

Odd-frequency superconductivity in Sr_2RuO_4 measured by Kerr effect

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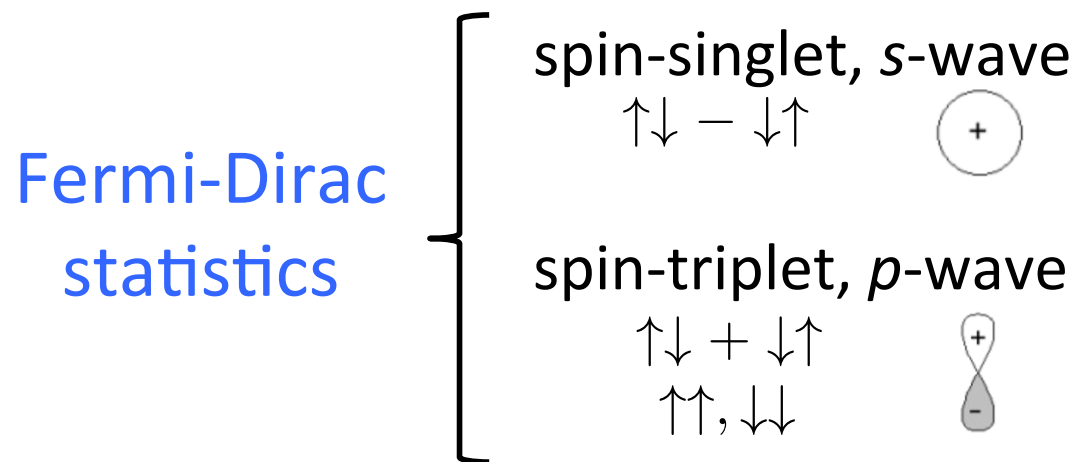
Outline

- Introduction to odd-frequency (ω) superconductivity
 - SF, SN junctions
 - Odd- ω bulk superconductivity in multiband systems
 - Simple two-band superconductors
 - Odd- ω superconductivity and Kerr effect in Sr_2RuO_4
-

Different Phases of Matter

Ordered states \rightarrow symmetry of order parameter, Δ

Superconductivity: $\Delta \sim F = \langle \psi_\alpha \psi_\beta \rangle$







Different Phases of Matter

Ordered states \rightarrow symmetry of order parameter, Δ

Superconductivity: $\Delta \sim F = \langle \psi_\alpha(t) \psi_\beta(0) \rangle$ ($t \leftrightarrow \omega$, frequency)

Fermi-Dirac
statistics

- even- ω , spin-singlet, s-wave
 $\uparrow\downarrow - \downarrow\uparrow$ 
- odd- ω** , spin-triplet, s-wave
 $\uparrow\downarrow + \downarrow\uparrow$
 $\uparrow\uparrow, \downarrow\downarrow$ 

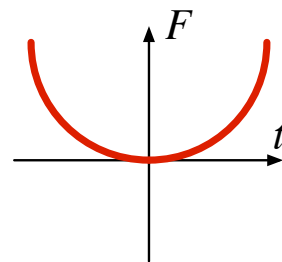


Odd- ω Pairing

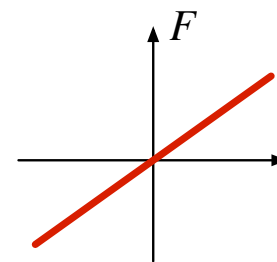
BCS order parameter: $F(\mathbf{r}, t; \mathbf{r}', t' \rightarrow t) = \langle \psi(\mathbf{r}, t) \psi(\mathbf{r}', t' \rightarrow t) \rangle$

vanishes for odd- ω pairing

Equal-time odd- ω order parameter: $\left. \frac{dF(\mathbf{r}, t; \mathbf{r}', t')}{dt} \right|_{t \rightarrow t'}$



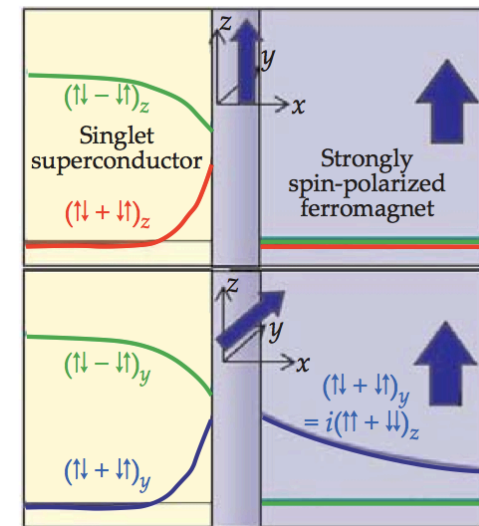
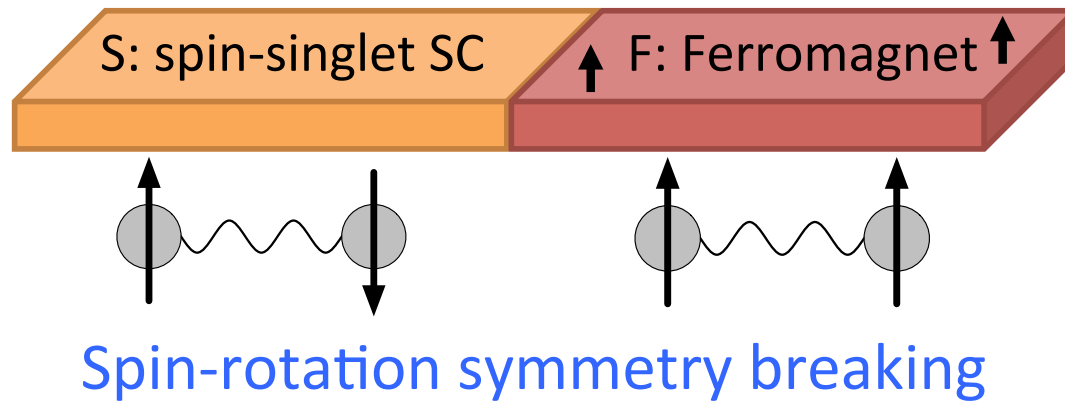
even- ω



odd- ω



SF Interface



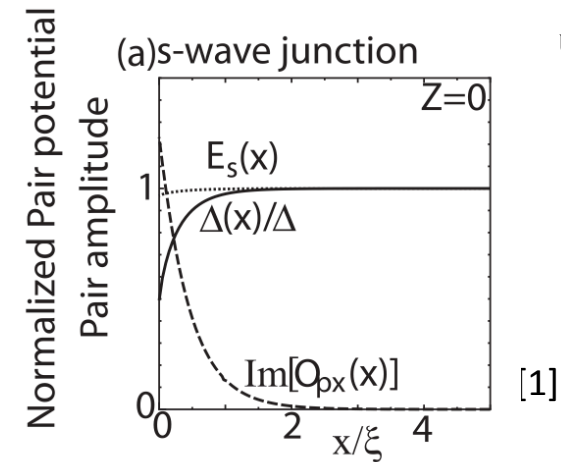
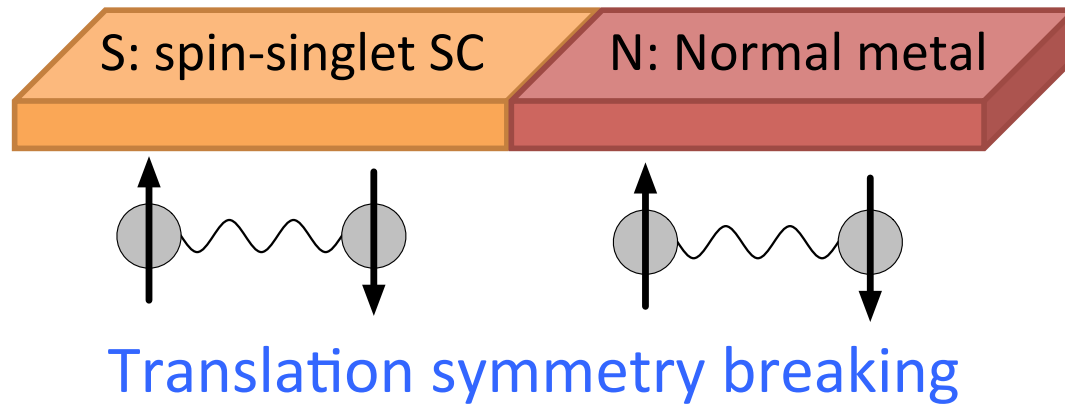
[1]

Spin-singlet *s*-wave SC → odd- ω spin-triplet *s*-wave pairing

- Long-range superconducting proximity effect in F
- *s*-wave = disorder robust



SN Interface



Spin-singlet *s*-wave SC → odd- ω spin-singlet *p*-wave pairing

- Only high-transparency junctions
- *p*-wave = only ballistic systems

[1]: Tanaka et al, PRL 99, 037005 (2007)



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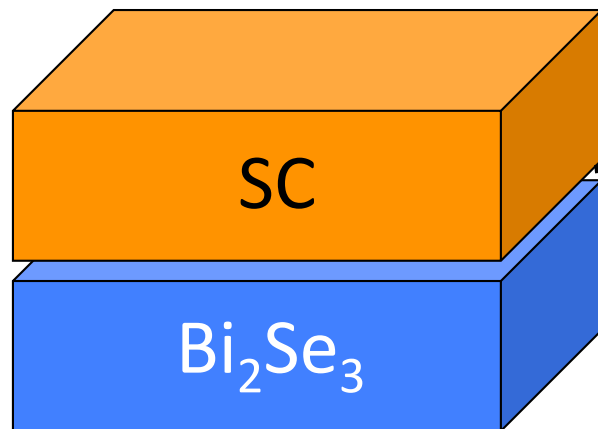
Odd- ω Bulk Superconductivity in Multiband Systems

- A clue from Bi_2Se_3
 - Simple two-band superconductors
-

Bi₂Se₃ – SC Hybrid Structure



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$$H_{SC} = \sum_{\mathbf{k}, \sigma} (-2 \cos(k_x a) - 2 \cos(k_y a) + \mu_{SC}) c_{\mathbf{k}\sigma}^\dagger c_{\mathbf{k}\sigma} + \frac{1}{2} \sum_{\mathbf{k}, \sigma, \sigma'} \Delta_{\sigma\sigma'}(\mathbf{k}) c_{\mathbf{k}\sigma}^\dagger c_{-\mathbf{k}\sigma'}^\dagger - \Delta_{\sigma\sigma'}^*(-\mathbf{k}) c_{-\mathbf{k}\sigma} c_{\mathbf{k}\sigma'}$$

2D superconductor

$$H_T = - \sum_{\mathbf{k}, \sigma} T_1 c_{\mathbf{k}\sigma}^\dagger b_{1\mathbf{k}\sigma} + T_2 c_{\mathbf{k}\sigma}^\dagger b_{2\mathbf{k}\sigma} + \text{H.c.}$$

Local tunneling

$$H_{TI} = \gamma_0 - 2 \sum_{\mathbf{k}, i} \gamma_i \cos(k_i a) + \sum_{\mathbf{k}, \mu} d_\mu \Gamma_\mu$$

Bi₂Se₃ (τ orbital Pauli matrix) [1]

$$\left[\begin{array}{l} d_0 = \epsilon - 2 \sum_i t_i \cos(k_i a), d_i = -2\lambda_i \sin(k_i a) \\ \Gamma_0 = \tau_x \otimes \sigma_0, \Gamma_x = -\tau_z \otimes \sigma_y, \Gamma_y = \tau_z \otimes \sigma_x, \Gamma_z = \tau_y \otimes \sigma_0 \end{array} \right]$$

Superconducting Symmetries in Bi_2Se_3



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Singlet/triplet, spatial ($s/p/d$), even/odd- ω , even/odd orbital

Superconductor			Even-frequency		Odd-frequency	
Γ	Basis function	J_z	Even-orbital	Odd-orbital	Even-orbital	Odd-orbital
A_{1g}	$\psi = 1$	0	A_{1g} singlet, A_{2u} triplet ($m_s = \pm 1$)	-	-	A_{1g} singlet, A_{2u} triplet ($m_s = \pm 1$)
B_{1g}	$\psi = k_x^2 - k_y^2$	± 2	B_{1g} singlet, B_{2u} triplet ($m_s = \pm 1$)	-	-	B_{1g} singlet, B_{2u} triplet ($m_s = \pm 1$)
B_{2g}	$\psi = 2k_x k_y$	± 2	B_{2g} singlet, B_{1u} triplet ($m_s = \pm 1$)	-	-	B_{2g} singlet, B_{1u} triplet ($m_s = \pm 1$)
A_{1u}	$\mathbf{d} = (k_x, k_y, 0)$	0	A_{1u} triplet ($m_s = \pm 1$)	A_{1g} triplet ($m_s = 0$)	A_{1g} triplet ($m_s = 0$)	A_{1u} triplet ($m_s = \pm 1$)
A_{2u}	$\mathbf{d} = (k_y, -k_x, 0)$	0	A_{2u} triplet ($m_s = \pm 1$), A_{1g} singlet	-	-	A_{2u} triplet ($m_s = \pm 1$), A_{1g} singlet
B_{1u}	$\mathbf{d} = (k_x, -k_y, 0)$	± 2	B_{1u} triplet ($m_s = \pm 1$), B_{2g} singlet	B_{1g} triplet ($m_s = 0$)	B_{1g} triplet ($m_s = 0$)	B_{1u} triplet ($m_s = \pm 1$), B_{2g} singlet
B_{2u}	$\mathbf{d} = (k_y, k_x, 0)$	± 2	B_{2u} triplet ($m_s = \pm 1$), B_{1g} singlet	B_{2g} triplet ($m_s = 0$)	B_{2g} triplet ($m_s = 0$)	B_{2u} triplet ($m_s = \pm 1$), B_{1g} singlet
E_{2u}^+	$\mathbf{d} = (0, 0, k_x + ik_y)$	1	E_{2u}^+ triplet ($m_s = 0$)	A_{1g} triplet ($m_s = 1$), $B_{1g} + iB_{2g}$ triplet ($m_s = -1$)	A_{1g} triplet ($m_s = 1$), $B_{1g} + iB_{2g}$ triplet ($m_s = -1$)	E_{2u}^+ triplet ($m_s = 0$)
E_{2u}^-	$\mathbf{d} = (0, 0, k_x - ik_y)$	-1	E_{2u}^- triplet ($m_s = 0$)	A_{1g} triplet ($m_s = -1$), $B_{1g} - iB_{2g}$ triplet ($m_s = 1$)	A_{1g} triplet ($m_s = -1$), $B_{1g} - iB_{2g}$ triplet ($m_s = 1$)	E_{2u}^- triplet ($m_s = 0$)

ABS and Balatsky, PRB 87, 220506(R) (2013)

Frequency and Orbital



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Complete reciprocity in oddness in frequency and orbital index

Superconductor			Even-frequency		Odd-frequency	
Γ	Basis function	J_z	Even-orbital	Odd-orbital	Even-orbital	Odd-orbital
A _{1g}	$\psi = 1$	0	A _{1g} singlet, A _{2u} triplet ($m_s = \pm 1$)	-	-	A _{1g} singlet, A _{2u} triplet ($m_s = \pm 1$)
B _{1g}	$\psi = k_x^2 - k_y^2$	± 2	B _{1g} singlet, B _{2u} triplet ($m_s = \pm 1$)	-	-	B _{1g} singlet, B _{2u} triplet ($m_s = \pm 1$)
B _{2g}	$\psi = 2k_x k_y$	± 2	B _{2g} singlet, B _{1u} triplet ($m_s = \pm 1$)	-	-	B _{2g} singlet, B _{1u} triplet ($m_s = \pm 1$)
A _{1u}	$\mathbf{d} = (k_x, k_y, 0)$	0	A _{1u} triplet ($m_s = \pm 1$)	A _{1g} triplet ($m_s = 0$)	A _{1g} triplet ($m_s = 0$)	A _{1u} triplet ($m_s = \pm 1$)
A _{2u}	$\mathbf{d} = (k_y, -k_x, 0)$	0	A _{2u} triplet ($m_s = \pm 1$) A _{1g} singlet	-	-	A _{2u} triplet ($m_s = \pm 1$), A _{1g} singlet
B _{1u}	$\mathbf{d} = (k_x, -k_y, 0)$	± 2	B _{1u} triplet ($m_s = \pm 1$), B _{2g} singlet	B _{1g} triplet ($m_s = 0$)	B _{1g} triplet ($m_s = 0$)	B _{1u} triplet ($m_s = \pm 1$), B _{2g} singlet
B _{2u}	$\mathbf{d} = (k_y, k_x, 0)$	± 2	B _{2u} triplet ($m_s = \pm 1$), B _{1g} singlet	B _{2g} triplet ($m_s = 0$)	B _{2g} triplet ($m_s = 0$)	B _{2u} triplet ($m_s = \pm 1$), B _{1g} singlet
E _{2u} ⁺	$\mathbf{d} = (0, 0, k_x + ik_y)$	1	E _{2u} ⁺ triplet ($m_s = 0$)	A _{1g} triplet ($m_s = 1$), B _{1g} +iB _{2g} triplet ($m_s = -1$)	A _{1g} triplet ($m_s = 1$), B _{1g} +iB _{2g} triplet ($m_s = -1$)	E _{2u} ⁺ triplet ($m_s = 0$)
E _{2u} ⁻	$\mathbf{d} = (0, 0, k_x - ik_y)$	-1	E _{2u} ⁻ triplet ($m_s = 0$)	A _{1g} triplet ($m_s = -1$), B _{1g} -iB _{2g} triplet ($m_s = 1$)	A _{1g} triplet ($m_s = -1$), B _{1g} -iB _{2g} triplet ($m_s = 1$)	E _{2u} ⁻ triplet ($m_s = 0$)

ABS and Balatsky, PRB 87, 220506(R) (2013)



Multiband Superconductors

- S: Spin (even: spin-triplet; odd: spin-singlet)
- P: Spatial parity (even: s, d -wave; odd: p, f -wave)
- O: Orbital or band parity (even; odd orbital)
- T: Time (even; odd-frequency)

SPOT = -1

S = 0	P	T	O	S = 1	P	T	O
even- ω	+	+	+	even- ω	-	+	+
even- ω	-	+	-	even- ω	+	+	-
odd- ω	+	-	-	odd- ω	+	-	+
odd- ω	-	-	+	odd- ω	-	-	-



Simple Two-Band Superconductor

$$H_{ab} = \sum_{k\sigma} \varepsilon_a(k) a_{k\sigma}^\dagger a_{k\sigma} + \varepsilon_b(k) b_{k\sigma}^\dagger b_{k\sigma} \\ + \sum_k \Delta_a(k) a_{k\uparrow}^\dagger a_{-k\downarrow}^\dagger + \Delta_b(k) b_{k\uparrow}^\dagger b_{-k\downarrow}^\dagger + \text{H.c.}$$

Interband
hybridization



$$H_{cd} = \sum_{k\sigma} \varepsilon_c(k) c_{k\sigma}^\dagger c_{k\sigma} + \varepsilon_d(k) d_{k\sigma}^\dagger d_{k\sigma} \quad \text{Diagonal bands} \\ + \sum_k \Delta_c(k) c_{k\uparrow}^\dagger c_{-k\downarrow}^\dagger + \Delta_d(k) d_{k\uparrow}^\dagger d_{-k\downarrow}^\dagger + \text{H.c.} \quad \text{Intraband pairing} \\ + \sum_k \Delta_{cd}(k) \left[c_{k\uparrow}^\dagger d_{-k\downarrow}^\dagger + d_{k\uparrow}^\dagger c_{-k\downarrow}^\dagger \right] + \text{H.c.}$$

$$\Delta_{cd} = \frac{(\Delta_b - \Delta_a) |\Gamma|}{\sqrt{(\varepsilon_a - \varepsilon_b)^2 + 4|\Gamma|^2}}$$



Time-Dependent Pairing

Time-ordered s-wave interband pairing:

$$F^{\pm}(\tau) = \frac{1}{2N_{\mathbf{k}}} \sum_{\mathbf{k}} \mathcal{T}_{\tau} \langle c_{-\mathbf{k}\downarrow}(\tau) d_{\mathbf{k}\uparrow}(0) \pm d_{-\mathbf{k}\downarrow}(\tau) c_{\mathbf{k}\uparrow}(0) \rangle$$

$$F^e = F^+(\tau \rightarrow 0^+) \quad \text{Even-}\omega, \text{ even-interband pairing}$$

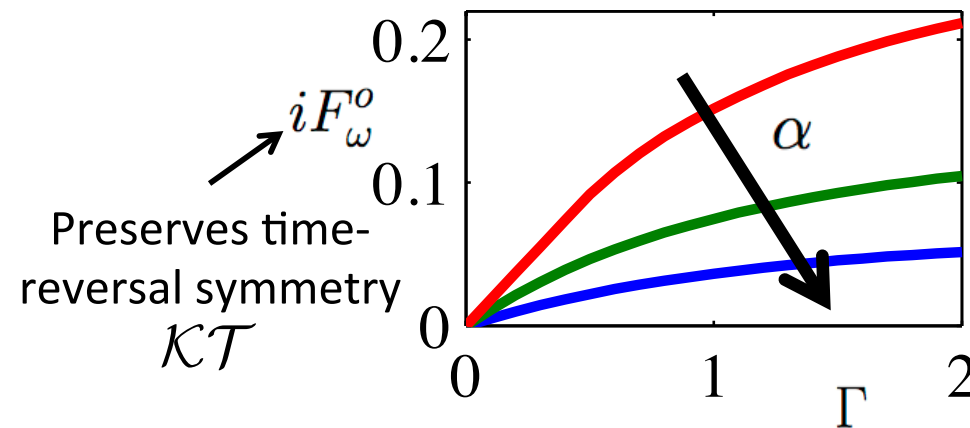
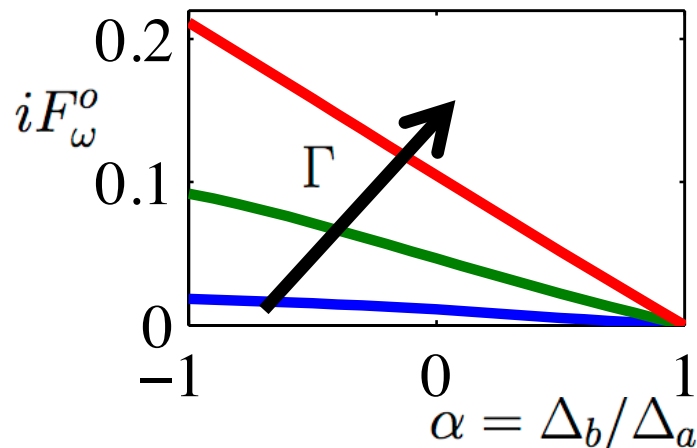
$$F_{\omega}^o = \left. \frac{\partial F^-}{\partial \tau} \right|_{\tau \rightarrow 0^+} \quad \text{Odd-}\omega, \text{ odd-interband pairing}$$

Odd- ω , Odd-Interband Pairing



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For $\begin{cases} \varepsilon_a = \varepsilon_b \\ \Delta_a = -\Delta_b \end{cases} \longrightarrow \begin{cases} \varepsilon_{c,d} = \varepsilon_a \mp \Gamma \\ \Delta_c = \Delta_d = 0 \\ \Delta_{cd} = \Delta_a \text{ Interband pairing} \end{cases} \xrightarrow{\Gamma < \Delta_a} \begin{cases} F^e = -\frac{1}{2N_k} \sum_k \frac{\Delta_a}{\sqrt{\varepsilon_a^2 + |\Delta_a|^2}} \\ \text{BCS equation} \\ F_\omega^o = i\Gamma F^e \quad \text{Odd-}\omega \end{cases}$

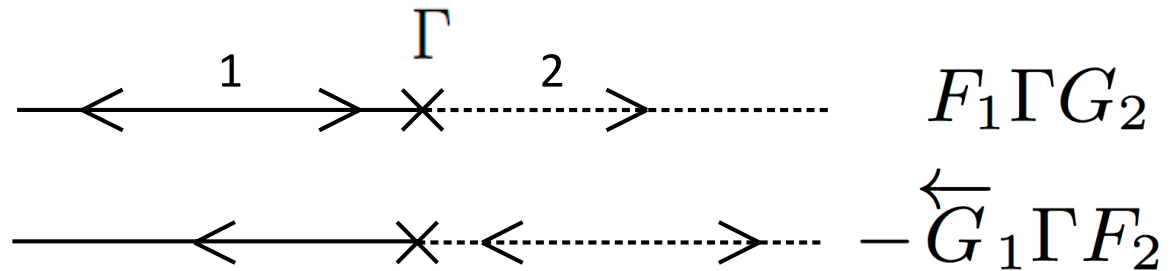




Using Perturbation Theory

$$H = \sum_{k\sigma} \varepsilon_1(k) a_{k\sigma}^\dagger a_{k\sigma} + \varepsilon_2(k) b_{k\sigma}^\dagger b_{k\sigma} + \sum_{k\sigma} \Gamma(k) a_{k\sigma}^\dagger b_{k\sigma} + \text{H.c.}$$
$$+ \sum_k \Delta_1(k) a_{k\uparrow}^\dagger a_{-k\downarrow}^\dagger + \Delta_2(k) b_{k\uparrow}^\dagger b_{-k\downarrow}^\dagger + \text{H.c.}$$

Interband pairing F_{12} :





Interband Pairing

Perturbation theory to infinite order in Γ

$$\text{Odd-interband: } F_{12}^{\text{odd}}(k, i\omega) = \frac{F_{12} - F_{21}}{2} = i\omega\Gamma(\Delta_1 - \Delta_2)/D$$

$$\text{Even-interband: } F_{12}^{\text{even}}(k, i\omega) = \frac{F_{12} + F_{21}}{2} = \Gamma(\Delta_1\varepsilon_2 - \Delta_2\varepsilon_1)/D$$

$$\left[\begin{array}{l} D = (\omega^2 + E_1^2)(\omega^2 + E_2^2) - \Gamma^2[2(\varepsilon_1\varepsilon_2 - \omega^2) - \Delta_2^*\Delta_1 - \Delta_1^*\Delta_2] + \Gamma^4 \\ E_j^2 = \varepsilon_j^2 + |\Delta_j|^2 \end{array} \right]$$

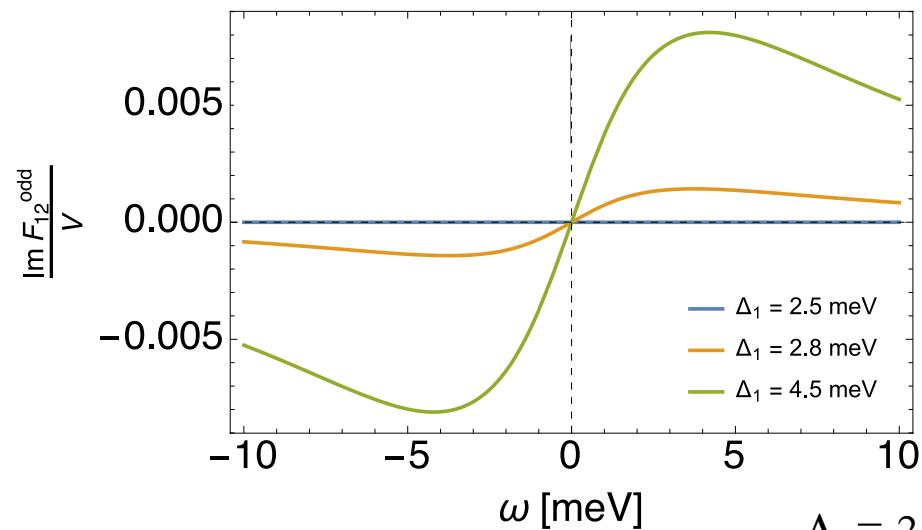
Odd-frequency pairing: $\Gamma \neq 0, \Delta_1 \neq \Delta_2$

Interband Frequency Dependence



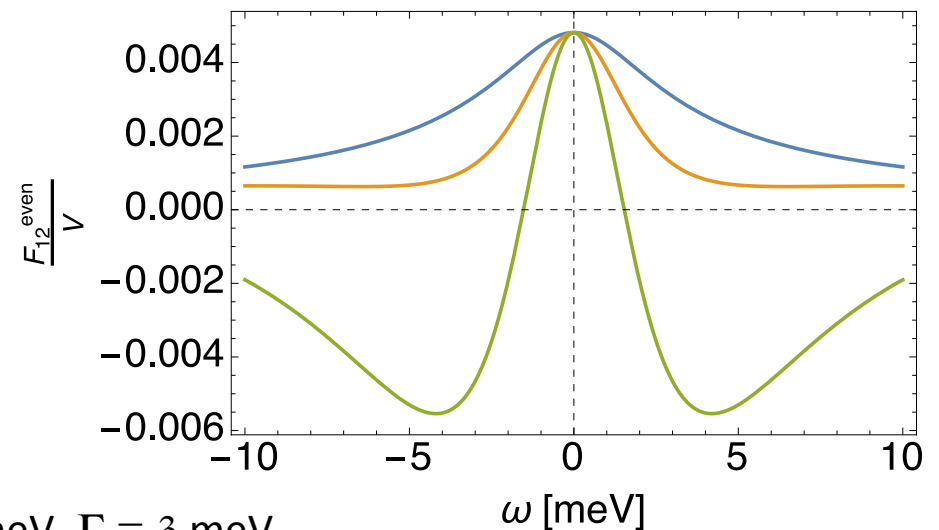
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Odd-frequency



$\Delta_2 = 2.5 \text{ meV}, \Gamma = 3 \text{ meV}$

Even-frequency



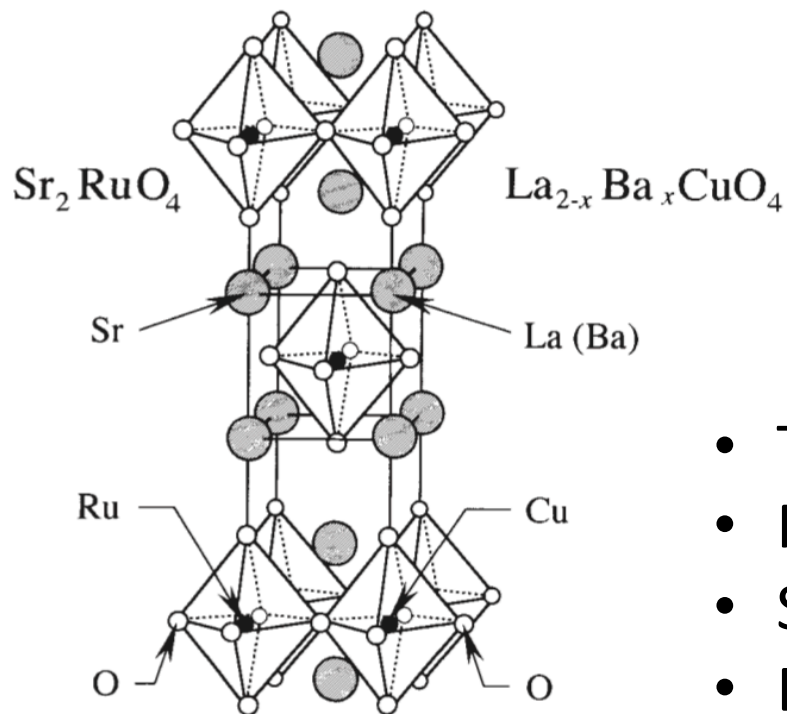
Odd-frequency pairing: $\Gamma \neq 0, \Delta_1 \neq \Delta_2$



Odd- ω Superconductivity in Sr_2RuO_4

- Superconductivity in Sr_2RuO_4
 - Two- and three-orbital models
 - Odd- ω pairing measured by Kerr effect
-

Strontium Ruthenate, Sr_2RuO_4



Superconductivity in a layered perovskite without copper

Y. Maeno*, H. Hashimoto*, K. Yoshida*,
S. Nishizaki*, T. Fujita*, J. G. Bednorz†
& F. Lichtenberg†‡

Nature 372, 532 (1994)

- $T_c = 1.5\text{K}$
- Non-s-wave (disorder sensitive)
- Spin-triplet (neutron scattering, Knight shift)
- Breaks time-reversal symmetry (Kerr effect)

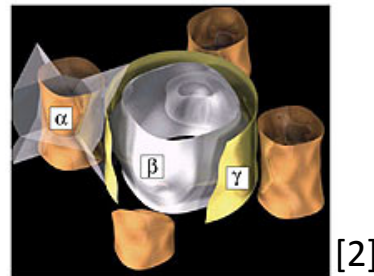
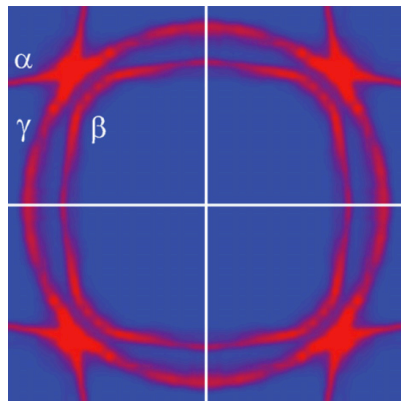
→ Spin-triplet chiral p_x+ip_y -wave symmetry

Properties of Sr_2RuO_4



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Three Fermi surfaces (FSs)

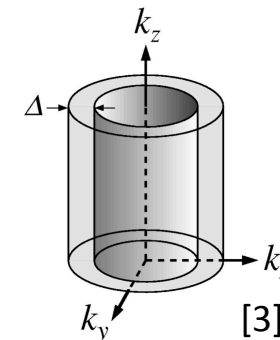


Superconducting state

$$\begin{pmatrix} \Delta_{\uparrow\uparrow}(\mathbf{k}) & \Delta_{\uparrow\downarrow}(\mathbf{k}) \\ \Delta_{\downarrow\uparrow}(\mathbf{k}) & \Delta_{\downarrow\downarrow}(\mathbf{k}) \end{pmatrix} = \begin{pmatrix} -d_x(\mathbf{k}) + id_y(\mathbf{k}) & \psi(\mathbf{k}) + d_z(\mathbf{k}) \\ -\psi(\mathbf{k}) + d_z(\mathbf{k}) & d_x(\mathbf{k}) + id_y(\mathbf{k}) \end{pmatrix}$$

Triplet \mathbf{d} -vector for chiral state

$$\mathbf{d}(\mathbf{k}) = (0, 0, k_x \pm ik_y)$$



Fully gapped on
cylindrical FS

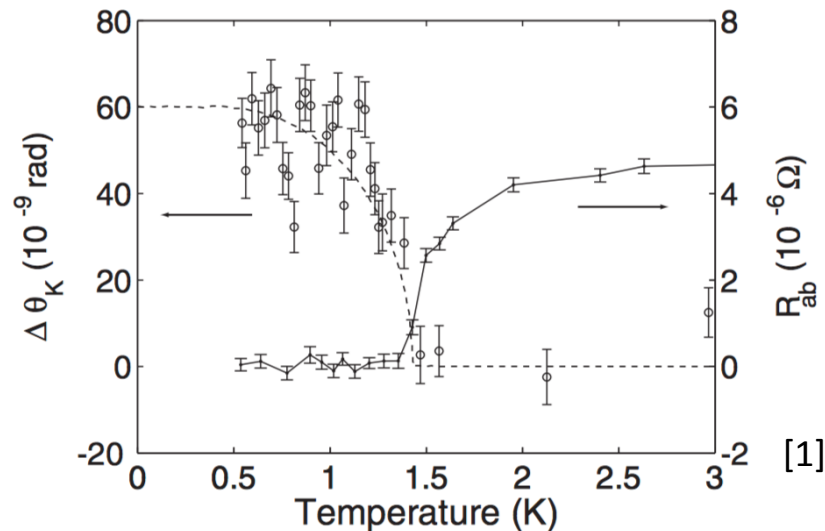
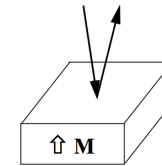
Three Ru 4d orbitals:

- $xy \rightarrow \gamma$ (electron-like)
- $xz, yz \rightarrow \beta$ (electron-like) and α (hole-like)



Kerr Effect in Sr_2RuO_4

Reflected light has a slightly rotated plane of polarization if material breaks time-reversal symmetry (TRS)



SC state in Sr_2RuO_4 breaks TRS

But ...

Clean **single-band** chiral SC has **zero** Kerr effect

→ **Interband pairing with relative superconducting phases** [2]

Electric-field driven interband transitions with relative SC phases
→ finite transverse Hall current response → finite Kerr effect



Two-Orbital Model for Sr_2RuO_4

Hamiltonian $\sum_{\mathbf{k}} \Psi_{\mathbf{k}}^\dagger \hat{H}_{\mathbf{k}} \Psi_{\mathbf{k}}$ with $\hat{H}_{\mathbf{k}} = \begin{pmatrix} \hat{H}_0(\mathbf{k}) & \check{\Delta}(\mathbf{k}) \\ \check{\Delta}^\dagger(\mathbf{k}) & -\hat{H}_0(-\mathbf{k}) \end{pmatrix}$

- xz (1), yz (2) orbitals
→ α, β bands
- Spin-triplet $\mathbf{d} \parallel \hat{\mathbf{z}}$

$$\Psi_{\mathbf{k}}^\dagger = (c_{\mathbf{k}\uparrow 1}^\dagger \ c_{\mathbf{k}\uparrow 2}^\dagger \ c_{-\mathbf{k}\downarrow 1} \ c_{-\mathbf{k}\downarrow 2})$$

$$\hat{H}_0(\mathbf{k}) = \begin{pmatrix} \xi_1 & \epsilon_{12} \\ \epsilon_{12} & \xi_2 \end{pmatrix}, \quad \check{\Delta}(\mathbf{k}) = \begin{pmatrix} \Delta_1 & \Delta_{12} \\ \Delta_{12} & \Delta_2 \end{pmatrix}$$

Intraorbital energy

Interorbital hybridization Interorbital pairing Intraorbital pairing



Odd- ω Pairing and Kerr Effect

Odd- ω , odd-interorbital pairing:

$$F_{12} - F_{21} = i\omega [(\Delta_2 - \Delta_1)\epsilon_{12} + \Delta_{12}(\xi_1 - \xi_2)] / D \quad (D \sim \omega^2)$$

Interorbital hybridization Interorbital pairing
+ gap asymmetry + dispersion asymmetry

Kerr effect: [1]

$$\sigma_H \propto \epsilon_{12} \text{Im}(\Delta_1^* \Delta_2) + \xi_1 \text{Im}(\Delta_2^* \Delta_{12}) - \xi_2 \text{Im}(\Delta_1^* \Delta_{12})$$

Interorbital hybridization Interorbital pairing
+ gap asymmetry + dispersion asymmetry

→ Intrinsic Kerr effect evidence of odd- ω superconductivity



Generic Three-Orbital Model

$$\hat{H}_0(\mathbf{k}) = \begin{pmatrix} \xi_1 & \epsilon_{12} & \epsilon_{13} \\ \epsilon_{12} & \xi_2 & \epsilon_{23} \\ \epsilon_{13} & \epsilon_{23} & \xi_3 \end{pmatrix}, \quad \check{\Delta}(\mathbf{k}) = \begin{pmatrix} \Delta_1 & \Delta_{12} & \Delta_{13} \\ \Delta_{12} & \Delta_2 & \Delta_{23} \\ \Delta_{13} & \Delta_{23} & \Delta_3 \end{pmatrix}$$

xy (1), xz (2), yz (3) orbitals $\rightarrow \gamma, \alpha, \beta$ bands

General interorbital pairing:

$$F_{AS} = \sum_{i,j,k=1,\dots,N} \epsilon_{ijk} F_{ij}$$

$$F_{AS}(-\omega) = -F_{AS}(\omega)$$

odd- ω , odd-interorbital

$$F_S = \sum_{i \neq j=1,\dots,N} F_{ij}$$

$$F_S(-\omega) = F_S(\omega)$$

even- ω , even-interorbital



Examples Odd- ω Pairing

- No interorbital pairing and $\epsilon_{ij} = \Gamma$:

$$F_{\text{odd}} = 2\Gamma i\omega [\Delta_1(\epsilon_2 - \epsilon_3)(\epsilon_2 + \epsilon_3 + \Gamma) + |\Delta_1|^2(\Delta_3 - \Delta_2) + \text{two cyclic permutations}] / D_3$$

- Only $\epsilon_{23} \neq 0, \Delta_{23} \neq 0$ (xz, yz \rightarrow α, β bands with hybridization, xy \rightarrow γ band)

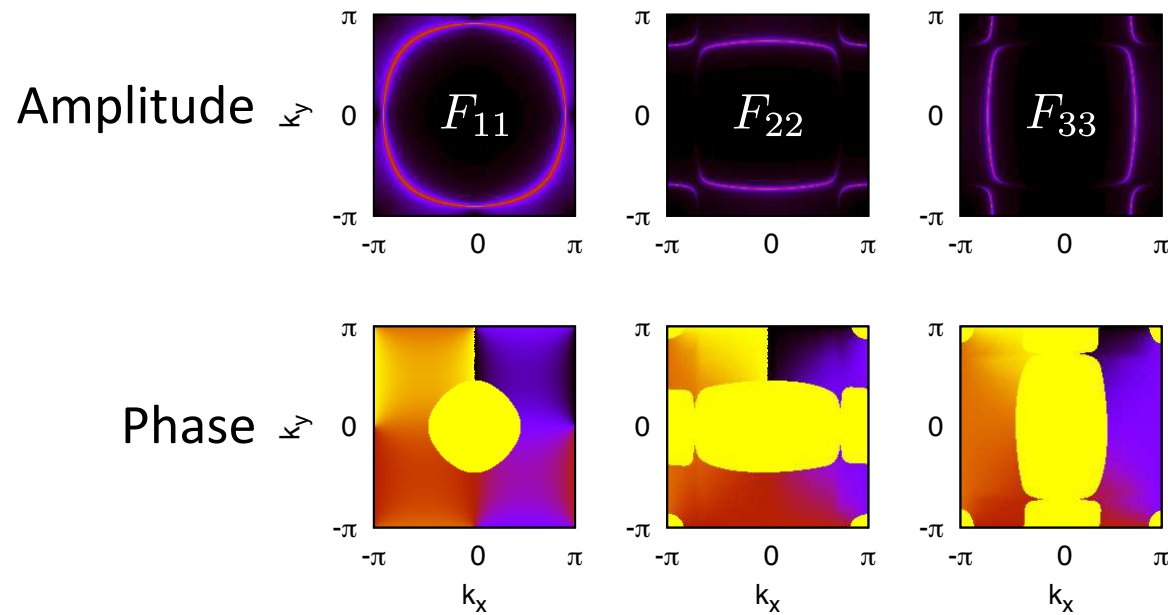
$$F_{\text{odd}} = 2i\omega [\Delta_{23}(\epsilon_3 - \epsilon_2) + \epsilon_{23}(\Delta_3 - \Delta_2)] / D'_3$$

\rightarrow Odd- ω pairing requires only limited interorbital processes, exactly same as finite Kerr rotation [1]



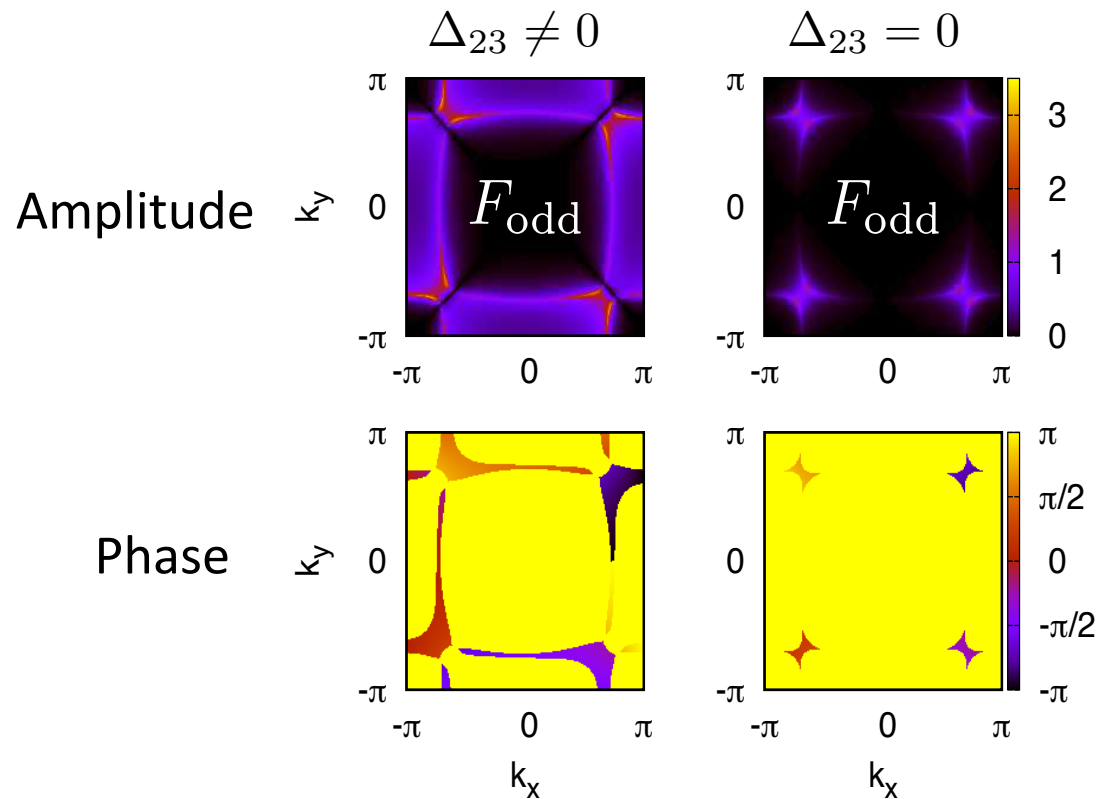
Pair Amplitudes in Sr_2RuO_4

Only $\epsilon_{23} \neq 0, \Delta_{23} \neq 0$



- α, β bands only
- Chiral p -wave symmetry

Odd- ω Pairing in k-Space



- Chiral p -wave symmetry
- Peaked at α, β hybridization
- Nodal structure dependent on exact parameters



Summary

- Bulk odd- ω superconductivity in multiband systems
 - Finite interband/orbital pairing or hybridization
 - Odd- ω pairing in Sr_2RuO_4 measured by Kerr effect
 - Future directions:
 - Other odd- ω bulk superconductors (UPt_3 , ...)
 - Consequences of odd- ω pairing?
 - Band/orbital \rightarrow layer, wires, ...
-

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STRATEGIC RESEARCH





Summary

- Bulk odd- ω superconductivity in multiband systems
 - Finite interband/orbital pairing or hybridization
 - Odd- ω pairing in Sr_2RuO_4 measured by Kerr effect
 - Future directions:
 - Other odd- ω bulk superconductors (UPt_3 , ...)
 - Consequences of odd- ω pairing?
 - Band/orbital \rightarrow layer, wires, ...
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