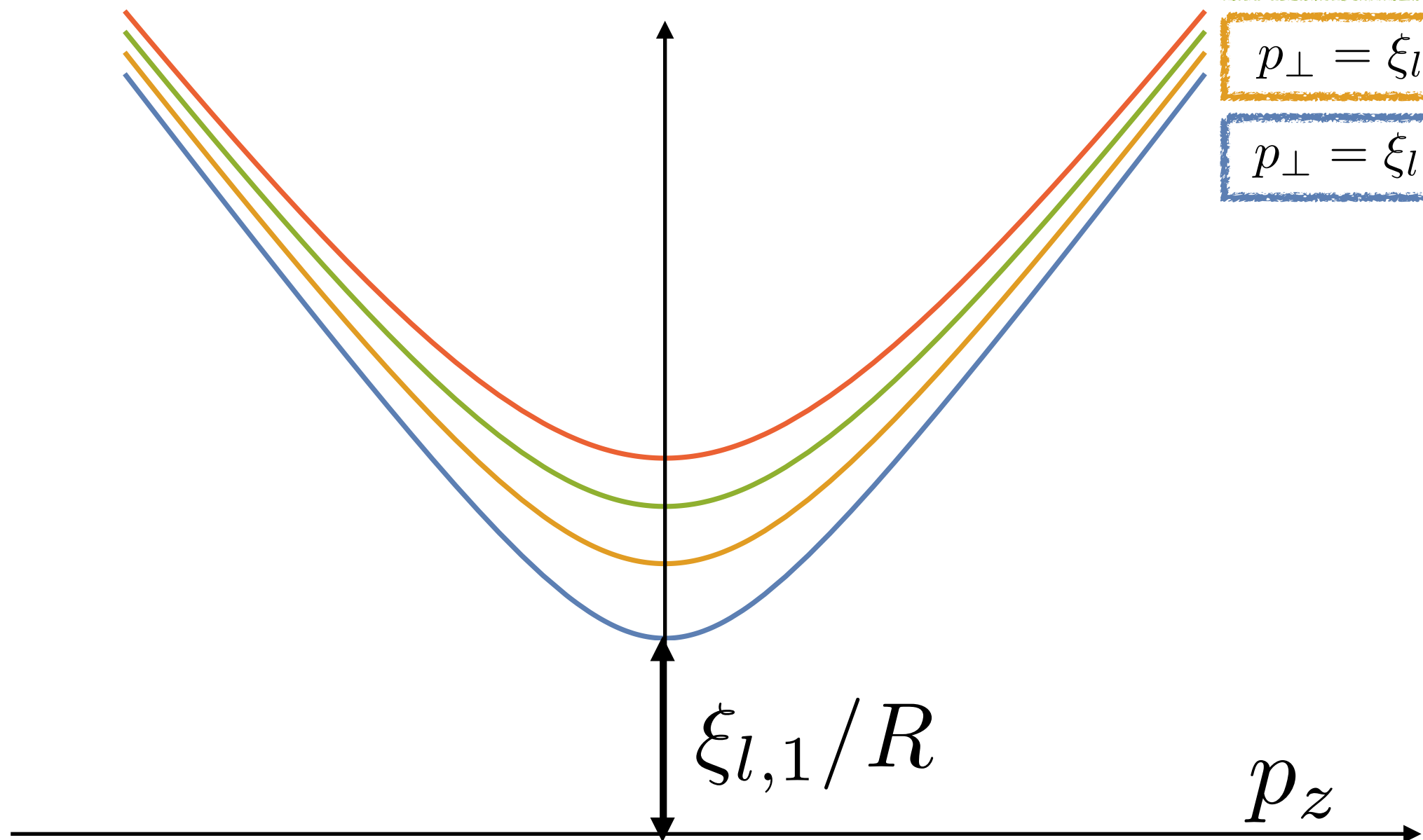
$$\varepsilon = \sqrt{p_{\perp}^2 + p_z^2}$$

$$p_{\perp} = \xi_{l,4}/R$$

$$p_{\perp} = \xi_{l,3}/R$$

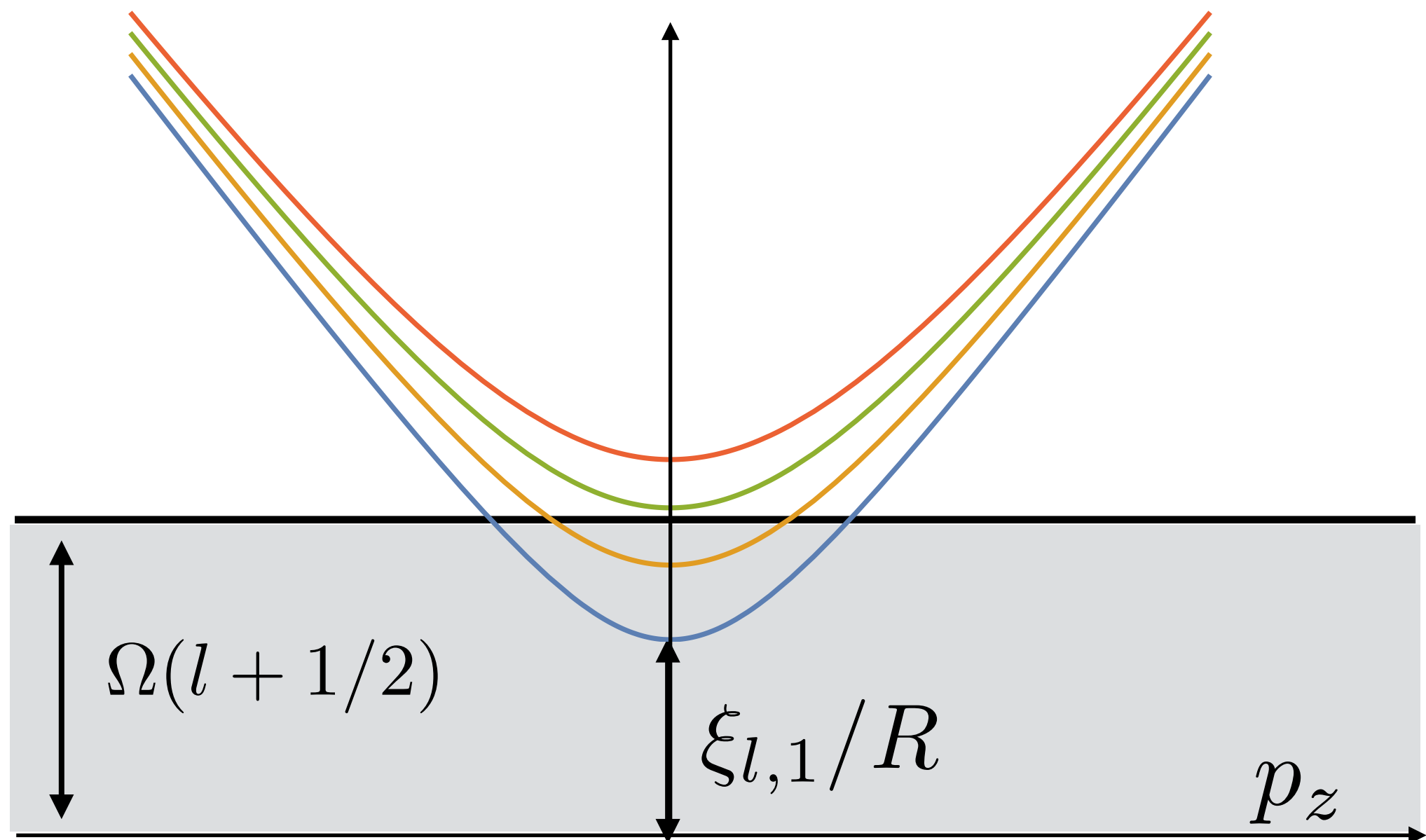
$$p_{\perp} = \xi_{l,2}/R$$

$$p_{\perp} = \xi_{l,1}/R$$



Rotational Effect at $T = 0$

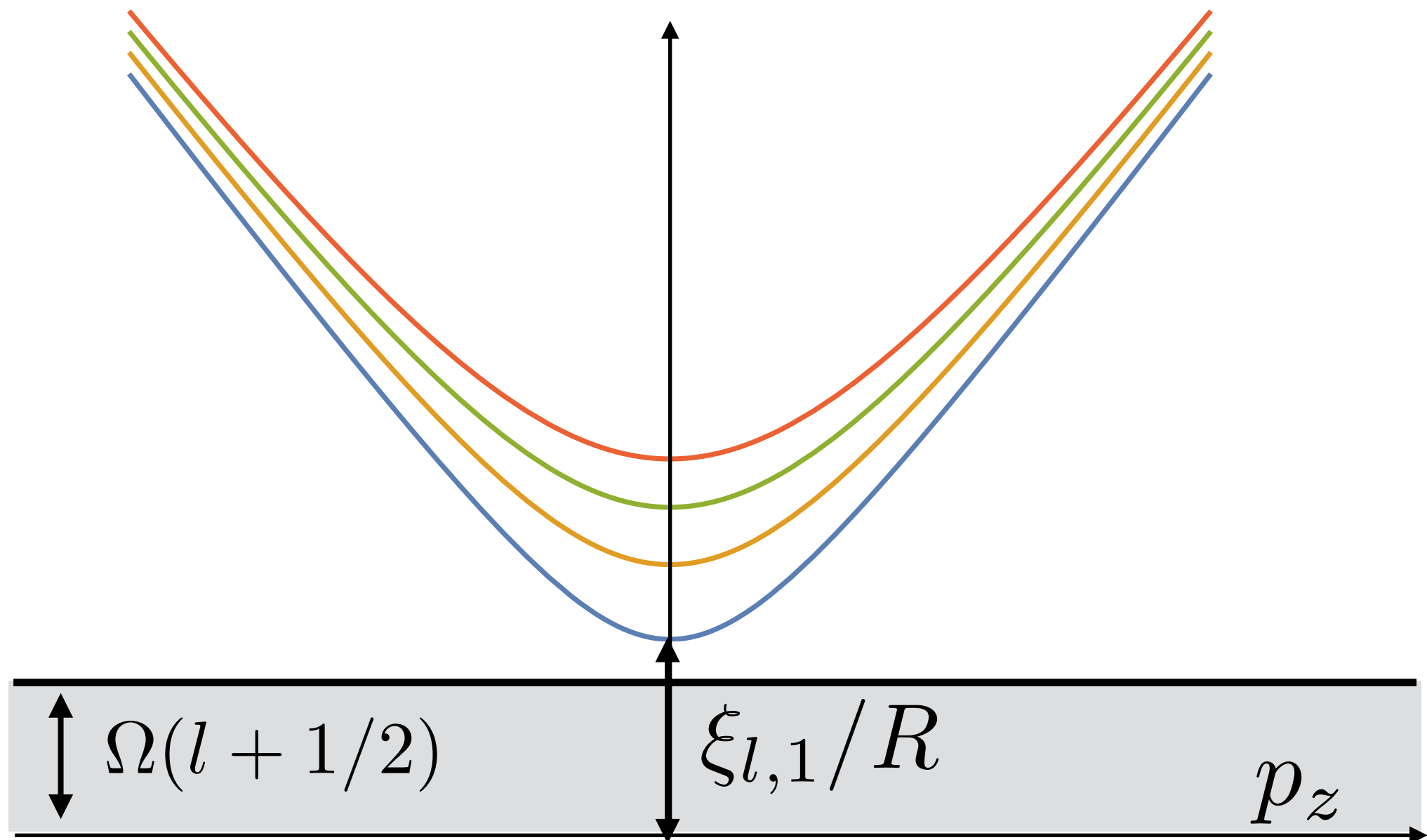
$$f(\varepsilon) = \frac{1}{e^{\beta\{\varepsilon - \Omega(l + 1/2)\}} + 1} \longrightarrow f(\varepsilon) = \theta(\Omega(l + 1/2) - \varepsilon)$$



visible

Rotational Effect at $T = 0$

$$f(\varepsilon) = \frac{1}{e^{\beta\{\varepsilon - \Omega(l + 1/2)\}} + 1} \longrightarrow f(\varepsilon) = \theta(\Omega(l + 1/2) - \varepsilon)$$



invisible

cf. Silver Blaze

Rotational Effect at $T = 0$

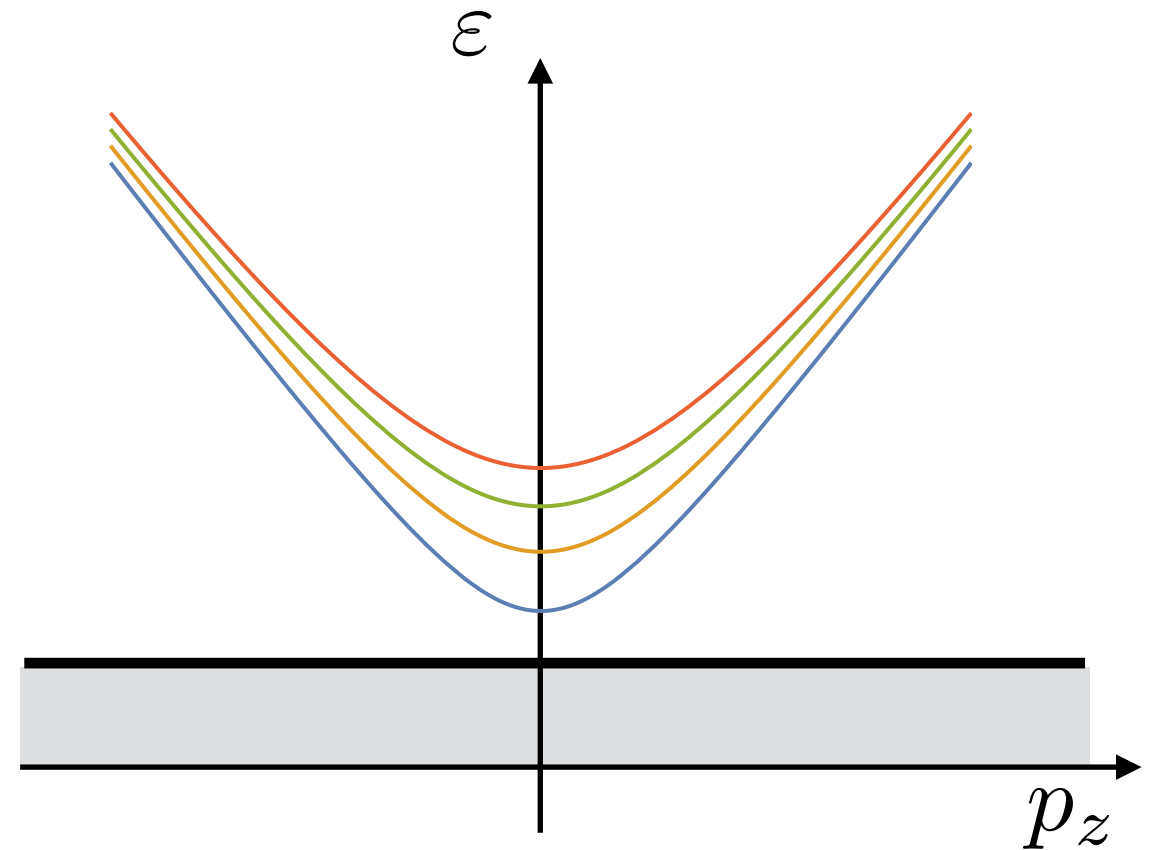
Ebihara, Fukushima, KM (2017)

causality $\Omega \leq 1/R$

$$\xi_{l,1}/R > \Omega(l + 1/2)$$

for arbitrary l

NO rotational effect at $T = 0$



* same discussion with the MIT b.c. Chernodub, Gongyo (2017)

Note : visible at finite temperature

$$\text{CVE} \quad j_5 = \frac{T^2}{12} \Omega$$

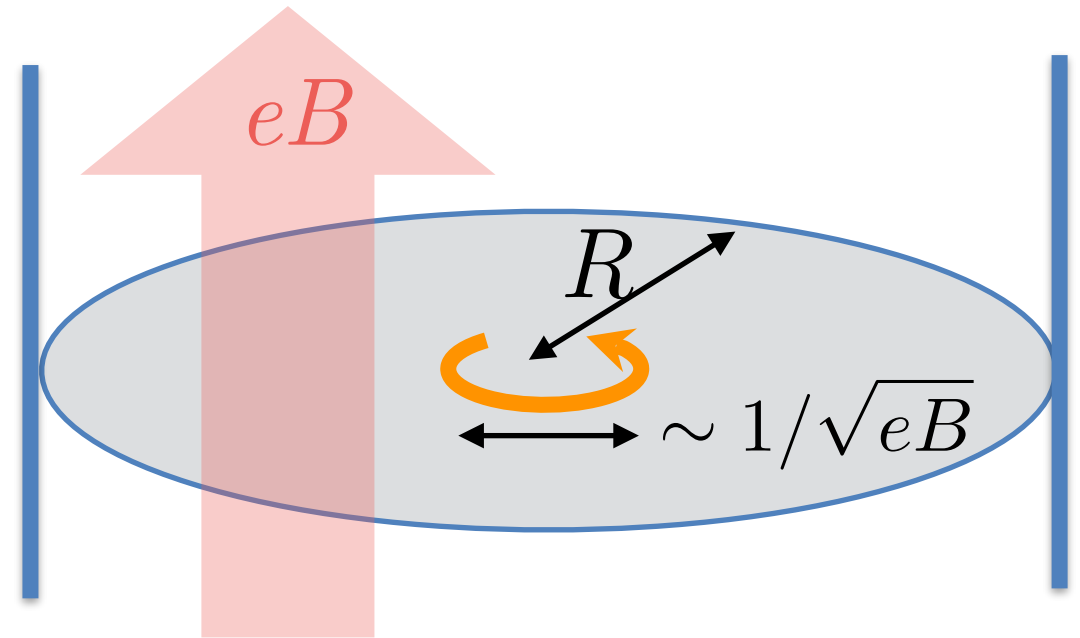
Vilenkin (1979)

2. Finite-size system with eB

Cyclotron Motion

(1) $1/\sqrt{eB} \ll R$

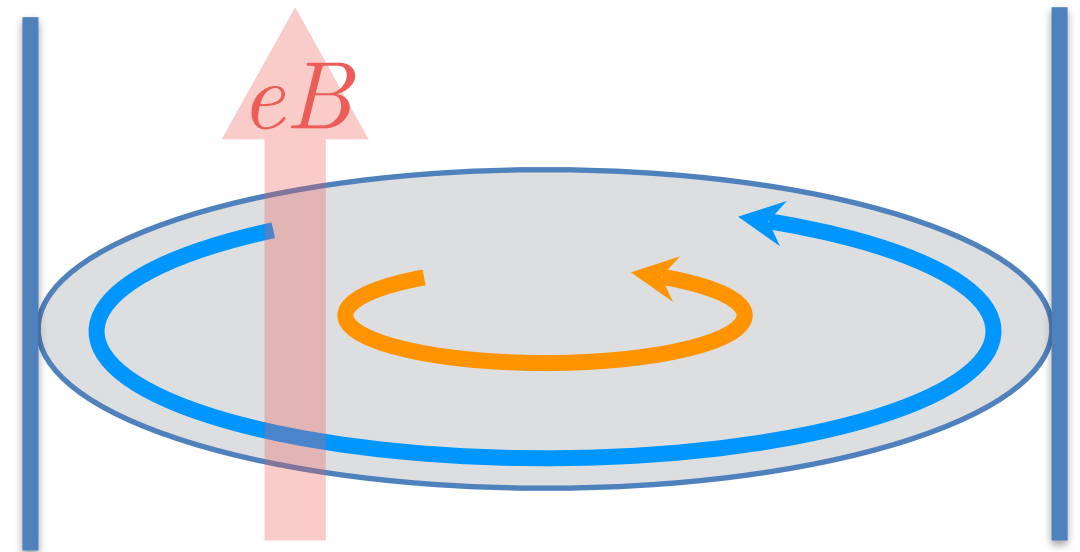
$\longrightarrow p_{\perp} = \sqrt{2neB}$
independent of l



(2) $1/\sqrt{eB} \lesssim R$

small $l \longrightarrow$ still $p_{\perp} \simeq \sqrt{2neB}$

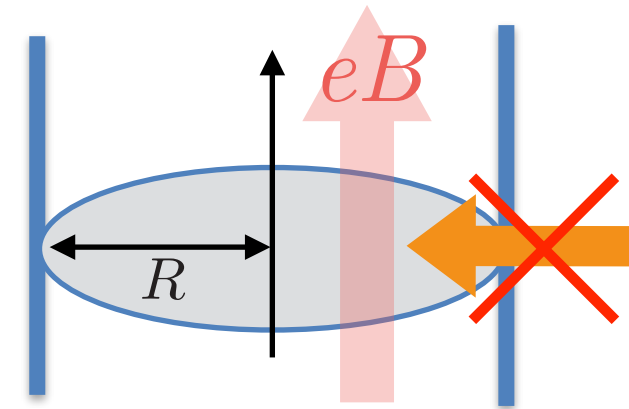
large $l \longrightarrow$ modified



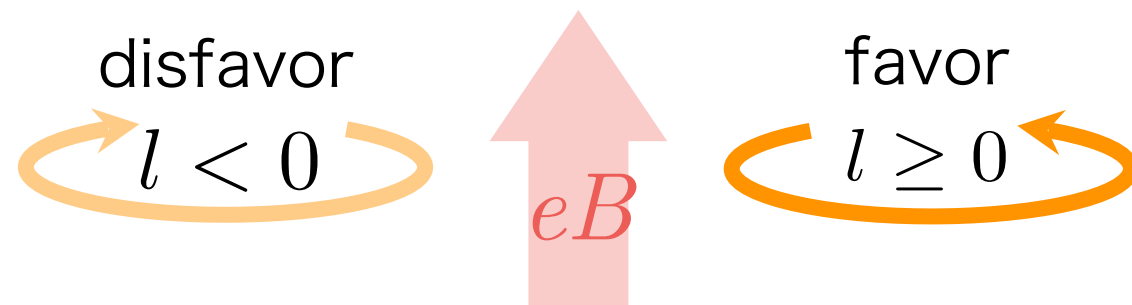
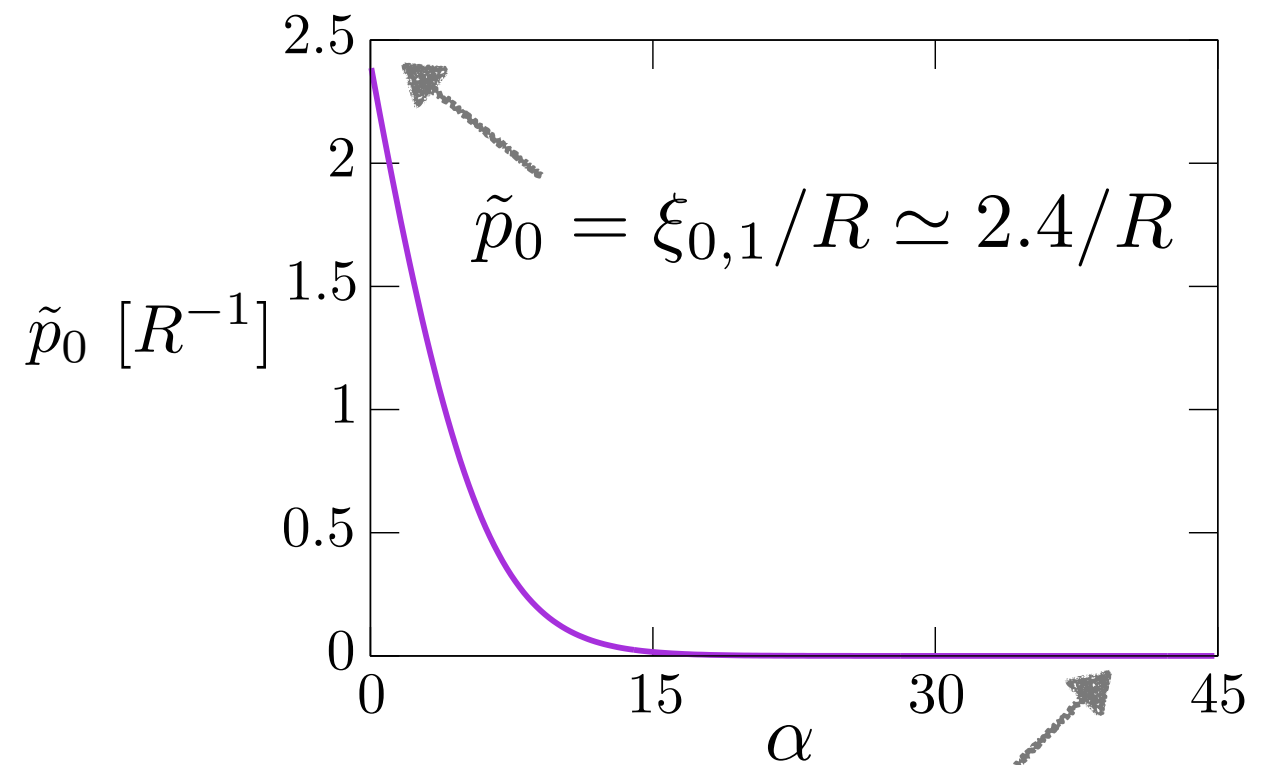
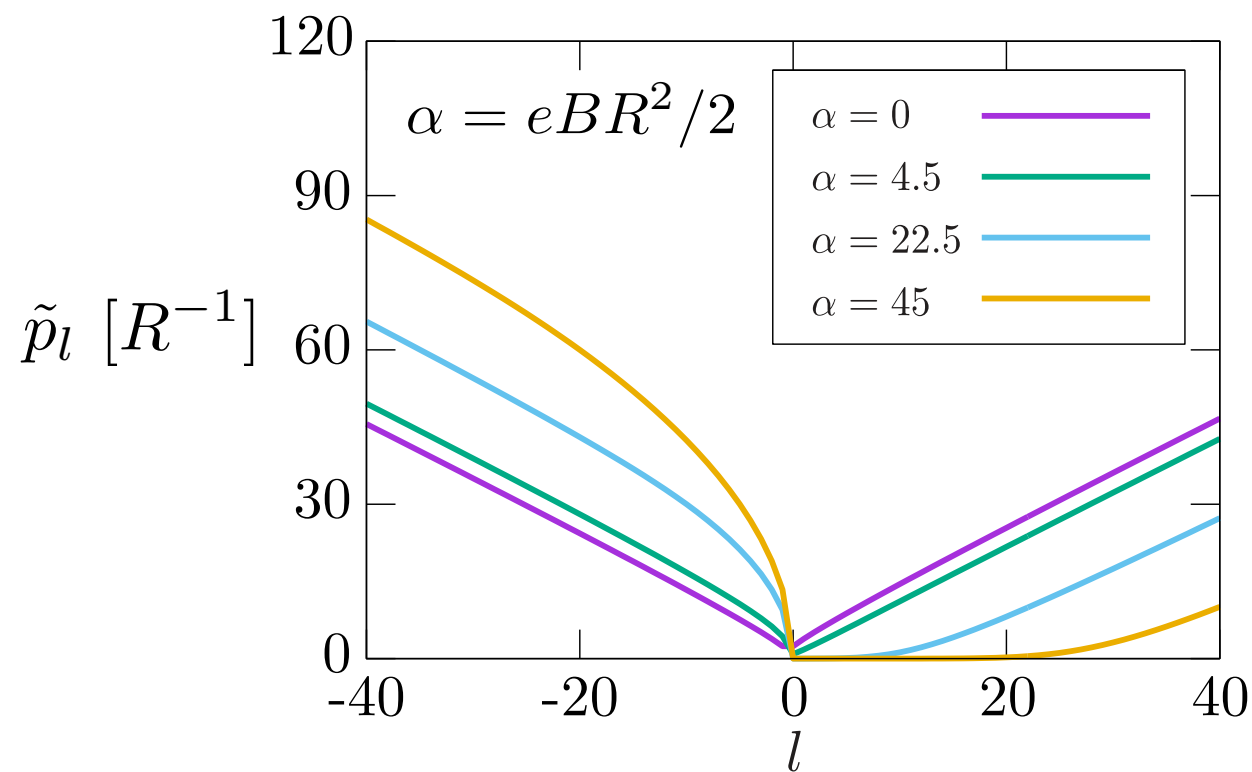
“Incomplete Landau quantization”

Incomplete Landau Level

$$[i\gamma^\mu(\partial_\mu + ieA_\mu) - m]\psi = 0 \text{ with}$$



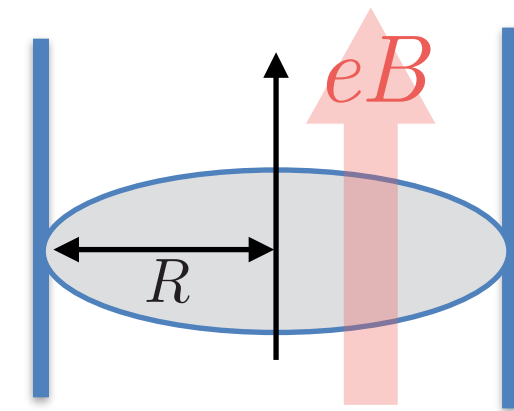
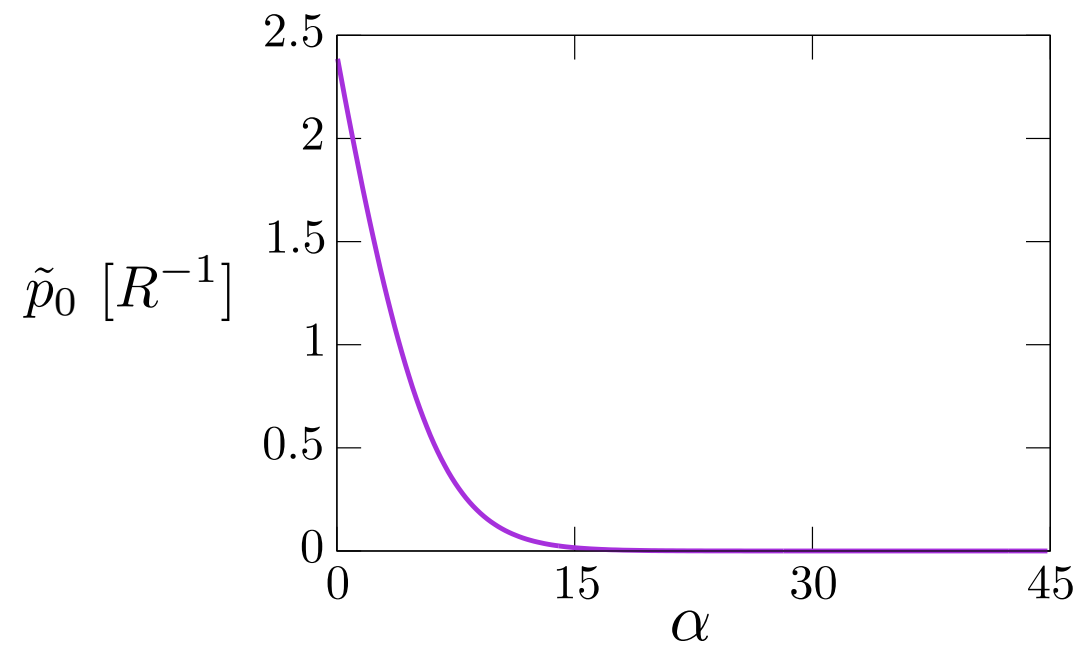
\tilde{p}_l = lowest transverse momentum for l



$\tilde{p}_0 = 0$
Landau zero mode

3. Finite-size system with Ω and eB

Gapped to Gapless

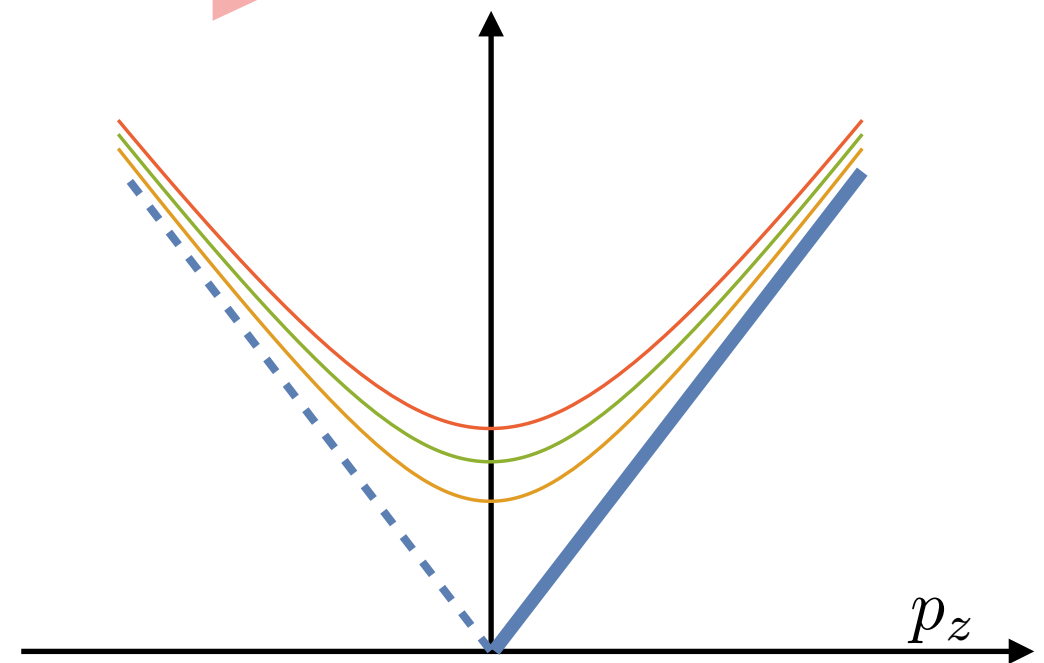
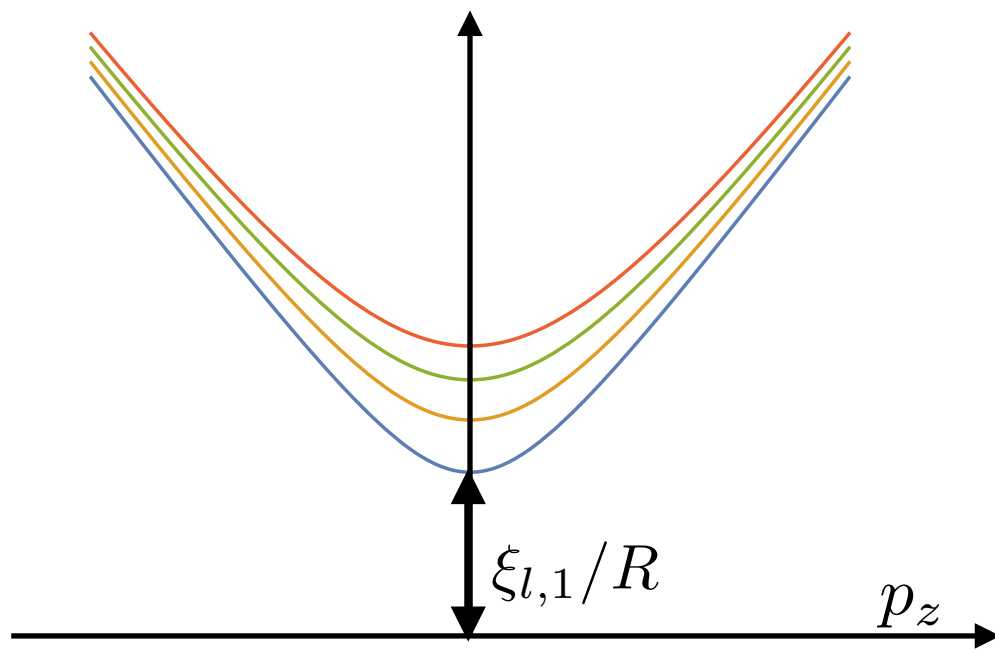


$$\alpha = eBR^2/2$$

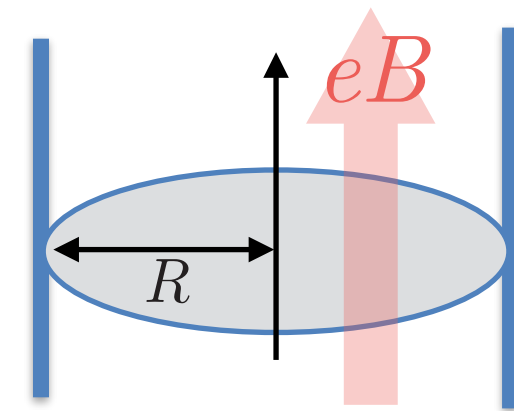
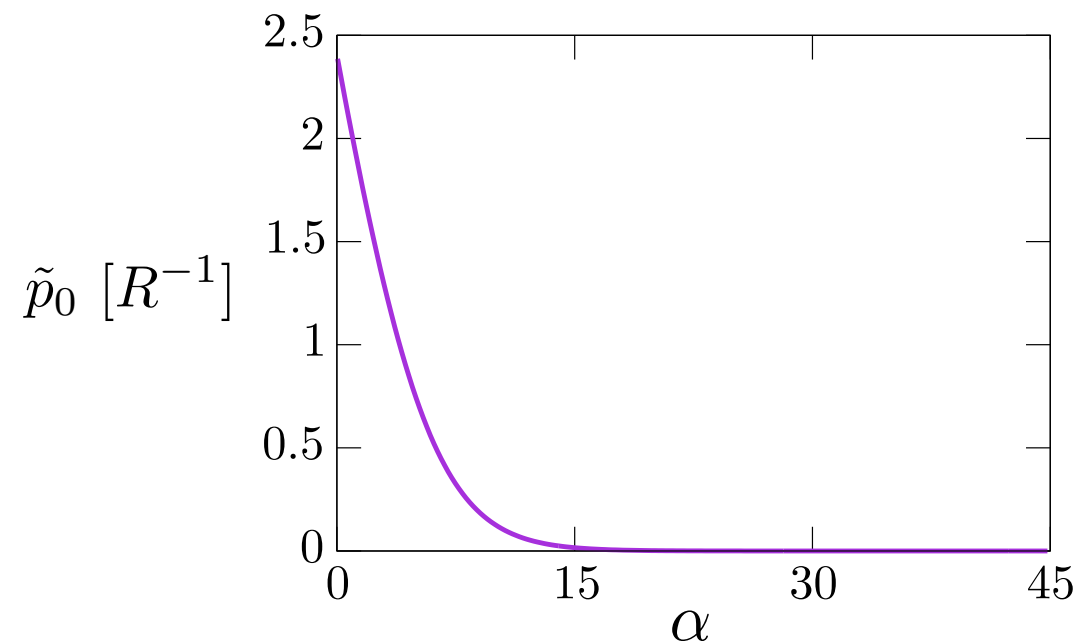
weak

magnetic field

strong

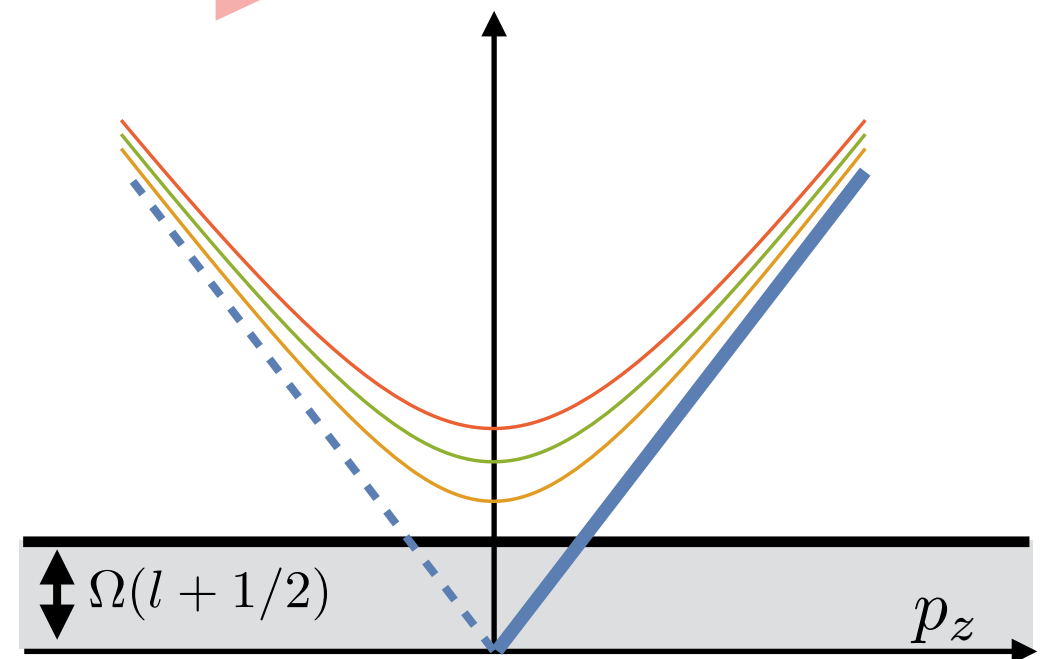
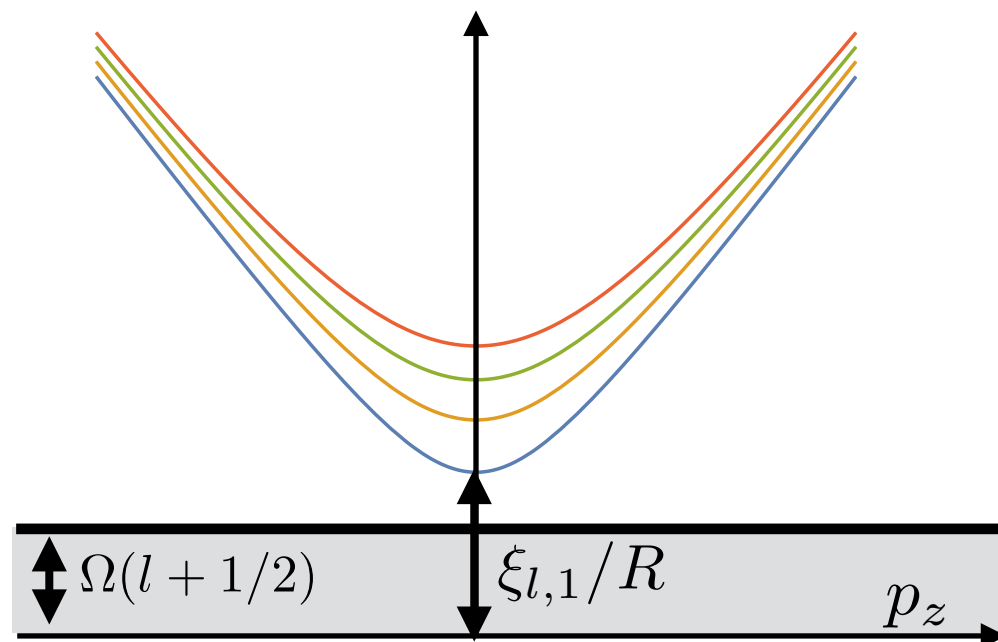


Gapped to Gapless



$$\alpha = eBR^2/2$$

weak magnetic field strong



visible rotational effect due to **magnetic field**

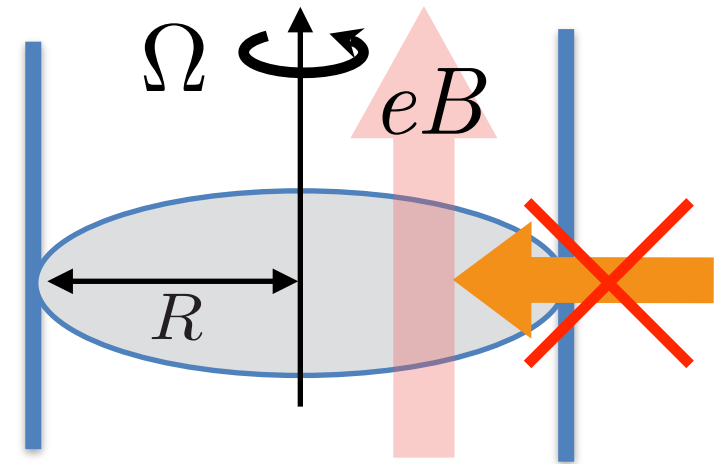
Ex.1) Anomalous Density

Ebihara, Fukushima, KM (2017)
Hattori, Huang, KM (in preparation)

magnetic field

rotation

$$[i\gamma^\mu (\partial_\mu + ieA_\mu + \Gamma_\mu)]\psi = 0 \quad \text{with}$$



$$n(r) = \langle \psi^\dagger(x)\psi(x) \rangle \quad f_\pm(\varepsilon) = \frac{1}{e^{\beta(\varepsilon \mp \Omega j)} + 1}$$

$$= \sum_{p_z, p_\perp} [f_+(\varepsilon) - f_-(\varepsilon)] \times (r\text{-dependence}) \longrightarrow n(r=0) \xrightarrow{\sqrt{eB} \gg \Omega} \frac{eB\Omega}{4\pi^2}$$

temperature independent

Cf. Hattori, Yin (2016)

imply a **tensorial** anomalous term

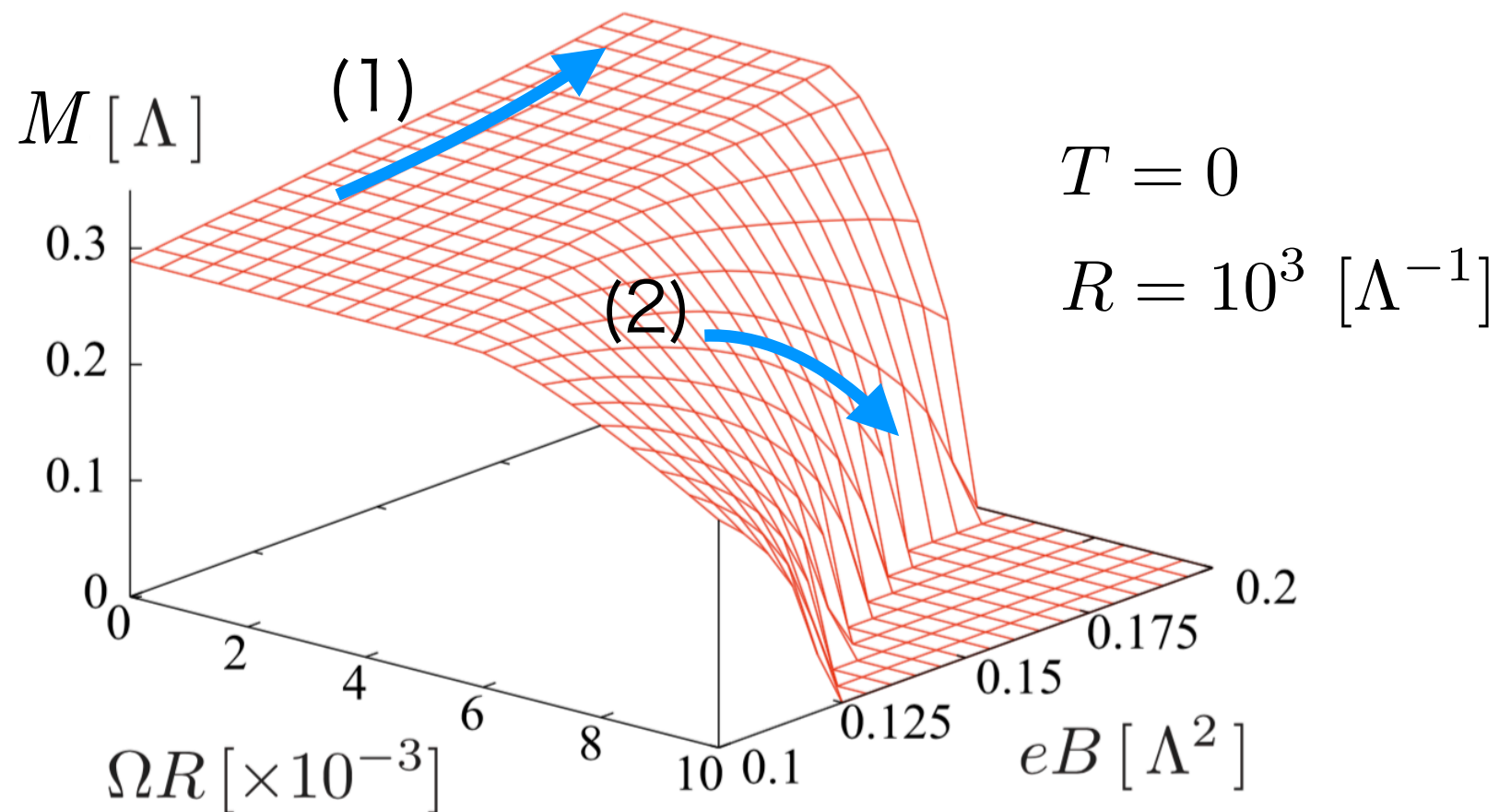
$$\Delta S \sim \mu \mathbf{B} \cdot \boldsymbol{\Omega} = \mu g_{ij} B^i \Omega^j \xrightarrow{\frac{\delta \Delta S}{\delta g_{ij}}} T^{ij} = \# e B^i \Omega^j$$

Cf. Hernandez, Kovtun (2017)

Ex.2) Chiral Symmetry

Chen, Huang, Fukushima, KM (2016)

NJL model (mean field approx.) + neglecting inhomogeneity



(1) eB increases \longrightarrow M increases

Magnetic Catalysis

(2) eB increases \longrightarrow M decreases

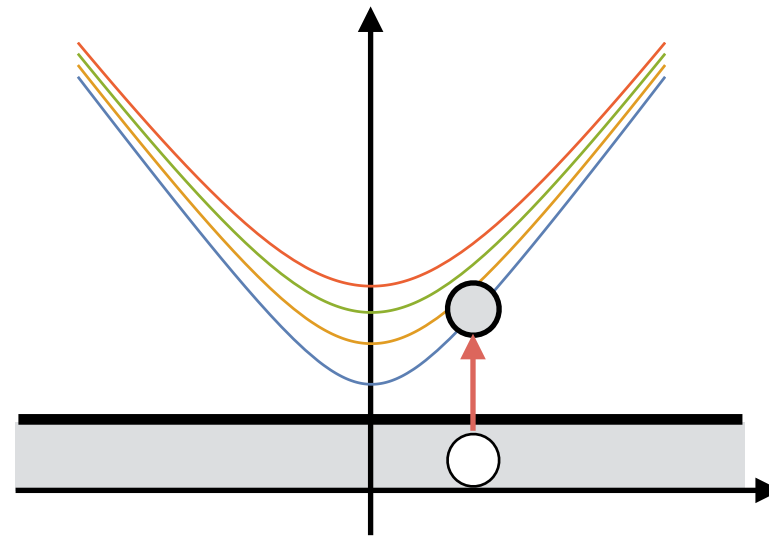
Inverse of MC

‘Rotational magnetic inhibition’

Summary: Rotation effect requires ...

- **temperature**

thermal excitation

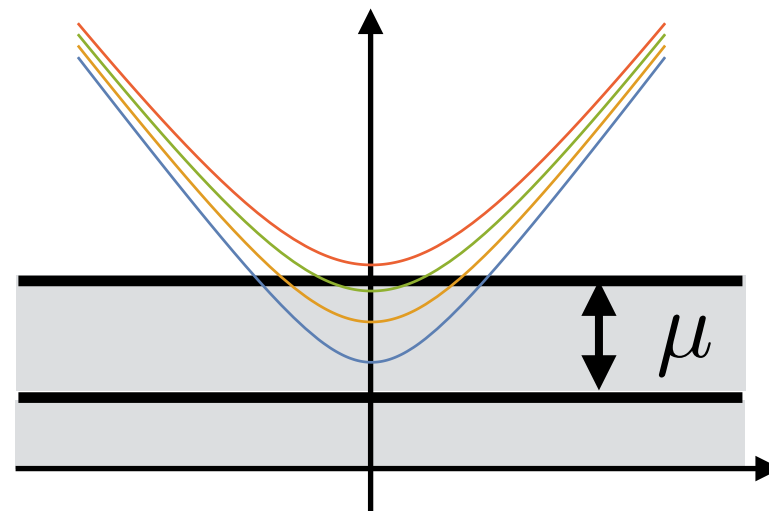


Jiang, Liao (2016)

Chernodub, Gongyo (2016)

- **density**

Fermi surface shift

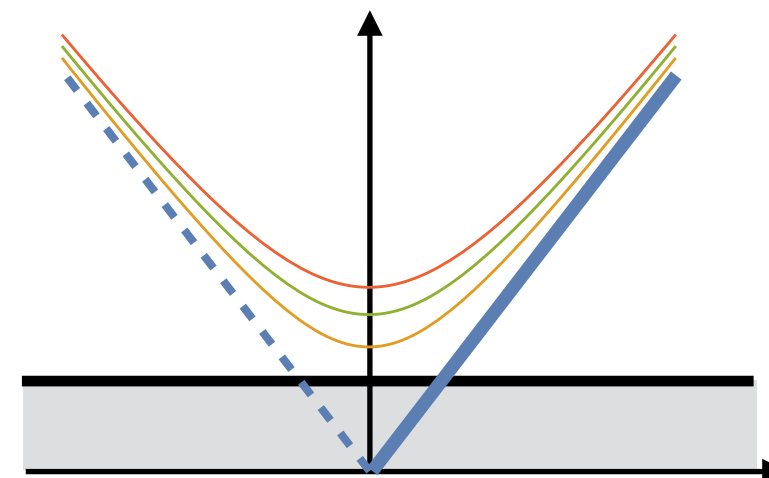


Huang, Nishimura,

Yamamoto (2017)

- **magnetic field**

Landau zero mode



Chen, Huang,

Fukushima, KM (2016)

Liu, Zahed (2017)

Outlook

- Rotation yields abundant phase structures
Chen, Huang, Fukushima, KM (2016)
Jiang, Liao (2016) Chernodub, Gongyo (2016)
Liu, Zahed (2017) Huang, Nishimura, Yamamoto (2017)
- Rotational magnetic inhibition : spacial dependent chiral condensate
- Novel anomalous (magneto-vortical) correction $T^{ij} = \# e B^i \Omega^j$
Hattori, Huang, KM (in preparation)
- Finite-size system under B -field **EdH effect for chiral fermions**
Fukushima, Hirono, Huang, Kharzeev, KM (in preparation)