

Mini-Symposium of Novel Quantum States in Condensed Matter 2017

22 November 2017
YITP, Kyoto

Topological Thouless pumping of ultracold fermions

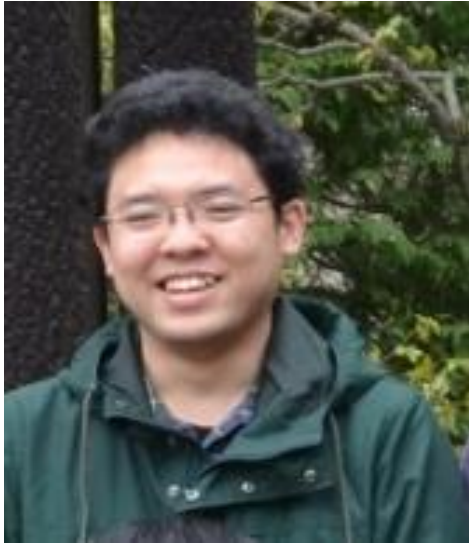
Kyoto University

Yoshiro Takahashi



Collaborators

S. Nakajima
(Kyoto Univ.)



Now in Hakubi

A. Sawada
(Kyoto Univ.)



Now in industry

Y. Kuno
(Kyoto Univ.)



L. Wang
(IOP, CAS)



Also

T. Tomita, S. Taie, H. Ozawa, T. Ichinose (Kyoto Univ.)

M. Troyer (ETH)

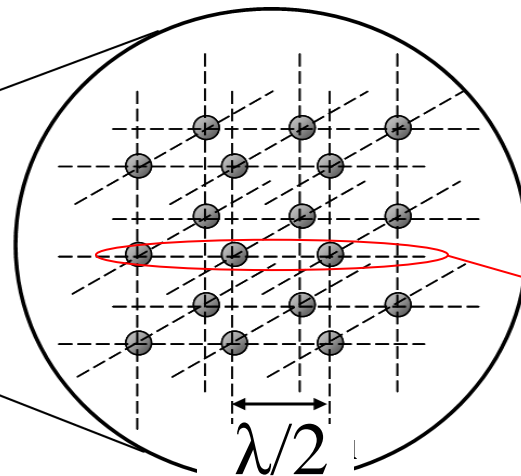
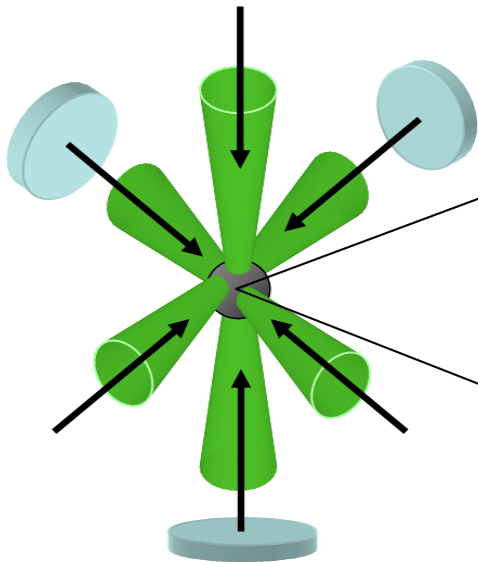
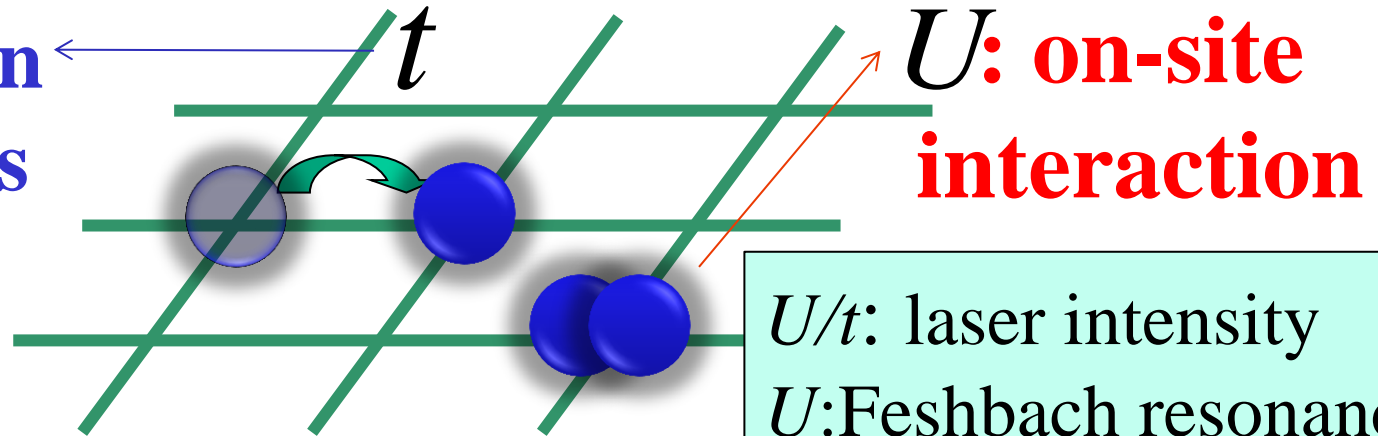
Outline

- I) Demonstration of Topological Thouless Pumping
- II) Effect of Disorder on Topological Thouless Pumping
- III) Prospects on Topological Thouless Pumping

Ultracold Atoms in an Optical Lattice

$$H = -t \sum_{\langle i,j \rangle} c_i^\dagger c_j + U \sum_i n_{i\uparrow} n_{i\downarrow}$$

hopping between adjacent lattices



“Optical lattice”
= periodic potential for atoms

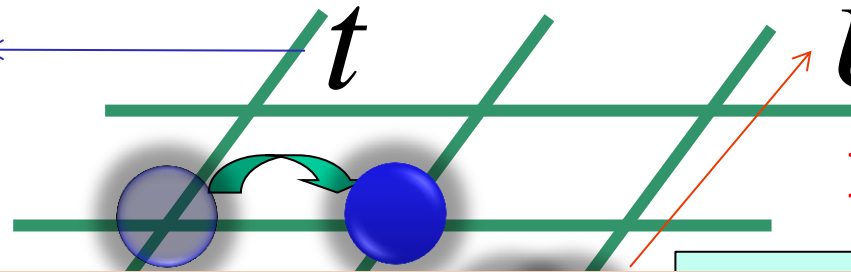
A graph showing a periodic potential V as a function of position x . The potential is represented by a black sine-squared wave. Grey spheres representing atoms are shown trapped in the potential wells. A red arrow points from the text above to the graph.

$$V = V_o \sin^2(kx)$$

Ultracold Atoms in an Optical Lattice

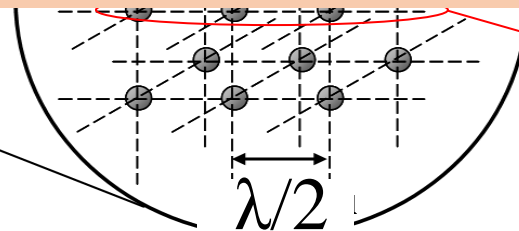
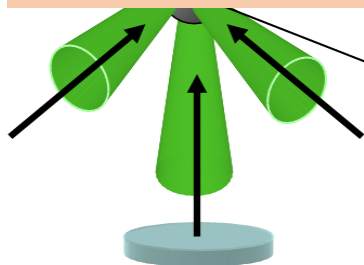
$$H = -t \sum_{\langle i,j \rangle} c_i^\dagger c_j + U \sum_i n_{i\uparrow} n_{i\downarrow}$$

hopping between adjacent lattices

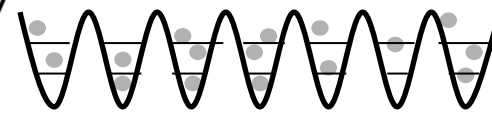


U : on-site interaction

Study of Topological Physics is an important direction of ultracold atoms in an optical lattice



= periodic potential for atoms



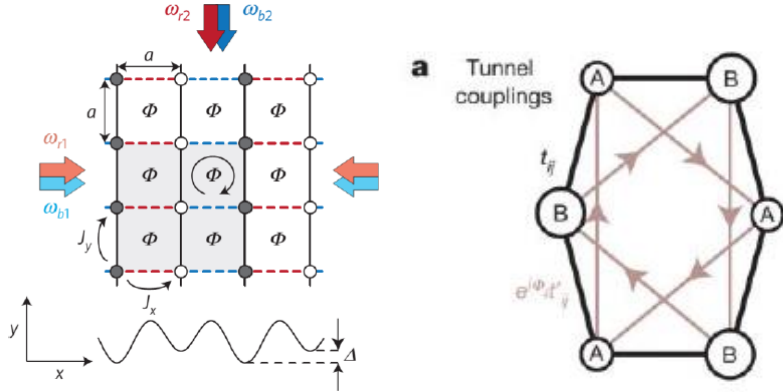
$$V = V_o \sin^2(kx)$$

Exploring topological physics using optical lattice

Spatial 2D:

Hofstadter Hamiltonian (MIT, MPQ)
 topological Haldane model (ETH)

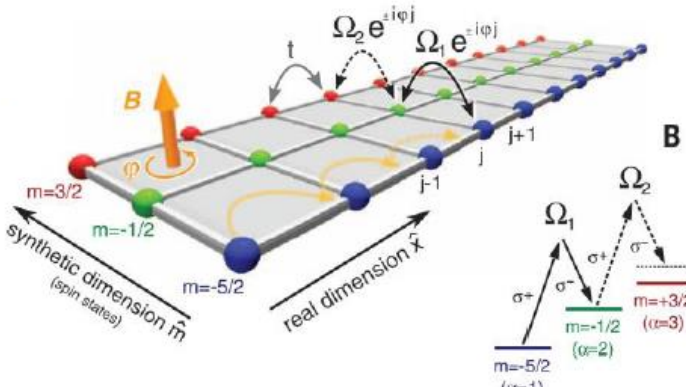
...



Spatial 1D + Synthetic dimension:

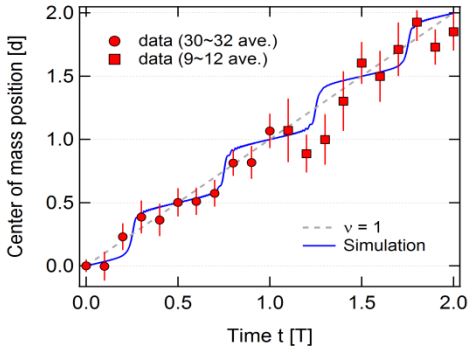
chiral edge state (NIST, LENS)

...



Spatial 1D + Temporal 1D (1+1):

Thouless topological charge pumping
 fermion (Kyoto)
 boson (MPQ)



Thouless Topological Charge Pumping

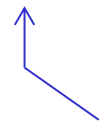
Thouless PRB 27, 6083(1983) “Quantization of Particle Transport”

Particle Transport in a 1D Periodically-driven lattice system:

$H(k+2\pi/L, t) = H(k, t)$: spatially periodic 1D potential
 $H(k, t+T) = H(k, t)$: temporally cyclic adiabatic evolution

pumped charge for one cycle T :

$$C = \frac{i}{2\pi} \sum_{\lambda} f_{\lambda} \int_0^T dt \int_0^{2\pi/L} dk \left[\left\langle \frac{\partial \psi_{\lambda k}}{\partial t} \left| \frac{\partial \psi_{\lambda k}}{\partial k} \right\rangle - \left\langle \frac{\partial \psi_{\lambda k}}{\partial k} \left| \frac{\partial \psi_{\lambda k}}{\partial t} \right\rangle \right]$$



topological invariant
: Chern number

$f_{\lambda} = 1$ for filled bands λ : band index
 $= 0$ for empty bands

Thouless Topological Charge Pumping

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Hall conductivity for Integer Quantum Hall Effect (TKNNformula)

$$\sigma_{xy} = \frac{e^2}{h} \sum_{m \leq m_0} \int_{\text{BZ}} \frac{d^2 \tilde{k}}{2\pi i} \left[\left\langle \frac{\partial u_{m\tilde{k}}}{\partial \tilde{k}_x} \middle| \frac{\partial u_{m\tilde{k}}}{\partial \tilde{k}_y} \right\rangle - \left\langle \frac{\partial u_{m\tilde{k}}}{\partial \tilde{k}_y} \middle| \frac{\partial u_{m\tilde{k}}}{\partial \tilde{k}_x} \right\rangle \right]$$
$$= \frac{e^2}{h} \underline{N_{\text{ch}}} \longleftarrow \text{topological invariant: Chern Number}$$

Thouless Topological Charge Pumping

Thouless PRB 27, 6083(1983) “Quantization of Particle Transport”

Particle Transport in a 1D Periodically-driven lattice system:

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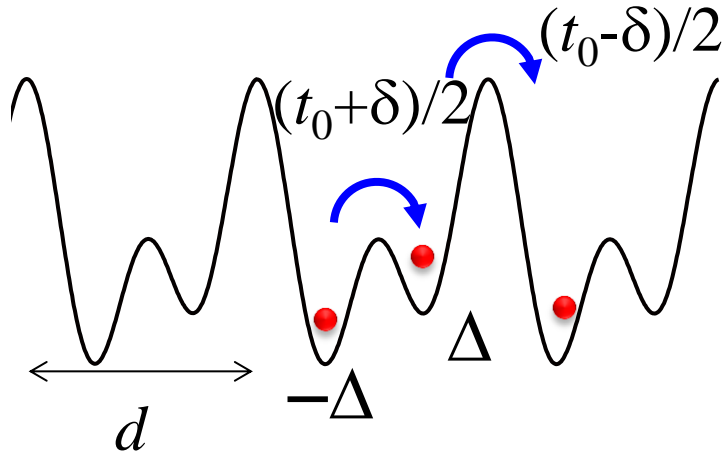
$$C = \frac{i}{2\pi} \sum_{\lambda} f_{\lambda} \int_0^T \underline{dt} \int_0^{2\pi/L} dk \left[\left\langle \frac{\partial \psi_{\lambda k}}{\partial t} \left| \frac{\partial \psi_{\lambda k}}{\partial k} \right\rangle - \left\langle \frac{\partial \psi_{\lambda k}}{\partial k} \left| \frac{\partial \psi_{\lambda k}}{\partial t} \right\rangle \right]$$

“1D quantum charge pumping” shares the same topological origin as “2D IQHE”

Topological Charge Pumping has never been realized in any system

Quantum Rice-Mele Charge Pumping

Rice & Mele, PRL (1982): One-Dimensional Dimerized Lattice Model



Su-Schrieffer-Heeger model

$$\mathcal{H} = \sum_i \left(\frac{t_0}{2} + (-1)^i \frac{\delta}{2} \right) (c_i^\dagger c_{i+1} + \text{h.c.})$$

$$+ \Delta \sum_i (-1)^i c_i^\dagger c_i$$

Staggered
Potential

“two-level description”

$$a_k = \frac{1}{\sqrt{N}} \sum_{j:\text{even}} c_j e^{-ikdj}$$

$$b_k = \frac{1}{\sqrt{N}} \sum_{j:\text{odd}} c_j e^{-ikdj}$$

$$\sigma_x = a_k^\dagger b_k + b_k^\dagger a_k$$

$$\sigma_y = -i(a_k^\dagger b_k - b_k^\dagger a_k)$$

$$\sigma_z = a_k^\dagger a_k - b_k^\dagger b_k$$

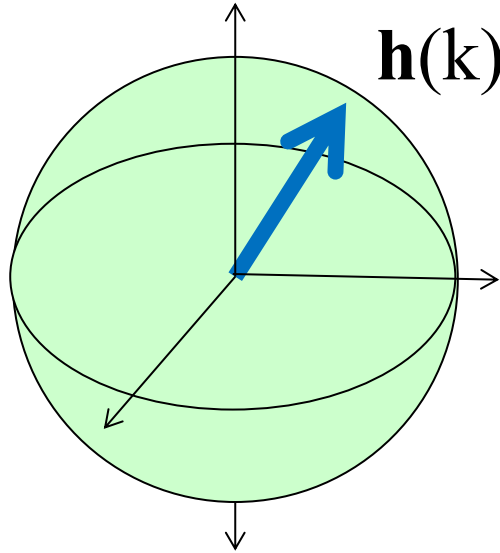
$$H = \sum_k \mathbf{h}(k) \cdot \boldsymbol{\sigma}(k)$$

$$\mathbf{h}(k) = (t_0 \cos(kd/2), -\delta \sin(kd/2), \Delta)$$

$$\mathbf{E}(k) = \pm \sqrt{t_0^2 \cos^2(kd/2) + \delta^2 \sin^2(kd/2) + \Delta^2}$$

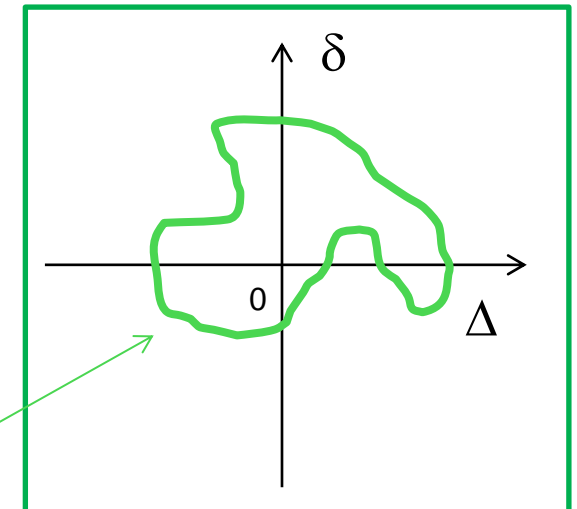
Quantum Rice-Mele Charge Pumping

Rice & Mele, PRL (1982): One-Dimensional Dimerized Lattice Model



$$\mathbf{h}(\mathbf{k}) = (t_0 \cos(kd/2), -\delta(t) \sin(kd/2), \Delta(t))$$

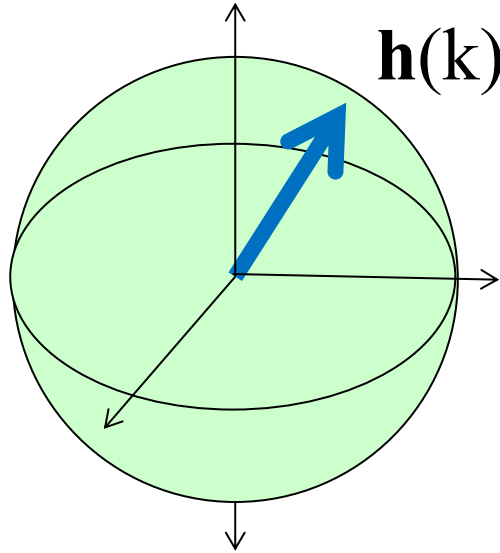
degeneracy point: $\delta(t) = 0, \Delta(t) = 0$



“charge pumping sequence”
↔ “trajectory in $\delta - \Delta$ parameter space”

Quantum Rice-Mele Charge Pumping

Rice & Mele, PRL (1982): One-Dimensional Dimerized Lattice Model

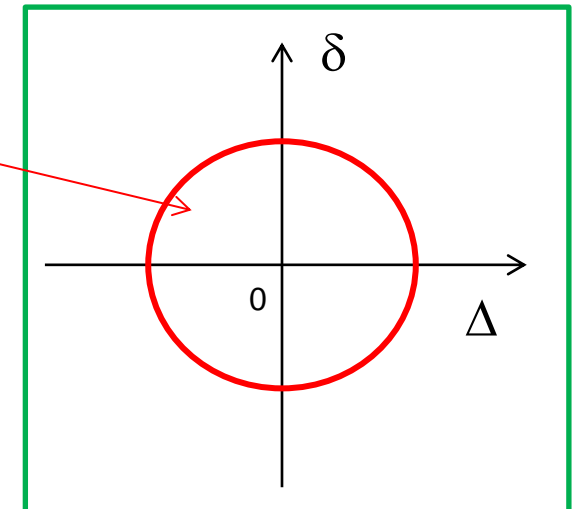


$$\mathbf{h}(\mathbf{k}) = (t_0 \cos(kd/2), -\delta(t) \sin(kd/2), \Delta(t))$$

degeneracy point: $\delta(t) = 0, \Delta(t) = 0$

$$\begin{cases} \delta(t) = \delta_0 \cos\left(\frac{2\pi t}{T}\right) \\ \Delta(t) = \Delta_0 \sin\left(\frac{2\pi t}{T}\right) \end{cases}$$

Pumped Charge: $n = -\text{sgn}(t_0 \delta_0 \Delta_0)$

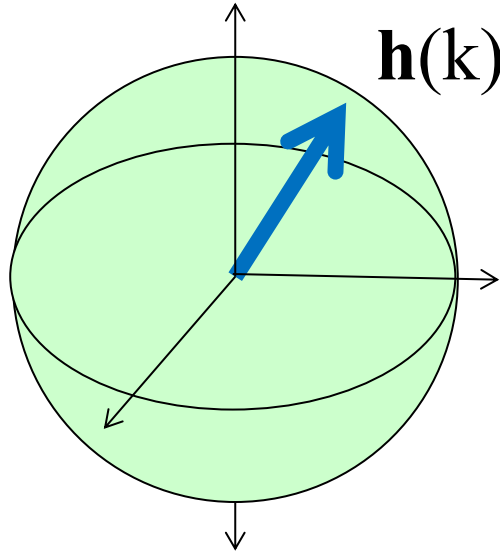


S. -Q, Shen, "Topological Insulators: Dirac Equation in Condensed matter"
D. Xiao *et.al.* RMP (2010)

“charge pumping sequence”
↔ “trajectory in $\delta - \Delta$ parameter space”

Quantum Rice-Mele Charge Pumping

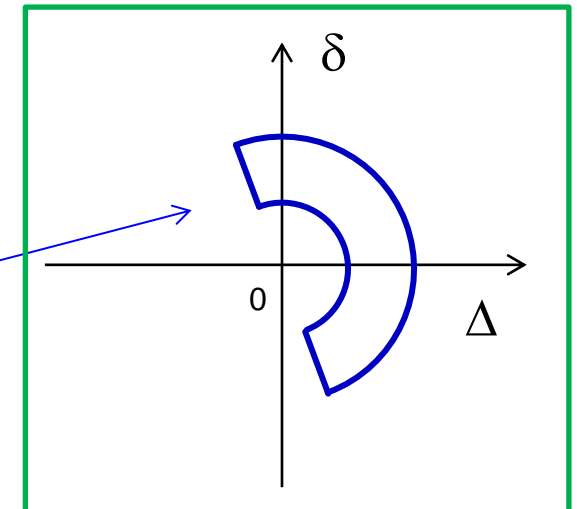
Rice & Mele, PRL (1982): One-Dimensional Dimerized Lattice Model



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degeneracy point: $\delta(t) = 0, \Delta(t) = 0$

Pumped Charge: $n=0$

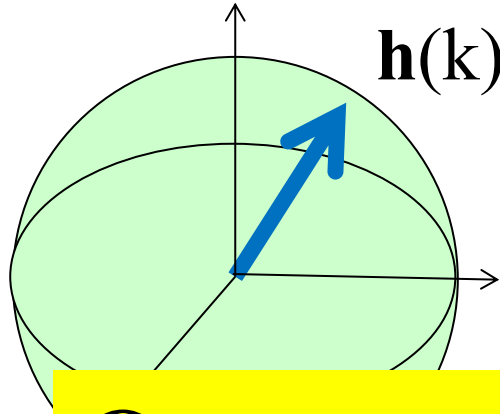


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Quantum Rice-Mele Charge Pumping

Rice & Mele, PRL (1982): One-Dimensional Dimerized Lattice Model

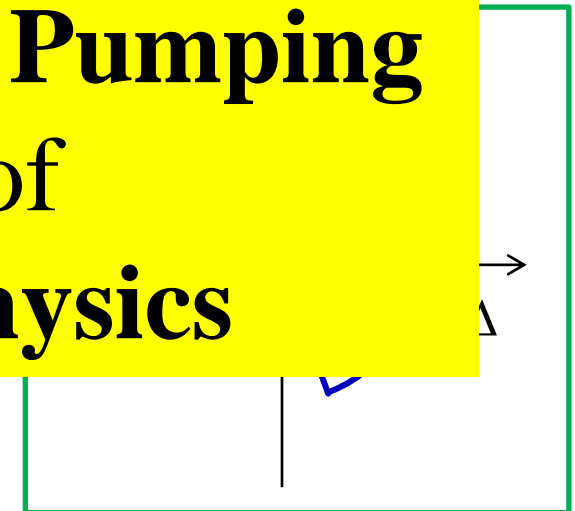


$$\mathbf{h}(\mathbf{k}) = (t_0 \cos(kd/2), -\delta(t) \sin(kd/2), \Delta(t))$$

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Quantum Rice-Mele Charge Pumping
captures the essence of
Topological Quantum Physics

Pumped Charge: $\nu = 0$

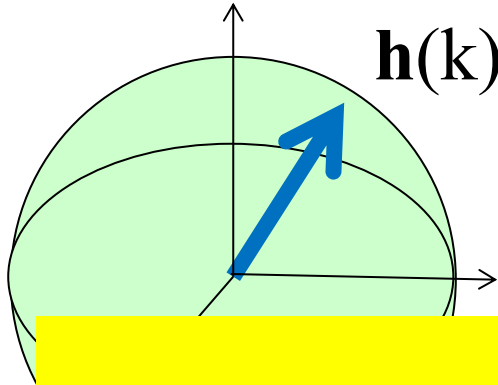


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Quantum Rice-Mele Charge Pumping

Rice & Mele, PRL (1982): One-Dimensional Dimerized Lattice Model



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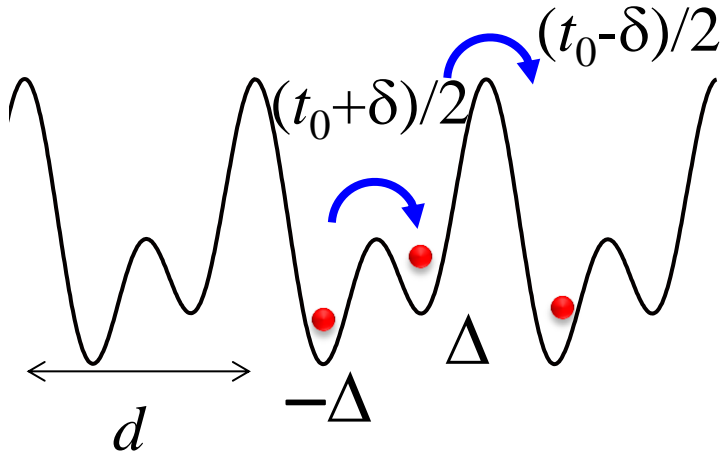
How can we realize
Quantum Rice-Mele Charge Pumping
by using
Ultracold Atoms in an Optical Lattice ?

S. -Q, Shen, "Topological Insulators: Dirac Equation in Condensed matter"
D. Xiao *et.al.* RMP (2010)

"charge pumping sequence"
 \leftrightarrow "trajectory in $\delta - \Delta$ parameter space"

Quantum Rice-Mele Charge Pumping

Rice & Mele, PRL (1982): One-Dimensional Dimerized Lattice Model



Su-Schrieffer-Heeger model

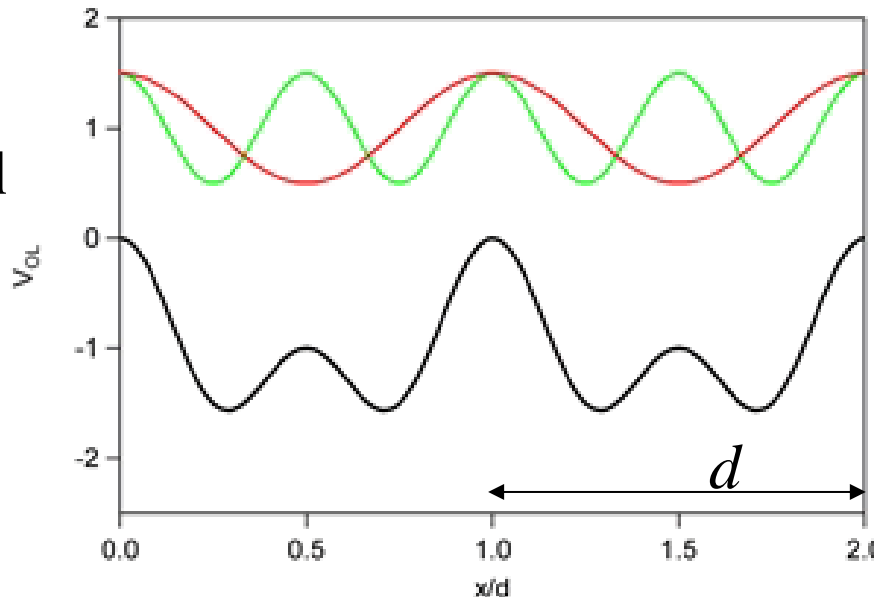
$$\mathcal{H} = \sum_i \left(\frac{t_0}{2} + (-1)^i \frac{\delta}{2} \right) (c_i^\dagger c_{i+1} + \text{h.c.})$$

$$+ \Delta \sum_i (-1)^i c_i^\dagger c_i$$

Staggered Potential

L. Wang *et al*, PRL111, 026802(2013)

proposed
“dynamical optical
super-lattice”
scheme



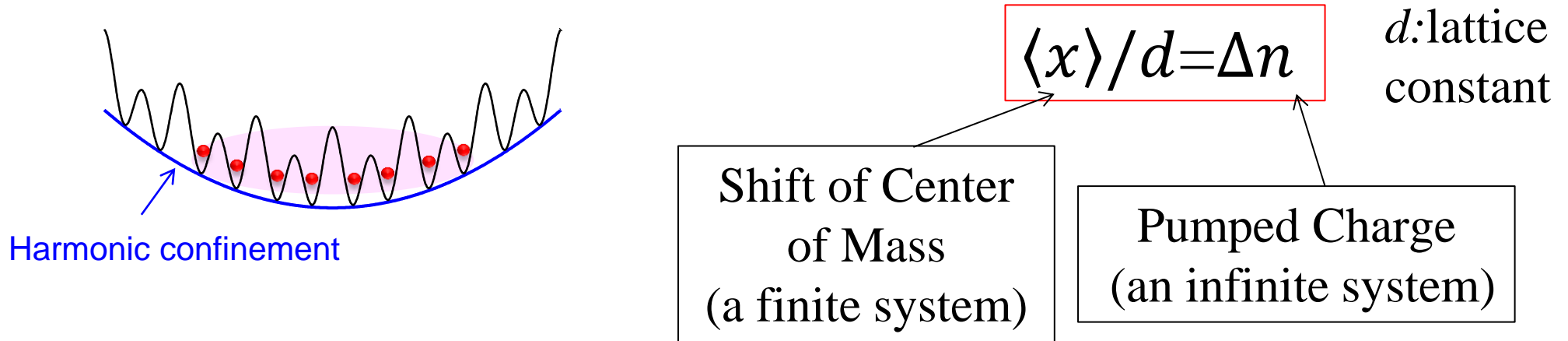
V_S : static short lattice
lattice constant $d/2$

V_L : dynamic long lattice
lattice constant d

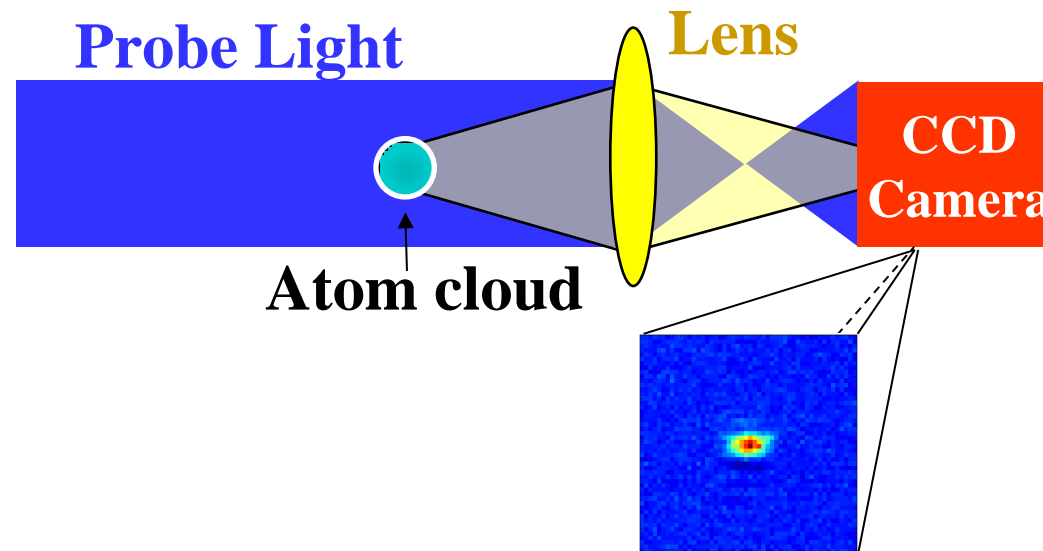
$$V_{OL} = V_S + V_L$$

Experimental Observable: shift of atom cloud

L. Wang *et al*, PRL111, 026802(2013)



"In Situ Absorption Imaging"

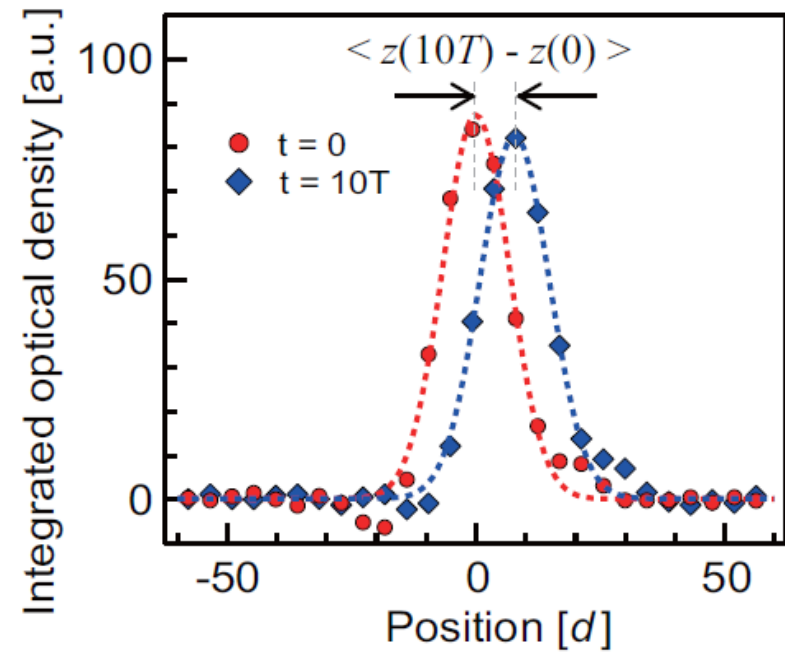
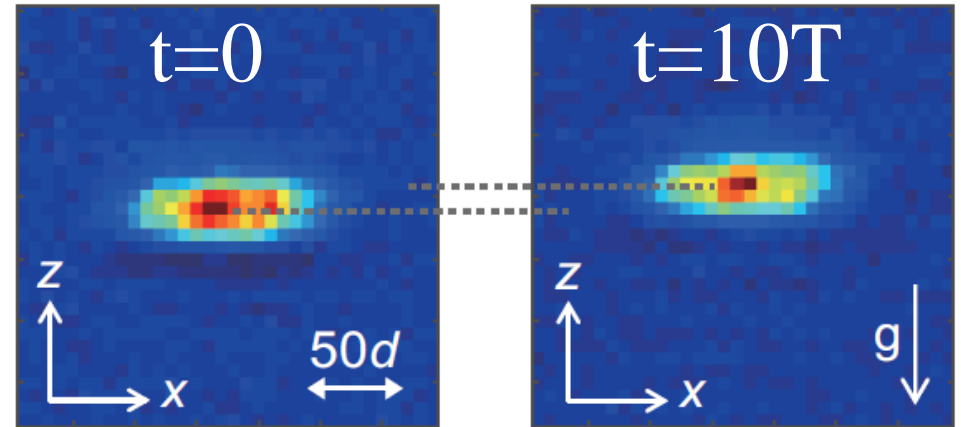
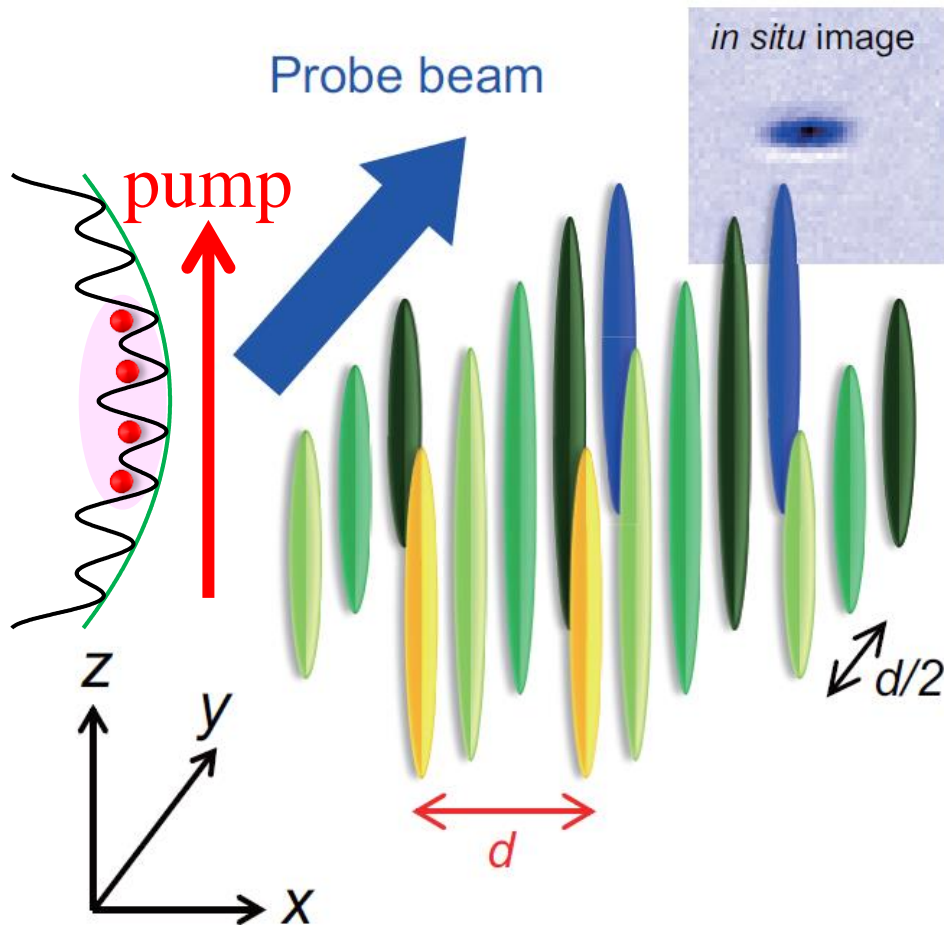


Experimental Observation of Charge Pumping

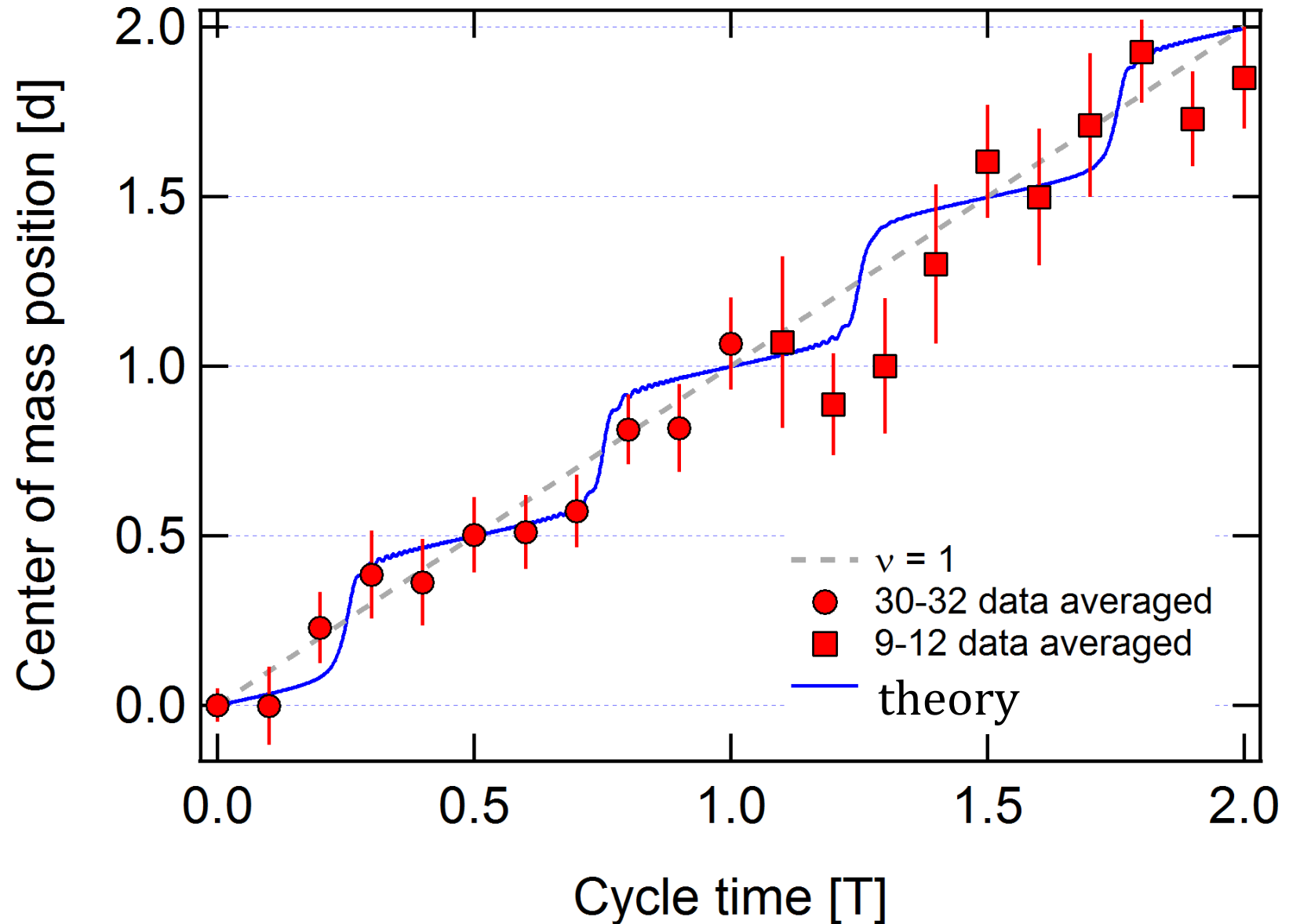
Atomic Species : ^{171}Yb (Fermion, $m_I = \pm 1/2$)

Number of atoms : $5 \times 10^3 \sim 2 \times 10^3$

Temperature : 30 nK



Experimental Observation of Charge Pumping

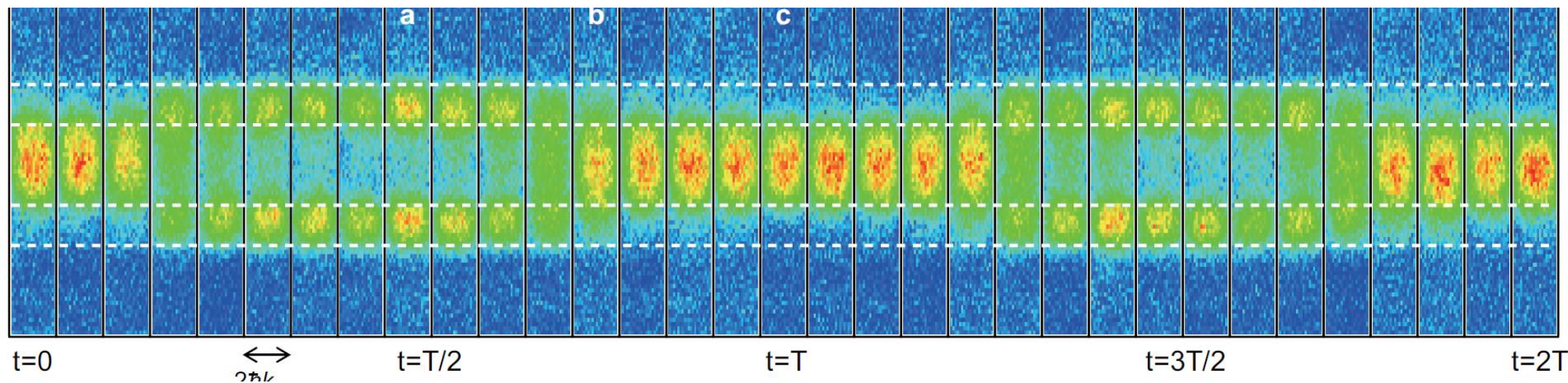
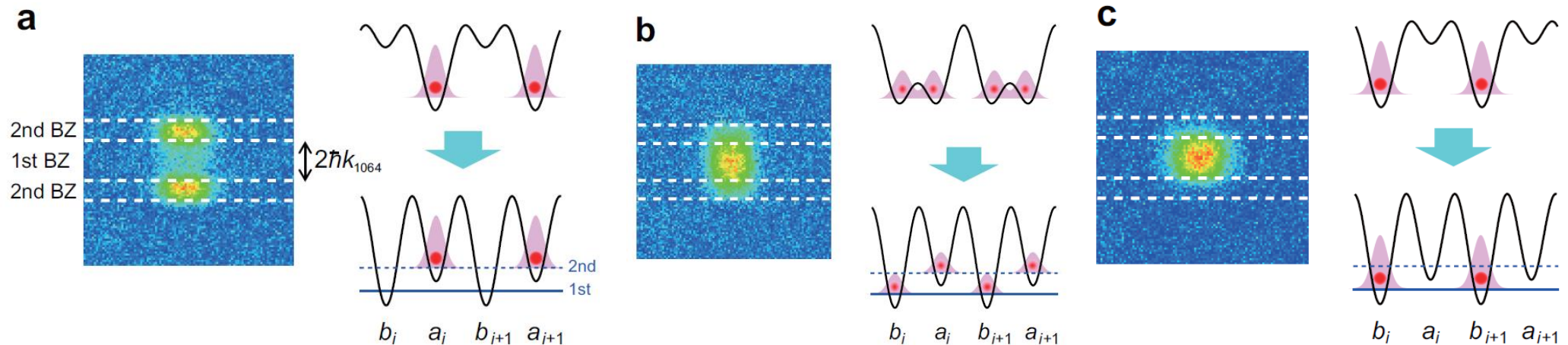


Sublattice Mapping

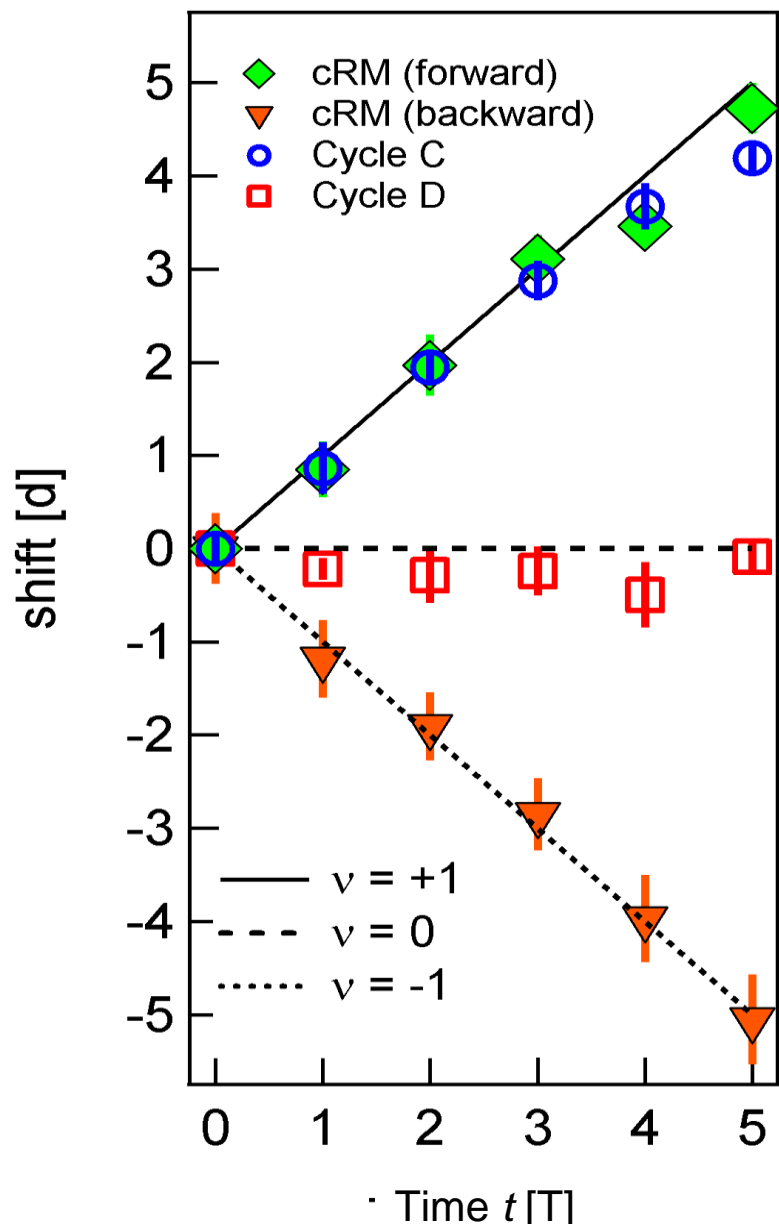
“Right”
in a double well

“superposition”
in a double well

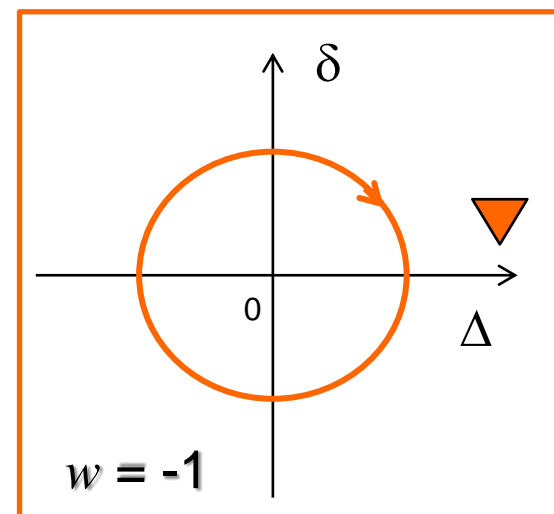
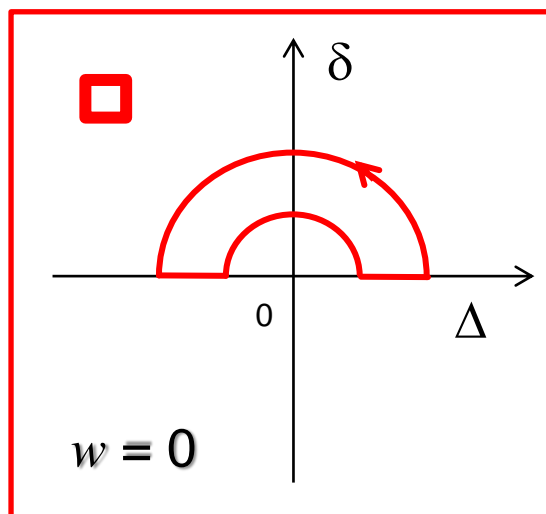
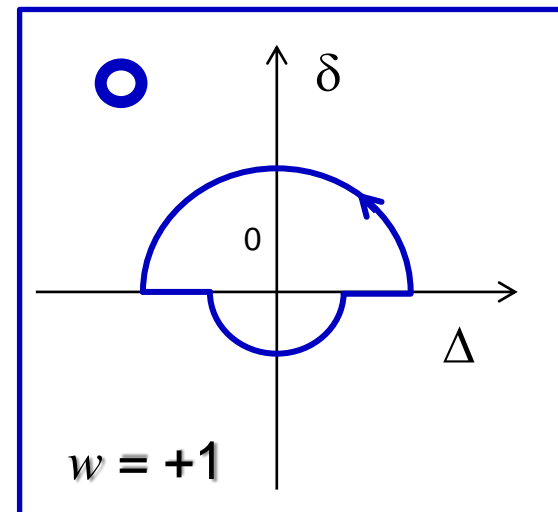
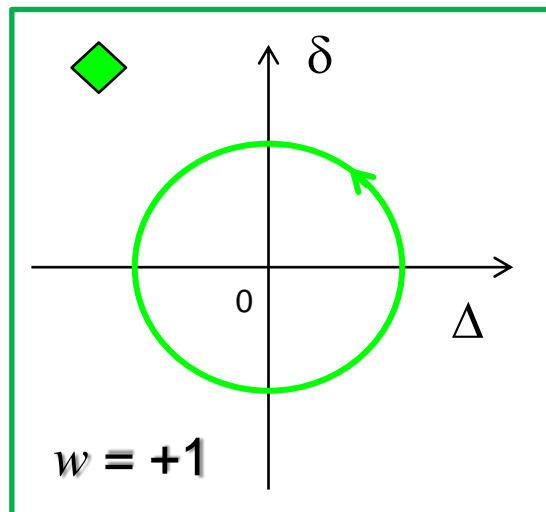
“Left”
in a double well



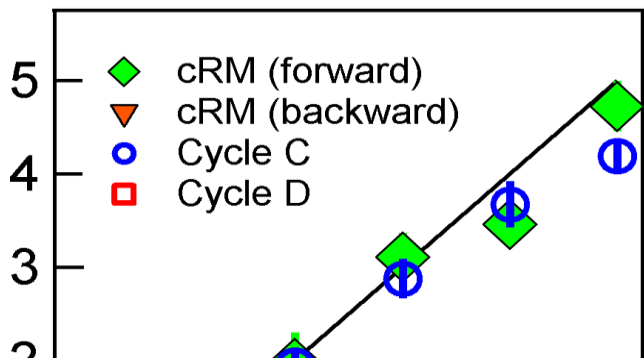
Demonstrating “Topological” Charge Pumping



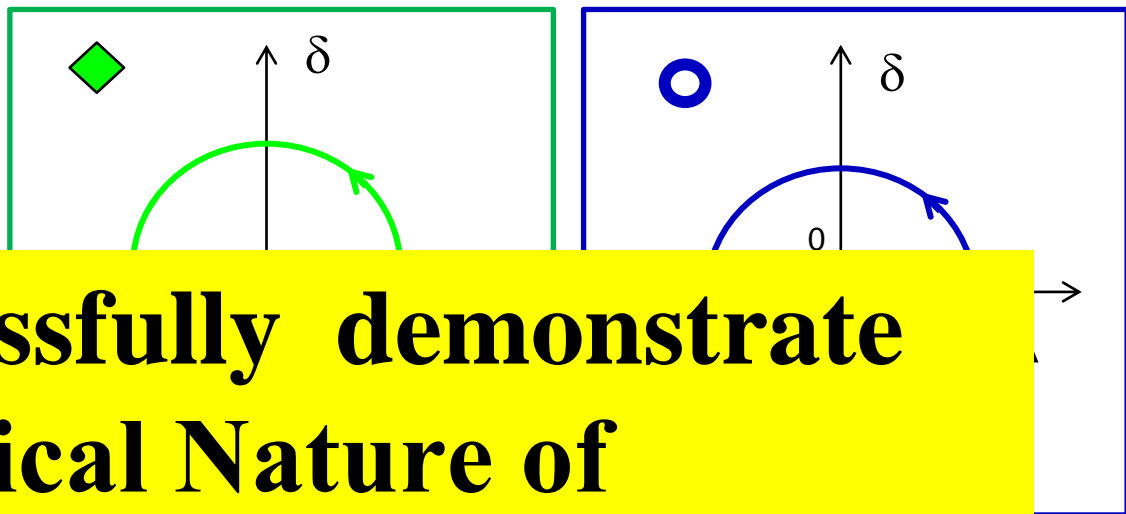
Degeneracy Point: $(\delta, \Delta) = (0, 0)$



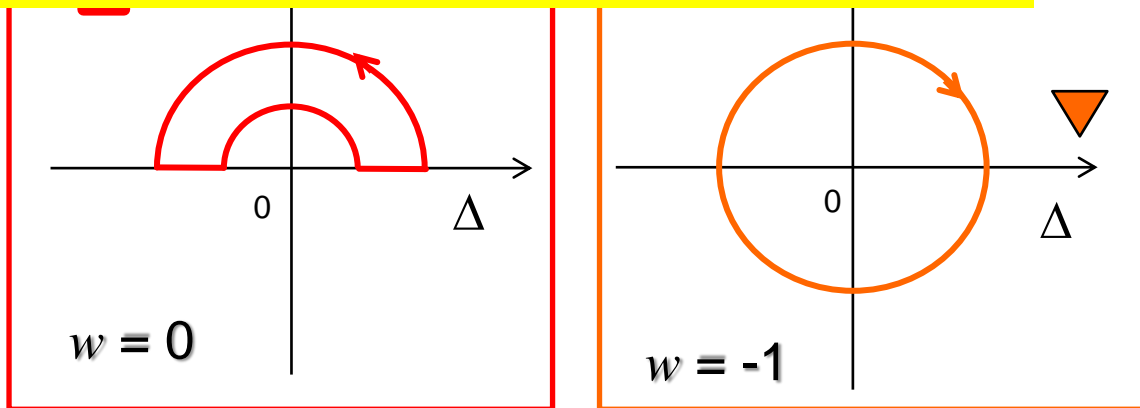
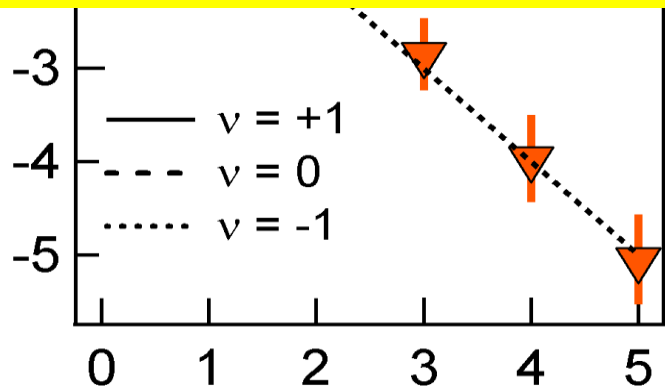
Demonstrating “Topological” Charge Pumping



Degeneracy Point: $(\delta, \Delta) = (0, 0)$



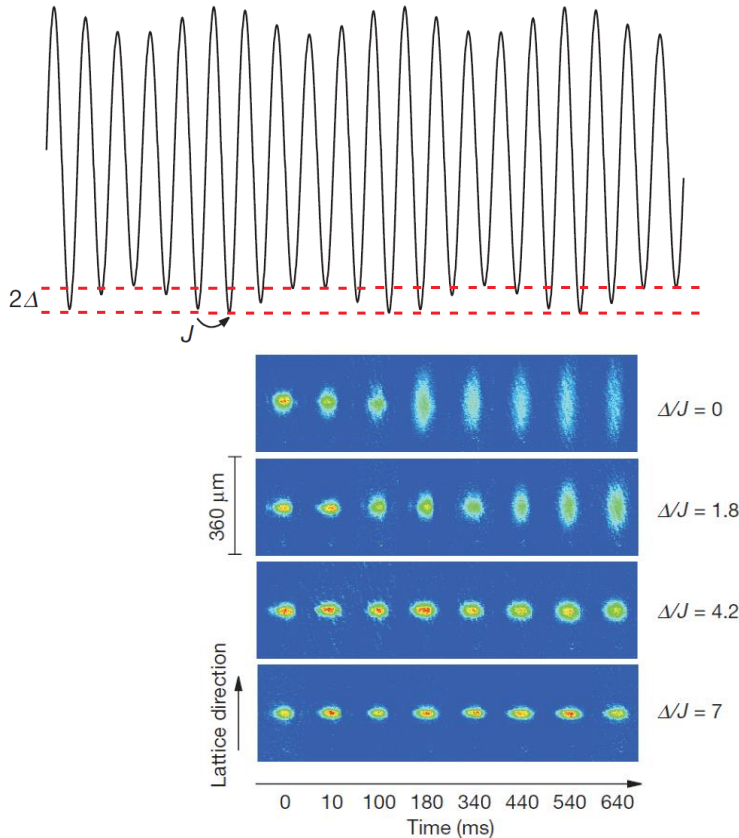
We could successfully demonstrate “Topological Nature of Quantum Rice-Mele Charge Pumping”



Time t [T]

transport is suppressed by disorder

Anderson localization



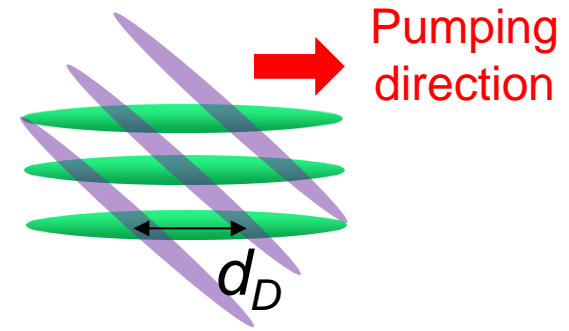
G. Roati *et al.*, Nature **453**, 895 (2008)



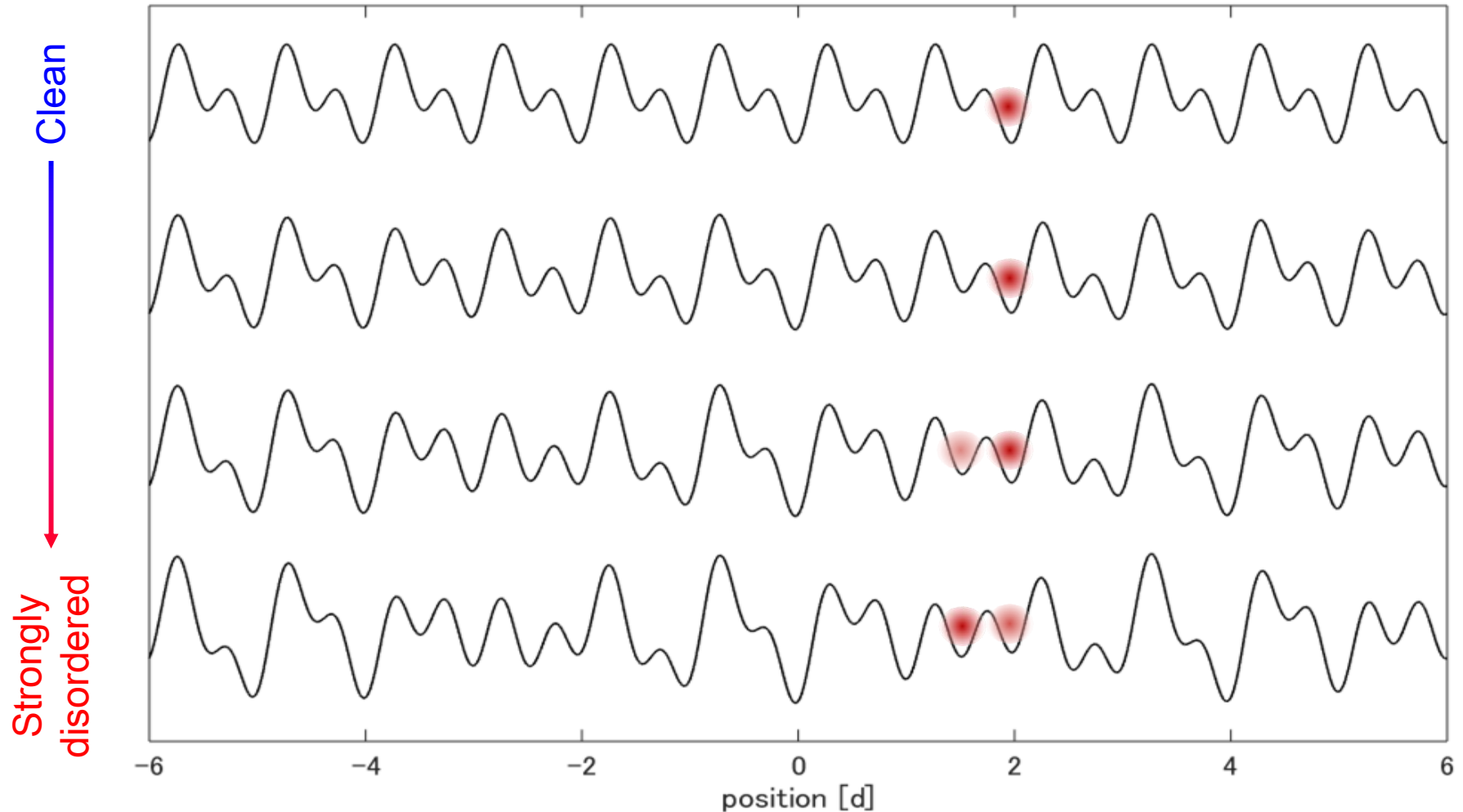
**Does Disorder
Suppress Thouless
Topological Pumping?**

Additional Lattice : V_D

$\lambda=798\text{nm}$, tilted by $45^\circ \Rightarrow d_D: \sim 564\text{nm} \sim 3\sqrt{2}d/4$
($d=532\text{nm}$)



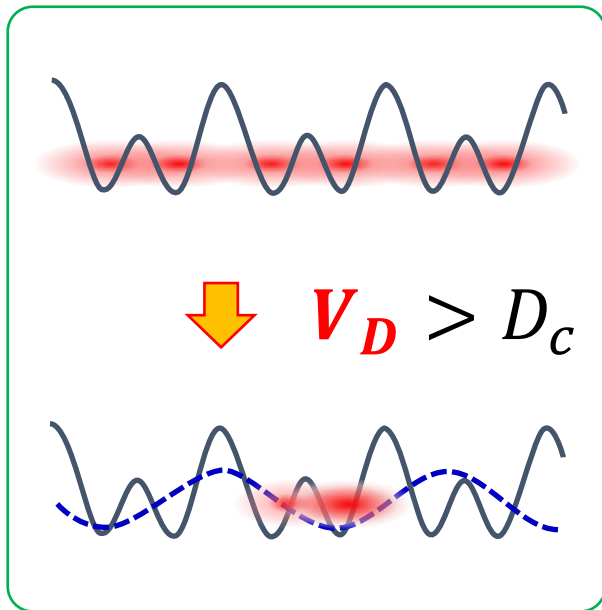
“Quasi-periodic disorder”



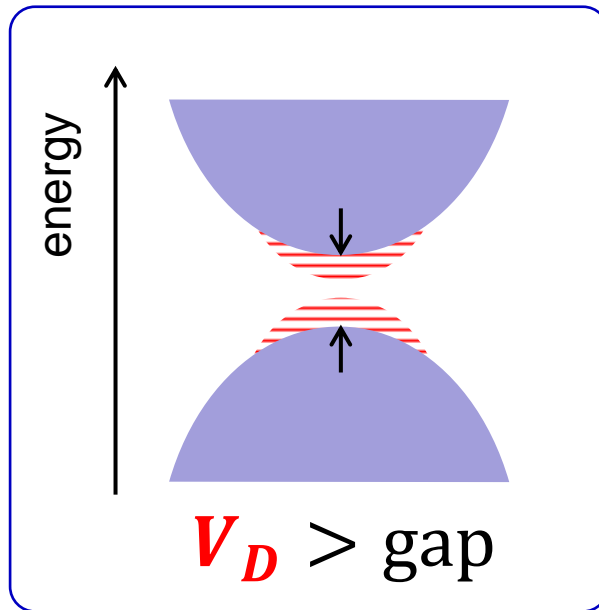
Effect of Disorder for *Topological Charge Pump*

What is the most relevant energy scale for V_D to influence the TCP?

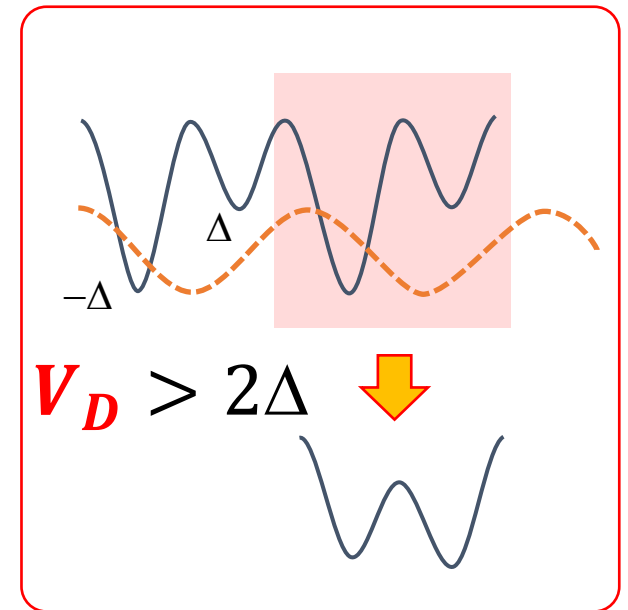
Anderson
Localization
transition point D_c



Band gap



On-site offset 2Δ



Effect of Disorder for *Topological Charge Pump*

Theory

Physics Letters A 380 (2016) 2317–2321



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Quantum pumping induced by disorder in one dimension

1D Wilson-Dirac Model

Jihong Qin^a, Huaiming Guo^{b,*}

^a Department of Physics, University of Science and Technology Beijing, Beijing 100083, China

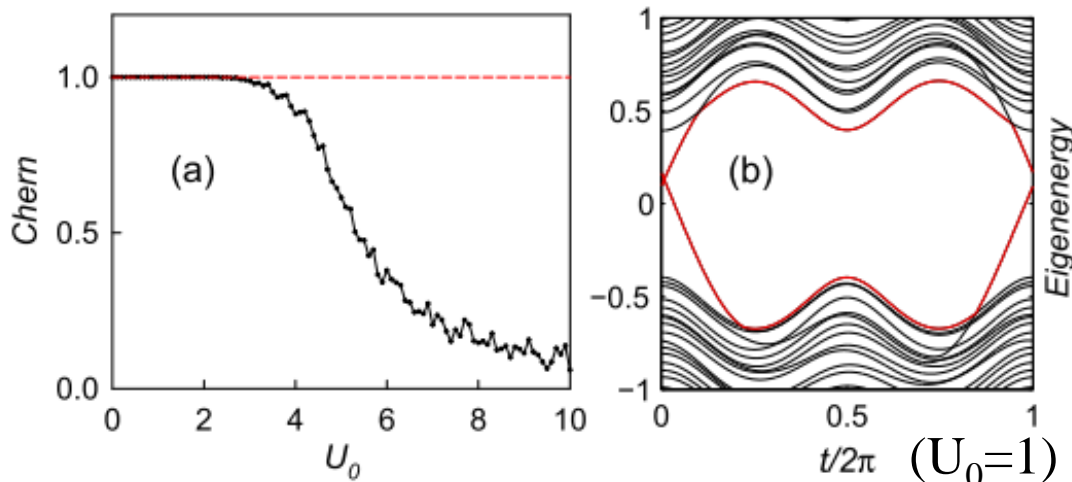
^b Department of Physics, Beihang University, Beijing 100191, China

$$H_0 = \sum_i (M + 2B) \Psi_i^\dagger \sigma_z \Psi_i - \sum_{i, \hat{x}} B \Psi_i^\dagger \sigma_z \Psi_{i+\hat{x}} - \sum_{i, \hat{x}} \text{sgn}(\hat{x}) IA \Psi_i^\dagger \sigma_x \Psi_{i+\hat{x}}$$

Disorder Term

$$H_{dis} = \sum_i U_i (c_i^\dagger c_i + d_i^\dagger d_i),$$

U_i uniformly distributed in $(-\frac{U_0}{2}, \frac{U_0}{2})$.



Prospects : effect of disorder

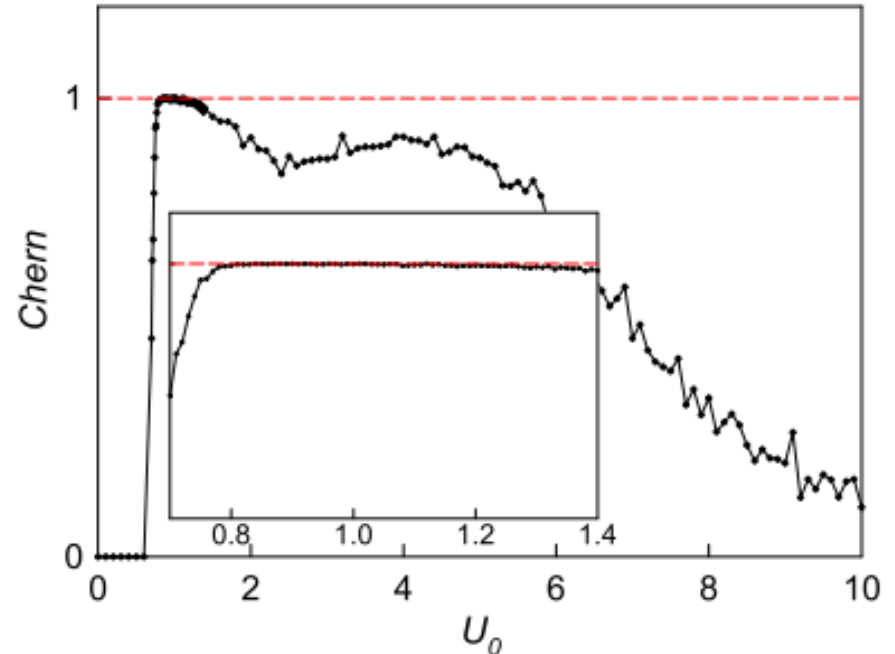
Disorder induced Topological Pumping

1D Wilson-Dirac Model

$$H_0 = \sum_i (M + 2B) \Psi_i^\dagger \sigma_z \Psi_i - \sum_{i, \hat{x}} B \Psi_i^\dagger \sigma_z \Psi_{i+\hat{x}} \\ - \sum_{i, \hat{x}} \text{sgn}(\hat{x}) A \Psi_i^\dagger \sigma_x \Psi_{i+\hat{x}}$$

Disorder Term $H_{dis} = \sum_i U_i (c_i^\dagger c_i + d_i^\dagger d_i),$

U_i uniformly distributed in $(-\frac{U_0}{2}, \frac{U_0}{2})$.



[Phys. Lett. A 380, 2317(2016)]

Time-dependent Disorder

Q Niu and D J Thouless J. Phys. A **17**, 2453 (1984)

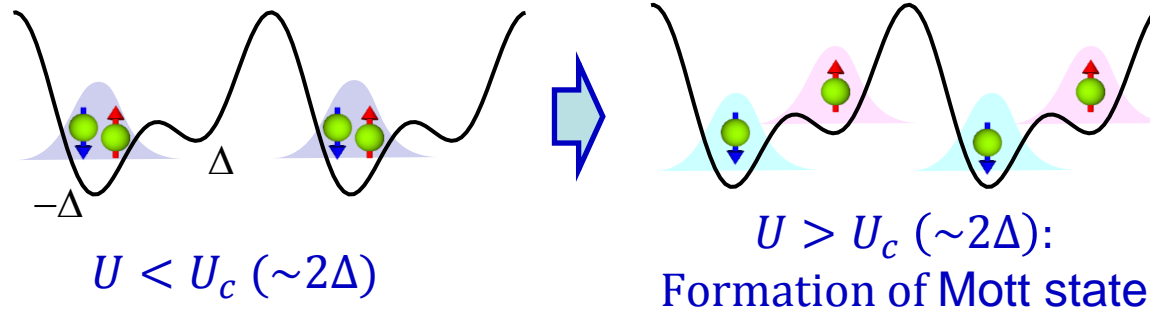
*“**a weak disorder**, which does not close a finite Fermi gap, is irrelevant to the charge transport no matter **whether it is static or varies periodically in time**”*

Non-Quasi-Periodic Disorder

Prospects :effect of interaction

Mott transition in Ionic Hubbard model (M. Nakagawa, private commun.)

(G. Ortiz et al., PRB 54, 13515 (1996))

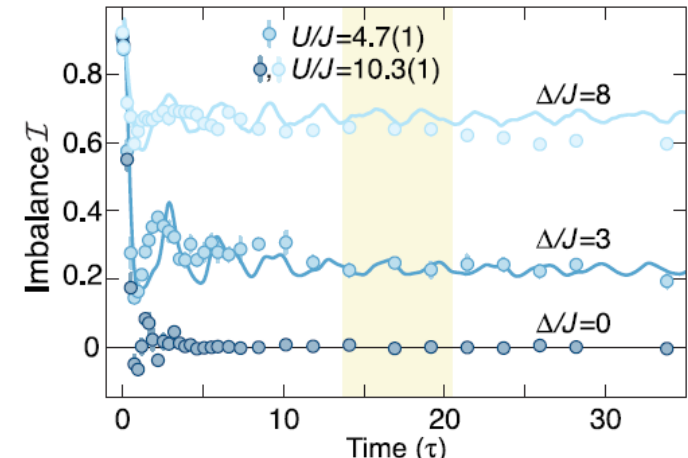
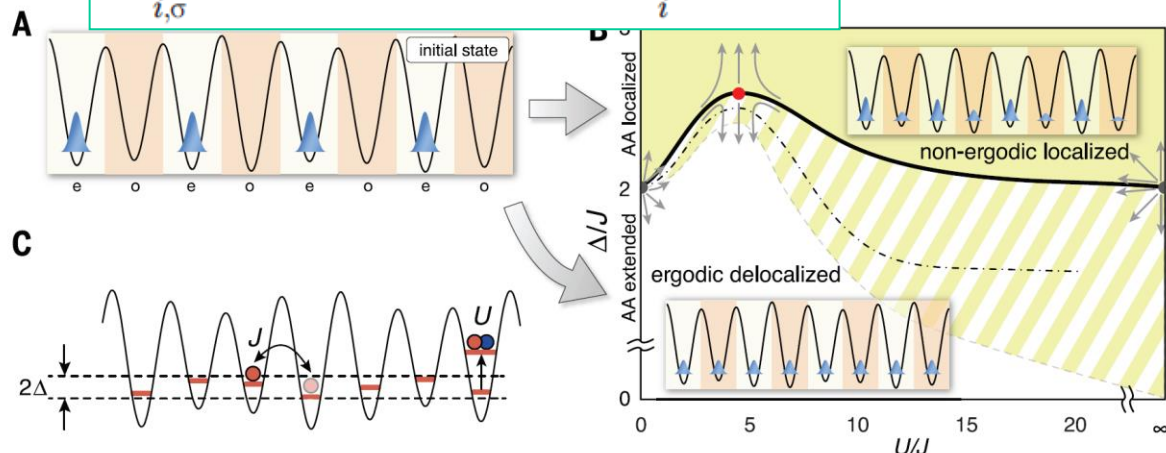


Many-body Localization

(M. Schreiber, et al, Science 349,842(2015))

$$\hat{H} = -J \sum_{i,\sigma} (\hat{c}_{i,\sigma}^\dagger \hat{c}_{i+1,\sigma} + \text{h.c.}) + \Delta \sum_{i,\sigma} \cos(2\pi\beta i + \phi) \hat{c}_{i,\sigma}^\dagger \hat{c}_{i,\sigma} + U \sum_i \hat{n}_{i,\uparrow} \hat{n}_{i,\downarrow}$$

1D, Fermion, Interaction, Disorder



Summary

I) Demonstration of Topological Thouless Pumping

[S. Nakajima, *et al*, Nature Physics, **12**, 296(2016)]

Realization of Rice-Mele model by optical super-lattice

Revealing the topological nature

II) Effect of Disorder on Topological Thouless Pumping

Successful implementation of quasi-periodic disorder potential

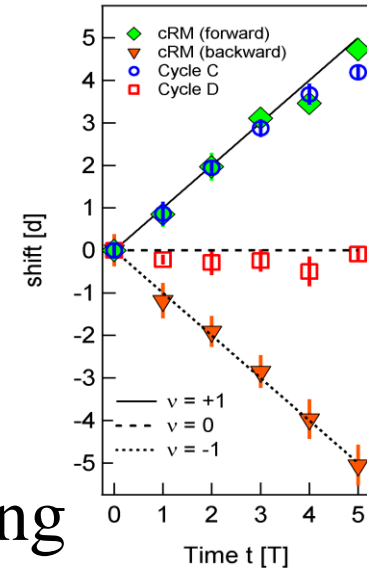
Demonstrating the robustness against the disorder

Suggesting the topological phase transition

III) Prospects on Topological Thouless Pumping

Various topics on disorder

Effect of interaction on topological Thouless pumping



Thank you very much for attention



16 August Mount Daimonji at Kyoto