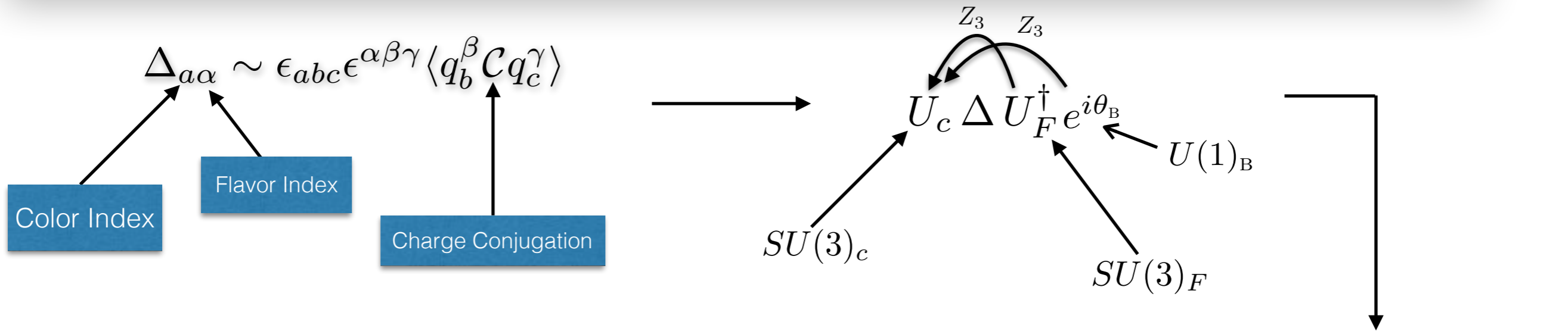


Non Abelian Vortices and Aharonov-Bohm Effect in Dense QCD

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$OPS = \frac{SU(3) \times U(1)}{Z_3} = U(3)$

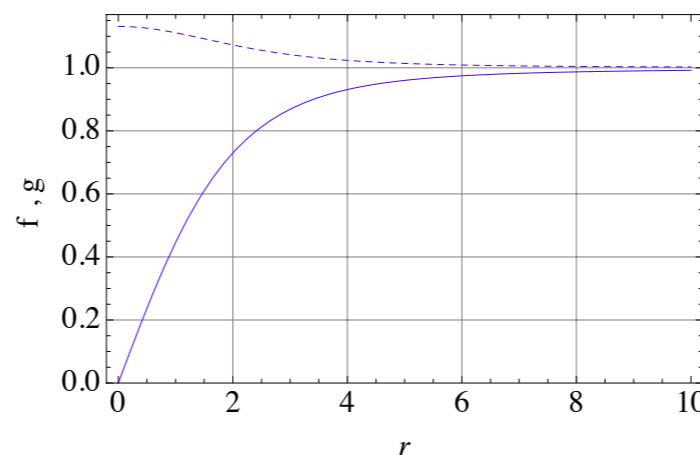
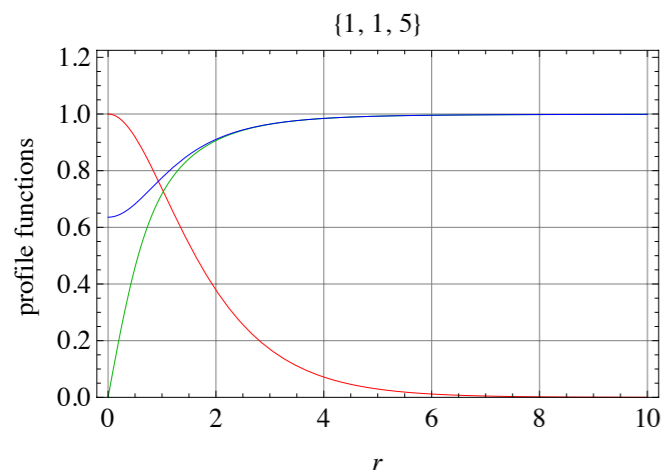
$\Pi_1(OPS) = Z$

$H = \frac{SU(3)_{C+F}}{Z_3}$

$G = \frac{U(1)_B \times SU(3)_C \times SU(3)_F}{Z_3 \times Z_3}$

$\Phi(r, \theta) = \Delta_{CFL} \begin{pmatrix} e^{i\theta} f(r) & 0 & 0 \\ 0 & g(r) & 0 \\ 0 & 0 & g(r) \end{pmatrix}, \quad A_i(r) = -\frac{1}{3g_s} \frac{\epsilon_{ij} x_j}{r^2} A(r) \begin{pmatrix} 2 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & -1 \end{pmatrix}. \quad A(0) = 0, \quad A(R) = 1.$

$f(0) = g'(0) = 0, \quad f(R) = g(R) = 1$



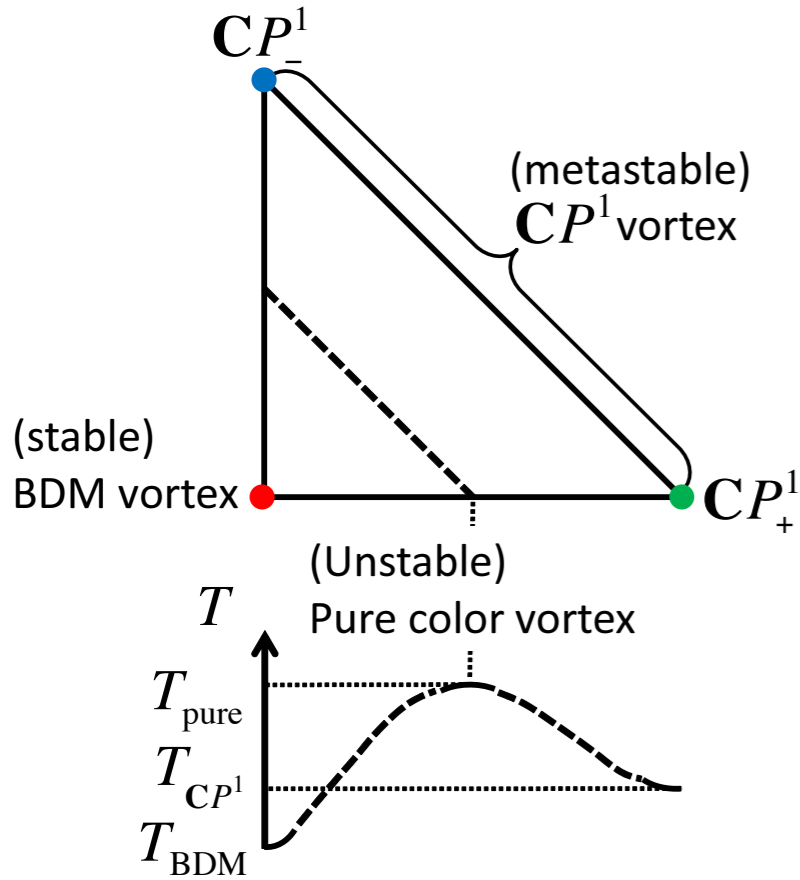
$SU(3)_{C+F} \xrightarrow{\Delta_{CFL}} SU(2) \times U(1)$

\downarrow

CP^2 Moduli

Electro Magnetic Interaction

$$D_\mu \Phi = \partial_\mu \Phi - ig_s A_\mu^a \tau^a \Phi - ie B_\mu \Phi Q, \tau^a \in SU(3)_C \quad D_\mu \Phi_{diag} = [\partial_\mu - i(g A_\mu^a H^a - e Q_{em} B_\mu)] \Phi_{diag}, H^a \in \{\tau^3, \tau^8\}$$



$$Q_{em} = \frac{1}{3} \begin{pmatrix} 2 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & -1 \end{pmatrix}$$

$U(1)_{em}$ generator

$$A_\mu^M = \frac{g_s}{g_M} A_\mu^8 - \frac{\eta e}{g_M} B_\mu, \quad A_\mu^q = \frac{\eta e}{g_M} A_\mu^8 + \frac{g_s}{g_M} B_\mu$$

$$\eta = \frac{2}{\sqrt{3}}, \quad g_M^2 = g_s^2 + \eta^2 e^2$$

Massive gauge field

Mass less Gauge field

$$B_\mu = \frac{g_s}{g_M} A_\mu^q - \frac{\eta e}{g_M} A_\mu^M$$

$$q_{eff} = \frac{q g_s}{\sqrt{g_s^2 + \eta^2 e^2}}$$

Vortex Configurations:

BDM Vortex:

$$\Phi(r, \theta) = \Delta_{CFL} \begin{pmatrix} e^{i\theta} f(r) & 0 & 0 \\ 0 & g(r) & 0 \\ 0 & 0 & g(r) \end{pmatrix}, \quad A_i^M(r) T^8 = -\frac{\epsilon_{ij} x_j}{3g_M r^2} [1 - h(r)] \begin{pmatrix} 2 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & -1 \end{pmatrix}.$$

CP¹ Vortex:

$$\Phi(r, \theta) = \Delta_{CFL} \begin{pmatrix} g(r) & 0 & 0 \\ 0 & e^{i\theta} f(r) & 0 \\ 0 & 0 & g(r) \end{pmatrix}, \quad A_i^M(r) T^8 = \frac{1}{6g_M} \frac{\epsilon_{ij} x_j}{r^2} [1 - h(r)] \begin{pmatrix} 2 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & -1 \end{pmatrix}, \quad A_i^3(r) T^3 = -\frac{1}{2g_s} \frac{\epsilon_{ij} x_j}{r^2} [1 - h(r)] \begin{pmatrix} 0 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & -1 \end{pmatrix}.$$

[Balachandran, et al. 2005]

[Nitta, et al. 2012]

Aharonov-Bohm Scattering

Electric Charge

$$(\not{\partial} - iq\not{B}_\mu + iM_{e/\mu_e}) \psi_{e/\mu_e} = 0$$

$$\frac{d\sigma}{d\vartheta} = \frac{\sin^2(\pi\varphi)}{2\pi k \sin^2(\frac{\vartheta}{2})}, \varphi = \frac{q}{2\pi} \times \text{Flux.}$$

$$\varphi = \left| \frac{q}{2\pi} \oint B_\mu \cdot dl \right| = \left| \frac{\eta e q}{2\pi g_M} \oint A^M \cdot dl \right|$$

Scattering Angle

$$\varphi_{e/\mu_e}^{\text{BDM}} = \frac{2e^2}{3g_s^2 + 2e^2}$$

$$\varphi_{e/\mu_e}^{\text{CP}^1} = \frac{e^2}{3g_s^2 + 2e^2}$$

$$\varphi_\Sigma^{\text{BDM}} = \frac{2e^2}{3g_s^2 + 2e^2}, \varphi_\Sigma^{\text{CP}^1} = \frac{e^2}{3g_s^2 + 2e^2}$$

$$\Sigma_{\text{cfl}}^{AB} = \Phi_{Aa}^{\dagger L} \Phi_{aB}^R$$

$$\Sigma_{\text{cfl}}^{AB} \sim \epsilon^{ACD} \epsilon^{BEF} \bar{q}_{L(a}^C \bar{q}_{Lb)}^D q_{R(a}^E q_{Rb)}^F$$

$$\Sigma'_{\text{cfl}} = e^{ieQ\alpha} \Sigma_{\text{cfl}} e^{-ieQ\alpha}$$

$$\begin{pmatrix} 0 & \Sigma_{\text{cfl}}^{1+} & \Sigma_{\text{cfl}}^{2+} \\ \Sigma_{\text{cfl}}^{1-} & 0 & 0 \\ \Sigma_{\text{cfl}}^{2-} & 0 & 0 \end{pmatrix}$$

Relaxation rate for particles scattering off vortices:

Vortex density:

$$n_v \simeq 1.9 \times 10^{19} \left(\frac{1 \text{ ms}}{P_{\text{rot}}} \right) \left(\frac{\mu/3}{300 \text{ MeV}} \right) \left(\frac{R}{10 \text{ km}} \right)^2$$

$$s_0 = \frac{1}{\tau_v} = \frac{n_v}{k_F} \sin^2 \pi\phi$$

Relaxation rate for Coulomb interaction:

$$\frac{1}{\tau_q} = \frac{6\zeta(3)}{\pi^2} \tilde{\alpha} T$$

The temperature T_f below which flux tubes dominate the relaxation of deviations from thermal equilibrium is determined as

$$T_f = \frac{\pi^2}{6\zeta(3)} \frac{\sin^2(\pi\phi)}{\tilde{\alpha}} \frac{n_v}{k_F}$$

Thank you