

Towards a cubic closed string field theory

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Outline: 1. Introduction

Introduction

Hubbard–Stratonovich action

IBL_∞ action

Conclusion

Motivations

String field theory:

- ▶ QFT with infinite number of spacetime fields
- ▶ split string amplitudes into Feynman diagrams
- ▶ motivations:
 - ▶ unified framework: off-shell amplitudes, renormalisation
 - ▶ classical solutions
 - ▶ consistency of string theory
 - ▶ path integral → non-perturbative effects

String field theory – definitions

- ▶ CFT Hilbert space \mathcal{H} , string field $\Psi \in \mathcal{H}$
- ▶ (bosonic) classical action

$$S = \sum_{n=2}^{\infty} \frac{g_s^{n-2}}{n!} \mathcal{V}_n(\Psi^n)$$

g_s : coupling constant

- ▶ string vertex $\mathcal{V}_n : \mathcal{H}^{\otimes n} \rightarrow \mathbb{C}$

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- ▶ BRST operator: $\ell_1 = Q_B$
- ▶ eom and gauge invariance (L_∞ structure)

$$\sum_{n=1}^{\infty} \frac{g_s^{n-1}}{n!} \ell_n(\Psi^n) = 0, \quad \delta \Psi = \sum_{n=1}^{\infty} \frac{g_s^{n-1}}{n!} \ell_n(\Lambda, \Psi^{n-1})$$

String field

- ▶ Fourier expansion

$$|\Psi\rangle = \sum_r \int \frac{d^D k}{(2\pi)^D} \psi_r(k) |k, r\rangle$$

- ▶ $\psi_r(k)$: spacetime fields
- ▶ $\{|k, r\rangle\}$: basis of \mathcal{H}
- ▶ k : spacetime momentum (non-compact dim.)
- ▶ r : discrete modes (Lorentz indices, compact dim., etc.)
- ▶ classical closed string

$$N_{\text{gh}}(\Psi) = 2, \quad L_0^- |\Psi\rangle = b_0^- |\Psi\rangle = 0$$

String field theory – difficulties

- ▶ infinite number of interactions → non-polynomial action
 - ▶ definition of vertex \mathcal{V}_n :
 - ▶ n holomorphic function $\{f_{n,i}\}$ (with constraints)
 - ▶ subspace of moduli space $\mathcal{V}_n \subset \mathcal{M}_{0,n}$
 - ▶ gluing with propagators and summing
= single covering of $\mathcal{M}_{0,n}$
- no explicit representation for $n > 5$
- [[hep-th/9412106](#), Belopolsky][[hep-th/0408067](#), Moeller][[hep-th/0609209](#),
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But:

- ▶ general properties understood, used to prove consistency of string theory [[Sen-Zwiebach '94, Sen '15-18, ...](#)]
- ▶ expect progress from [[1806.00449, Headrick-Zwiebach](#)]
[\[1806.00450, Headrick-Zwiebach\]](#)

The open string miracle

- ▶ open string: \exists truncation to cubic interactions [Witten '86]

$$\ell_2(A, B) \sim A * B, \quad \forall n > 2 : \quad \ell_n = 0$$

non-commutative $*$ -product

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- ▶ many explicit computations:
 - ▶ analytic and level-truncation solutions (tachyon condensation, brane decay, marginal deformations, defects, time-dependent backgrounds...)
 - ▶ effective actions
 - ▶ etc.

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There is no covariant BRST invariant single covering of $\mathcal{M}_{0,4}$ built from a symmetric cubic vertex assuming standard plumbing fixture.

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Way out: use auxiliary fields to decompose interactions

1. Hubbard–Stratonovich transformation
2. IBL_∞ structure

Outline: 2. Hubbard–Stratonovich action

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Hubbard–Stratonovich transformation

Hubbard–Stratonovich (intermediate field) representation:

- ▶ originally: many-body physics (condensed matter and nuclear physics)
- ▶ decompose interactions to lower-order interactions through auxiliary fields
- ▶ optimal representation: cubic interactions, quadratic in the physical field, linear in the auxiliary fields

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- ▶ (vector) ϕ^4 , ϕ^{2n} [[1003.1037, Rivasseau-Wang](#)][[1601.02805, Llonni-Rivasseau](#)], matrix and tensor models [[1609.05018, Llonni-Rivasseau](#)]

Scalar ϕ^4 model

- ▶ ϕ^4 action

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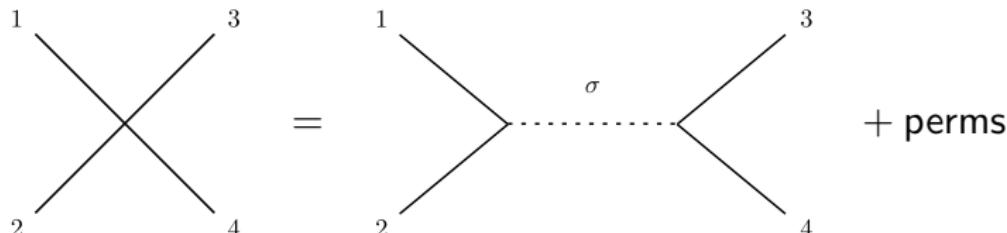
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- ▶ Hubbard–Stratonovich action

$$S_{\text{HS}} = - \int d^d x \left(\frac{1}{2} \phi K \phi - 2\lambda \frac{\sigma^2}{2} + 2\lambda \sigma \phi^2 \right)$$

- ▶ Feynman graphs



Scalar ϕ^4 vector model

- ▶ vector ϕ^4 [[Weinberg vol. 2](#)]

$$S = \int dk_1 \cdots dk_4 V_{i_1 i_2 i_3 i_4}(k_1, \dots, k_4) \phi_{i_1}(k_1) \cdots \phi_{i_4}(k_4)$$
$$+ \frac{1}{2} \int dk \phi_i(k) K_{ij}(k) \phi_j(-k)$$

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$$\delta S = - \int dk_1 \cdots dk_4 V_{i_1 \dots i_4}(k_1, \dots, k_4) \left(\sigma_{i_1 i_2}(k_1, k_2) - \phi_{i_1}(k_1) \phi_{i_2}(k_2) \right)$$
$$\times \left(\sigma_{i_3 i_4}(k_3, k_4) - \phi_{i_3}(k_3) \phi_{i_4}(k_4) \right)$$

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$$+ 2 \int dk_1 \cdots dk_4 V_{i_1 \dots i_4} \sigma_{i_1 i_2} \phi_{i_3} \phi_{i_4}$$

String field theory

- ▶ string field action to $O(g_s^2)$

$$S = \frac{1}{2} \mathcal{V}_2(\Psi^2) + \frac{g_s}{3!} \mathcal{V}_3(\Psi^3) + \frac{g_s^2}{4!} \mathcal{V}_4(\Psi^4)$$

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$$S_{HS} = \frac{1}{2} \mathcal{V}_2(\Psi^2) + \frac{g_s}{3!} \mathcal{V}_3(\Psi^3) - \frac{g_s^2}{4!} \mathcal{V}_4(\Sigma, \Sigma) + \frac{2g_s^2}{4!} \mathcal{V}_4(\Sigma, \Psi^2)$$

Properties

- ▶ new product $m_2 : \mathcal{H}^{\otimes 2} \rightarrow \mathcal{H}^{\otimes 2}$

$$\mathcal{V}_4(A_1, \dots, A_4) = \langle A_1 \otimes A_2 | m_2(A_3, A_4) \rangle$$

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- ▶ Fourier expansion \rightarrow entangled states

$$|\Sigma\rangle = \sum_{r,s} \int \frac{d^D k_1}{(2\pi)^D} \frac{d^D k_2}{(2\pi)^D} \sigma_{rs}(k_1, k_2) |\phi_r(k_1)\rangle \otimes |\phi_s(k_2)\rangle$$

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 - ▶ cubic in the field
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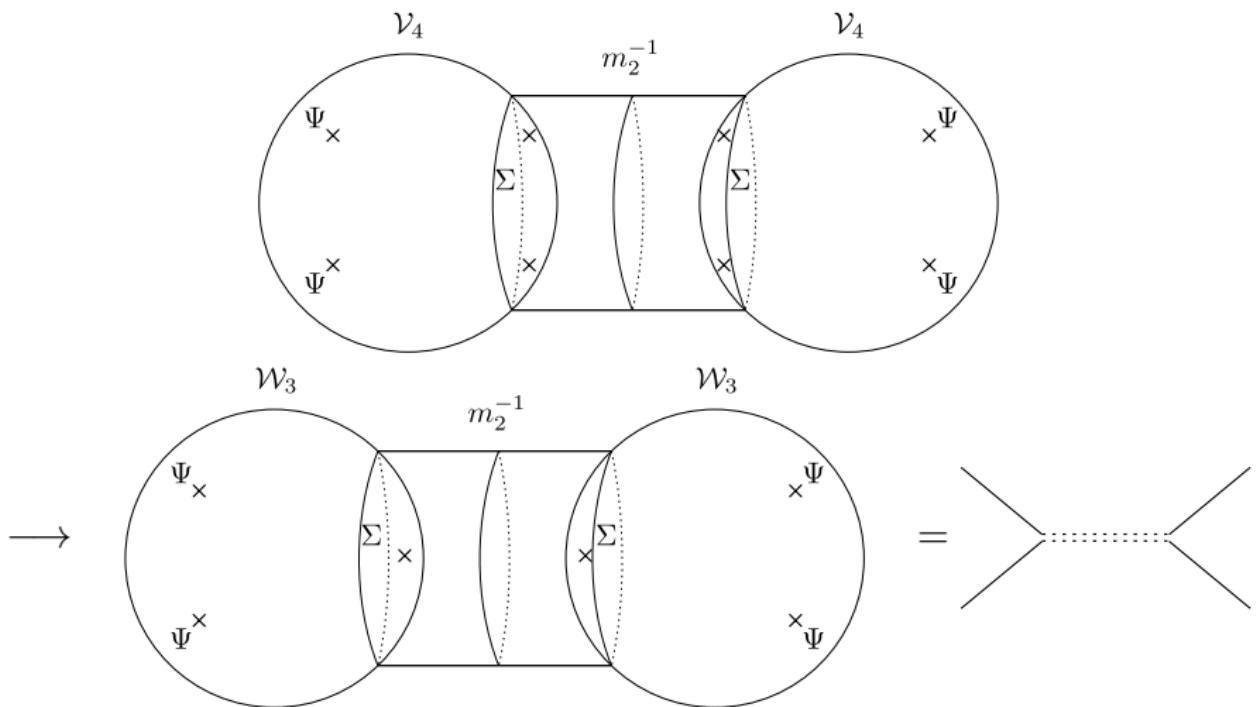
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- ▶ worldsheet: insert states of tensor product at different punctures \rightarrow merge them together (no singularities)
 - ▶ \mathcal{V}_4 : glue two punctures with two punctures
 - ▶ \mathcal{W}_2 : glue one puncture with one puncture

Worldsheet interpretation (2)



Generalisation

Questions:

- ▶ how to generalize to higher orders?
→ will encounter products with more outputs
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- ▶ which structure will have the action?

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⇒ IBL_∞ algebra

Outline: 3. IBL_∞ action

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IBL_∞ algebra

- ▶ symmetric tensor algebra (implicit symmetrization)

$$S\mathcal{H} = \bigoplus_{n \geq 1} \mathcal{H}^{\otimes n}$$

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- ▶ IBL_∞ products $p_{m,n} : \mathcal{H}^{\otimes m} \rightarrow \mathcal{H}^{\otimes n}$

$$\hat{p} = \sum_{m,n \geq 1} g_s^{m+n-2} \hat{p}_{m,n}$$

[[1508.02741](#), Cieliebak-Fukaya-Latschev] (see also [[1109.4101](#),
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$$\hat{p} \circ \hat{p} = 0 \quad \Rightarrow \quad \sum_{\substack{m_1+m_2=m+1 \\ n_1+n_2=n+1}} \hat{p}_{m_1,n_1} \circ \hat{p}_{m_2,n_2} = 0$$

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- ▶ induces a BV algebra [[1511.01591](#), [Markl-Voronov](#)]

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$$\omega = \sum_{n \geq 1} \omega_n, \quad \omega_n = \omega_1^{\otimes n} + \text{perms}, \quad \omega_1 : \mathcal{H} \times \mathcal{H} \rightarrow \mathbb{C}$$

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- ▶ gauge invariance

$$\delta |\Psi\rangle = \hat{p}(\Lambda, e^\Psi), \quad \Lambda = \sum_{n \geq 1} \Lambda_n, \quad \Lambda_n \in \mathcal{H}^{\otimes n}$$

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$$\delta |\Psi\rangle = \hat{p}(\Lambda, e^\Psi), \quad \Lambda = \sum_{n \geq 1} \Lambda_n, \quad \Lambda_n \in \mathcal{H}^{\otimes n}$$

- ▶ note: generalize at quantum level (new index $g \geq 0$)

IBL_∞ string field theory

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Conjecture 2

There is a parametrization such that the action is cubic in the fields.

Comments

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Outline: 4. Conclusion

Introduction

Hubbard–Stratonovich action

IBL_∞ action

Conclusion

Outlook

Summary: path for constructing a cubic closed string field theory

- ▶ SFT with IBL_∞ algebra
- ▶ algebraic description of SFT with auxiliary fields

If correct:

- ▶ classical solutions (consistent truncation of auxiliary fields)
- ▶ thermal effects (σ saddle point = mean field theory)
- ▶ explicit computations of vertices and amplitudes
- ▶ non-perturbative effects

If not correct:

- ▶ learn more about the structure of SFT
- ▶ if holds up to some order, still useful for perturbative computations