Precision studies of AdS/CFT

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Based on works with: Chi-Ming Chang, Martin Fluder, Ying-Hsuan Lin, Mark Mezei, Silviu Pufu

- M. Mezei, S. S. Pufu and Y. Wang, *A 2d/1d Holographic Duality*, 1703.08749.
- C.-M. Chang, M. Fluder, Y.-H. Lin and Y. Wang, Spheres, Charges, Instantons, and Bootstrap: A Five-Dimensional Odyssey, JHEP 03 (2018) 123, [1710.08418].
- C.-M. Chang, M. Fluder, Y.-H. Lin and Y. Wang, Romans Supergravity from Five-Dimensional Holograms, 1712.10313.
- M. Mezei, S. S. Pufu and Y. Wang, 5d Super-Yang-Mills, 3d Chern-Simons, and 6d (2,0) Theories, 1807.XXXXX.

Techniques to study superconformal fixed points

 $\mathsf{SUSY} \to \mathsf{Control}$ over strongly-coupled dynamics

supersymmetric partition function from localization

 $S^1 imes S^{d-1}$, S^d ...

protected operator algebra

chiral ring, chiral algebra, topological quantum mechanics, defects ...

superconformal bootstrap

stress-tensor bootstrap, flavor bootstrap, Mellin amplitudes

anomaly

conformal anomalies, t'Hooft anomalies ...

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Q1: Can we use these techniques in the large N limit to do precision studies of AdS/CFT?

- Q2: Gain more insight about String/M-theory?
- Q3: Non-SUSY AdS/CFT from embedding in dual pairs with SUSY?
- Q4: Confirmations/Lessons for strongly coupled field theory phenomena?

 AdS_3 Examples with $\mathcal{N} = (2, 2)$ and (4, 4) SUSY well-studied. Little known about $\mathcal{N} = (1, 1), (3, 3)$ etc. See Massimo's and Matthias's talks for $\mathcal{N} = (4, 4)$.

 \textit{AdS}_{4} Examples with $\mathcal{N}=2,3,4,6,8$ studied extensively. Though $\mathcal{N}=1$ is little studied.

 AdS_5 Lots of works for $\mathcal{N} = 1, 2$ and $\mathcal{N} = 4$, little for $\mathcal{N} = 3$. See Robert's and Charlotte's talks for $\mathcal{N} = 4$.

 AdS_6 Only $\mathcal{N} = 1$ is possible See also Oren's talk. AdS_7 Many works for (2,0) but few for (1,0). See also Eric's talk for (2,0).

Our focus today

 AdS_3 Examples with $\mathcal{N} = (2, 2)$ and (4, 4) SUSY well-studied. Little known about $\mathcal{N} = (1, 1), (3, 3)$ etc. See Massimo's and Matthias's talks for large N = 4.

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 AdS_6 Only $\mathcal{N} = 1$ is possible See also Oren's talk. AdS_7 Many works for (2,0) but fewer for (1,0). See also Eric's talk for (2,0).

AdS_4/CFT_3

Mezei-Pufu-YW: arXiv:1703.08749

- 3d N = 4 SCFT has a 1d topological sector with nontrivial operator algebra [Chester-Lee-Pufu-Yacoby, Beem-Rastelli-Peelaers, Dedushenko-Pufu-Yacoby]
- Defined by **cohomology** with respect to nilpotent $Q_{1,2}$. Correlations function only depend on ordering on a line in \mathbb{R}^3 .
- Higgs branch operators twisted by $SU(2)_H$ R-symmetry polarizations
- Higgs branch chiral ring → noncommutative associative deformation (also for Coulomb branch).
- Stereographic map: $\mathbb{R} \subset \mathbb{R}^3 \to S^1 \subset S^3$, the TQM can be formulated as a **gauged** QM (for certain 3d SCFTs).
- Goal: Obtain a nontrivial 1d subsector of the large N 3d SCFTs

In [Dedushenko-Pufu-Yacoby]

• Localizing supercharge $Q = Q_1 + Q_2$ satisfies

$$Q^2 \sim iP_{\tau} + R_C$$

- BPS locus is **covariantly constant** along τ , and parametrized by matter fields Q, \tilde{Q} on S^1 at $\theta = \pi/2$
- Weighted by a gauged QM action where the gauge field *A* is **emergent**

$$Z_{S^{3}} \sim \int D\mathcal{A}DQ \, D\tilde{Q} \, \exp\left[4\pi r \int d\varphi \, \tilde{Q}(\partial_{\varphi} + i\mathcal{A}_{\varphi})Q\right]$$



2d protected sector of the N = 4 SUGRA on AdS_4

Goal: 1d TQM in 3d SCFT \rightarrow 2d protected sector in AdS_4 SUGRA? The strategy in [Mezei-Pufu-YW'17]

- Focus on matter multiplets in $\mathcal{N} = 4$ SUGRA
- Action with off-shell SUSY on manifold with **boundary**
- Show that on the BPS locus the fields are covariantly constant along τ
- Reduce the theory on AdS_4 to $\frac{1}{2}AdS_3$ and write the action as a total derivative
- Obtain 2d Yang-Mills action on AdS₂ with appropriate boundary terms on S¹. Gauge group determined by the matter content. Yifan Wang



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Focus on a familiar $AdS_4/{
m CFT}_3$ dual pair

 $\mathcal{N}=8 \text{ ABJM at } k=1 \xleftarrow{\text{Holographic dual}} \text{M theory on } AdS_4 \times S^7$

The k = 1 ABJM theory has an equivalent/mirror description in terms of $\mathcal{N} = 4 \ U(N)$ vector multiplet coupled to one adjoint hyper X and one fundamental hyper Q [Kapustin, Bashkirov '10]

$$\begin{array}{l} \text{In [Mezei-Pufu-YW'17]} \\ \text{1d TQM} \xleftarrow{\text{Holographic dual}} \text{2d YM with } \frac{1}{g_{\text{YM}}^2} = \frac{N^{3/2}}{3\sqrt{2}} \text{ and } \mathfrak{g} = \text{Sdiff}(S^2) \end{array}$$

Field theory side:

- TQM determined by matter content of 3d SCFT
- Solve large N limit analytically
- 2,3 and some 4-point correlators in the planar limit.

Gravity side:

- Lowest KK modes: $\mathcal{N}=4$ SYM subsector of 4d $\mathcal{N}=8$ SUGRA
- Localization to obtain bulk 2d action
- Explicit check of the 2d/1d proposal at the 2-derivative order
- Predictions for bulk higher derivatives interactions from field theory

Similar analysis for $\mathcal{N} = 4$ higher spin type AdS/CFT duals in the paper.

 AdS_6/CFT_5

Chang-Fluder-Lin-YW: arXiv:1710.08418, arXiv:1712.10313

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Local correlators in 5d N = 1 SCFT

- 5d SCFTs do not have weakly coupled UV descriptions but rather emerge as UV fixed points of IR free gauge theories. $g_{YM}^2 \sim \frac{1}{m_l}$. [Seiberg, Morrision, Intriligator ...]
- Global symmetries at the fixed point not manifest in the IR phase
- Localization computation for superconformal index and S⁵ free energy [Kim-Kim-Lee, Kim-Kim, Hosomichi-Seong-Terashimam, Källén-Qiu-Zabine, Lockhart-Vafa ...]

Goal: Develop a method to compute local correlators for 5d SCFTs How do we compute local correlators? Any? Simplest ones:

$$\left\langle T_{\mu\nu}(x)T_{\rho\sigma}(0)\right\rangle = \frac{C_T}{V_{\hat{\mathrm{S}}^4}^2}\frac{\mathcal{I}_{\mu\nu,\sigma\rho}(x)}{x^{10}}, \quad \left\langle J_{\mu}^{\mathsf{a}}(x)J_{\nu}^{\mathsf{b}}(0)\right\rangle = \frac{C_J}{V_{\hat{\mathrm{S}}^4}^2}\frac{\delta^{\mathsf{ab}}I_{\mu\nu}(x)}{x^8}.$$

Local correlators in 5d N = 1 SCFT

In [Chang-Fluder-Ying-YW: 1710.08418]

- Look at quadratic **deformation** of S^5 partition function by squashing parameters a_i and flavor masses M^a .
- Classify supergravity counter-terms to make sure coefficients of deformations in F_{S^5} do not suffer from **ambiguities**
- Do conformal perturbation theory and solve superconformal Ward identities

$$\begin{aligned} F_{\mathrm{S}^{5}}|_{a_{i}^{2}} &= -\frac{\pi^{2}C_{T}}{1920} \left(\sum_{i=1}^{3} a_{i}^{2} - \sum_{i < j} a_{i} a_{j} \right) \\ F_{\mathrm{S}^{5}}|_{M^{2}} &= \frac{3\pi^{2}r^{2}C_{J}}{256} \delta_{ab} M^{a} M^{b} \end{aligned}$$

• New 5d superconformal anomaly, conformal bootstrap, non-BPS spectrum, symmetric enhancement, testing F-theorem, C-theorem ...

KK-reduction of type I' SUGRA on $AdS_6 \times_w M_4$

For AdS/CFT duals

$$C_T \leftrightarrow \frac{1}{G_N}, \quad C_J \leftrightarrow \frac{1}{e^2}$$

Goal: Perform the KK reduction of the 10d type I' SUGRA solution describing AdS_6/CFT_5 duals \rightarrow fix G_N, e . Then compare with field theory answers C_T, C_J .

- Complication: warped AdS_6 solutions. singularity at the boundary of M_4
- Missing gauge fields for flavor symmetries: from D8 flavor branes
- Integrate over M₄ gives finite answer for 2-derivative couplings! (SUSY non-renormalization theorems)

$$\frac{1}{\kappa^2} \int_{AdS_6 \times M^4} \sqrt{G} (R + F^2 + \dots) \rightarrow \int_{AdS_6} \sqrt{g_6} (\frac{1}{G_N} R_6 + \frac{1}{e^2} F_6^2 + \dots)$$

In [Chang-Fluder-Ying-YW: 1712.10313]

- A class of AdS_6/CFT_5 duals (Seiberg-Morrison theories)
- SCFT described by USp(2N) $\mathcal{N} = 1$ SYM, 1 antisymmetric and $N_f \leqslant 7$ fundamental hypers. It has type I' SUGRA dual.
- Instanton contributions in F_{S^5} suppressed at large N. Solve matrix model to get analytic results

$$C_T = \frac{1152\sqrt{2}}{\pi\sqrt{8-N_f}} N^{\frac{5}{2}}, \quad C_J(E_{N_f+1}) = \frac{256\sqrt{2}}{3\pi\sqrt{8-N_f}} N^{\frac{3}{2}}.$$
(2.1)

- Matches with supergravity action on AdS₆ from massive IIA reduced on HS⁴ with D8 flavor branes (tricky warped geometry).
- More general *AdS*₆/*CFT*₅ pairs [Bergman-Rodriguez Gomez;...] were also analyzed. Symmetry enhancement explained.

AdS_7/CFT_6

Mezei-Pufu-YW: arXiv:1807.XXXXX

3d Chern-Simons on S^3 from 6d (2,0) SCFT on $S^1 imes S^5$

• The mysterious 6d (2,0) SCFT is intrinsically non-Lagrangian. But on S^1 it is believed to reduce to 5d $\mathcal{N} = 2$ SYM (ADE gauge group, higher derivative terms). In particular

6d (2,0) ADE SCFT on $S^1 \times S^5 \leftrightarrow$ 5d ADE MSYM on S^5

• In [Kim-Kim: 1206.6339], they noticed a "coincidence": the S^5 partition function of 5d U(N) MSYM contains a U(N) CS matrix model

$$Z_{S^5} \sim \int d^N \lambda \, e^{i\pi k \sum_{i=1}^N \lambda_i^2} \prod_{i < j} \left[2 \sinh \left[\pi (\lambda_i - \lambda_j) \right] \right]^2 \, .$$

Imaginary CS level

$$k = i \frac{4\pi^2 R}{g_{\rm YM}^2}$$

 $\mbox{Goal}:$ Take the coincidence seriously: find the 3d CS sector of 6d (2,0) SCFT

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3d Chern-Simons on S^3 from 6d (2,0) SCFT on $S^1 imes S^5$

In [Mezei-Pufu-YW: arXiv:1807.XXXXX]

 \bullet Localizing supercharge ${\mathcal Q}$ such that

$$\mathcal{Q}^2 = iP_\tau + R_{45}$$

 S^5 as $S^1 \times S^3$ fibered over an interval $\theta = \left[0, \frac{\pi}{2}\right]$.

- BPS locus parametrized by emergent gauge fields A on S³ at θ = 0 and 5d instantons located at the θ = π/2
- Weighted by CS action

$$Z_{S^5} \sim \int D\mathcal{A} \exp \frac{ik}{4\pi} \int_{S^3} \operatorname{Tr} \left(\mathcal{A} d\mathcal{A} + \frac{2}{3} \mathcal{A}^3 \right)$$

with **imaginary** level $k \sim i \frac{R_6}{R_{c5}}$.



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 $rac{1}{8}$ BPS surface operators and knots on S^3

CS sector of 6d (2,0) on $S^1 \times S^5$

• New SUSY observables beyond the superconformal index (SCI):



- CS interpretation → These surface operators are topological (invariant under infinitesimal shape changes).
- Refining the **chiral algebra** limit of 6d SCI [Beem-Rastelli-van Rees, Bullimore-Kim].

- Surface operator on Σ in 6d SCFT are dual to M2-branes on \mathcal{M} in $AdS_7 \times S^4$ such that $\partial \mathcal{M} = \Sigma$ (winds in both ∂AdS_7 and S^4 , complicated shape!)
- Need to minimize M2 brane action

$$S_{\mathrm{M2}} = au_{\mathrm{M2}} \int_{\mathcal{M}} d^3 \sigma \sqrt{-g} + i au_{\mathrm{M2}} \int_{\mathcal{M}} A_3 + i S_{\mathrm{WZW}}$$

- Difficult. SUSY helps → calibrated submanifolds M (1st order differential equations). No log ambiguities.
- Need appropriate counter-terms
- Answer matches with field theory results computed from the 3d CS theory (topological!).

Outlook

 AdS_4/CFT_3 :

- Extract predictions for higher derivative couplings in M-theory (recall Eric's talk).
- Understanding the 2d protected sector of M-theory from topological twist [Kostello...]
- Generalization to other large N 3d SCFTs.
- embedding of SYK/tensor models?

 AdS_6/CFT_5 :

- Beyond two point functions
- Large N limit of other 5d SCFTs

 AdS_7/CFT_6

- Relation to other 3d CS sectors of (2,0) theory [Dimofte-Gaiotto-Gukov, Cordova-Jafferis, Lee-Yamazaki ...]
- Identify the dual 4d protected sector of M-theory
- Extend the 6d analysis to (1,0) theories

Thank you!