Neutron stars and stellar mergers as a laboratory for dense QCD matter

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Main reference: Annala, Gorda, Kurkela, Vuorinen, PRL 120 (2018), 1711.02644





 10^8 years ago, 10^{24} m away from us, two stars of $R \sim 10^4$ m and $M \sim 10^{30}$ kg crossed paths.

On August 17 2017, LIGO measured a 10^{-17} m oscillation in the length of its 10^{3} m arms.

Today: Implications on physics of scale 10^{-15} m.

- I. Neutron star basics
- II. Observational information
- III. Theoretical limits for the neutron star matter Equation of State
- IV. Interpolating the EoS from low to high densities
- V. Final thoughts

What are neutron stars?



When a hydrogen burning star runs out of fuel:

- M $\leq 9M_{sun} \Rightarrow$ White dwarf
- $M \gtrsim 9M_{sun} \Rightarrow$ Supernova explosion $\circ M \gtrsim 20M_{sun} \Rightarrow$ Gravitational collapse into BH $\circ M \lesssim 20M_{sun} \Rightarrow$ Gravitational collapse into...



NS characteristics:

- Masses $\leq 2M_{\odot}$
- Radii $\approx 11 13$ km
- Spin frequencies \lesssim kHz
- Temperatures \lesssim keV
- \bullet Strong magnetic fields up to $10^{15} {\rm G}$

Unique laboratory for strong interaction physics: Density in NS cores high enough to probe nuclear matter well beyond saturation density



Physics picture: Hydrostatic equilibrium resulting from fierce competition between gravity and the pressure of QCD matter

GR description via Tolman-Oppenheimer-Volkov eqs:



$$\frac{dM(r)}{dr} = 4\pi r^2 \varepsilon(r),$$

$$\frac{dp(r)}{dr} = -\frac{G\varepsilon(r)M(r)}{r^2} \frac{(1+p(r)/\varepsilon(r))\left(1+4\pi r^3 p(r)/M(r)\right)}{1-2GM(r)/r}$$

$$\varepsilon(p) \Rightarrow M(R)$$

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Particle/nuclear theory challenge: Find Equation of State of strongly interacting matter that is

- Cold and dense
- Electrically neutral: $2/3n_u - n_d/3 - n_s/3 + n_e = 0$
- In beta equilibrium: $\mu_B/3 = \mu_d = \mu_s = \mu_u + \mu_e$



Big open questions:

- Can QCD theorists predict neutron star measurements?
- Can we infer the QCD matter EoS from observations?
- Can deconfined matter be found inside the stars?

What do we know from observations?



By now, two accurate Shapiro delay measurements of twosolar-mass stars: Demorest et al., Nature 467 (2010) Antoniadis et al., Science 340 (2013)

 $\therefore M_{\rm max} > 2M_{\odot}$



Radius measurements more problematic, but progress through observation of X-ray emission:

- Cooling of thermonuclear X-ray bursts provide radii to ~400m [Nättilä et al., Astronomy & Astrophysics 608 (2017), ...]
- With NICER mission, launched 3 June 2017, X-ray pulse profiling \rightarrow Radius of a single star (perhaps) to \sim 200m

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Gravitational wave breakthrough: LIGO and Virgo observation of NS merger 130 million ly away!

Three types of potential inputs:

- Tidal deformabilities of the NSs during inspiral – good measure of stellar compactness
- 2) EM signatures present if no immediate collapse to a BH
- 3) Ringdown pattern sensitive to
 EoS (also at T ≠ 0), but freq.
 too high for LIGO

LIGO and Virgo collaborations, PRL 119 (2017); 1805.11581





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Tidal deformability: How large a quadrupolar moment a star's gravitational field develops due to an external quadrupolar field

$$Q_{ij} = -\Lambda \mathcal{E}_{ij}$$

Substantial effect on observed GW waveform during inspiral phase



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$$Q_{ij} = -\Lambda \mathcal{E}_{ij}$$

Recent LIGO bound $70 < \Lambda(1.4M_{\odot}) < 580$ at 90% credence using low spin prior [LIGO and Virgo, arXiv:1805.11581]



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EM counterpart: short gamma ray burst detected 1.7s after GW measurement, followed by an optical signal

- Kilonova: Decay of heavy r-process elements
- GRB \rightarrow Proposed upper limit for the maximal mass of NSs: $M_{\text{max}} \leq 2.16^{+0.17}_{-0.15} M_{\odot}$ [Rezzolla, Most, Weih, ApJ 852 (2018)]

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Ringdown pattern: Unlike in BH mergers, expect a complex period of relaxation characterized by GW spectrum sensitive to both initial NS masses and the EoS



Baiotti, Rezzolla, Rept.Prog.Phys. 80 (2017) ²⁴

Post-merger dynamics can be studied with relativistic hydrodynamics, showing marked sensitivity to EoS, but frequency range (currently) too high for LIGO and Virgo



Takami, Rezzolla, Baiotti, PRD 91 (2015)

EoS – theoretical limits



Low-density behavior of EoS well known from nuclear theory side. Challenges begin close to saturation density:

- At $1.1n_s$, current errors in Chiral Effective Theory EoS $\pm 24\%$ mostly due to uncertainties in effective theory parameters
- State-of-the-art EoS NNNLO in chiral perturbation theory power counting [Tews et al., PRL 110 (2013), Hebeler et al., ApJ 772 (2013)]



Asymptotic freedom of QCD \Rightarrow High-density limit from a non-interacting theory. However,...

- At interesting densities $(1 10)n_s$ system strongly interacting but no nonperturbative methods available
- Naïve expectation: Weak coupling methods only useful at very high densities



Recent improvement: First part of four-loop pressure at T = 0derived: $p_{4-\text{loop}} \ni -\frac{11}{12} \frac{N_c d_A}{(2\pi)^3} \alpha_s m_{\infty}^4 \ln^2 \alpha_s$ [Gorda, Kurkela, Romatschke, Säppi, Vuorinen, arXiv:1806.xxxx] – cf. talk by Matias Säppi



Three-loop result with nonzero quark masses [Kurkela, Romatschke, Vuorinen, PRD 81 (2009)]

- Uncertainty of result at $\pm 24\%$ level around $40n_s$
- Main uncertainty from renormalization scale dependence
- Pairing contributions to EoS subdominant at relevant densities



Conclusion: Sizable no man's land extending from outer core to densities not realized inside physical neutron stars

Options: Use models, deform theory, or interpolate EoS between known limits and use astrophysical constraints

Interpolation – with and without observational constraints

Interpolate EoS using piecewise polytropic form, $p_i(n) = \kappa_i n^{\gamma_i}$, varying all parameters (γ_i , μ_i^{match})

Require:

- Smooth matching to nuclear and quark matter EoSs
- Continuity of p and n with at most one exception (1st order transition)
- 3) Subluminality
- 4) Optional: astrophysical constraints

[Kurkela et al., ApJ 789 (2014)]





Quadrutropic interpolation, using close to 200.000 randomly generated EoSs

Figures mostly from Annala, Gorda, Kurkela, Vuorinen, PRL 120 (2018)



Implement then twosolar-mass constraint: Accept only EoSs that fulfill $M_{\rm max} > 2M_{\odot}$

Assumption here and in the following: All stars considered main seq. NSs

 Excluded: twin stars [e.g. Alvarez-Castillo, Blaschke, PRC96
 (2017)], strange quark stars [e.g. Weber et al., IAU 291 (2013)]





Next, determine tidal deformabilities for each EoS and compare to LIGO results for $1.4M_{\odot}$ stars:

• $70 < \Lambda(1.4M_{\odot}) < 580$ at 90% credence







Recent result of Nättilä et al., Astronomy & Astrophysics 608 (2017):

 $R(1.9M_{\odot}) = 12.4 \pm 0.4$ km with 1σ credence

 10^{4}

16

Caveat: mass measurement of the star in question still rather uncertain







Lesson 1: $2M_{\odot}$ and pQCD constraints force EoS to be hard at low density and soft at high density \rightarrow Efficient bracketing of $p(\epsilon)$, $R(1.4M_{\odot}) \geq 10$ km





Lesson 2: Stringent limits from only one tidal deformab. measurement: EoS cannot be overly stiff at low density, $R(1.4M_{\odot}) \leq 13$ km





Lesson 3: Accurate radius measurements of stars with well-known (large) masses very valuable for EoS determination especially at low densities



One stellar merger later – where are we?

Big open questions:

- Can QCD theorists predict neutron star measurements?
 Not there yet need fundamentally new machinery
- Can we infer the QCD matter EoS from observations?
 Looks very promising, fast progress with GWs
- Can deconfined matter be found inside the stars?
 > Tough question, but we're on the right path!