

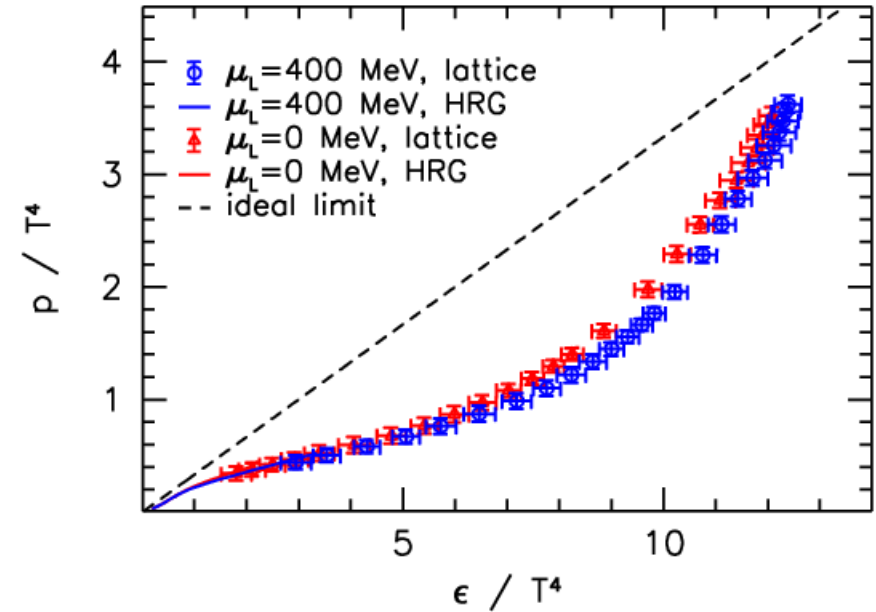
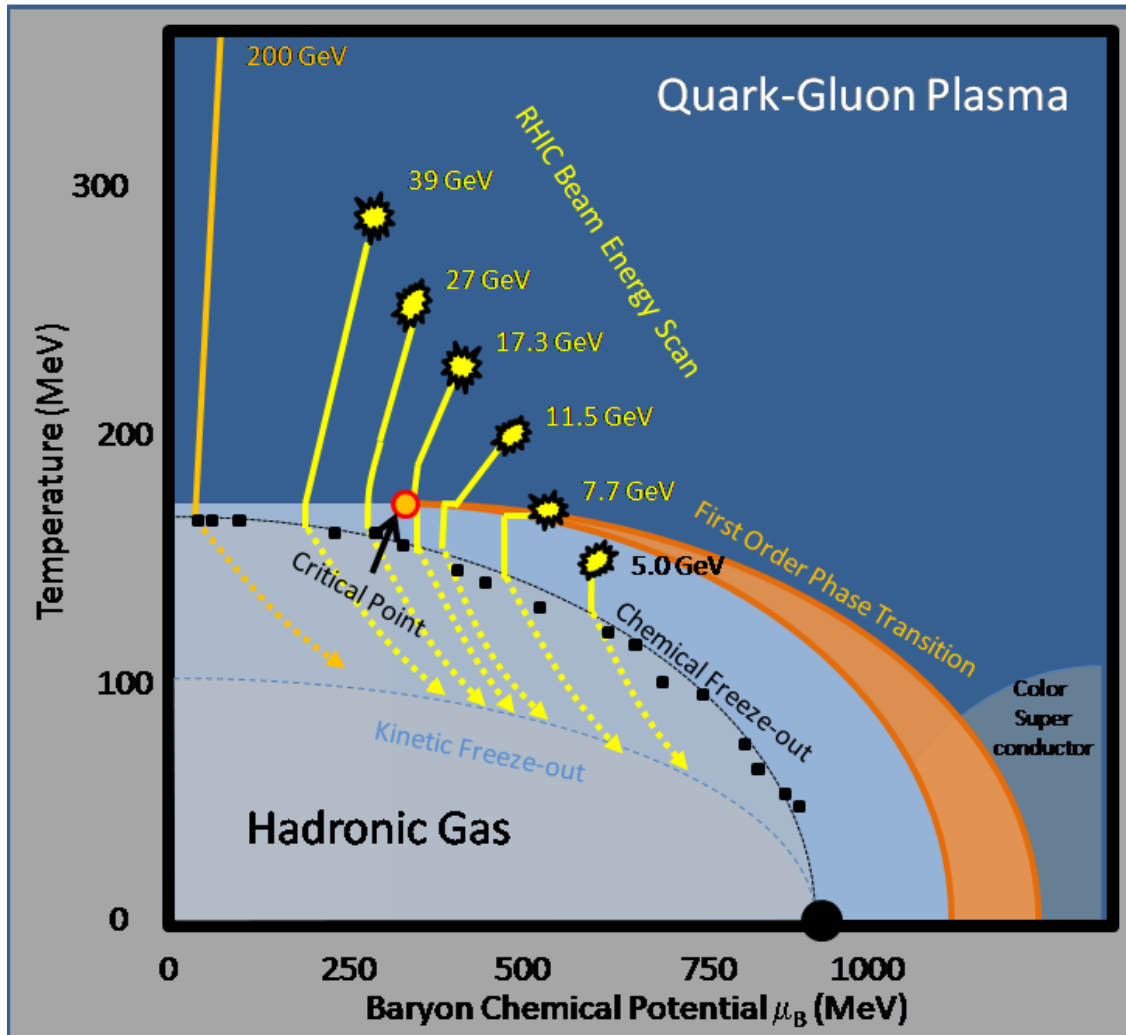
K/pi ratios from a dynamically integrated transport approach

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- Introduction
- JAM+hydro: A new dynamically integrated transport model
- Equation of state (EoS) controlled collision term

This talk is based on arXiv:1805.09024

Search for the QCD equation of state (EoS) by the beam energy scan



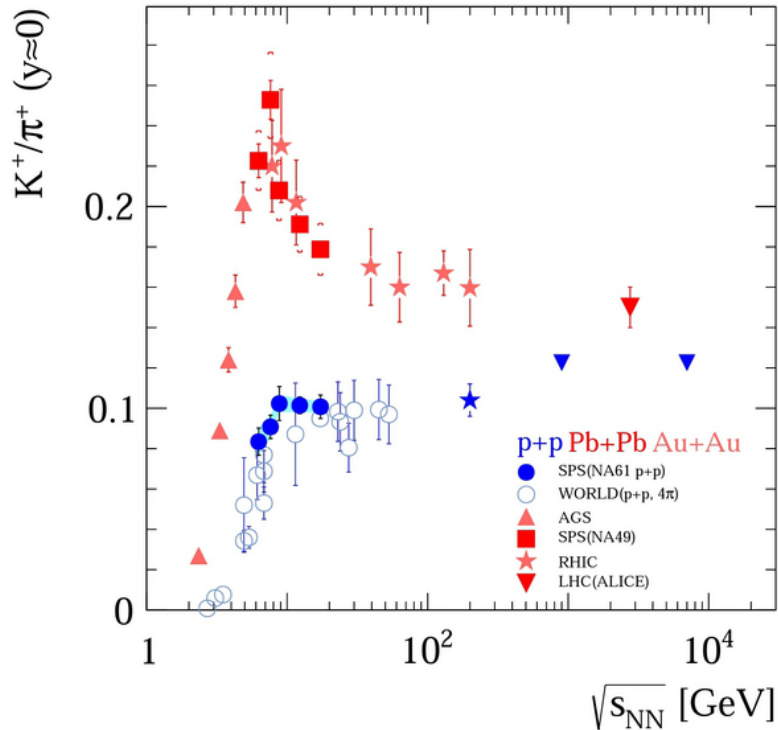
EoS from lattice QCD
Sz.Borsanyi, et.al JHEP 1208(2013)053

Effective models:

NJL, PNJL, PQM,
Quasi-particle model.....

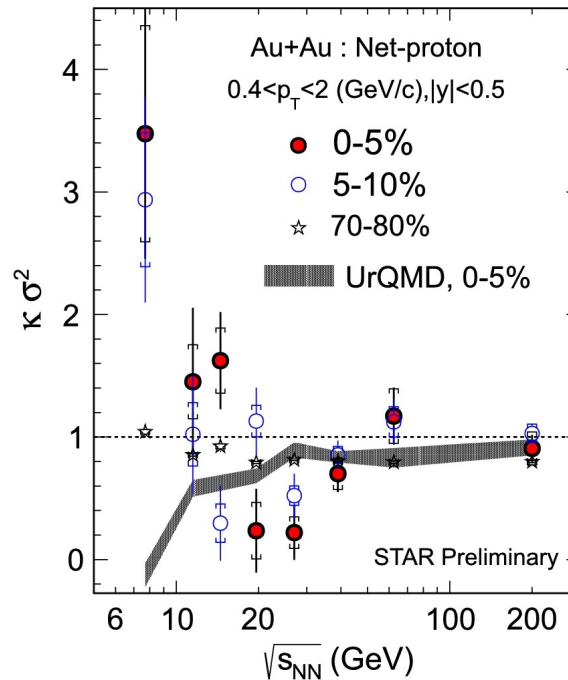
Location of the critical point?

Non-monotonic structures in beam energy dependence



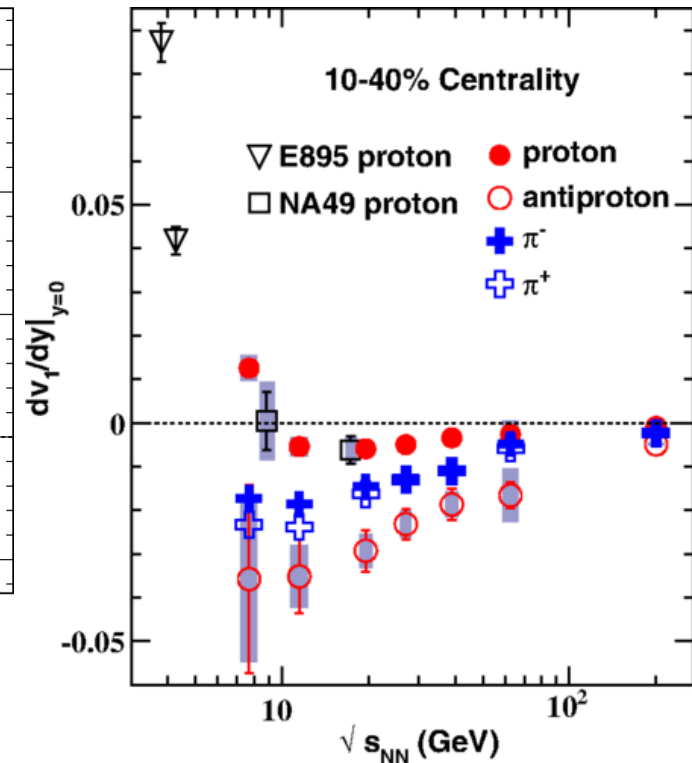
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Onset of de-confinement?



X. Luo QM15

First-order phase transition? End point?



Modeling at RHIC/LHC hybrid approach works

Initial conditions: Glauber, CGC, event generators

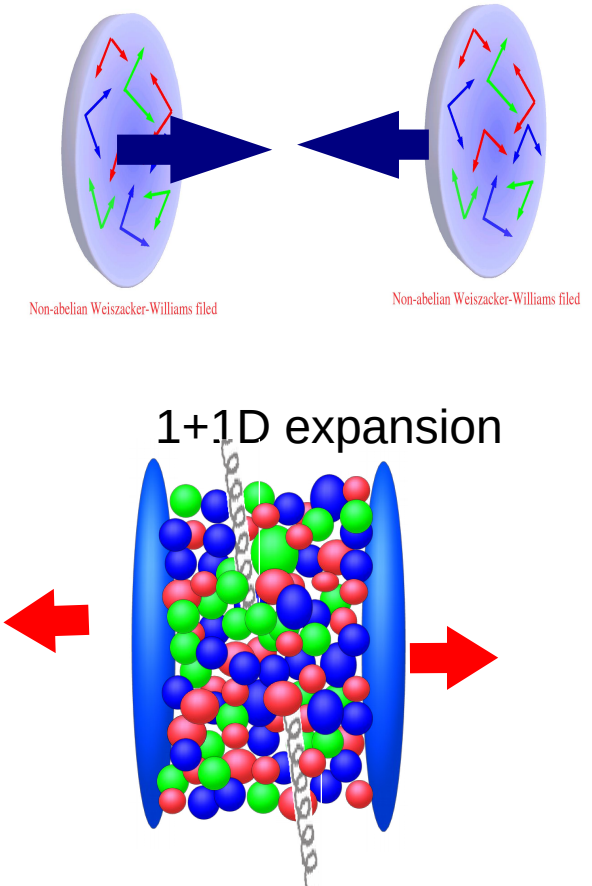
Hydrodynamics: viscous-hydro, anisotropic hydro

Hadron gas: hadron transport models

EoS: crossover from lattice QCD

Jet production and its Energy loss: pQCD

Single thermalization time at $\tau = 0.6 \text{ fm}/c$

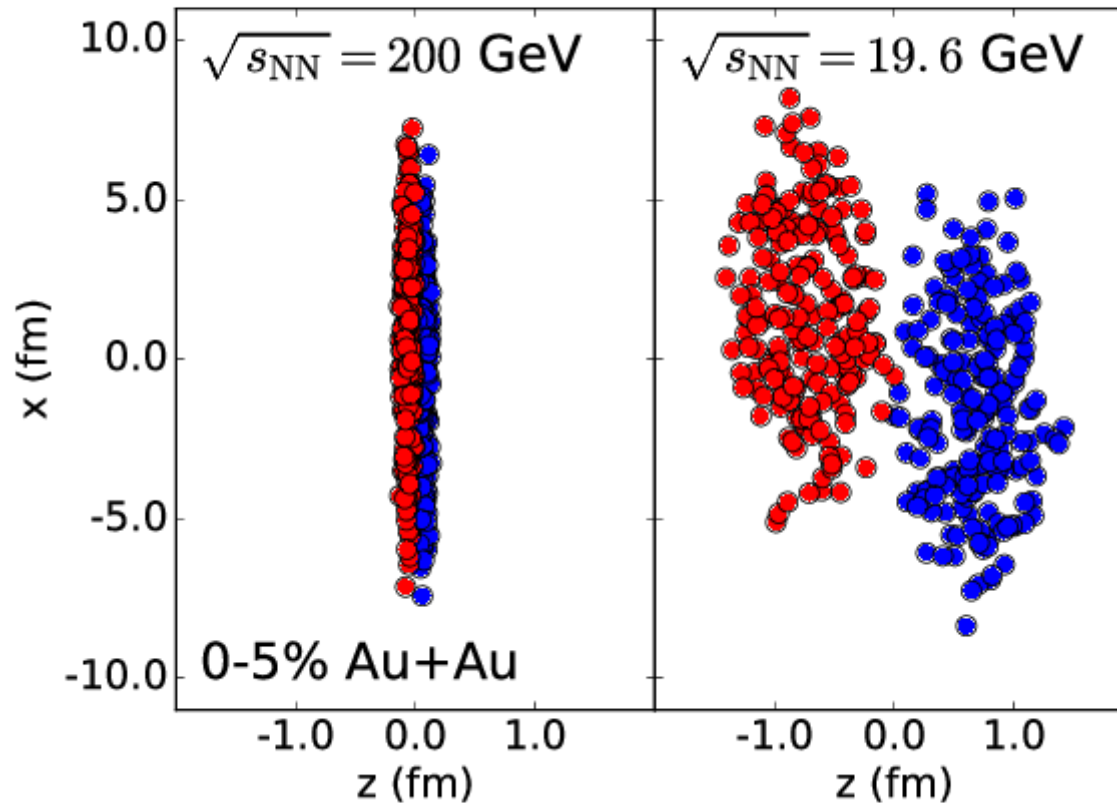


At high energies, factorization in time and energy works:

e.g. CGC + hydrodynamics + energy loss of jets + hadron transport model

Initial nucleon positions

Chun Shen, Bjorn Schenke arXiv:1710.00881 [nucl-th]

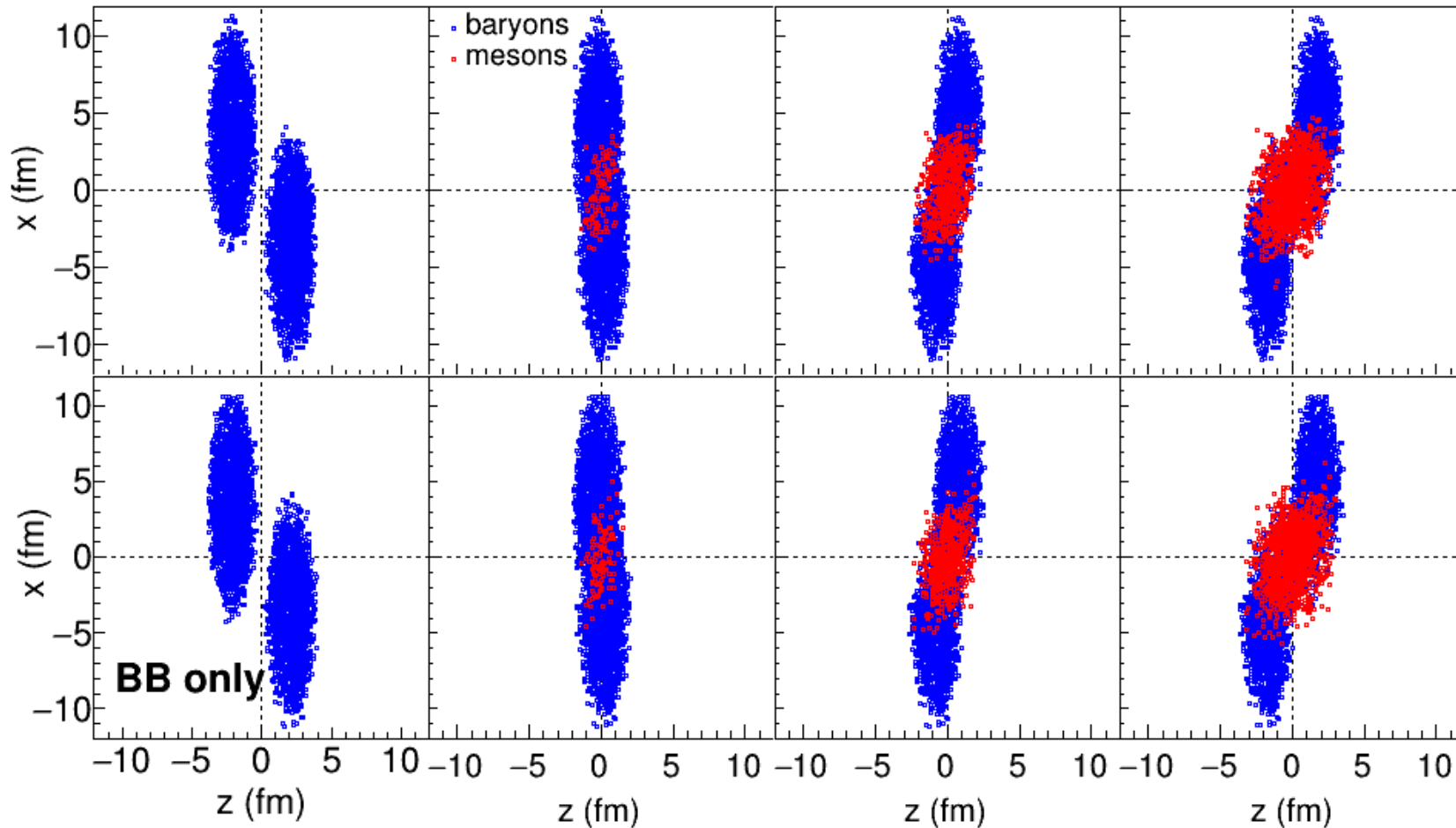


Assumption of single thermalization time breaks down at low beam energies

Time evolution

Red:mesons. Meson-baryon interactions are important before two Nuclei pass each other.

$$\sqrt{s_{NN}} = 7.7 \text{ GeV}$$



What is needed for HI at high baryon density

Cannot apply factorization in time;
models at RHIC/LHC does not work

1) Hybrid approach:

coupling of Hydro + Boltzmann

2) Initial stages of collisions have to be described by the non-equilibrium transport models

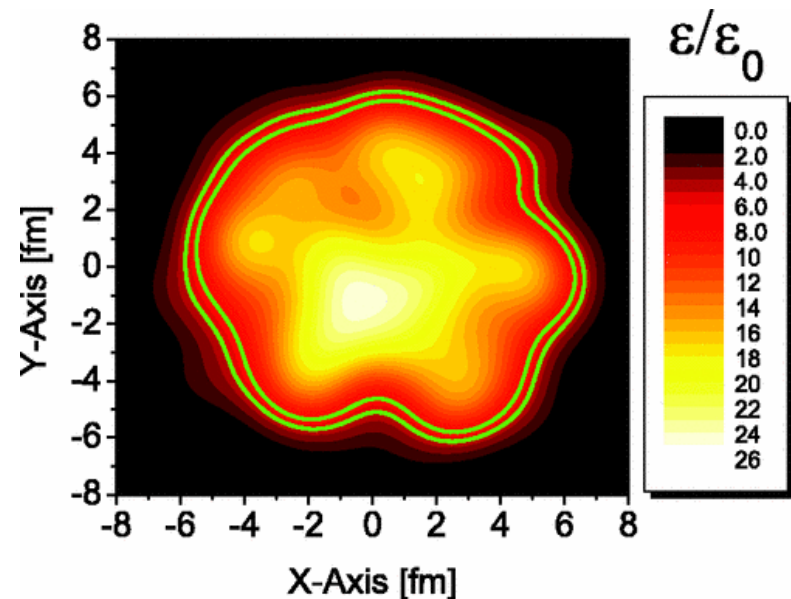
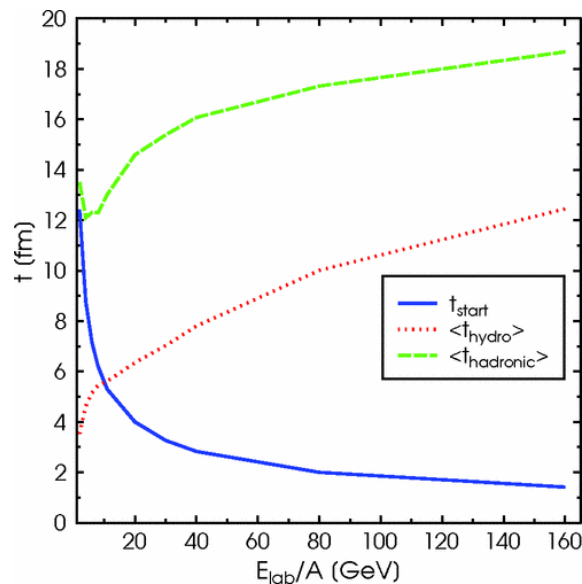
3) How to model EoS with non-vanishing chemical potentials?

Progresses in hybrid dynamical models at high baryon density region

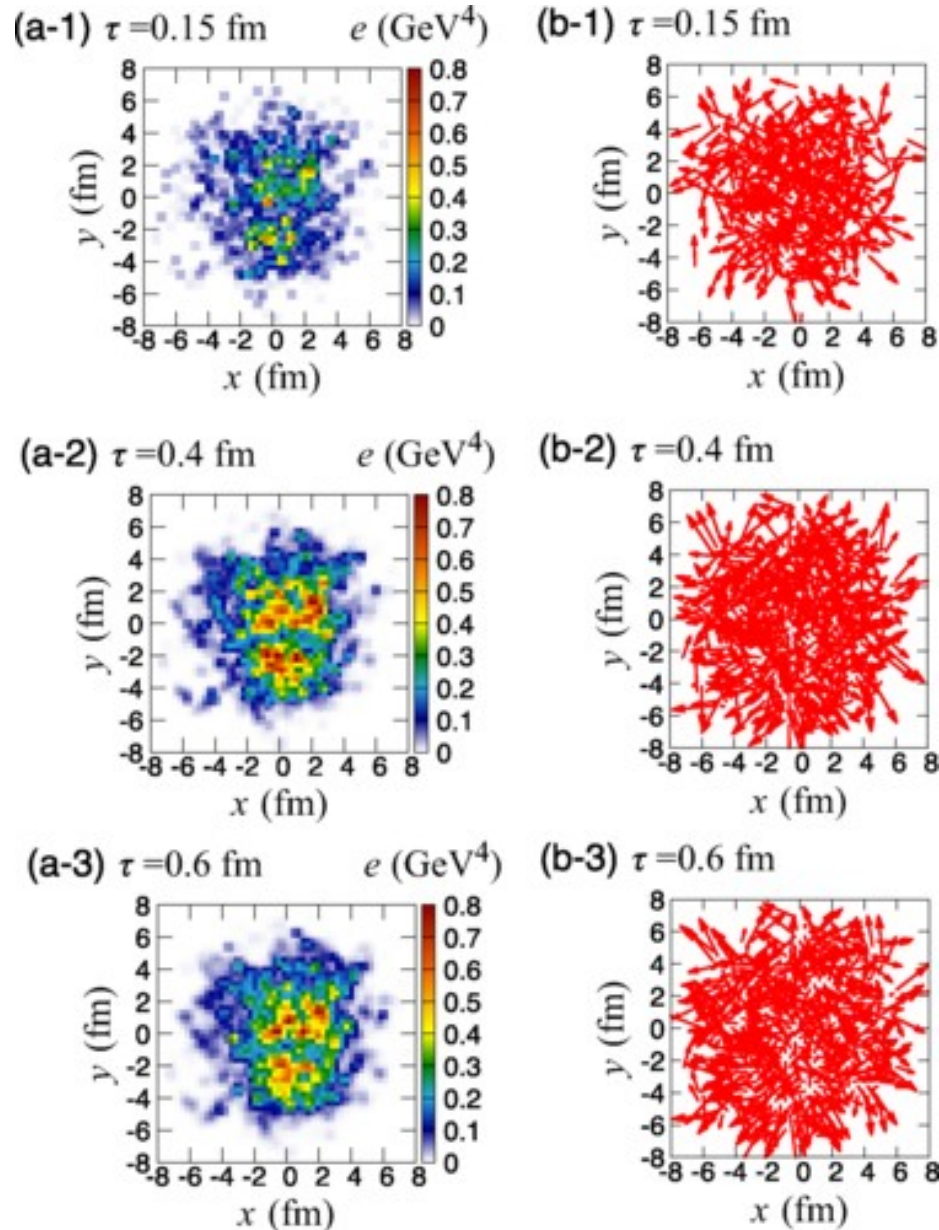
- UrQMD+hydro (2008) Petersen et.al.

initial condition from UrQMD → hydrodynamics → UrQMD

- Core-corona separation (K, Werner, 2007) in space configurations Steinheimer and Bleicher (2011)
- Dynamical initialization (2018) Shen and Schenke core-corona separation in time direction.



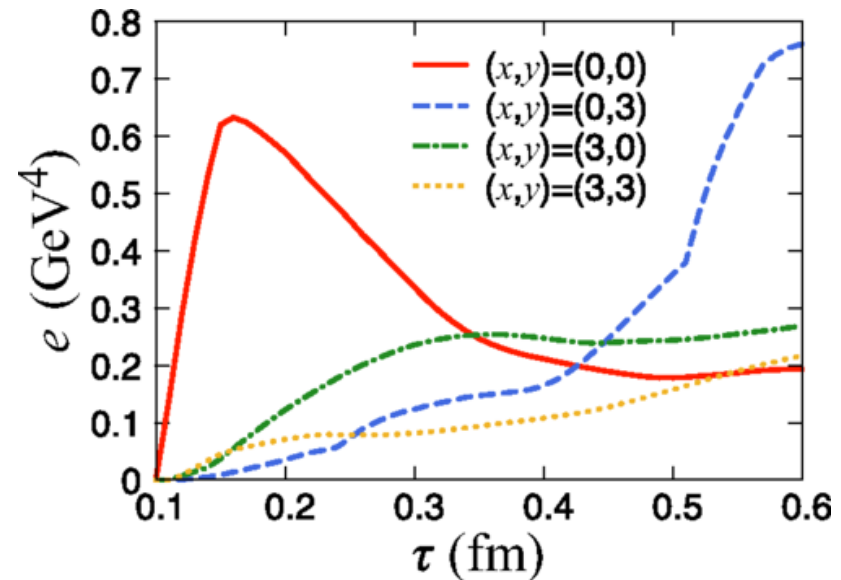
Dynamical Initialization of fluid at LHC



Dynamical initial condition for hydrodynamics as introduced by M. Okai, et. al Phys. Rev C 95, 054914 (2017)

$$\partial_{\mu} T_f^{\mu\nu} = J^{\nu}$$

$$J^{\mu} = - \sum_i \frac{dp^{\mu}}{dt} \delta(\mathbf{r} - \mathbf{r}_i)$$



Motivations

Dynamically integrated transport model:
microscopic transport model + hydrodynamics

Dynamical initial condition for hydrodynamics through source terms

$$\partial_{\mu} T_f^{\mu\nu} = J^{\nu}, \quad \partial_{\mu} N_f^{\mu} = n_B$$

Some particles are converted into fluid elements with Gaussian profile.

$$J^{\mu} = \sum_i \frac{p_i^{\mu}(t)}{\Delta t} G(\mathbf{r} - \mathbf{r}_i(t)), \quad n_B = \sum_i \frac{B_i}{\Delta t} G(\mathbf{r} - \mathbf{r}_i(t))$$

JAM microscopic transport model

- space-time propagation of particles based on cascade method
- Resonance (up to 2GeV) and string excitation and decays
- Re-scattering among all hadrons
- DPM type string excitation law as in HIJING.
- Use Pythia6 for string fragmentation
- Propagation by the hadronic mean-fields within RQMD/S formulation
- EoS controlled collision term

$$\dot{\mathbf{r}}_i = \frac{\mathbf{p}_i}{p_i^0} + \sum_j \frac{m_j}{p_j^0} \frac{\partial V_j}{\partial \mathbf{p}_i} \quad \dot{\mathbf{p}}_i = - \sum_j \frac{m_j}{p_j^0} \frac{\partial V_j}{\partial \mathbf{r}_i} \quad p_i^0 = \sqrt{\mathbf{p}_i^2 + m_i^2 + 2m_i V_i}$$

Arguments of potential $\mathbf{r}_i - \mathbf{r}_j$ and $\mathbf{p}_i - \mathbf{p}_j$ are replaced by the distances in the two-body c.m.

A new approach: JAM+hydro model

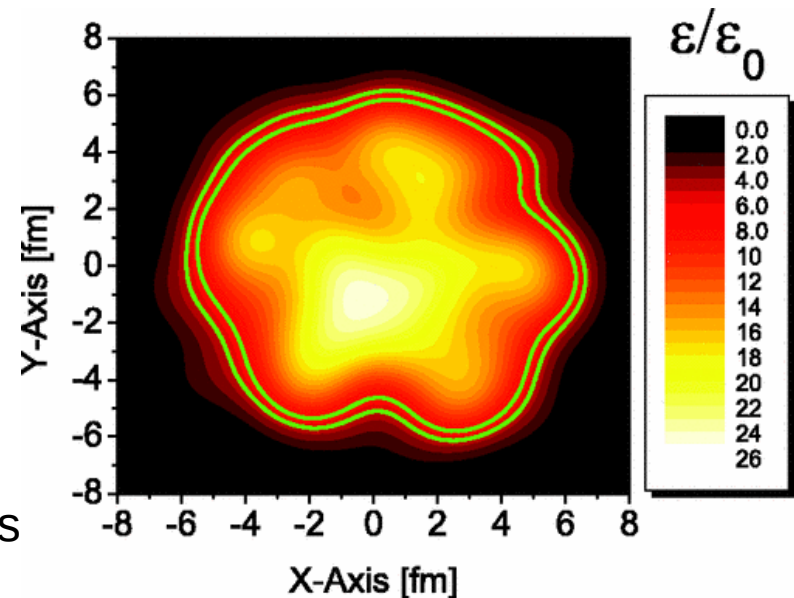
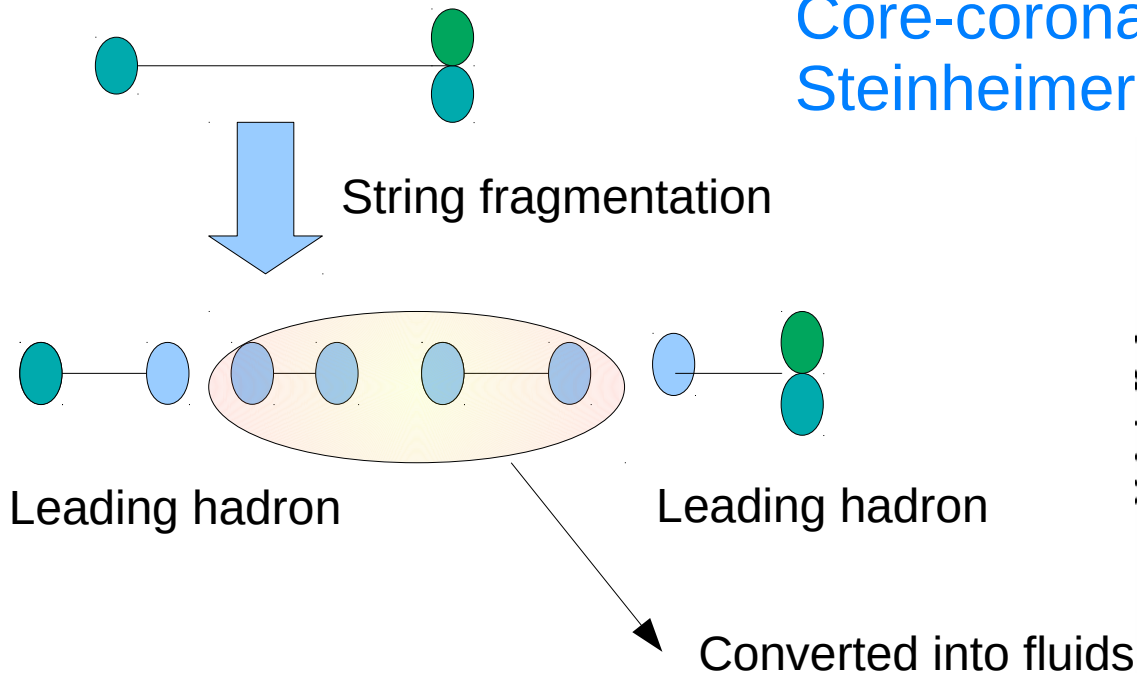
Dynamical coupling of fluids through source terms

$$\partial_{\mu} T_f^{\mu\nu} = J^{\nu}, \quad \partial_{\mu} N_B^{\mu} = \rho_B$$

Time dependent Core-corona separation

Put Hadrons from string or resonance decay into fluids after their formation time except leading hadrons when local energy density exceeds a hadronization energy density

Core-corona separation (K. Werner, 2007)
Steinheimer and Bleicher (2011)



Model parameters

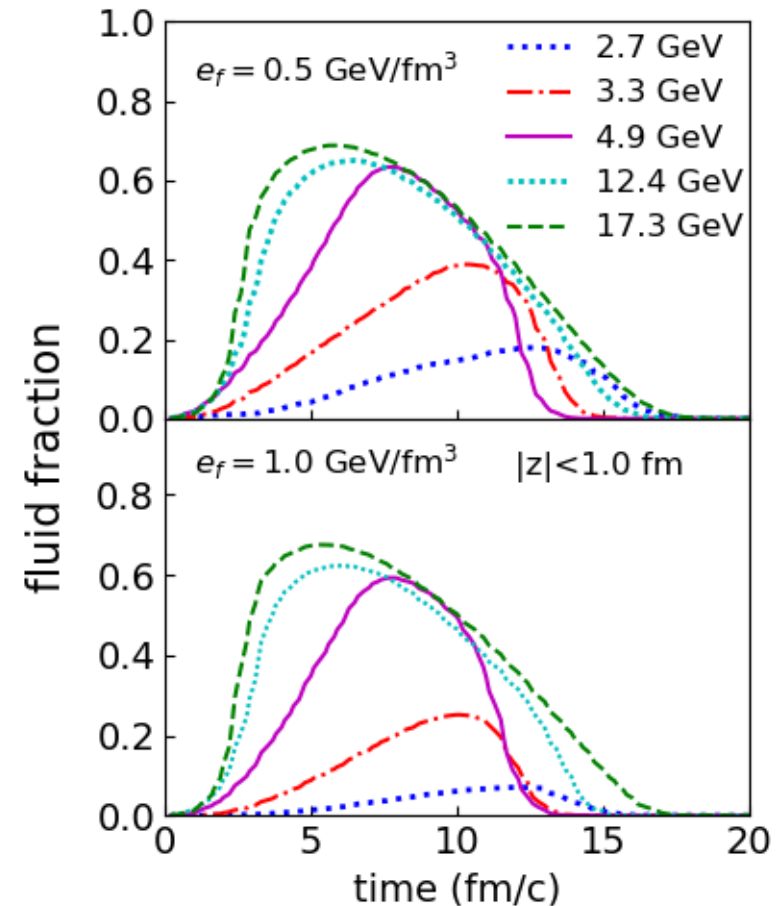
1) fluidization energy density

$$e_f = 0.5 - 1.0 \text{ GeV}/\text{fm}^3$$

2) particlization energy density

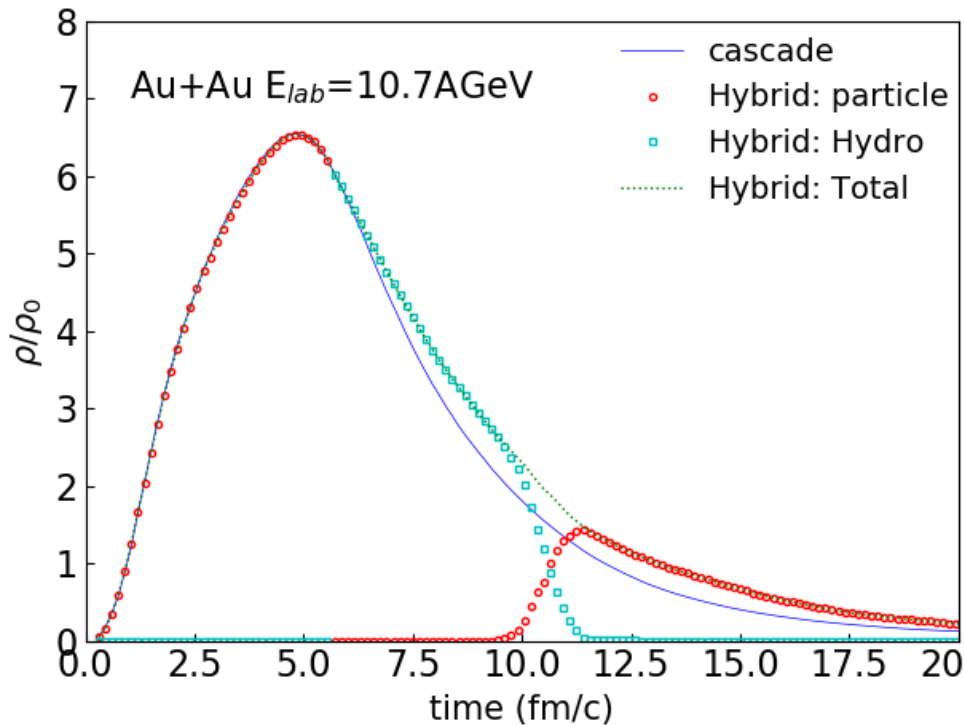
$$e_p = 0.5 \text{ GeV}/\text{fm}^3$$

- ## 3) equation of state: EOS-Q
- first-order phase transition
 - Bag model $B=235\text{MeV}^4$
 - hadronic resonances up to 2GeV
 - baryon density dependent
 - repulsive potential for baryons



Fraction of fluid energy at central region is about 70% at top SPS energy.

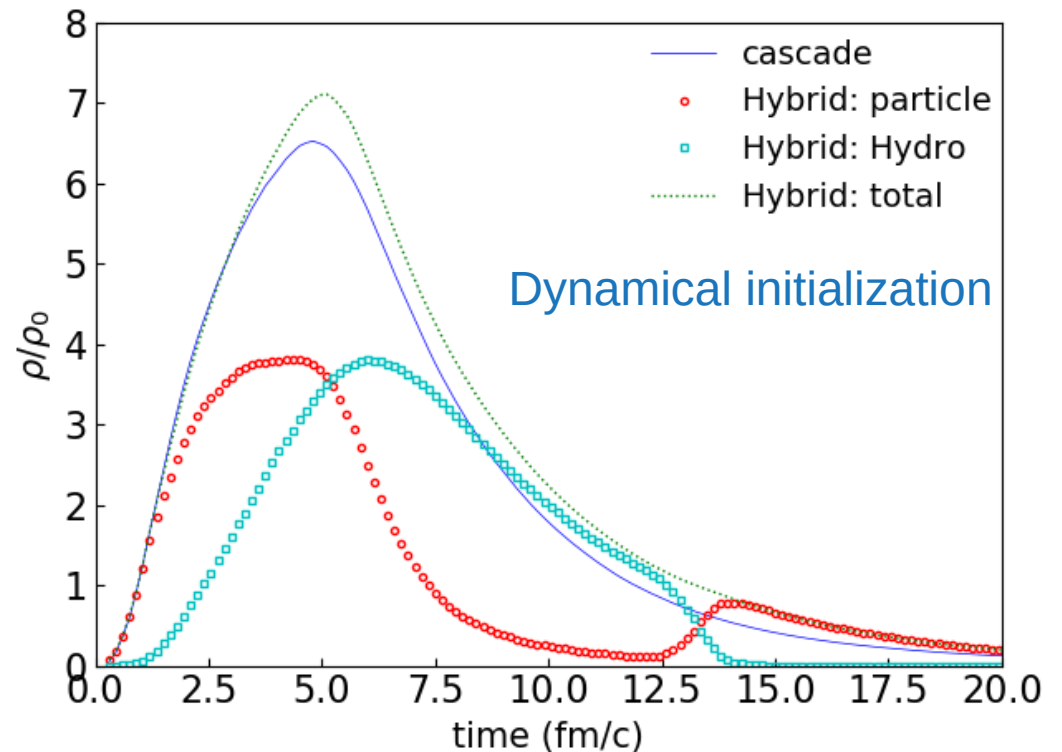
Hybrid model for AGS and SPS energies



Switch to hydro evolution
after two nuclei pass each other.

Switch to hadron transport
below a critical energy density.

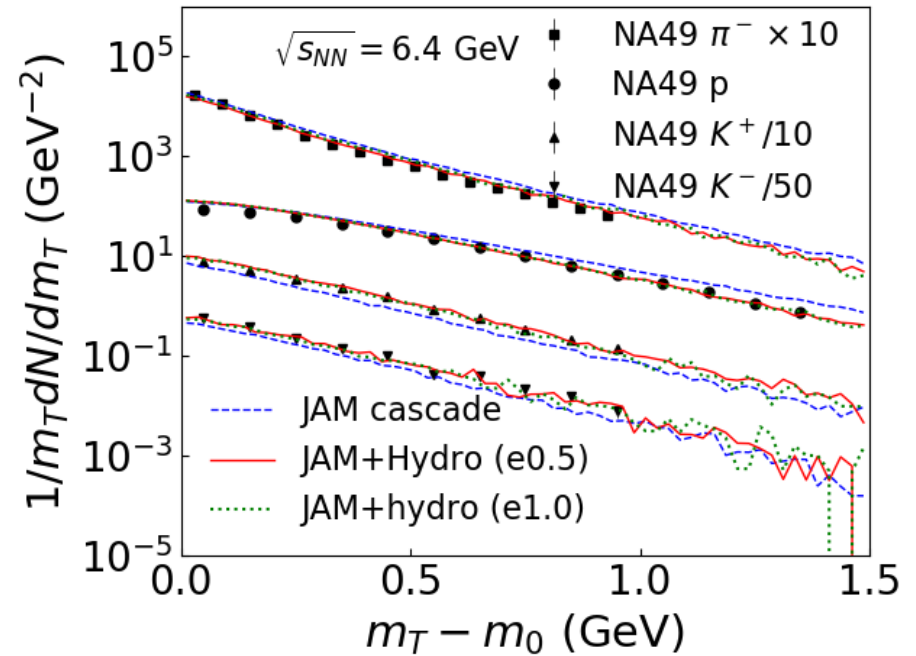
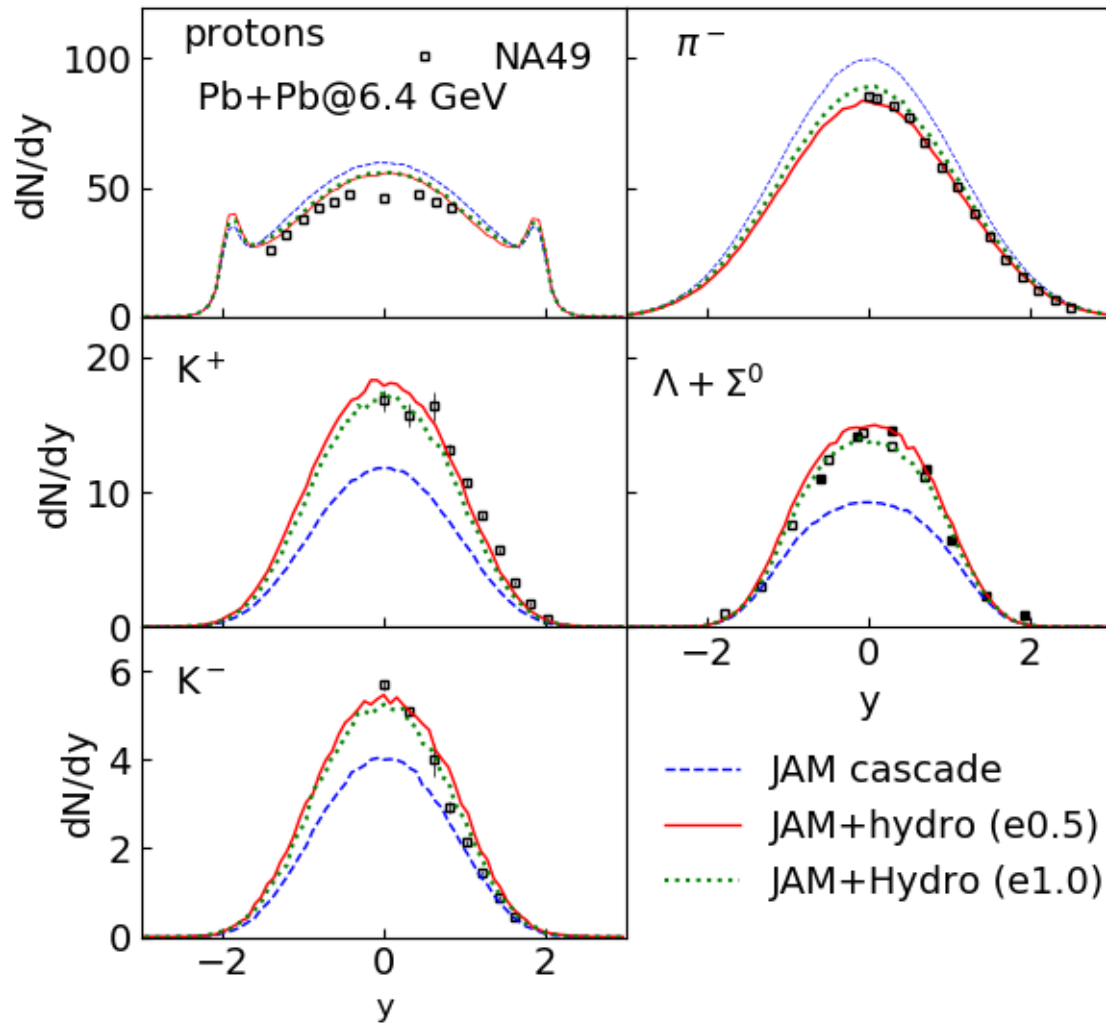
It is important to take into account potential effect in the Cooper-Fry formula
to ensure smooth transition from fluid to particles.



Hadronic EoS is used

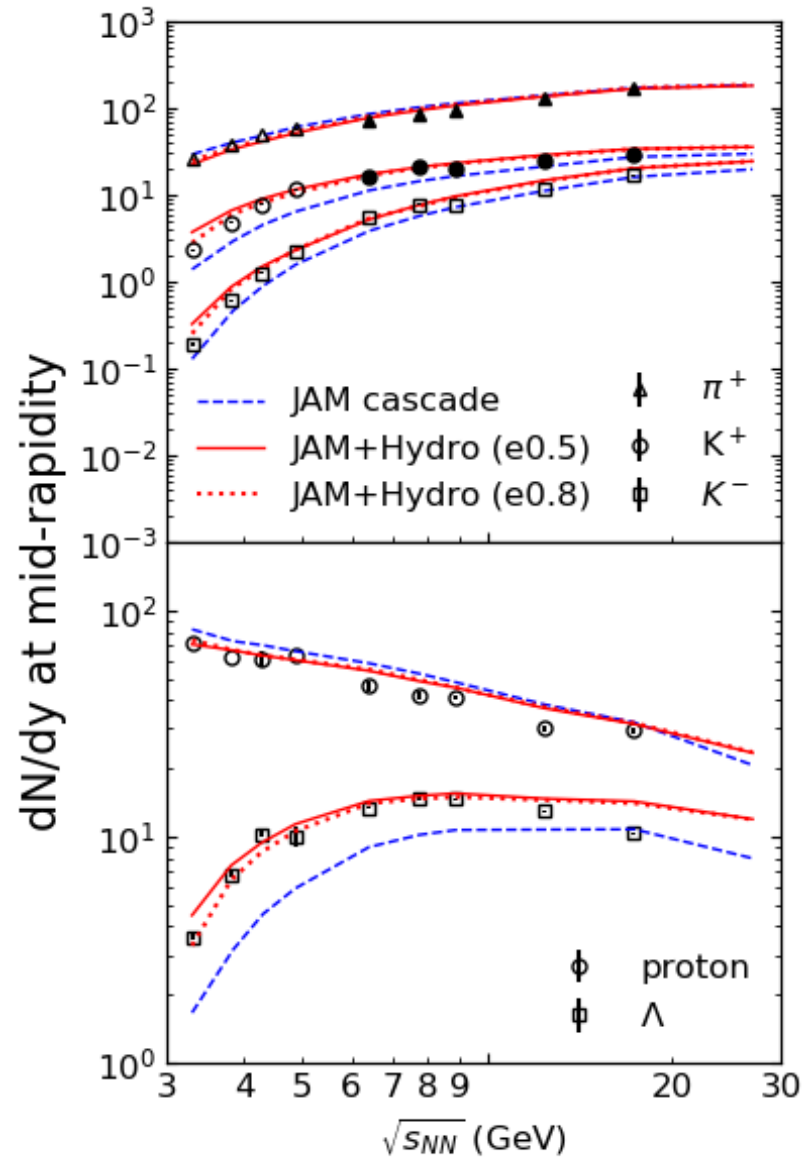
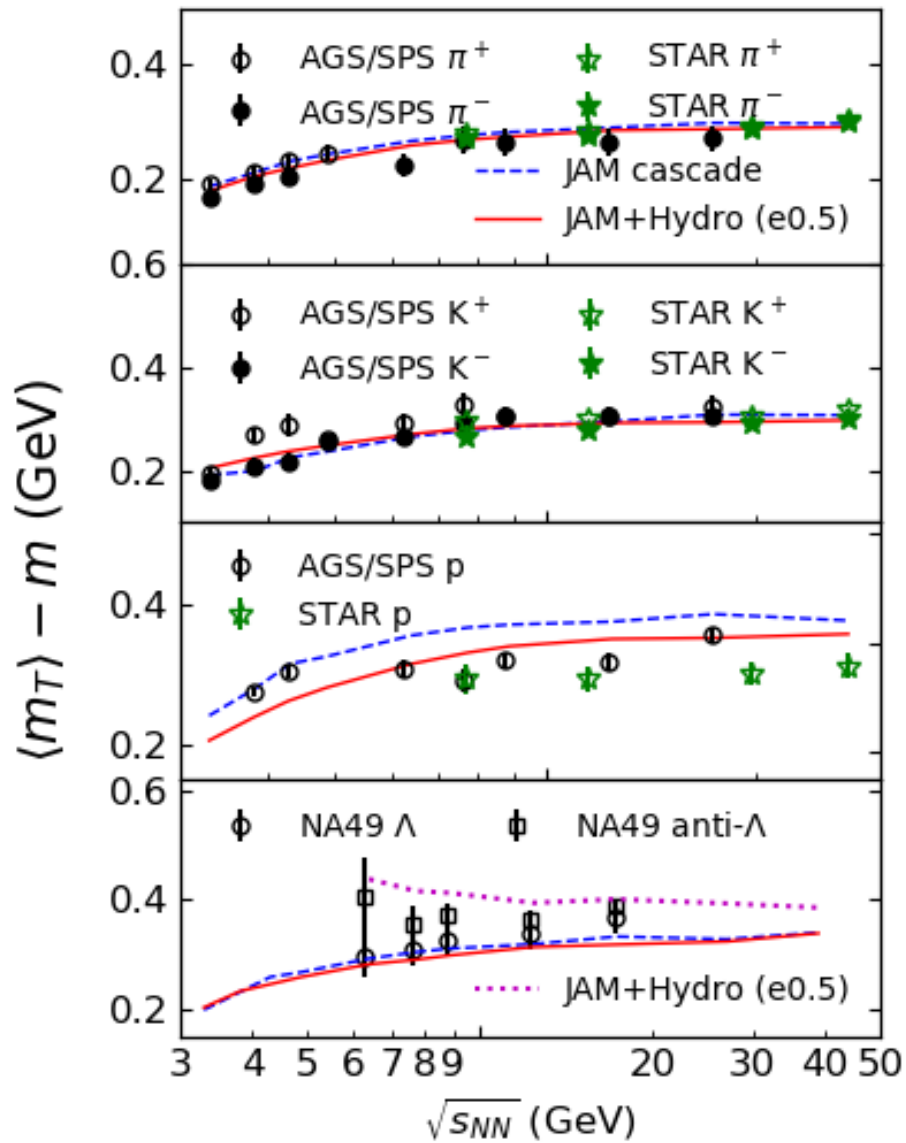
$$\mu = B\mu_B + S\mu_S \rightarrow B(\mu_B - V(\rho_B)) + S\mu_S$$

Particle spectra from a new hybrid model in Pb+Pb at $E_{lab}=20A\text{GeV}$

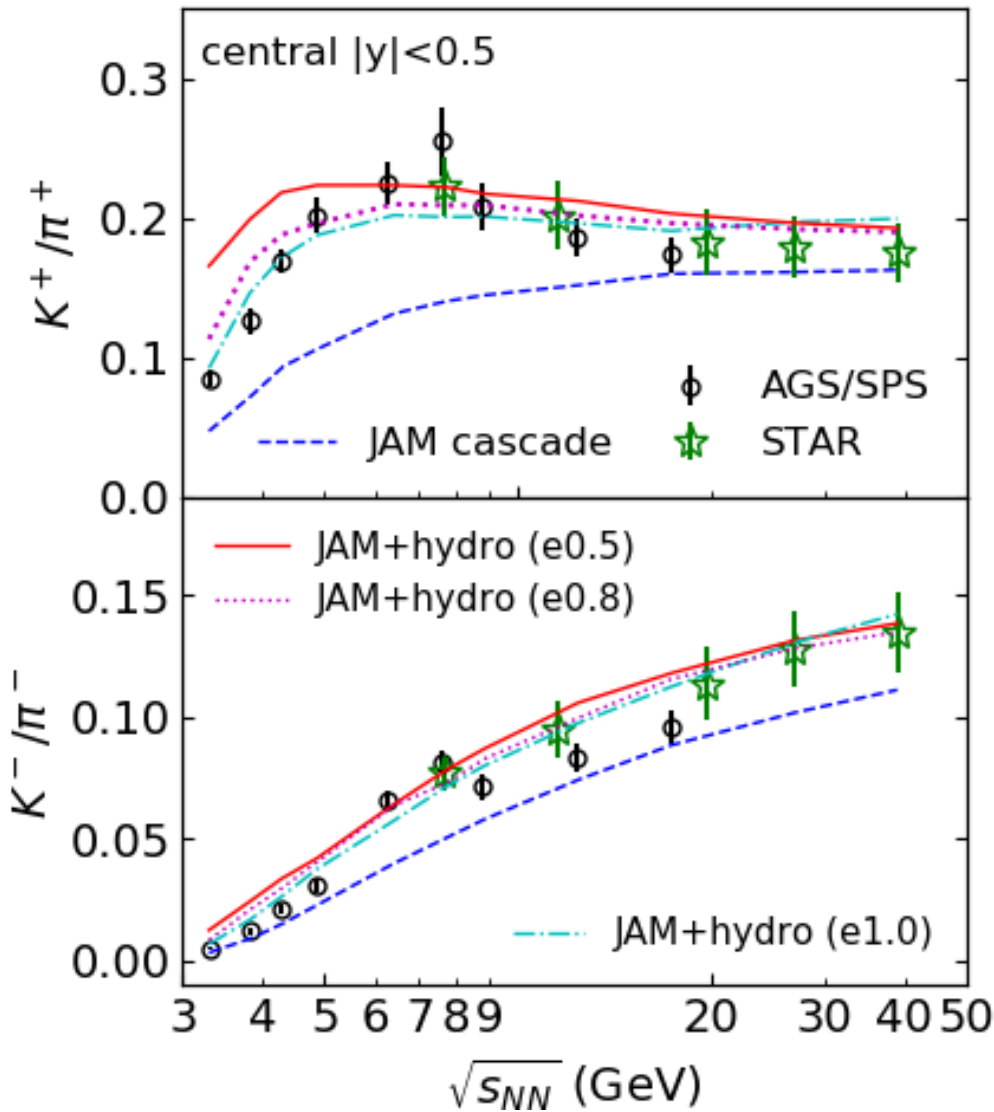


Fluidization energy density 0.5 or 1.0

Beam energy dependence of transverse mass and multiplicities from a new hybrid model.

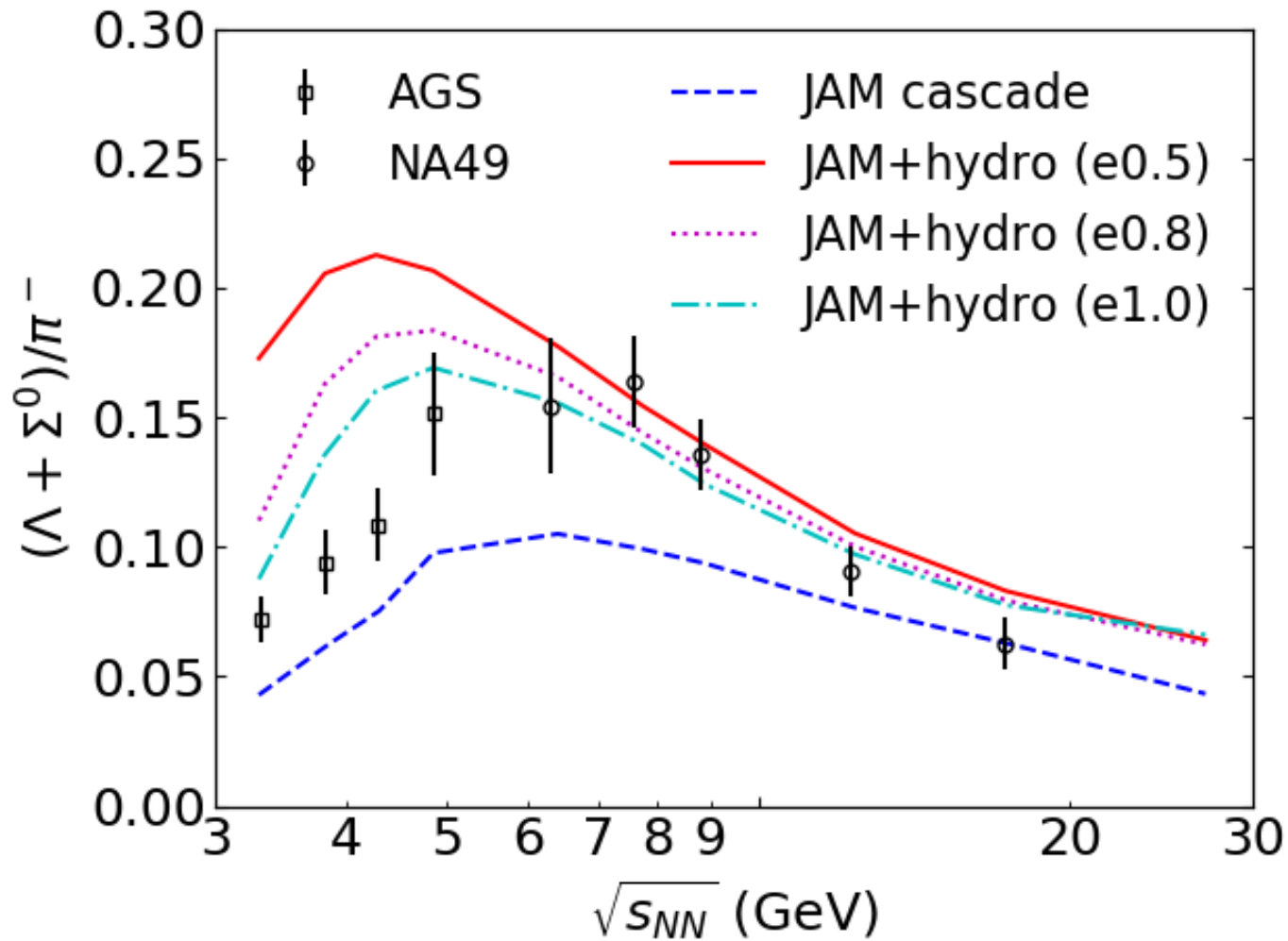


Beam energy dependence of K/π ratios from hybrid model.



Incomplete thermalization of the system is important for the description of K/π ratio.

Beam energy dependence of Λ/π ratios from a new hybrid model.



EOS modified collision term

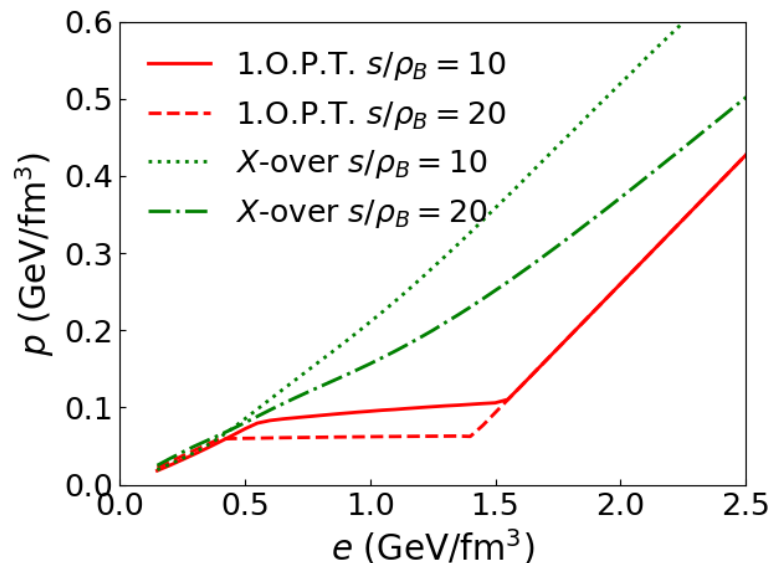
H. Sorge, Phys. Rev. Lett. 82,2048 (1999) Virial Theorem for two body collisions

$$P = P_{free} + \frac{1}{3TV} \sum_{(i,j)} [(\mathbf{p}'_i - \mathbf{p}_i) \cdot \mathbf{r}_i + (\mathbf{p}'_j - \mathbf{p}_j) \cdot \mathbf{r}_j]$$

The momentum change is constrained by

$$(\mathbf{p}'_i - \mathbf{p}_i) \cdot (\mathbf{r}_i - \mathbf{r}_j) = 3 \frac{(P - P_{free})}{\rho} (\Delta t_i + \Delta t_j)$$

When $P < P_{free}$: attractive orbit in the collision.

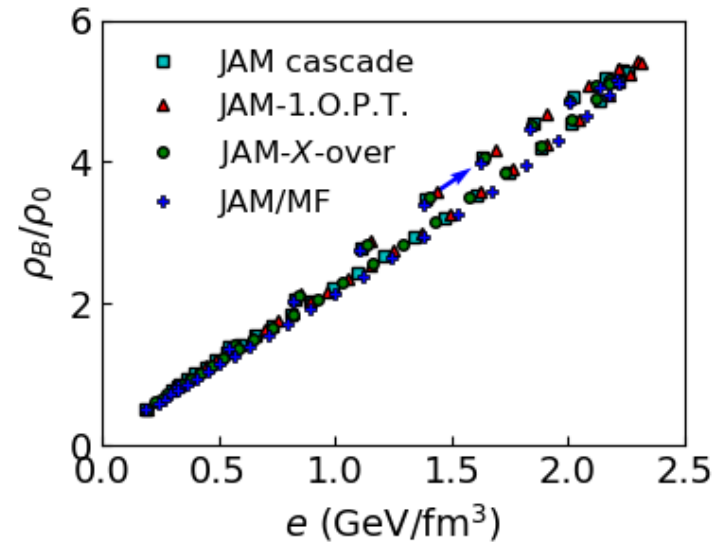
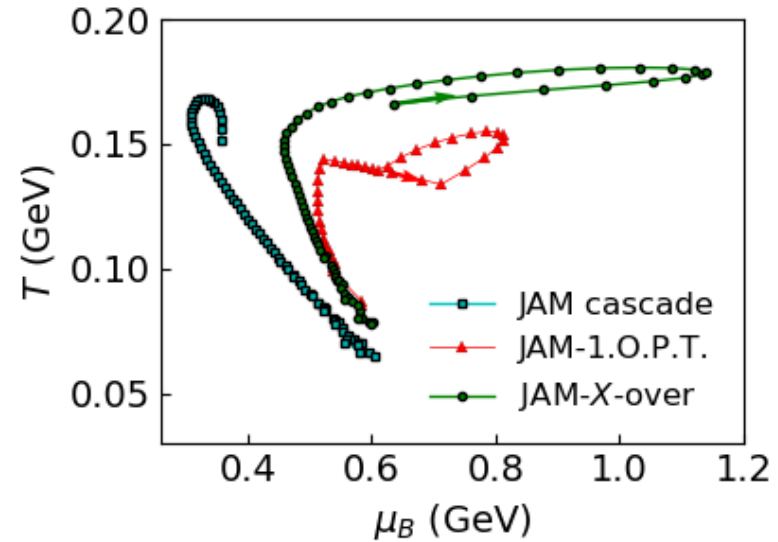
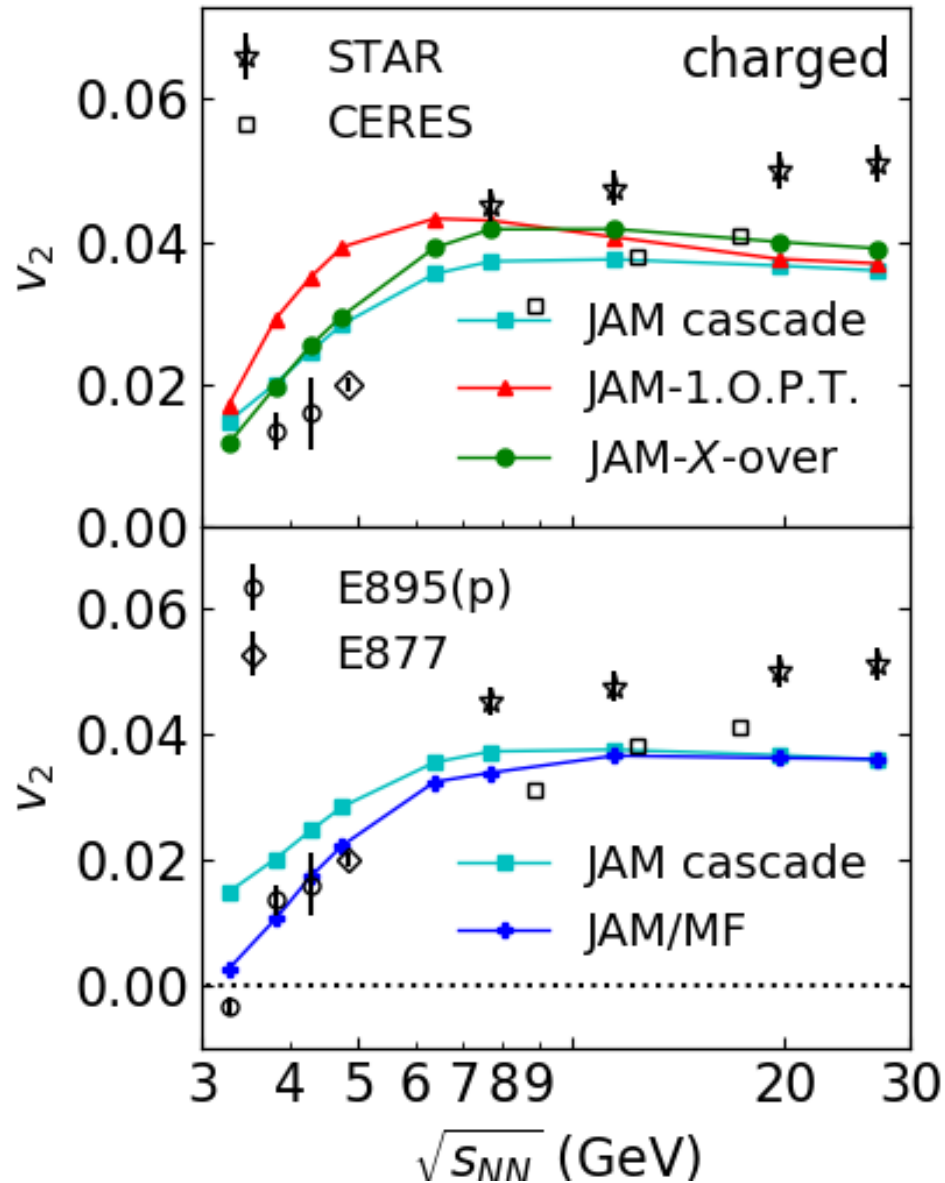


Fully baryon density dependent EoSs are implemented.
 Cross-over EOS: J. Steinheimer
 EOS-Q: Kolb, Sollfrank, Heinz

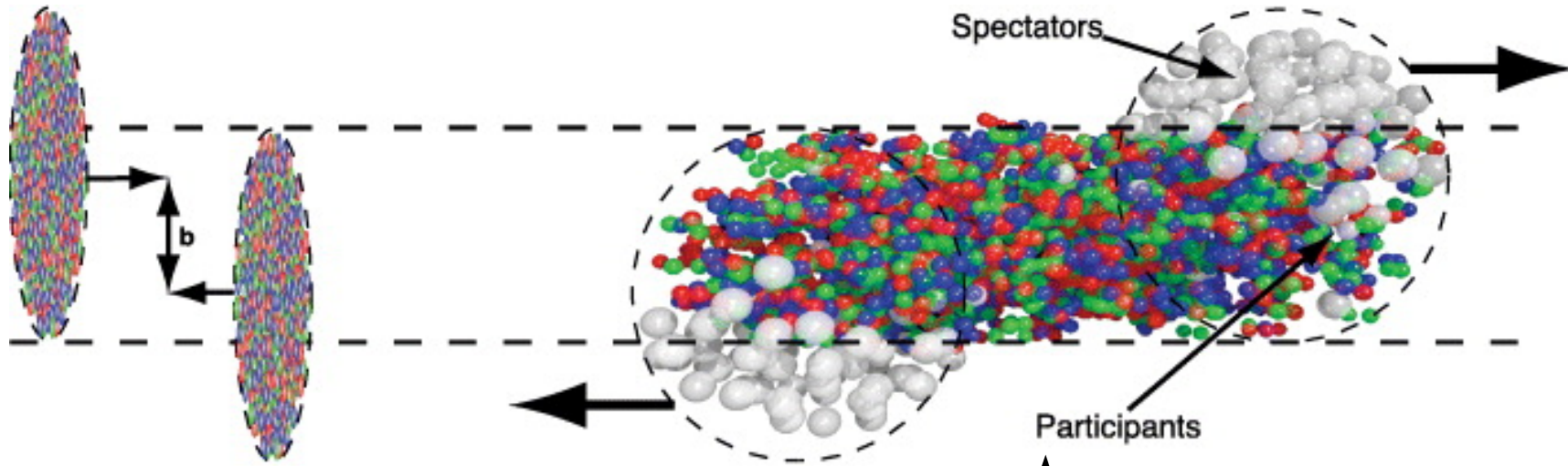
- Any EoS can be incorporated
- CPU time is as fast as standard cascade simulation
- Fully non-equilibrium transport approach

$$\frac{dN}{dyd^2p_T} = \frac{d^2N}{2\pi dp_T dy} (1 + 2v_1 \cos(\phi) + 2v_2 \cos(2\phi) + \dots)$$

V2 excitation functions $v_n = \langle \cos(n\phi) \rangle$



Effects of interaction with spectator

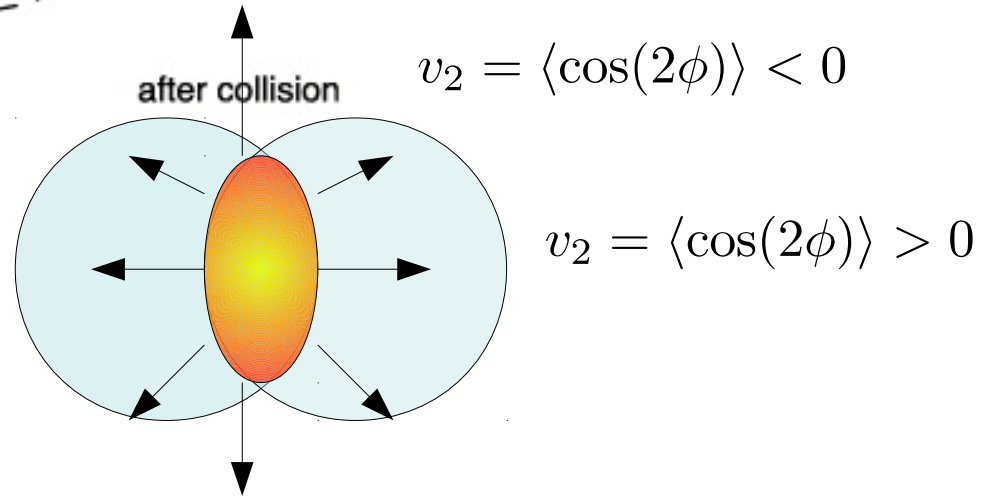
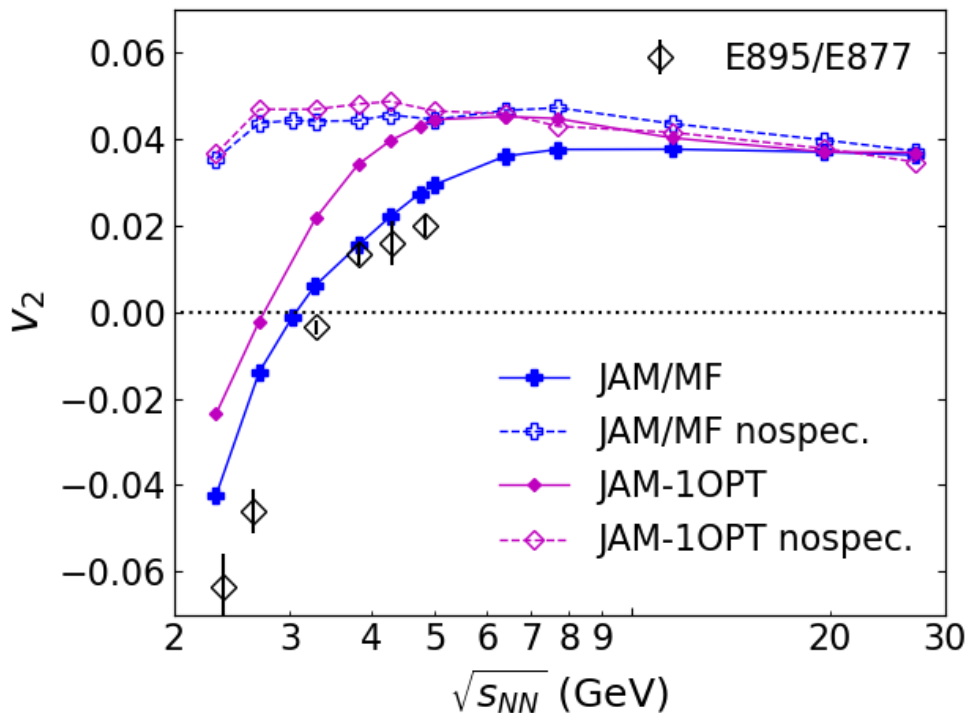


before collision

after collision

$$v_2 = \langle \cos(2\phi) \rangle < 0$$

$$v_2 = \langle \cos(2\phi) \rangle > 0$$



Comparison of the results without spectator interactions

In the case of first-order phase transition, shadowing effects are weak.

Summary

- JAM is a microscopic transport model for high energy nuclear collisions based on strings and hadronic resonances
- We have developed a new hybrid approach by **dynamical initialization of hydrodynamics** which takes into account **time dependent core-corona picture** in order to simulate heavy ion collision at high baryon energy region.
- K/pi ratios from JAM+hydro approach is in good agreement with the data, which is explained by the incomplete thermalization of the system.